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Asymptotic behaviour for a diffusion equation governed by nonlocal interactions

Andami Ovono Armel

March 5, 2010

Abstract

In this paper we study the asymptotic behaviour of a nonlocal non-linear parabolic equation governed by a parameter. After giving the existence of unique branch of solutions composed by stable solutions in stationary case, we gives for the parabolic problem L^∞ estimates of solution based on using the Moser iterations and existence of global attractor. We finish our study by the issue of asymptotic behaviour in some cases when $t \rightarrow \infty$.

1 Introduction

The non-local issues are important in studying the behavior of certain physical phenomena and population dynamics. A major difficulty in studying these problems often lie in the absence of well-known properties as maximum principle, regularity and properties of Lyapunov (see [5], [6]) and also the difficulty to characterize and determine the stationary solutions associated thus making study the asymptotic behavior of these solutions very difficult.

In this paper we study the solution $u(t, x)$ to the nonlocal equation

$$\begin{cases} u_t - \operatorname{div}(a(l_r(u(t)))\nabla u) = f & \text{dans } \mathbb{R}^+ \times \Omega \\ u(x, t) = 0 & \text{sur } \mathbb{R}^+ \times \partial\Omega \\ u(\cdot, 0) = u_0 & \text{dans } \Omega. \end{cases} \quad (1)$$

In the above problem u_0 and f are such that

$$u_0 \in L^2(\Omega), \quad f \in L^2(0, T, L^2(\Omega)), \quad (2)$$

with T a arbitrary positive number, a is a continuous function such that

$$\exists m, M \quad \text{such that} \quad 0 < m \leq a(\epsilon) \leq M \quad \forall \epsilon \in \mathbb{R}. \quad (3)$$

The nonlocal functional l_r is defined such that

$$l_r(\cdot)(x) : L^2(\Omega) \rightarrow \mathbb{R}, \quad u \rightarrow l_r(u(t))(x) = \int_{\Omega \cap B(x, r)} g(y)u(t, y)dy. \quad (4)$$

Here $B(x, r)$ is the closed ball of \mathbb{R}^n with radius r and $g \in L^2(\Omega)$. It is sometimes possible to consider g more generally, especially when one is interested in the study of stationary solutions of (see [3]).

In physical point of view problem (1) gives many applications especially where $g = 1$ in population dynamics. Indeed, in this situation u may represent a population density and $l_r(u)$ the total mass of the subdomain $\Omega \cap B(x, r)$ of Ω . Hence (1) can describe the evolution of a population whose diffusion velocity depends on the total mass of a subdomain of Ω . For more details of modelisation we refer the reader to [7]. This type of equation can be applied more generally to other models including the study of propagation of mutant gene (see [11],[12],[13]). A very recent study of this propagation was made by Bendahmane and Sepúlveda [4] in which they analyze using a finite volume scheme adapted, the transmission of this gene through 3 types of people: susceptible, infected and recovered.

In mathematical point of view, when $r = d$ where d is the diameter of Ω problem (1) has been studied in various forms (see [6],[8],[9],[15]).

However when $0 < r < d$, several questions from the theory of bifurcations have arisen concerning the structure of stationary solutions including the existence of a principle of comparison of different solutions depending on the parameter r and the existence of branches (local and global) of solutions. A large majority of these issues has been resolved in [3]. It shows that when a is decreasing the existence of a unique global branch of solutions and existence of branch of solutions that are purely local. Some questions may then arise:

- (i) The unique branch described in [3] it is composed of stable solutions?
- (ii) What about stability properties of the corresponding parabolic problem?

The plan for this work is the following. In section 2 we give some existence and uniqueness results. Section 3 is devoted to stationary problem corresponding to (1). In particular, we study in a radial case, a generalisation of Chipot-Lovat results about determination of the number of solutions. We also establish that the unique global branch of solutions described in [3] is composed by stable solutions (theorem 3.8). In section 4 firstly we address a L^∞ estimate taking to account L^p estimate based on Moser iterations. Secondly we prove existence of absorbing set in H_0^1 , which allows us to prove the existence of a global attractor associated to (1) (see remark 5). Finally we obtain a result of stability properties of the corresponding parabolic problem.

2 Existence and uniqueness results

In this section we show a result of existence . We set $V = H_0^1(\Omega)$ and V' its dual, we take the norm in V , $\|\cdot\|_V$ such that

$$\|u\|_V^2 = \int_{\Omega} |\nabla u|^2 dx$$

$\langle \cdot, \cdot \rangle$ means the duality bracket of V' and V .

Then we have

Theorem 2.1. *Let $T > 0$, $f \in L^2(0, T, V')$ and $u_0 \in L^2(\Omega)$, we assume that a is a continuous function and the assumption (3) checked then for every r fixed,*

$r \in [0, \text{diam}(\Omega)]$, there exists a function u such that

$$\begin{cases} u \in L^2(0, T, V), & u_t \in L^2(0, T, V') \\ u(0, \cdot) = u_0 & \text{in } \Omega \\ \frac{d}{dt}(u, \phi) + \int_{\Omega} a(l_r(u(t))) \nabla u \nabla \phi dx = \langle f, \phi \rangle & \text{in } D'(0, T) \quad \forall \phi \in H_0^1(\Omega). \end{cases} \quad (5)$$

Moreover if a is locally Lipschitz i.e

$$\forall c \quad \exists \gamma_c \quad \text{such that} \quad |a(\epsilon) - a(\epsilon')| \leq \gamma_c |\epsilon - \epsilon'| \quad \forall \epsilon, \epsilon' \in [-c, c], \quad (6)$$

then the solution of (5) is unique.

Remark 1. Before to do the proof, it is necessary to see that for $r = 0$ problem (5) is linear and the proof follows a well-known result (see [10]), it is even when $r = \text{diam}(\Omega)$ (see [7]). We will focus therefore in the following where $r \in]0, \text{diam}(\Omega)[$.

Proof. For the existence proof we will use the Schauder fixed point theorem. Let $w \in L^2(0, T, L^2(\Omega))$ we get

$$t \longrightarrow l_r(w(t)),$$

is measurable as a is continuous then

$$t \longrightarrow a(l_r(w(t))),$$

is too. The problem of finding $u = u(t, x)$ solution of

$$\begin{cases} u \in L^2(0, T, V) \cap C([0, T], L^2(\Omega)) & u_t \in L^2(0, T, V') \\ u(0, \cdot) = u_0 \\ \frac{d}{dt}(u, \phi) + \int_{\Omega} a(l_r(w(t))) \nabla u \nabla \phi dx = \langle f, \phi \rangle & \text{in } D'(0, T) \quad \forall \phi \in H_0^1(\Omega), \end{cases} \quad (7)$$

is linear, besides (7) admits a unique solution $u = F_r(w)$ (see [10], [7]). Thus we show that the application

$$w \longrightarrow F_r(w) = u, \quad (8)$$

admits a fixed point. Taking $w = u$ in (7) we get using (3) and the Cauchy-Schwarz inequality

$$\frac{1}{2} \frac{d}{dt} |u|_2^2 + m \|u\|_V^2 \leq |f|_{\star} \|u\|_V, \quad (9)$$

$\|\cdot\|_V$ is the usual norm in V and $|f|_{\star}$ is the dual norm of f . We take

$$|u|_{L^2(0, T, V)} = \left\{ \int_0^T \|u\|_V^2 dt \right\}^{\frac{1}{2}}.$$

Using Young's inequality to the right-hand side of (9), it follows that

$$\frac{1}{2} \frac{d}{dt} |u|_2^2 + \frac{m}{2} \|u\|_V^2 \leq \frac{1}{2m} |f|_{\star}^2. \quad (10)$$

By integrating (10) on $(0, t)$ for $t \leq T$ we obtain

$$\frac{1}{2} |u(t)|_2^2 + \frac{m}{2} \int_0^t \|u\|_V^2 dt \leq \frac{1}{2} |u_0|_2^2 + \frac{1}{2m} \int_0^t |f|_{\star}^2. \quad (11)$$

We deduce that there exists a constant $C = C(m, u_0, f)$ such that

$$|u|_{L^2(0,T,V)} \leq C \quad (12)$$

Moreover

$$\langle u_t, v \rangle + \langle -\operatorname{div}(a(l_r(u(t)))\nabla u), v \rangle = \langle f, v \rangle \quad \forall v \in V,$$

This gives us

$$|u_t|_* \leq M\|u\|_V + |f|_* \quad (13)$$

By raising (13) squared and using the Young inequality we have that

$$|u_t|_*^2 \leq 2M^2\|u\|_V^2 + 2|f|_*^2 \quad (14)$$

By integrating (14) on $(0, t)$ and assuming (12) we obtain

$$|u_t|_{L^2(0,T,V')} \leq C' \quad (15)$$

with $C' = C'(m, M, f, u_0)$ and C' is independent to w . It follows from (12) and (15)

$$|u_t|_{L^2(0,T,V')}^2 + |u|_{L^2(0,T,V)}^2 \leq R, \quad (16)$$

with $R = C^2 + C'^2$. From (12) and the Poincaré inequality it follows that

$$|u|_{L^2(0,T,L^2(\Omega))} \leq R', \quad (17)$$

By setting

$$R_1 = \max(R', R), \quad (18)$$

and associating (17) and (18), it follows that the application F maps the ball $B(0, R_1)$ of $L^2(0, T, L^2(\Omega))$ into itself. Moreover the balls of $H^1(0, T, V, V')$ are relatively compact in $L^2(0, T, L^2(\Omega))$ (see [10] for more details), (16) clearly shows us that $F(B(0, R_1))$ is relatively compact in $B(0, R_1)$ with

$$B(0, R_1) = \{u \in L^2(0, T, L^2(\Omega)); \quad |u|_{L^2(0,T,L^2(\Omega))} \leq R_1\}.$$

In order to apply the Schauder fixed point theorem, as announced, we just need to show that F is continuous from $B(0, R_1)$ to itself. This is actually the case and completes the proof of existence.

We will now discuss the uniqueness assuming of course that assumption (6) be verified. Consider u_1 and u_2 two solutions (5), by subtracting one obtains in $D'(0, T)$

$$\frac{d}{dt}(u_1 - u_2, v) + \int_{\Omega} (a(l_r(u_1(t)))\nabla u_1(t) - a(l_r(u_2(t)))\nabla u_2(t))\nabla \phi dx = 0 \quad \forall \phi \in H_0^1(\Omega). \quad (19)$$

Since

$$\begin{aligned} a(l_r(u_1(t)))\nabla u_1 - a(l_r(u_2(t)))\nabla u_2(t) &= (a(l_r(u_1(t))) - a(l_r(u_2(t))))\nabla u_1(t) \\ &\quad + a(l_r(u_2(t)))\nabla(u_1(t) - u_2(t)), \end{aligned} \quad (20)$$

we get

$$\begin{aligned} \frac{d}{dt}(u_1 - u_2, v) + \int_{\Omega} a(l_r(u_2(t))) \nabla(u_1(t) - u_2(t)) \nabla \phi dx \\ = - \int_{\Omega} (a(l_r(u_1(t))) - a(l_r(u_2(t)))) \nabla u_1 \nabla \phi dx \quad \forall \phi \in H_0^1(\Omega). \end{aligned} \quad (21)$$

Moreover $u_1, u_2 \in C([0, T], L^2(\Omega))$ there exist $z > 0$ such that

$$l_r(u_1(t)), l_r(u_2(t)) \in [-z, z]. \quad (22)$$

Taking $v = u_1 - u_2$ in (21), it comes easily by Cauchy-Schwartz inequality and (6)

$$\frac{1}{2} \frac{d}{dt} |u_1 - u_2|_2^2 + m \|u_1 - u_2\|_V^2 \leq \gamma |l_r(u_1(t)) - l_r(u_2(t))| \|u_1\|_V \|u_1 - u_2\|_V. \quad (23)$$

We get in [3]

$$|l_r(u(t))| \leq C |B(x, r) \cap \Omega|^{\frac{1}{n \vee 3}} |g|_2 |u(t)|_2 \leq |\Omega|^{\frac{1}{n \vee 3}} |g|_2 |u(t)|_2, \quad (24)$$

where C a constant, $|\Omega|$ represents the measure of Ω and $n \vee 3$ the maximum between the dimension n of Ω and 3. By using (24), (23) and the Young inequality

$$ab \leq \frac{1}{2m} b^2 + \frac{m}{2} a^2.$$

We deduce

$$\frac{d}{dt} |u_1 - u_2|_2^2 + m \|u_1 - u_2\|_V^2 \leq p(t) |u_1 - u_2|_2^2, \quad (25)$$

with

$$p(t) = \frac{1}{m} (\gamma C |\Omega|^{\frac{1}{n \vee 3}} |g|_2 \|u_1\|_V)^2 \in L^1(0, T),$$

which leads to

$$\frac{d}{dt} |u_1 - u_2|_2^2 \leq p(t) |u_1 - u_2|_2^2. \quad (26)$$

Multiplying (26) by $e^{-\int_0^t p(s) ds}$ it follows that

$$e^{-\int_0^t p(s) ds} \frac{d}{dt} |u_1 - u_2|_2^2 - p(t) e^{-\int_0^t p(s) ds} |u_1 - u_2|_2^2 \leq 0. \quad (27)$$

Hence

$$\frac{d}{dt} \{e^{-\int_0^t p(s) ds} |u_1 - u_2|_2^2\} \leq 0. \quad (28)$$

This shows that $t \mapsto e^{-\int_0^t p(s) ds} |u_1 - u_2|_2^2$ is nonincreasing. Since for $t = 0$,

$$u_1(0, \cdot) = u_2(0, \cdot) = u_0.$$

This function vanishes at 0 and nonnegative, we conclude that it is identically zero. This concludes the proof. \square

3 Stationary solutions

Consider the weak formulation to the stationary problem associated to (1)

$$(P_r) \begin{cases} -\operatorname{div}(a(l_r(u))\nabla u) = f & \text{dans } \Omega \\ u \in H_0^1(\Omega). \end{cases} \quad (29)$$

3.1 The case $r = d$

By taking ϕ the weak solution of the problem

$$\begin{cases} -\Delta\phi = f & \text{dans } \Omega \\ \phi \in H_0^1(\Omega), \end{cases} \quad (30)$$

we get due to a Chipot-Lovat [8] results that

Theorem 3.1. *Let a be a mapping from \mathbb{R} into $(0, \infty)$. The problem (P_d) has many solutions as the problem in \mathbb{R}*

$$\mu a(\mu) = l_d(\phi), \quad (31)$$

with $\mu = l_d(u_d)$.

Remark 2. Theorem 3.1 allows us to see where a is increasing that the problem P_d admits a unique solution and determine for a given a the exact number of solutions (P_d) . However it is difficult or impossible to adapt the proof of the theorem 3.1 in case $0 < r < d$.

3.2 The case $0 < r < d$

As announced in the introduction we focus our study to the case of radial solutions of (P_d) . We will assume Ω is the open ball of \mathbb{R}^n with radius $d/2$ centered at zero. We set

$$L_r^2(\Omega) = \{u \in L^2(\Omega) \quad \exists \tilde{u} \in L^2(]0, d/2[) \quad \text{such that } u(x) = \tilde{u}(\|x\|)\},$$

and we also assume that

$$\begin{aligned} f &\in L_r^2(\Omega) \\ g &\in L_r^2(\Omega) \\ a &\in W^{1,\infty}(\mathbb{R}), \inf_{\mathbb{R}} a > 0 \\ f &\geq 0 \quad \text{a.e in } \Omega \\ g &\geq 0 \quad \text{a.e in } \Omega. \end{aligned} \quad (32)$$

We start by giving in some sense in a linear case a result that will be used later to explain the asymptotic behavior.

Proposition 3.2. *Let $A, B \in C(\overline{\Omega})$ be positive radial functions such that $A \leq B$ in $\overline{\Omega}$ and also $f, h \in L^2(\Omega)$ two positive radial functions. Let $u \in H_0^1(\Omega)$ the radial solution to*

$$-\operatorname{div}(A(x)\nabla u) = f \quad \text{in } \Omega, \quad (33)$$

and

$$-\operatorname{div}(B(x)\nabla u) = h \quad \text{in } \Omega. \quad (34)$$

Then $f \leq h$ a.e in Ω .

Proof. We proved in [3] that if u is a the radial solution of (33) then for a.e t in $[0, d/2]$,

$$\tilde{u}'(t) = -\frac{1}{\tilde{A}(t)} \int_0^t \left(\frac{s}{t}\right)^{n-1} \tilde{f}(s) ds. \quad (35)$$

From (33), (34) and (35) we obtain

$$\frac{\tilde{B}(t)}{\tilde{A}(t)} \int_0^t \left(\frac{s}{t}\right)^{n-1} \tilde{f}(s) ds = \int_0^t \left(\frac{s}{t}\right)^{n-1} \tilde{h}(s) ds.$$

Since $A \leq B$ in $\overline{\Omega}$ and $f, h \geq 0$ with $f \not\equiv 0, h \not\equiv 0$ hence $f \leq g$. □

In a nonlocal case, some results of existence of radial solutions and comparison principle between u_r, u_d and u_0 has been demonstrated in [3]. It is also proved that if we set for all $r \in [0, d]$

$$I_r := [\inf_{\Omega} l_r(\phi), \sup_{\Omega} l_r(\phi)]. \quad (36)$$

Here ϕ denotes the solution of

$$\begin{cases} -\Delta \phi = f & \text{dans } \Omega \\ \phi \in H_0^1(\Omega). \end{cases} \quad (37)$$

By the inclusion or not of I_r at an interval of \mathbb{R} we somehow generalize the theorem 3.1.

Lemma 3.3. *Let $r \in [0, d]$. Assume that (32) holds true and there exist $0 \leq m_1 \leq m_2$ such that*

$$a(m_1) = \max_{[m_1, m_2]} a \quad a(m_2) = \min_{[m_1, m_2]} a \quad (38)$$

$$I_r \subset [m_1 a(m_1), m_2 a(m_2)]. \quad (39)$$

Then (P_r) admits a radial solution u and

$$m_1 \leq l_r(u) \leq m_2 \quad \text{a.e in } \Omega. \quad (40)$$

For the proof, we refer the reader to [3].

Generalizing this construction type of the diffusion coefficient a we obtain

Proposition 3.4. *Let $r \in [0, d]$. Assume that (32) holds true and there exist an odd integer n_1 and $n_1 + 1$ positive real numbers $\{m_i\}_{i=0 \dots n_1}$, with $m_0 = 0$ and for all $i \in \{0, \dots, n_1 - 1\}$ we have $m_i < m_{i+1}$. Moreover*

$$\begin{aligned} a(m_i) &= \max_{[m_i, m_{i+1}]} a; \quad a(m_{i+1}) = \min_{[m_i, m_{i+1}]} a \quad \forall i \in \{0, 2, \dots, n_1 - 3, n_1 - 1\} \\ I_r &\subset \bigcap_{i=0, 2, \dots, n_1 - 3, n_1 - 1} [m_i a(m_i), m_{i+1} a(m_{i+1})] \end{aligned} \quad (41)$$

Then (P_r) admits at least $\frac{n_1+1}{2}$ radial solutions $\{u_i\}_{i \in \{0, 2, \dots, n_1 - 1\}}$ such that

$$m_i \leq l_r(u_i) \leq m_{i+1} \quad \forall i \in \{0, 2, \dots, n_1 - 3, n_1 - 1\}.$$

Proof. The proof here is by induction. Indeed we set

$$\mathcal{P}_{n_1} = \{ \text{If condition (41) is satisfied then } (P_r) \text{ admits at least } \frac{n_1 + 1}{2} \text{ solutions.} \}$$

By using lemma 3.3 with $m_1 = 0$ and $m_2 = m_1$, it is easy to prove for $n_1 = 1$ that \mathcal{P}_{n_1} is true. For $n_1 > 1$, This procedure can be repeated to prove that if \mathcal{P}_{n_1-2} holds true then \mathcal{P}_{n_1} holds too. \square

Example 1. Let us see a function a satisfying proposition 3.4. For this, we consider the case $n_1 = 3$ and $r \in (0, d]$. Considering (32) and the strong maximum principle we get $\min I_r > 0$. Taking

$$m_1 := 2 \frac{\max I_r}{a(0)}, \quad a(m_1) := \frac{a(0)}{2}$$

with $a(0) > 0$ and also a decreasing on $[0, m_1]$ then we prove lemma 3.3 conditions.

By repeating this process with $m_2 > m_1$ and setting

$$a(m_2) := \frac{\min I_r}{m_2}, \quad m_3 := 2 \frac{\max I_r}{a(m_2)}$$

with $a(m_3) := \frac{a(m_2)}{2}$ and also a is decreasing on $[m_2, m_3]$. This shows the existence of such a .

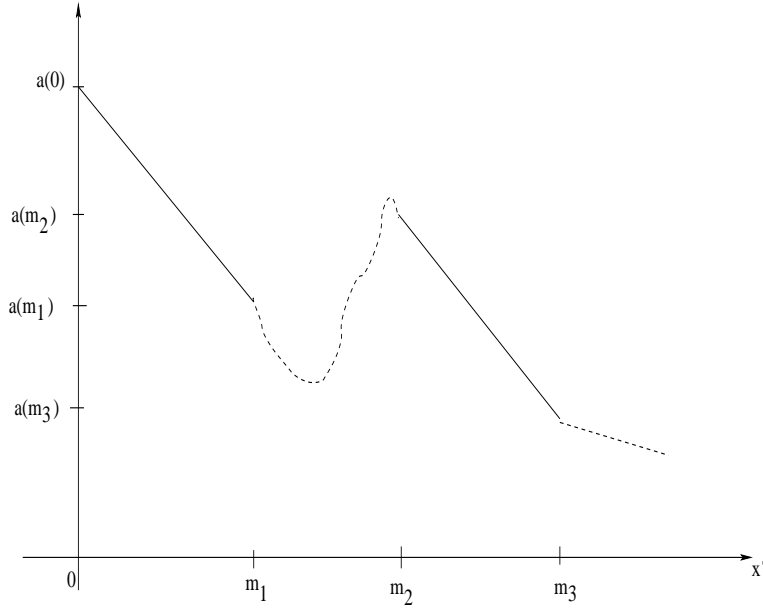


Figure 1: The case $n_1 = 3$

In the representation of a we have deliberately left, on solid line parts of the curve satisfying the conditions of proposition 3.4 and dotted line one without constraints. This situation are explain in the figure 1.

Remark 3. As previously announced, the proposition 3.4 generalizes a certain point of view Theorem 3.1. However it does not accurately determine the exact number of solutions of (P_r) and the bifurcation points of branch of solutions. We have shown in [2] way to solve this problem by using the linearized problem, the principle of comparisons obtained in [3] and the Krein-Rutman theorem.

3.3 Stable solutions of (P_r)

Definition 3.5. Given a domain $\Omega \subset \mathbb{R}^n$, a solution $u_r \in H_0^1(\Omega)$ of (P_r) is stable if:

$$\forall \phi \in H_0^1(\Omega) \quad G_{u_r}(\phi) := \int_{\Omega} a(l_r(u_r)) |\nabla \phi|^2 - \int_{\Omega} a'(l_r(u_r)) l_r(\phi) \nabla u_r \nabla \phi \geq 0. \quad (42)$$

Definition 3.6. Given $u : [0, d] \rightarrow H_0^1(\Omega)$, the graph of u is called a (global) branch of solutions if

- (i) $u \in C([0, d], H_0^1(\Omega))$,
- (ii) $u(r)$ is solution to (P_r) for all r in $[0, d]$.

u is called a local branch if it's defined only on a subinterval of $[0, d]$ with positive measure.

Before concluding this section, we will focus into the case a nonincreasing to prove the stability of the global branch of solutions. Assume for all $r \in [0, d]$, u_r is a solution to (P_r) and

$$0 \leq l_r(u_r)(x) \leq \mu_d \quad \text{for a.e. } x \in \Omega. \quad (43)$$

Assume that there exists a solution μ_d to (31) such that

$$a(\mu_d) = \min_{[0, \mu_d]} a \quad \text{and} \quad a(0) = \max_{[0, \mu_d]} a. \quad (44)$$

We prove in [3]

Theorem 3.7. Assume (32), (43), (44) and (31) holds. Assume in addition that $a \in W^{1, \infty}(\mathbb{R})$ and for some positive constant ϵ , it holds that

$$C_1 |g|_2 |f|_2 |a'|_{\infty, [-\epsilon, \mu_d + \epsilon]} \frac{1}{a(\mu_d)^2} < 1, \quad (45)$$

where C_1 is a constant dependent to Ω . Then

- (i) For all r in $[0, d]$, (P_r) possesses a unique radial solution u_r in $[u_0, u_d]$;
- (ii) $\{(r, u_r) : r \in [0, d]\}$ is a branch of solutions without bifurcation point;
- (iii) it is only global branch of solutions;
- (iv) if in addition, a is nonincreasing on $[0, \mu_d]$ then $r \mapsto u_r$ is nondecreasing.

Remark 4. It is very difficult to obtain property (iv) for any a . However when a is nonincreasing provide us important information for studying the stability of this branch of solutions.

Corollary 3.8. Let u_d^1 the smallest solution to (P_d) . Assume (32) and (31) holds true and there exists a solution μ_d to (31) satisfied (44). Assume in addition that $a \in W^{1,\infty}(\mathbb{R})$, u_d^1 satisfied (43) and for some positive constant ϵ , it holds that

$$C_1 |g|_2 |f|_2 |a'|_{\infty, [-\epsilon, \mu_d + \epsilon]} \frac{1}{a(\mu_d)^2} < 1, \quad (46)$$

where C_1 is a constant dependent to Ω .

Then $\{(r, u_r) : r \in [0, d]\}$ is the only global branch of solutions starting to u_d^1 .

Proof. The fact that $\{(r, u_r) : r \in [0, d]\}$ is the only global branch of solutions results from theorem 3.7. We will now show that this unique branch of solutions is stable and start at $r = d$ by u_d^1 . For this we consider without loss of generality (P_d) admits two solutions u_d^1 and u_d^2 such that $u_d^1 \leq u_d^2$. We denote by μ_1 and μ_2 respectively solutions of (31) corresponding to u_d^1 and u_d^2 (see figure 2). It is easy to see that μ_1 and μ_2 satisfied (44).

Assume $\{(r, u_r) : r \in [0, d]\}$ is the only global branch of solutions starting to u_d^2 . Then we get $C_1 |g|_2 |f|_2 |a'|_{\infty, [-\epsilon, \mu_2 + \epsilon]} \frac{1}{a(\mu_2)^2} < 1$. In this case, using theorem 3.7 we get (P_r) possesses a unique radial solution u_r in $[u_0, u_d^2]$ and the mapping $r \mapsto u_r$ is nondecreasing. By continuity of this mapping, we can find a $r_0 \in]0, d[$ such that $u_{r_0} = u_d^1$ for a.e $x \in \Omega$. This means that u_d^1 is a solution of (P_{r_0}) . This gives us an absurdity and concludes the proof. \square

We are now able to prove:

Proposition 3.9. Under assumptions and notation of corollary 3.8, the global branch of solutions described in theorem 3.7 is composed by stable solutions.

Proof. For all $r \in [0, d]$, let u_r be a solution belonging to the global branch of solutions described in theorem 3.7. By using the linearized problem of (P_r) , we get $\forall \phi \in H_0^1(\Omega)$

$$\begin{aligned} \int_{\Omega} a(l_r(u_r)) |\nabla \phi|^2 - \int_{\Omega} a'(l_r(u_r)) l_r(\phi) \nabla u_r \nabla \phi &\geq \\ \inf_{\Omega} a(l_r(u_r)) |\nabla \phi|_2^2 - C |g|_2 |a'|_{\infty, [-\epsilon, \mu_1 + \epsilon]} |\nabla u_r|_2 |\nabla \phi|_2^2. & \end{aligned} \quad (47)$$

Taking into account that $|\nabla u_r|_2 \leq C(\Omega) \frac{|f|_2}{\inf_{\Omega} a(l_r(u_r))}$ where $C(\Omega)$ designed the Poincaré Sobolev constant. We obtain

$$\begin{aligned} \int_{\Omega} a(l_r(u_r)) |\nabla \phi|^2 - \int_{\Omega} a'(l_r(u_r)) l_r(\phi) \nabla u_r \nabla \phi &\geq \\ |\nabla \phi|_2^2 \left(\inf_{\Omega} a(l_r(u_r)) - C_1 |g|_2 |a'|_{\infty, [-\epsilon, \mu_1 + \epsilon]} \frac{|f|_2}{\inf_{\Omega} a(l_r(u_r))} \right). & \end{aligned} \quad (48)$$

Moreover by assumptions (43) and (44) we get $a(\mu_1) \leq \inf_{\Omega} a(l_r(u_r))$.

Thus (46) becomes

$$C_1 |g|_2 |f|_2 |a'|_{\infty, [-\epsilon, \mu_d + \epsilon]} \frac{1}{\inf_{\Omega} a(l_r(u_r))^2} < 1. \quad (49)$$

We deduces

$$\int_{\Omega} a(l_r(u_r)) |\nabla \phi|^2 - \int_{\Omega} a'(l_r(u_r)) l_r(\phi) \nabla u_r \nabla \phi \geq 0. \quad (50)$$

This concluded the proof. \square

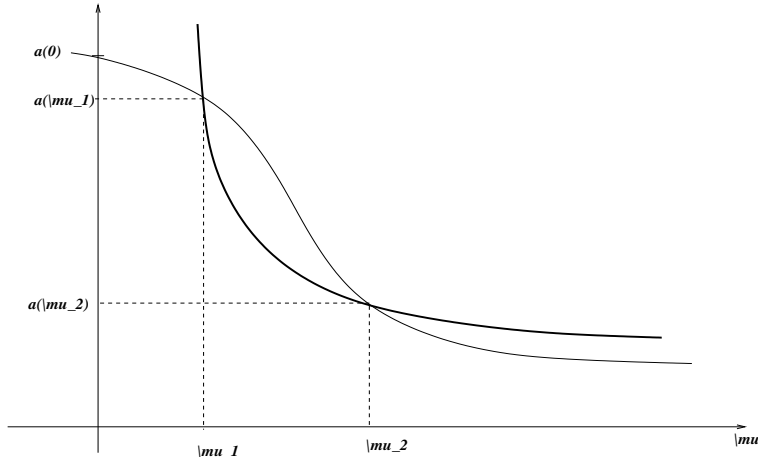


Figure 2: case of 2 solutions

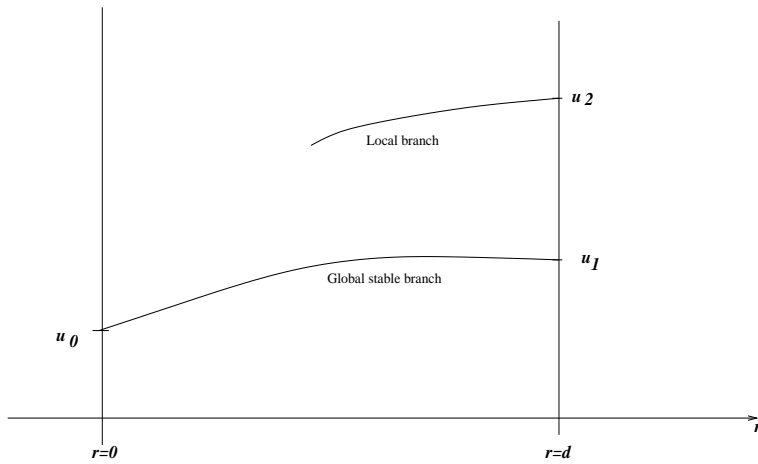


Figure 3: Branch of solutions

4 Parabolic problem

4.1 L^∞ estimate

In what follows we obtain L^∞ estimate of the solution (1) from L^q estimate. The method we use is based on iterations Moser type, for more details on the method see [14].

We get

Theorem 4.1. *Let $n \geq 3$ and u a classical solution of (1) defined on $[0, T]$. Assume that $p > 1$ and $q > 1$ such that $\frac{1}{p} + \frac{1}{q} = 1$. Suppose further that $U_q = \sup_{t < T} |u(t)|_q < \infty$, $f \in L^\infty(0, \infty, L^q(\Omega))$. If $p < \frac{n}{n-2}$ then $U_\infty < \infty$.*

To prove this theorem we need the following proposition:

Lemma 4.2. *Consider u a classical solution of (1) on $[0, T]$, $r \geq 1$ and $p > 1$ such that $\frac{1}{p} + \frac{1}{q} = 1$ with $p < \frac{n}{n-2}$. We take $\tilde{U}_r = \max\{1, |u_0|_\infty, U_r = \sup_{t < T} |u(t)|_r\}$ and let*

$$\sigma(r) = \frac{p(n+2)}{2[r(2p-pn+n)+np]}.$$

Then there exists a constant $C_2 = C_2(\Omega, m)$ such that

$$\tilde{U}_{2r} \leq [C_2 \|f\|_{L^\infty(0, \infty, L^q(\Omega))}]^{\sigma(r)} r^{\sigma(r)} \tilde{U}_r.$$

Proof. Multiplying (1) by u^{2r-1} and then using the Hölder inequality yields

$$\frac{1}{2r} \frac{d}{dt} \int_{\Omega} u^{2r} dx + m \frac{2r-1}{r^2} \int_{\Omega} |\nabla(u^r)|^2 dx \leq |f|_q |u^{2r-1}|_p. \quad (51)$$

As

$$|u^{2r-1}|_p = |u^r|_{\frac{2r-1}{p}}, \quad (52)$$

by taking $w = u^r$ in (51) and (52), we get easily

$$\frac{1}{2r} \frac{d}{dt} |w|_2^2 + m \frac{2r-1}{r^2} |\nabla w|_2^2 \leq |f|_q |w|_{\alpha p}^\alpha, \quad (53)$$

with $\alpha = \frac{2r-1}{r}$. Let β such that

$$\frac{1}{\alpha p} = \beta + \frac{1-\beta}{2^*}, \quad (54)$$

with $2^* = \frac{2n}{n-2}$. We claim that $\beta \in (0, 1)$.

In fact

$$\beta = \frac{2nr - (n-2)(2r-1)p}{(n+2)(2r-1)p}.$$

Since $p < \frac{2r}{2r-1} \frac{n}{n-2}$ then $\beta > 0$. As well as $(n+2)(2r-1)p > 2nr - (n-2)(2r-1)p$ implies that $\beta < 1$ this prove that $\beta \in (0, 1)$.

Using an interpolation inequality (see [14]) in (53) and (54), we get

$$\frac{1}{2r} \frac{d}{dt} |w|_2^2 + m \frac{2r-1}{r^2} |\nabla w|_2^2 \leq |f|_q \left(|w|_1^\beta |w|_{2^*}^{1-\beta} \right)^\alpha. \quad (55)$$

Applying Sobolev injections in (55), we have

$$\frac{1}{2r} \frac{d}{dt} |w|_2^2 + m \frac{2r-1}{r^2} |\nabla w|_2^2 \leq \left[|f|_q \left(\frac{2r}{m} \right)^{\frac{\alpha(1-\beta)}{2}} |w|_1^{\beta\alpha} C^{(1-\beta)\alpha} \right] \left[\left(\frac{m}{2r} \right)^{\frac{\alpha(1-\beta)}{2}} |\nabla w|_2^{(1-\beta)\alpha} \right], \quad (56)$$

and also

$$\frac{1}{2r} \frac{d}{dt} |w|_2^2 + m \frac{2r-1}{r^2} |\nabla w|_2^2 \leq \left[|f|_q \left(\frac{2r}{m} \right)^{\frac{\alpha(1-\beta)}{2}} |w|_1^{\beta\alpha} C^{(1-\beta)\alpha} \right] \left[\left(\frac{m}{2r} \right) |\nabla w|_2^2 \right]^{\frac{\alpha(1-\beta)}{2}}. \quad (57)$$

Since $\beta \in (0, 1)$ and $\frac{\alpha}{2} \in (0, 1)$ it is clear that $\frac{\alpha(1-\beta)}{2} \in (0, 1)$. Applying Young's inequality to (57) with $\frac{\alpha(1-\beta)}{2} + 1 - \frac{\alpha(1-\beta)}{2} = 1$. We obtain

$$\frac{1}{2r} \frac{d}{dt} |w|_2^2 + m \frac{2r-1}{r^2} |\nabla w|_2^2 \leq \delta \left[|f|_q^{\frac{1}{\delta}} \left(\frac{2r}{m} \right)^{\frac{\alpha(1-\beta)}{2\delta}} |w|_1^{\frac{\beta\alpha}{\delta}} C^{\frac{2}{\delta}} \right] + \frac{\alpha(1-\beta)}{2} \left[\left(\frac{m}{2r} \right) |\nabla w|_2^2 \right], \quad (58)$$

with $\delta = 1 - \frac{\alpha(1-\beta)}{2}$.

Joining the fact that $\frac{\alpha(1-\beta)}{2} \in (0, 1)$ and $\delta < 1$ to (58), we deduce

$$\frac{1}{2r} \frac{d}{dt} |w|_2^2 + m \frac{3r-2}{2r^2} |\nabla w|_2^2 \leq |f|_q^{\frac{1}{\delta}} \left(\frac{2r}{m} \right)^{\frac{\alpha(1-\beta)}{2\delta}} |w|_1^{\frac{\beta\alpha}{\delta}} C^{\frac{2}{\delta}}. \quad (59)$$

We set

$$2r\sigma(r) - 1 = \frac{\alpha(1-\beta)}{2\delta} \quad \text{and} \quad 2\rho(r) = \frac{\beta\alpha}{\delta},$$

(59) becomes

$$\frac{1}{2r} \frac{d}{dt} |w|_2^2 + m \frac{3r-2}{2r^2} |\nabla w|_2^2 \leq |f|_q^{\frac{1}{\delta}} \left(\frac{2r}{m} \right)^{2r\sigma(r)-1} |w|_1^{2\rho(r)} C^{\frac{2}{\delta}}. \quad (60)$$

This gives us taking into account that $\frac{3r-2}{r} > 1$

$$\frac{d}{dt} |w|_2^2 + m |\nabla w|_2^2 \leq |f|_q^{\frac{1}{\delta}} \left(\frac{2r}{m} \right)^{2r\sigma(r)} |w|_1^{2\rho(r)} m C^{\frac{2}{\delta}}. \quad (61)$$

By a calculation we can verify that

$$\rho(r) = \frac{2nr - (n-2)(2r-1)p}{2r(p(n+2)+n) - 2n(2r-1)p},$$

and also that $\rho(r) \in (0, 1)$.

Using the Poincaré Sobolev inequality and that $\rho(r) < 1$ in (61), yields

$$\frac{d}{dt} |w|_2^2 + \frac{m}{C_1(\Omega)} |w|_2^2 \leq |f|_q^{\frac{1}{\delta}} \left(\frac{2r}{m} \right)^{2r\sigma(r)} |w|_1^2 m C^{\frac{2}{\delta}}, \quad (62)$$

where $C_1(\Omega)$ designed the Poincaré Sobolev constant. Noticing that

$$e^{-\frac{m}{C_1(\Omega)}t} \frac{d}{dt} \left(e^{\frac{m}{C_1(\Omega)}t} |w|_2^2 \right) = \frac{d}{dt} |w|_2^2 + \frac{m}{C_1(\Omega)} |w|_2^2 \leq |f|_q^{\frac{1}{\delta}} \left(\frac{2r}{m} \right)^{2r\sigma(r)} |w|_1^2 m C^{\frac{2}{\delta}}. \quad (63)$$

and integrating (63) on $[0, t)$ we get

$$|w(t)|_2^2 \leq |w(0)|_2^2 + \|f\|_{L^\infty(0, \infty, L^q(\Omega))}^{\frac{1}{\delta}} \left(\frac{2r}{m}\right)^{2r\sigma(r)} m C^{\frac{2}{\delta}} |w|_1^2. \quad (64)$$

Since

$$|w(0)|_2^2 = \int_{\Omega} w(0)^2 dx = \int_{\Omega} u(0)^{2r} dx \leq |\Omega| |u(0)|_{\infty}^{2r} \leq |\Omega| \tilde{U}_r^{2r}, \quad (65)$$

(64) and (65) gives us

$$\tilde{U}_{2r}^{2r} \leq |\Omega| \tilde{U}_r^{2r} + \|f\|_{L^\infty(0, \infty, L^q(\Omega))}^{\frac{1}{\delta}} \left(\frac{2r}{m}\right)^{2r\sigma(r)} m C^{\frac{2}{\delta}} \tilde{U}_r^{2r}. \quad (66)$$

Whereas $\frac{1}{\delta} > 1$, $2r\sigma(r) > 0$ and $\sigma(r) = \frac{1}{2r\delta}$ it follows that

$$\tilde{U}_{2r} \leq C_2^{\sigma(r)} \|f\|_{L^\infty(0, \infty, L^q(\Omega))}^{\sigma(r)} r^{\sigma(r)} \tilde{U}_r, \quad (67)$$

with $C_2 = C_2(\Omega, m)$. This completes the proof of Lemma. \square

We have also

Lemma 4.3. *Let $r > 1$, $n \geq 3$, $p < \frac{n}{n-2}$ and $\sigma(r) = \frac{p(n+2)}{2[r(2p-pn+n)+np]}$ then we get*

$$\sigma(2^k r) \leq \theta^k \sigma(r) \quad \forall k \in \mathbb{N},$$

with $\theta \in (0, 1)$.

Proof. By asking $c_1 = \frac{p(n+2)}{2}$, $c_2 = (2p - pn + n)$ and $c_3 = np$ yields $\sigma(r) = \frac{c_1}{rc_2+c_3}$ with $c_1, c_2, c_3 \in \mathbb{R}_+^*$. By taking $\theta = 1 - \frac{c_2}{2c_2+c_3}$ the proof of this lemma is deduced by reasoning by induction. \square

Returning now to the proof of the theorem, by lemma 4.2 we get

Proof.

$$\tilde{U}_{2r} \leq [C_2 \|f\|_{L^\infty(0, \infty, L^q(\Omega))}]^{\sigma(r)} r^{\sigma(r)} \tilde{U}_r.$$

By iterating this equation by taking $r = h, r = 2h, r = 2^2h, \text{etc}$, we obtain

$$\tilde{U}_{2^{k+1}h} \leq [C_2 \|f\|_{L^\infty(0, \infty, L^q(\Omega))}]^{\lambda_1} 2^{\lambda_2} h^{\lambda_1} \tilde{U}_h,$$

with

$$\lambda_1 := \sigma(h) + \sigma(2h) + \sigma(2^2h) + \dots + \sigma(2^{k-1}h) + \sigma(2^k h)$$

et

$$\lambda_2 := \sigma(2h) + 2\sigma(2^2h) + 3\sigma(2^3h) + \dots + (k-1)\sigma(2^{k-1}h) + k\sigma(2^k h).$$

To complete the proof we just need to show that $\lambda_1, \lambda_2 < +\infty$. Indeed by lemma 4.3

$$\lambda_1 \leq \sum_{\mu=0}^k \alpha^\mu \sigma(h) \leq \sum_{\mu=0}^{\infty} \alpha^\mu \sigma(h) = \frac{\sigma(h)}{(1-\alpha)} < \infty.$$

Noting also that

$$\sigma(2^k h) \leq \theta^{k-1} \sigma(2h) \quad \forall k \in \mathbb{N}^*,$$

it follows

$$\lambda_2 \leq \sum_{\mu=1}^k \mu \alpha^{\mu-1} \sigma(2h) \leq \sum_{\mu=1}^{\infty} \mu \alpha^{\mu-1} \sigma(2h) = \frac{\sigma(2h)}{(1-\alpha)^2} < \infty.$$

This completes the proof of the theorem. \square

4.2 Uniform estimate in time

We prove in what follows an estimate of u in $L^\infty(\mathbb{R}^+, H_0^1(\Omega))$. We get

Theorem 4.4. *Assume that $f \in L^2(\Omega)$, $g \in H^1(\Omega)$, $u_0 \in H_0^1(\Omega)$ and $a \in W^{1,\infty}(\mathbb{R})$ with $\inf_{\mathbb{R}} a > 0$. Then a solution u of (1) is such that $u \in L^\infty(\mathbb{R}^+, H_0^1(\Omega))$.*

Taking a spectral basis related to the Laplace operator in the Galerkin approximation (see [16]) we find that $-\Delta u$ can be regarded as test function in $L^2(0, T, L^2(\Omega))$ for all $T > 0$. By multiplying (1) by $-\Delta u(t)$ and integrating over Ω

$$(u_t, -\Delta u) + (-\operatorname{div}(a(l_r(u))\nabla u), -\Delta u) = (f, -\Delta u), \quad (68)$$

and also

$$\frac{1}{2} \frac{d}{dt} \|u\|_V^2 + (-a(l_r(u))\Delta u, -\Delta u) + (-a'(l_r(u))\nabla l_r(u) \cdot \nabla u, -\Delta u) = (f, -\Delta u). \quad (69)$$

Here (\cdot, \cdot) is the usual scalar product on $L^2(\Omega)$. Taking to account

$$|\nabla l_r(u)|_2 \leq K \|g\|_{H^1(\Omega)} |\nabla u|_2, \quad (70)$$

where K is a constant depending of Ω . It comes

$$|(-a'(l_r(u))\nabla l_r(u) \cdot \nabla u, -\Delta u)| \leq K \|g\|_{H^1(\Omega)} |a'|_\infty \|u\|_V^2 |\Delta u|_2 \quad (71)$$

Now from (71) and (69) we have

$$\frac{1}{2} \frac{d}{dt} \|u\|_V^2 + m |\Delta u|_2^2 - K \|g\|_{H^1(\Omega)} |a'|_\infty \|u\|_V^2 |\Delta u|_2 \leq |f|_2 |\Delta u|_2. \quad (72)$$

By using Young's inequality $ab \leq \frac{1}{2m} a^2 + \frac{m}{2} b^2$, we get

$$\frac{d}{dt} \|u\|_V^2 \leq \frac{1}{m} |f|_2^2 + \frac{1}{m} (K \|g\|_{H^1(\Omega)})^2 \|a'\|_\infty^2 \|u\|_V^4. \quad (73)$$

In order to apply the uniform Gronwall lemma to (73) we start with a small estimate. Remember that

$$\frac{d}{dt} |u|_2^2 + m \|u\|_V^2 \leq \frac{1}{\lambda m} |f|_2^2, \quad (74)$$

where λ is the principal eigenvalue of the Laplacian operator with Dirichlet boundary conditions.

By integrating on $[t, t_0]$ we get

$$|u(t+t_0)|_2^2 + m \int_t^{t+t_0} \|u\|_V^2 ds \leq \int_t^{t+t_0} \frac{1}{\lambda m} |f|_2^2 ds + |u(t)|_2^2, \quad (75)$$

and also

$$\int_t^{t+t_0} \|u\|_V^2 ds \leq \frac{t_0}{\lambda m^2} |f|_2^2 + \frac{1}{m} |u(t)|_2^2. \quad (76)$$

Let $\rho_0 > 0$ such that $|u(t)|_2^2 \leq \rho_0^2$. By setting

$$a_1 = \frac{1}{m} c_1(\Omega)^2 |a'|_\infty^2 a_3 \quad a_2 = \frac{t_0}{m} |f|_2^2 \quad a_3 = \frac{t_0 \lambda}{m^2} |f|_2^2 + \frac{1}{m} \rho_0^2,$$

we obtain by using uniform Gronwall lemma to (73)

$$\|u(t+t_0)\|_V \leq \left(\frac{a_3}{t_0} + a_2\right) \exp(a_1) \quad \forall t \geq 0, \quad t_0 > 0. \quad (77)$$

Hence $u \in L^\infty(t_0, +\infty, H_0^1(\Omega))$. By using (73) and the classical Gronwall lemma it is easy to see that $u \in L^\infty(0, t_0, H_0^1(\Omega))$. This completes the proof of the theorem.

Remark 5. This theorem show us the existence of absorbing set in $H_0^1(\Omega)$. By considering $S(t)$ the semigroup associated to the equation (1) defined by

$$\begin{aligned} S(t) : L^2(\Omega) &\rightarrow L^2(\Omega) \\ u_0 &\mapsto u(t), \end{aligned}$$

with $u(t)$ a solution of (1). As a result of the theorem 4.4 and the compact embedding of $H_0^1(\Omega)$ into $L^2(\Omega)$ we deduce that the semigroup $S(t)$ possesses a global attractor. Indeed it is easy to show the existence of absorbing set in $L^2(\Omega)$, the main difficulty here is to show that $S(t)$ is such that

$$\begin{aligned} \forall B \subset L^2(\Omega) \text{ bounded, } \exists t_0 = t_0(B) \\ \text{such that } \bigcap_{t \geq t_0} \bigcup S(t)B \text{ is relatively compact in } L^2(\Omega). \end{aligned} \quad (78)$$

This property known as $S(t)$ is uniformly compact for t large can be proved by using theorem 4.4 and the compact embedding of $H_0^1(\Omega)$ into $L^2(\Omega)$.

4.3 Asymptotic behaviour

In this part we are interested in asymptotic behaviour of a weak solution of (1). Our main interest here is the radial solutions. By radial solutions we means $\tilde{u}(t, |x|) = u(t, x)$. As in the stationary case Ω is a open ball of \mathbb{R}^n . Remember that

$$L_r^2(\Omega) = \{v \in L^2(\Omega) \quad \exists \tilde{v} \in L^2(]0, d/2[) \text{ such that } v(x) = \tilde{v}(\|x\|)\}.$$

In order to not make confusion between u_0 the solution to (P_0) and the initial value of (1), we will take u^0 the initial value of (1).

Theorem 4.5. Assume that $f, g \in L_r^2(\Omega)$, a is a continuous function and the assumption (3) checked then (1) admits a radial solution.

Proof. Let $w \in L^2(0, t, L_r^2(\Omega))$ it is clear that $l_r(w)$ is radial and also $a(l_r(w))$. Thus by (8) F_r maps $L^2(0, t, L_r^2(\Omega))$ into itself. The proof now follows by using arguments similar to those used in theorem 2.1. \square

Assume now

$$f, g \geq 0 \quad \text{in } \Omega \quad (79)$$

and

$$u_0 \leq u^0 \leq u_d, \quad (80)$$

with u^0 the initial value to (1) and u_0 and u_d respectively the solution of (P_0) and (P_d) .

We can now give a stability result assuming that (1) admits a unique solution.

Theorem 4.6. *Assume (79) and $f, g \in L^2_r(\Omega)$. Let u , u_d and u_0 respectively the solution of (1), (P_d) and (P_d) . If*

$$u_0 \leq u^0 \leq u_d,$$

then

$$u_0 \leq u \leq u_d \quad \forall t.$$

Proof. Let

$$\mathcal{S} = \{t \mid l(u(s)) \in [0, l_d(u_d)] \quad \forall s \leq t\}. \quad (81)$$

It is easy to prove that \mathcal{S} contains 0 (see 80). By setting

$$t^* = \sup\{t \mid t \in \mathcal{S}\}. \quad (82)$$

By continuity of the mapping $t \rightarrow l_d(u(t))$, we get

$$l_d(u(t^*)) \in [0, l_d(u_d)]. \quad (83)$$

By using (1) and (P_d) we get in $\mathcal{D}(0, t^*)$

$$\frac{d}{dt}(u_d - u, \phi) + \int_{\Omega} a(l_d(u)) \nabla(u_d - u) \nabla \phi = - \int_{\Omega} (a(l_d(u_d)) - a(l_d(u))) \nabla u_d \nabla \phi \quad \forall \phi \in H_0^1(\Omega). \quad (84)$$

Choising $\phi = (u_d - u)^-$, (84) becomes

$$\frac{1}{2} \frac{d}{dt} |(u_d - u)^-|_2^2 + \int_{\Omega} a(l_d(u_d)) |\nabla(u_d - u)^-|^2 = \int_{\Omega} (a(l_d(u)) - a(l_d(u_d))) \nabla u_d \nabla (u_d - u)^-. \quad (85)$$

Since a is nonincreasing $(a(l_d(u)) - a(l_d(u_d))) \geq 0 \quad \forall t \leq t^*$ hence proposition 3.2 yields

$$\int_{\Omega} (a(l_d(u)) - a(l_d(u_d))) \nabla u_d \nabla (u_d - u)^- \leq 0. \quad (86)$$

Thus

$$\frac{1}{2} \frac{d}{dt} |(u_d - u)^-|_2^2 + a(l_d(u_d)) |\nabla(u_d - u)^-|_2^2 \leq 0 \quad (87)$$

Applying Poincaré Sobolev inequality we get

$$\frac{1}{2} \frac{d}{dt} |(u_d - u)^-|_2^2 + C_2 |(u_d - u)^-|_2^2 \leq 0, \quad (88)$$

this prove

$$\frac{d}{dt} \{e^{2t C_2} |(u_d - u)^-|_2^2\} \leq 0. \quad (89)$$

Moreover $(u_d - u)^-(0) = (u_d - u^0)^- = 0$ it follows that $u_d \geq u \quad \forall t \in [0, t^*]$. In the same way we can also prove $u_0 \leq u \quad \forall t \in [0, t^*]$. It follows

$$u_0 \leq u \leq u_d \quad \forall t \in [0, t^*] \quad (90)$$

To finish we just need to prove that t^* is very large, this is typically the case. Indeed if $t^* < \infty$ then

$$l(u(t^*)) = 0 \quad \text{or} \quad l_d(u_d). \quad (91)$$

From (79) and (90) we deduce

$$u(t^*) = u_0 \quad \text{or} \quad u(t^*) = u_d. \quad (92)$$

Due to the uniqueness of (1), we deduce that $t = \infty$. This shows that

$$u_0 \leq u \leq u_d \quad \forall t,$$

and achieve the proof. □

Remark 6. The fact that $|u(t)|_2^2$ is not a Lyapunov function that is to say decreases in time, makes very complex the study of certain asymptotic properties of our problem. Indeed under our study it is tempting to show that for r fixed $r \in]0, d[$

$$u(t) \rightarrow u_r^1 \quad \text{in} \quad L^2(\Omega),$$

where u is the solution of (1) and u_r^1 the solution belonging to the stable global branch described previously. A numerical study would be a great contribution to try to carry out some of our theoretical intuitions.

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