open science

# Almost sure asymptotics for the maximum local time in Brownian environment <br> Roland Diel 

## To cite this version:

Roland Diel. Almost sure asymptotics for the maximum local time in Brownian environment. 2010. hal-00460810v1

## HAL Id: hal-00460810 <br> https://hal.science/hal-00460810v1

Preprint submitted on 2 Mar 2010 (v1), last revised 15 Jun 2011 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Almost sure asymptotics for the maximum local time in Brownian environment 

Roland DIEL *


#### Abstract

We study the asymptotic behaviour of the maximum local time $\mathbf{L}^{*}(t)$ of the Brox's process, the diffusion in Brownian environment. Shi 15 proved that the maximum speed of $\mathbf{L}^{*}(t)$ is surprisingly, at least $t \log (\log (\log t))$ whereas in the discrete case it is $t$. We show here that $t \log (\log (\log t))$ is the proper rate and we prove that for the minimum speed the rate is the same as in the discrete case (Dembo, Gantert, Peres and Shi 5]) namely $t(\log (\log (\log t)))^{-1}$.


## 1 Introduction

Let $(W(x), x \in \mathbf{R})$ be a two-sided one-dimensional Brownian motion on $\mathbf{R}$ with $W(0)=0$. We call diffusion process in the environment $W$ a process $\left(\mathbf{X}(t), t \in \mathbf{R}^{+}\right)$whose infinitesimal generator given $W$ is

$$
\frac{1}{2} e^{W(x)} \frac{d}{d x}\left(e^{-W(x)} \frac{d}{d x}\right)
$$

Notice that if we take $W$ differentiable, $\left(\mathbf{X}(t), t \in \mathbf{R}^{+}\right)$would be the solution of the following stochastic differential equation

$$
\left\{\begin{array}{l}
\mathrm{d} \mathbf{X}(t)=\mathrm{d} \beta(t)-\frac{1}{2} W^{\prime}(\mathbf{X}(t)) \mathrm{d} t \\
\mathbf{X}(0)=0
\end{array}\right.
$$

in which $\beta$ is a standard one-dimensional Brownian motion independent of $W$. Of course as we choose for $W$ a Brownian motion, the previous equation does not have any rigorous sense but it explains the denomination environment for $W$.

[^0]This process was first introduced by S. Schumacher [14] and T. Brox [3]. The diffusion $\mathbf{X}$ is recurrent and Brox shows in [3] that it is sub-diffusive with asymptotic behaviour in $(\log t)^{2}$. Moreover for a given instant $t, \mathbf{X}$ has the property to be localized in the neighborhood of a random point $m_{\log t}$ only depending on $t$ and $W$. The limit law of $m_{\log t} /(\log t)^{2}$ and therefore of $\mathbf{X}_{t} /(\log t)^{2}$ is independently made explicit by H. Kesten [1]] and A. O. Golosov [8].
H. Kesten and A. O. Golosov's aim was to determine the limit law of the discrete time and space analogue of Brox's model introduced by F. Solomon [18] and then studied by Ya. G. Sinai [17]. This random walk in random environment, called Sinai's walk, $\left(S_{n}, n \in \mathbf{N}\right)$ actually has the same limit distribution as Brox's.

Turning back to Brox's diffusion, one can note that H. Tanaka 21, 20] obtained a deeper localization and later Y. Hu and Z. Shi [9] got the almost sure rates of convergence of the limsup and $\lim \inf$ of $X$. It appears that these rates are the same as the ones for Sinai's walk. The question of an invariance principle between these two processes remains open (see Z. Shi (16] for a survey).

In the present paper, we are interested in the asymptotics of the maximum of the local time of $\mathbf{X}$. Indeed, to the diffusion $\mathbf{X}$ corresponds a local time process ( $\mathbf{L}(t, x), t \geq 0, x \in \mathbf{R}$ ) defined by the occupation time formula: $\mathbf{L}$ is the unique a.s. jointly continuous process such that for any bounded Borel function $f$ and for any $t \geq 0$,

$$
\int_{0}^{t} f\left(\mathbf{X}_{s}\right) \mathrm{d} s=\int_{\mathbf{R}} f(x) \mathbf{L}(t, x) \mathrm{d} x
$$

We shall see later that such a process exists. The first results on the behavior of $\mathbf{L}$ can be found in [15] and [10]. In particular it is proven in (10] that for any $x \in \mathbf{R}$,

$$
\frac{\log (\mathbf{L}(t, x))}{\log t} \xrightarrow{\mathcal{L}} U \wedge \hat{U}, t \rightarrow+\infty
$$

where $U$ and $\hat{U}$ are independent variables uniformely distributed in $(0,1)$ and $\xrightarrow{\mathcal{L}}$ is the convergence in law. Note that in the same paper Hu and Shi also prove that Sinai's walk local time has the same behaviour.

We now consider the maximum of the local time of $\mathbf{X}$,

$$
\mathbf{L}^{*}(t):=\sup _{x \in \mathbf{R}} \mathbf{L}(t, x) .
$$

Shi was the first one to be interested in this process; in (15) he gives a lower bound on the lim sup behaviour,

$$
\limsup _{t \rightarrow \infty} \frac{\mathbf{L}^{*}(t)}{t \log _{3}(t)} \geq \frac{1}{32}
$$

where for any $i \in \mathbf{N}, \log _{i+1}=\log \circ \log _{i}$ and $\log _{1}=\log$. In the same paper, he computed the similar rate in the discrete case: there is constant $c \in(0, \infty)$ such that

$$
\limsup _{n \rightarrow \infty} \frac{L_{S}^{*}(n)}{n}=c
$$

Thereby he highlighted a different behaviour for the discrete model and for the continuous one. The limit law of $\mathbf{L}^{*}(t)$ was then determined in [1]):

$$
\frac{\mathbf{L}^{*}(t)}{t} \xrightarrow{\mathcal{L}} \frac{1}{\int_{-\infty}^{\infty} e^{-\tilde{W}(x)} \mathrm{d} x}, t \rightarrow+\infty
$$

where $\tilde{W}$ has the same law as the environment conditioned to stay positive. But unlike the discrete case (see (7]), this result does not allow to give an upper bound on the almost sure behaviour.

Here we prove that the quantity $\log _{3}(t)$ is the correct renormalization for the limsup and an analog result for the liminf:

Theorem 1.1 Almost surely,

$$
\begin{aligned}
& \frac{1}{32} \leq \limsup _{t \rightarrow \infty} \frac{\mathbf{L}^{*}(t)}{t \log _{3}(t)} \leq 2 e \text { and } \\
& \frac{j_{0}^{2}}{64} \leq \liminf _{t \rightarrow \infty} \frac{\log _{3}(t) \mathbf{L}^{*}(t)}{t} \leq e \pi^{2}
\end{aligned}
$$

where $j_{0}$ is the smallest strictly positive root of Bessel function $J_{0}$.
We can compare these results to the ones of the discrete case, see (15] (and $\sqrt[7]{ }$ for the value of the constant) for the limsup and (5) for the liminf: there exists two constants $c, c^{\prime} \in(0, \infty)$ such that almost surely,

$$
\limsup _{n \rightarrow \infty} \frac{L_{S}^{*}(n)}{n}=c \quad \text { and } \quad \liminf _{n \rightarrow \infty} \frac{\log _{3}(n) L_{S}^{*}(n)}{n}=c^{\prime} .
$$

As Shi noted it, the limsup in the continuous case and in the discrete case have a different normalization, but we see with Theorem 1.1 that the normalization is the same for the liminf. There is the following heuristic interpretation: in the discrete case like in the continuous one, $\mathbf{L}^{*}(t) / t$ behaves approximately in the best case like the integral of the exponential of the "steepest" environment and in the worst case like the integral of the exponential of the "flattest" one. However if in discrete case there is a steepest environment, in the continuous case it is possible to be as steep as we want, and in the two cases, the environment can be as flat as we want. This explains the difference for the lim sup. To compute the exact rate, we need to study carefully both the behaviour of the "steepest" and "flattest" environment and the frequency of the change of the favourite point.

We now present a useful representation of $\mathbf{X}$ we use throughout the article to study the process : the martingale representation theorem says
that, given the environment $W, \mathbf{X}$ is a Brownian motion rescaled in time and space. Precisely, there is a Brownian motion $B$ started at 0 , independent of $W$ such that if we define the scale function

$$
\begin{equation*}
\forall x \in \mathbb{R}, S_{W}(x):=\int_{0}^{x} e^{W(y)} \mathrm{d} y \tag{1}
\end{equation*}
$$

and the random time change

$$
\begin{equation*}
\forall t \geq 0, T_{W, B}(t):=\int_{0}^{t} e^{-2 W\left(S_{W}^{-1}(B(s))\right)} \mathrm{d} s \tag{2}
\end{equation*}
$$

then

$$
\begin{equation*}
\mathbf{X}=S_{W}^{-1} \circ B \circ T_{W, B}^{-1} \tag{3}
\end{equation*}
$$

To simplify notations, we write when there is no possible confusion $S$ and $T$ for respectively $S_{W}$ and $T_{W, B}$. We also denote by $P$ the total probability and expectation and by $P^{W}$ the probability and expectation given the environment $W$. Using (3), we easily obtain that the local time has to be defined by, for any $x \in \mathbf{R}$ and $t \geq 0$,

$$
\begin{equation*}
\mathbf{L}(t, x)=e^{-W(x)} L\left(T^{-1}(t), S(x)\right) \tag{4}
\end{equation*}
$$

where $L$ is the local time process of the Brownian motion $B$.
The rest of the paper is organized as follows: in Section 2 , we present some technical estimates used in the article, in Section 3 , we show that environment $W$ has "good" properties with a large enough probability, Section $4{ }^{4}$ is centered on the study of the local time of the process $\mathbf{X}$, then in Section 5, we look at the behaviour of the maximum of the local time process $\mathbf{L}^{*}$ and finally in Section 6, Theorem 1.1 is proved.

## 2 Three useful theorems and some technical estimates

Here are presented some results used throughout the article. We begin this section with some notations :

$$
\begin{array}{rll}
\forall x \geq 0, & \boldsymbol{\tau}(x) & :=\inf \{t \geq 0 / \mathbf{X}(t)=x\}, \\
\forall x \geq 0, & \tau(x) & :=\inf \{t \geq 0 / B(t)=x\}, \\
\forall x \geq 0, \forall r \geq 0, & \boldsymbol{\sigma}(r, x) & :=\inf \{t \geq 0 / \mathbf{L}(t, x) \geq r\} \text { and } \\
\forall x \geq 0, \forall r \geq 0, & \sigma(r, x) & :=\inf \{t \geq 0 / L(t, x) \geq r\} .
\end{array}
$$

More generally, the quantities related to $\mathbf{X}$ are written in bold font and the ones related to $B$ in normal font. Furthermore, in the rest of the paper, letter $K$ stands for a universal constant whose value can change from one line to another and for any process $M$, we denote by $\tau_{M}(x)$ the hitting time of height $x$ by $M$.

As we saw in the previous section, the study of the process $\mathbf{X}$ is reduced by scaling in time and space to the study of a Brownian motion. We therefore use in our proofs the following two theorems that describe the law of the local time of a Brownian motion stopped at some properly chosen random times:

Theorem 2.1 (Ray-Knight) Let a be a positive real number.
The process $\{L(\tau(a), a-y), y \in[0, a]\}$ is a square of a 2 -dimensional Bessel process started at 0 . And conditionally on $L(\tau(a), 0),\{L(\tau(a),-y), y \geq 0\}$ is a square of a 0 -dimensional Bessel process started at $L(\tau(a), 0)$ independent of $\{L(\tau(a), y), y \geq 0\}$.

Theorem 2.2 (Ray-Knight) Let $r$ be a positive real number.
The processes $\left\{L(\sigma(r, 0), y), y \in \mathbf{R}^{+}\right\}$and $\left\{L(\sigma(r, 0),-y), y \in \mathbf{R}^{+}\right\}$are two independent squares of 0 -dimensional Bessel processes started at $r$.

In the rest of the article, we denote by $Z$ a square of a 0 -dimensional Bessel process started at 1 and by $Q$ a square of a 2-dimensional Bessel process started at 0 . To use the previous theorems, we must estimate the behaviour of these two processes. This is done in the following lemma.

Lemma 2.3 For all $v, \delta, M>0$, we have

1. $P\left(\sup _{0 \leq t \leq v}|Z(t)-1| \geq \delta\right) \leq 4 \frac{\sqrt{(1+\delta) v}}{\delta} \exp \left(-\frac{\delta^{2}}{8(1+\delta) v}\right)$,
2. $P\left(\sup _{t \geq 0} Z(t) \geq M\right)=\frac{1}{M}$,
3. $P\left(\sup _{0 \leq t \leq v} Q(t) \geq M\right) \leq 4 e^{-\frac{M}{2 v}}$,
4. There is a constant $K>0$ such that for any $0<a<b$,

$$
P\left(\sup _{a \leq t \leq b} \frac{1}{t} Q(t) \geq M\right) \leq K e^{-\frac{M}{2 \log (8 b / a)}} .
$$

## Proof:

11 and 2. These are proved in [19] Lemma 3.1.
3. Denote by $B$ and $\tilde{B}$ two independent Brownian motions. The processes $Q$ and $B^{2}+\tilde{B}^{2}$ have the same law, so

$$
\begin{aligned}
P\left(\sup _{0 \leq t \leq v} Q(t) \geq M\right) & \leq P\left(\left(\sup _{0 \leq t \leq v}|B|(t)\right)^{2}+\left(\sup _{0 \leq t \leq v}|\tilde{B}|(t)\right)^{2} \geq M\right) \\
& \leq 4 P\left(\left(\sup _{0 \leq t \leq v} B(t)\right)^{2}+\left(\sup _{0 \leq t \leq v} \tilde{B}(t)\right)^{2} \geq M\right)
\end{aligned}
$$

where the second inequality comes from reflection principle. And thus,

$$
P\left(\sup _{0 \leq t \leq v} Q(t) \geq M\right) \leq 4 P(Q(v) \geq M)=4 e^{-\frac{M}{2 v}}
$$

as $Q(v)$ has exponential distribution of parameter $2 v$.
4. Thanks to the scaling property of $Q$,

$$
\begin{aligned}
P\left(\sup _{a \leq t \leq b} \frac{1}{t} Q(t) \geq M\right) & =P\left(\sup _{a / b \leq t \leq 1} \frac{1}{t} Q(t) \geq M\right) \\
& \leq P\left(\sup _{a / b \leq t \leq 1} \frac{Q(t)}{t \log (8 / t)} \geq \frac{M}{\log (8 b / a)}\right) \\
& \leq P\left(\sup _{0 \leq t \leq 1} \sqrt{\frac{Q(t)}{t \log (8 / t)}} \geq \sqrt{\frac{M}{\log (8 b / a)}}\right) .
\end{aligned}
$$

We then conclude with Lemma 6.1 in [9] which says that there exist a constant $K>0$ such that for all $x>0$,

$$
P\left(\sup _{0 \leq t \leq 1} \sqrt{\frac{Q(t)}{t \log (8 / t)}} \geq x\right) \leq K e^{-\frac{x^{2}}{2}}
$$

We also need to study the behaviour of the environment $W$. Let us start with a notation for the minimum of $W$ on an interval

$$
\underline{W}(x, y):=\left\{\begin{array}{lc}
\min _{z \in[x, y]} W(z) & \text { if } x \leq y \\
+\infty & \text { if not }
\end{array}\right.
$$

and for the maximum

$$
\bar{W}(x, y):=\left\{\begin{array}{lc}
\max _{z \in[x, y]} W(z) & \text { if } x \leq y \\
-\infty & \text { if not }
\end{array}\right.
$$

We can now define

$$
H_{v}:=\inf \{x \geq 0 ; W(x)-\underline{W}(0, x) \geq v\}
$$

and

$$
m_{v}:=\inf \left\{x \geq 0 ; W(x)=\underline{W}\left(0, H_{v}\right)\right\}
$$

Brox showed in [3] that at time $e^{v}$, with high probability, the process $\mathbf{X}$ has spent much of its time in the neighborhood of $m_{v}$ or of the symmetric point on the negative real halfline. For our study, we need to know the law of the environment in the neighborhood of these points ; it is given by the following theorem due to Tanaka:

Theorem 2.4 (Tanaka [21]) Let $R$ be a Bessel process of dimension 3 started at 0 and

$$
\begin{aligned}
\tau_{R}(v) & :=\inf \{x \leq 0 / R(x) \geq v\} \\
\zeta_{R}(v) & :=\inf \left\{x \leq 0 / R(x)-\inf _{y \geq x} R(y) \geq v\right\} \text { and } \\
\rho_{R}(v) & :=\sup \left\{x \leq \zeta_{R}(v) / R(x)-\inf _{y \geq x} R(y)=0\right\}
\end{aligned}
$$

Under $P$, the processes $\left(W\left(-x+m_{v}\right)-W\left(m_{v}\right)\right)_{x \in\left[0, m_{v}\right]}$ and $\left(W\left(x+m_{v}\right)-\right.$ $\left.W\left(m_{v}\right)\right)_{x \in\left[0, H_{v}\right]}$ are independent and the following equality in law is true :

$$
\left(W\left(-x+m_{v}\right)-W\left(m_{v}\right)\right)_{x \in\left[0, m_{v}\right]} \stackrel{\mathcal{L}}{=}(R(x))_{x \in\left[0, \rho_{R}(v)\right]}
$$

and

$$
\left(W\left(x+m_{v}\right)-W\left(m_{v}\right)\right)_{x \in\left[0, H_{v}\right]} \stackrel{\mathcal{L}}{=}(R(x))_{x \in\left[0, \tau_{R}(v)\right]} .
$$

Therefore, to use this theorem it is necessary to have informations on the behaviour of Bessel process of dimension 3. It is done in the following lemma:

Lemma 2.5 Let $R$ be a 3-dimensional Bessel process started in 0. There is a positive real number $K$ such that for every $a, t>0$,

$$
\frac{a}{\sqrt{t}} e^{-a^{2} / 2 t} \leq P\left(\sup _{[0, t]} R>a\right) \leq K\left(\frac{a}{\sqrt{t}}+\frac{\sqrt{t}}{a}\right) e^{-a^{2} / 2 t}
$$

Proof: According to the reflection principle for Brownian motion, one can find a $K>0$ such that

$$
P\left(\sup _{[0, t]} R>a\right) \leq K P(R(t)>a)
$$

Moreover, $P\left(\sup _{[0, t]} R>a\right) \geq P(R(t)>a)$. So the estimates of the lemma are a consequence of usual estimates for 3-dimensional Bessel process.

## 3 Estimates on the environment

We start with the definition of the environment reversed in time

$$
(\widehat{W}(x) ; \quad x \in \mathbf{R}):=(W(-x) ; \quad x \in \mathbf{R})
$$

which is still a Brownian motion. This allows us to restrict the study to $W$ on $\mathbf{R}^{+}$and to get similar results on $\mathbf{R}^{-}$by symmetry.

We study the environment on $\left[0, H_{v}\right]$ and particularly in the neighborhood of $m_{v}$ as it is the place where the diffusion spends most of its time.

Unfortunately, the probability that, at time $e^{v}$, the process has reached the bottom of the valley $m_{v}$ and, secondly it has not left the valley, tends to 1 but is not growing fast enough to derive almost sure results. So we rather study valleys of height $v-c_{1} \log v$, where $c_{1}$ is a positive real number whose value will be determined later, so that, with high probability, at time $e^{v}$, the process has reached the bottom of the valley. We are also interested in the valley of height $v+c_{3} \log v$ so that the process has not left this valley at time $e^{v}$. We therefore fix three constants $c_{1}, c_{2}, c_{3}>0$ and define recursively for any $v>1$,

$$
\begin{array}{rlrl}
b_{v, 0}^{-} & :=0 \\
\forall i \geq 0, b_{v, i+1}^{-} & :=\inf \left\{x \geq b_{v, i}^{-} ;\right. & & \left.W(x)-\underline{W}\left(b_{v, i}^{-}, x\right) \geq v-c_{1} \log v\right\} \\
m_{v, i+1}^{-} & :=\inf \left\{x \geq b_{v, i}^{-} ;\right. & \left.W(x)=\underline{W}\left(b_{v, i}^{-}, b_{v, i+1}^{-}\right)\right\} \text {and }
\end{array}
$$

We then have to specify for any $i \in \mathbf{N}^{*}$,

$$
\begin{aligned}
a_{v, i}^{-} & :=\sup \left\{x \leq m_{v, i}^{-} ; W(x)-W\left(m_{v, i}^{-}\right) \geq v-c_{2} \log v\right\} \vee b_{v, i-1}^{-} \\
c_{v}^{+} & :=\inf \left\{x \geq 0 ; W(x)-W(0, x) \geq v+c_{3} \log v\right\} \\
m_{v}^{+} & :=\inf \left\{x \geq 0 ; W(x)=\underline{W}\left(0, c_{v}^{+}\right)\right\} \\
b_{v}^{+} & :=\inf \left\{x \geq m_{v}^{+} ; W(x)-W\left(m_{v}^{+}\right) \geq v-c_{1} \log v\right\} \text { and } \\
a_{v}^{+} & :=\sup \left\{x \leq m_{v}^{+} ; W(x)-W\left(m_{v}^{+}\right) \geq v-c_{2} \log v\right\} \vee 0
\end{aligned}
$$



Figure 1: A sample path of $W$
There is a $i \in \mathbf{N}^{*}$ such that $m_{v}^{+}=m_{v, i}^{-}$. We want to prove that, with a sufficiently high probability, $m_{v}^{+} \in\left\{m_{v, 1}^{-} ; m_{v, 2}^{-}\right\}$and moreover that, after $m_{v}^{+}$, there is no point of depth close to $W\left(m_{v}^{+}\right)$in the valley $\left[0, c_{v}^{+}\right]$. This can be expressed formally as follows:

$$
\Gamma_{v}^{1}:=\left\{c_{v}^{+} \leq b_{v, 3}^{-} ; \underline{W}\left(b_{v}^{+}, c_{v}^{+}\right)-W\left(m_{v}^{+}\right) \geq\left(c_{1}+c_{3}\right) \log v\right\}
$$

We also need many more technical conditions to ensure that the environment does not stray too far from its average behaviour:

$$
\begin{gathered}
\Gamma_{v}^{2}:=\left\{b_{v, 2}^{-} \leq v^{6} ; W\left(m_{v, 1}^{-}\right) \geq-v^{2} ; W\left(m_{v, 2}^{-}\right)-W\left(b_{v, 1}^{-}\right) \geq-v^{2}\right\} \\
\bigcap\left\{m_{v, 1}^{-}-a_{v, 1}^{-} \geq \frac{1}{v^{2}} ; m_{v, 2}^{-}-a_{v, 2}^{-} \geq \frac{1}{v^{2}}\right\} \\
\bigcap\left\{c_{v}^{+}-m_{v}^{+} \geq v ; W\left(c_{v}^{+}\right)-\underline{W}\left(\left(c_{v}^{+}-\log v\right) \vee m_{v}^{+}, c_{v}^{+}\right) \leq 2 \log v\right\} \\
\bigcap\left\{\bar{W}\left(\left(m_{v, 1}^{-}-\log v\right) \vee a_{v, 1}^{-}, m_{v, 1}^{-}\right)-W\left(m_{v, 1}^{-}\right) \leq 2 \log v\right\} \\
\bigcap\left\{\bar{W}\left(\left(m_{v, 2}^{-}-\log v\right) \vee a_{v, 2}^{-}, m_{v, 2}^{-}\right)-W\left(m_{v, 2}^{-}\right) \leq 2 \log v\right\}
\end{gathered}
$$

We would also wish sometimes that $m_{v}^{+}=m_{v, 1}^{-}$(which of course is obtained with a probability lower than the previous one). For that we use the event:

$$
\Gamma_{v}^{3}:=\left\{c_{v}^{+} \leq b_{v, 2}^{-} ; \underline{W}\left(b_{v}^{+}, c_{v}^{+}\right)-W\left(m_{v}^{+}\right) \geq\left(c_{1}+c_{3}\right) \log v\right\}
$$

We define $\widehat{\Gamma}_{v}^{1}, \widehat{\Gamma}_{v}^{2}$ and $\widehat{\Gamma}_{v}^{3}$ similarly from $\widehat{W}$.
We therefore work on

$$
\begin{equation*}
\Gamma_{v}=\Gamma_{v}^{1} \cap \Gamma_{v}^{2} \cap \widehat{\Gamma}_{v}^{1} \cap \widehat{\Gamma}_{v}^{2} \tag{5}
\end{equation*}
$$

or

$$
\begin{equation*}
\Gamma_{v}^{\prime}=\Gamma_{v}^{3} \cap \Gamma_{v}^{2} \cap \widehat{\Gamma}_{v}^{3} \cap \widehat{\Gamma}_{v}^{2} \tag{6}
\end{equation*}
$$

The first step as stated before is to show that these events occur with a high enough probability. We denote for every event $A, \bar{A}:=\Omega \backslash A$.

Proposition 3.1 There exists a constant $K>0$ such that for any $v \geq 2$,

$$
P\left(\bar{\Gamma}_{v}\right) \leq K\left(\frac{\log v}{v-c_{1} \log v}\right)^{2} \text { and } P\left(\bar{\Gamma}_{v}^{\prime}\right) \leq \frac{K \log v}{v-c_{1} \log v}
$$

Proof: We start with an upper bound for $P\left(\bar{\Gamma}_{v}^{1}\right)$. Define

$$
\begin{aligned}
& W_{1}:=\left(W\left(b_{v, 1}^{-}+x\right)-W\left(b_{v, 1}^{-}\right)\right)_{x \in\left[0, b_{v, 2}^{-}-b_{v, 1}^{-}\right]} \text {and } \\
& W_{2}:=\left(W\left(b_{v, 2}^{-}+x\right)-W\left(b_{v, 2}^{-}\right)\right)_{x \in\left[0, b_{v, 3}^{-}-b_{v, 2}^{-}\right]} .
\end{aligned}
$$

The event $\left\{c_{v}^{+}>b_{v, 3}^{-}\right\}$can then be rewritten

$$
\left\{\sup _{\left[0, b_{v, 2}^{-}-b_{v, 1}^{-}\right]} W_{1} \leq\left(c_{1}+c_{3}\right) \log v ; \sup _{\left[0, b_{v, 3}^{-}-b_{v, 2}^{-}\right]} W_{2} \leq\left(c_{1}+c_{3}\right) \log v\right\}
$$

The processes $W_{1}$ et $W_{2}$ are independent and have the same law as $\left(W(x), x \in\left[0, b_{v, 1}^{-}\right]\right)$, then

$$
\begin{aligned}
P\left(c_{v}^{+}>b_{v, 3}^{-}\right) & =\left(P\left(\sup _{\left[0, b_{v, 1}^{-}\right]} W \leq\left(c_{1}+c_{3}\right) \log v\right)\right)^{2} \\
& =\left(P\left(W \text { hits }-v+\left(2 c_{1}+c_{3}\right) \log v \text { before }\left(c_{1}+c_{3}\right) \log v\right)\right)^{2} \\
& =\left(\frac{\left(c_{1}+c_{3}\right) \log v}{v-c_{1} \log v}\right)^{2}
\end{aligned}
$$

For the second part of $\bar{\Gamma}_{v}^{1}$, according to Theorem 2.4 , if we denote by $R$ a Bessel process of dimension 3 strated at $v-c_{1} \log v$,

$$
\begin{aligned}
P\left(\underline{W}\left(b_{v}^{+}, c_{v}^{+}\right)-W\left(m_{v}^{+}\right)<\left(c_{1}\right.\right. & \left.\left.+c_{3}\right) \log v\right) \\
& =P\left(\underline{R}\left(0, \tau_{R}\left(v+c_{3} \log v\right)\right)<\left(c_{1}+c_{3}\right) \log v\right) \\
& =\left(\frac{\left(c_{1}+c_{3}\right) \log v}{v-c_{1} \log v}\right)^{2}
\end{aligned}
$$

(See for example Property 2.2.2 of part II, chap 5 in [2].) We obtain in the same way

$$
P\left(\bar{\Gamma}_{v}^{3}\right) \leq \frac{\left(c_{1}+c_{3}\right) \log v}{v-c_{1} \log v}
$$

We continue with an upper bound for $P\left(\bar{\Gamma}_{v}^{2}\right)$. As $W-\underline{W}$ has the same law as $|W|$,

$$
\begin{aligned}
P\left(b_{v, 2}^{-}>v^{6}\right) & \leq P\left(b_{v, 2}^{-}-b_{v, 1}^{-}>\frac{v^{6}}{2}\right)+P\left(b_{v, 1}^{-}>\frac{v^{6}}{2}\right) \\
& =2 P\left(\tau_{|W|}\left(v-c_{1} \log v\right)>\frac{v^{6}}{2}\right) \\
& \leq 2 P\left(\tau_{W}(v)>\frac{v^{6}}{2}\right) \leq \frac{K}{v^{2}}
\end{aligned}
$$

Moreover, $-W\left(m_{v, 1}^{-}\right)$and $W\left(b_{v, 1}^{-}\right)-W\left(m_{v, 2}^{-}\right)$are exponentially distributed with mean $v-c_{1} \log v$ (see for example the first lemma of 12]) and are independent. Thus

$$
\begin{aligned}
P\left(W\left(b_{v, 1}^{-}\right)-W\left(m_{v, 2}^{-}\right)<-v^{2}\right) & =P\left(W\left(m_{v, 1}^{-}\right)<-v^{2}\right) \\
& \leq P\left(W\left(m_{v, 1}^{-}\right)<-\left(v-c_{1} \log v\right)^{2}\right) \\
& =v^{c_{1}} e^{-v}
\end{aligned}
$$

According to Theorem 2.4, still denoting by $R$ a Bessel process of dimension 3 but now started at 0 ,

$$
\begin{aligned}
P\left(m_{v, i}^{-}-a_{v, i}^{-}<\frac{1}{v^{2}}\right) \leq & P\left(m_{v, i}^{-}-a_{v, i}^{-}<\frac{1}{\left(v-c_{1} \log v\right)^{2}}\right) \\
\leq & P\left(\tau_{R}\left(v-c_{1} \log v\right)<\frac{1}{\left(v-c_{1} \log v\right)^{2}}\right) \\
& +P\left(m_{v, i}^{-}<\frac{1}{\left(v-c_{1} \log v\right)^{2}}\right)
\end{aligned}
$$

Yet, according to the scaling property of Brownian motion and a lemma proved by Cheliotis in (4] (claim at the end of the proof of Lemma 13), there is a constant $K>0$, such that

$$
P\left(m_{v, i}^{-}<\frac{1}{\left(v-c_{1} \log v\right)^{2}}\right)=P\left(m_{1}<\frac{1}{\left(v-c_{1} \log v\right)^{4}}\right) \leq \frac{K}{\left(v-c_{1} \log v\right)^{2}}
$$

Moreover, according to Lemma 2.5,

$$
P\left(\tau_{R}\left(v-c_{1} \log v\right)<\frac{1}{\left(v-c_{1} \log v\right)^{2}}\right) \leq K\left(v-c_{1} \log v\right)^{2} e^{-\left(v-c_{1} \log v\right)^{4} / 2}
$$

We also obtain the following upper bound:

$$
\begin{aligned}
P\left(c_{v}^{+}-m_{v}^{+}<v\right) & =P\left(\tau_{R}\left(v+c_{3} \log v\right)<v\right) \\
& \leq P\left(\tau_{R}(v)<v\right) \\
& \leq K \sqrt{v} e^{-v / 2}
\end{aligned}
$$

It remains to control $P\left(W\left(c_{v}^{+}\right)-\underline{W}\left(\left(c_{v}^{+}-\log v\right) \vee m_{v}^{+}, c_{v}^{+}\right)>2 \log v\right)$. Let $\beta=v+c_{3} \log v$, using one more time Theorem 2.4, we see that

$$
\begin{aligned}
W\left(c_{v}^{+}\right) & -\underline{W}\left(\left(c_{v}^{+}-\log v\right) \vee m_{v}^{+}, c_{v}^{+}\right) \\
& =W\left(c_{v}^{+}\right)-W\left(m_{v}^{+}\right)-\left(\min _{t \in\left[0, \log v \wedge\left(c_{v}^{+}-m_{v}^{+}\right)\right]} W\left(\left(c_{v}^{+}-t\right)-W\left(m_{v}^{+}\right)\right)\right.
\end{aligned}
$$

has the same law as

$$
\beta-\min _{t \in\left[0, \log v \wedge \tau_{R}(\beta)\right]} R\left(\tau_{R}(\beta)-t\right)=\max _{t \in\left[0, \log v \wedge \tau_{R}(\beta)\right]}\left(\beta-R\left(\tau_{R}(\beta)-t\right)\right)
$$

And according to Proposition 4.8, Chapter VII of (13], the processes

$$
\left(\beta-R\left(\tau_{R}(\beta)-t\right) ; t \in\left[0, \tau_{R}(\beta)\right]\right) \text { and }\left(R(t) ; t \in\left[0, \tau_{R}(\beta)\right]\right)
$$

have the same law. Therefore

$$
\begin{aligned}
P\left(W\left(c_{v}^{+}\right)-\underline{W}\left(\left(c_{v}^{+}-\log v\right) \vee m_{v}^{+}\right.\right. & \left.\left.c_{v}^{+}\right)>2 \log v\right) \\
& =P\left(\max _{t \in\left[0, \log v \wedge \tau_{R}(\beta)\right]} R(t)>2 \log v\right) \\
& \leq P\left(\max _{t \in[0, \log v]} R(t)>2 \log v\right) \\
& \leq K \frac{\sqrt{\log v}}{v^{2}}
\end{aligned}
$$

where the last inequality comes from Lemma 2.5. Finally, we just have to obtain an upper bound for

$$
P\left(\bar{W}\left(\left(m_{v, 1}^{-}-\log v\right) \vee a_{v, 1}^{-}, m_{v, 1}^{-}\right)-W\left(m_{v, 1}^{-}\right)>2 \log v\right)
$$

to prove the lemma. This can be done in the same way as the previous one.

We can now proceed with the local time of the diffusion $\mathbf{X}$.

## 4 Asymptotic behaviour of $L$ at particular times

Let $r$ be a positive real number. As for the $c_{i}$, its value will be fixed later. We define $\boldsymbol{\sigma}_{v, 1}^{-}:=\boldsymbol{\sigma}\left(r e^{v}, m_{v, 1}^{-}\right)$, and in the same way $\boldsymbol{\sigma}_{v, 2}^{-}, \widehat{\boldsymbol{\sigma}}_{v, 1}^{-}, \widehat{\boldsymbol{\sigma}}_{v, 2}^{-}$and $\boldsymbol{\sigma}_{v}^{+}$. We show in this section that at these times, it is possible to estimate the local time of $\mathbf{X}$ in the neighborhood of the corresponding point. We first give an estimate at fixed environment in Proposition 4.1, then in Proposition 4.4, we give an estimate independent of the environment provided that this one belongs to $\Gamma_{v}$.

## Proposition 4.1 (Estimates depending on the environment)

We define for $i \in\{1,2\}$,

$$
\begin{aligned}
\mathcal{A}_{v}^{i} & :=\left\{\forall x \in\left[a_{v, i}^{-}, b_{v, i}^{-}\right]\left|\frac{\mathbf{L}\left(\boldsymbol{\sigma}_{v, i}^{-}, x\right)}{r e^{v-W(x)+W\left(m_{v, i}^{-}\right)}}-1\right| \leq \delta\right\}, \\
\mathcal{B}_{v}^{i} & :=\left\{\forall x \in\left[b_{v, i-1}^{-}, a_{v, i}^{-}\right], \mathbf{L}\left(\boldsymbol{\sigma}_{v, i}^{-}, x\right) \leq \delta r e^{v}\right\}, \\
\mathcal{C}_{v} & :=\left\{\forall x \in\left[b_{v}^{+}, c_{v}^{+}\right], \mathbf{L}\left(\boldsymbol{\sigma}_{v}^{+}, x\right) \leq \delta r e^{v}\right\} \text { and } \\
\mathcal{D}_{v} & :=\left\{\forall x>c_{v}^{+}, \mathbf{L}\left(\boldsymbol{\sigma}_{v}^{+}, x\right)=0\right\} .
\end{aligned}
$$

We also define $\widehat{\mathcal{A}_{v}^{i}}, \widehat{\mathcal{B}_{v}^{i}}, \widehat{\mathcal{C}_{v}}$ and $\widehat{\mathcal{D}_{v}}$ from $\widehat{W}$ in the same way. Then there is
a constant $K>0$ such that for any $v \geq 2$, any $0 \leq \delta \leq 1$ and any $r>0$,

$$
\begin{aligned}
P^{W}\left(\overline{\mathcal{A}_{v}^{i}}\right) & \leq \frac{K}{\delta} \sqrt{\frac{m_{v, i}^{-}}{r v^{c_{2}}}} \exp \left(-\frac{\delta^{2} r v^{c_{2}}}{K m_{v, i}^{-}}\right), \\
P^{W}\left(\overline{\mathcal{B}_{v}^{i}}\right) & \leq K \exp \left(-\frac{\delta r v^{c_{1}}}{2\left(m_{v, i}^{-}-b_{v, i-1}^{-}\right) \log \left(8 \frac{S\left(m_{v, i}^{-}\right)-S\left(b_{v, i-1}^{-}\right)}{\left(S\left(m_{v, i}^{-}\right)-S\left(a_{v, i}\right)\right)}\right)}\right)+\frac{1}{\delta v^{c_{1}-c_{2}}}, \\
P^{W}\left(\overline{\mathcal{C}_{v}}\right) & \leq \frac{1}{\delta} e^{-\underline{W}\left(b_{v}^{+}, c_{v}^{+}\right)+W\left(m_{v}^{+}\right)} \text {and } \\
P^{W}\left(\overline{\mathcal{D}_{v}}\right) & \leq \frac{r e^{v+W\left(m_{v}^{+}\right)}}{2\left(S\left(c_{v}^{+}\right)-S\left(m_{v}^{+}\right)\right)} .
\end{aligned}
$$

We have similar estimates for $P^{W}\left(\widehat{\mathcal{A}_{v}^{i}}\right), P^{W}\left(\widehat{\mathcal{B}_{v}^{i}}\right), P^{W}\left(\widehat{\mathcal{C}_{v}}\right)$ and $P^{W}\left(\widehat{\mathcal{D}_{v}}\right)$.

## Proof:

We estimates probabilities of the various events relative to $W$, we obtain results relative to $\widehat{W}$ by symmetry. To simplify the notations, all along the proof, we shall not mark the index $v$ in variables and events. Let us begin with events $\mathcal{A}^{1}$ and $\mathcal{B}^{1}\left(\mathcal{A}^{2}\right.$ and $\mathcal{B}^{2}$ can be studied in the same way $)$.

We decompose the local time in two terms, the first one representing the contribution of the local time before $\boldsymbol{\tau}\left(m_{1}^{-}\right)$(the first time where $\mathbf{X}$ reaches $m_{1}^{-}$) is negligible compared to the second one representing the contribution of local time between $\boldsymbol{\tau}\left(m_{1}^{-}\right)$and $\boldsymbol{\sigma}_{1}^{-}=\boldsymbol{\sigma}\left(r e^{v}, m_{1}^{-}\right)$(the inverse of local time in $m_{1}^{-}$):

$$
\mathbf{L}\left(\boldsymbol{\sigma}_{1}^{-}, x\right)=\mathbf{L}\left(\boldsymbol{\tau}\left(m_{1}^{-}\right), x\right)+\left(\mathbf{L}\left(\boldsymbol{\sigma}_{1}^{-}, x\right)-\mathbf{L}\left(\boldsymbol{\tau}\left(m_{1}^{-}\right), x\right)\right) .
$$

Let us begin with a lemma describing the behaviour of the first process.
Lemma 4.2 For any $v, \delta, r>0$,

$$
\begin{aligned}
P_{\mathcal{A}^{1}, 1}^{W} & :=P^{W}\left(\sup _{x \in\left[a_{1}^{-}, m_{1}^{-}\right]} \frac{\mathbf{L}\left(\boldsymbol{\tau}\left(m_{1}^{-}\right), x\right)}{r e^{v-\left(W(x)-W\left(m_{1}^{-}\right)\right)}}>\delta\right) \\
& \leq K \exp \left(-\frac{\delta v^{c}}{2 m_{1}^{-}}\right) \\
P_{\mathcal{B}^{1}, 1}^{W} & :=P^{W}\left(\sup _{x \in\left[b_{0}^{-}, a_{1}^{-}\right]} \mathbf{L}\left(\boldsymbol{\tau}\left(m_{1}^{-}\right), x\right)>\delta r e^{v}\right) \\
& \leq K \exp \left(-\frac{\delta r v^{c_{1}}}{2\left(m_{i}^{-}-b_{i-1}^{-}\right) \log \left(8 \frac{S\left(m_{i}^{-}\right)-S\left(b_{i-1}^{-}\right)}{\left(S\left(m_{i}^{-}\right)-S\left(a_{i}^{-}\right)\right)}\right)}\right) .
\end{aligned}
$$

Proof: Thanks to (3), we easily verify that $\boldsymbol{\tau}\left(m_{1}^{-}\right)=T\left(\tau\left(S\left(m_{1}^{-}\right)\right)\right.$. So with (4), we obtain for any $x \geq 0$,

$$
\mathbf{L}\left(\boldsymbol{\tau}\left(m_{1}^{-}\right), x\right)=e^{-W(x)} L\left(\tau\left(S\left(m_{1}^{-}\right)\right), S(x)\right)
$$

Thus,

$$
P_{\mathcal{A}^{1}, 1}^{W}=P^{W}\left(\sup _{x \in\left[a_{1}^{-}, m_{1}^{-}\right]} L\left(\tau\left(S\left(m_{1}^{-}\right)\right), S(x)\right)>\delta r e^{v+W\left(m_{1}^{-}\right)}\right)
$$

According to the first Ray-Knight theorem, Theorem 2.1,

$$
\left\{L\left(\tau\left(S\left(m_{1}^{-}\right)\right), S\left(m_{1}^{-}\right)-y\right), y \in\left[0, S\left(m_{1}^{-}\right)\right]\right\}
$$

has the same law as a square of a Bessel process of dimension 2 started at 0 . Therefore, with Item 3 of Lemma 2.3 ,

$$
P_{\mathcal{A}^{1}, 1}^{W} \leq K \exp \left(-\frac{\delta r e^{v+W\left(m_{1}^{-}\right)}}{2\left(S\left(m_{1}^{-}\right)-S\left(a_{1}^{-}\right)\right)}\right)
$$

and by definition of $a_{1}^{-}$,

$$
\begin{aligned}
S\left(m_{1}^{-}\right)-S\left(a_{1}^{-}\right) & =\int_{a_{1}^{-}}^{m_{1}^{-}} e^{W(x)} \mathrm{d} x \leq m_{1}^{-} e^{\bar{W}\left(a_{1}^{-}, m_{1}^{-}\right)} \\
& \leq m_{1}^{-} e^{v-c_{2} \log v+W\left(m_{1}^{-}\right)}
\end{aligned}
$$

Hence we obtain the initial upper bound of the lemma.
Let us continue with the second one, to highlight the fact that the computations are identical for $\mathcal{B}^{1}$ and for $\mathcal{B}^{2}$ we make the quantity $b_{0}^{-}$appears, although it is zero. Using a similar argument as above,

$$
\begin{aligned}
P_{\mathcal{B}^{1}, 1}^{W} & =P^{W}\left(\sup _{x \in\left[b_{0}^{-}, a_{1}^{-}\right]} e^{-W(x)} Q\left(S\left(m_{1}^{-}\right)-S(x)\right)>\delta r e^{v}\right) \\
& =P^{W}\left(\sup _{x \in\left[b_{0}^{-}, a_{1}^{-}\right]} \frac{e^{-W(x)}\left(S\left(m_{1}^{-}\right)-S(x)\right)}{S\left(m_{1}^{-}\right)-S(x)} Q\left(S\left(m_{1}^{-}\right)-S(x)\right)>\delta r e^{v}\right)
\end{aligned}
$$

By definition of $m_{1}^{-}, P$-a.s, for any $x \in\left[b_{0}^{-}, a_{1}^{-}\right]$,

$$
\begin{aligned}
e^{-W(x)}\left(S\left(m_{1}^{-}\right)-S(x)\right) & =e^{-W(x)} \int_{x}^{m_{1}^{-}} e^{W(y)} \mathrm{d} y \\
& \leq\left(m_{1}^{-}-b_{0}^{-}\right) e^{-W(x)+\bar{W}\left(x, m_{1}^{-}\right)} \\
& \leq\left(m_{1}^{-}-b_{0}^{-}\right) e^{v-c_{1} \log v}
\end{aligned}
$$

where the last inequality is due to $m_{1}^{-}$being the first positive real $x$ such that $W(x)-\underline{W}(0, x) \geq v-c_{1} \log v$. Thus, coming back to the probability $P_{\mathcal{B}^{1}, 1}^{W}$, we obtain

$$
P_{\mathcal{B}^{1}, 1}^{W} \leq P^{W}\left(\sup _{u \in\left[S\left(m_{1}^{-}\right)-S\left(a_{1}^{-}\right), S\left(m_{1}^{-}\right)-S\left(b_{0}^{-}\right)\right]} \frac{1}{u} Q(u)>\frac{\delta r v^{c_{1}}}{m_{1}^{-}-b_{0}^{-}}\right) .
$$

According to Item $\begin{aligned} & \text { of Lemma 2.3, we finally have }\end{aligned}$

$$
P_{\mathcal{B}^{1}, 1}^{W} \leq K \exp \left(-\frac{\delta r v^{c_{1}}}{2\left(m_{1}^{-}-b_{0}^{-}\right) \log \left(8 \frac{S\left(m_{1}^{-}\right)-S\left(b_{0}^{-}\right)}{S\left(m_{1}^{-}\right)-S\left(a_{1}^{-}\right)}\right)}\right) .
$$

This concludes the proof of the lemma.

Lemma 4.3 For any $v, \delta, r>0$,

$$
\begin{aligned}
P_{\mathcal{A}^{1}, 2}^{W} & :=P^{W}\left(\sup _{x \in\left[a_{1}^{-}, b_{1}^{-}\right]}\left|\frac{\mathbf{L}\left(\boldsymbol{\sigma}_{1}^{-}, x\right)-\mathbf{L}\left(\boldsymbol{\tau}\left(m_{1}^{-}\right), x\right)}{r e^{v-\left(W(x)-W\left(m_{1}^{-}\right)\right)}}-1\right|>\delta\right) \\
& \leq \frac{8}{\delta} \sqrt{\frac{(1+\delta) m_{1}^{-}}{r v^{c_{2}}}} \exp \left(-\frac{\delta^{2} r v^{c_{2}}}{8(1+\delta) m_{1}^{-}}\right), \\
P_{\mathcal{B}^{1}, 2}^{W} & :=P^{W}\left(\sup _{x \in\left[b_{0}^{-}, a_{1}^{-}\right]} \mathbf{L}\left(\boldsymbol{\sigma}_{1}^{-}, x\right)-\mathbf{L}\left(\boldsymbol{\tau}\left(m_{1}^{-}\right), x\right)>\delta r e^{v}\right) \\
& \leq \frac{1}{\delta v^{c_{1}-c_{2}}} .
\end{aligned}
$$

Proof: It is easy to verify that the inverse of local time $\boldsymbol{\sigma}$ satisfies the following equality for every $r>0$ and $y \in \mathbf{R}$,

$$
\begin{equation*}
\boldsymbol{\sigma}(r, y)=T\left(\sigma\left(r e^{W(y)}, S(y)\right)\right) \tag{7}
\end{equation*}
$$

Thus, ( (1) implies that, for any $r>0$ and $y \in \mathbf{R}$,

$$
\mathbf{L}\left(\boldsymbol{\sigma}_{1}^{-}, x\right)=e^{-W(x)} L\left(\sigma\left(r e^{W\left(m_{1}^{-}\right)+v}, S\left(m_{1}^{-}\right)\right), S(x)\right) .
$$

And we obtain the following expression for the local time,

$$
\begin{aligned}
\mathbf{L}\left(\boldsymbol{\sigma}_{1}^{-}, x\right) & -\mathbf{L}\left(\boldsymbol{\tau}\left(m_{1}^{-}\right), x\right) \\
& =e^{-W(x)}\left(L\left(\sigma\left(r e^{W\left(m_{1}^{-}\right)+v}, S\left(m_{1}^{-}\right)\right), S(x)\right)-L\left(\tau\left(S\left(m_{1}^{-}\right)\right), S(x)\right)\right) \\
& \stackrel{\mathcal{L}}{=} e^{-W(x)} r e^{W\left(m_{1}^{-}\right)+v} L(\sigma(1,0), s(x))
\end{aligned}
$$

where

$$
s(x):=\left(S(x)-S\left(m_{1}^{-}\right)\right) \frac{e^{-W\left(m_{1}^{-}\right)-v}}{r} .
$$

Denoting by $Z$ a square of a 2 -dimensional Bessel process started at 1 , according to the second Ray-Knight theorem (Theorem 2.2), we have

$$
\begin{aligned}
P_{\mathcal{A}^{1}, 2}^{W} & \leq P^{W}\left(\sup _{0 \leq y \leq\left|s\left(a_{1}^{-}\right)\right|}|Z(y)-1|>\delta\right) \\
& +P^{W}\left(\sup _{0 \leq y \leq s\left(b_{1}^{-}\right)}|Z(y)-1|>\delta\right) .
\end{aligned}
$$

Therefore, using Item 1 of Lemma 2.3 we conclude that

$$
\begin{aligned}
P_{\mathcal{A}^{1}, 2}^{W} & \leq \frac{4}{\delta} \sqrt{(1+\delta)\left|s\left(a_{1}^{-}\right)\right|} \exp \left(-\frac{\delta^{2}}{8(1+\delta)\left|s\left(a_{1}^{-}\right)\right|}\right) \\
& +\frac{4}{\delta} \sqrt{(1+\delta) s\left(b_{1}^{-}\right)} \exp \left(-\frac{\delta^{2}}{8(1+\delta) s\left(b_{1}^{-}\right)}\right)
\end{aligned}
$$

Moreover, the definition of $a_{1}^{-}$implies

$$
\begin{aligned}
\left|s\left(a_{1}^{-}\right)\right| & =\frac{e^{-W\left(m_{1}^{-}\right)-v}}{r} \int_{a_{1}^{-}}^{m_{1}^{-}} e^{W(x)} \mathrm{d} x \\
& \leq m_{1}^{-} \frac{e^{W\left(a_{1}^{-}\right)-W\left(m_{1}^{-}\right)-v}}{r} \leq \frac{m_{1}^{-}}{r v^{c}} .
\end{aligned}
$$

This last upper bound and the similar one for $s\left(b_{1}^{-}\right)$lead us to the first inequality of the lemma. Now we continue with the proof of the second inequality. Still reasoning in the same way, we can have

$$
P_{\mathcal{B}^{1}, 2}^{W}=P^{W}\left(\sup _{x \in\left[b_{0}^{-}, a_{1}^{-}\right]} e^{-W(x)+W\left(m_{1}^{-}\right)} L(\sigma(1,0), s(x))>\delta\right) .
$$

One more time, thanks to the second Ray-Knight theorem and denoting by $Z$ a Bessel process of dimension 0 started at 1 , we obtain:

$$
\begin{aligned}
P_{\mathcal{B}^{1}, 2}^{W} & \leq P^{W}\left(e^{-\underline{W}\left(b_{0}^{-}, a_{1}^{-}\right)+W\left(m_{1}^{-}\right)} \sup _{u \geq 0} Z(u)>\delta\right) \\
& =\frac{1}{\delta} e^{-\underline{W}\left(b_{0}^{-}, a_{1}^{-}\right)+W\left(m_{1}^{-}\right)} .
\end{aligned}
$$

The second line is a consequence of Item 2 of Lemma 2.3. Let us denote by $n^{1}$ the unique real number in $\left[b_{0}^{-}, a_{1}^{-}\right]$such that $W\left(n^{1}\right)=\underline{W}\left(b_{0}^{-}, a_{1}^{-}\right)$then

$$
\begin{aligned}
W\left(m_{1}^{-}\right)-W\left(n^{1}\right) & =\bar{W}\left(n^{1}, a_{1}^{-}\right)-W\left(n^{1}\right)-\left(\bar{W}\left(n^{1}, a_{1}^{-}\right)-W\left(m_{1}^{-}\right)\right) \\
& \leq v-c_{1} \log v-\left(v-c_{2} \log v\right)=\left(c_{2}-c_{1}\right) \log v .
\end{aligned}
$$

Therefore,

$$
P_{\mathcal{B}^{1}, 2}^{W} \leq \frac{1}{\delta v^{c_{1}-c_{2}}}
$$

and the lemma is proved.
Combining the results of Lemmas 4.2 and 4.3 yields to the upper bounds of $P^{W}\left(\overline{\mathcal{A}^{1}}\right)$ and $P^{W}\left(\overline{\mathcal{B}^{1}}\right)$ of Proposition 4.1.

We now continue with the estimate of $P^{W}(\mathcal{C})$. Reducing once again process $\mathbf{X}$ to a Brownian motion by a time and space change, we obtain

$$
P^{W}(\overline{\mathcal{C}})=P^{W}\left(\sup _{x \in\left[b^{+}, c^{+}\right]} e^{-W(x)+W\left(m^{+}\right)} L(\sigma(1,0), s(x))>\delta\right) .
$$

As before, if $Z$ denotes a squared 0-dimensional Bessel process starting at 1, the second Ray-Knight theorem gives:

$$
\begin{aligned}
P^{W}(\overline{\mathcal{C}}) & \leq P^{W}\left(e^{-\underline{W}\left(b^{+}, c^{+}\right)+W\left(m^{+}\right)} \sup _{u \geq 0} Z(u)>\delta\right) \\
& \leq \frac{1}{\delta} e^{-\underline{W}\left(b^{+}, c^{+}\right)+W\left(m^{+}\right)}
\end{aligned}
$$

where for the second inequality, we used Item 2 of Lemma 2.3. Finally we show that with high probability at time $\boldsymbol{\sigma}_{1}^{-}$, diffusion $\mathbf{X}$ does not hit $c^{+}$, that is to say, we find an upper bound for $P^{W}(\overline{\mathcal{D}})$. The scale change in time and space and the usual properties of Brownian motion give us (see eg [2] Formula 4.1.2 page 185)

$$
\begin{aligned}
P^{W}(\overline{\mathcal{D}}) & =P^{W}\left(\boldsymbol{\tau}\left(c^{+}\right)<\boldsymbol{\sigma}\left(r e^{v}, m^{+}\right)\right) \\
& =P^{W}\left(\tau\left(S\left(c^{+}\right)\right)<\sigma\left(r e^{v+W\left(m^{+}\right)}, S\left(m^{+}\right)\right)\right) \\
& =1-\exp \left(-\frac{r e^{v+W\left(m^{+}\right)}}{2\left(S\left(c^{+}\right)-S\left(m^{+}\right)\right)}\right) \leq \frac{r e^{v+W\left(m^{+}\right)}}{2\left(S\left(c^{+}\right)-S\left(m^{+}\right)\right)} .
\end{aligned}
$$

This completes the proof of the proposition.
We now give upper bounds independent of the medium while it is in the event $\Gamma_{v}$ defined in 5 .

Proposition 4.4 We use the same notations as in the previous proposition. There is a constant $K>0$ such that for any $v \geq 2$, any $0 \leq \delta \leq 1$, and any
$r>0$, if $W \in \Gamma_{v}$,

$$
\begin{aligned}
P^{W}\left(\overline{\mathcal{A}_{v}^{i}}\right) & \leq \frac{K}{\delta \sqrt{r v^{c_{2}-6}}} \exp \left(-\frac{\delta^{2} r v^{c_{2}-6}}{K}\right) \\
P^{W}\left(\overline{\mathcal{B}_{v}^{i}}\right) & \leq K \exp \left(-\frac{\delta r v^{c_{1}-8}}{K}\right)+\frac{1}{\delta v^{c_{1}-c_{2}}} \\
P^{W}\left(\overline{\mathcal{C}_{v}}\right) & \leq \frac{1}{\delta v^{c_{1}+c_{3}}} \text { and } \\
P^{W}\left(\overline{\mathcal{D}_{v}}\right) & \leq \frac{r}{v^{c_{3}-2} \log v}
\end{aligned}
$$

We have again similar estimates for $P^{W}\left(\widehat{\mathcal{A}_{v}^{i}}\right), P^{W}\left(\widehat{\mathcal{B}_{v}^{i}}\right), P^{W}\left(\widehat{\mathcal{C}_{v}}\right), P^{W}\left(\widehat{\mathcal{D}_{v}}\right)$.
Proof: We only have to control the values of the upper bounds of Proposition 4.1 when $W \in \Gamma_{v}$.

For $P^{W}\left(\overline{\mathcal{A}_{v}^{i}}\right)$, it is enough to notice that $m_{v, i}^{-} \leq b_{v, 2}^{-}$and that, on $\Gamma_{v}$, the variable $b_{v, 2}^{-}$is smaller than $v^{6}$. This also allows us to obtain the upper bound $m_{v, i}^{-}-b_{v, i-1}^{-} \leq b_{v, 2}^{-} \leq v^{6}$. To estimate $P^{W}\left(\overline{\mathcal{B}_{v}^{i}}\right)$, it remains to study

$$
\frac{S\left(m_{v, i}^{-}\right)-S\left(b_{v, i-1}^{-}\right)}{S\left(m_{v, i}^{-}\right)-S\left(a_{v, i}^{-}\right)} \leq \frac{\left(m_{v, i}^{-}-b_{v, i-1}^{-}\right) e^{\bar{W}\left(b_{v, i-1}^{-}, m_{v, i}^{-}\right)}}{\left(m_{v, i}^{-}-a_{v, i}^{-}\right) e^{W\left(m_{v, i}^{-}\right)}}
$$

But, on one hand $\bar{W}\left(b_{v, i-1}^{-}, m_{v, i}^{-}\right)-W\left(b_{v, i-1}^{-}\right) \leq v-c_{1} \log v \leq v$ and on the other hand, as we work on $\Gamma_{v}$, it is easy to find upper bound for the other quantities involved in the computation. And we finally obtain,

$$
\frac{S\left(m_{v, i}^{-}\right)-S\left(b_{v, i-1}^{-}\right)}{S\left(m_{v, i}^{-}\right)-S\left(a_{v, i}^{-}\right)} \leq v^{8} e^{v+v^{2}}
$$

which implies the result of the proposition.
As on $\Gamma_{v}$, we have $\underline{W}\left(b_{v}^{+}, c_{v}^{+}\right)-W\left(m_{v}^{+}\right) \geq\left(c_{1}+c_{3}\right) \log v$, the third estimate is obtained immediately. It remains only the upper bound of $P^{W}\left(\overline{\mathcal{D}_{v}}\right)$. For this, it is enough to found a lower bound to

$$
\begin{aligned}
S\left(c_{v}^{+}\right)-S\left(m_{v}^{+}\right) & =\int_{m_{v}^{+}}^{c_{v}^{+}} e^{W(x)} \mathrm{d} x \geq \int_{\left(c_{v}^{+}-\log v\right) \vee m_{v}^{+}}^{c_{v}^{+}} e^{W(x)} \mathrm{d} x \\
& \geq\left(\log v \wedge\left(c_{v}^{+}-m_{v}^{+}\right)\right) e^{\underline{W}\left(\left(c_{v}^{+}-\log v\right) \vee m_{v}^{+}, c_{v}^{+}\right)}
\end{aligned}
$$

Thus, as on $\Gamma_{v}$ we have $c_{v}^{+}-m_{v}^{+} \geq v$ and $\underline{W}\left(\left(c_{v}^{+}-\log v\right) \vee m^{+}, c_{v}^{+}\right) \geq$ $W\left(c_{v}^{+}\right)-2 \log v$, the following is true:

$$
P^{W}\left(\overline{\mathcal{D}_{v}}\right) \leq \frac{r v^{2}}{\log v} e^{v+W\left(m_{v}^{+}\right)-W\left(c_{v}^{+}\right)}
$$

Moreover $W\left(c_{v}^{+}\right)-W\left(m_{v}^{+}\right)=v+c_{3} \log v$, we therefore obtain the last result of the proposition.

## 5 Asymptotics of the maximum local time

Thanks to Proposition 4.4, we can now study the process $\mathbf{L}$ at the time $\boldsymbol{\sigma}_{v}:=\widehat{\boldsymbol{\sigma}}_{v} \wedge \widehat{\boldsymbol{\sigma}}_{v, 1}^{-} \wedge \boldsymbol{\sigma}_{v, 1}^{-} \wedge \boldsymbol{\sigma}_{v}^{+}$the first instant the diffusion has "passed a time" $r e^{v}$ in one of the points $m_{v, 1}^{-}, m_{v}^{+}, \widehat{m}_{v, 1}^{-}$or $\widehat{m}_{v}^{+}$. Let us begin by fixing the values of constants $c_{i}$ : we take a real number $c>0$, then we choose $c_{1}, c_{2}$ and $c_{3}$ as follows $c_{1}:=2 c+8, c_{2}:=c+6$ and $c_{3}:=c+2$. This yields to the following proposition:

Proposition 5.1 Let $v \geq 2$. We define

$$
I_{v}^{-}:=\int_{a_{v, 1}^{-}}^{b_{v, 1}^{-}} e^{-W(x)+W\left(m_{v, 1}^{-}\right)} \mathrm{d} x, I_{v}^{+}:=\int_{a_{v}^{+}}^{b_{v}^{+}} e^{-W(x)+W\left(m_{v}^{+}\right)} \mathrm{d} x
$$

We define similarly $\widehat{I}_{v}^{-}$and $\widehat{I}_{v}^{+}$from $\widehat{W}$ and

$$
j_{v}:=\widehat{I}_{v}^{-} \wedge \widehat{I}_{v}^{+} \wedge I_{v}^{-} \wedge I_{v}^{+} \text {et } J_{v}:=\widehat{I}_{v}^{-}+\widehat{I}_{v}^{+}+I_{v}^{-}+I_{v}^{+}
$$

Then there is a constant $K>0$ such that for any $0<\delta<1$ and any $r>0$, if $W \in \Gamma_{v}$,

$$
\begin{aligned}
P^{W} & \left(r e^{v} \leq \mathbf{L}^{*}\left(\boldsymbol{\sigma}_{v}\right) \leq r e^{v}(1+\delta) ; j_{v}(1-\delta) \leq \frac{\boldsymbol{\sigma}_{v}}{r e^{v}} \leq J_{v}+2 v^{6} \delta\right) \\
& \geq 1-K\left(\frac{1}{\delta \sqrt{r v^{c}}} \exp \left(-\frac{\delta^{2} r v^{c}}{K}\right)+\exp \left(-\frac{\delta r v^{c}}{K}\right)+\frac{1}{\delta v^{c}}+\frac{r}{v^{c}}\right) .
\end{aligned}
$$

Note that in the previous section we used the four points $m_{v, 1}^{-}, m_{v, 2}^{-}, \widehat{m}_{v, 1}^{-}$ and $\widehat{m}_{v, 2}^{-}$whereas in the last proposition we use the points $m_{v, 1}^{-}, m_{v}^{+}, \widehat{m}_{v, 1}^{-}$ and $\widehat{m}_{v}^{+}$. The former ones were interesting because they simplify the computations of Proposition 4.4, but the latter ones allow us to reduce the study of $I_{v}^{-}, I_{v}^{+}, \widehat{I}_{v}^{-}$and $\widehat{I}_{v}^{+}$to the study of the integral

$$
I_{v}=\int_{a_{v}}^{b_{v}} e^{-W(x)+W\left(m_{v}\right)} \mathrm{d} x
$$

where

$$
\begin{aligned}
& b_{v}:=\inf \left\{x \geq m_{v} ; W(x)-W\left(m_{v}\right) \geq v-c_{1} \log v\right\} \text { and } \\
& a_{v}:=\sup \left\{x \leq m_{v} ; W(x)-W\left(m_{v}\right) \geq v-c_{2} \log v\right\} \vee 0 .
\end{aligned}
$$

Proof: We only have to prove that on the intersection of all the events of Proposition 4.4 the event

$$
\left\{r e^{v} \leq \mathbf{L}^{*}\left(\boldsymbol{\sigma}_{v}\right) \leq r e^{v}(1+\delta) ; j_{v}(1-\delta) \leq \frac{\boldsymbol{\sigma}_{v}}{r e^{v}} \leq J_{v}+2 v^{6} \delta\right\}
$$

is realized. As $\boldsymbol{\sigma}_{v}$ is the first time the diffusion has "passed a time" $r e^{v}$ in one of the points $m_{v, 1}^{-}, m_{v}^{+}, \widehat{m}_{v, 1}^{-}$or $\widehat{m}_{v}^{+}$we already have $r e^{v} \leq \mathbf{L}^{*}\left(\boldsymbol{\sigma}_{v}\right)$.

Moreover, the local time is non-decreasing, so for every $x \in \mathbf{R}$,

$$
\mathbf{L}\left(\boldsymbol{\sigma}_{v}, x\right)=\mathbf{L}\left(\widehat{\boldsymbol{\sigma}}_{v}, x\right) \wedge \mathbf{L}\left(\widehat{\boldsymbol{\sigma}}_{v, 1}^{-}, x\right) \wedge \mathbf{L}\left(\boldsymbol{\sigma}_{v, 1}^{-}, x\right) \wedge \mathbf{L}\left(\boldsymbol{\sigma}_{v}^{+}, x\right)
$$

On $\Gamma_{v}$ we have $\boldsymbol{\sigma}_{v}^{+} \in\left\{\boldsymbol{\sigma}_{v, 1}^{-}, \boldsymbol{\sigma}_{v, 2}^{-}\right\}$and $\widehat{\boldsymbol{\sigma}}_{v} \in\left\{\widehat{\boldsymbol{\sigma}}_{v, 1}^{-}, \widehat{\boldsymbol{\sigma}}_{v, 2}^{-}\right\}$. We then obtain the upper bound $\mathbf{L}^{*}\left(\boldsymbol{\sigma}_{v}\right) \leq r e^{v}(1+\delta)$.

Let us continue with the estimate of $\boldsymbol{\sigma}_{v}$ : by definition of the local time,

$$
\boldsymbol{\sigma}_{v}=\int_{-\infty}^{+\infty} \mathbf{L}\left(\boldsymbol{\sigma}_{v}, x\right) \mathrm{d} x, P \text {-a.s. }
$$

If $m_{v}^{+}=m_{v, 1}^{-}$, then $b_{v}^{+}=b_{v, 1}^{-}$and $\boldsymbol{\sigma}_{v}^{+}=\boldsymbol{\sigma}_{v, 1}^{-}$, therefore

$$
\begin{aligned}
& \int_{0}^{+\infty} \mathbf{L}\left(\boldsymbol{\sigma}_{v}, x\right) \mathrm{d} x \\
& \quad \leq \int_{0}^{a_{v, 1}^{-}} \mathbf{L}\left(\boldsymbol{\sigma}_{v, 1}^{-}, x\right) \mathrm{d} x+\int_{a_{v, 1}^{-}}^{b_{v, 1}^{-}} \mathbf{L}\left(\boldsymbol{\sigma}_{v, 1}^{-}, x\right) \mathrm{d} x+\int_{b_{v}^{+}}^{c_{v}^{+}} \mathbf{L}\left(\boldsymbol{\sigma}_{v}^{+}, x\right) \mathrm{d} x \\
& \quad \leq\left(I_{v}^{-}+\delta c_{v}^{+}\right) r e^{v}
\end{aligned}
$$

else if $m_{v}^{+}=m_{v, 2}^{-}$,

$$
\begin{aligned}
& \int_{0}^{+\infty} \mathbf{L}\left(\boldsymbol{\sigma}_{v}, x\right) \mathrm{d} x \leq \int_{0}^{a_{v, 1}^{-}} \mathbf{L}\left(\boldsymbol{\sigma}_{v, 1}^{-}, x\right) \mathrm{d} x+\int_{a_{v, 1}^{-}}^{b_{v, 1}^{-}} \mathbf{L}\left(\boldsymbol{\sigma}_{v, 1}^{-}, x\right) \mathrm{d} x \\
&+\int_{b_{v, 1}^{-}}^{a_{v, 2}^{-}} \mathbf{L}\left(\boldsymbol{\sigma}_{v, 2}^{-}, x\right) \mathrm{d} x+\int_{a_{v, 2}^{-}}^{b_{v, 2}^{-}} \mathbf{L}\left(\boldsymbol{\sigma}_{v, 2}^{-}, x\right) \mathrm{d} x+\int_{b_{v}^{+}}^{c_{v}^{+}} \mathbf{L}\left(\boldsymbol{\sigma}_{v}^{+}, x\right) \mathrm{d} x \\
& \leq\left(I_{v}^{-}+I_{v}^{+}+\delta c_{v}^{+}\right) r e^{v}
\end{aligned}
$$

The integral $\int_{-\infty}^{0} \mathbf{L}\left(\boldsymbol{\sigma}_{v}, x\right) \mathrm{d} x$ has a similar upper bound. As on $\Gamma_{v}$ we have $c_{v}^{+}+\widehat{c}_{v}^{+} \leq 2 v^{6}$, the upper bound of the proposition follows immediately.

For the lower bound, if $\boldsymbol{\sigma}_{v}=\boldsymbol{\sigma}_{v, 1}^{-}$, we have

$$
\begin{aligned}
\frac{\boldsymbol{\sigma}_{v, 1}^{-}}{r e^{v}} & \geq \int_{a_{v, 1}^{-}}^{b_{v, 1}^{-}} \mathbf{L}\left(\boldsymbol{\sigma}_{v, 1}^{-}, x\right) \mathrm{d} x \\
& \geq I_{v}^{-}(1-\delta)
\end{aligned}
$$

Doing the same computation if $\boldsymbol{\sigma}_{v}$ takes one of the three other possible values, the lower bound stated in the proposition is obtained immediately.

We are now interested in the behaviour of the local time in deterministic time:

Proposition 5.2 Let $v \geq 2$. For any $0<\delta \leq 1 / 2$, if $W \in \Gamma_{v}$,

$$
\begin{aligned}
& P^{W}\left(\frac{e^{v}}{J_{v}+2 v^{6} \delta} \leq \mathbf{L}^{*}\left(e^{v}\right) \leq \frac{e^{v}(1+\delta)}{j_{v}(1-\delta)}\right) \\
& \quad \geq 1-K\left(\frac{1}{\delta \sqrt{v^{c-6}}} \exp \left(-\frac{\delta^{2} v^{c-6}}{K}\right)+\exp \left(-\frac{\delta v^{c-6}}{K}\right)+\frac{1}{\delta v^{c}}+\frac{1}{v^{c-4}}\right) .
\end{aligned}
$$

Proof: We now use the real number $r$ which appears in all propositions since the beginning. Furthermore we use the result of the previous proposition. Let us define

$$
\rho(v):=\frac{1}{J_{v}+2 v^{6} \delta} \text { and } r(v):=\frac{1}{j_{v}(1-\delta)} .
$$

We also write $\boldsymbol{\sigma}_{v}^{\rho}$ for the time $\boldsymbol{\sigma}_{v}$ associated with $\rho$ and $\boldsymbol{\sigma}_{v}^{r}$ for the one associated with $r$. We now consider the events

$$
\begin{aligned}
\Omega^{\rho} & :=\left\{\frac{\boldsymbol{\sigma}_{v}^{\rho}}{\rho(v) e^{v}} \leq J_{v}+v^{3} \delta\right\}=\left\{\boldsymbol{\sigma}_{v}^{\rho} \leq e^{v}\right\}, \\
\Omega^{r} & :=\left\{\mathbf{L}^{*}\left(\boldsymbol{\sigma}_{v}^{r}\right) \leq r(v) e^{v}(1+\delta) ; j_{v}(1-\delta) \leq \frac{\boldsymbol{\sigma}_{v}^{r}}{r(v) e^{v}}\right\} \\
& =\left\{\mathbf{L}^{*}\left(\boldsymbol{\sigma}_{v}^{r}\right) \leq r(v) e^{v}(1+\delta) ; e^{v} \leq \boldsymbol{\sigma}_{v}^{r}\right\}
\end{aligned}
$$

As the maximum of the local time is a non decreasing function,

$$
\begin{aligned}
\text { on } \Omega^{r}, \mathbf{L}^{*}\left(e^{v}\right) & \leq \mathbf{L}^{*}\left(\boldsymbol{\sigma}_{v}^{r}\right) \leq r(v) e^{v}(1+\delta) \\
\text { and on } \Omega^{\rho}, \rho(v) e^{v} & \leq \mathbf{L}^{*}\left(\boldsymbol{\sigma}_{v}^{\rho}\right) \leq \mathbf{L}^{*}\left(e^{v}\right) .
\end{aligned}
$$

Therefore it is enough to find a lower bound of $P\left(\Omega^{r} \cap \Omega^{\rho}\right)$. According to the previous proposition, we only have to estimate $r$ and $\rho$. On $\Gamma_{v}$ it is true that

$$
j_{v} \leq J_{v} \leq c_{v}^{+}-\widehat{c}_{v}^{+} \leq 2 v^{6} .
$$

Moreover we have $m_{v, 1}^{-}-a_{v, 1}^{-} \geq \frac{1}{v^{2}}$ and

$$
\bar{W}\left(m_{v, 1}^{-}-\log v \vee a_{v, 1}^{-}, m_{v, 1}^{-}\right)-W\left(m_{v, 1}^{-}\right) \leq 2 \log v,
$$

thus

$$
\begin{aligned}
I_{v}^{-} & =\int_{a_{v, 1}^{-}}^{b_{v, 1}^{-}} e^{-W(x)+W\left(m_{v, 1}^{-}\right)} \mathrm{d} x \\
& \geq\left(\left(m_{v, 1}^{-}-a_{v, 1}^{-}\right) \wedge(\log v)\right) e^{-\bar{W}\left(m_{v, 1}^{-}-\log v \vee a_{v, 1}^{-}, m_{v, 1}^{-}\right)+W\left(m_{v, 1}^{-}\right)} \\
& \geq \frac{1}{v^{4}} .
\end{aligned}
$$

The lower bounds for $I_{v}^{+}, \widehat{I}_{v}^{-}$and $\widehat{I}_{v}^{+}$are found in the same way. Finally,

$$
\frac{K}{v^{6}} \leq \rho(v) \leq r(v) \leq 2 v^{4}
$$

and it is now easy to obtain the estimate of the proposition.
Using the result of the previous proposition with $\delta=v^{-7}$ and $c>20$ and the upper bound for $P\left(\Gamma_{v}\right)$ of Proposition 3.1, this yields for any $v \geq 2$ to

$$
\begin{equation*}
P\left(\frac{e^{v}}{J_{v}+2 v^{-1}} \leq \mathbf{L}^{*}\left(e^{v}\right) \leq \frac{e^{v}\left(1+v^{-1}\right)}{j_{v}\left(1-v^{-1}\right)}\right) \geq 1-K\left(\frac{\log v}{v-c_{1} \log v}\right)^{2} \tag{8}
\end{equation*}
$$

We also obtain, using this time the event $\Gamma_{v}^{\prime}$,

$$
\begin{equation*}
P\left(\frac{e^{v}}{I_{v}+\widehat{I}_{v}+2 v^{-1}} \leq \mathbf{L}^{*}\left(e^{v}\right) \leq \frac{e^{v}\left(1+v^{-1}\right)}{I_{v} \wedge \widehat{I}_{v}\left(1-v^{-1}\right)}\right) \geq 1-\frac{K \log v}{v-c_{1} \log v} \tag{9}
\end{equation*}
$$

## 6 Proof of Theorem (1.1)

As shown by Proposition 5.2, the asymptotic behaviour of $\mathbf{L}^{*}$ has a direct link with the ones of $j_{v}, J_{v}$ and more generally with the behaviour of integral $I_{v}=\int_{a_{v}}^{b_{v}} e^{-W(x)+W\left(m_{v}\right)} \mathrm{d} x$. So to prove the theorem we need to study this integral.

### 6.1 Maximum and minimum speed of process $I_{v}$

Let us begin with a lower bound for the maximum speed:
Lemma 6.1 Let $v_{n}=e^{n}$. Then P-a.s.,

$$
\lim \sup _{n \rightarrow \infty} \frac{I_{v_{n}} \wedge \widehat{I}_{v_{n}}}{\log _{2} v_{n}} \geq \frac{1}{e \pi^{2}}
$$

Proof: We first define the sequence of events

$$
E_{n}:=\left\{m_{v_{n}}>m_{v_{n-1}} ; \int_{m_{v_{n}}}^{b_{v_{n}}} e^{-W(x)+W\left(m_{v_{n}}\right)} \mathrm{d} x \geq \frac{\log n}{e \pi^{2}}\right\}
$$

Denote by $\left(\mathcal{F}_{n}\right)_{n \in \mathbf{N}}$ the filtration generated by $\left(W(t) ; 0 \leq t \leq b_{v_{n}}\right)$, we also put $W_{n}=W\left(\cdot+b_{v_{n}}\right)-W\left(b_{v_{n}}\right)$. The events

$$
\left\{m_{v_{n}}>m_{v_{n-1}}\right\}=\left\{W_{n-1} \text { hits }-v_{n-1} \text { before } v_{n}-v_{n-1}\right\}
$$

and

$$
\left\{\int_{m_{v_{n}}}^{b_{v_{n}}} e^{-W(x)+W\left(m_{v_{n}}\right)} \mathrm{d} x \geq \frac{\log n}{e \pi^{2}}\right\}
$$

are independent of one another and of $\mathcal{F}_{n-1}$. Moreover,

$$
P\left(W_{n} \text { hits }-v_{n-1} \text { before } v_{n}-v_{n-1}\right)=\frac{v_{n}-v_{n-1}}{v_{n}}=\left(1-e^{-1}\right)
$$

and according to Theorem 2.4,

$$
\begin{aligned}
P\left(\int_{m_{v_{n}}}^{b_{v_{n}}} e^{-W(x)+W\left(m_{v_{n}}\right)} \mathrm{d} x \geq \frac{\log n}{e \pi^{2}}\right) & \geq P\left(\int_{0}^{T_{R}(1)} e^{-R(x)} \mathrm{d} x \geq \frac{\log n}{e \pi^{2}}\right) \\
& \geq P\left(T_{R}(1) \geq \frac{\log n}{\pi^{2}}\right) .
\end{aligned}
$$

According to Theorem 2 of [22] (recall that $j_{1 / 2,0}=\pi$.) there is a real number $K>0$ such that

$$
P\left(T_{R}(1) \geq x\right) \sim \frac{1}{K} e^{-\pi^{2} \frac{x}{2}} .
$$

Then we have, (the value of $K$ can change)

$$
P\left(T_{R}(1) \geq x\right) \geq \frac{1}{K} e^{-\pi^{2} \frac{x}{2}}
$$

and so

$$
P\left(E_{n}\right) \geq \frac{1-e^{-1}}{K \sqrt{n}}
$$

We now define the similar event for $\widehat{W}$ :

$$
\widehat{E}_{n}:=\left\{\widehat{m}_{v_{n}}>\widehat{m}_{v_{n-1}} ; \int_{\widehat{m}_{v_{n}}}^{\widehat{b}_{v_{n}}} e^{-W(x)+W\left(\widehat{m}_{v_{n}}\right)} \mathrm{d} x \geq \frac{\log n}{e \pi^{2}}\right\} .
$$

The events $E_{n}, \widehat{E}_{n}$ are pairwise independent, thus

$$
P\left(E_{n} \cap \widehat{E}_{n}\right)=P\left(E_{n}\right) P\left(\widehat{E}_{n}\right) \geq \frac{\left(1-e^{-1}\right)^{2}}{K^{2} n} .
$$

We can now conclude thanks to the second Borel-Cantelli lemma.
We are not interested in an upper bound of the minimum speed because this would lead, except for the value of the constant, to the result obtained by Shi in (15]. We now search almost sure bounds of the integral. To this end,
we start first by studying the successive values $\mu_{n}$ the process $\left\{m_{v}, v \geq 1\right\}$ can take. This is precisely defined as follows:

$$
\begin{aligned}
\beta_{1} & :=\inf \{x \geq 0 ; W(x)-\underline{W}(0, x)=1\} \\
\mu_{1} & :=\inf \left\{x \geq 0 ; W(x)=\underline{W}\left(0, \beta_{1}\right)\right\} \\
\gamma_{1} & :=\inf \left\{x \geq \beta_{1} ; W(x)=W\left(\mu_{1}\right)\right\} \\
M_{1} & :=\inf \left\{x \geq \beta_{1} ; W(x)=\bar{W}\left(\beta_{1}, \gamma_{1}\right)\right\} \\
h_{1} & :=W\left(M_{1}\right)-W\left(\mu_{1}\right)
\end{aligned}
$$

and recursively for any $n \in \mathbf{N}^{*}$,

$$
\begin{aligned}
\beta_{n+1} & :=\inf \left\{x \geq \gamma_{n} ; W(x)-\underline{W}\left(\gamma_{n}, x\right)=h_{n-1}\right\} \\
\mu_{n+1} & :=\inf \left\{x \geq \gamma_{n} ; W(x)=\underline{W}\left(\gamma_{n}, \beta_{n+1}\right)\right\} \\
\gamma_{n+1} & :=\inf \left\{x \geq \beta_{n+1} ; W(x)=W\left(\mu_{n}\right)\right\} \\
M_{n+1} & :=\inf \left\{x \geq \beta_{n+1} ; W(x)=\bar{W}\left(\beta_{n+1}, \gamma_{n+1}\right)\right\}, \\
h_{n} & :=W\left(M_{n}\right)-W\left(\mu_{n}\right) \text { and } \\
\mathcal{F}_{n} & :=\sigma\left(W(x), 0 \leq x \leq \gamma_{n}\right) .
\end{aligned}
$$



Figure 2: The variables for a sample path of $W$

Lemma 6.2 There is a real $K>0$ such that for any $n>0$ and any $\lambda>0$,

$$
\begin{aligned}
& P\left(\int_{\gamma_{n-1}}^{M_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \geq \lambda\right) \leq K e^{-j_{0}^{2} \frac{\lambda}{16}} \text { and } \\
& P\left(\int_{\mu_{n}}^{\beta_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \leq \lambda\right) \leq K\left(\frac{1}{\sqrt{e \lambda}}+\sqrt{e \lambda}\right) e^{-1 /(2 e \lambda)}
\end{aligned}
$$

where $j_{0}$ is the smallest strictly positive root of Bessel function $J_{0}$.
Proof: The process $\left(W\left(\gamma_{n-1}+x\right)-W\left(\gamma_{n-1}\right)\right)_{x \geq 0}$ is a Brownian motion independent of $\mathcal{F}_{n-1}$. Therefore given $h_{n-1}=h$, Theorem 2.4 gives us the law of the process

$$
\left(W\left(\mu_{n}+x\right)-W\left(\mu_{n}\right),-\mu_{n}+\gamma_{n-1} \leq x \leq \beta_{n}-\mu_{n}\right)
$$

So denoting by $R$ and $\tilde{R}$ two independent Bessel processes of dimension 3 started at 0, we have
$P\left(\int_{\gamma_{n-1}}^{M_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \geq \lambda\right) \leq P\left(\int_{0}^{\infty} e^{-R(x)} \mathrm{d} x+\int_{0}^{\infty} e^{-\tilde{R}(x)} \mathrm{d} x \geq \lambda\right)$.
Moreover using Le Gall's Ray-Knight Theorem (Proposition 1.1 [6]), it can be proved that $\int_{0}^{\infty} e^{-R(x)} \mathrm{d} x$ has the same law as $4 T_{Q}(1)$ where $T_{Q}(1)$ is the hitting time of height 1 by a a square of a 2 -dimensional Bessel process started in 0 . Thereby, we denote by $\tilde{T}_{Q}(1)$ a independent copy of $T_{Q}(1)$ and we have

$$
\begin{aligned}
P\left(\int_{\gamma_{n-1}}^{M_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \geq \lambda\right) & \leq P\left(T_{Q}(1)+\tilde{T}_{Q}(1) \geq \frac{\lambda}{4}\right) \\
& \leq 2 P\left(T_{Q}(1) \geq \frac{\lambda}{8}\right)
\end{aligned}
$$

Moreover according to Theorem 2 of [22], as in the proof of the previous lemma,

$$
\begin{aligned}
P\left(\int_{\gamma_{n-1}}^{M_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \geq \lambda\right) & \leq P\left(T_{Q}(1)+\tilde{T}_{Q}(1) \geq \frac{\lambda}{4}\right) \\
& \leq K e^{-j_{0}^{2} \frac{\lambda}{16}}
\end{aligned}
$$

We still have thanks to Theorem 2.4, denoting $T_{R}(1):=\inf \{x \geq 0, R(x) \geq$ $1\}$,

$$
\begin{aligned}
P\left(\int_{\mu_{n}}^{\beta_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \leq \lambda\right) & \leq P\left(\int_{0}^{T_{R}(1)} e^{-R(x)} \mathrm{d} x \leq \lambda\right) \\
& \leq P\left(T_{R}(1) \leq e \lambda\right) \\
& \leq K\left(\frac{1}{\sqrt{e \lambda}}+\sqrt{e \lambda}\right) e^{-1 /(2 e \lambda)}
\end{aligned}
$$

This concludes the proof.
We also need the following lemma:
Lemma 6.3 Let $a<e<b$. Then P-a.s., for $n$ large enough,

$$
a^{n}<h_{n}<b^{n}, a^{n}<W\left(\mu_{n}\right)-W\left(\mu_{n+1}\right)<b^{n} \text { and } a^{2 n}<\gamma_{n}<b^{2 n}
$$

Proof: We begin with the law of the sequence $\left(h_{n}\right)$. For any $h \geq 1$, any $n \in \mathbf{N}$ and any $x \geq 1$,

$$
\begin{aligned}
P\left(\left.\frac{h_{n+1}}{h_{n}} \leq h \right\rvert\, h_{n}=x\right) & =P\left(\left.\frac{h_{n+1}-h_{n}}{h_{n}} \leq h-1 \right\rvert\, h_{n}=x\right) \\
& =P\left(\tau_{W}((h-1) x) \geq \tau_{W}(-x)\right)=1-\frac{1}{h}
\end{aligned}
$$

Thus the variables $r_{n}:=h_{n+1} / h_{n}$ are independent and so $\log r_{n}$ is exponentially distributed with mean 1 . Therefore $\log h_{n}=\sum \log r_{k}$ has the law $\Gamma(n, 1)$ and for any $1<a<e, n \in \mathbf{N}$,

$$
\begin{aligned}
P\left(h_{n} \leq a^{n}\right) & =P\left(\log h_{n} \leq n \log a\right)=\int_{0}^{n \log a} \frac{x^{n-1}}{(n-1)!} e^{-x} \mathrm{~d} x \\
& \leq \frac{(n \log a)^{n}}{a^{n}(n-1)!}
\end{aligned}
$$

because the function $x \rightarrow x^{n-1} e^{-x}$ is non decreasing on $[0, n \log a]$. Stirling Formula $n!\sim\left(\frac{n}{e}\right)^{n} \sqrt{2 \pi n}$ gives us

$$
\frac{(n \log a)^{n}}{a^{n}(n-1)!} \sim \sqrt{\frac{n}{2 \pi}}\left(\frac{e \log a}{a}\right)^{n} .
$$

As for any $a \in] 1, e\left[, 0<\frac{e \log a}{a}<1\right.$, the serie $\sum P\left(h_{n} \leq a^{n}\right)$ converges. Then Borel-Cantelli lemma yields directly the first part of the lemma. We proove the upper bound in the same way.

For the second result, we begin by noting that, given $h_{n-1}=x$, the process $W\left(\mu_{n-1}\right)-W\left(\mu_{n}\right)$ has the same law as $W\left(m_{x}\right)$. Therefore,

$$
P\left(W\left(\mu_{n-1}\right)-W\left(\mu_{n}\right)<h \mid h_{n-1}=x\right)=1-e^{-h / x} \leq \frac{h}{x} .
$$

Taking $1<c<a<e$, we obtain

$$
\begin{aligned}
& P\left(W\left(\mu_{n-1}\right)-W\left(\mu_{n}\right)<c^{n-1}\right) \\
& \quad \leq P\left(W\left(\mu_{n-1}\right)-W\left(\mu_{n}\right)<c^{n-1} ; h_{n-1}>a^{n-1}\right)+P\left(h_{n-1} \leq a^{n-1}\right) \\
& \quad \leq\left(\frac{c}{a}\right)^{n-1}+P\left(h_{n-1} \leq a^{n-1}\right) .
\end{aligned}
$$

The result of the previous proof implies that the sum of

$$
P\left(W\left(\mu_{n-1}\right)-W\left(\mu_{n}\right)<c^{n-1}\right)
$$

converges and Borel-Cantelli Lemma shows that for large $n$,

$$
W\left(m_{n-1}\right)-W\left(m_{n}\right) \geq c^{n-1} .
$$

The second inequality can be proved in the same way.
We reason similarly to find an upper bound of $\gamma_{n}$ : first, we show that for $c>0$,

$$
\begin{equation*}
P\left(\frac{\gamma_{n}-\gamma_{n-1}}{h_{n}^{2}}>c\right) \leq 2 P\left(\tau_{W}(1) \geq c\right) \leq \frac{K}{\sqrt{c}} . \tag{10}
\end{equation*}
$$

Indeed, as before, conditionning by $W\left(\mu_{n}\right)$,

$$
P\left(\frac{\gamma_{n}-\beta_{n}}{h_{n}^{2}}>c\right) \leq P\left(\frac{\gamma_{n}-\beta_{n}}{W\left(\mu_{n}\right)^{2}}>c\right)=P\left(\tau_{W}(1) \geq c\right)
$$

and conditionning by $h_{n-1}$,

$$
P\left(\frac{\beta_{n}-\gamma_{n-1}}{h_{n}^{2}}>c\right) \leq P\left(\frac{\beta_{n}-\gamma_{n-1}}{h_{n-1}^{2}}>c\right) \leq P\left(\tau_{W}(1) \geq c\right) .
$$

The bound of (10) follows immediately. Now, let $\epsilon>0$ and $\gamma_{0}=0$, we have for any $n \geq 1$,

$$
\begin{aligned}
& P\left(\frac{\gamma_{n}}{h_{n}^{2}}>(1+\epsilon)^{2 n}\right) \leq P\left(\sum_{k=1}^{n} \frac{\gamma_{k}-\gamma_{k-1}}{h_{k}^{2}}>(1+\epsilon)^{2 n}\right) \\
& \quad \leq \sum_{k=1}^{n} P\left(\frac{\gamma_{k}-\gamma_{k-1}}{h_{k}^{2}}>\frac{(1+\epsilon)^{2 n}}{n}\right) \leq \frac{K n^{3 / 2}}{(1+\epsilon)^{n}} .
\end{aligned}
$$

Therefore Borel-Cantelli Lemma implies that $P$-a.s., for $n$ large enough,

$$
\gamma_{n} \leq(1+\epsilon)^{2 n} h_{n}^{2} .
$$

It is then easy to deduce the upper bound for $\gamma_{n}$. And the lower bound can be obtain easily using same techniques as before.

It is now possible to control the asymptotic behaviour of the integral:
Proposition 6.4 $P$-almost surely,

$$
\begin{aligned}
& \liminf _{v \rightarrow \infty} \log _{2} v \int_{a_{v}}^{b_{v}} e^{-W(x)+W\left(m_{v}\right)} \mathrm{d} x \geq \frac{1}{2 e} \text { and } \\
& \limsup _{v \rightarrow \infty} \frac{1}{\log _{2} v} \int_{a_{v}}^{b_{v}} e^{-W(x)+W\left(m_{v}\right)} \mathrm{d} x \leq \frac{16}{j_{0}^{2}}
\end{aligned}
$$

Proof: Let us begin with the lower bound. Fix $c>2 e$, according to Lemma 6.2 , for any $n \in \mathbf{N}^{*}$,

$$
P\left(\int_{\mu_{n}}^{\beta_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \leq \frac{1}{c \log (n-1)}\right) \leq K \frac{\sqrt{\log (n-1)}}{(n-1)^{c / 2 e}} .
$$

According to the first Borel-Cantelli Lemma, $P$-a.s. for $n$ large enough,

$$
\int_{\mu_{n}}^{\beta_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x>\frac{1}{c \log (n-1)} .
$$

We come back to the $m_{v}$. We fix $v>1$. $P$-a.s. there is a unique $n \in \mathbf{N}^{*}$ such that $h_{n-1}<v \leq h_{n}$. Thus necessarily $m_{v}=\mu_{n}$ and $b_{v} \geq \beta_{n}$. Therefore

$$
\int_{m_{v}}^{b_{v}} e^{-W(x)+W\left(m_{v}\right)} \mathrm{d} x \geq \int_{\mu_{n}}^{\beta_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x>\frac{1}{c \log (n-1)}
$$

Let $a \in(1, e)$, according to Lemma 6.3, if $n$ is large enough, $h_{n-1}>a^{n-1}$. Thereby $P$-a.s., for $v$ large enough,

$$
\int_{m_{v}}^{b_{v}} e^{-W(x)+W\left(m_{v}\right)} \mathrm{d} x \geq \frac{1}{c\left(\log _{2} v-\log _{2} a\right)}
$$

When $c$ goes to $2 e$, we obtain the result of the proposition.
Now we continue with the upper bound. Fix $c>16 / j_{0}^{2}$. One more time thanks to Lemma 6.2 and Borel-Cantelli Lemma, $P$-a.s. for $n$ large enough,

$$
\int_{\gamma_{n-1}}^{M_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x<c \log (n-1)
$$

Fix $v>1, P$-a.s., there is a unique $n \in \mathbf{N}^{*}$ such that $h_{n-1}<v \leq h_{n}$. Therefore $m_{v}=\mu_{n}, b_{v} \leq M_{n}$ and taking $1<a<e<b$,

$$
\begin{aligned}
\int_{a_{v}}^{b_{v}} e^{-W(x)+W\left(m_{v}\right)} \mathrm{d} x & \leq \int_{0}^{M_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \\
& \leq c_{n} e^{-W\left(\mu_{n-1}\right)+W\left(\mu_{n}\right)}+\int_{\gamma_{n}}^{M_{n}} e^{-W(x)+W\left(\mu_{n}\right)} \mathrm{d} x \\
& \leq b^{n} e^{-a^{n}}+c \log (n-1) \\
& \leq b^{n} e^{-a^{n}}+c\left(\log _{2} v-\log _{2} a\right)
\end{aligned}
$$

Thus,

$$
\limsup \frac{1}{\log _{2} v} \int_{a_{v}}^{b_{v}} e^{-W(x)+W\left(m_{v}\right)} \mathrm{d} x \leq c
$$

and when $c$ goes to $16 / j_{0}^{2}$, this yields to the upper bound.
We can now come back to the process $\mathbf{L}^{*}$.

### 6.2 End of the proof of Theorem 1.1

The previous results allow us to know the asymptotic behaviour of $\mathbf{L}^{*}$. Using (8) with $v_{n}=n^{2 / 3}$ and Borel-Cantelli Lemma, we obtain that, $P$-almost surely for $n$ large enough,

$$
\frac{e^{v_{n}}}{J_{v_{n}}+1 / v_{n}} \leq \mathbf{L}^{*}\left(e^{v_{n}}\right) \leq \frac{e^{v_{n}}\left(1+1 / v_{n}\right)}{j_{v_{n}}\left(1-1 / v_{n}\right)}
$$

Thereby Proposition 6.4 gives us

$$
\begin{aligned}
& \limsup _{n \rightarrow \infty} \frac{\mathbf{L}^{*}\left(e^{v_{n}}\right)}{e^{v_{n}} \log _{2} v_{n}} \leq \frac{1}{\liminf j_{v_{n}} \log _{2} v_{n}} \leq 2 e \text { and } \\
& \liminf _{n \rightarrow \infty} \frac{\log _{2} v_{n} \mathbf{L}^{*}\left(e^{v_{n}}\right)}{e^{v_{n}}} \geq \liminf _{n \rightarrow \infty} \frac{\log _{2} v_{n}}{J_{v_{n}}} \geq \frac{j_{0}^{2}}{64}
\end{aligned}
$$

Moreover as $\mathbf{L}^{*}$ is non decreasing, denoting by $[x]$ the integer part of $x$,

$$
\begin{aligned}
\frac{j_{0}^{2}}{64} \leq \liminf _{v \rightarrow \infty} \frac{\mathbf{L}^{*}\left(e^{\left[v^{3 / 2}\right]^{2 / 3}}\right)}{e^{\left[v^{3 / 2}\right]^{2 / 3}}} & \leq \liminf _{v \rightarrow \infty} \frac{\mathbf{L}^{*}\left(e^{v}\right)}{e^{v}} \text { et } \\
\limsup _{v \rightarrow \infty} \frac{\mathbf{L}^{*}\left(e^{v}\right)}{e^{v}} & \leq \limsup _{v \rightarrow \infty} \frac{\mathbf{L}^{*}\left(e^{\left(\left[v^{3 / 2}\right]+1\right)^{2 / 3}}\right)}{e^{\left(\left[v^{3 / 2}\right]+1\right)^{2 / 3}}} \leq 2 e
\end{aligned}
$$

and we have shown the first part of Theorem 1.1.
For the second part, we use this time (9) with $v_{n}=e^{n}$ and Borel-Cantelli Lemma. We then have, $P$-almost surely for $n$ large enough,

$$
\frac{e^{v_{n}}}{I_{v_{n}}^{-}+\widehat{I}_{v_{n}}^{-}+1 / n} \leq \mathbf{L}^{*}\left(e^{v_{n}}\right) \leq \frac{e^{v_{n}}(1+1 / n)}{I_{v_{n}}^{-} \wedge \widehat{I}_{v_{n}}^{-}(1-1 / n)}
$$

Lemma 6.1 leads directly to

$$
\liminf _{v \rightarrow \infty} \frac{\mathbf{L}^{*}\left(e^{v}\right)}{e^{v}} \leq e \pi^{2}
$$

And the proof of the theorem is completed.

## 7 What can be added

The methods we used allow us to obtain an other result on the position of the favorite point of the diffusion. We denote by $m^{*}(t)$ the favorite point with the smaller modulus for the diffusion $\mathbf{X}$. Using a method similar to the one in the proof of Proposition 5.2, it can be shown that $P$-a.s., for $v$ large enough, there exists a real $m_{v}^{\#} \in\left\{m_{v, 1}^{R,-}, m_{v}^{R,+}, \hat{m}_{v, 1}^{R,-}, \hat{m}_{v}^{R,+}\right\}$, such that with natural notations, $m_{e^{v}}^{*} \in\left[a_{v}^{\#}, b_{v}^{\#}\right]$ and

$$
0 \leq W\left(m_{e^{v}}^{*}\right)-W\left(m_{v}^{\#}\right) \leq 2 \log \frac{r(v)}{R(v)} \leq 2 \log _{3} v
$$

We can also prove that for any $\epsilon>0, P$-a.s., for $v$ large enough, for any $t \in\left[a_{v}, m_{v}-(\log v)^{2+\epsilon}\right] \cup\left[m_{v}+(\log v)^{2+\epsilon}, b_{v}\right]$,

$$
W(t)-W\left(m_{v}\right)>3 \log _{3} v
$$

Hence we obtain the following:
Theorem 7.1 Let c be a real number strictly larger than 2. P-almost surely,

$$
\limsup _{t \rightarrow \infty} \frac{\left|m^{*}(t)-m_{\log t}^{\#}\right|}{\left(\log _{2} t\right)^{c}}=0
$$

This theorem has to be compared with the corollary in the introduction of [4], note that accuracy gained in time is lost in space.

Acknowledgment: I thank Romain Abraham and Pierre Andreoletti for their helpful remarks and comments.

## References

[1] P. Andreoletti and R. Diel. Limit law of the local time for Brox's diffusion. Preprint, 2008.
[2] A. Borodin and P. Salminen. Handbook of Brownian Motion-Facts and Formulae. Birkhäuser, 1996.
[3] T. Brox. A one-dimensional diffusion process in a Wiener medium. Ann. Probab., 14(4): 1206-1218, 1986.
[4] D. Cheliotis. Localization of favorite points for diffusion in random environment. Stochastic Process. Appl., 118(7): 1159-1189, 2008.
[5] A. Dembo, N. Gantert, Y. Peres, and Z. Shi. Valleys and the maximal local time for random walk in random environment. Probability Theory and Related Fields, 137: 443-473, 2007.
[6] J.F. Le Gall. Sur la mesure de Hausdorff de la courbe brownienne. Sém. Probability. XIX. Lect. Notes Math., 1123: 297-313, 1985.
[7] N. Gantert, Y. Peres, and Z. Shi. The infinite valley for a recurrent random walk in random environment. Preprint.
[8] A. O. Golosov. On limiting distribution for a random walk in a critical one-dimensional random environment. Com. of the Mosc. Math. Soc., pages 199-200, 1986.
[9] Y. Hu and Z. Shi. The limits of Sinai's simple random walk in random environment. Ann. Probab., 26(4): 1477-1521, 1998.
[10] Y. Hu and Z. Shi. The local time of simple random walk in random environment. J. of Theoret. Probab., 11(3), 1998.
[11] H. Kesten. The limit distribution of Sinai's random walk in random environment. Physica, 138A: 299-309, 1986.
[12] J. Neveu and J. Pitman. Renewal property of the extrema and tree property of the excursion of a one-dimensional Brownian motion. séminaire de probabilitées xxiii. Lecture Notes Math., 1372:239-247, 1989.
[13] D. Revuz and M. Yor. Continuous Martingales and Brownian Motion. Springer, 1999.
[14] S. Schumacher. Diffusions with random coefficients. Contemp. Math., 41: 351-356, 1985.
[15] Z. Shi. A local time curiosity in random environment. Stoch. Proc. Appl., 76(2): 231-250, 1998.
[16] Z. Shi. Sinai's walk via stochastic calculus. Panoramas et Synthèses, 12: 53-74, 2001.
[17] Ya. G. Sinai. The limit behaviour of a one-dimensional random walk in a random medium. Theory Probab. Appl., 27(2): 256-268, 1982.
[18] F. Solomon. Random walks in random environment. Ann. Probab., 3(1): 1-31, 1975.
[19] M. Talet. Annealed tail estimates for a Brownian motion in a drifted Brownian potential. Ann. Probab., 35(1): 32-67, 2007.
[20] H. Tanaka. Limit theorem for one-dimensional diffusion process in Brownian environment. Lecture Notes in Mathematics, 1322, 1988.
[21] H. Tanaka. Localization of a diffusion process in a one-dimensional Brownian environmement. Comm. Pure Appl. Math., 17: 755-766, 1994.
[22] S.J. Taylor Z. Ciesielski. First passage times and sojourn times for Brownian motion in space and the exact Hausdorff measure of the sample path. Trans. Amer. Math. Soc., 103: 434-450, 1962.


[^0]:    *Laboratoire MAPMO - C.N.R.S. UMR 6628 - Fédération Denis-Poisson, Université d'Orléans, (Orléans France).

    MSC 2000 60J25; 60J55.
    Key words : Diffusion in Brownian environment, Local time

