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Lower Bounds for the Eigenvalues of the Dirac Operator on Spin^c Manifolds

Roger NAKAD

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Institut Élie Cartan, Université Henri Poincaré, Nancy I, B.P 239
54506 Vandœuvre-Lès-Nancy Cedex, France.

nakad@iecn.u-nancy.fr

Abstract

In this paper, we extend the Hijazi inequality, involving the Energy-Momentum tensor, for the eigenvalues of the Dirac operator on Spin^c manifolds without boundary. The limiting case is then studied and examples are given.

Key words: Spin^c structures, Dirac operator, eigenvalues, Energy-Momentum tensor, perturbed Yamabe operator, conformal geometry.

1 Introduction

On a compact Riemannian spin manifold (M^n, g) of dimension $n \geq 2$, Th. Friedrich [5] showed that any eigenvalue λ of the Dirac operator satisfies

$$\lambda^2 \geq \lambda_1^2 := \frac{n}{4(n-1)} \inf_M S_g, \quad (1)$$

where S_g denotes the scalar curvature of M . The limiting case of (1) is characterized by the existence of a special spinor called real Killing spinor. This is a section ψ of the spinor bundle satisfying for every $X \in \Gamma(TM)$,

$$\nabla_X \psi = -\frac{\lambda_1}{n} X.\psi,$$

where $X.\psi$ denotes the Clifford multiplication and ∇ is the spinorial Levi-Civita connection [20]. In 1995, O. Hijazi [16] established that, on such manifolds any eigenvalue of the Dirac operator satisfies the following

$$\lambda^2 \geq \inf_M \left(\frac{1}{4} S_g + |T^\psi|^2 \right), \quad (2)$$

where T^ψ is the field of symmetric endomorphisms associated with the field of quadratic forms (denoted also by T^ψ) called the Energy-Momentum tensor, defined on the complement set of zeroes of the eigenspinor ψ , for any vector field X by

$$T^\psi(X) = \operatorname{Re} \left\langle X.\nabla_X \psi, \frac{\psi}{|\psi|^2} \right\rangle.$$

The limiting case of (2) is characterized by the existence of a spinor field ψ satisfying for all $X \in TM$,

$$\nabla_X \psi = -T^\psi(X).\psi. \quad (3)$$

The trace of T^ψ being equal to λ , Inequality (2) improves Inequality (1) since by the Cauchy-Schwarz inequality, $|T^\psi|^2 \geq \frac{(\operatorname{tr}(T^\psi))^2}{n}$, where tr denotes the trace of T^ψ . N. Ginoux and G. Habib showed in [9] that the Heisenberg manifold is a limiting manifold for (2) but equality in (1) cannot occur.

Using the conformal covariance of the Dirac operator, O. Hijazi [14] showed that, on a compact Riemannian spin manifold (M^n, g) of dimension $n \geq 3$, any eigenvalue of the Dirac operator satisfies

$$\lambda^2 \geq \frac{n}{4(n-1)} \mu_1, \quad (4)$$

where μ_1 is the first eigenvalue of the Yamabe operator given by

$$L := 4 \frac{n-1}{n-2} \Delta_g + S_g,$$

Δ_g is the Laplacian acting on functions. In dimension 2, C. Bär [2] proved that any eigenvalue of the Dirac operator on M satisfies

$$\lambda^2 \geq \frac{2\pi\chi(M)}{\operatorname{Area}(M, g)}, \quad (5)$$

where $\chi(M)$ is the Euler-Poincaré characteristic of M . The limiting case of (4) and (5) is also characterized by the existence of a Killing spinor. In

terms of the Energy-Momentum tensor, O. Hijazi [16] proved that, on such manifolds any eigenvalue of the Dirac operator satisfies the following

$$\lambda^2 \geq \begin{cases} \frac{1}{4}\mu_1 + \inf_M |T^\psi|^2 & \text{if } n \geq 3, \\ \frac{\pi\chi(M)}{\text{Area}(M,g)} + \inf_M |T^\psi|^2 & \text{if } n = 2. \end{cases} \quad (6)$$

Again, the trace of T^ψ being equal to λ , Inequality (6) improves Inequalities (4) and (5). The limiting case of (6) is characterized by the existence of a spinor field $\bar{\varphi}$ satisfying for all $X \in TM$,

$$\bar{\nabla}_X \bar{\varphi} = -T^{\bar{\varphi}}(X)\bar{\varphi}, \quad (7)$$

where $\bar{\varphi} = e^{-\frac{n-1}{2}u}\bar{\psi}$, ψ is an eigenspinor associated with the first eigenvalue of the Dirac operator and $\bar{\psi}$ is the image of ψ under the isometry between the spinor bundles of (M^n, g) and $(M^n, \bar{g} = e^{2u}g)$. Suppose that on a spin manifold M , there exists a spinor field ψ such that for all $X \in TM$,

$$\nabla_X \psi = -E(X).\psi, \quad (8)$$

where E is a symmetric 2-tensor defined on TM . It is easy to see that E must be equal to T^ψ . If the dimension of M is equal to 2, Th. Friedrich [6] proved that the existence of a pair (ψ, E) satisfying (8) is equivalent to the existence of a local immersion of M into the euclidean space \mathbb{R}^3 with Weingarten tensor equal to E . In [21], B. Morel showed that if M^n is a hypersurface of a manifold N carrying a parallel spinor, then the Energy-Momentum tensor (associated to the restriction of the parallel spinor) appears, up to a constant, as the second fundamental form of the hypersurface. G. Habib [11] studied Equation (8) for an endomorphism E not necessarily symmetric. He showed that the symmetric part of E is T^ψ and the skew-symmetric part of E is the tensor defined on the complement set of zeroes of ψ by

$$Q^\psi(X, Y) = \frac{1}{2} \text{Re} \langle Y.\nabla_X \psi - X.\nabla_Y \psi, \frac{\psi}{|\psi|^2} \rangle,$$

for all $X, Y \in TM$. Then he modifies the connection in the direction of E and gets that any eigenvalue of the Dirac operator satisfies

$$\lambda^2 \geq \inf_M \left(\frac{1}{4}S_g + |T^\psi|^2 + |Q^\psi|^2 \right). \quad (9)$$

The Heisenberg group and the solvable group are examples of limiting manifolds [11]. For a better understand of the tensor Q^ψ , he studied Riemannian flows and proved that if the normal bundle carries a parallel spinor, the tensor Q^ψ plays the role of the O'Neill tensor of the flow. In this paper we prove the corresponding inequalities for Spin^c manifolds:

Theorem 1.1 *Let (M^n, g) be a compact Riemannian Spin^c manifold of dimension $n \geq 2$, and denote by $i\Omega$ the curvature form of the connection A on the \mathbb{S}^1 -principal fibre bundle $(\mathbb{S}^1 M, \pi, M)$. Then any eigenvalue of the Dirac operator to which is attached an eigenspinor ψ satisfies*

$$\lambda^2 \geq \inf_M \left(\frac{1}{4} S_g - \frac{c_n}{4} |\Omega|_g + |T^\psi|^2 + |Q^\psi|^2 \right), \quad (10)$$

where $c_n = 2 \lfloor \frac{n}{2} \rfloor^{\frac{1}{2}}$ and $|\Omega|_g$ is the norm of Ω with respect to g .

In particular for $Q^\psi \equiv 0$, we have

$$\lambda^2 \geq \inf_M \left(\frac{1}{4} S_g - \frac{c_n}{4} |\Omega|_g + |T^\psi|^2 \right). \quad (11)$$

In 1999, A. Moroianu and M. Herzlich [13] proved that on Spin^c manifolds of dimension $n \geq 3$, any eigenvalue of the Dirac operator satisfies

$$\lambda^2 \geq \lambda_1^2 := \frac{n}{4(n-1)} \mu_1, \quad (12)$$

where μ_1 is the first eigenvalue of the perturbed Yamabe operator defined by

$$L^\Omega = L - c_n |\Omega|_g.$$

The limiting case of (12) is characterized by the existence of a real Killing spinor ψ satisfying $\Omega \cdot \psi = i \frac{c_n}{2} |\Omega|_g \psi$. In terms of the Energy-Momentum tensor we prove:

Theorem 1.2 *Under the same conditions as Theorem 1.1, any eigenvalue λ of the Dirac operator to which is attached an eigenspinor ψ satisfies*

$$\lambda^2 \geq \begin{cases} \frac{1}{4} \mu_1 + \inf_M |T^\psi|^2 & \text{if } n \geq 3, \\ \frac{\pi \chi(M)}{\text{Area}(M, g)} - \frac{1}{2} \frac{\int_M |\Omega|_g v_g}{\text{Area}(M, g)} + \inf_M |T^\psi|^2 & \text{if } n = 2, \end{cases} \quad (13)$$

where μ_1 is the first eigenvalue of the perturbed Yamabe operator.

Using the Cauchy-Schwarz inequality in dimension $n \geq 3$, we have that Inequality (13) implies Inequality (12). As a corollary of Theorem 1.2, we compare the lower bound to a conformal invariant (the Yamabe number) and to a topological invariant, in case of 4-dimensional manifolds whose associated line bundle has self dual curvature (see Corollary (4.1) and Corollary (4.2)). Finally, we study the limiting case of (11) and (13), and we give examples.

Even though the number $\inf_M |T^\psi|^2$ is not a nice geometric invariant, it appears naturally in some situations. For example, on hypersurfaces of certain limiting Spin^c manifolds it is easy to see, with the help of the Spin^c Gauss formula, that it is precisely the second fundamental form. Also, when deforming the Riemannian metric in the direction of the Energy-Momentum tensor, the eigenvalues of the Dirac operator on a Spin^c manifold are then critical (see [24]). The author would like to thank Oussama Hijazi for his support and encouragements.

2 Spin^c geometry and the Dirac operator

In this section, we briefly introduce basic notions concerning Spin^c manifolds and the Dirac operator. Details can be found in [7], [20] and [22].

Let (M^n, g) be a compact connected oriented Riemannian manifold of dimension $n \geq 2$ without boundary. Furthermore, let SOM be the SO_n -principal bundle over M of positively oriented orthonormal frames. A Spin^c structure of M is a Spin_n^c -principal bundle $(\text{Spin}^c M, \pi, M)$ and a \mathbb{S}^1 -principal bundle $(\mathbb{S}^1 M, \pi, M)$ together with a double covering given by $\theta : \text{Spin}^c M \rightarrow SOM \times_M \mathbb{S}^1 M$ such that

$$\theta(ua) = \theta(u)\xi(a),$$

for every $u \in \text{Spin}^c M$ and $a \in \text{Spin}_n^c$, where ξ is the 2-fold covering of Spin_n^c over $SO_n \times \mathbb{S}^1$. A Riemannian manifold that admits a Spin^c structure is called a Riemannian Spin^c manifold.

Let $\Sigma^c M := \text{Spin}^c M \times_{\rho_n} \Sigma_n$ be the associated spinor bundle where $\Sigma_n = \mathbb{C}^{2^{\lfloor \frac{n}{2} \rfloor}}$ and $\rho_n : \text{Spin}_n^c \rightarrow \text{End}(\Sigma_n)$ the complex spinor representation. A section of $\Sigma^c M$ will be called a spinor and the set of all spinors will be denoted by $\Gamma(\Sigma^c M)$. The spinor bundle $\Sigma^c M$ is equipped with a natural Hermitian scalar product, denoted by $\langle \cdot, \cdot \rangle$ and satisfies

$$\langle X.\psi, \varphi \rangle = - \langle \psi, X.\varphi \rangle \quad \text{for every } X \in \Gamma(TM) \text{ and } \psi, \varphi \in \Gamma(\Sigma^c M),$$

where $X.\psi$ denotes the Clifford multiplication of X and ψ . With this Hermitian scalar product we define an L^2 -scalar product

$$(\psi, \phi) = \int_M \langle \psi, \phi \rangle v_g,$$

for any spinors ψ and ϕ . Additionally, given a connection 1-form A on \mathbb{S}^1M , $A : T(\mathbb{S}^1M) \rightarrow i\mathbb{R}$ and the connection 1-form ω^M on SOM for the Levi-Civita connection ∇^M , induce a connection on the principal bundle $SOM \times_M \mathbb{S}^1M$, and hence a covariant derivative ∇ on $\Gamma(\Sigma^c M)$ [7], given by

$$\nabla_{e_i} \psi = \left[b, e_i(\sigma) + \frac{1}{4} \sum_{j=1}^n e_j \cdot \nabla_{e_i}^M e_j \cdot \sigma + \frac{1}{2} A(s_*(e_i))\sigma \right], \quad (14)$$

where $\psi = [b, \sigma]$ is a locally defined spinor field, (e_1, \dots, e_n) is a local oriented orthonormal tangent frame and $s : U \rightarrow \mathbb{S}^1M$ is a local section of \mathbb{S}^1M .

The curvature of A is an imaginary valued 2-form denoted by $F_A = dA$, i.e., $F_A = i\Omega$, where Ω is a real valued 2-form on \mathbb{S}^1M . We know that Ω can be viewed as a real valued 2-form on M [7]. In this case $i\Omega$ is the curvature form of the associated line bundle L . It's the complex line bundle associated with the \mathbb{S}^1 -principal bundle via the standard representation of the unit circle. The spinorial curvature \mathcal{R} associated with the connection ∇ , is given by

$$\mathcal{R}_{X,Y} = \frac{1}{4} \sum_{i,j=1}^n g(R_{X,Y} e_i, e_j) e_i \cdot e_j \cdot + \frac{i}{2} \Omega(X, Y).$$

In the Spin^c case, the Ricci identity translates to

$$\sum_j e_j \cdot \mathcal{R}_{e_j, X} \psi = \frac{1}{2} Ric(X) \cdot \psi - \frac{i}{2} (X \lrcorner \Omega) \cdot \psi, \quad (15)$$

where \lrcorner denotes the interior product. For every spinor ψ , the Dirac operator is locally defined by

$$D\psi = \sum_{i=1}^n e_i \cdot \nabla_{e_i} \psi.$$

It is an elliptic, self-adjoint operator with respect to the L^2 -scalar product and verifies the Schrödinger-Lichnerowicz formula

$$D^2 = \nabla^* \nabla + \frac{1}{4} S_g \text{Id}_{\Gamma(\Sigma^c M)} + \frac{i}{2} \Omega.,$$

where $\Omega.$ is the extension of the Clifford multiplication to differential forms given by $(e_i^* \wedge e_j^*) \cdot \psi = e_i \cdot e_j \cdot \psi$.

3 Eigenvalue estimates on Spin^c manifolds

In this section, we prove the lower bound (10). This proof is based on the following Lemma given by A. Moroianu and M. Herzlich in [13]:

Lemma 3.1 [13]. *Let (M^n, g) be a Spin^c manifold. For any spinor $\psi \in \Gamma(\Sigma^c M)$ and a real 2-form Ω , we have*

$$\langle i\Omega.\psi, \psi \rangle \geq -\frac{c_n}{2}|\Omega|_g|\psi|^2, \quad (16)$$

where $|\Omega|_g$ is the norm of Ω , with respect to g given by $|\Omega|_g^2 = \sum_{i < j} (\Omega_{ij})^2$, in any orthonormal local frame. Moreover, if equality holds in (16), then

$$\Omega.\psi = i\frac{c_n}{2}|\Omega|_g\psi. \quad (17)$$

Proof of Theorem 1.1: The modified connection

$$\tilde{\nabla}_X\psi := \nabla_X\psi + T^\psi(X).\psi + Q^\psi(X).\psi,$$

satisfies $|\tilde{\nabla}\psi|^2 = |\nabla\psi|^2 - |T^\psi|^2|\psi|^2 - |Q^\psi|^2|\psi|^2$. After integration on M , the Schrödinger-Lichnerowicz formula gives

$$\begin{aligned} \int_M |\tilde{\nabla}\psi|^2 v_g &= \int_M |D\psi|^2 v_g - \int_M \frac{1}{4} S_g |\psi|^2 v_g - \int_M (|T^\psi|^2 + |Q^\psi|^2) |\psi|^2 v_g \\ &\quad - \int_M \langle \frac{i}{2} \Omega.\psi, \psi \rangle v_g. \end{aligned}$$

Let ψ be an eigenspinor corresponding to the eigenvalue λ of D . By Lemma 3.1, it follows

$$\begin{aligned} \lambda^2 \int_M |\psi|^2 v_g &\geq \frac{1}{4} \int_M S_g |\psi|^2 v_g + \int_M (|T^\psi|^2 + |Q^\psi|^2) |\psi|^2 v_g \\ &\quad + \int_M \langle \frac{i}{2} \Omega.\psi, \psi \rangle v_g \\ &\geq \int_M \left(\frac{1}{4} S_g - \frac{c_n}{4} |\Omega|_g + |T^\psi|^2 + |Q^\psi|^2 \right) |\psi|^2 v_g. \end{aligned}$$

Finally,

$$\lambda^2 \geq \inf_M \left(\frac{1}{4} S_g - \frac{c_n}{4} |\Omega|_g + |T^\psi|^2 + |Q^\psi|^2 \right).$$

4 Conformal geometry and eigenvalue estimates

Before proving Theorem 1.2, we give some preliminaries about conformal geometry. The conformal class of g is the set of metrics $\bar{g} = e^{2u}g$, for a real function u on M . At a given point x of M , we consider a g -orthonormal basis

$\{e_1, \dots, e_n\}$ of $T_x M$. The corresponding \bar{g} -orthonormal basis is denoted by $\{\bar{e}_1 = e^{-u}e_1, \dots, \bar{e}_n = e^{-u}e_n\}$. This correspondence extends to the Spin^c level to give an isometry between the corresponding spinor bundles. We put a $(-)$ above every object which is naturally associated with the metric \bar{g} . Then, for any spinor fields ψ and φ , one has

$$\langle \bar{\psi}, \bar{\varphi} \rangle = \langle \psi, \varphi \rangle,$$

where $\langle \cdot, \cdot \rangle$ denotes the natural Hermitian scalar products on $\Gamma(\Sigma^c M)$, and on $\Gamma(\Sigma^c \bar{M})$. The corresponding Dirac operators satisfy

$$\bar{D}(e^{-\frac{(n-1)}{2}u} \bar{\psi}) = e^{-\frac{n+1}{2}u} \bar{D}\psi.$$

The norm of any real 2-form Ω with respect to g and \bar{g} are related by

$$|\Omega|_{\bar{g}} = e^{-2u} |\Omega|_g.$$

O. Hijazi [16] showed that on a spin manifold the Energy-Momentum tensor verifies

$$|T^{\bar{\varphi}}|^2 = e^{-2u} |T^{\varphi}|^2 = e^{-2u} |T^{\psi}|^2,$$

where $\varphi = e^{-\frac{(n-1)}{2}u} \psi$. We extend the result to a Spin^c manifold and get the same relation.

Lemma 4.1 *Under the same conditions as Theorem 1.1, any eigenvalue λ of the Dirac operator satisfies*

$$\lambda^2 \geq \frac{1}{4} \sup_u \inf_M (S_u e^{2u} - c_n |\Omega|_g) + \inf_M |T^{\psi}|^2.$$

Proof: The modified connection introduced in [16]:

$$\nabla_X^{T^{\psi}} \psi = \nabla_X \psi + T^{\psi}(X) \cdot \psi,$$

verifies $|\nabla^{T^{\psi}} \psi|^2 = |\nabla \psi|^2 - |T^{\psi}|^2 |\psi|^2$. Using the Schrödinger-Lichnerowicz formula on M , applied to the spinor field $\bar{\varphi}$ with respect to the metric \bar{g} , yields

$$\begin{aligned} \int_M |\bar{\nabla}^{T^{\bar{\varphi}}} \bar{\varphi}|^2 v_{\bar{g}} &= \int_M |\bar{D}\bar{\varphi}|^2 v_{\bar{g}} - \int_M \frac{1}{4} S_u |\bar{\varphi}|^2 v_{\bar{g}} - \int_M |T^{\bar{\varphi}}|^2 |\bar{\varphi}|^2 v_{\bar{g}} \\ &\quad - \int_M \langle \frac{i}{2} \Omega \cdot \bar{\varphi}, \bar{\varphi} \rangle v_{\bar{g}}. \end{aligned} \quad (18)$$

For $\varphi = e^{-\frac{(n-1)}{2}u} \psi$ and $D\psi = \lambda\psi$, one gets $\bar{D}\bar{\varphi} = \lambda e^{-u} \bar{\varphi}$, and hence by Lemma 3.1

$$0 \leq \int_M |\bar{\nabla}^{T^{\bar{\varphi}}} \bar{\varphi}|^2 v_{\bar{g}} \leq \int_M \left[\lambda^2 - \left(\frac{1}{4} S_u e^{2u} + |T^{\psi}|^2 - \frac{c_n}{4} |\Omega|_g \right) \right] e^{-2u} |\bar{\varphi}|^2 v_{\bar{g}}. \quad (19)$$

Lemma 4.2 *Let (M^n, g) be a compact Riemannian Spin^c manifold of dimension $n \geq 2$ and S_g (resp. S_u or S_h) the scalar curvature associated with the metric g (resp. with a conformal metric $\bar{g} = e^{2u}g = h^{\frac{4}{n-2}}g$). The 2-form $i\Omega$ denotes the curvature tensor on the \mathbb{S}^1 -principal bundle associated with the Spin^c structure. We have*

$$\sup_u \inf_M (S_u e^{2u} - c_n |\Omega|_g) = \begin{cases} \mu_1 & \text{if } n \geq 3, \\ \frac{4\pi\chi(M) - 2 \int_M |\Omega|_g v_g}{\text{Area}(M, g)} & \text{if } n = 2, \end{cases} \quad (20)$$

where μ_1 is the first eigenvalue of the perturbed Yamabe operator.

Proof: For $n \geq 3$, let $\mathbf{h} > 0$ be an eigenfunction of L^Ω associated with the eigenvalue μ_1 such that $\int_M \mathbf{h}^2 v_g = 1$. For a conformal metric $\bar{g} = e^{2u}g = h^{\frac{4}{n-2}}g$, we have

$$S_h h^{\frac{4}{n-2}} - c_n |\Omega|_g = S_u e^{2u} - c_n |\Omega|_g = h^{-1} L^\Omega h.$$

So $\mu_1 = \mathbf{h}^{-1} L^\Omega \mathbf{h} = S_h \mathbf{h}^{\frac{4}{n-2}} - c_n |\Omega|_g$. For any positive function H , we write $fH = \mathbf{h}$, where f is a positive function, and referring to [15] we get

$$\mu_1 = \int (H^{-1} L H) f^2 H^2 v_g - c_n \int_M |\Omega|_g f^2 H^2 v_g + \int_M H^2 |df|^2 v_g.$$

Finally,

$$\mu_1 \geq \inf_M (H^{-1} L^\Omega H) = \inf_M (S_v e^{2v} - c_n |\Omega|_g),$$

where $e^{2v} = H^{\frac{4}{n-2}}$, then $\mu_1 = \sup_u \inf_M (S_u e^{2u} - c_n |\Omega|_g)$. For $n = 2$ and for every u we have $S_u e^{2u} = S_g + 2 \Delta_g u$. The Stokes and Gauß-Bonnet theorems yield

$$\inf_M (S_u e^{2u} - 2|\Omega|_g) \leq \frac{\int_M (S_u e^{2u} - 2|\Omega|_g) v_g}{\text{Area}(M, g)} = \frac{4\pi\chi(M) - 2 \int_M |\Omega|_g v_g}{\text{Area}(M, g)}.$$

Let u_0 be a solution of the following equation [1]

$$2 \Delta_g u = \frac{\int_M (S_g - 2|\Omega|_g) v_g}{\text{Area}(M, g)} - S_g + 2|\Omega|_g, \quad (21)$$

hence,

$$S_{u_0} e^{2u_0} - 2|\Omega|_g = 2 \Delta_g u_0 + S_g - 2|\Omega|_g = \frac{4\pi\chi(M) - 2 \int_M |\Omega|_g v_g}{\text{Area}(M, g)}.$$

Proof of Theorem 1.2: Combining Lemma 4.2 and Lemma 4.1, Theorem 1.2 follows.

Remark 4.1 *Inequality (11) improves Inequality (12), which itself implies the Friedrich Spin^c inequality given by*

$$\lambda^2 \geq \frac{n}{4(n-1)} \inf_M (S_g - c_n |\Omega|_g). \quad (22)$$

Equality holds in (22) if and only if equality holds in (12), i.e. if and only if the eigenspinor ψ associated with the first eigenvalue of D is a Killing spinor and $\Omega \cdot \psi = i \frac{c_n}{2} |\Omega|_g \psi$.

Corollary 4.1 *Any eigenvalue of the Dirac operator on a compact Riemannian Spin^c manifold of dimension $n \geq 3$, satisfies*

$$\lambda^2 \geq \frac{1}{4} \text{vol}(M, g)^{-\frac{2}{n}} \left(Y(M, [g]) - c_n \|\Omega\|_{\frac{n}{2}} \right) + \inf_M |T^\psi|^2,$$

where $Y(M, [g])$ is the Yamabe number given by

$$Y(M, [g]) = \inf_{\eta \neq 0} \frac{\int_M 4 \frac{n-1}{n-2} |d\eta|^2 + S_g \eta^2}{\left(\int_M |\eta|^{\frac{2n}{n-2}} \right)^{\frac{n-2}{n}}}.$$

Proof : Using the Hölder inequality, it follows

$$\mu_1 = \inf_{\eta \neq 0} \frac{\int_M 4 \frac{n-1}{n-2} |d\eta|^2 + (S_g - c_n |\Omega|_g) \eta^2}{\int_M \eta^2} \geq \inf_{\eta \neq 0} \frac{\int_M 4 \frac{n-1}{n-2} |d\eta|^2 + (S_g - c_n |\Omega|) \eta^2}{\left(\int_M |\eta|^{\frac{2n}{n-2}} \right)^{\frac{n-2}{n}} \text{vol}(M, g)^{\frac{2}{n}}}.$$

Using the Hölder inequality one more time, we deduce

$$\mu_1 \text{vol}(M, g)^{\frac{2}{n}} \geq \inf_{\eta \neq 0} \frac{\int_M 4 \frac{n-1}{n-2} |d\eta|^2 + S_g \eta^2}{\left(\int_M |\eta|^{\frac{2n}{n-2}} \right)^{\frac{n-2}{n}}} - c_n \left(\int_M |\Omega|^{\frac{n}{2}} \right)^{\frac{2}{n}} = Y(M, [g]) - c_n \|\Omega\|_{\frac{n}{2}}.$$

Finally, replacing in (13), we get the result.

Corollary 4.2 *On a compact 4-dimensional Spin^c manifold with self-dual curvature $i\Omega$, any eigenvalue of the Dirac operator satisfies*

$$\lambda^2 \geq \frac{1}{4} \text{vol}(M, g)^{-\frac{1}{2}} \left(Y(M, [g]) - 4\pi\sqrt{2}\sqrt{c_1(L)^2} \right) + \inf_M |T^\psi|^2,$$

where $c_1(L)$ is the Chern number of the line bundle L associated with the Spin^c structure.

Proof: It follows directly from Corollary 4.1 and the fact that if $n = 4$ and Ω self-dual, then $\int_M |\Omega|_g^2 v_g = 4\pi^2 c_1(L)^2$ (see [7]).

5 Equality case

In this section, we study the limiting case of (11) and (13). Examples are then given.

Proposition 5.1 *Under the same conditions as Theorem 1.1,*

$$\text{Equality in (11) holds} \iff \begin{cases} \nabla^{T^\psi} \psi \equiv 0, \\ \Omega \cdot \psi = i \frac{c_n}{2} |\Omega|_g \psi, \end{cases} \quad (23)$$

where ψ is an eigenspinor associated with the first eigenvalue of the Dirac operator.

Proof: If equality in (11) is achieved, the two conditions follow directly. Now, suppose that $\nabla^{T^\psi} \psi \equiv 0$ and $\Omega \cdot \psi = i \frac{c_n}{2} |\Omega|_g \psi$. The condition $\nabla^{T^\psi} \psi \equiv 0$ implies that $|\psi|^2$ is constant. The curvature tensor on the Spin^c bundle associated with the connection ∇ (resp. ∇^{T^ψ}) is denoted by \mathcal{R} (resp. \mathcal{R}^{T^ψ}). One easily gets the following relation

$$\mathcal{R}_{X,Y}^{T^\psi} \psi = \mathcal{R}_{X,Y} \psi + dT^\psi(X, Y) \cdot \psi + [T^\psi(X), T^\psi(Y)] \cdot \psi,$$

where dT^ψ is a 2-form with values in $\Gamma(TM)$ given by

$$dT^\psi(X, Y) = (\nabla_X T^\psi)Y - (\nabla_Y T^\psi)X.$$

Taking $Y = e_j$ and performing its Clifford multiplication by e_j yields by the Ricci identity (15) on a Spin^c manifold

$$\begin{aligned} -\frac{1}{2} \text{Ric}(X) \cdot \psi + \frac{i}{2} (X \lrcorner \Omega) \cdot \psi &+ \sum_j e_j \cdot dT^\psi(X, e_j) \cdot \psi \\ &+ \sum_j e_j \cdot [T^\psi(X), T^\psi(e_j)] \cdot \psi = 0 \end{aligned} \quad (24)$$

We then decompose the last two terms in (24) using that $X \cdot \alpha = X \wedge \alpha - X \lrcorner \alpha$ for any form α , it follows

$$\begin{aligned} \sum_j e_j \cdot dT^\psi(X, e_j) \cdot \psi &= \sum_j [e_j \wedge dT^\psi(X, e_j)] \cdot \psi - [X(\text{tr } T^\psi) + \text{div } T^\psi(X)] \psi. \\ \sum_j e_j \cdot [T^\psi(X), T^\psi(e_j)] \cdot \psi &= 2(\text{tr } T^\psi) T^\psi(X) \cdot \psi - 2 \sum_j g(X, T^\psi(e_j)) T^\psi(e_j) \cdot \psi. \end{aligned}$$

Taking the scalar product of (24) with ψ , and after separating real and imaginary parts, yields for every vector field X the relation

$$\left(X(\text{tr } T^\psi) + \text{div } T^\psi(X) \right) |\psi|^2 = \frac{i}{2} \langle (X \lrcorner \Omega) \cdot \psi, \psi \rangle. \quad (25)$$

But since Equality (17) holds we compute

$$\begin{aligned} \langle (X \lrcorner \Omega). \psi, \psi \rangle &= \langle (X \wedge \Omega). \psi, \psi \rangle - \langle X. \Omega. \psi, \psi \rangle \\ &= \langle (X \wedge \Omega). \psi, \psi \rangle - i \left[\frac{n}{2} \right]^{\frac{1}{2}} |\Omega|_g \langle X. \psi, \psi \rangle. \end{aligned}$$

After separating real and imaginary parts, $\langle (X \lrcorner \Omega). \psi, \psi \rangle$ must vanish. Using this and $\sum_{j=1}^n e_j. (e_j \lrcorner \Omega) = 2\Omega$, Clifford multiplication of (24) with e_k , and for $X = e_k$, gives

$$-\frac{1}{2} S_g \psi - i \Omega. \psi = \sum_{k,j} e_j \cdot (e_k \wedge dT^\psi(e_j, e_k)) \cdot \psi - 2(\text{tr } T^\psi)^2 \psi + 2|T^\psi|^2 \psi.$$

An easy computation implies that $\sum_{k,j} e_j \cdot (e_k \wedge dT^\psi(e_j, e_k)) \cdot \psi = 0$, hence

$$-\frac{1}{2} S_g + \left[\frac{n}{2} \right]^{\frac{1}{2}} |\Omega|_g = -2(\text{tr } T^\psi)^2 + 2|T^\psi|^2, \quad (26)$$

which implies Equality in (11).

Proposition 5.2 *On a compact Riemannian Spin^c manifold (M^n, g) of dimension $n \geq 3$, assume that the first eigenvalue λ_1 of the Dirac operator to which is attached an eigenspinor ψ satisfies the equality case in (13). Then, $|T^\psi|$ is constant and if \mathbf{h} denotes an eigenfunction of the Yamabe operator corresponding to μ_1 , then for any vector field X*

$$g(X, T^\psi(d\mathbf{h}) - \lambda_1 d\mathbf{h}) = g(\lambda_1 X - T^\psi(X), d\mathbf{h}) = 0. \quad (27)$$

Proof: If $n \geq 3$ and equality holds in (13), Equation (19) with $h = \mathbf{h}$ gives $|T^\psi|$ is constant, $\overline{\nabla}^{T^\psi} \overline{\varphi} = 0$ and $\Omega. \overline{\varphi} = i \frac{c_n}{2} |\Omega|_{\overline{g}} \overline{\varphi}$. By Proposition 5.1, Equality (26) and (25) can be considered for a conformal metric \overline{g} to get

$$\begin{aligned} (\text{tr } T^\psi)^2 &:= f^2 = \frac{1}{4} S_{\mathbf{h}} - \frac{c_n}{4} |\Omega|_{\overline{g}} + |T^\psi|^2, \\ \text{grad } f &= -\text{div } T^\psi. \end{aligned}$$

It is straightforward to see that these two equalities give (27).

Example 1: Note that as every spin manifold is Spin^c with trivial line bundle, the Heisenberg manifold is a Spin^c manifold. In addition if we consider the trivial flat connection on this trivial line bundle, then the Heisenberg manifold is a limiting manifold for (11) but equality in (22) cannot occur.

Example 2: Let M^4 be a compact oriented Riemannian Spin^c 4-dimensional manifold. Let $\Phi \in \Sigma^+ M$ and A , a connection on the S^1 -principal bundle associated with the Spin^c structure, be a solution of the Seiberg-Witten equations

$$D\Phi = 0 \quad \text{and} \quad i\Omega^+ = -\frac{1}{4}\omega^\Phi,$$

where $\omega^\Phi(X, Y) = \langle X.Y.\Phi, \Phi \rangle$ for every $X, Y \in TM$ and Ω^+ is the self-dual part of Ω (the curvature form of A). Referring to [7], we know that

$$|\Omega^+|_g = \frac{\sqrt{2}}{4}|\Phi|^2 \leq -\frac{\sqrt{2}}{4} \inf_M S_g.$$

Using the fact that $\Omega.\psi^+ = \Omega^+.\psi^+$ for any spinor $\psi^+ \in \Sigma^+ M$, Inequality (16) gives

$$\langle i\Omega.\psi^+, \psi^+ \rangle \geq \frac{1}{2} \inf_M S_g |\psi^+|^2. \quad (28)$$

If $\psi = \psi^+ + \psi^-$ is an eigenspinor associated with the eigenvalue λ of D , then ψ^+ is an eigenspinor of D^2 associated with the eigenvalue λ^2 . Using Inequality (28) in the proof of Theorem 1.1 applied to the spinor ψ^+ , the lower bound (11) becomes

$$\lambda^2 \geq \frac{1}{2} \inf_M S_g + \inf_M |T^{\psi^+}|^2. \quad (29)$$

Equality in (29) holds if and only if (23) holds and $|\Phi|^2 = -S_g = \text{constant}$.

Example 3: If the lower bound (22) is achieved, automatically equality holds in (11). Here we will give an example where equality holds in (11) but not in (22).

Let $(M^3, g) = (S^3, \text{can})$ be endowed with its unique spin structure and consider a Killing spinor ψ with Killing constant $\frac{1}{2}$. As the norm of ψ is constant, we may suppose that $|\psi| = 1$. Let ξ be the Killing vector field on M defined by

$$ig(\xi, X) = \langle X.\psi, \psi \rangle.$$

In [13], it is shown that:

1. $id\xi(X, Y) = -\langle X \wedge Y.\psi, \psi \rangle$ for any $X, Y \in TM$.
2. $d|\xi|^2 = -2d\xi(\xi, \cdot) = -2g(\nabla_\xi \xi, \cdot) \simeq -2\nabla_\xi \xi = 0$.
3. $\xi.\psi = i\psi$ and $|\xi| = 1$.

4. $\xi.\psi = -e_1.e_2.\psi$, where $\{\xi/|\xi|, e_1, e_2\}$ is an oriented local orthonormal frame.

Let h be a real constant such that $h > 1$. We define the metric g^h on M , by:

$$\begin{cases} g^h(\xi, X) = g(\xi, X) \text{ pour tout } X \in TM, \\ g^h(X, Y) = h^{-2}g(X, Y) \text{ pour } X, Y \perp \xi. \end{cases}$$

Using the following isomorphism

$$\begin{aligned} (TM, g) &\longrightarrow (TM, g^h) \\ Z &\longrightarrow Z^h = \begin{cases} Z & \text{si } Z = \xi, \\ hZ & \text{si } Z \perp \xi, \end{cases} \end{aligned}$$

if $u = \{\xi, e_1, e_2\}$ is a positive local g -orthonormal frame defined in a neighborhood U of x , then $u^h = \{\xi^h = \xi, e_1^h = he_1, e_2^h = he_2\}$ is a positive local g^h -orthonormal frame defined in a neighborhood U of x .

There exists an isomorphism of vector bundles (see [13]) given by:

$$\begin{aligned} \Sigma_g M &\longrightarrow \Sigma_{g^h} M \\ \psi = [\tilde{u}, \phi] &\longrightarrow \psi^h = [\tilde{u}^h, \phi], \end{aligned}$$

satisfying,

$$\langle \psi_1, \psi_2 \rangle_{\Sigma_g M} = \langle \psi_1^h, \psi_2^h \rangle_{\Sigma_{g^h} M} \text{ and } (X.\psi)^h = X^h.\psi^h \text{ for any } X \in TM.$$

The covariant derivative of the spinor $\psi^h = [\tilde{u}^h, \phi]$ is given by (see [13]):

$$\nabla_{X^h}^h \psi^h = \frac{h^2}{2} X^h.\psi^h + i((1 - h^2)\xi)(X^h)\psi^h.$$

Let $\alpha = (1 - h^2)\xi$ be a 1-form on M . We may view $i\alpha$ as a connection 1-form on the trivial S^1 bundle. Let $L = M \times \mathbb{C}$ be the induced trivial line bundle over M . We denote by σ the global section of L and by ∇^0 the covariant derivative on L induced by the above connection. It satisfies

$$\nabla_X^0 \sigma = i\alpha(X)\sigma, \text{ for any } X \in TM.$$

On the twisted bundle $\Sigma_{g^h} M \otimes L$, we consider the connection $\bar{\nabla} = \nabla^h \otimes \nabla^0$ and we calculate

$$\begin{aligned} \bar{\nabla}_{e_1^h}(\psi^h \otimes \sigma) &= \frac{h^2}{2} e_1^h.(\psi^h \otimes \sigma), \\ \bar{\nabla}_{e_2^h}(\psi^h \otimes \sigma) &= \frac{h^2}{2} e_2^h.(\psi^h \otimes \sigma), \end{aligned}$$

$$\bar{\nabla}_\xi(\psi^h \otimes \sigma) = \left(\frac{-3h^2}{2} + 2\right)\xi.(\psi^h \otimes \sigma).$$

The spinor $\psi^h \otimes \sigma$ is a section of $\Sigma_{g^h}M \otimes L$, which is, of course, the spinor bundle associated to the Spin^c structure with auxiliary line bundle L^2 . It is easy to see that $\psi^h \otimes \sigma$ is an eigenspinor associated with the eigenvalue $\frac{h^2}{2} - 2$, and it is clear that $\psi^h \otimes \sigma$ is not a Killing spinor since $h \neq 1$, so (M, g^h) is not a limiting manifold for the Friedrich Spin^c inequality. But it is a limiting manifold for the lower bound (11), in fact we will prove that (23) holds.

The complex 2-form $id\alpha$ is the curvature form associated with the connection ∇^0 on L . We have:

$$d\alpha.(\psi^h \otimes \sigma) = (1 - h^2)d\xi.(\psi^h \otimes \sigma) = i(h^2 - 1)h^2\psi^h \otimes \sigma.$$

The norm of $d\alpha$ with respect to the metric g^h is given by

$$|d\alpha|_{g^h}^2 = (1 - h^2)^2|d\xi|_{g^h}^2 = (1 - h^2)^2(d\xi(e_1^h, e_2^h))^2 = h^4(1 - h^2)^2.$$

Since $h > 1$, $|d\alpha|_{g^h} = h^2(h^2 - 1)$, then the second equation of (23) is verified. Furthermore, it is easy to check that

$$T^{\psi^h \otimes \sigma}(e_1^h, e_1^h) = T^{\psi^h \otimes \sigma}(e_2^h, e_2^h) = -\frac{h^2}{2},$$

$$T^{\psi^h \otimes \sigma}(e_1^h, \xi) = T^{\psi^h \otimes \sigma}(e_2^h, \xi) = T^{\psi^h \otimes \sigma}(e_1^h, e_2^h) = 0,$$

$$T^{\psi^h \otimes \sigma}(\xi, \xi) = \frac{3h^2}{2} - 2.$$

Finally, it is straightforward to verify that the first equation of (23) holds:

$$-T^{\psi^h \otimes \sigma}(e_1^h).(\psi^h \otimes \sigma) = \frac{h^2}{2}e_1^h.(\psi^h \otimes \sigma) = \bar{\nabla}_{e_1^h}(\psi^h \otimes \sigma),$$

$$-T^{\psi^h \otimes \sigma}(e_2^h).(\psi^h \otimes \sigma) = \frac{h^2}{2}e_2^h.(\psi^h \otimes \sigma) = \bar{\nabla}_{e_2^h}(\psi^h \otimes \sigma),$$

$$-T^{\psi^h \otimes \sigma}(\xi).(\psi^h \otimes \sigma) = \left(\frac{-3h^2}{2} + 2\right)\xi.(\psi^h \otimes \sigma) = \bar{\nabla}_\xi(\psi^h \otimes \sigma).$$

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