



**HAL**  
open science

# Thermodynamic form of the equation of motion for perfect fluids of grade $n$

Henri Gouin

► **To cite this version:**

Henri Gouin. Thermodynamic form of the equation of motion for perfect fluids of grade  $n$ . Comptes Rendus de l'Academie des Sciences Serie II, 1987, 305 (II), pp.833-838. hal-00488912

**HAL Id: hal-00488912**

**<https://hal.science/hal-00488912>**

Submitted on 3 Jun 2010

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Thermodynamic form of the equation of motion for perfect fluids of grade $n$

*Translation of C. R. Acad. Sci. Paris t. 305, II, p. 833-838 (1987)*

Henri Gouin

*Université d'Aix-Marseille & C.N.R.S. U.M.R. 6181,  
Case 322, Av. Escadrille Normandie-Niemen, 13397 Marseille Cedex 20 France*

---

## Abstract

We propose a thermodynamic form of the equation of motion for perfect fluids of grade  $n$  which generalizes the one given by J. Serrin in the case of perfectly compressible fluids ([1], p. 171). First integrals and circulation theorems are deduced and a classification of the flows is given.

*Key words:* Conservative fluid motions ; Thermodynamics ; Fluids of grade  $n$  .  
*PACS:* 45.10.Db ; 47.10.+g ; 47.15.Ki ; 47.17.+e

---

## 1 Introduction

In continuum mechanics, first gradient media cannot give a model for fluids with strong variations of density. Material surfaces need their own characteristic behavior and properties of energy [2]. D.J. Korteweg has been the first to point out the convenience of fluids of grade upper than one and not to model interfacial layers by means of discontinuity surfaces in liquid-vapor interfaces [3]. Recently, this approach was used in the case of dynamic changes of phases [4].

Till now, the thermodynamics of fluids was neglected and it was not possible to model flows with strong variations of temperature such as those associated with combustion phenomena or non isothermal interfaces. In fact, it is difficult to take the gradients of temperature into account (it is not possible to consider

---

*Email address:* [henri.gouin@univ-cezanne.fr](mailto:henri.gouin@univ-cezanne.fr) (Henri Gouin).

a virtual displacement of temperature) but we can use the specific entropy through internal energy density.

To this aim, we usefully describe conservative flows of a compressible fluid by the means of its internal energy depending on both entropy and density. This will only be a limit mathematical model, however the study of mathematical structure of the equations of motion is fundamentally necessary. So, to improve the model for strong variations of entropy and density in interfacial layers we consider an internal energy function of the two quantities and their spatial gradients up to a  $n - 1$  convenient order: the fluid will be said *of grade  $n$*  [5].

The paper aims to show that the equations of motion for perfect fluids of any grade can be written in an universal thermodynamic form structurally similar to the one given by J. Serrin in the case of conservative perfect fluids [1]. When the thermodynamic form is applied to the second gradient fluids [6,7], it leads to three results: first integrals associated with circulation theorems (such as Kelvin theorems), potential equations representing the motion of the fluid [8,9] and a classification of the flows similar to the one of compressible perfect fluids [10,11].

## 2 Fluids of grade $n$ .

Perfect fluids of grade  $n$  ( $n$  is any integer) are continuous media with an internal energy per unit mass  $\varepsilon$  which is a function of the specific entropy  $s$  and the density  $\rho$  in the form :

$$\varepsilon = \varepsilon(s, \text{grad } s, \dots, (\text{grad})^{n-1} s, \rho, \text{grad } \rho, \dots, (\text{grad})^{n-1} \rho)$$

where  $(\text{grad})^p$ ,  $p \in \{1, \dots, n - 1\}$ , denotes the successive gradients in the space  $D_t$  occupied by the fluid at present time. We easily may consider the case of an inhomogeneous fluid but, for the sake of simplicity we will not do it. So, the material is supposed to have infinitely short memory and the motion history until an arbitrarily chosen past does not affect the determination of the stresses at present time.

## 3 Equation of motion written in thermodynamic form.

The virtual works principle (or the virtual powers principle) is a convenient way to find the equation of motion. For conservative motions, it writes as the Hamilton principle [8].

A particle is identified in Lagrange representation by the position  $\mathbf{X}(X_1, X_2, X_3)$

occupied in the reference space  $D_0$ . At time  $t$ , its position is given in  $D_t$  by the Eulerian representation  $\mathbf{x}(x_1, x_2, x_3)$ .

The variations of particles motion are deduced from the function family :

$$\mathbf{X} = \psi(\mathbf{x}, t; \alpha) \quad (1)$$

where  $\alpha$  denotes the parameter defined in the vicinity of 0 associated with a family of virtual motions of the fluid. The real motion corresponds with  $\alpha = 0$  [1].

Virtual displacements associated with any variation of the real motion can be written in the form :

$$\delta \mathbf{X} = \left. \frac{\partial \psi}{\partial \alpha}(\mathbf{x}, t; \alpha) \right|_{\alpha=0}.$$

This variation is *dual* and mathematically equivalent to Serrin's one ([1], p. 145, [8]). Let  $L$  be the Lagrangian of the fluid of grade  $n$  :

$$L = \rho \left( \frac{1}{2} \mathbf{V}^* \mathbf{V} - \varepsilon - \Omega \right)$$

where  $\mathbf{V}$  denotes the velocity of particles,  $\Omega$  the potential of mass forces defined on  $D_t$  and  $*$  the transposition in  $D_t$ . Between times  $t_1$  and  $t_2$ , the *Hamilton action* writes [1,8] :

$$a = \int_{t_1}^{t_2} \int_{D_t} L \, dv \, dt.$$

where  $dv$  denotes the volume element.

The density satisfies the conservation of mass :

$$\rho \det \mathbf{F} = \rho_0(\mathbf{X}) \quad (2)$$

where  $\rho_0$  is defined on  $D_0$  and  $\mathbf{F}$  is the gradient of deformation. The motion is supposed to be conservative, then the specific entropy is constant along each trajectory :

$$s = s_0(\mathbf{X}). \quad (3)$$

Classical calculus of variations yields the variation of Hamilton action :

From

$$\delta a = a'(\alpha)|_{\alpha=0},$$

we deduce,

$$\begin{aligned} \delta a = \int_{t_1}^{t_2} \int_{D_t} & \left[ \left( \frac{L}{\rho} - \rho \varepsilon'_\rho \right) \delta \rho + \rho V_i \delta V_i - \rho \varepsilon'_s \delta s \right. \\ & - \rho (\varepsilon_{,\rho,i} \delta \rho_{,i} + \dots + \varepsilon_{,\rho,i_1 \dots i_{n-1}} \delta \rho_{,i_1 \dots i_{n-1}} + \varepsilon_{,s,i} \delta s_{,i} + \dots \\ & \left. + \varepsilon_{,s,i_1 \dots i_{n-1}} \delta s_{,i_1 \dots i_{n-1}}) \right] dx_1 dx_2 dx_3 dt. \end{aligned} \quad (4)$$

The definition of *dual* virtual motions yields :

$$\delta(\text{grad})^p \rho = (\text{grad})^p \delta \rho \quad \text{and} \quad \delta(\text{grad})^p s = (\text{grad})^p \delta s.$$

By using the Stokes formula, let us integrate by parts. Virtual displacements are supposed to be null in the vicinity of the edge of  $D_t$  and integrated terms are null on the edge. We deduce :

$$\begin{aligned} \delta a = \int_{t_1}^{t_2} \int_{D_t} & \left\{ \left[ \left( \frac{L}{\rho} - \rho \varepsilon'_\rho - \sum_{p=1}^{n-1} (-1)^p (\rho \varepsilon_{,\rho,i_1 \dots i_p})_{,i_1 \dots i_p} \right) \delta \rho \right. \right. \\ & \left. \left. - \left[ \rho \varepsilon'_s + \sum_{p=1}^{n-1} (-1)^p (\rho \varepsilon_{,s,i_1 \dots i_p})_{,i_1 \dots i_p} \right] \delta s - \rho V_i \delta V_i \right\} dx_1 dx_2 dx_3 dt. \end{aligned} \quad (5)$$

With  $\text{div}_p$  denoting the divergence operator iterated  $p$  times on the edge of  $D_t$ , we obtain Rel. (5) in tensorial form :

$$\begin{aligned} \delta a = \int_{t_1}^{t_2} \int_{D_t} & \left\{ \left[ \frac{L}{\rho} - \rho \varepsilon'_\rho - \sum_{p=1}^{n-1} (-1)^p \text{div}_p \left( \rho \frac{\partial \varepsilon}{\partial (\text{grad})^p \rho} \right) \right] \delta \rho \right. \\ & \left. - \left[ \rho \varepsilon'_s + \sum_{p=1}^{n-1} (-1)^p \text{div}_p \left( \rho \frac{\partial \varepsilon}{\partial (\text{grad})^p s} \right) \right] \delta s - \rho \mathbf{V}^* \delta \mathbf{V} \right\} dv dt. \end{aligned}$$

By taking (2) into account, we obtain :

$$\delta \rho = \rho \text{div}_0 \delta \mathbf{X} + \frac{1}{\det \mathbf{F}} \frac{\partial \rho_0}{\partial \mathbf{X}} \delta \mathbf{X}$$

where  $\text{div}_0$  denotes the divergence operator relatively to Lagrange variables in  $D_0$ .

We also get :

$$\delta s = \frac{\partial s_0}{\partial \mathbf{X}} \delta \mathbf{X}.$$

The definition of velocity implies :

$$\frac{\partial \mathbf{X}}{\partial \mathbf{x}}(\mathbf{x}, t) \mathbf{V} + \frac{\partial \mathbf{X}}{\partial t}(\mathbf{x}, t) = 0,$$

therefore

$$\frac{\partial \delta \mathbf{X}}{\partial \mathbf{x}} \mathbf{V} + \frac{\partial \mathbf{X}}{\partial \mathbf{x}} \delta \mathbf{V} + \frac{\partial \delta \mathbf{X}}{\partial t} = 0.$$

Let us consider

$$\delta \mathbf{V} = -F \delta \dot{\mathbf{X}},$$

where  $\dot{\phantom{x}}$  denotes the material derivative. Denoting

$$p = \rho^2 \varepsilon'_\rho + \rho \sum_{p=1}^{n-1} (-1)^p \operatorname{div}_p \left( \rho \frac{\partial \varepsilon}{\partial (\operatorname{grad})^p \rho} \right)$$

$$\theta = \varepsilon'_s + \frac{1}{\rho} \sum_{p=1}^{n-1} (-1)^p \operatorname{div}_p \left( \rho \frac{\partial \varepsilon}{\partial (\operatorname{grad})^p s} \right)$$

$$h = \varepsilon + \frac{p}{\rho} \quad \text{and} \quad m = \frac{1}{2} \mathbf{V}^* \mathbf{V} - h - \Omega,$$

then, Rel. (5) yields :

$$\begin{aligned} \delta a &= \int_{t_1}^{t_2} \int_{D_t} \left[ m \delta \rho - \rho \theta \delta s + \rho (\mathbf{V}^* \mathbf{F}) \delta \mathbf{X} \right] dv dt \\ &= \int_{t_1}^{t_2} \int_{D_0} \rho_0 \left[ (\mathbf{V}^* \dot{\mathbf{F}}) - \theta \operatorname{grad}_0^* s - \operatorname{grad}_0^* m \right] \delta \mathbf{X} dv_0 dt \end{aligned}$$

where  $\operatorname{grad}_0$  denotes the gradient operator in  $D_0$ .

The principle for any displacement  $\delta \mathbf{X}$  null on the edge of  $D_0$ ,  $\delta a = 0$  implies:

$$\mathbf{V}^* \dot{\mathbf{F}} = \theta \operatorname{grad}_0^* s_0 + \operatorname{grad}_0^* m. \quad (6)$$

Noting that  $(\mathbf{\Gamma}^* + \mathbf{V}^* \frac{\partial \mathbf{V}}{\partial \mathbf{x}}) \mathbf{F} = \mathbf{V}^* \dot{\mathbf{F}}$ , we get :

$$\mathbf{\Gamma} = \theta \operatorname{grad} s - \operatorname{grad}(h + \Omega). \quad (7)$$

Taking Rel. (3) into account, we obtain :

$$\dot{s} = 0.$$

Relation (7) is the generalization of Rel. (29.8) in [1]. This is a thermodynamic form of the equation of motion of perfect fluids of grad  $n$ .

Obviously, term  $p$  has the same dimension as pression,  $\theta$  has the same dimension as temperature and  $h$  has the same dimension as specific enthalpy. It

seems natural to call them *pression, temperature and enthalpy of the fluid of grad  $n$* , respectively.

#### 4 Conservative properties of perfect fluids of grad $n$ .

Relation (7) leads to the same conclusions as those obtained in [1], [7,8,9,10,11] but for fluids of grad  $n$ . Let us remind the most important results.

With  $J$  denoting the circulation of velocity vector along a closed fluid curve  $\mathcal{C}$  convected by the flow,

$$\frac{dJ}{dt} = \int_{\mathcal{C}} \theta ds.$$

The *Kelvin theorems* are deduced: the circulation of the velocity vector along a closed, isentropic (or isothermal) fluid curve is constant.

For any motion of fluids of grad  $n$ , we can write the velocity field in the form:

$$\mathbf{V} = \text{grad } \varphi + \psi \text{ grad } s + \tau \text{ grad } \chi, \quad (8)$$

the scalar potentials  $\varphi, \psi, s, \tau$  and  $\chi$  verifying :

$$\dot{\varphi} = \frac{1}{2} \mathbf{V}^* \mathbf{V} - h - \Omega, \quad \dot{\tau} = 0, \quad \dot{\psi} = \theta, \quad \dot{\chi} = 0, \quad \dot{s} = 0. \quad (9)$$

Equations (8) and (9) induce the same classification as for conservative flows of compressible perfect fluids [8,9] :

*Oligotropic motions.* - They are motions for which surfaces of equal entropy are vortex surfaces. The circulation of the velocity vector along a closed, isentropic fluid curve is null. Equation (8) of the motion yields :

$$\mathbf{V} = \text{grad } \varphi + \psi \text{ grad } s.$$

*Homentropic motions.* - In the whole fluid  $s$  is constant and Eq. (8) yields :

$$\mathbf{V} = \text{grad } \varphi + \tau \text{ grad } \chi.$$

The Cauchy theorem can be easily written :

$$\frac{d}{dt} \left( \frac{\text{rot } \mathbf{V}}{\rho} \right) = \frac{\partial \mathbf{V}}{\partial \mathbf{x}} \frac{\text{rot } \mathbf{V}}{\rho}.$$

Denoting  $\mathcal{H} = \frac{1}{2}\mathbf{V}^*\mathbf{V} + h + \Omega$ , Eq. (7) yields the Crocco-Vazsonyi equation generalized to stationary motions of perfect fluids of grad  $n$  :

$$\text{rot } \mathbf{V} \times \mathbf{V} = \theta \text{ grad } s - \text{grad } \mathcal{H}.$$

The laws of conservation expressed by the Kelvin theorems correspond to the group of permutations of particles of equal entropy.

This group keeps the equations of motion invariant. It is associated to an expression of Noether's theorem as in [12]. So, it is natural to conjecture such results for a general fluid whose internal energy is a functional of the density and the entropy.

## References

- [1] J. Serrin, Encyclopedia of Physics, VIII/1, Springer, Berlin, 1959.
- [2] M. Barrere and R. Prud'homme, Equations fondamentales de l'Aerothermochimie, Masson, Paris, 1973.
- [3] D.J. Korteweg, Archives Neerlandaises, XXVIII, p. 1-24, 1901.
- [4] J. Serrin, ed., New Perspectives in Thermodynamics, Springer, New York, 1986.
- [5] C. Truesdell and W. Noll, Encyclopedia of Physics, III/3, Springer, Berlin, 1965.
- [6] P. Germain, J. de Mécanique, 12, p. 235-274, 1973.
- [7] P. Casal and H. Gouin, C.R. Acad. Sci. Paris, 300, série II, p. 231-234 et p. 301-304, 1985.
- [8] P. Casal, J. de Mécanique, 5, p. 149-161, 1966.
- [9] P. Casal and H. Gouin, Lecture Notes in Physics 344, p. 85-98, 1989, and arxiv:0803.3160.
- [10] H. Gouin, Research Notes in Mathematics, 46, Pitman, London, p. 128-136, 1981.
- [11] H. Gouin, J. de Mécanique, 20, p. 273-287, 1981.
- [12] H. Gouin, Mech. Res. Comm., 3, p. 151-156, 1976.