Critical exponent for damped wave equations with nonlinear memory

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Abstract

We consider the Cauchy problem in $\mathbb{R}^n$, $n \geq 1$, for a semilinear damped wave equation with nonlinear memory. Global existence and asymptotic behavior as $t \to \infty$ of small data solutions have been established in the case when $1 \leq n \leq 3$. Moreover, we derive a blow-up result under some positive data for in any dimensional space. It turns out that the critical exponent indeed coincides with the one to the corresponding semilinear heat equation.

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1. Introduction

This paper concerns with the Cauchy problem for the damped wave equation with nonlinear memory

$$\begin{cases} u_{tt} - \Delta u + u_t = \displaystyle\int_0^t (t-s)^{-\gamma} |u(s)|^p \, ds & t > 0, x \in \mathbb{R}^n, \\ u(0,x) = u_0(x), \quad u_t(0,x) = u_1(x) & x \in \mathbb{R}^n, \end{cases} \tag{1.1}$$

where the unknown function u is real-valued, $n \geq 1$, $0 < \gamma < 1$ and $p > 1$. Throughout this paper, we assume that

$$(u_0, u_1) \in H^1(\mathbb{R}^n) \times L^2(\mathbb{R}^n) \tag{1.2}$$

and

$$\mathrm{supp} u_i \subset B(K) := \{x \in \mathbb{R}^n : \ |x| < K\}, \quad K > 0, \ i = 0, 1. \tag{1.3}$$

For the simplicity of notations, $\|\cdot\|_q$ and $\|\cdot\|_{H^1}$ $(1 \leq q \leq \infty)$ stand for the usual $L^q(\mathbb{R}^n)$-norm and $H^1(\mathbb{R}^n)$-norm, respectively.

In recent years, questions of global existence and blow-up of solutions for nonlinear hyperbolic equations with a damping term have been studied by many mathematicians, see [9, 10, 13, 18, 20] and the references therein. To focus on our motivation, we shall mention below only some results related to Todorova and Yordanov [20]. For the Cauchy problem for the semilinear damped wave equation with the forcing term

$$u_{tt} - \Delta u + u_t = |u|^p, \qquad u(0) = u_0, \quad u_t(0) = u_1, \tag{1.4}$$

it has been conjectured that the damped wave equation has the diffuse structure as $t \to \infty$ (see e.g. [1, 12]). This suggests that problem (1.4) should have $p_c(n) := 1 + 2/n$ as critical exponent which is called the Fujita exponent named after Fujita [6], in general space dimension. Indeed, Todorova and Yordanov [20] have showed that the critical exponent is exactly $p_c(n)$, that is, if $p > p_c(n)$ then all small initial data solutions of (1.4) are global, while if $1 < p < p_c(n)$ then all solutions of (1.4) with initial data having positive average value blow-up in finite time regardless of the

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smallness of the initial data. Moreover, they showed that in the case of $p > p_c(n)$, the support of the solution of (1.4) is strongly suppressed by the damping, so that the solution is concentrated in a ball much smaller than $|x| < t + K$, namely

$$\|Du(t,\cdot)\|_{L^2(\mathbb{R}^n\setminus B(t^{1/2+\delta}))} = O(e^{-t^{2\delta/4}}), \quad \text{as } t \to \infty,$$

where $D := (\partial_t, \nabla_x)$. Furthermore, they proved that the total energy of the solutions of (1.4) decays at the rate of the linear equation, namely

$$\|Du(t,\cdot)\|_{L^2(\mathbb{R}^n)} = O(t^{-n/4-1/2}), \quad \text{as } t \to \infty.$$

Our goal is to generalize the above properties founded by Todorova and Yordanov to our problem (1.1) with the same assumptions on the initial data. The method used to prove the global existence is the weighted energy method developed in [20]. On the other hand, the test function method (see [3, 4, 5, 14, 15] and the references therein) is the key to prove the blow-up result. We denote that our global existence and asymptotic behavior as $t \to \infty$ for small data solutions are obtained in the case when $1 \le n \le 3$, due to the nonlocal in time nonlinearity. While the blow-up result is done in any dimensional space. Let us present our main results.

First, the following local well-posedness result is needed.

Proposition 1. *Let $1 < p \le n/(n-2)$ for $n \ge 3$, and $p \in (1,\infty)$ for $n = 1, 2$. Under the assumptions* (1.2)-(1.3) *and $\gamma \in (0,1)$, the problem* (1.1) *possesses a unique maximal mild solution u, i.e. satisfies the integral equation* (3.26) *below, such that*

$$u \in C([0,T_{\max}), H^1(\mathbb{R}^n)) \cap C^1([0,T_{\max}), L^2(\mathbb{R}^n)),$$

where $0 < T_{\max} \le \infty$. Moreover, $u(t,\cdot)$ is supported in the ball $B(t+K)$. In addition:

$$\text{either } T_{\max} = \infty \quad \text{or else} \quad T_{\max} < \infty \text{ and } \|u(t)\|_{H^1} + \|u_t(t)\|_2 \to \infty \text{ and } t \to T_{\max}. \tag{1.5}$$

Remark 1. We say that u is a global solution of (1.1) if $T_{\max} = \infty$, while in the case of $T_{\max} < \infty$, we say that u blows up in finite time.

Now, set

$$p_\gamma := 1 + \frac{2(2-\gamma)}{(n-2+2\gamma)_+} \quad \text{and} \quad p^* = \max\left\{\frac{1}{\gamma}; p_\gamma\right\}.$$

Our global existence result is the following

Theorem 1. *Let $p^* < p \le n/(n-2)$ for $n = 3$, and $p^* < p < \infty$ for $n = 1, 2$. Assume that $1/3 < \gamma < 1/2$ and the initial data satisfy* (1.2)-(1.3) *such that $\|u_0\|_{H^1} + \|u_1\|_{L^2}$ is sufficiently small. Then the problem* (1.1) *admits a unique global mild solution*

$$u \in C([0,\infty), H^1(\mathbb{R}^n)) \cap C^1([0,\infty), L^2(\mathbb{R}^n)).$$

The second result is the finite time blow-up of the solution under some positive data which shows that the assumption on the exponent in the above theorem is critical and it is exactly the same critical exponent to the corresponding semilinear heat equation $u_t - \Delta u = \int_0^t (t-s)^{-\gamma} u(s)\, ds$ founded by Cazenave, Dickistein and Weissler [2] and Fino and Kirane [5].

Theorem 2. *Let $1 < p \le n/(n-2)$ for $n \ge 3$, and $p \in (1,\infty)$ for $n = 1, 2$. Assume that $(n-2)/n < \gamma < 1$ and (u_0, u_1) satisfy* (1.2) *such that*

$$\int_{\mathbb{R}^n} u_i(x)\, dx > 0, \quad i = 0, 1. \tag{1.6}$$

If $p \le p*$, then the mild solution of the problem (1.1) *blows up in finite time.*

Since $1/\gamma \le p^*$, so we have the following

Corollary 1. *Let $n \ge 3$ and $1 < p \le n/(n-2)$. Assume that $\gamma < (n-2)/n$ and (u_0, u_1) satisfy* (1.2) *and* (1.6), *then the mild solution of the problem* (1.1) *blows up in finite time.*

As the by-product of our analysis in Theorem 1, we have the following result concerning the asymptotic behavior as $t \to \infty$ of solutions.

Theorem 3. *Under the assumptions of Theorem* 1, *the asymptotic behavior of the small data global solution u of* (1.1) *is given by*

$$\|Du(t,\cdot)\|_{L^2(\mathbb{R}^n \setminus B(t^{1/2+\delta}))} = O(e^{-t^{2\delta/4}}), \quad t \to \infty, \tag{1.7}$$

that is the solution decays exponentially outside every ball $B(t^{1/2+\delta}), \delta > 0$. *Moreover, the total energy decays at the rate of the linear equation*

$$\|Du(t,\cdot)\|_{L^2(\mathbb{R}^n)} = O(t^{-n/4-1/2}), \quad t \to \infty, \tag{1.8}$$

where $D := (\partial t, \nabla_x)$.

As we have seen, we are restricted ourselves in the case of compactly supported data. This restriction leads us to the finite propagation speed property of the wave which plays an important role in the proof of the global solvability together with the weighted energy method. The blow-up result and the local existence theorem could be proved removing the requirement for the compactness assumptions on the support of the initial data. For the global existence without assuming the compactness of support on the initial data, we refer the reader to [7, 8, 9, 16, 17]. In that case, we have to take $u_0 \in H^1(\mathbb{R}^n) \cap L^1(\mathbb{R}^n)$ and $u_1 \in L^2(\mathbb{R}^n) \cap L^1(\mathbb{R}^n)$.

This paper is organized as follows: in Section 2, we present some definitions and properties concerning the fractional integrals and derivatives. Section 3 contains the proofs of the global existence theorem (Theorem 1) and the asymptotic behavior of solution (Theorem 3). Section 4 is devoted to the proof of the blow-up result (Theorem 2). Finally, to make this paper self-contained, we shall sketch the proof of the local existence of solution (Proposition 1) in Appendix.

2. Preliminaries

In this section, we give some preliminary properties on the fractional integrals and fractional derivatives that will be used in the proof of Theorem 2.

If $AC[0,T]$ is the space of all functions which are absolutely continuous on $[0,T]$ with $0 < T < \infty$, then, for $f \in AC[0,T]$, the left-handed and right-handed Riemann-Liouville fractional derivatives $D^{\alpha}_{0|t} f(t)$ and $D^{\alpha}_{t|T} f(t)$ of order $\alpha \in (0,1)$ are defined by

$$D^{\alpha}_{0|t} f(t) := \partial_t J^{1-\alpha}_{0|t} f(t) \quad \text{and} \quad D^{\alpha}_{t|T} f(t) := -\frac{1}{\Gamma(1-\alpha)} \partial_t \int_t^T (s-t)^{-\alpha} f(s)\, ds, \quad t \in [0,T], \tag{2.1}$$

where Γ is the Euler gamma function, and

$$J^{\alpha}_{0|t} f(t) := \frac{1}{\Gamma(\alpha)} \int_0^t (t-s)^{\alpha-1} f(s)\, ds \tag{2.2}$$

is the Riemann-Liouville fractional integral, for all $f \in L^q(0,T)$ $(1 \le q \le \infty)$. We refer the reader to [11] for the definitions above. Furthermore, for every $f, g \in C([0,T])$ such that $D^{\alpha}_{0|t} f(t), D^{\alpha}_{t|T} g(t)$ exist and are continuous, for all $t \in [0,T]$, $0 < \alpha < 1$, we have the formula of integration by parts (see (2.64) p. 46 in [19])

$$\int_0^T \left(D^{\alpha}_{0|t} f\right)(s) g(s)\, ds \;=\; \int_0^T f(s) \left(D^{\alpha}_{t|T} g\right)(s)\, ds. \tag{2.3}$$

Note also that, for all $f \in AC^{n+1}[0,T]$ and all integer $n \ge 0$, we have (see (2.2.30) in [11])

$$(-1)^n \partial^n_t . D^{\alpha}_{t|T} f = D^{n+\alpha}_{t|T} f, \tag{2.4}$$

where

$$AC^{n+1}[0,T] := \{ f : [0,T] \to \mathbb{R} \text{ and } \partial^n_t f \in AC[0,T] \}$$

and ∂_t^n is the usual n times derivative. Moreover, for all $1 \le q \le \infty$, the following formula (see [Lemma 2.4 p.74][11])

$$D^{\alpha}_{0|t} J^{\alpha}_{0|t} = Id_{L^q(0,T)} \tag{2.5}$$

holds almost everywhere on $[0, T]$.
In the proof of Theorem 2, the following results are useful: if $w_1(t) = (1 - t/T)_+^{\sigma}$, $t \ge 0$, $T > 0$, $\sigma >> 1$, then

$$D^{\alpha}_{t|T} w_1(t) = CT^{-\sigma}(T-t)_+^{\sigma-\alpha}, \quad D^{\alpha+1}_{t|T} w_1(t) = CT^{-\sigma}(T-t)_+^{\sigma-\alpha-1}, \quad D^{\alpha+2}_{t|T} w_1(t) = CT^{-\sigma}(T-t)_+^{\sigma-\alpha-1}, \tag{2.6}$$

for all $\alpha \in (0, 1)$; so

$$\left(D^{\alpha}_{t|T} w_1\right)(T) = 0, \quad \left(D^{\alpha}_{t|T} w_1\right)(0) = C\ T^{-\alpha}, \quad \left(D^{\alpha+1}_{t|T} w_1\right)(T) = 0 \quad \text{and} \quad \left(D^{\alpha+1}_{t|T} w_1\right)(0) = C\ T^{-\alpha-1}. \tag{2.7}$$

For the proof of this results, see [3, Preliminaries]. Furthermore, throughout this paper, positive constants will be denoted by C and will change from line to line.

3. Global existence and asymptotic behavior

In view of the Proposition 1, global existence of a solution follows from the boundedness of its energy at all times. To obtain such a priori estimates, we shall proceed our proof based on the weighted energy method recently developed in Todorova and Yordanov [20]. We begin with the identity

$$\begin{aligned} e^{2\psi(x,t)} u_t(u_{tt} - \Delta u + u_t) &= \frac{d}{dt}\Big(\frac{e^{2\psi(x,t)}}{2}(|u_t|^2 + |\nabla u|^2)\Big) - \operatorname{div}(e^{2\psi(x,t)} u_t \nabla u) \\ &\quad - \frac{e^{2\psi(x,t)}}{\psi_t}(\psi_t \nabla u - u_t \nabla \psi)^2 + \frac{e^{2\psi(x,t)}}{\psi_t}(\psi_t + |\nabla\psi|^2 - \psi_t^2), \end{aligned} \tag{3.8}$$

which holds for each function $\psi(x, t)$ and all $x \in \mathbb{R}^n$, $t \ge 0$. In order to get good estimates, we would like to have

$$\psi_t < 0, \quad \psi_t + |\nabla\psi|^2 - \psi_t^2 = 0. \tag{3.9}$$

For this purpose we choose

$$\psi(x,t) = \frac{1}{2}(t + K - \sqrt{(t+K)^2 - |x|^2}), \quad |x| < t + K. \tag{3.10}$$

It is easily checked that the following phase function satisfies (3.9) and has the advantage of being regular on the support of the solution. We also note that, since

$$\sqrt{(t+K)^2 - |x|^2} \le t + K - |x|^2/[2(t+K)],$$

the function ψ satisfies the inequality

$$\psi(x,t) \ge \frac{|x|^2}{4(t+K)}. \tag{3.11}$$

Proof of Theorem 1. Let u be the local solution of the problem (1.1) in $[0, T_{\max})$. Let us introduce the weighted energy functional

$$W(t) := (1+t)^{n/4+1/2} \|Du(t,\cdot)\|_2 + \|e^{\psi(t,\cdot)} Du(t,\cdot)\|_2. \tag{3.12}$$

We will show that $W(t) \le CI_0$, where $I_0 := \|u_0\|_{H^1} + \|u_1\|_2$ is small enough. This not only gives the global existence but also shows that the solution decays at least as fast as that of the linear part $u_{tt} - \Delta u + u_t = 0$. For the rate of the linear problem, see (3.27) below.

The estimate of the first term in (3.12) will be done by the following lemmas.

Lemma 1. *For all $\beta > n/4p + (2-\gamma)/p$, $\delta > 0$ and all $t \in [0, T_{\max})$, the following weighted energy estimate holds*

$$(1+t)^{n/4+1/2} \|Du(t,\cdot)\|_2 \le CI_0 + C(\max_{[0,t]}(1+\tau)^{\beta} \|e^{\delta\psi(\tau,\cdot)} u(\tau,\cdot)\|_{2p})^p. \tag{3.13}$$

Lemma 2. (*[20, Proposition 2.4]*) *Let* $\theta(q) = n(1/2 - 1/q)$ *and* $0 \le \theta(q) \le 1$, *and let* $0 < \sigma \le 1$. *If* $u \in H^1(\mathbb{R}^n)$ *with* $suppu \subset B(t+K)$, $t \ge 0$. *Then*

$$\|e^{\sigma\psi(t,\cdot)}u\|_q \le C_K(1+t)^{(1-\theta(q))/2}\|\nabla u\|_2^{1-\sigma}\|e^{\psi(t,\cdot)}\nabla u\|_2^{\sigma},. \tag{3.14}$$

where $\psi(t,x)$ *is the weight function from* (3.10).

The estimate of the second term in (3.12) is implied by the following lemma.

Lemma 3. *For all* $t \in [0, T_{\max})$, *the following estimate holds*

$$\|e^{\psi(t,\cdot)}Du(t,\cdot)\|_2 \le CI_0 + C\max_{[0,t]} W(\tau)(\max_{[0,t]}(1+\tau)^{(2-\gamma)/p}\|e^{\frac{1}{p}\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p})^p. \tag{3.15}$$

We postpone the proofs of Lemmas 1 and 3 to the end of this section.
It follows from Lemmas 1 and 3 that

$$W(t) \le CI_0 + C(\max_{[0,t]}(1+\tau)^{\beta}\|e^{\delta\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p})^p + C\max_{[0,t]} W(\tau)(\max_{[0,t]}(1+\tau)^{(2-\gamma)/p}\|e^{\frac{1}{p}\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p})^p. \tag{3.16}$$

On the other hand, Lemma 2 with $q = 2p$ and $\sigma = \delta \le 1$ gives

$$\begin{aligned}\|e^{\delta\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p} &\le C(1+\tau)^{1-\theta(2p))/2}\|\nabla u\|_2^{1-\delta}\|e^{\psi(t,\cdot)}\nabla u\|_2^{\delta}\\ &\le C(1+\tau)^{(1-\theta(2p))/2}(1+\tau)^{-(1-\delta)(n/4+1/2)}W(\tau)^{1-\delta}W(\tau)^{\delta}\\ &= C(1+\tau)^{(1-\theta(2p))/2-(1-\delta)(n/4+1/2)}W(\tau).\end{aligned} \tag{3.17}$$

Similarly, we have

$$\|e^{\frac{1}{p}\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p} \le C(1+\tau)^{(1-\theta(2p))/2-(1-1/p)(n/4+1/2)}W(\tau). \tag{3.18}$$

Using (3.17) and (3.18), we obtain from (3.16)

$$\begin{aligned}W(t) \le\ & CI_0 + C\left(\max_{[0,t]}(1+\tau)^{\beta+(1-\theta(2p))/2-(1-\delta)(n/4+1/2)}W(\tau)\right)^p\\ &+C\max_{[0,t]} W(\tau)\left(\max_{[0,t]}(1+\tau)^{(2-\gamma)/p+(1-\theta(2p))/2-(1-1/p)(n/4+1/2)}W(\tau)\right)^p.\end{aligned} \tag{3.19}$$

Set $\beta = n/4p + (2-\gamma)/p + \nu$, $\nu > 0$, so if we compute the exponent of $(\tau+1)$ in the right side of (3.19), we obtain

$$\beta + (1-\theta(2p))/2 - (1-\delta)(n/4+1/2) = [\nu + \delta(n/4+1/2)] - n(p-1-2(2-\gamma)/n)/(2p) \tag{3.20}$$

and

$$(2-\gamma)/p + (1-\theta(2p))/2 - (1-1/p)(n/4+1/2) = -n(p-1-1/n-2(2-\gamma)/n)/(2p). \tag{3.21}$$

So, as $\gamma < 1/2 < (7-\sqrt{9+4n})/4$ implies $p > p_\gamma > 1+1/n+2(2-\gamma)/n$ for $1 \le n \le 3$, and as $p > p^* > 1+2(2-\gamma)/n$, we can deduce, choosing ν small enough, that the quantity in (3.20)-(3.21) are negative. Hence, we can rewrite (3.19) like

$$\max_{[0,t]} W(\tau) \le CI_0 + C(\max_{[0,t]} W(\tau))^p + C(\max_{[0,t]} W(\tau))^{p+1}. \tag{3.22}$$

Now, write $I_0 = \|u_0\|_{H^1} + \|u_1\|_2 = C\varepsilon$, for small $\varepsilon > 0$ which is determined later, and put

$$T^* = \sup\{t \ge 0 :\ W(t) \le 2C\varepsilon\}.$$

Then, (3.22) implies $W(t) \le C\varepsilon + C\varepsilon^p$. Therefore, taking small ε such that $C\varepsilon + C\varepsilon^p + C\varepsilon^{p+1} < 2C\varepsilon$ we conclude that $T^* = \infty$ (For details we refer the reader to [9, Proposition 2.1] and [18, Proposition 2.1]), i.e.

$$W(t) = (1+t)^{n/4+1/2}\|Du(t,\cdot)\|_2 + \|e^{\psi(t,\cdot)}Du(t,\cdot)\|_2 \le C\varepsilon, \quad t \ge 0. \tag{3.23}$$

Thus we have completed the proof of Theorem 1. □

Proof of Theorem 3. The estimate (1.8) follows directly from (3.23). Next, it follows from inequality (3.11) and estimate (3.23) that

$$C\varepsilon \geq \|e^{\psi(t,\cdot)}Du(t,\cdot)\|_{L^2(\mathbb{R}^n)} \geq \|e^{|\cdot|^2/4(t+K)}Du(t,\cdot)\|_{L^2(\mathbb{R}^n\setminus B(t^{1/2+\delta}))} \geq e^{t^{1+2\delta}/4(t+K)}\|Du(t,\cdot)\|_{L^2(\mathbb{R}^n\setminus B(t^{1/2+\delta}))},$$

which implies (1.7). □

To show Lemma 1, we need a linear estimates for the fundamental solution of the following linear damped wave equation

$$w_{tt} - \Delta w + w_t = 0, \quad w(0,x) = u_0(x), \quad w_t(0,x) = u_1(x), \tag{3.24}$$

for $t \in (0,\infty)\times\mathbb{R}^n$. Let $K_0(t), K_1(t)$ be

$$K_0(t) := e^{-\frac{t}{2}}\cos\{ta(|\nabla|)\}, \qquad K_1(t) := e^{-\frac{t}{2}}\frac{\sin\{ta(|\nabla|)\}}{a(|\nabla|)},$$

where

$$\mathcal{F}[a(|\nabla|)](\xi) = a(\xi) = \begin{cases} \sqrt{|\xi|^2 - 1/4}, & |\xi| > 1/2, \\ \\ i\sqrt{1/4 - |\xi|^2}, & |\xi| \leq 1/2. \end{cases}$$

Note that $K_0(t) + 1/2K_1(t) = \partial_t K_1(t)$. Then the solution of (3.24) is given (cf. [13]) through the Fourier transform by $K_0(t)$ and $K_1(t)$ as

$$w(t,x) = K_0(t) * u_0 + K_1(t) * \left(\frac{1}{2}u_0 + u_1\right). \tag{3.25}$$

The Duhamel principle implies that the solution $u(t,x)$ of nonlinear equation (1.1) solves the integral equation

$$u(t,x) = w(t,x) + \Gamma(\alpha)\int_0^t K_1(t-\tau) * J^{\alpha}_{0|\tau}(|u|^p)(\tau)\,d\tau, \tag{3.26}$$

where $\alpha := 1-\gamma$ and $J^{\alpha}_{0|t}$ is given by (2.2). We can now state Matsumura's result, on the estimate of $K_0(t)$ and $K_1(t)$, as follows:

Proposition 2. *([13]) If $f \in L^m(\mathbb{R}^n) \cap H^{k+|\nu|-1}(\mathbb{R}^n)$ $(1 \leq m \leq 2)$, then*

$$\|\partial_t^k \nabla_x^{\nu} K_1(t) * f\|_2 \leq C(1+t)^{n/4-n/(2m)-|\nu|/2-k}(\|f\|_m + \|f\|_{H^{k+|\nu|-1}(\mathbb{R}^n)}).$$

Proof of Lemma 1. We begin to estimate the linear term $\|Dw(t,\cdot)\|_2$. It is not difficult to see, using Proposition 2 with $m = 1$, that

$$\begin{aligned} \|Dw(t,\cdot)\|_2 &\leq C(1+t)^{-n/4-1/2}(\|u_0\|_{H^1} + \|u_0\|_1 + \|u_1\|_2 + \|u_1\|_1) \\ &\leq CI_0(1+t)^{-n/4-1/2}. \end{aligned} \tag{3.27}$$

For the nonlinear term in (3.26), we apply again Proposition 2 with $m = 1$, and obtain

$$\begin{aligned} \int_0^t \|DK_1(t-\tau) * J^{\alpha}_{0|\tau}(|u|^p)(\tau)\|_2\,ds &\leq \int_0^t (t-\tau+1)^{-n/4-1/2}\left(\|J^{\alpha}_{0|\tau}(|u|^p)(\tau)\|_1 + \|J^{\alpha}_{0|\tau}(|u|^p)(\tau)\|_2\right)d\tau \\ &\leq \int_0^t (t-\tau+1)^{-n/4-1/2}\left(J^{\alpha}_{0|\tau}\|u(\tau)\|_p^p + J^{\alpha}_{0|\tau}\|u(\tau)\|_{2p}^p\right)d\tau. \end{aligned} \tag{3.28}$$

To transform the L^p-norm into a weighted L^{2p}-norm, we use the Cauchy inequality

$$\begin{aligned} \|u(\tau,\cdot)\|_p^p &\equiv \int_{B(\tau+K)} |u(\tau,x)|^p\,dx \\ &\leq \left(\int_{B(\tau+K)} e^{-2p\delta\psi(\tau,x)}\,dx\right)^{1/2}\left(\int_{B(\tau+K)} e^{2p\delta\psi(\tau,x)}|u(\tau,x)|^{2p}\,dx\right)^{1/2}, \end{aligned}$$

for $\delta > 0$. From (3.11), we have $\psi(\tau, x) \geq |x|^2/4(\tau + K)$ for $x \in B(\tau + K)$, so the first integral is estimated as follows

$$\int_{B(\tau+K)} e^{-2p\delta\psi(\tau,x)}\,dx \leq \int_{B(\tau+K)} e^{-p\delta|x|^2/2(\tau+k)}\,dx \leq \int_{\mathbb{R}^n} e^{-p\delta|x|^2/2(\tau+k)}\,dx \equiv \left(\frac{2\pi}{p\delta}\right)^{n/2}(\tau+K)^{n/2}.$$

Thus, for the norm $\|u(\tau,\cdot)\|_p$ in (3.28) we obtain the weighted estimate

$$\|u(\tau,\cdot)\|_p^p \leq C_{K,\delta}(\tau+1)^{n/4}\|e^{\delta\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p}^p, \quad \delta > 0. \tag{3.29}$$

Next, as $\psi > 0$, the norm $\|u(\tau,\cdot)\|_{2p}$ in (3.28) can obviously be estimated by

$$\|u(\tau,\cdot)\|_{2p}^p \leq C_{\delta}(\tau+1)^{n/4}\|e^{\delta\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p}^p. \tag{3.30}$$

It follows from (3.28)-(3.30) that

$$\begin{aligned}
&\int_0^t \|DK_1(t-\tau) * J_{0|\tau}^{\alpha}(|u|^p)(\tau)\|_2\,ds \\
&\leq C\int_0^t (t-\tau+1)^{-n/4-1/2}\int_0^{\tau}(\tau-\sigma)^{-\gamma}\left((\sigma+1)^{n/(4p)}\|e^{\delta\psi(\sigma,\cdot)}u(\sigma,\cdot)\|_{2p}\right)^p\,d\sigma\,d\tau \\
&\leq C\left(\max_{[0,t]}(\tau+1)^{\beta}\|e^{\delta\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p}\right)^p\int_0^t (t-\tau+1)^{-n/4-1/2}\int_0^{\tau}(\tau-\sigma)^{-\gamma}(\sigma+1)^{-(2-\gamma)}\,d\sigma\,d\tau \qquad (3.31)\\
&\qquad (3.32)
\end{aligned}$$

for any $\beta > n/(4p) + (2-\gamma)/p$. Since

$$\begin{aligned}
&\int_0^t (t-\tau+1)^{-n/4-1/2}\int_0^{\tau}(\tau-\sigma)^{-\gamma}(\sigma+1)^{-(2-\gamma)}\,d\sigma\,d\tau \\
&\leq \int_0^t (t-\tau+1)^{-n/4-1/2}\left(\int_0^{\tau}(\tau-\sigma)^{-2\gamma}\,d\sigma\right)^{1/2}\left(\int_0^{\tau}(\sigma+1)^{-2(2-\gamma)}\,d\sigma\right)^{1/2}\,d\tau \\
&\leq C(1+t)^{1/2-\gamma}\int_0^t (t-\tau+1)^{-n/4-1/2}(1+\tau)^{1/2-(2-\gamma)}\,d\tau \\
&\leq C(1+t)^{1/2-\gamma}\left(\int_0^t (t-\tau+1)^{-n/2-1}\,d\tau\right)^{1/2}\left(\int_0^t (1+\tau)^{-3+2\gamma}\,d\tau\right)^{1/2} = C(1+t)^{-n/4-1/2}, \qquad (3.33)
\end{aligned}$$

we find

$$\int_0^t \|DK_1(t-\tau) * J_{0|\tau}^{\alpha}(|u|^p)(\tau)\|_2\,ds \leq C(1+t)^{-n/4-1/2}\left(\max_{[0,t]}(\tau+1)^{\beta}\|e^{\delta\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p}\right)^p. \tag{3.34}$$

Finally, (3.26), (3.27) and (3.34) imply the desired estimate. □

Proof of Lemma 3. The proof of this lemma will be done along the line as in [20, Proposition 2.3] except for some changing on the estimate of the nonlinear part. Indeed, let us assume, for the moment, that $u \in C^2([0, T_{\max}), L^2(\mathbb{R}^n))$. Multiplying (1.1) by $e^{2\psi}u_t$, using (3.8) and (3.9), we get

$$\frac{d}{dt}\left(\frac{e^{2\psi(x,t)}}{2}(|u_t|^2 + |\nabla u|^2)\right) - \operatorname{div}(e^{2\psi(x,t)}u_t\nabla u) - \frac{e^{2\psi(x,t)}}{\psi_t}(\psi_t\nabla u - u_t\nabla\psi)^2 = \Gamma(\alpha)J_{0|t}^{\alpha}(|u|^p)e^{2\psi}u_t.$$

Integrating over $[0,t]\times\mathbb{R}^n$, we have

$$\begin{aligned}
\|e^{\psi(t,\cdot)}Du(t,\cdot)\|_2 &\leq CI_0 + C\int_0^t\int_{\mathbb{R}^n} J_{0|\tau}^{\alpha}(|u|^p)(\tau)e^{2\psi(\tau,x)}u_t(\tau,x)\,dx\,d\tau \\
&\leq CI_0 + C\int_0^t \|J_{0|\tau}^{\alpha}(|u|^p)(\tau)e^{\psi(\tau,\cdot)}\|_2 \times \|e^{\psi(\tau,\cdot)}u_t(\tau,\cdot)\|_2\,d\tau \\
&\leq CI_0 + C\max_{[0,t]}W(\tau)\int_0^t \|J_{0|\tau}^{\alpha}(|u|^p)(\tau)e^{\psi(\tau,\cdot)}\|_2\,d\tau. \qquad (3.35)
\end{aligned}$$

Actually, the above estimate holds for any solution from the energy space by the density argument.

Next, let us estimate the integral term in the right side of (3.35). Since $\psi_t < 0$, we conclude from the definition of $J^\alpha_{0|t}$ that

$$\begin{aligned}\int_0^t \|J^\alpha_{0|\tau}(|u|^p)(\tau)e^{\psi(\tau,\cdot)}\|_2\,d\tau &\leq \int_0^t\int_0^\tau (\tau-\sigma)^{-\gamma}\|e^{\psi(\sigma,\cdot)}|u(\sigma,\cdot)|^p\|_2\,d\sigma\,d\tau\\ &= \int_0^t\int_0^\tau (\tau-\sigma)^{-\gamma}(\sigma+1)^{-(2-\gamma)}\left((\sigma+1)^{(2-\gamma)/p}\|e^{\frac{1}{p}\psi(\sigma,\cdot)}u(\sigma,\cdot)\|_{2p}\right)^p\,d\sigma\,d\tau\\ &\leq \left(\max_{[0,t]}(1+\tau)^{(2-\gamma)/p}\|e^{\frac{1}{p}\psi(\tau,\cdot)}u(\tau,\cdot)\|_{2p}\right)^p \int_0^t\int_0^\tau (\tau-\sigma)^{-\gamma}(\sigma+1)^{-(2-\gamma)}\,d\sigma\,d\tau. \quad (3.36)\end{aligned}$$

Moreover, it follows from Hölder's inequality that

$$\begin{aligned}\int_0^t\int_0^\tau (\tau-\sigma)^{-\gamma}(\sigma+1)^{-(2-\gamma)}\,d\sigma\,d\tau &\leq \int_0^t\left(\int_0^\tau (\tau-\sigma)^{-2\gamma}\,d\sigma\right)^{1/2}\left(\int_0^\tau (\sigma+1)^{-2(2-\gamma)}\,d\sigma\right)^{1/2}d\tau\\ &\leq C(1+t)^{1/2-\gamma}\int_0^t(\tau+1)^{1/2-(2-\gamma)}\,d\tau = C. \quad (3.37)\end{aligned}$$

Finally (3.35)-(3.37) imply (3.15), which completes the proof of Lemma 3.

4. Blow-up result

In this section we devote ourselves to the proof of Theorem 2. We start by introducing the definition of the weak solution of (1.1).

Definition 1. (*Weak solution*) *Let $T > 0$, $\gamma \in (0,1)$ and $u_0, u_1 \in L^1_{loc}(\mathbb{R}^n)$. We say that u is a weak solution if $u \in L^p((0,T), L^p_{loc}(\mathbb{R}^n))$ and satisfies*

$$\begin{aligned}&\Gamma(\alpha)\int_0^T\int_{\mathbb{R}^n} J^\alpha_{0|t}(|u|^p)\varphi\,dx\,dt + \int_{\mathbb{R}^n} u_1(x)\varphi(0,x)\,dx + \int_{\mathbb{R}^n} u_0(x)(\varphi(0,x)-\varphi_t(0,x))\,dx\\ &= \int_0^T\int_{\mathbb{R}^n} u\varphi_{tt}\,dx\,dt - \int_0^T\int_{\mathbb{R}^n} u\varphi_t\,dx\,dt - \int_0^T\int_{\mathbb{R}^n} u\Delta\varphi\,dx\,dt, \quad (4.1)\end{aligned}$$

for all compactly supported function $\varphi \in C^2([0,T]\times\mathbb{R}^n)$ such that $\varphi(\cdot\,,T) = 0$ and $\varphi_t(\cdot\,,T) = 0$.

Next, the following lemma is useful for the proof of Theorem 2. The proof of this lemma is much the same procedure as in the proof of [3, Lemma 2].

Lemma 4. (*Mild → Weak*) *Let $T > 0$ and $\gamma \in (0,1)$. Suppose that $1 < p \leq n/(n-2)$, if $n \geq 3$, and $p \in (1,\infty)$, if $n = 1, 2$. If $u \in C([0,T], H^1(\mathbb{R}^n)) \cap C^1([0,T], L^2(\mathbb{R}^n))$ is the mild solution of* (1.1)*, then u is a weak solution of* (1.1)*.*

Remark. We need the mild solution to use, in the proof of Theorem 2, the alternative (1.5). Without this properties, we obtain just a nonexistence of global solution and not a blow-up result.

Proof of Theorem 2. We assume on the contrary, using (1.5), that u is a global mild solution of (1.1). So, from Lemma 4 we have

$$\begin{aligned}&\Gamma(\alpha)\int_0^T\int_{\mathrm{supp}\varphi} J^\alpha_{0|t}(|u|^p)\,\varphi\,dx\,dt + \int_{\mathrm{supp}\varphi} u_1(x)\varphi(0,x)\,dx + \int_{\mathrm{supp}\varphi} u_0(x)(\varphi(0,x)-\varphi_t(0,x))\,dx\\ &= \int_0^T\int_{\mathrm{supp}\varphi} u\,\varphi_{tt}\,dx\,dt - \int_0^T\int_{\mathrm{supp}\varphi} u\,\varphi_t\,dx\,dt - \int_0^T\int_{\mathrm{supp}\Delta\varphi} u\,\Delta\varphi\,dx\,dt, \quad (4.2)\end{aligned}$$

for all $T > 0$ and all compactly supported test function $\varphi \in C^2([0,T]\times\mathbb{R}^n)$ such that $\varphi(\cdot\,,T) = 0$ and $\varphi_t(\cdot\,,T) = 0$. Let $\varphi(x,t) = D^{\alpha}_{t|T}\left(\tilde{\varphi}(x,t)\right) := D^{\alpha}_{t|T}\left(\varphi_1^{\ell}(x)\varphi_2(t)\right)$ with $\varphi_1(x) := \Phi\left(|x|/B\right)$, $\varphi_2(t) := (1-t/T)_+^{\eta}$, where $D^{\alpha}_{t|T}$ is given by (2.1), $\ell, \eta \gg 1$ and $\Phi \in C^{\infty}(\mathbb{R}_+)$ be a cut-off non-increasing function such that

$$\Phi(r) = \begin{cases} 1 & \text{if } 0 \le r \le 1 \\ 0 & \text{if } r \ge 2, \end{cases}$$

and $0 \le \Phi \le 1$, $|\Phi'(r)| \le C_1/r$ for all $r > 0$. The constant $B > 0$ in the definition of φ_1 is fixed and will be chosen later. In the following, we denote by $\Omega(B)$ the support of φ_1 and by $\Delta(B)$ the set containing the support of $\Delta\varphi_1$ which are defined as follows:

$$\Omega(B) = \{x \in \mathbb{R}^n :\ |x| \le 2B\}, \quad \Delta(B) = \{x \in \mathbb{R}^n :\ B \le |x| \le 2B\}.$$

We return to (4.2), which actually reads

$$\begin{aligned} &\Gamma(\alpha)\int_0^T\int_{\Omega(B)} J^{\alpha}_{0|t}(|u|^p) D^{\nu}_{t|T}\tilde{\varphi}\,dx\,dt + \int_{\Omega(B)} u_1(x) D^{\alpha}_{t|T}\tilde{\varphi}(0,x)\,dx + \int_{\Omega(B)} u_0(x)(D^{\alpha}_{t|T}\tilde{\varphi}(0,x) - \partial_t D^{\alpha}_{t|T}\tilde{\varphi}(0,x))\,dx \\ &= \int_0^T\int_{\Omega(B)} u\,\partial_t^2 D^{\alpha}_{t|T}\tilde{\varphi}\,dx\,dt - \int_0^T\int_{\Omega(B)} u\,\partial_t D^{\alpha}_{t|T}\tilde{\varphi}\,dx\,dt - \int_0^T\int_{\Delta(B)} u\,\Delta D^{\alpha}_{t|T}\tilde{\varphi}\,dx\,dt. \end{aligned} \tag{4.3}$$

From (2.3), (2.4) and (2.7), we conclude that

$$\begin{aligned} &\int_0^T\int_{\Omega(B)} D^{\alpha}_{0|t} J^{\alpha}_{0|t}(|u|^p)\tilde{\varphi}\,dx\,dt + C\,T^{-\alpha}\int_{\Omega(B)} u_1(x)\varphi_1^{\ell}(x)\,dx + C(T^{-\alpha} + T^{-\alpha-1})\int_{\Omega(B)} u_0(x)\varphi_1^{\ell}(x)\,dx \\ &= C\int_0^T\int_{\Omega(B)} u(D^{2+\alpha}_{t|T}\tilde{\varphi} + D^{1+\alpha}_{t|T}\tilde{\varphi})\,dx\,dt - C\int_0^T\int_{\Delta(B)} u\,\Delta(\varphi_1^{\ell}) D^{\alpha}_{t|T}\varphi_2\,dx\,dt, \end{aligned} \tag{4.4}$$

where $D^{\alpha}_{0|t}$ is defined in (2.1). Moreover, using (2.5) and the fact that (1.6) implies $\int_{\Omega(B)}\varphi_1^{\ell}(x)u_i(x) \ge 0, i = 0,1$, it follows

$$\begin{aligned} \int_0^T\int_{\Omega(B)} |u|^p\tilde{\varphi}\,dx\,dt &\le C\int_0^T\int_{\Omega(B)} |u|\varphi_1^{\ell}(D^{2+\alpha}_{t|T}\varphi_2 + D^{1+\alpha}_{t|T}\varphi_2)\,dx\,dt \\ &\quad + C\int_0^T\int_{\Delta(B)} |u|\varphi_1^{\ell-2}(|\Delta\varphi_1| + |\nabla\varphi_1|^2) D^{\alpha}_{t|T}\varphi_2\,dx\,dt \\ &=: I_1 + I_2, \end{aligned} \tag{4.5}$$

where we have used the formula $\Delta(\varphi_1^{\ell}) = \ell\varphi_1^{\ell-1}\Delta\varphi_1 + \ell(\ell-1)\varphi_1^{\ell-2}|\nabla\varphi_1|^2$ and $\varphi_1 \le 1$. Next we observe that by introducing the term $\tilde{\varphi}^{1/p}\tilde{\varphi}^{-1/p}$ in the right side of (4.5) and applying Young's inequality we have

$$I_1 \le \frac{1}{2p}\int_0^T\int_{\Omega(B)} |u|^p\tilde{\varphi}\,dx\,dt + C\int_0^T\int_{\Omega(B)} \varphi_1^{\ell}\varphi_2^{-1/(p-1)}((D^{2+\alpha}_{t|T}\varphi_2)^{p'} + (D^{1+\alpha}_{t|T}\varphi_2)^{p'})\,dx\,dt, \tag{4.6}$$

where $p' = p/(p-1)$. Similarly,

$$I_2 \le \frac{1}{2p}\int_0^T\int_{\Omega(B)} |u|^p\tilde{\varphi}\,dx\,dt + C\int_0^T\int_{\Omega(B)} \varphi_1^{\ell-2p'}\varphi_2^{-1/(p-1)}(|\Delta\varphi_1|^{p'} + |\nabla\varphi_1|^{2p'})(D^{\alpha}_{t|T}\varphi_2)^{p'}\,dx\,dt. \tag{4.7}$$

Combining (4.6) and (4.7), it follows from (4.5) that

$$\begin{aligned} \int_0^T\int_{\Omega(B)} |u|^p\tilde{\varphi}\,dx\,dt &\le C\int_0^T\int_{\Omega(B)} \varphi_1^{\ell}\varphi_2^{-1/(p-1)}((D^{2+\alpha}_{t|T}\varphi_2)^{p'} + (D^{1+\alpha}_{t|T}\varphi_2)^{p'})\,dx\,dt \\ &\quad + C\int_0^T\int_{\Omega(B)} \varphi_1^{\ell-2p'}\varphi_2^{-1/(p-1)}(|\Delta\varphi_1|^{p'} + |\nabla\varphi_1|^{2p'})(D^{\alpha}_{t|T}\varphi_2)^{p'}\,dx\,dt. \end{aligned} \tag{4.8}$$

At this stage, for the choice of B, we have to distinguishes 3 cases.

• Case of $p < p_\gamma$: in this case, we take $B = T^{1/2}$. So, using (2.6) and the change of variables: $s = T^{-1}t$, $y = T^{-1/2}x$, we get from (4.8) that

$$\int_0^T \int_{\Omega(T^{1/2})} |u|^p \tilde{\varphi}\, dx\, dt \le C(T^{-(\alpha+2)p'+n/2+1} + T^{-(\alpha+1)p'+n/2+1}), \tag{4.9}$$

where C is independent of T. Letting $T \to \infty$ in (4.9), thanks to $p < p_\gamma$ and the Lebesgue dominated convergence theorem, it is yielded that

$$\int_0^\infty \int_{\mathbb{R}^n} |u|^p\, dx\, dt = 0,$$

which implies $u(x,t) = 0$ for all t and a.e. x. This contradicts our assumption (1.6).

• Case of $p = p_\gamma$: Let $B = R^{-1/2}T^{1/2}$, where $1 \ll R < T$ is such that when $T \to \infty$ we don't have $R \to \infty$ at the same time. Moreover, from the last case and the fact that $p = p_\gamma$, there exist a positive constant D independent of T such that

$$\int_0^\infty \int_{\mathbb{R}^n} |u|^p\, dx\, dt \le D,$$

which implies that

$$\int_0^T \int_{\Delta(R^{-1/2}T^{1/2})} |u|^p \tilde{\varphi}\, dx\, dt \to 0 \quad \text{as} \quad T \to \infty. \tag{4.10}$$

On the other hand, using Hölder's inequality instead of Young's one, we estimate the integral I_2 in (4.5) as follows:

$$I_2 \le C\left(\int_0^T \int_{\Delta(R^{-1/2}T^{1/2})} |u|^p \tilde{\varphi}\right)^{1/p} \left(\int_0^T \int_{\Omega(R^{-1/2}T^{1/2})} \varphi_1^{\ell-2p'} \varphi_2^{-1/(p-1)} (|\Delta\varphi_1|^{p'} + |\nabla\varphi_1|^{2p'})(D_{t|T}^{\alpha}\varphi_2)^{p'}\, dx\, dt\right)^{1/p'}. \tag{4.11}$$

Similarly to the last case, substituting (4.6) and (4.11) into (4.5), taking account of $p = p_\gamma$ and the scaled variable $s = T^{-1}t$, $y = R^{1/2}T^{-1/2}x$, we get

$$\int_0^T \int_{\Omega(R^{-1/2}T^{1/2})} |u|^p\, dx\, dt \le C(T^{-p'}R^{-n/2} + R^{-n/2}) + CR^{1-n/(2p')}\left(\int_0^T \int_{\Delta(R^{-1/2}T^{1/2})} |u|^p \tilde{\varphi}\right)^{1/p}.$$

Letting $T \to \infty$, using (4.10), we get

$$\int_0^\infty \int_{\mathbb{R}^N} |u|^p\, dx\, dt \le CR^{-n/2},$$

which implies a contradiction, when $R \to \infty$, with (1.6).

• Case of $p \le (1/\gamma)$: Let $B = R$ with the same R introduced in the case $p = p_\gamma$. Then, by taking the new variables $s = T^{-1}t$, $y = R^{-1}x$, it follows from (4.8) that

$$\int_0^T \int_{\Omega(R)} |u|^p \tilde{\varphi}\, dx\, dt \le CR^n(T^{-(2+\alpha)p'+1} + T^{-(1+\alpha)p'+1}) + CR^{n-2p'}T^{-\alpha p'+1}.$$

Finally, we get a contradiction with (1.6) by letting the following limits: first $T \to \infty$, next $R \to \infty$. Precisely, in the case of $p = 1/\gamma$, we have to use the condition $\gamma > (n-2)/n$ to obtain the desired convergence. This completes the proof of Theorem 2.

Appendix

In this appendix let us sketch the proof of Proposition 1. Let us define a semigroup $S(t) : H^1(\mathbb{R}^n) \times L^2(\mathbb{R}^n) \to H^1(\mathbb{R}^n) \times L^2(\mathbb{R}^n)$ by

$$S(t) : \begin{bmatrix} u_0 \\ u_1 \end{bmatrix} \mapsto \begin{bmatrix} w \\ w_t \end{bmatrix},$$

where $w \in C([0,\infty), H^1(\mathbb{R}^n)) \cap C^1([0,\infty), L^2(\mathbb{R}^n))$ is the linear solution of (3.24) given by (3.25). So, view of (3.26), a mild solution of the nonlinear problem (1.1) is equivalent to following integral equation:

$$U(t) = S(t)U_0 + \int_0^t S(t-s)F(s)\,ds, \tag{4.12}$$

where

$$U(t) = \begin{bmatrix} u(t,\cdot) \\ u_t(t,\cdot) \end{bmatrix}, \quad U_0 = \begin{bmatrix} u_0 \\ u_1 \end{bmatrix}, \quad F(s) = \begin{bmatrix} 0 \\ J^{\alpha}_{0|s}(|u|^p)(s) \end{bmatrix}.$$

It sufficient now to prove the local existence of a solution of (4.12) in $H^1(\mathbb{R}^n) \times L^2(\mathbb{R}^n)$. Let $T > 0$ and consider the following Banach space

$$E := \{U = {}^t(u,v) : \ (u,v) \in C([0,T], H^1(\mathbb{R}^n) \times L^2(\mathbb{R}^n)), \ \text{supp}u(t,\cdot) \subset B(K+t) \text{ and } \|U\|_E \le CM\},$$

where

$$\|U\|_E := \|u\|_{C([0,T];H^1(\mathbb{R}^n))} + \|v\|_{C([0,T];L^2(\mathbb{R}^n))} \quad \text{and} \quad M := \|u_0\|_{H^1} + \|u_1\|_2.$$

In order to use the Banach fixed point theorem, we introduce the following map Φ on E defined by

$$\Phi[U](t) := S(t)U_0 + \int_0^t S(t-s)F(s)\,ds.$$

Now, for $U = (u,v) \in E$, we have

$$\|J^{\alpha}_{0|t}(|u|^p)(t)\|_2 \le Ct^{1-\gamma}\|u(t,\cdot)\|^p_{2p} \le Ct^{1-\gamma}\|u(t,\cdot)\|^p_{H^1} \le Ct^{1-\gamma}\|U\|^p_E, \quad t \in [0,T],$$

where we have used the Sobolev imbedding $H^1(\mathbb{R}^n) \subset L^{2p}(\mathbb{R}^n)$. Next, using Matsumura's result (Proposition 2) with $m = 2$ and the finite propagation speed phenomena, we deduce via the Banach fixed point theorem that there exists a local solution $U \in E$ on a small interval $[0,T]$ satisfies (4.12). For details, we refer the reader to [5, Theorem 1] and [3, Theorem 2]. By consequence, there exist a local solution $u \in C([0,T], H^1(\mathbb{R}^n)) \cap C^1([0,T], L^2(\mathbb{R}^n))$ satisfies (3.26) and $\text{supp}u(t,\cdot) \subset B(t+K)$. However, since our equation (1.1) is nonautonomous, we prefer apply Gronwall's inequality to get the uniqueness (cf. [2, Theorem 3.1]). Indeed, if $u, v \in C([0,T], H^1(\mathbb{R}^n)) \cap C^1([0,T], L^2(\mathbb{R}^n))$ are two mild solutions (i.e. satisfy (3.26)) for some $T > 0$, we have

$$\begin{aligned} \|u(t) - v(t)\|_{H^1} &\le C\int_0^t \|K_1(t-\tau) * J^{\alpha}_{0|\tau}(|u|^p - |v|^p)(\tau)\|_{H^1}\,d\tau \\ &\le C\int_0^t (1+t-\tau)^{-1/2}\|J^{\alpha}_{0|\tau}(|u|^p - |v|^p)(\tau)\|_2\,d\tau \\ &\le C\int_0^t \|J^{\alpha}_{0|\tau}(|u|^p - |v|^p)(\tau)\|_2\,d\tau, \end{aligned} \tag{4.13}$$

where we have used again Matsumura's result (Proposition 2) with $m = 2$. As $||u|^p - |v|^p| \le C|u-v|(|u|^p + |v|^p)$, so by Hölder's inequality ($\|ab\|_2 \le \|a\|_{2p}\|b\|_{2p'}$) with $p' = p/(p-1)$ and Sobolev's imbedding ($H^1 \subset L^{2p}$), we obtain

$$\begin{aligned} \int_0^t \|J^{\alpha}_{0|\tau}(|u|^p - |v|^p)(\tau)\|_2\,d\tau &\le C\int_0^t J^{\alpha}_{0|\tau}(\|u-v\|_{H^1}(\|u\|^{p-1}_{H^1} + \|v\|^{p-1}_{H^1}))(\tau)\,d\tau \\ &\le C\int_0^t\int_0^{\tau} (\tau-s)^{-\gamma}\|u(s,\cdot) - v(s,\cdot)\|_{H^1}\,ds\,d\tau \end{aligned}$$

$$= C \int_0^t \int_s^t (\tau - s)^{-\gamma} \|u(s,\cdot) - v(s,\cdot)\|_{H^1} \, d\tau \, ds$$
$$= C \int_0^t (t-s)^{1-\gamma} \|u(s,\cdot) - v(s,\cdot)\|_{H^1} \, ds. \tag{4.14}$$

Combining (4.13) and (4.14), we get

$$\|u(t) - v(t)\|_{H^1} \leq C \int_0^t (t-s)^{1-\gamma} \|u(s,\cdot) - v(s,\cdot)\|_{H^1} \, ds.$$

Finally, using a variation type of Gronwall's inequality, it follows that $u(t) \equiv v(t)$. As a consequence of this uniqueness result, we can extend our solution u on a maximal interval $[0, T_{\max})$. Moreover, if $T_{\max} < \infty$, then $\|u(t,\cdot)\|_{H^1} + \|u_t(t,\cdot)\|_2 \to \infty$ as $t \to T_{\max}$. For details, see [2, Theorem 3.1] or [5, Theorem 1].

Acknowledgements

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