# Decorrelation estimates for the eigenlevels of the discrete Anderson model in the localized regime 

Frédéric Klopp

## To cite this version:

Frédéric Klopp. Decorrelation estimates for the eigenlevels of the discrete Anderson model in the localized regime. 2010. hal-00467011v2

## HAL Id: hal-00467011 <br> https://hal.science/hal-00467011v2

Preprint submitted on 5 Sep 2010

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# DECORRELATION ESTIMATES FOR THE EIGENLEVELS OF THE DISCRETE ANDERSON MODEL IN THE LOCALIZED REGIME 

FRÉDÉRIC KLOPP


#### Abstract

The purpose of the present work is to establish decorrelation estimates for the eigenvalues of the discrete Anderson model localized near two distinct energies inside the localization region. In dimension one, we prove these estimates at all energies. In higher dimensions, the energies are required to be sufficiently far apart from each other. As a consequence of these decorrelation estimates, we obtain the independence of the limits of the local level statistics at two distinct energies.

Résumé. Dans ce travail, nous établissons des inégalités de décorrélation pour les valeurs propres proches de deux énergies distinctes. En dimension 1, nous démontrons que ces inégalités sont vraies quel que soit le choix de ces deux énergies. En dimension supérieure, il nous faut supposer que les deux énergies sont suffisamment éloignées l'une de l'autre. Comme conséquence de ces inégalités de décorrélation, nous démontrons que les limites des statistiques locales des valeurs propres sont indépendantes pour deux énergies distinctes.


## 1. Introduction

On $\ell^{2}\left(\mathbb{Z}^{d}\right)$, consider the random Anderson model

$$
H_{\omega}=-\Delta+V_{\omega}
$$

where $-\Delta$ is the free discrete Laplace operator

$$
\begin{equation*}
(-\Delta u)_{n}=\sum_{|m-n|=1} u_{m} \quad \text { for } u=\left(u_{n}\right)_{n \in \mathbb{Z}^{d}} \in \ell^{2}\left(\mathbb{Z}^{d}\right) \tag{1.1}
\end{equation*}
$$

and $V_{\omega}$ is the random potential

$$
\begin{equation*}
\left(V_{\omega} u\right)_{n}=\omega_{n} u_{n} \quad \text { for } u=\left(u_{n}\right)_{n \in \mathbb{Z}^{d}} \in \ell^{2}\left(\mathbb{Z}^{d}\right) . \tag{1.2}
\end{equation*}
$$

We assume that the random variables $\left(\omega_{n}\right)_{n \in \mathbb{Z}^{d}}$ are independent identically distributed and that their common distribution admits a compactly supported bounded density, say $g$. It is then well known (see e.g. [12]) that

- let $\Sigma:=[-2 d, 2 d]+\operatorname{supp} g$ and $S_{-}$and $S_{+}$be the infimum and supremum of $\Sigma$; for almost every $\omega=\left(\omega_{n}\right)_{n \in \mathbb{Z}^{d}}$, the spectrum of $H_{\omega}$ is equal to $\Sigma$;
- for some $S_{-}<s_{-} \leq s_{+}<S_{+}$, the intervals $I_{-}=\left[S_{-}, s_{-}\right)$and $I_{+}=\left(s_{+}, S_{+}\right]$ are contained in the region of localization for $H_{\omega}$ i.e. the region of $\Sigma$ where the finite volume fractional moment criteria of [1] are verified for restrictions of $H_{\omega}$ to sufficiently large cubes (see also Proposition 2.1). In particular, $I:=I_{-} \cup I_{+}$contains only pure point spectrum associated to exponentially decaying eigenfunctions; for the precise meaning of the region of localization, we refer to section 2.1.2; if the disorder is sufficiently large or if the dimension $d=1$ then, one can pick $I=\Sigma$;

[^0]- there exists a bounded density of states, say $\lambda \mapsto \nu(E)$, such that, for any continuous function $\varphi: \mathbb{R} \rightarrow \mathbb{R}$, one has

$$
\begin{equation*}
\int_{\mathbb{R}} \varphi(E) \nu(E) d E=\mathbb{E}\left(\left\langle\delta_{0}, \varphi\left(H_{\omega}\right) \delta_{0}\right\rangle\right) . \tag{1.3}
\end{equation*}
$$

Here, and in the sequel, $\mathbb{E}(\cdot)$ denotes the expectation with respect to the random parameters, and $\mathbb{P}(\cdot)$ the probability measure they induce.
Let $N$ be the integrated density of states of $H_{\omega}$ i.e. $N$ is the distribution function of the measure $\nu(E) d E$. The function $\nu$ is only defined $E$-almost everywhere. In the sequel, when we speak of $\nu(E)$ for some $E$, we mean that the non decreasing function $N$ is differentiable at $E$ and that $\nu(E)$ is its derivative at $E$.
1.1. The results. For $L \in \mathbb{N}$, let $\Lambda=\Lambda_{L}=[-L, L]^{d}$ be a large box and $N:=\# \Lambda_{L}=$ $(2 L+1)^{d}$ be its cardinality. Let $H_{\omega}(\Lambda)$ be the operator $H_{\omega}$ restricted to $\Lambda$ with periodic boundary conditions. The notation $|\Lambda| \rightarrow+\infty$ is a shorthand for considering $\Lambda=\Lambda_{L}$ in the limit $L \rightarrow+\infty$. Let us denote the eigenvalues of $H_{\omega}(\Lambda)$ ordered increasingly and repeated according to multiplicity by $E_{1}(\omega, \Lambda) \leq E_{2}(\omega, \Lambda) \leq \cdots \leq E_{N}(\omega, \Lambda)$.

Let $E$ be an energy in $I$ such that $\nu(E)>0$. The local level statistics near $E$ is the point process defined by

$$
\begin{equation*}
\Xi(\xi, E, \omega, \Lambda)=\sum_{n=1}^{N} \delta_{\xi_{n}(E, \omega, \Lambda)}(\xi) \tag{1.4}
\end{equation*}
$$

where

$$
\begin{equation*}
\xi_{n}(E, \omega, \Lambda)=|\Lambda| \nu(E)\left(E_{n}(\omega, \Lambda)-E\right), \quad 1 \leq n \leq N . \tag{1.5}
\end{equation*}
$$

One of the most striking results describing the localization regime for the Anderson model is
Theorem 1.1 ([15]). Assume that $E \in I$ be such that $\nu(E)>0$.
When $|\Lambda| \rightarrow+\infty$, the point process $\Xi(\cdot, E, \omega, \Lambda)$ converges weakly to a Poisson process on $\mathbb{R}$ with intensity the Lebesgue measure i.e. for $\left(U_{j}\right)_{1 \leq j \leq J}, U_{j} \subset \mathbb{R}$ bounded measurable and $U_{j^{\prime}} \cap U_{j}=\emptyset$ if $j \neq j^{\prime}$ and $\left(k_{j}\right)_{1 \leq j \leq J} \in \mathbb{N}^{J}$, one has

$$
\mathbb{P}\left(\left\{\omega ;\left\{\begin{array}{c}
\#\left\{j ; \xi_{n}(E, \omega, \Lambda) \in U_{1}\right\}=k_{1} \\
\vdots \\
\vdots \\
\#\left\{j ; \xi_{n}(E, \omega, \Lambda) \in U_{J}\right\}=k_{J}
\end{array}\right\}\right) \underset{\Lambda \rightarrow \mathbb{Z}^{d}}{\rightarrow} \prod_{j=1}^{J} e^{-\left|U_{j}\right|} \frac{\left|U_{j}\right|^{k_{j}}}{k_{j}!} .\right.
$$

An analogue of Theorem 1.1 was first proved in [17] for a different one-dimensional random operator.
Once Theorem 1.1 is known, a natural question arises:

- for $E \neq E^{\prime}$, are the limits of $\Xi(\xi, E, \omega, \Lambda)$ and $\Xi\left(\xi, E^{\prime}, \omega, \Lambda\right)$ stochastically independent?
This question has arisen and has been answered for other types of random operators like random matrices (see e.g. [14]); in this case, the local statistics are not Poissonian.
For the Anderson model, this question has been open (see e.g. $[16,19]$ ) and to the best of our knowledge, the present paper is the first to bring an answer. The conjecture is also open for the continuous Anderson model and random CMV matrices where the local statistics have also been proved to be Poissonian (see e.g. [4, 7, 19, 20]).
The main result of the present paper is

Theorem 1.2. Assume that the dimension $d=1$. Pick $E \in I$ and $E^{\prime} \in I$ such that $E \neq E^{\prime}$, $\nu(E)>0$ and $\nu\left(E^{\prime}\right)>0$.
When $|\Lambda| \rightarrow+\infty$, the point processes $\Xi(E, \omega, \Lambda)$ and $\Xi\left(E^{\prime}, \omega, \Lambda\right)$, defined in (1.4), converge weakly respectively to two independent Poisson processes on $\mathbb{R}$ with intensity the Lebesgue measure. That is, for $\left(U_{j}\right)_{1 \leq j \leq J}, U_{j} \subset \mathbb{R}$ bounded measurable and $U_{j^{\prime}} \cap U_{j}=\emptyset$ if $j \neq j^{\prime}$ and $\left(k_{j}\right)_{1 \leq j \leq J} \in \mathbb{N}^{J}$ and $\left(U_{j}^{\prime}\right)_{1 \leq j \leq J^{\prime}}, U_{j}^{\prime} \subset \mathbb{R}$ bounded measurable and $U_{j^{\prime}}^{\prime} \cap U_{j}^{\prime}=\emptyset$ if $j \neq j^{\prime}$ and $\left(k_{j}^{\prime}\right)_{1 \leq j \leq J} \in \mathbb{N}^{J^{\prime}}$ one has

When $d \geq 2$, we also prove
Theorem 1.3. Assume that $d$ is arbitrary. Pick $E \in I$ and $E^{\prime} \in I$ such that $\left|E-E^{\prime}\right|>2 d$, $\nu(E)>0$ and $\nu\left(E^{\prime}\right)>0$.
When $|\Lambda| \rightarrow+\infty$, the point processes $\Xi(E, \omega, \Lambda)$ and $\Xi\left(E^{\prime}, \omega, \Lambda\right)$, defined in (1.4), converge weakly respectively to two independent Poisson processes on $\mathbb{R}$ with intensity the Lebesgue measure.
In section 3, we show that Theorems 1.2 and 1.3 follow from Theorem 1.1 and the decorrelation estimates that we present now. They are the main technical results of the present paper.
Lemma 1.1. Assume $d=1$ and pick $\beta \in(1 / 2,1)$. For $\alpha \in(0,1)$ and $\left\{E, E^{\prime}\right\} \subset I$ s.t. $E \neq E^{\prime}$, for any $c>0$, there exists $C>0$ such that, for $L \geq 3$ and $c L^{\alpha} \leq \ell \leq L^{\alpha} / c$, one has

$$
\mathbb{P}\left(\left\{\begin{array}{l}
\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap\left(E+L^{-d}(-1,1)\right) \neq \emptyset,  \tag{1.7}\\
\left.\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right)\right) \cap\left(E^{\prime}+L^{-d}(-1,1)\right) \neq \emptyset
\end{array}\right\}\right) \leq C(\ell / L)^{2 d} e^{(\log L)^{\beta}} .
$$

This lemma shows that, up to sub-polynomial errors, the probability to obtain simultaneously an eigenvalue near $E$ and another one near $E^{\prime}$ is bounded by the product of the estimates given for each of these events by Wegner's estimate (see section 2.1.1). In this sense, (1.7) is similar to Minami's estimate for two distinct energies.
Lemma 1.1 proves a result conjectured in $[16,19]$ in dimension 1.
In arbitrary dimension, we prove (1.7), actually a somewhat stronger estimate, only when the two energies $E$ and $E^{\prime}$ are sufficiently far apart.
Lemma 1.2. Assume $d$ is arbitrary. Pick $\beta \in(1 / 2,1)$. For $\alpha \in(0,1)$ and $\left\{E, E^{\prime}\right\} \subset I$ s.t. $\left|E-E^{\prime}\right|>2 d$, for any $c>0$, there exists $C>0$ such that, for $L \geq 3$ and $c L^{\alpha} \leq \ell \leq L^{\alpha} / c$, one has

$$
\mathbb{P}\left(\left\{\begin{array}{l}
\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap\left(E+L^{-d}(-1,1)\right) \neq \emptyset,  \tag{1.8}\\
\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap\left(E^{\prime}+L^{-d}(-1,1)\right) \neq \emptyset
\end{array}\right\}\right) \leq C(\ell / L)^{2 d}(\log L)^{C} .
$$

This e.g. proves the independence of the processes for energies in opposite edges of the almost sure spectrum.
The estimate (1.8) in Lemma 1.2 is somewhat stronger than (1.7); one can obtain an analogous estimate in dimension 1 if one restricts oneself to energies $E$ and $E^{\prime}$ such that $E-E^{\prime}$ does not belong to some set of measure 0 (see Lemma 2.11 in Remark 2.2 at the end of section 2.3).

Remark 1.1. As the proof of Theorems 1.2 and 1.3 shows, the estimates (1.7) are (1.8) are stronger than what it needed. It suffices to show that the probabilities in (1.7) are (1.8) are $o\left((\ell / L)^{d}\right)$.
In [7] (see also [8]), the authors provide another proof of Theorems 1.1 and of Theorems 1.2 and 1.3 under the assumption that the probabilities in (1.7) are (1.8) are $o\left((\ell / L)^{d}\right)$. The analysis done in [8] deals with both discrete and continuous models. It yields a stronger version of Theorem 1.1 and Theorems 1.2 and 1.3 in essentially the same step.
Whereas in the proof of Lemma 1.1, we explicitly use the fact that $H_{\omega}=H_{0}+V_{\omega}$ where $H_{0}$ is the free Laplace operator (1.1), the proof we give of Lemma 1.2 still works if $H_{0}$ is any convolution matrix with exponentially decaying off diagonal coefficients if one replaces the condition $\left|E-E^{\prime}\right|>2 d$ with the condition $\left|E-E^{\prime}\right|>\sup \sigma\left(H_{0}\right)-\inf \sigma\left(H_{0}\right)$.

## 2. Proof of the decorrelation estimates

Before starting with the proofs of Lemma 1.1 and 1.2, let us recall additional properties for the discrete Anderson model known to be true under the assumptions we made on the distribution of the random potential.
2.1. Some facts on the discrete Anderson model. Basic estimates on the distribution of the eigenvalues of the Anderson model are the Wegner and Minami estimates.

### 2.1.1. The Wegner and Minami estimates. One has

Theorem 2.1 ([22]). There exists $C>0$ such that, for $J \subset \mathbb{R}$, and $\Lambda$, a cube in $\mathbb{Z}^{d}$, one has

$$
\begin{equation*}
\mathbb{E}\left[\operatorname{tr}\left(\mathbf{1}_{J}\left(H_{\omega}(\Lambda)\right)\right)\right] \leq C|J||\Lambda| \tag{2.1}
\end{equation*}
$$

where

- $\left.H_{\omega}(\Lambda)\right)$ is the operator $H_{\omega}$ restricted to $\Lambda$ with periodic boundary conditions,
- $\mathbf{1}_{J}(H)$ is the spectral projector of the operator $H$ on the energy interval $J$.

We refer to $[10,13,21]$ for simple proofs and more details on the Wegner estimate.
Another crucial estimate is the Minami estimate.
Theorem 2.2 ( $[15,2,9,5])$. There exists $C>0$ such that, for $J \subset K$, and $\Lambda$, a cube in $\mathbb{Z}^{d}$, one has

$$
\begin{equation*}
\mathbb{E}\left[\operatorname{tr}\left(\mathbf{1}_{J}\left(H_{\omega}(\Lambda)\right)\right) \cdot \operatorname{tr}\left(\mathbf{1}_{K}\left(H_{\omega}(\Lambda)\right)-1\right)\right] \leq C|J||K||\Lambda|^{2} . \tag{2.2}
\end{equation*}
$$

For $J=K$, the estimate (2.2) was proved in $[15,2,9,5]$; for $J \neq K$, it can be found in [5]. In their nature, (1.7) or (1.8) and (2.2) are quite similar: the Minami estimate can be interpreted as a decorrelation estimate for close together eigenvalues. It can be used to obtain the counterparts of Theorems 1.2 and 1.3 when $E$ and $E^{\prime}$ tend to each other as $|\Lambda| \rightarrow+\infty$ (see [7]).

Note that the Minami estimate (2.2) has been proved for the discrete Anderson model on intervals $I$ irrelevant of the spectral type of $H_{\omega}$ in $I$. Our proof of the decorrelation estimates (1.7) and (1.8) makes use of the fact that $I$ lies in the localized region.
2.1.2. The localized regime. Let us now give a precise description of what we mean with the region of localization or the localized regime. We prove

Proposition 2.1. Recall that $I=I_{+} \cup I_{-}$is the region of $\Sigma$ where the finite volume fractional moment criteria of [1] for $H_{\omega}(\Lambda)$ are verified for $\Lambda$ sufficiently large.
Then,
(Loc): there exists $\nu>0$ such that, for any $p>0$, there exists $q>0$ and $L_{0}>0$ such that, for $L \geq L_{0}$, with probability larger than $1-L^{-p}$, if
(1) $\varphi_{n, \omega}$ is a normalized eigenvector of $H_{\omega}\left(\Lambda_{L}\right)$ associated to an energy $E_{n, \omega} \in I$,
(2) $x_{n, \omega} \in \Lambda_{L}$ is a maximum of $x \mapsto\left|\varphi_{n, \omega}(x)\right|$ in $\Lambda_{L}$,
then, for $x \in \Lambda_{L}$, one has

$$
\begin{equation*}
\left|\varphi_{n, \omega}(x)\right| \leq L^{q} e^{-\nu\left|x-x_{n, \omega}\right|} \tag{2.3}
\end{equation*}
$$

The point $x_{n, \omega}$ is called a localization center for $\varphi_{n, \omega}$ or $E_{n, \omega}$.
Note that, by Minami's estimate, the eigenvalues of $H_{\omega}(\Lambda)$ are almost surely simple. Thus, we can associate a localization center to an eigenvalue as it is done in Proposition 2.1.
In its spirit, this result is not new (see e.g. [1, 6, 7]). We state it in a form convenient for our purpose. We prove Proposition 2.1 in section 4
2.2. The proof of Lemmas 1.1 and 1.2. The basic idea of the proof is to show that, when $\omega$ varies, two eigenvalues of $H_{\omega}(\Lambda)$ cannot vary in a synchronous manner, or, put in another way, locally in $\omega$, if $E(\omega)$ and $E^{\prime}(\omega)$ denote the two eigenvalues under consideration, for some $\gamma$ and $\gamma^{\prime}$, the mapping $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \mapsto\left(E(\omega), E^{\prime}(\omega)\right)$ is a local diffeomorphism when all the other random variables, that is $\left(\omega_{\alpha}\right)_{\alpha \notin\left\{\gamma, \gamma^{\prime}\right\}}$, are fixed.
As we are in the localized regime, we will exploit this by noting that eigenvalues of $H_{\omega}(\Lambda)$ can only depend significantly of $(\log L)^{d}$ random variables i.e. we can study what happens in cubes that are of side-length $\log L$ while the energy interval where we want to control things are of size $L^{-d}$. This is the essence of Lemma 2.1 below. This lemma is proved under the general assumptions (2.1), (2.2) and (Loc). In particular, it is valid for if one replaces the discrete Laplacian with any convolution matrix with exponentially decaying off diagonal coefficients.
The second step consists in analyzing the mapping $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \mapsto\left(E(\omega), E^{\prime}(\omega)\right)$ on these smaller cubes. The main technical result is Lemma 2.4 that shows that, under the conditions of Lemmas 1.1 and 1.2 , with a large probability, eigenvalues away from each other cannot move synchronously as functions of the random variables. Of course, this will not be correct for all random models: constructing artificial degeneracies, one can easily coin up random models where this is not the case.
Lemmas 1.1 and 1.2 will be proved in essentially the same way; the only difference will be in Lemma 2.4 that controls the joint dependence of two distinct eigenvalues on the random variables.

Let $J_{L}=E+L^{-d}[-1,1]$ and $J_{L}^{\prime}=E^{\prime}+L^{-d}[-1,1]$. Pick $L$ sufficiently large so that $J_{L} \subset I$ and $J_{L}^{\prime} \subset I$ are contained in $I$ where (Loc) holds true.
Pick $c L^{\alpha} \leq \ell \leq L^{\alpha} / c$ where $c>0$ is fixed. By (2.2), we know that

$$
\mathbb{P}\left(\#\left[\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap J_{L}\right] \geq 2 \text { or } \#\left[\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap J_{L}^{\prime}\right] \geq 2\right) \leq C(\ell / L)^{2 d}
$$

where $\#[\cdot]$ denotes the cardinality of $\cdot$.
So if we define

$$
\mathbb{P}_{0}=\mathbb{P}\left(\#\left[\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap J_{L}\right]=1, \#\left[\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap J_{L}^{\prime}\right]=1\right)
$$

it suffices to show that

$$
\mathbb{P}_{0} \leq C(\ell / L)^{2 d} \cdot \begin{cases}e^{(\log L)^{\beta}} & \text { if the dimension } d=1  \tag{2.4}\\ (\log L)^{C} & \text { if the dimension } d>1\end{cases}
$$

First, using the assumption (Loc), we are going to reduce the proof of (2.4) to the proof of a similar estimate where the cube $\Lambda_{\ell}$ will be replaced by a much smaller cube, a cube of side length of order $\log L$. We prove

Lemma 2.1. There exists $C>0$ such that, for $L$ sufficiently large,

$$
\mathbb{P}_{0} \leq C(\ell / L)^{2 d}+C(\ell / \tilde{\ell})^{d} \mathbb{P}_{1}
$$

where $\tilde{\ell}=C \log L$ and

$$
\left.\mathbb{P}_{1}:=\mathbb{P}\left(\#\left[\sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap \tilde{J}_{L}\right] \geq 1\right) \text { and } \#\left[\sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap \tilde{J}_{L}^{\prime}\right] \geq 1\right)
$$

where $\tilde{J}_{L}=E+L^{-d}(-2,2)$ and $\tilde{J}_{L}^{\prime}=E^{\prime}+L^{-d}(-2,2)$.
Proof of Lemma 2.1. Fix $C>0$ large so that $e^{-C \gamma \log L / 2} \leq L^{-2 d-q}$ where $q$ and $\gamma$ are given by assumption (Loc) where we choose $p=d$. Let $\Omega_{0}$ be the set of probability $1-L^{-p}$ where (1) and (2) in assumption (Loc) are satisfied. Define $\tilde{\ell}=C \log L$. We prove

Lemma 2.2. There exists a covering of $\Lambda_{\ell}$ by cubes, say $\Lambda_{\ell}=\cup_{\gamma \in \Gamma}\left[\gamma+\Lambda_{\tilde{\ell}}\right]$, such that $\# \Gamma \asymp(\ell / \tilde{\ell})^{d}$, and, if $\omega \in \Omega_{0}$ is such that $H_{\omega}\left(\Lambda_{\ell}\right)$ has exactly one eigenvalue in $J_{L}$ and exactly one eigenvalue in $J_{L}^{\prime}$, then
(1) either, there exists $\gamma$ and $\gamma^{\prime}$ such that $\gamma+\Lambda_{\tilde{\ell}} \cap \gamma^{\prime}+\Lambda_{\tilde{\ell}}=\emptyset$ and

- $H_{\omega}\left(\gamma+\Lambda_{\tilde{\ell}}\right)$ has exactly one e.v. in $\tilde{J}_{L}$
- $H_{\omega}\left(\gamma^{\prime}+\Lambda_{\tilde{\ell}}\right)$ has exactly one e.v. in $\tilde{J}_{L}^{\prime}$.
(2) or $H_{\omega}\left(\Lambda_{5 \tilde{\ell}}(\gamma)\right)$ has exactly one e.v. in $\tilde{J}_{L}$ and exactly one e.v. in $\tilde{J}_{L}^{\prime}$.

We postpone the proof of Lemma 2.2 to complete that of Lemma 2.1. Using the estimate on $\mathbb{P}\left(\Omega_{0}\right)$, the independence of $H_{\omega}\left(\gamma+\Lambda_{\tilde{\ell}}\right)$ and $H_{\omega}\left(\gamma^{\prime}+\Lambda_{\tilde{\ell}}\right)$ when alternative (1) is the case in Lemma 2.2, Wegner's estimate (2.1) and the fact the random variables are identically distributed, we compute

$$
\begin{aligned}
\mathbb{P}_{0} \leq L^{-2 d}+ & C(\ell \mid \tilde{\ell})^{d} \mathbb{P}\left(\left\{\begin{array}{l}
\sigma\left(H_{\omega}\left(\Lambda_{3 \tilde{\ell}}(0)\right)\right) \cap \tilde{J}_{L} \neq \emptyset \\
\sigma\left(H_{\omega}\left(\Lambda_{3 \tilde{\ell}}(0)\right)\right) \cap \tilde{J}_{L}^{\prime} \neq \emptyset
\end{array}\right\}\right) \\
& +C(\ell / \tilde{\ell})^{2 d} \mathbb{P}\left(\#\left[\sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}(0)\right)\right) \cap \tilde{J}_{L}\right] \geq 1\right) \mathbb{P}\left(\#\left[\sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}(0)\right)\right) \cap \tilde{J}_{L}^{\prime}\right] \geq 1\right) \\
\leq C L^{-2 d}+ & C(\ell / \tilde{\ell})^{2 d}(\tilde{\ell} / L)^{2 d}+C\left(\ell / \tilde{\ell}^{d} \mathbb{P}_{1} \leq C(\ell / L)^{2 d}+C(\ell / \tilde{\ell})^{d} \mathbb{P}_{1}\right.
\end{aligned}
$$

where $\mathbb{P}_{1}$ is defined in Lemma 2.1 for $5 \tilde{\ell}$ replaced with $\tilde{\ell}$. This completes the proof of Lemma 2.1.

Proof of Lemma 2.2. For $\gamma \in \tilde{\ell} \mathbb{Z}^{d} \cap \Lambda_{\ell}$, consider the cubes $\left(\gamma+\Lambda_{\tilde{\ell}}\right)_{\gamma \in \tilde{\ell} \mathbb{Z}^{d} \cap \Lambda_{\ell}}$. They cover $\Lambda_{\ell}$. Recall that we are taking periodic boundary conditions. If the localization centers associated to the two eigenvalues of $H_{\omega}\left(\Lambda_{\ell}\right)$ assumed to be respectively in $\tilde{J}_{L}$ and $\tilde{J}_{L}^{\prime}$ are at a distance less than $3 \tilde{\ell}$ from one another, then we can find $\gamma \in \tilde{\ell} \mathbb{Z}^{d}$ such that both localization centers belong $\gamma+\Lambda_{4 \tilde{\ell}}$ (for $\tilde{\ell}=C \log L$ and $C>0$ sufficiently large). Thus, by the localization property (Loc), we are in case (2).
If the distance is larger than $3 \tilde{\ell}$, we can find $\gamma \in \tilde{\ell} \mathbb{Z}^{d}$ and $\gamma^{\prime} \in \tilde{\ell} \mathbb{Z}^{d}$ such that each of the cubes $\gamma+\Lambda_{\tilde{\ell} / 2}$ and $\gamma^{\prime}+\Lambda_{\tilde{\ell} / 2}$ contains exactly one of the localization centers and $\left(\gamma+\Lambda_{\tilde{\ell} / 2}\right) \cap$ $\left(\gamma^{\prime}+\Lambda_{\tilde{\ell} / 2}\right)=\emptyset$. So for $\tilde{\ell}=C \log L$ and $C>0$ sufficiently large, by the localization property (Loc), we are in case (1).
This completes the proof of Lemma 2.2.

We now proceed with the proof of (2.4). Therefore, by Lemma 2.1, it suffices to prove that $\mathbb{P}_{1}$, defined in Lemma 2.1, satisfies, for some $C>0$,

$$
\mathbb{P}_{1} \leq C(\tilde{\ell} / L)^{2 d} \cdot \begin{cases}e^{\tilde{\beta}^{\beta}} & \text { if the dimension } d=1  \tag{2.5}\\ \tilde{\ell}^{C} & \text { if the dimension } d>1\end{cases}
$$

Let $\left(E_{j}(\omega, \tilde{\ell})\right)_{1 \leq j \leq(2 \tilde{\ell}+1)^{d}}$ be the eigenvalues of $H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)$ ordered in an increasing way and repeated according to multiplicity.
Assume that $\omega \mapsto E(\omega)$ is the only eigenvalue of $H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)$ in $J_{L}$. In this case, by standard perturbation theory arguments (see e.g. [11, 18]), we know that
(1) $E(\omega)$ being simple, $\omega \mapsto E(\omega)$ is real analytic, and if $\omega \mapsto \varphi(\omega)=(\varphi(\omega ; \gamma))_{\gamma \in \Lambda_{\tilde{\ell}}}$ denotes the associated normalized real eigenvector, it is also real analytic in $\omega$;
(2) one has $\partial_{\omega_{\gamma}} E(\omega)=\varphi^{2}(\omega ; \gamma) \geq 0$ which, in particular, implies that

$$
\begin{equation*}
\left\|\nabla_{\omega} E(\omega)\right\|_{\ell^{1}}=1 ; \tag{2.6}
\end{equation*}
$$

(3) the Hessian of $E$ is given by $\operatorname{Hess}_{\omega} E(\omega)=\left(\left(h_{\gamma \beta}\right)\right)_{\gamma, \beta}$ where

- $h_{\gamma, \beta}=-2 \operatorname{Re}\left\langle\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)-E(\omega)\right)^{-1} \psi_{\gamma}(\omega), \psi_{\beta}(\omega)\right\rangle$,
- $\psi_{\gamma}=\varphi(\omega ; \gamma) \Pi(\omega) \delta_{\gamma}$
- $\Pi(\omega)$ is the orthogonal projector on the orthogonal to $\varphi(\omega)$.

We prove
Lemma 2.3. There exists $C>0$ such that

$$
\left\|\operatorname{Hess}_{\omega}(E(\omega))\right\|_{\ell \infty \rightarrow \ell^{1}} \leq \frac{C}{\operatorname{dist}\left(E(\omega), \sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \backslash\{E(\omega)\}\right)}
$$

Proof of Lemma 2.3. First, note that, by definition, $H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)$ depends on $(2 \tilde{\ell}+1)^{d}$ random variables so that $\operatorname{Hess}_{\omega} E(\omega)$ is a $(2 \tilde{\ell}+1)^{d} \times(2 \tilde{\ell}+1)^{d}$ matrix. Hence, for $a=\left(a_{\gamma}\right)_{\gamma \in \Lambda_{\tilde{\ell}}} \in \mathbb{C}^{\Lambda_{\tilde{\ell}}}$ and $b=\left(b_{\gamma}\right)_{\gamma \in \Lambda_{\tilde{\ell}}} \in \mathbb{C}^{\Lambda_{\tilde{\ell}}}$, we compute

$$
\left\langle\operatorname{Hess}_{\omega} E a, b\right\rangle=-2\left\langle\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)-E(\omega)\right)^{-1} \psi_{a}, \psi_{b}\right\rangle
$$

where

$$
\psi_{a}=\Pi(\omega)\left(\sum_{\gamma \in \Lambda_{\tilde{\ell}}} a_{\gamma}\left|\delta_{\gamma}\right\rangle\left\langle\delta_{\gamma}\right|\right) \varphi(\omega)=\sum_{\gamma \in \Lambda_{\tilde{\ell}}} a_{\gamma} \varphi(\omega ; \gamma) \Pi(\omega) \delta_{\gamma} .
$$

Hence, $\left\|\psi_{a}\right\|_{2} \leq C\|a\|_{\infty}$ and, for some $C>0$,

$$
\left\|\operatorname{Hess}_{\omega}(E(\omega))\right\|_{\ell \infty \rightarrow \ell^{1}} \leq \frac{C}{\operatorname{dist}\left(E(\omega), \sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \backslash\{E(\omega)\}\right)}
$$

This completes the proof of Lemma 2.3.
Note that, using (2.2), Lemma 2.3 yields, for $\varepsilon \in\left(4 L^{-d}, 1\right)$,

$$
\mathbb{P}\left(\left\{\omega ; \begin{array}{c}
\sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap \tilde{J}_{L}=\{E(\omega)\} \\
\left\|\operatorname{Hess}_{\omega}(E(\omega))\right\|_{\ell^{\infty} \rightarrow \ell^{1}} \geq \varepsilon^{-1}
\end{array}\right\}\right) \leq C \varepsilon \tilde{\ell}^{\tilde{L}^{2}} L^{-d} .
$$

Hence, for $\varepsilon \in\left(4 L^{-d}, 1\right)$, one has

$$
\begin{equation*}
\mathbb{P}_{1} \leq C \varepsilon \tilde{\ell}^{2 d} L^{-d}+\mathbb{P}_{\varepsilon} \tag{2.7}
\end{equation*}
$$

where

$$
\begin{gather*}
\mathbb{P}_{\varepsilon}=\mathbb{P}\left(\Omega_{0}(\varepsilon)\right)  \tag{2.8}\\
7
\end{gather*}
$$

and

We will now estimate $\mathbb{P}_{\varepsilon}$. The basic idea is to prove that the eigenvalues $E(\omega)$ and $E^{\prime}(\omega)$ depend effectively on at least two independent random variables. A simple way to guarantee this is to ensure that their gradients with respect to $\omega$ are not co-linear. In the present case, the gradients have non negative components and their $\ell^{1}$-norm is 1 ; hence, it suffices to prove that they are different to ensure that they are not co-linear.
We prove
Lemma 2.4. Let $L \geq 1$. For the discrete Anderson model, one has
(1) in any dimension d: for $\Delta E>2 d$, if the random variables $\left(\omega_{\gamma}\right)_{\gamma \in \Lambda_{L}}$ are bounded by $K$, for $E_{j}(\omega)$ and $E_{k}(\omega)$, simple eigenvalues of $H_{\omega}\left(\Lambda_{L}\right)$ such that $\left|E_{k}(\omega)-E_{j}(\omega)\right| \geq \Delta E$, one has

$$
\begin{equation*}
\left\|\nabla_{\omega}\left(E_{j}(\omega)-E_{k}(\omega)\right)\right\|_{2} \geq \frac{\Delta E-2 d}{K}(2 L+1)^{-d / 2} \tag{2.10}
\end{equation*}
$$

(2) in dimension 1: fix $E<E^{\prime}$ and $\beta>1 / 2$; let $\mathbb{P}$ denote the probability that there exists $E_{j}(\omega)$ and $E_{k}(\omega)$, simple eigenvalues of $H_{\omega}\left(\Lambda_{L}\right)$ such that $\left|E_{k}(\omega)-E\right|+\left|E_{j}(\omega)-E^{\prime}\right| \leq$ $e^{-L^{\beta}}$ and such that

$$
\begin{equation*}
\left\|\nabla_{\omega}\left(E_{j}(\omega)-E_{k}(\omega)\right)\right\|_{1} \leq e^{-L^{\beta}} \tag{2.11}
\end{equation*}
$$

then, there exists $c>0$ such that

$$
\begin{equation*}
\mathbb{P} \leq e^{-c L^{2 \beta}} \tag{2.12}
\end{equation*}
$$

We postpone the proof of Lemma 2.4 for a while to estimate $\mathbb{P}_{\varepsilon}$. Set

$$
\lambda=\lambda_{L}= \begin{cases}e^{-\tilde{\ell}^{\beta}} & \text { if the dimension } d=1  \tag{2.13}\\ \frac{\Delta E-2 d}{K} \tilde{\ell}^{-d / 2} & \text { if the dimension } d>1\end{cases}
$$

For $\gamma$ and $\gamma^{\prime}$ in $\Lambda_{\tilde{\ell}}$, define

$$
\begin{equation*}
\Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)=\Omega_{0}(\varepsilon) \cap\left\{\omega ;\left|J_{\gamma, \gamma^{\prime}}\left(E(\omega), E^{\prime}(\omega)\right)\right| \geq \lambda\right\} \tag{2.14}
\end{equation*}
$$

where $J_{\gamma, \gamma^{\prime}}\left(E(\omega), E^{\prime}(\omega)\right)$ is the Jacobian of the mapping $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \mapsto\left(E(\omega), E^{\prime}(\omega)\right)$ i.e.

$$
J_{\gamma, \gamma^{\prime}}\left(E(\omega), E^{\prime}(\omega)\right)=\left|\begin{array}{cc}
\partial_{\omega_{\gamma}} E(\omega) & \partial_{\omega_{\gamma^{\prime}}} E(\omega) \\
\partial_{\omega_{\gamma}} E^{\prime}(\omega) & \partial_{\omega_{\gamma^{\prime}}} E^{\prime}(\omega)
\end{array}\right|
$$

In section 2.4 , we prove
Lemma 2.5. Pick $(u, v) \in\left(\mathbb{R}^{+}\right)^{2 n}$ such that $\|u\|_{1}=\|v\|_{1}=1$. Then

$$
\max _{j \neq k}\left|\begin{array}{ll}
u_{j} & u_{k} \\
v_{j} & v_{k}
\end{array}\right|^{2} \geq \frac{1}{4 n^{5}}\|u-v\|_{1}^{2}
$$

We apply Lemma 2.4 with $L=\tilde{\ell}$ and Lemma 2.5 to obtain that

$$
\begin{equation*}
\mathbb{P}_{\varepsilon} \leq \sum_{\gamma \neq \gamma^{\prime}} \mathbb{P}\left(\Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)\right)+\mathbb{P}_{r} \tag{2.15}
\end{equation*}
$$

where
(1) in dimension 1 , we have $\mathbb{P}_{r} \leq C \tilde{\ell}^{2 d} e^{-c \tilde{\ell}^{2 \beta^{\prime}}}$ for any $1 / 2<\beta^{\prime}<\beta$; thus, for $L$ sufficiently large, as $\tilde{\ell} \geq c \log L$ and $\beta>1 / 2$, we have

$$
\begin{equation*}
\mathbb{P}_{r} \leq L^{-2 d} \tag{2.16}
\end{equation*}
$$

(2) in dimension $d$, as by assumption $\Delta E>2 d$, one has $\mathbb{P}_{r}=0$, thus, (2.16) still holds.

In the sequel, we will write $\omega=\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}, \omega_{\gamma, \gamma^{\prime}}\right)$ where $\omega_{\gamma, \gamma^{\prime}}=\left(\omega_{\beta}\right)_{\beta \notin\left\{\gamma, \gamma^{\prime}\right\}}$.
To estimate $\mathbb{P}\left(\Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)\right)$, we use
Lemma 2.6. Pick $\varepsilon=L^{-d} \lambda^{-3}$. For any $\omega_{\gamma, \gamma^{\prime}}$, if there exists $\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right) \in \mathbb{R}^{2}$ such that $\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}, \omega_{\gamma, \gamma^{\prime}}\right) \in \Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)$, then, for $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \in \mathbb{R}^{2}$ such that $\left|\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right|_{\infty} \geq L^{-d} \lambda^{-2}$, one has $\left(E_{j}(\omega), E_{j^{\prime}}(\omega)\right) \notin \tilde{J}_{L} \times \tilde{J}_{L}^{\prime}$.
Recall that $g$ is the density of the random variables $\left(\omega_{\gamma}\right)_{\gamma}$; it is assumed to be bounded and compactly supported. Hence, the probability $\mathbb{P}\left(\Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)\right)$ is estimated as follows

$$
\begin{align*}
& \mathbb{P}\left(\Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)\right)=\mathbb{E}_{\gamma, \gamma^{\prime}}\left(\int_{\mathbb{R}^{2}} \mathbf{1}_{\Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)}(\omega) g\left(\omega_{\gamma}\right) g\left(\omega_{\gamma^{\prime}}\right) d \omega_{\gamma} d \omega_{\gamma^{\prime}}\right) \\
& \leq \mathbb{E}_{\gamma, \gamma^{\prime}}\left(\int_{\left|\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right| \infty<L^{-d} \lambda^{-2}} g\left(\omega_{\gamma}\right) g\left(\omega_{\gamma^{\prime}}\right) d \omega_{\gamma} d \omega_{\gamma^{\prime}}\right)  \tag{2.17}\\
& \leq C L^{-2 d} \lambda^{-4}
\end{align*}
$$

where $\mathbb{E}_{\gamma, \gamma^{\prime}}$ denotes the expectation with respect to all the random variables except $\omega_{\gamma}$ and $\omega_{\gamma^{\prime}}$.
Summing (2.17) over $\left(\gamma, \gamma^{\prime}\right) \in \Lambda_{\tilde{\ell}}^{2}$, using (2.15) and (2.16), we obtain

$$
\mathbb{P}_{\varepsilon} \leq C L^{-2 d} \lambda^{-4}
$$

We now plug this into (2.7) and use the fact that $\varepsilon=L^{-d} \lambda^{-3}$ to complete the proof of (2.5). This completes the proofs of Lemmas 1.1 and 1.2.

Proof of Lemma 2.6. Recall that, for any $\gamma, \omega_{\gamma} \mapsto E_{j}(\omega)$ and $\omega_{\gamma} \mapsto E_{j^{\prime}}(\omega)$ are non decreasing. Hence, to prove Lemma 2.6, it suffices to prove that, for $\left|\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right|_{\infty}=L^{-d} \lambda^{-2}$, one has $\left(E_{j}(\omega), E_{j^{\prime}}(\omega)\right) \notin \tilde{J}_{L} \times \tilde{J}_{L}^{\prime}$.
Let $\mathcal{S}_{\beta}$ denote the square $\mathcal{S}_{\beta}=\left\{\left|\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right|_{\infty} \leq L^{-d} \lambda^{-2}\right\}$.
Recall that $\varepsilon=L^{-d} \lambda^{-3}$. Pick $\omega_{\gamma, \gamma^{\prime}}$ such that there exists $\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right) \in \mathbb{R}^{2}$ for which one has $\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}, \omega_{\gamma, \gamma^{\prime}}\right) \in \Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)$. To shorten the notations, in the sequel, we write only the variables $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)$ as $\omega_{\gamma, \gamma^{\prime}}$ stays fixed throughout the proof; e.g. we write $E\left(\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)\right)$ instead of $E\left(\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}, \omega_{\gamma, \gamma^{\prime}}\right)\right)$.
Consider the mapping $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \mapsto \varphi\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right):=\left(E(\omega), E^{\prime}(\omega)\right)$. We will show that $\varphi$ defines an analytic diffeomorphism form $\mathcal{S}_{\beta}$ to $\varphi\left(\mathcal{S}_{\beta}\right)$.
By (2.14) and (2.9), the definitions of $\Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)$ and $\Omega_{0}(\varepsilon)$, we know that

$$
\begin{aligned}
& \sigma\left(H_{\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap(E-C \varepsilon, E+C \varepsilon)=\{E(\omega)\} \subset\left(E-C L^{-d}, E+C L^{-d}\right), \\
& \sigma\left(H_{\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{\prime}\right)}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap[(E-C \varepsilon, E-C \varepsilon / 2) \cup(E+C \varepsilon / 2, E+C \varepsilon)]=\emptyset, \\
& \sigma\left(H_{\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap\left(E^{\prime}-C \varepsilon, E^{\prime}+C \varepsilon\right)=\left\{E^{\prime}(\omega)\right\} \subset\left(E^{\prime}-C L^{-d}, E^{\prime}+C L^{-d}\right), \\
& \left.\sigma\left(H_{\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)}\right), \Lambda_{\tilde{\ell})}\right) \cap\left[\left(E^{\prime}-C \varepsilon, E^{\prime}-C \varepsilon / 2\right) \cup\left(E^{\prime}+C \varepsilon / 2, E^{\prime}+C \varepsilon\right)\right]=\emptyset . \\
& 9
\end{aligned}
$$

By (2.6), as $L^{-d} \lambda^{-2} \leq \lambda \varepsilon$, for $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \in \mathcal{S}_{\beta}$, one has

$$
\begin{aligned}
& \sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap(E-C \varepsilon / 2, E+C \varepsilon / 2)=\{E(\omega)\} \subset(E-C \varepsilon / 4, E+C \varepsilon / 4), \\
& \sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap[(E-C \varepsilon / 2, E-C \varepsilon / 4) \cup(E+C \varepsilon / 4, E+C \varepsilon / 2)]=\emptyset, \\
& \sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap\left(E^{\prime}-C \varepsilon / 2, E^{\prime}+C \varepsilon / 2\right)=\{E(\omega)\} \subset\left(E^{\prime}-C \varepsilon / 4, E^{\prime}+C \varepsilon / 4\right), \\
& \sigma\left(H_{\omega}\left(\Lambda_{\tilde{\ell}}\right)\right) \cap\left[\left(E^{\prime}-C \varepsilon / 2, E^{\prime}-C \varepsilon / 4\right) \cup\left(E^{\prime}+C \varepsilon / 4, E^{\prime}+C \varepsilon / 2\right)\right]=\emptyset .
\end{aligned}
$$

Hence, by Lemma 2.3, for $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \in \mathcal{S}_{\beta}$, one has

$$
\left\|\operatorname{Hess}_{\omega}(E(\omega))\right\|_{\ell^{\infty} \rightarrow \ell^{1}}+\left\|\operatorname{Hess}_{\omega}\left(E^{\prime}(\omega)\right)\right\|_{\ell^{\infty} \rightarrow \ell^{1}} \leq C \varepsilon^{-1} \leq C L^{d} \lambda^{3} .
$$

By (2.6) and the Fundamental Theorem of Calculus, for $\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \in \mathcal{S}_{\beta}$, we get that,

$$
\begin{align*}
& \left\|\nabla \varphi\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\nabla \varphi\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right\| \\
& \leq\left(\left\|\operatorname{Hess}_{\omega}(E(\omega))\right\|_{\ell \infty \rightarrow \ell^{1}}+\left\|\operatorname{Hess}_{\omega}\left(E^{\prime}(\omega)\right)\right\|_{\ell \infty \rightarrow \ell^{1}}\right) L^{-d} \lambda^{-1} \leq C \lambda^{2} . \tag{2.18}
\end{align*}
$$

Let us show that $\varphi$ is one-to-one on the square $\mathcal{S}_{\beta}$. Using (2.18), we compute

$$
\left|\varphi\left(\omega_{\gamma}^{\prime}, \omega_{\gamma^{\prime}}^{\prime}\right)-\varphi\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\nabla \varphi\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right) \cdot\binom{\omega_{\gamma}^{\prime}-\omega_{\gamma}}{\omega_{\gamma^{\prime}}^{\prime}-\omega_{\gamma^{\prime}}}\right| \leq \lambda^{2}\left\|\binom{\omega_{\gamma}^{\prime}-\omega_{\gamma}}{\omega_{\gamma^{\prime}}^{\prime}-\omega_{\gamma^{\prime}}}\right\|
$$

As $\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}, \omega_{\gamma, \gamma^{\prime}}\right) \in \Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)$, we have

$$
\left|\operatorname{Jac} \varphi\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right| \geq \lambda .
$$

Hence, for $\tilde{\ell}$ large, we have

$$
\left|\varphi\left(\omega_{\gamma}^{\prime}, \omega_{\gamma^{\prime}}^{\prime}\right)-\varphi\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)\right| \geq \frac{1}{2} \lambda\left\|\binom{\omega_{\gamma}^{\prime}-\omega_{\gamma}}{\omega_{\gamma^{\prime}}^{\prime}-\omega_{\gamma^{\prime}}}\right\|
$$

so $\varphi$ is one-to-one. The estimate (2.18) yields

$$
\left|\operatorname{Jac} \varphi\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\operatorname{Jac} \varphi\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right| \leq \lambda^{2}
$$

As $\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}, \omega_{\gamma, \gamma^{\prime}}\right) \in \Omega_{0, \beta}^{\gamma, \gamma^{\prime}}(\varepsilon)$, for $L$ sufficiently large, this implies that

$$
\begin{equation*}
\forall\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right) \in \mathcal{S}_{\beta}, \quad\left|J_{\gamma, \gamma^{\prime}}\left(E(\omega), E^{\prime}(\omega)\right)\right| \geq \frac{1}{2} \lambda . \tag{2.19}
\end{equation*}
$$

The Local Inversion Theorem then guarantees that $\varphi$ is an analytic diffeomorphism from $\mathcal{S}_{\beta}$ onto $\varphi\left(\mathcal{S}_{\beta}\right)$. By (2.19), the Jacobian matrix of its inverse is bounded by $C \tilde{\ell}^{\beta}$ for some $C>0$ independent of $L$. Hence, if for some $\left|\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right|_{\infty}=L^{-d} \lambda^{-2}$, one has $\left(E(\omega), E^{\prime}(\omega)\right) \in \tilde{J}_{L} \times \tilde{J}_{L}^{\prime}$, then

$$
L^{-d} \lambda^{-2}=\left|\left(\omega_{\gamma}, \omega_{\gamma^{\prime}}\right)-\left(\omega_{\gamma}^{0}, \omega_{\gamma^{\prime}}^{0}\right)\right|_{\infty}=\left|\varphi^{-1}\left(E(\omega), E^{\prime}(\omega)\right)-\varphi^{-1}\left(E, E^{\prime}\right)\right|_{\infty} \leq C L^{-d} \lambda^{-1}
$$

which is absurd when $L \rightarrow+\infty$ as $\lambda=\lambda_{L} \rightarrow 0$ (see (2.13)). This completes the proof of Lemma 2.6.
2.3. Proof of Lemma 2.4. A fundamental difference between the points (1) and (2) in Lemma 2.4 is that to prove point (2), we will the fact that $H_{0}$ is the discrete Laplacian. In the proof of point (1), we can take $H_{0}$ to be any convolution matrix with exponentially decaying off diagonal coefficients if one replaces the condition $\left|E-E^{\prime}\right|>2 d$ with the condition $\left|E-E^{\prime}\right|>\sup \sigma\left(H_{0}\right)-\inf \sigma\left(H_{0}\right)$.
As it is simpler, we start with the proof of point (1).
2.3.1. The proof of point (1). Let $E_{j}(\omega)$ and $E_{k}(\omega)$ be simple eigenvalues of $H_{\omega}\left(\Lambda_{L}\right)$ such that $\left|E_{k}(\omega)-E_{j}(\omega)\right| \geq \Delta E>2 d$. Then, $\omega \mapsto E_{j}(\omega)$ and $\omega \mapsto E_{k}(\omega)$ are real analytic functions. Let $\omega \mapsto \varphi_{j}(\omega)$ and $\omega \mapsto \varphi_{k}(\omega)$ be normalized eigenvectors associated respectively to $E_{j}(\omega)$ and $E_{k}(\omega)$. Differentiating the eigenvalue equation in $\omega$, one computes

$$
\begin{aligned}
\omega \cdot \nabla_{\omega}\left(E_{j}(\omega)\right. & \left.-E_{k}(\omega)\right)=\left\langle V_{\omega} \varphi_{j}(\omega), \varphi_{j}(\omega)\right\rangle-\left\langle V_{\omega} \varphi_{k}(\omega), \varphi_{k}(\omega)\right\rangle \\
& =E_{j}(\omega)-E_{k}(\omega)+\left\langle-\Delta \varphi_{k}(\omega), \varphi_{k}(\omega)\right\rangle-\left\langle-\Delta \varphi_{j}(\omega), \varphi_{j}(\omega)\right\rangle
\end{aligned}
$$

As $0 \leq-\Delta \leq 2 d$ and as $\varphi_{j}(\omega)$ and $\varphi_{k}(\omega)$ are normalized, we get that

$$
\Delta E-2 d \leq\left|E_{j}(\omega)-E_{k}(\omega)\right|-2 d \leq\left|\omega \cdot \nabla_{\omega}\left(E_{j}(\omega)-E_{k}(\omega)\right)\right|
$$

Hence, as the random variables $\left(\omega_{\gamma}\right)_{\gamma \in \Lambda}$ are bounded, the Cauchy Schwartz inequality yields

$$
\left\|\nabla_{\omega}\left(E_{j}(\omega)-E_{k}(\omega)\right)\right\|_{2} \geq \frac{\Delta E-2 d}{K}(2 L+1)^{-d / 2}
$$

which completes the proof of (2.10).
2.3.2. The proof of point (2). Let us now assume $d=1$. Fix $E<E^{\prime}$. Pick $E_{j}(\omega)$ and $E_{k}(\omega)$, simple eigenvalues of $H_{\omega}\left(\Lambda_{L}\right)$ such that $\left|E_{k}(\omega)-E\right|+\left|E_{j}(\omega)-E^{\prime}\right| \leq e^{-L^{\beta}}$. Then, $\omega \mapsto E_{j}(\omega)$ and $\omega \mapsto E_{k}(\omega)$ are real analytic functions. Let $\omega \mapsto \varphi^{j}(\omega)$ and $\omega \mapsto \varphi^{k}(\omega)$ be normalized eigenvectors associated respectively to $E_{j}(\omega)$ and $E_{k}(\omega)$. One computes

$$
\nabla_{\omega} E_{j}(\omega)=\left(\left[\varphi^{j}(\omega ; \gamma)\right]^{2}\right)_{\gamma \in \Lambda_{L}} \quad \text { and } \quad \nabla_{\omega} E_{k}(\omega)=\left(\left[\varphi^{k}(\omega ; \gamma)\right]^{2}\right)_{\gamma \in \Lambda_{L}}
$$

Hence, if

$$
\begin{equation*}
e^{-L^{\beta}} \geq\left\|\nabla_{\omega}\left(E_{j}(\omega)-E_{k}(\omega)\right)\right\|_{1}=\sum_{\gamma \in \Lambda_{L}}\left|\varphi^{j}(\omega ; \gamma)-\varphi^{k}(\omega ; \gamma)\right| \cdot\left|\varphi^{j}(\omega ; \gamma)+\varphi^{k}(\omega ; \gamma)\right| \tag{2.20}
\end{equation*}
$$

as $\left\|\nabla_{\omega} E_{j}(\omega)\right\|=\left\|\nabla_{\omega} E_{k}(\omega)\right\|=1$, there exists a partition of $\Lambda_{L}=\{-L, \cdots, L\}$, say $\mathcal{P} \subset \Lambda_{L}$ and $\mathcal{Q} \subset \Lambda_{L}$ such that $\mathcal{P} \cup \mathcal{Q}=\Lambda_{L}$ and $\mathcal{P} \cap \mathcal{Q}=\emptyset$ and such that

- for $\gamma \in \mathcal{P},\left|\varphi^{j}(\omega ; \gamma)-\varphi^{k}(\omega ; \gamma)\right| \leq e^{-L^{\beta} / 2}$;
- for $\gamma \in \mathcal{Q},\left|\varphi^{j}(\omega ; \gamma)+\varphi^{k}(\omega ; \gamma)\right| \leq e^{-L^{\beta} / 2}$.

Introduce the orthogonal projectors $P$ and $Q$ defined by

$$
P=\sum_{\gamma \in \mathcal{P}}|\gamma\rangle\langle\gamma| \quad \text { and } \quad Q=\sum_{\gamma \in \mathcal{Q}}|\gamma\rangle\langle\gamma| .
$$

One has

$$
\left\|P \varphi^{j}-P \varphi^{k}\right\|_{2} \leq \sqrt{L} e^{-L^{\beta} / 2} \quad \text { and } \quad\left\|Q \varphi^{j}+Q \varphi^{k}\right\|_{2} \leq \sqrt{L} e^{-L^{\beta} / 2}
$$

Clearly $\left\|P \varphi^{j}\right\|^{2}+\left\|Q \varphi^{j}\right\|^{2}=\left\|\varphi^{j}\right\|^{2}=1$. As $\left\langle\varphi^{j}, \varphi^{k}\right\rangle=0$, one has

$$
\begin{aligned}
0 & =\left\langle(P+Q) \varphi^{j},(P+Q) \varphi^{k}\right\rangle=\left\langle P \varphi^{j}, P \varphi^{k}\right\rangle+\left\langle Q \varphi^{j}, Q \varphi^{k}\right\rangle \\
& =\left\|P \varphi^{j}\right\|^{2}-\left\|Q \varphi^{j}\right\|^{2}+O\left(\sqrt{L} e^{-L^{\beta} / 2}\right)
\end{aligned}
$$

Hence

$$
\left\|P \varphi^{j}\right\|^{2}=\frac{1}{2}+O\left(\sqrt{L} e^{-L^{\beta} / 2}\right) \text { and }\left\|Q \varphi^{j}\right\|^{2}=\frac{1}{2}+O\left(\sqrt{L} e^{-L^{\beta} / 2}\right)
$$

This implies that

$$
\begin{equation*}
\mathcal{P} \neq \emptyset \text { and } \mathcal{Q} \neq \emptyset \tag{2.21}
\end{equation*}
$$

We set $h_{-}=P \varphi^{j}-P \varphi^{k}$ and $h_{+}=Q \varphi^{j}+Q \varphi^{k}$. The eigenvalue equations for $E_{j}(\omega)$ and $E_{k}(\omega)$ yields

$$
\left(-\Delta+W_{\omega}\right) \varphi^{j}=\Delta E(\omega) \varphi^{j} \underset{11}{\operatorname{and}\left(-\Delta+W_{\omega}\right) \varphi^{k}=-\Delta E(\omega) \varphi^{k} .}
$$

where

$$
\Delta E(\omega)=\left(E_{j}(\omega)-E_{k}(\omega)\right) / 2, \quad W_{\omega}=V_{\omega}-\bar{E}(\omega), \quad \bar{E}(\omega)=\left(E_{j}(\omega)+E_{k}(\omega)\right) / 2
$$

To simplify the notation, from now on, we write $u=\varphi^{j}$; then, one has $\varphi^{k}=P u-Q u+$ $O\left(\sqrt{L} e^{-L^{\beta} / 2}\right)$. This yields

$$
\begin{cases}\left(-\Delta+W_{\omega}\right)(P u+Q u) & =\Delta E(\omega)(P u+Q u) \\ \left(-\Delta+W_{\omega}\right)\left(P u-Q u+h_{-}-h_{+}\right) & =-\Delta E(\omega)\left(P u-Q u+h_{-}-h_{+}\right)\end{cases}
$$

that is

$$
\left\{\begin{aligned}
\left(-\Delta+W_{\omega}\right)(P u) & =\Delta E(\omega) Q u-h \\
\left(-\Delta+W_{\omega}\right)(Q u) & =\Delta E(\omega) P u+h
\end{aligned}\right.
$$

where $h:=\left(-\Delta+W_{\omega}-\Delta E(\omega)\right)\left(h_{-}-h_{+}\right) / 2$. As $P W_{\omega} Q=0$, this can also be written as

$$
\begin{cases}{[-(P \Delta Q+Q \Delta P)-\Delta E] u} & =h_{1}  \tag{2.22}\\ {\left[-(P \Delta P+Q \Delta Q)+V_{\omega}-\bar{E}\right] u} & =h_{2}\end{cases}
$$

where

$$
\begin{array}{cl}
h_{1}:=(P-Q) h+(\Delta E(\omega)-\Delta E) u, & h_{2}:=(Q-P) h+(\bar{E}(\omega)-\bar{E}) u, \\
\Delta E=\left(E^{\prime}-E\right) / 2, & \bar{E}=\left(E+E^{\prime}\right) / 2 .
\end{array}
$$

By our assumption on $E_{j}(\omega)$ and $E_{k}(\omega)$, we know that

$$
|\Delta E(\omega)-\Delta E| \leq 2 e^{-L^{\beta}}, \quad|\bar{E}(\omega)-\bar{E}| \leq e^{-L^{\beta}}, \quad\|h\| \leq C \sqrt{L} e^{-L^{\beta} / 2}
$$

Hence, we get that

$$
\begin{equation*}
\left\|h_{1}\right\|+\left\|h_{2}\right\| \leq C \sqrt{L} e^{-L^{\beta} / 2} \tag{2.23}
\end{equation*}
$$

So the above equations imply that

- $\Delta E$ is at a distance at most $\sqrt{L} e^{-L^{\beta} / 2}$ to the spectrum of the deterministic operator $-(P \Delta Q+Q \Delta P)$,
- $u$ is close to being in the eigenspace associated to the eigenvalues close to $\Delta E$,
- finally, $u$ is close to being in the kernel of the random operator $-(P \Delta P+Q \Delta Q)+$ $V_{\omega}-\bar{E}$.
The firsts conditions will be used to describe $u$. The last condition will be interpreted as a condition determining the random variables $\omega_{\gamma}$ for sites $\gamma$ such that $\left|u_{\gamma}\right|$ is not too small. We will show that the number of these sites is of size the volume of the cube $\Lambda_{L}$; so, the probability that the second equation in $(2.22)$ be satisfied should be very small.
To proceed, we first study the operator $-P \Delta Q-Q \Delta P$. As we consider periodic boundary conditions, we compute

$$
\begin{equation*}
-P \Delta Q-Q \Delta P=\sum_{\gamma \in \partial \mathcal{P}}(|\gamma+1\rangle\langle\gamma|+|\gamma\rangle\langle\gamma+1|)+\sum_{\gamma \in \partial \mathcal{Q}}(|\gamma+1\rangle\langle\gamma|+|\gamma\rangle\langle\gamma+1|) \tag{2.24}
\end{equation*}
$$

where $\partial \mathcal{P}=\{\gamma \in \mathcal{P} ; \gamma+1 \in \mathcal{Q}\} \subset \mathcal{P}$ and $\partial \mathcal{Q}=\{\gamma \in \mathcal{Q} ; \gamma+1 \in \mathcal{P}\} \subset \mathcal{Q}$. By (2.21), we know that $\partial \mathcal{P} \neq \emptyset$ and $\partial \mathcal{Q} \neq \emptyset$.
We first note that $\partial \mathcal{P} \cap \partial \mathcal{Q}=\emptyset$. Here, as we are considering the operators with periodic boundary conditions on $\Lambda_{L}$, we identify $\Lambda_{L}$ with $\mathbb{Z} / L \mathbb{Z}$.
For $\mathcal{A} \subset \Lambda_{L}$ we define $\mathcal{A}+1=\{p+1 ; p \in \mathcal{A}\}$ to be the shift by one of $\mathcal{A}$. By definition, $(\partial \mathcal{P}+1) \subset \mathcal{Q}$ and $(\partial \mathcal{Q}+1) \subset \mathcal{P}$. Hence, $(\partial \mathcal{P}+1) \cap \partial \mathcal{P}=\emptyset$ and $(\partial \mathcal{Q}+1) \cap \partial \mathcal{Q}=\emptyset$.
Consider the set $\mathcal{C}:=\partial \mathcal{P} \cup \partial \mathcal{Q}$. We can partition it into its "connected components" i.e.
$\mathcal{C}$ can be written as a disjoint union of intervals of integers, say $\mathcal{C}=\cup_{l=1}^{l_{0}} \mathcal{C}_{l}^{c}$. Then, by the definition of $\partial \mathcal{P}$ and $\partial \mathcal{Q}$, for $l \neq l^{\prime}$, one has,

$$
\begin{equation*}
\mathcal{C}_{l}^{c} \cap \mathcal{C}_{l^{\prime}}^{c}=\mathcal{C}_{l}^{c} \cap\left(\mathcal{C}_{l^{\prime}}^{c}+1\right)=\emptyset . \tag{2.25}
\end{equation*}
$$

Define $\mathcal{C}_{l}=\mathcal{C}_{l}^{c} \cup\left(\mathcal{C}_{l}^{c}+1\right)$. (2.25) implies that, for $l \neq l^{\prime}$,

$$
\begin{equation*}
\mathcal{C}_{l} \cap \mathcal{C}_{l^{\prime}}=\emptyset . \tag{2.26}
\end{equation*}
$$

Note that one may have $\cup_{l=1}^{l_{0}} \mathcal{C}_{l}=\Lambda_{L}$. The representation (2.24) then implies that the following block decomposition

$$
\begin{equation*}
-P \Delta Q-Q \Delta P=-\sum_{l=1}^{l_{0}} C_{l} \Delta C_{l} \tag{2.27}
\end{equation*}
$$

where $C_{l}$ is the projector $C_{l}=\sum_{\gamma \in \mathcal{C}_{j}}|\gamma\rangle\langle\gamma|$.
Note that, by (2.26), the projectors $C_{l}$ and $C_{l^{\prime}}$ are orthogonal to each other for $l \neq l^{\prime}$. So the spectrum of the operator $-P \Delta Q-Q \Delta P$ is given by the union of the spectra of $\left(C_{l} \Delta C_{l}\right)_{1 \leq l \leq l_{0}}$. Each of these operators is the Dirichlet Laplacian on an interval of length $\# \mathcal{C}_{l}$. Its spectral decomposition can be computed explicitly. We will use some facts from this decomposition that we state now.

Lemma 2.7. On a segment of length $n$, the Dirichlet Laplacian $\Delta_{n}$ i.e. the $n \times n$ matrix

$$
\Delta_{n}=\left(\begin{array}{cccccc}
0 & 1 & 0 & \cdots & \cdots & 0 \\
1 & 0 & 1 & 0 & & \\
0 & 1 & 0 & \ddots & \ddots & \vdots \\
\vdots & 0 & \ddots & \ddots & 1 & 0 \\
\vdots & & \ddots & 1 & 0 & 1 \\
0 & \cdots & \cdots & 0 & 1 & 0
\end{array}\right)
$$

satisfies

- its eigenvalues are simple and are given by $(2 \cos (k \pi /(n+1)))_{1 \leq k \leq n}$;
- for $k \in\{1, \cdots, n\}$, the eigenspace associated to $2 \cos (k \pi /(n+1))$ is generated by the vector $(\sin [k j \pi /(n+1)])_{1 \leq j \leq n}$.
Moreover, there exists $K_{1}>0$ such that, for any $n \geq 1$, one has

$$
\begin{equation*}
\inf _{1 \leq k<k^{\prime} \leq n}\left|2 \cos \left(\frac{k \pi}{n+1}\right)-2 \cos \left(\frac{k^{\prime} \pi}{n+1}\right)\right| \geq \frac{1}{K_{1} n^{2}} . \tag{2.28}
\end{equation*}
$$

Proof of Lemma 2.7. The first statement follows immediately from the identity

$$
\sin \left(\frac{k(j+1) \pi}{n+1}\right)+\sin \left(\frac{k(j-1) \pi}{n+1}\right)=2 \cos \left(\frac{k \pi}{n+1}\right) \sin \left(\frac{k j \pi}{n+1}\right) .
$$

The estimate (2.28) is an immediate consequence of

$$
\cos \left(\frac{k \pi}{n+1}\right)-\cos \left(\frac{k^{\prime} \pi}{n+1}\right)=-2 \sin \left(\frac{\left(k+k^{\prime}\right) \pi}{2(n+1)}\right) \sin \left(\frac{\left(k-k^{\prime}\right) \pi}{2(n+1)}\right) .
$$

We now solve the first equation in (2.22) that is describe $u$ solution to this equation.

Lemma 2.8. Let $u$ be a solution to (2.22) such that $\|u\|=1$. Then, for $L$ sufficiently large, one has

$$
\begin{equation*}
\left\|u-\sum_{l=1}^{l_{0}} C_{l} u\right\| \leq e^{-L^{\beta} / 3} \tag{2.29}
\end{equation*}
$$

where, if for $1 \leq l \leq l_{0}$, we write $\mathcal{C}_{l}=\left\{\gamma_{l}^{-}, \cdots, \gamma_{l}^{+}\right\} \quad\left(n_{l}=\gamma_{l}^{+}-\gamma_{l}^{-}+1\right)$, then,

- either there exists a unique $k_{l} \in\left\{1, \cdots, n_{l}\right\}$ satisfying

$$
\begin{equation*}
\left|2 \cos \left(\frac{k_{l} \pi}{n_{l}+1}\right)-\Delta E\right|<\frac{1}{K_{1} n^{2}} \tag{2.30}
\end{equation*}
$$

and $\alpha^{l} \in \mathbb{R}$ such that

$$
\begin{equation*}
\left\|C_{l} u-\alpha^{l} u^{l}\right\| \leq e^{-L^{\beta} / 3} \tag{2.31}
\end{equation*}
$$

where

$$
u_{\gamma}^{l}= \begin{cases}\sin \left(\frac{k_{l}\left(\gamma-\gamma_{l}^{-}+1\right) \pi}{n_{l}+1}\right) & \text { if } \gamma \in C_{l}, \\ 0 & \text { if } \gamma \notin C_{l} .\end{cases}
$$

- there exists no $k_{l} \in\left\{1, \cdots, n_{l}\right\}$ satisfying (2.30) then

$$
\left\|C_{l} u\right\| \leq e^{-L^{\beta} / 3} .
$$

Proof of Lemma 2.8. By Lemma 2.7, the spacing between consecutive eigenvalues of $-C_{l} \Delta C_{l}$ is bounded below by $1 /\left(K_{1} n^{2}\right)$.
Let $C^{\perp}=1-\sum_{l=1}^{l_{0}} C_{l}$. Hence, $u=\sum_{l=1}^{l_{0}} C_{l} u+C^{\perp} u$, the terms in this sums being two by two orthogonal to each other. As $\Delta E>0$, the first equation in (2.22) then yields
(2.32) $\forall 1 \leq l \leq l_{0}, \quad\left\|-C_{l} \Delta C_{l} u-\Delta E C_{l} u\right\| \leq C \sqrt{L} e^{-L^{\beta} / 2} \quad$ and $\quad\left\|C^{\perp} u\right\| \leq C \sqrt{L} e^{-L^{\beta} / 2}$.

Write $\mathcal{C}_{l}=\left\{\gamma_{l}^{-}, \gamma_{l}^{-}+1, \cdots, \gamma_{l}^{+}\right\}$where one may have $\gamma_{l}^{-}=\gamma_{l}^{+}$. We assume that the $\left(\mathcal{C}_{l}\right)_{1 \leq l \leq l_{0}}$ are ordered so that $\gamma_{l}^{+}<\gamma_{l+1}^{-}$.
By the characterization of the spectrum of $-C_{l} \Delta C_{l}$,

- if $2 \cos \left(k_{l} \pi /(n+1)\right)$ is an eigenvalue of $-C_{l} \Delta C_{l}$ closer to $\Delta E$ than a distance $L^{-2} / 4 K_{1}$ (by the remark made above, such an eigenvalue is unique), then, for some $\alpha^{l}$ real, one has

$$
\left\|C_{l} u-\alpha^{l} u^{l}\right\| \leq C L^{5 / 2} e^{-L^{\beta} / 2} .
$$

- if there is no such eigenvalue, then

$$
\begin{equation*}
\left\|C_{l} u\right\| \leq C L^{5 / 2} e^{-L^{\beta} / 2} . \tag{2.33}
\end{equation*}
$$

This completes the proof of Lemma 2.8.
We now prove that $\left|u_{\gamma}\right|$ cannot be really small for too many $\gamma$.
Lemma 2.9. There exists $c>0$ such that, for $L$ sufficiently large,
(1) either $\# \mathcal{C} \geq L / 3$ and, for $\gamma \in \mathcal{C},\left|u_{\gamma}\right| \geq e^{-L^{\beta} / 6}$,
(2) or $l_{0} \geq 2 c L^{\beta}$ and there exists $l^{*} \in\left\{1, \cdots, l_{0}\right\}$ such that, for $\left|l-l^{*}\right| \leq c L^{\beta}$, and $\gamma \in \mathcal{C}_{l}$, one has $\left|u_{\gamma}\right| \geq e^{-L^{\beta} / 6}$.

Proof of Lemma 2.9. To prove Lemma 2.9, we compare the values of $u$ on $\mathcal{C}_{l}$ and $\mathcal{C}_{l+1}$, that is, the vectors $C_{l} u$ and $C_{l+1} u$ given by Lemma 2.8.
First, notice that up to an error of size at most $e^{-L^{\beta} / 3}, u$ on $\mathcal{C}_{l}$ is determined by its coefficient $u_{\gamma_{l}^{+}}$, or equivalently, by its coefficient $u_{\gamma_{l}^{-}}$; in particular as $\sin \left(k_{l} \pi /\left(n_{l}+1\right)\right)=$ $(-1)^{k_{l}-1} \sin \left(k_{l} n_{l} \pi /\left(n_{l}+1\right)\right)$, the representations (2.29) and (2.31) yields

$$
\begin{equation*}
\left|\left|u_{\gamma_{l}^{-}}\right|-\left|u_{\gamma_{l}^{+}}\right|\right|+\left|u_{\gamma_{l}^{+}}-\alpha^{l} \sin \left(k_{l} \pi /\left(n_{l}+1\right)\right)\right| \leq C e^{-L^{\beta} / 3} . \tag{2.34}
\end{equation*}
$$

Notice also that, as $2 \leq n_{l} \leq 2 L+1$ is fixed, for $\rho^{*}:=\sqrt{\frac{n_{l}}{2}+\frac{1}{2} \cos \left(\frac{2 k_{l} \pi}{n_{l}+1}\right)}$, one has

$$
\begin{equation*}
\sup _{1 \leq l \leq l_{0}}\left|\left\|C_{l} u\right\|-\rho_{l}\right| \alpha^{l} \| \mid \leq C e^{-L^{\beta} / 3} \tag{2.35}
\end{equation*}
$$

To compare the values of $u$ on $\mathcal{C}_{l}$ and $\mathcal{C}_{l+1}$, we use the second equation of (2.22) or, equivalently, the eigenvalue equation for $u$ that reads (see (2.22))

$$
\begin{equation*}
\left(-\Delta+V_{\omega}-\bar{E}\right) u=\Delta E u+e \tag{2.36}
\end{equation*}
$$

where $e=h_{1}+h_{2}\left(\right.$ see $(2.22)$ ); hence, $\|e\| \leq C \sqrt{L} e^{-L^{\beta} / 2}$.
We will discuss three cases depending on how far $\gamma_{l}^{+}$and $\gamma_{l+1}^{-}$are from one another:
(1) if $\operatorname{dist}\left(\mathcal{C}_{l}, \mathcal{C}_{l+1}\right) \geq 3$, that is, if $\gamma_{l}^{+}<\gamma_{l}^{+}+1<\gamma_{l+1}^{-}-1<\gamma_{l+1}^{-}:$as $\left\{\gamma_{l}^{+}+1, \cdots, \gamma_{l+1}^{-}-\right.$ $1\} \cap\left[\cup_{l=1}^{l_{0}} \mathcal{C}_{l}\right]=\emptyset$, by (2.29), we know that $\left|u_{n}\right| \leq C L^{2-\alpha}$ for $n \in\left\{\gamma_{l}^{+}+1, \cdots, \gamma_{l+1}^{-}-1\right\}$. The eigenvalue equation (2.36) at the points $\gamma_{l}^{+}+1$ and $\gamma_{l+1}^{-}-1$ then tells us that

$$
\left|u_{\gamma_{l}^{+}}\right|+\left|u_{\gamma_{l+1}^{-}}\right| \leq C e^{-L^{\beta} / 3} .
$$

Thus, by (2.34) and (2.35)

$$
\begin{equation*}
\left\|C_{l} u\right\|+\left\|C_{l+1} u\right\| \leq C e^{-L^{\beta} / 4} . \tag{2.37}
\end{equation*}
$$

(2) if $\operatorname{dist}\left(\mathcal{C}_{l}, \mathcal{C}_{l+1}\right)=2$, that is, if $\gamma_{l}^{+}<\gamma_{l}^{+}+1=\gamma_{l+1}^{-}-1<\gamma_{l+1}^{-}:$as $\gamma_{l}^{+}+1 \notin \cup_{l=1}^{l_{0}} \mathcal{C}_{l}$, by (2.29), we know that $\left|u_{\gamma_{l}^{+}+1}\right| \leq C L^{2-\alpha}$. Hence, in the same way as above, the eigenvalue equation (2.36) at the point $\gamma_{l}^{+}+1$ tells us that

$$
\left|u_{\gamma_{l}^{+}}+u_{\gamma_{l+1}^{-}}\right| \leq C e^{-L^{\beta} / 3} .
$$

Thus, by (2.34) and (2.35)

$$
\begin{equation*}
\left|\left\|C_{l} u\right\|-\left\|C_{l+1} u\right\|\right| \leq C e^{-L^{\beta} / 4} . \tag{2.38}
\end{equation*}
$$

(3) if $\operatorname{dist}\left(\mathcal{C}_{l}, \mathcal{C}_{l+1}\right)=1$, that is, if $\gamma_{l}^{+}+1=\gamma_{l+1}^{-}$: then, the first equation in (2.22) and the decomposition (2.27) yield

$$
\left|u_{\gamma_{l}^{+}-1}-\Delta E u_{\gamma_{l}^{+}}\right|+\left|u_{\gamma_{l+1}^{-}+1}-\Delta E u_{\gamma_{l+1}^{-}}\right| \leq C e^{-L^{\beta} / 3} .
$$

The eigenvalue equation (2.36) at the points $\gamma_{l}^{+}$and $\gamma_{l+1}^{-}$yields

$$
\begin{aligned}
\mid u_{\gamma_{l}^{+}-1}+u_{\gamma_{l+1}^{-}}+\left(\omega_{\gamma_{l}^{+}}-\bar{E}-\Delta E\right) & u_{\gamma_{l}^{+}} \mid \\
& +\left|u_{\gamma_{l}^{+}}+u_{\gamma_{l+1}^{-}+1}+\left(\omega_{\gamma_{l+1}^{-}}-\bar{E}-\Delta E\right) u_{\gamma_{l+1}^{-}}\right| \leq C e^{-L^{\beta} / 3}
\end{aligned}
$$

Summing these two equations, we obtain

$$
\left|u_{\gamma_{l+1}^{-}}+\left(\omega_{\gamma_{l}^{+}}-\bar{E}\right) u_{\gamma_{l}^{+}}\right|+\left|u_{\gamma_{l}^{+}}+\left(\omega_{\gamma_{l+1}^{-}}-\bar{E}\right) u_{\gamma_{l+1}^{-}}\right| \leq C e^{-L^{\beta} / 3} .
$$

Then, as the random variables $\left(\omega_{n}\right)_{n \in \mathbb{Z}}$ are bounded, using (2.34) and (2.35), there exists $C>1$ such that

$$
\begin{equation*}
\frac{1}{C}\left(\left\|C_{l} u\right\|-C e^{-L^{\beta} / 4}\right) \leq\left\|C_{l+1} u\right\| \leq C\left(\left\|C_{l} u\right\|+e^{-L^{\beta} / 4}\right) . \tag{2.39}
\end{equation*}
$$

Notice that (2.38) and (2.37) also imply that (2.39) (at the expense of possibly changing the constant $C$ ) also holds in case (1) and case (2). Hence, for $1 \leq l, l^{\prime} \leq l_{0}$, we have

$$
\begin{equation*}
C^{-\left|l^{\prime}-l\right|}\left\|C_{l^{\prime}} u\right\|-C^{\left|l^{\prime}-l\right|} e^{-L^{\beta} / 4} \leq\left\|C_{l} u\right\| \leq C^{\left|l^{\prime}-l\right|}\left\|C_{l^{\prime}} u\right\|+C^{\left|l^{\prime}-l\right|} e^{-L^{\beta} / 4} \tag{2.40}
\end{equation*}
$$

If case (1) in the above alternative never holds i.e. if for $1 \leq l \leq l_{0}$, one has $\operatorname{dist}\left(\mathcal{C}_{l}, \mathcal{C}_{l+1}\right) \leq 2$, then, one has $\# \mathcal{C} \geq L / 3$.
We know that $\|C u\|=1+O\left(e^{-L^{\beta} / 3}\right)$. So, for $L$ sufficiently large, there exists $1 \leq l^{*} \leq l_{0}$ such that

$$
\left\|C_{l^{*} u}\right\| \geq\left(2 \sqrt{\ell_{0}}\right)^{-1} \geq(4 \sqrt{L})^{-1}
$$

Hence, by (2.40), either of two things occur

- for some $l$, one has $\left\|C_{l} u\right\| \leq e^{-L^{\beta} / 5}$, then $\left|l-l^{*}\right| \geq \tilde{c} L^{\beta}$ for some $\tilde{c}>0$; thus, $l_{0} \geq 2 \tilde{c} L^{\beta}$; and for some $0<c<\tilde{c}$, for $\left|l-l^{*}\right| \leq c L^{\beta}$, one has $\left\|C_{l} u\right\| \geq e^{-L^{\beta} / 5}$.
- for $1 \leq l \leq l_{0}$, one has $\left\|C_{l} u\right\| \geq e^{-L^{\beta} / 5}$; then, case (1) never occurs, thus, by the observation made above, $\# \mathcal{C} \geq L / 3$
Finally, notice that, by (2.35), (2.34) and the form of $u^{l}$ (see Lemma 2.8), $\left\|C_{l} u\right\| \geq e^{-L^{\beta} / 5}$ implies that $\left|u_{n}\right| \geq e^{-L^{\beta} / 6}$ for $n \in \mathcal{C}_{l}$.
This completes the proof of Lemma 2.9.
We now show that our characterization of $u$, a solution of (2.22), imposes very restrictive conditions on the random variables $\left(\omega_{\gamma}\right)_{-L \leq \gamma \leq L}$.
If $\gamma$ is inside one of the connected components of $\mathcal{C}$, say $\mathcal{C}_{l}$, that is, if $\{\gamma-1, \gamma, \gamma+1\} \subset \mathcal{C}_{l}$, then, by the first equation in (2.22), we know that

$$
\left|u_{\gamma+1}+u_{\gamma-1}-\Delta E u_{\gamma}\right| \leq C e^{-L^{\beta} / 3} .
$$

Plugging this into (2.36), the eigenvalue equation for $u$, we get

$$
\left|\left(\omega_{\gamma}-\bar{E}\right) u_{\gamma}\right| \leq e^{-L^{\beta} / 4} .
$$

Hence, if $\gamma$ belongs to one of the $\left(\mathcal{C}_{l}\right)_{l}$ singled out in Lemma 2.9, the lower bound for $\left|u_{\gamma}\right|$ given in Lemma 2.9 yields

$$
\begin{equation*}
\left|\omega_{\gamma}-\bar{E}\right| \leq C e^{-L^{\beta} / 12} \tag{2.41}
\end{equation*}
$$

Now, if $n_{l}>2$, there exists $\gamma \in \mathcal{C}_{l}$ such that $\{\gamma-1, \gamma, \gamma+1\} \subset \mathcal{C}_{l}$. On the other hand, if $n_{l}=2$, then, the approximate eigenvalue equation on $\mathcal{C}_{l}$ reads

$$
\left\|\left(\begin{array}{cc}
E & \omega_{\gamma_{l}^{-}} \\
\omega_{\gamma_{l}^{+}} & E
\end{array}\right)\binom{u_{\gamma_{l}^{+}}}{u_{\gamma_{l}^{-}}}\right\| \leq C \sqrt{L} e^{-L^{\beta} / 2} .
$$

So, if $\left\|C_{l} u\right\| \geq e^{-L^{\beta} / 6}$, one has

$$
\begin{equation*}
\left|1-\left(\omega_{\gamma_{l}^{-}}-E\right)\left(\omega_{\gamma_{l}^{+}}-E\right)\right| \leq C e^{-L^{\beta} / 3} . \tag{2.42}
\end{equation*}
$$

Hence, we see that the random variables must satisfy at least $c L^{\beta}$ distinct conditions of the type (2.41) or (2.42). As the random variables are supposed to be independent, identically distributed with a bounded density, these condition imply that (2.20) can occur with a given partition $\mathcal{P}$ and $\mathcal{Q}$ with a probability at most, $e^{-c L^{2 \beta}}$ for some $c>0$. As the total number of
partitions is bounded by $2^{L}$ and as $\beta>1 / 2$, we obtain that, $\mathbb{P}$, the probability that (2.20) holds, is bounded by (2.12). This completes the proof of Lemma 2.4.

Remark 2.1. The estimate (2.12) can be improved as, actually, not all partitions are allowed as we saw in the course of the proof. Moreover, it is sufficient to assume that the distribution function of the random variables be Hölder continuous for the method to work.

Remark 2.2. We now present a natural weaker analogue of point (2) in Lemma 2.4. Fix $\rho>0$ and define

$$
\Delta \mathcal{E}_{L}^{c}=\bigcup_{l=0}^{L} \sigma\left(-C_{l} \Delta C_{l}\right)+\left[-L^{-\rho}, L^{-\rho}\right] .
$$

then, for $\rho>3$, one has $\left|\Delta \mathcal{E}_{L}^{c}\right| \leq 2 L^{2-\rho}$, thus,

$$
\left|\bigcap_{n \geq 1} \bigcup_{L \geq n} \Delta \mathcal{E}_{L}^{c}\right|=0
$$

Define the set of total measure

$$
\Delta \mathcal{E}=\mathbb{R} \backslash\left(\bigcap_{N \geq 1} \bigcup_{L \geq N} \Delta \mathcal{E}_{L}^{c}\right)
$$

Hence, if $E-E^{\prime}=\Delta E \in \Delta \mathcal{E}$, for $L$ sufficiently large, as

$$
\inf _{1 \leq l \leq L} \operatorname{dist}\left(\Delta E, \sigma\left(-C_{l} \Delta C_{l}\right)\right) \geq L^{-\rho},
$$

by the decomposition (2.27), a solution $u$ to the first equation in (2.22) must satisfy $\|u\| \leq$ $L^{-(\nu-\rho)}$ if $\left\|h_{1}\right\| \leq L^{-\nu}$. Hence, we obtain

Lemma 2.10. Fix $\nu>4$. For the discrete Anderson model in dimension 1, for $E-E^{\prime} \in \Delta \mathcal{E}$, for $L$ sufficiently large, if $E_{j}(\omega)$ and $E_{k}(\omega)$ are simple eigenvalues of $H_{\omega}\left(\Lambda_{L}\right)$ such that $\left|E_{k}(\omega)-E\right|+\left|E_{j}(\omega)-E^{\prime}\right| \leq L^{-\nu}$ then $\left\|\nabla_{\omega}\left(E_{j}(\omega)-E_{k}(\omega)\right)\right\|_{1} \geq L^{-\nu}$.

This can then be used as Lemma 2.4 is used in the proof of Lemma 1.1 to prove the following variant of the decorrelation estimates in dimension 1

Lemma 2.11. Assume $d=1$. For $\alpha \in(0,1)$ and $E-E^{\prime} \in \Delta \mathcal{E}$ s.t. $\left\{E, E^{\prime}\right\} \subset I$, for any $c>0$, there exists $C>0$ such that, for $L \geq 3$ and $c L^{\alpha} \leq \ell \leq L^{\alpha} / c$, one has

$$
\mathbb{P}\left(\left\{\begin{array}{l}
\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap\left(E+L^{-d}(-1,1)\right) \neq \emptyset, \\
\sigma\left(H_{\omega}\left(\Lambda_{\ell}\right)\right) \cap\left(E^{\prime}+L^{-d}(-1,1)\right) \neq \emptyset
\end{array}\right\}\right) \leq C(\ell / L)^{2 d}(\log L)^{C} .
$$

Comparing with Lemma 1.1, we improved the bound on the probability at the expense of reducing the set of validity in $\left(E, E^{\prime}\right)$.
2.4. Proof of Lemma 2.5. Pick $(u, v) \in\left(\mathbb{R}^{+}\right)^{2 n}$ such that $\|u\|_{1}=\|v\|_{1}=1$. At the expense of exchanging $u$ and $v$, we may assume that $\|v\|_{2} \geq\|u\|_{2}$. Write $u=\alpha v+v^{\perp}$ where $\left\langle v, v^{\perp}\right\rangle=0$. Note that, as all the coefficient of both $u$ and $v$ are non negative, $v^{\perp}=0$ is equivalent $u=v$. Let us now assume $u \neq v$ that is $v^{\perp} \neq 0$. One computes

$$
\begin{equation*}
\|u\|_{2}^{2}=\alpha^{2}\|v\|_{2}^{2}+\left\|v^{\perp}\right\|_{2}^{2} \text { and }\|u-v\|_{2}^{2}=(\alpha-1)^{2}\|v\|_{2}^{2}+\left\|v^{\perp}\right\|_{2}^{2} . \tag{2.43}
\end{equation*}
$$

Moreover, as all the coefficients of $v$ are non negative, $v^{\perp}$ admits at least one negative coefficient. As all the coefficients of $u$ are non negative, the decomposition $u=\alpha v+v^{\perp}$
implies that $\alpha>0$. The first equation in (2.43) and the condition $\|v\|_{2} \geq\|u\|_{2}$ then imply $\alpha \in(0,1)$. Combining this with $u=\alpha v+v^{\perp}$ and $\|u\|_{1}=\|v\|_{1}=1$ yields

$$
0<1-\alpha \leq\left\|v^{\perp}\right\|_{1} .
$$

Hence, by the second equation in (2.43) and the Cauchy-Schwartz inequality, we get

$$
\begin{equation*}
\frac{1}{\sqrt{n}}\|u-v\|_{1} \leq\|u-v\|_{2} \leq\|v\|_{2}\left\|v^{\perp}\right\|_{1}+\left\|v^{\perp}\right\|_{2} \leq 2 \sqrt{n}\left\|v^{\perp}\right\|_{2} \tag{2.44}
\end{equation*}
$$

For any $(j, k)$, one has

$$
\left|\begin{array}{cc}
u_{j} & u_{k} \\
v_{j} & v_{k}
\end{array}\right|=\left|\begin{array}{cc}
v_{j}^{\perp} & v_{k}^{\perp} \\
v_{j} & v_{k}
\end{array}\right| .
$$

As $\left\langle v, v^{\perp}\right\rangle=0$, one computes

$$
\begin{aligned}
\sum_{j, k}\left|\begin{array}{cc}
u_{j} & u_{k} \\
v_{j} & v_{k}
\end{array}\right|^{2} & =\sum_{j, k}\left(\left(v_{j} v_{k}^{\perp}\right)^{2}+\left(v_{k} v_{j}^{\perp}\right)^{2}-2 v_{j} v_{k}^{\perp} v_{k} v_{j}^{\perp}\right) \\
& =2\left(\sum_{j} v_{j}^{2}\right)\left(\sum_{k}\left(v_{k}^{\perp}\right)^{2}\right)-2\left(\sum_{j} v_{j} v_{j}^{\perp}\right)\left(\sum_{k} v_{k} v_{k}^{\perp}\right) \\
& =2\|v\|_{2}^{2}\left\|v^{\perp}\right\|_{2}^{2} \geq \frac{1}{2 n^{3}}\|u-v\|_{1}^{2} .
\end{aligned}
$$

Thus,

$$
\max _{j \neq k}\left|\begin{array}{ll}
u_{j} & u_{k} \\
v_{j} & v_{k}
\end{array}\right|^{2} \geq \frac{1}{4 n^{5}}\|u-v\|_{1}^{2}
$$

which completes the proof of Lemma 2.5.

## 3. The proofs of Theorems 1.2 and 1.3

In [7], the authors extensively study the distribution of the energy levels of random systems in the localized phase. Their results apply also to the discrete Anderson model; in particular, they provide a proof of Theorems 1.2 and 1.3 once the decorrelation estimates obtained in Lemmas 1.1 and 1.2 are known. We provide an alternate proof. The proof in [7] relies on a construction that also proves Theorem 1.1 (actually a stronger uniform result). Here, we only prove Theorems 1.2 and 1.3 independently of the values of the limits in Theorem 1.1. The localization centers of Proposition 2.1 are not defined uniquely. One can easily check that, under the assumptions of Proposition 2.1, all the localization centers for a given eigenvalue or eigenfunction are contained in a disk of radius at most $C \log L$ (for some $C>0$ ). To define a unique localization center, we order the centers lexicographically and let the localization center associated to the eigenvalue or eigenfunction be the largest one (i.e. the one most upper left in dimension 2).
We prove
Lemma 3.1. Pick $\alpha \in(0,1)$ and $c>0$. Let $\nu$ be defined by (Loc). Assume $\ell=\ell(L)$ satisfies $c L^{\alpha} \leq \ell \leq L^{\alpha} / c$.
If (Loc) (see Proposition 2.1) is satisfied then, for any $p>0$ and $\varepsilon>0$, there exists $L_{0}>0$ such that, for $L \geq L_{0}$, with probability larger than $1-L^{-p}$,
(1) if $\left(E_{j}\right)_{1 \leq j \leq J} \in I^{J}$ are eigenvalues of $H_{\omega}\left(\Lambda_{L}\right)$ with localization center in $\gamma+\Lambda_{\ell}$, then the operator $H_{\omega}\left(\gamma+\Lambda_{\ell(1+\varepsilon)}\right)$ has $J$ eigenvalues, say $\left(\tilde{E}_{j}\right)_{1 \leq j \leq J}$, with localization center in $\gamma+\Lambda_{\ell(1+\varepsilon / 2)}$ and such that $\sup _{1 \leq j \leq J}\left|E_{j}-\tilde{E}_{j}\right| \leq e^{-\nu \varepsilon \ell / 4}$.
(2) if $\left(E_{j}\right)_{1 \leq j \leq J} \in I^{J}$ are eigenvalues of $H_{\omega}\left(\gamma+\Lambda_{\ell(1+\varepsilon)}\right)$ with localization center in $\gamma+\Lambda_{\ell}$, then the operator $H_{\omega}\left(\Lambda_{L}\right)$ has $J$ eigenvalues, say $\left(\tilde{E}_{j}\right)_{1 \leq j \leq J}$, with localization center in $\gamma+\Lambda_{\ell(1+\varepsilon / 2)}$ and such that $\sup _{1 \leq j \leq J}\left|E_{j}-\tilde{E}_{j}\right| \leq e^{-\nu \varepsilon \ell / 4}$.
(3) if $\left(E_{j}\right)_{1 \leq j \leq J} \in I^{J}$ are eigenvalues of $H_{\omega}\left(\gamma+\Lambda_{\ell(1+\varepsilon)}\right)$ with localization center in $\gamma+$ $\left(\Lambda_{\ell(1+\varepsilon / 2)} \backslash \Lambda_{\ell}\right)$, then there exists $\left(\beta_{j}\right)_{1 \leq j \leq J}$ such that, for $1 \leq j \leq J$, one has

- $\beta_{j} \in \frac{\varepsilon \ell}{16} \mathbb{Z}^{d} \cap\left[\gamma+\left(\Lambda_{\ell(1+\varepsilon / 2)} \backslash \Lambda_{\ell}\right)\right]$,
- the operator $H_{\omega}\left(\beta_{j}+\Lambda_{\varepsilon \ell / 4}\right)$ has an eigenvalue, say $\tilde{E}_{j}$, satisfying $\left|E_{j}-\tilde{E}_{j}\right| \leq$ $e^{-\nu \varepsilon \ell / 8}$.
The number $\nu>0$ is given by (Loc).
Similar results can be found in [7].
Proof. With probability at least $1-L^{-p}$, the conclusions of Proposition 2.1 hold which we assume from now on.
To prove (1), let $\left(\varphi_{j}\right)_{1 \leq j \leq J}$ be normalized eigenfunctions associated to $\left(E_{j}\right)_{1 \leq j \leq J}$. Then, setting $\tilde{\varphi}_{j}=1_{\gamma+\Lambda_{\ell(1+\varepsilon)}} \varphi_{j}$ and using (2.3) from (Loc) and the assumption that the localization center are in $\gamma+\Lambda_{\ell}$, one obtains

$$
\begin{gathered}
\left\|\left(\left(\left\langle\tilde{\varphi}_{j}, \tilde{\varphi}_{k}\right\rangle_{\ell^{2}\left(\gamma+\Lambda_{\ell(1+\varepsilon)}\right)}\right)\right)_{\substack{1 \leq j \leq J \\
1 \leq k \leq J}}-\mathrm{Id}\right\| \leq J^{2} e^{-\nu \varepsilon \ell / 4}, \\
\sup _{1 \leq j \leq J}\left\|\mathbf{1}_{\gamma+\left(\Lambda_{\ell(1+\varepsilon)} \backslash \Lambda_{\ell(1+\varepsilon / 2)}\right)} \tilde{\varphi}_{j}\right\|_{\ell^{2}\left(\gamma+\Lambda_{\ell(1+\varepsilon)}\right)} \leq e^{-\nu \varepsilon \ell / 6}, \\
\sup _{1 \leq j \leq J}\left\|\left(H_{\omega}\left(\gamma+\Lambda_{\ell(1+\varepsilon)}\right)-E_{j}\right) \tilde{\varphi}_{j}\right\|_{\ell^{2}\left(\gamma+\Lambda_{\ell(1+\varepsilon)}\right)} \leq e^{-\nu \varepsilon \ell / 4} .
\end{gathered}
$$

This immediately yields (1) for $L$ sufficiently large as

- $J \leq(2 L+1)^{d}$ and $c L^{\alpha} \leq \ell \leq L^{\alpha} / c$,
- at a localization center, the modulus of an eigenfunction is at least of order $L^{-d / 2}$.

Points (2) is proved in the same way. We omit further details.
To prove (3), we set $\tilde{\varphi}_{j}=\mathbf{1}_{\beta_{j}+\Lambda_{\ell \varepsilon / 4}} \varphi_{j}$ where $\beta_{j}$ is the point in $\frac{\varepsilon \ell}{16} \mathbb{Z}^{d}$ closest to the localization center of $\varphi_{j}$. The conclusion then follows from the same reasoning as above.
This completes the proof of Lemma 3.1.
Pick $\left(U_{j}\right)_{1 \leq j \leq J},\left(k_{j}\right)_{1 \leq j \leq J} \in \mathbb{N}^{J},\left(U_{j}^{\prime}\right)_{1 \leq j \leq J^{\prime}}$ and $\left(k_{j}^{\prime}\right)_{1 \leq j \leq J} \in \mathbb{N}^{J^{\prime}}$ as in Theorem 1.2. To prove Theorems 1.2 and 1.3, it suffices to prove (1.6) for $\left(U_{j}\right)_{1 \leq j \leq J}$ and $\left(U_{j}^{\prime}\right)_{1 \leq j \leq J^{\prime}}$ non empty compact intervals which we assume from now on.
Pick $L$ and $\ell$ such that $(2 L+1)=(2 \ell+1)\left(2 \ell^{\prime}+1\right), c L^{\alpha} \leq \ell \leq L^{\alpha} / c$ for some $\alpha \in(0,1)$ and $c>0$. Pick $\varepsilon>0$ small. Partition $\Lambda_{L}=\bigcup_{|\gamma| \leq \ell^{\prime}} \Lambda_{\ell}(\gamma)$ where $\Lambda_{\ell}(\gamma)=(2 \ell+1) \gamma+\Lambda_{\ell}$. For $\Lambda^{\prime} \subset \Lambda$ and $U \subset \mathbb{R}$, consider the random variables

$$
X\left(E, U, \Lambda, \Lambda^{\prime}\right):= \begin{cases}1 & \text { if } H_{\omega}(\Lambda) \text { has at least one eigenvalue in } \\ & E+(\nu(E)|\Lambda|)^{-1} U \text { with localization center in } \Lambda^{\prime}, \\ 0 & \text { if not; }\end{cases}
$$

if $\Lambda^{\prime}=\Lambda$, we write $X(E, U, \Lambda):=X(E, U, \Lambda, \Lambda)$, and

$$
\Sigma(E, U):=\sum_{|\gamma| \leq \ell^{\prime}} X\left(E, U, \Lambda, \Lambda_{\ell}(\gamma)\right), \quad \Sigma(E, U, \ell):=\sum_{|\gamma| \leq \ell^{\prime}} X\left(E, U, \Lambda_{\ell}(\gamma)\right) .
$$

We prove

## Lemma 3.2.

$$
\begin{aligned}
& \left|\mathbb{P}\left(\left\{\begin{array}{c}
\#\left\{j ; \xi_{n}(E, \omega, \Lambda) \in U_{1}\right\}=k_{1} \\
\left.\omega ; \quad \begin{array}{c}
\# \\
\#\left\{j ; \xi_{n}(E, \omega, \Lambda) \in U_{J}\right\}=k_{J}
\end{array}\right\}
\end{array}\right\}\right)-\mathbb{P}\left(\left\{\begin{array}{c}
\Sigma\left(E, U_{1}, \ell\right)=k_{1}, \\
\omega ; \\
\vdots\left(E, U_{J}, \ell\right)=k_{J}
\end{array}\right\}\right)\right| \underset{L \rightarrow+\infty}{\rightarrow} 0
\end{aligned}
$$

and

$$
\left|\mathbb{P}\left(\left\{\begin{array}{c}
\left.\begin{array}{c}
\#\left\{j ; \xi_{n}\left(E^{\prime}, \omega, \Lambda\right) \in U_{1}^{\prime}\right\}=k_{1}^{\prime} \\
\omega ; \\
\vdots \\
\#\left\{j ; \xi_{n}\left(E^{\prime}, \omega, \Lambda\right) \in U_{J^{\prime}}\right\}
\end{array}\right\}=k_{J^{\prime}}
\end{array}\right\}\right)-\mathbb{P}\left(\left\{\begin{array}{c}
\Sigma\left(E^{\prime}, U_{1}^{\prime}, \ell\right)=k_{1}^{\prime} \\
\omega ; \\
\vdots \\
\Sigma\left(E^{\prime}, U_{J^{\prime}}, \ell\right)=k_{J^{\prime}}^{\prime}
\end{array}\right\}\right)\right| \underset{L \rightarrow+\infty}{\rightarrow} 0 .
$$

Proof. We first prove
Lemma 3.3. For any $p>0$ and $\varepsilon>0$, there exists $C>0$ such that, for $U$ a compact interval and $L$ sufficiently large, one has

$$
\begin{equation*}
\mathbb{P}\left(\left\{\omega ; \#\left\{n ; \xi_{n}(E, \omega, \Lambda) \in U\right\} \neq \Sigma(E, U)\right\}\right) \leq C \ell^{d} L^{-d}(|U|+1)^{2}+L^{-p} . \tag{3.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathbb{P}(\{\omega ; \Sigma(E, U) \neq \Sigma(E, U, \ell)\}) \leq L^{-p}+C \varepsilon|U| . \tag{3.2}
\end{equation*}
$$

Proof of Lemma 3.3. As $\Lambda=\bigcup_{|\gamma| \leq \ell^{\prime}} \Lambda_{\ell}(\gamma)$ and these sets are two by two disjoint, the quantities $\#\left\{n ; \xi_{n}(E, \omega, \Lambda) \in U\right\}$ and $\Sigma(E, U)$ differ if and only if, for some $|\gamma| \leq \ell^{\prime}, H_{\omega}(\Lambda)$ has at least two eigenvalues in $E+(\nu(E)|\Lambda|)^{-1} U$ with localization center in $\Lambda_{\ell}(\gamma)$. By Lemma 3.1, this implies that, except on a set of probability at most $L^{-p}, H_{\omega}\left((2 \ell+1) \gamma+\Lambda_{2 \ell}\right)$ has at least two eigenvalues in $U+\left[-e^{-\nu \ell / 4}, e^{-\nu \ell / 4}\right]$. Thus, by Minami's estimate (2.2), this happens with a probability at most $C \ell^{2 d} L^{-2 d}(|U|+1)^{2}+L^{-p}$. Summing this estimate over all the possible $\gamma$ 's, we complete the proof of (3.1).
The proof of (3.2) is split into two steps. Define

$$
\Sigma(E, U, \varepsilon)=\sum_{|\gamma| \leq \ell^{\prime}} X\left(E, U,(2 \ell+1) \gamma+\Lambda_{\ell(1+\varepsilon)},(2 \ell+1) \gamma+\Lambda_{\ell(1-\varepsilon)}\right) .
$$

Then, we successively prove

$$
\begin{equation*}
\mathbb{P}(\{\omega ; \Sigma(E, U) \neq \Sigma(E, U, \varepsilon)\}) \leq L^{-p}+C \varepsilon|U| \tag{3.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathbb{P}(\{\omega ; \Sigma(E, U, \varepsilon) \neq \Sigma(E, U, \ell)\}) \leq L^{-p}+C \varepsilon|U| \tag{3.4}
\end{equation*}
$$

which implies (3.2).
To prove (3.3), we note that, by Lemma 3.1, except on a set of probability at most $L^{-p}$, $\Sigma(E, U)$ and $\Sigma(E, U, \varepsilon)$ differ if and only if, for some $|\gamma| \leq \ell^{\prime}$, one has
(1) either $\sigma\left(H_{\omega}(\Lambda)\right) \cap \delta \tilde{U} \neq \emptyset$,
(2) or $\sigma\left(H_{\omega}\left((2 \ell+1) \gamma+\Lambda_{\ell(1+\varepsilon)}\right)\right) \cap \delta \tilde{U} \neq \emptyset$,
(3) or $H_{\omega}\left((2 \ell+1) \gamma+\Lambda_{\ell(1+\varepsilon)}\right)$ has an eigenvalue in $\tilde{U}$ with a localization center in the cube $(2 \ell+1) \gamma+\left(\Lambda_{\ell(1+\varepsilon)} \backslash \Lambda_{\ell(1-\varepsilon)}\right)$
where $\tilde{U}=E+(\nu(E)|\Lambda|)^{-1} U+e^{-\nu \varepsilon \ell / 8}[-1,1]$ and $\delta \tilde{U}=\tilde{U} \backslash\left(E+(\nu(E)|\Lambda|)^{-1} U\right)$.
The probability of alternatives (1) and (2) is estimated using the Wegner estimate (2.1). It is bounded by $2 L^{d} e^{-\nu \varepsilon \ell / 8} \leq L^{-p}$ for $L$ sufficiently large.
By point (3) of Lemma 3.1, except on a set of probability at most $L^{-p}$, alternative (3) implies that, for some $\beta \in \gamma+\left(\Lambda_{\ell(1+\varepsilon / 2)} \backslash \Lambda_{\ell}\right)$, the operator $H_{\omega}\left(\beta+\Lambda_{\varepsilon \ell / 4}\right)$ has an eigenvalue in $E+(\nu(E)|\Lambda|)^{-1} U+e^{-\nu \varepsilon \ell / 8}[-1,1]$. The number of possible $\beta$ 's is bounded by $C \varepsilon \ell^{d} \varepsilon^{-d} \ell^{-d}=$ $C \varepsilon^{1-d}$. Using Wegner's estimate (2.1) and summing over the possible $\beta$ 's, this probability is bounded by $C \varepsilon^{1-d}(\varepsilon \ell / L)^{d}|U|+L^{-p} \leq C \varepsilon(\ell / L)^{d}|U|+L^{-p}$. Finally, we sum this over all possible $\gamma$ 's to obtain that the probability that alternative (3) holds for some $\gamma$ is bounded by $C \varepsilon|U|+L^{-p}$. This yields (3.3).
To prove (3.8), the reasoning is similar. By Lemma 3.1, except on a set of probability at most $L^{-p}, \Sigma(E, U, \ell)$ and $\Sigma(E, U, \varepsilon)$ differ if and only if, for some $|\gamma| \leq \ell^{\prime}$, one has
(1) either $\sigma\left(H_{\omega}\left((2 \ell+1) \gamma+\Lambda_{\ell(1+\varepsilon)}\right)\right) \cap \tilde{U} \neq \emptyset$,
(2) or $\sigma\left(H_{\omega}\left(\Lambda_{\ell}(\gamma)\right)\right) \cap \tilde{U} \neq \emptyset$,
(3) or $H_{\omega}\left((2 \ell+1) \gamma+\Lambda_{\ell(1+\varepsilon)}\right)$ has an eigenvalue in $\tilde{U}$ with localization center in the cube $(2 \ell+1) \gamma+\left(\Lambda_{\ell(1+\varepsilon)} \backslash \Lambda_{\ell(1-\varepsilon)}\right)$.
(4) or $H_{\omega}\left(\Lambda_{\ell}(\gamma)\right)$ has an eigenvalue in $\tilde{U}$ with localization center in $(2 \ell+1) \gamma+\left(\Lambda_{\ell} \backslash \Lambda_{\ell(1-\varepsilon)}\right)$.

Following the same steps as in the proof of (3.3), we obtain (3.8). We omit further details. This completes the proof of Lemma 3.3.

As $\varepsilon>0$ can be chosen arbitrarily small and $J$ and $J^{\prime}$ are finite and fixed, Lemma 3.3 clearly implies Lemma 3.2.

In view of Theorem 1.1 and Lemma 3.2, to prove (1.6), it suffices to prove that, in the limit $L \rightarrow+\infty$, the difference between the following quantities vanishes

$$
\mathbb{P}\left(\left\{\omega ; \begin{array}{c}
\Sigma\left(E, U_{1}, \ell\right)=k_{1}, \cdots, \Sigma\left(E, U_{J}, \ell\right)=k_{J}  \tag{3.5}\\
\Sigma\left(E^{\prime}, U_{1}^{\prime}, \ell\right)=k_{1}^{\prime}, \cdots, \Sigma\left(E^{\prime}, U_{J^{\prime}}, \ell\right)=k_{J^{\prime}}^{\prime}
\end{array}\right\}\right)
$$

and

$$
\mathbb{P}\left(\left\{\begin{array}{cc}
\Sigma\left(E, U_{1}, \ell\right)=k_{1}  \tag{3.6}\\
\omega ; & \vdots \\
\Sigma\left(E, U_{J}, \ell\right)=k_{J}
\end{array}\right\}\right) \mathbb{P}\left(\left\{\begin{array}{c}
\Sigma\left(E^{\prime}, U_{1}^{\prime}, \ell\right)=k_{1}^{\prime} \\
\omega ; \\
\vdots \\
\Sigma\left(E^{\prime}, U_{J^{\prime}}, \ell\right)=k_{J^{\prime}}^{\prime}
\end{array}\right\}\right) .
$$

Both terms in (3.5) and (3.6) define probability measures on $\mathbb{N}^{J+J^{\prime}}$. By Theorems 1.1 and Lemma 3.2, we know that the limit of the term in (3.6) also defines a probability measure on $\mathbb{N}^{J+J^{\prime}}$. Thus, by standard results on the convergence of probability measures (see e.g. [3]), the difference of (3.5) and (3.6) vanishes in the limit $L \rightarrow+\infty$ if and only if, for any $\left(t_{j}\right)_{1 \leq j \leq J}$ and $\left(t_{j^{\prime}}\right)_{1 \leq j^{\prime} \leq J^{\prime}}$ real, in the limit $L \rightarrow+\infty$, the following quantity vanishes

$$
\begin{aligned}
\mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} \Sigma\left(E, U_{j}, \ell\right)-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} \Sigma\left(E^{\prime}, U_{j^{\prime}}, \ell\right)}\right) & \\
& -\mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} \Sigma\left(E, U_{j}, \ell\right)}\right) \mathbb{E}\left(e^{-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} \Sigma\left(E^{\prime}, U_{j^{\prime}}, \ell\right)}\right)
\end{aligned}
$$

Note that, as the sets $\left(\Lambda_{\ell}(\gamma)\right)_{|\gamma| \leq \ell^{\prime}}$ are two by two disjoint and translates of each other, for a fixed $U$, the random variables $\left(X\left(E, U, \Lambda_{\ell}(\gamma)\right)_{|\gamma| \leq \ell^{\prime}}\right.$ are i.i.d. Bernoulli random variables. Thus,

$$
\begin{aligned}
& \mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} \Sigma\left(E, U_{j}, \ell\right)-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} \Sigma\left(E^{\prime}, U_{j^{\prime}}, \ell\right)}\right) \\
&=\prod_{|\gamma| \leq \ell^{\prime}} \mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)}\right)
\end{aligned}
$$

The Minami estimate (2.2) and the decorrelation estimates (1.7) and (1.8) of Lemmas 1.1 and 1.2 guarantee that, for any $\rho \in(0,1)$, one has, for some $C>0$ independent of $\gamma$,

$$
\begin{array}{rl}
\sup _{1 \leq j<\tilde{j} \leq J} & \mathbb{P}\binom{X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)=1}{X\left(E, U_{\tilde{j}}, \Lambda_{\ell}(\gamma)\right)=1} \\
+ & \sup _{1 \leq j^{\prime}<\dot{j}^{\prime} \leq J^{\prime}} \mathbb{P}\binom{X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)=1}{X\left(E^{\prime}, U_{\tilde{j}^{\prime}}, \Lambda_{\ell}(\gamma)\right)=1}  \tag{3.7}\\
& \quad+\sup _{\substack{1 \leq j \leq J \\
1 \leq j^{\prime} \leq J^{\prime}}} \mathbb{P}\binom{X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)=1}{X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)=1} \leq C\left(\frac{\ell}{L}\right)^{d(1+\rho)} .
\end{array}
$$

Using this, we compute

$$
\begin{align*}
& \mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)}\right) \\
&= 1+\sum_{j=1}^{J}\left(e^{-t_{j}}-1\right) \cdot \mathbb{P}\left(X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)=1\right)  \tag{3.8}\\
&+\sum_{j^{\prime}=1}^{J^{\prime}}\left(e^{-t_{j^{\prime}}}-1\right) \cdot \mathbb{P}\left(X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)=1\right)+O\left(\left(\frac{\ell}{L}\right)^{d(1+\rho)}\right)
\end{align*}
$$

Here, the term $O\left((\ell / L)^{d(1+\rho)}\right)$ is uniform in $\gamma$.
On the other hand, one has

$$
\begin{array}{r}
\mathbb{E}\left(e^{-t_{j} X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)}\right)=1+\left(e^{-t_{j}}-1\right) \cdot \mathbb{P}\left(X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)=1\right),  \tag{3.9}\\
\mathbb{E}\left(e^{-t_{j^{\prime}} X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)}\right)=1+\left(e^{-t_{j^{\prime}}}-1\right) \cdot \mathbb{P}\left(X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)=1\right) .
\end{array}
$$

By the Wegner estimate (2.1), we know that

$$
\begin{equation*}
\sup _{\substack{1 \leq j \leq J \\ 1 \leq j^{\prime} \leq J^{\prime}}}\left[\mathbb{P}\left(X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)=1\right)+\mathbb{P}\left(X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)=1\right)\right] \leq C\left(\frac{\ell}{L}\right)^{d} \tag{3.10}
\end{equation*}
$$

Thus, by (3.8), we have

$$
\begin{aligned}
& \mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)}\right) \\
& \quad=\prod_{j=1}^{J} \mathbb{E}\left(e^{-t_{j} X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)}\right) \prod_{j^{\prime}=1}^{J^{\prime}} \mathbb{E}\left(e^{-t_{j^{\prime}} X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)}\right)\left[1+O\left(\left(\ell L^{-1}\right)^{d(1+\rho)}\right)\right]
\end{aligned}
$$

In the same way, one proves

$$
\begin{array}{r}
\mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)}\right)=\prod_{j=1}^{J} \mathbb{E}\left(e^{-t_{j} X\left(E, U_{j}, \Lambda_{\ell}(\gamma)\right)}\right)\left[1+O\left(\left(\ell L^{-1}\right)^{d(1+\rho)}\right)\right] \\
\mathbb{E}\left(e^{-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)}\right)=\prod_{j^{\prime}=1}^{J^{\prime}} \mathbb{E}\left(e^{-t_{j^{\prime}} X\left(E^{\prime}, U_{j^{\prime}}, \Lambda_{\ell}(\gamma)\right)}\right)\left[1+O\left(\left(\ell L^{-1}\right)^{d(1+\rho)}\right)\right] \tag{3.11}
\end{array}
$$

As $\#\left\{|\gamma| \leq \ell^{\prime}\right\} \leq C\left(L \ell^{-1}\right)^{d}$, we obtain that

$$
\begin{align*}
& \mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} \Sigma\left(E, U_{j}, \ell\right)-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} \Sigma\left(E^{\prime}, U_{j^{\prime}}, \ell\right)}\right)  \tag{3.12}\\
& \quad=\mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} \Sigma\left(E, U_{j}, \ell\right)}\right) \mathbb{E}\left(e^{-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} \Sigma\left(E^{\prime}, U_{j^{\prime}}, \ell\right)}\right)\left[1+O\left(\left(\ell L^{-1}\right)^{d \rho}\right)\right]
\end{align*}
$$

Finally, note that (3.9), (3.10) and (3.11) imply that, for any $\left(t_{j}\right)_{1 \leq j \leq J}$ and $\left(t_{j^{\prime}}\right)_{1 \leq j^{\prime} \leq J^{\prime}}$, one has

$$
\sup _{L \geq 1}\left[\mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} \Sigma\left(E, U_{j}, \ell\right)}\right)+\mathbb{E}\left(e^{-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} \Sigma\left(E^{\prime}, U_{j^{\prime}}, \ell\right)}\right)\right]<+\infty
$$

Hence, by (3.12), as $c L^{\alpha} \leq \ell \leq L^{\alpha} / c$ for some $\alpha \in(0,1)$, we obtain that

$$
\begin{aligned}
& \mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} \Sigma\left(E, U_{j}, \ell\right)-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} \Sigma\left(E^{\prime}, U_{j^{\prime}}, \ell\right)}\right) \\
&-\mathbb{E}\left(e^{-\sum_{j=1}^{J} t_{j} \Sigma\left(E, U_{j}, \ell\right)}\right) \mathbb{E}\left(e^{-\sum_{j^{\prime}=1}^{J^{\prime}} t_{j^{\prime}} \Sigma\left(E^{\prime}, U_{j^{\prime}}, \ell\right)}\right) \underset{L \rightarrow+\infty}{\rightarrow} 0 .
\end{aligned}
$$

This completes the proof of Theorems 1.2 and 1.3.
Remark 3.1. The basic idea we used here is to split the cube $\Lambda$ into smaller two-by-two disjoint cubes $\left(\Lambda_{\gamma}(\ell)\right)_{\gamma}$ in such a way that, up to exponentially small errors, the eigenvalues of $H_{\omega}(\Lambda)$ can be represented as eigenvalues for $H_{\omega}\left(\Lambda_{\gamma}(\ell)\right)$ and that they are independent of each other. In [7] (see also [8] for a review of the results), this idea is exploited thoroughly to study the eigenvalue statistics for random operators in the localized regime.

## 4. Proof of Proposition 2.1

Let $I$ be a compact subset of the region of localization i.e. the region of $\Sigma$ where the finite volume fractional moment criteria of [1] for $H_{\omega}(\Lambda)$ are verified for $\Lambda$ sufficiently large. Then, by (A.6) of [1], we know that there exists $\alpha>0$ such that, for any $F \subset I, \forall(x, y) \subset \Lambda^{2}$, one has

$$
\begin{equation*}
\mathbb{E}\left(\left|\mu_{\omega, \Lambda}^{x, y}\right|(F)\right) \leq C e^{-\alpha|x-y|} \tag{4.1}
\end{equation*}
$$

where $\mu_{\omega, \Lambda}^{x, y}$ denotes the spectral measures of $H_{\omega}(\Lambda)$ associated to the vector $\delta_{x}$ and $\delta_{y}$. In particular, if $F$ contains a single eigenvalue of $H_{\omega}(\Lambda)$, say $E$, that is simple and associated to the normalized eigenvector, say, $\varphi$ then

$$
\begin{equation*}
\left|\mu_{\omega, \Lambda}^{x, y}\right|(F)=|\varphi(x)||\varphi(y)| \tag{4.2}
\end{equation*}
$$

Pick $\varepsilon$ and $\delta$ positive such that $\varepsilon|\Lambda|^{2}=\delta / K$ for some large $K$ to be chosen below. Then, partition $I=\cup_{1 \leq n \leq N} I_{n}$ into intervals $\left(I_{n}\right)_{n}$ of length $\varepsilon$. By Minami's estimate, one has

$$
\mathbb{P}\left(\left\{\omega ; \exists n \text { s.t. } I_{n} \text { contains } 2 \text { e.v. of } H_{\omega}(\Lambda)\right\}\right) \leq C \delta|I| / K \leq \delta / 2
$$

if $C|I| / K \leq 1 / 2$. Pick $K$ so that this be satisfied.
We now apply (4.1) to $F=I_{n}$ for $1 \leq n \leq N$ and sum the results for $s<\alpha$ to get

$$
\forall y \in \Lambda, \quad \mathbb{E}\left(\sum_{n} \sum_{x \in \Lambda} e^{s|x-y|}\left|\mu_{\omega, \Lambda}^{x, y}\right|\left(I_{n}\right)\right) \leq C|I| \varepsilon^{-1}
$$

Hence, by Markov's inequality,

$$
\mathbb{P}\left(\sum_{n} \sum_{(x, y) \in \Lambda^{2}} e^{s|x-y|}\left|\mu_{\omega, \Lambda}^{x, y}\right|\left(I_{n}\right) \geq \frac{\tilde{C}|\Lambda||I|}{\delta \varepsilon}\right) \leq \delta / 2
$$

Thus, using the relation between $\delta$ and $\varepsilon$, with a probability larger than $1-\delta$, we know that
(1) each interval $I_{n}$ contain at most a single eigenvalue, say, $E_{n}$ associated to the normalized eigenfunction, say, $\varphi_{n}$;
(2) by (4.2), one has

$$
\forall(x, y) \in \Lambda^{2}, \quad\left|\varphi_{n}(x)\right|\left|\varphi_{n}(y)\right| \leq \frac{C|\Lambda|^{3} e^{-s|x-y|}}{\delta^{2}}
$$

As $\varphi_{n}$ is normalized, if $x_{n}$ is a maximum of $x \mapsto\left|\varphi_{n}(x)\right|$, one has

$$
\left|\varphi_{n}\left(x_{n}\right)\right| \geq|\Lambda|^{-1 / 2}
$$

thus,

$$
\forall x \in \Lambda, \quad\left|\varphi_{n}(x)\right| \leq C|\Lambda|^{7 / 2} \delta^{-2} e^{-s\left|x-x_{n}\right|}
$$

This yields Proposition 2.1 if one picks $\delta=L^{-p}$ when $\Lambda=\Lambda_{L}$.

## References

[1] Michael Aizenman, Jeffrey H. Schenker, Roland M. Friedrich, and Dirk Hundertmark. Finite-volume fractional-moment criteria for Anderson localization. Comm. Math. Phys., 224(1):219-253, 2001. Dedicated to Joel L. Lebowitz.
[2] Jean V. Bellissard, Peter D. Hislop, and Günter Stolz. Correlation estimates in the Anderson model. J. Stat. Phys., 129(4):649-662, 2007.
[3] Patrick Billingsley. Convergence of probability measures. Wiley Series in Probability and Statistics: Probability and Statistics. John Wiley \& Sons Inc., New York, second edition, 1999. A Wiley-Interscience Publication.
[4] Jean-Michel Combes, François Germinet, and Abel Klein. Poisson statistics for eigenvalues of continuum random Schrödinger operators. Preprint.
[5] Jean-Michel Combes, François Germinet, and Abel Klein. Generalized eigenvalue-counting estimates for the Anderson model. J. Stat. Phys., 135(2):201-216, 2009.
[6] Francois Germinet and Abel Klein. New characterizations of the region of complete localization for random Schrödinger operators. J. Stat. Phys., 122(1):73-94, 2006.
[7] François Germinet and Frédéric Klopp. Spectral statistics for random Schrödinger operators in the localized regime. In progress.
[8] François Germinet and Frédéric Klopp. Spectral statistics for the discrete Anderson model in the localized regime. ArXiv http://arxiv.org/abs/1004.1261.
[9] Gian Michele Graf and Alessio Vaghi. A remark on the estimate of a determinant by Minami. Lett. Math. Phys., 79(1):17-22, 2007.
[10] Peter D. Hislop. Lectures on random Schrödinger operators. In Fourth Summer School in Analysis and Mathematical Physics, volume 476 of Contemp. Math., pages 41-131. Amer. Math. Soc., Providence, RI, 2008.
[11] Tosio Kato. Perturbation theory for linear operators. Classics in Mathematics. Springer-Verlag, Berlin, 1995. Reprint of the 1980 edition.
[12] Werner Kirsch. An invitation to random Schrödinger operators. In Random Schrödinger operators, volume 25 of Panor. Synthèses, pages 1-119. Soc. Math. France, Paris, 2008. With an appendix by Frédéric Klopp.
[13] Werner Kirsch and Bernd Metzger. The integrated density of states for random Schrödinger operators. In Spectral theory and mathematical physics: a Festschrift in honor of Barry Simon's 60th birthday, volume 76 of Proc. Sympos. Pure Math., pages 649-696. Amer. Math. Soc., Providence, RI, 2007.
[14] Madan Lal Mehta. Random matrices, volume 142 of Pure and Applied Mathematics (Amsterdam). Elsevier/Academic Press, Amsterdam, third edition, 2004.
[15] Nariyuki Minami. Local fluctuation of the spectrum of a multidimensional Anderson tight binding model. Comm. Math. Phys., 177(3):709-725, 1996.
[16] Nariyuki Minami. Energy level statistics for random operators. In Friedrich Götze, Werner Kirsch, Frédéric Klopp, and Thomas Kriecherbauer, editors, Disordered systems: random Schrödinger operators and random matrices, volume 5 of Oberwolfach Reports, pages 842-844, 2008.
[17] Stanislav Molchanov. The local structure of the spectrum of a random one-dimensional Schrödinger operator. Trudy Sem. Petrovsk., (8):195-210, 1982.
[18] Michael Reed and Barry Simon. Methods of modern mathematical physics. IV. Analysis of operators. Academic Press [Harcourt Brace Jovanovich Publishers], New York, 1978.
[19] Barry Simon. Fine structure of the zeros of orthogonal polynomials. I. A tale of two pictures. Electron. Trans. Numer. Anal., 25:328-368 (electronic), 2006.
[20] Mihai Stoiciu. The statistical distribution of the zeros of random paraorthogonal polynomials on the unit circle. J. Approx. Theory, 139(1-2):29-64, 2006.
[21] Ivan Veselić. Existence and regularity properties of the integrated density of states of random Schrödinger operators, volume 1917 of Lecture Notes in Mathematics. Springer-Verlag, Berlin, 2008.
[22] Franz Wegner. Bounds on the density of states in disordered systems. Z. Phys. B, 44(1-2):9-15, 1981.
(Frédéric Klopp) LAGA, U.M.R. 7539 C.N.R.S, Institut Galilée, Université Paris-Nord, 99 Avenue J.-B. Clément, F-93430 Villetaneuse, France et Institut Universitaire de France

E-mail address: klopp@math.univ-paris13.fr


[^0]:    2000 Mathematics Subject Classification. 81Q10,47B80,60H25,82D30.
    Key words and phrases. random Schrödinger operators, renormalized local eigenvalues, decorrelation estimates.

    The author is supported by the grant ANR-08-BLAN-0261-01.

