

# Magnetic resonance imaging method based on magnetic susceptibility effects to estimate bubble size in alveolar products: application to bread dough during proving.

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1	Magnetic resonance imaging method based on magnetic
2	susceptibility effects to estimate bubble size in alveolar products:
3	Application to bread dough during proving
4	
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### 18 Abstract

19 Magnetic resonance imaging has proven its potential application in bread dough and gas cell monitoring studies, and dynamic processes such as dough proving and baking can be 20 21 monitored. However, undesirable magnetic susceptibility effects often affect quantification 22 studies, especially at high fields. A new low field method is presented based on local assessment of porosity in spin-echo imaging, local characterization of signal loss in gradient-23 24 echo imaging and prediction of relaxation times by simulation to estimate bubble radii in 25 bread dough during proving. Maps of radii showed different regions of dough constituting 26 networks which evolved during proving. Mean radius and bubble distribution were assessed 27 during proving.

28

Key Words: magnetic susceptibility; MRI simulation; field inhomogeneity; gas bubble; gas
cell; bread; dough; proving; fermentation

#### 31 **1. Introduction**

32 Magnetic susceptibility effects are of increasing interest in Magnetic Resonance 33 Imaging (MRI). Traditionally, as field inhomogeneities induce geometry and intensity 34 distortions in spin-echo (SE) imaging and signal loss in gradient-echo (GE) imaging [1, 2], 35 susceptibility artefacts have been considered undesirable and therefore subjected to correction 36 [3]. However, there are number of emerging techniques and applications using magnetic 37 susceptibility effects and T<sub>2</sub>\*-weighted imaging. The difference in magnetic susceptibility 38 between deoxygenated and oxygenated blood, termed the BOLD effect, led to functional MRI 39 which is widely used in cognitive neuroscience [4]. Susceptibility Weighted Imaging (SWI) 40 uses the original phase image to enhance contrast between tissues with different 41 susceptibilities [5]. This technique has been applied to studies of brain tumors, trauma, 42 vascular malformations and for quantification of brain iron [6]. Superparamagnetic iron oxide 43 (SPIO) particles create intense magnetic field distortions within and around cells, leading to 44 irreversible signal dephasing in GE sequences. SPIO particles have been exploited as contrast 45 agents for non-invasive cell tracking to determine their biodistribution in different organs [7-46 10]. Susceptibility studies have been applied to material characterization [11, 12] and particle 47 identification in industrial systems [13].

48 The potential of gas microbubbles as an in vivo intravascular susceptibility contrast 49 agent for MRI has been demonstrated [14]. Moreover, differences in magnetic susceptibility 50  $(\Delta \gamma)$  between gas bubbles and a water environment create field inhomogeneities which induce 51 intravoxel dephasing and associated signal loss characterized by a T<sub>2</sub>\* shortening in GE 52 imaging. However, the presence of air bubbles in several food products has not to date been 53 the focus of studies of susceptibility effects. Understanding bubble mechanics is very 54 important with a wide range of foamed food products. Bubble studies make possible the optimization of process design, contribute to the development of strategies for deaeration and 55

are important in the area of texture and sensory analysis [15]. It can be anticipated that air or gas bubbles would be natural markers or contrast agents. Moreover, bubbles are part of a product and thus could be considered as a signature of the medium.

59 Proving flour dough is an essential stage in breadmaking. As the final character of most bakery products depends on the creation and control of gas bubble structures in the unbaked 60 61 matrix, improving the understanding of the nucleation and growth of bubbles is of major 62 interest [16]. MRI is a particularly suitable tool to study bread dough as it is non-invasive 63 (essential for dough monitoring during proving as the alveolar structure is fragile and may 64 collapse if intrusive measurement techniques are used), and provides a good trade-off between 65 spatial and time resolution [17-20]. Dough samples can thus be quite large and imaged with high resolution during proving [20]. Even the entire baking process can be monitored with 66 67 MRI [21]. Bonny et al. [17] used high-field (9.4 T) magnetic resonance microscopy along 68 with routine mathematical morphology to characterize the proving process noninvasively. 69 They faced critical susceptibility effects responsible for severe geometrical distortions due to 70 the heterogeneous structure characterized by many interfaces. Different stages of fermentation 71 could be identified, with a general increase in bubbles, but no quantitative assessment of 72 bubble size could be made. Van Duynhoven et al. [20] illustrated the ability of MRI and 73 image analysis to assess gas cell development in the growth of the dough during proving. 74 Basically, the size of each cell was determined by counting the number of pixels, thus 75 presupposing high resolution and hence small fields of view. Moreover, this method does not 76 account for susceptibility effects that are especially disrupting at 4.7 T and are responsible for 77 geometry and intensity distortions in SE imaging. X-ray tomography has also proved its 78 ability in analyzing bubble growth during bread making [22]. However, this technique is only 79 sensitive to density, and imaged samples have to be small (< 10 mm).

80 The aim of the study reported here was to characterize the alveolar structure of bread 81 dough and its evolution during proving using susceptibility effects. Grenier et al. [18] 82 assessed local dough porosity (volume of gas per volume of dough) during proving. However, 83 nothing can be inferred about bubble size for a given value of porosity, as different cell 84 distributions can be obtained [23, 24]. We therefore developed an original method to estimate 85 bubble size based on the local porosity assessed in SE imaging and on the susceptibility 86 effects in GE imaging. MRI simulation was used to predict signal loss for virtual networks of 87 gas bubbles embedded in a water matrix. A low magnetic field ( $B_0 = 0.2$  T) was used in order 88 to avoid the severe geometric distortions encountered at high field in SE imaging and critical 89 signal losses in GE imaging as the heterogeneous structures were characterized by many 90 susceptibility interfaces [17]. The dynamic side of bread proving requires acquisition to be 91 quite fast, leading to a limited number of acquisitions, limited resolution and an appropriate 92 hypothesis to compute maps of interest.

93 Section 2 reports how relaxation times affected by susceptibility were predicted with 94 simulation. Section 3 describes the materials and methods used with bread dough 95 preparations, the MRI protocol, computation of the maps and the principle of the estimation 96 algorithm. Section 4 presents the results on non-yeasted bread dough and yeasted dough 97 during proving, followed by a discussion in section 5.

## 98 **2.** Prediction of relaxation times by simulation

99 The algorithm developed for estimation of bubble size was partly based on simulation 100 results. This section introduces theoretical considerations and describes the simulation 101 protocol. We demonstrate how relaxation times can be modeled at 0.2 T and what can happen 102 at 1.5 T.

#### 103 2.1. Theoretical considerations

We first verified that the static dephasing regime conditions for spherical magnetic perturbers of radius *r* as described by Yung et al. [25] were largely met, as  $1/\tau \ll \delta \omega$  with  $\tau = r^2/D$  and  $\delta \omega = \gamma \Delta \chi B_0/3$ , *D* being the translational diffusion coefficient of spins taken as the water diffusion coefficient and  $\gamma$  the proton gyromagnetic ratio. In fact, *D* is even smaller in bread dough.

109 Microscopic field inhomogeneities (i.e. magnetic field inhomogeneities over distances 110 with orders of magnitude smaller than the voxel size) are responsible for irreversible signal decay  $(R_2 = 1/T_2)$  and mesoscopic field inhomogeneities (from perturbers smaller than the 111 voxel size but greater than the diffusion length) contribute to  $R_2^{\#}$ , the reversible portion of  $R_2^{*}$ 112  $(R_2^* = R_2 + R_2^{\#} = 1/T_2^*)$  [3, 26]. Macroscopic field inhomogeneities are not present in  $T_2^*$ 113 definition and induce non-exponential signal decay [3, 27]. Macroscopic field 114 115 inhomogeneities can typically arise from air inclusions or ferromagnetic objects. Finally, 116 assuming a constant proton density across the voxel ( $\rho(\mathbf{r}) = \rho$ ), signal decay in a voxel of volume V is expressed in Eq. (1): 117

118 
$$S(t)_{voxel} = \rho / V.e^{-R_2 t}. sinc(\mu g_x l_x/2). sinc(\mu g_y l_y/2). sinc(\mu g_z l_z/2)$$
(1)

119 where  $l_x$ ,  $l_y$ ,  $l_z$  are the voxel dimension, and field inhomogeneities along one direction *i* are 120 expressed with a linear gradient ( $\Delta B_i = g_{i.}i$ ). In this study, we defined  $T_{2\chi}$  as the constant of 121 the exponential curve fitting the initial part of the susceptibility-induced signal decay. 122 Microscopic, mesoscopic and macroscopic inhomogeneities were thus taken into account in 123 the approximation of  $T_{2\chi}$ .

#### 124 2.2. Simulation protocol

125 The purpose of the simulation was to investigate quantitatively the signal loss in GE 126 MR images induced by networks of gas bubbles embedded in water. We used the SIMRI MRI

127 simulator as described by Benoit-Cattin et al [28]. In a previous study [27], intravoxel 128 modeling and associated signal decay were quantitatively assessed and experimentally 129 validated in the case of a single well-resolved air-filled cylinder and in the case of a network 130 of small interacting air-filled cylinders. An overview of the simulation framework is 131 presented in Fig. 1. A 3D virtual object was defined by radius (r) and center to center distance between bubbles (*ccd*) expressed in pix<sub>obi</sub> (i.e. the pixel in the 256<sup>3</sup> object space), difference in 132 133 magnetic susceptibility ( $\Delta \chi$ ) and number of bubbles in each dimension (Nb). The additive 134 property of the magnetic field, and the analytical formulation of field inhomogeneities created by one sphere [1], were combined to compute the relative field inhomogeneities ( $\Delta B/B_0$ ). 135 136 From this map, SIMRI provides simulated MR images with susceptibility effects. At this stage, object parameters such as spin-lattice relaxation time  $T_{I_1}$  spin-spin relaxation time  $T_{2_2}$ 137 138 and proton density  $\rho$  could be defined. NMR experiments (Minispec PC 120, Bruker SA, 139 Wissembourg, France) on the bread dough showed several  $T_1$  and  $T_2$  components leading to 140 multi-exponential decay, as previously described [29, 30]. The first components were smaller 141 than the echo times (TE) used in this study. As a single value for relaxation times is necessary 142 in simulation, weighted averaging of the relaxation times of the different components was performed, resulting in  $T_1 = 100$  ms and  $T_2 = 20$  ms. Lodi et al. [31] assessed  $T_2$  values from 143 spin-echo images by fitting time-series points pixel-by-pixel, resulting in  $T_2$  maps of soy 144 145 bread (mean values around 18 ms) consistent with this averaging. Simulation parameters such as main magnetic field ( $B_0 = 0.2$  T), pixel bandwidth (BW = 279 Hz/pixel) and repetition time 146 (TR = 300 ms) were set to match real experimental conditions. By defining the number of 147 148 pixels constituting the simulated image, we could define the relative proportion of the bubble 149 radius in the simulated image pixel (pixima). For instance, considering the high-resolution object used as an example in Fig. 1 ( $r = 4 \text{ pix}_{obj}$ ,  $ccd = 12 \text{ pix}_{obj}$ ,  $\Delta \chi = -9.05 \text{ ppm}$ , Nb = 17 in150 the 256<sup>3</sup> object space), a 32x32 slice simulated image resulted in r = 0.5 pix<sub>ima</sub>, ccd = 1.5151

152  $pix_{ima}$  ( $r = 0.125 pix_{ima}$ ,  $ccd = 0.375 pix_{ima}$  with a 8x8 simulated image). For each object, GE 153 images were simulated at different increasing TE from 4 to 12 ms. The mean gray level was 154 then computed on regions of interest (ROI) enclosing several bubbles (visible on Fig. 1). This 155 centred ROI made it possible to avoid boundary effects and to account for the interactions between several bubbles. Good representation of the medium was thus assured with such an 156 157 ROI as bubbles were regularly spaced.  $T_{2\chi}$  was calculated as the constant of the exponential 158 curve fitting the initial part of the mean gray level decay computed in the ROI. In fact, objects 159 with large *ccd* gave quasi-exponential decay even for longer *TE* (12 ms). As *ccd* decreased, signal decay was hardly exponential because of high gradients  $g_i$  in Eq. (1) and  $T_{2\chi}$  was 160 161 assessed on the very first TE values (between 4 and 8 ms).

#### 162 2.3. Evolution of $T_{2\chi}$ as a function of bubble radius

163 Our aim was to link the object configuration with  $T_{2\chi}$  values so as to characterize the 164 medium. The initial question was to find out whether small or large intravoxel bubbles 165 engendered a similar signal loss for the same porosity (i.e. the gas volume fraction in the voxel). We therefore gathered all  $T_{2\chi}$  values relative to the different object configurations as a 166 167 function of bubble radius at the different porosities shown in Fig. 2.a. The same curve shape 168 was observed for each porosity. The first plateau indicated that a similar initial signal loss was 169 measured for small radii, i.e.  $T_{2\chi}$  was only dependent on porosity. Then, a rapid increase in 170  $T_{2\chi}$  leveled off on a second plateau; note that the lower the porosity the smaller the radii corresponding to  $T_{2\chi}$  values as they started to increase. This was appropriate to our study as 171 172 the initial mean porosity of bread dough was around 10% and overall porosity and bubbles 173 increased during proving. One significant feature was that  $T_{2\chi}$  was sensitive to porosity and 174 intravoxel structure according to bubble size.

#### 175 2.4. Modeling of $T_{2\chi}$ at 0.2 T

176  $T_{2\chi}$  as a function of bubble radius was modeled with a sigmoid curve for each porosity 177 value; a typical sigmoid function being defined in Eq. (2).

178 
$$f(x) = \frac{1}{1 + e^{-\lambda x}}$$
 (2)

As curves were not centered but shifted along the radius axis, we defined  $r_{break}$  as the radius corresponding to the symmetrical point of the curve. As the amplitude between low and high radius values was not 1, *K* was introduced as multiplying factor. Finally,  $T_{2\chi_{init}}$  was defined as  $T_{2\chi}$  offset, i.e. the value of  $T_{2\chi}$  for the smallest radius. Using these new variables and Eq. (2),  $T_{2\chi}$  as a function of radius is expressed in Eq. (3):

184 
$$T_{2\chi}(r) = T_{2\chi init} + K \cdot \frac{1}{1 + e^{-\lambda(r - r_{break})}}$$
 (3)

For several porosity values, the parameters of Eq. (3) ( $T_{2\chi_{init}}$ , K,  $r_{break}$ ,  $\lambda$ ) were assessed to best fit the data according to the least mean square criterion (using Marquardt's algorithm in TableCurve2D® software, Jandel Scientific, version 5). Each parameter separately was then studied as a function of porosity.  $T_{2\chi_{init}}$  and K (Fig. 3.a),  $r_{break}$  and  $\lambda$  (Fig. 3.b) were drawn up and found to be closely related to porosity with second order polynomial functions ( $R^2(T_{2\chi_{init}}) = 0.995$ ,  $R^2(K) = 0.988$ ,  $R^2(r_{break}) = 0.994$  and  $R^2(\lambda) = 0.993$ ).

191 Finally, a complete description of  $T_{2\chi}$  as a function of radius and porosity ( $\varepsilon$ ) is 192 expressed in Eqs. (4,5).

193 
$$T_{2\chi}(r,\varepsilon) = T_{2\chi init}(\varepsilon) + K(\varepsilon) \cdot \frac{1}{1 + e^{-\lambda(\varepsilon)(r - r_{break}(\varepsilon))}}$$
(4)

$$T_{2\chi init}(\varepsilon) = 27.1\varepsilon^{2} - 37.1\varepsilon + 17.2$$

$$K(\varepsilon) = -52.9\varepsilon^{2} + 49.3\varepsilon - 4.2$$

$$\lambda(\varepsilon) = 133.9\varepsilon^{2} - 127.5\varepsilon + 63.6$$

$$r_{break}(\varepsilon) = -1.1\varepsilon^{2} + 1.9\varepsilon + 0.2$$
(5)

Note that the polynomial fitting on *K* as a function of  $\varepsilon$  was accurate and determined for porosity values greater than 10%. It was obvious that a homogeneous object without bubbles ( $\varepsilon = 0$ ) would normally result in a zero multiplying factor (*K*). The same was valid for  $r_{break}$ , whose limit for low porosities should be zero. Moreover, zero porosity in the  $T_{2\chi_{init}}$  expression corresponded to the  $T_2^*$  value, i.e. the exponential signal decay constant representing the signal loss due to the mesoscopic field inhomogeneities.

#### 201 2.5. Signal losses at 1.5 T

202 Although a low magnetic field was used, we felt that it could be of interest to predict 203 signal losses in simulation at a higher field. The same simulations were therefore undertaken 204 at 1.5 T and the results are summarized in Fig. 2.b. Susceptibility was obviously stronger as 205 field inhomogeneities were proportional to the main magnetic field. When porosity was equal, 206  $T_{2\chi_{init}}$  times were considerably shortened compared to the situation at  $B_0 = 0.2$  T. Indeed, if  $\varepsilon$ 207 = 0.12, simulations gave the followings:  $T_{2\chi_{init}}$  = 12.65 ms at  $B_0$  = 0.2 T and  $T_{2\chi_{init}}$  = 3.8 ms at 208  $B_0 = 1.5$  T. The value of  $B_0$  did not seem to impact on  $r_{break}$  value. However, the K factor was 209 significantly increased (1.35 at  $B_0 = 0.2$  T against 5.11 at  $B_0 = 1.5$  T if  $\varepsilon = 0.12$ ). Signal decay 210 was dramatically decreased for higher porosities ( $\varepsilon > 0.2$ ), preventing GE studies at high 211 fields.

212

We have shown in this section that  $T_{2\chi}$  relaxation time could be modeled with a sigmoid function dependent on porosity and radius of bubbles. We therefore used Eqs. (4,5) in the next section to build an estimation algorithm to assess bubble size in bread dough.

#### 216 **3. Materials and methods**

#### 217 *3.1.* Sample preparation and experimental procedure

218 Bread dough samples were obtained by mixing 2000 g of wheat flour (Type 55, 219 Moulins Soufflet Pantin), 1140 g of water, 40 g of salt, 20 g of improver and 100 g of yeast 220 (optional) in a Moretti Forni Grain Spiry 8 dough mixer for 17 min at 100 rotations per 221 minute. The advantages of using non-veasted dough were to provide a stable object with 222 regard to the acquisition time, with possible comparison of structure with yeasted dough at the 223 initial time of proving. A fraction of gas (mainly air) is incorporated at the mixing stage and 224 porosity at the end of mixing was estimated at  $10 \pm 2\%$  depending on the mixer used for a 225 given recipe [32], with a mean gas bubble size of approximately 0.05 mm to 0.3 mm [22].

226 Dough temperature and water content were checked for the evaluation of the 227 reproducibility between batches as both are known to affect the relaxation signal of dough 228 [29, 30]. The temperature was  $24.5 \pm 0.5$  °C, and water content was  $45 \pm 3$  g of water per 100 229 g of dough. Cylindrical flasks ( $\emptyset = 50 \text{ mm}$ , 70 mm height) were filled with 50 g of yeasted 230 dough or 100 g of non-yeasted dough. A lid was placed over the flask to limit dehydration 231 during measurement. Flasks were then placed in a tunnel within the magnet, equipped with 232 thermal regulation set at 24.5°C. Internal tunnel temperature was monitored with 233 thermocouples. MRI acquisitions were begun at approximately t= 7 min, t=0 referring to the 234 end of mixing.

Cylindrical flasks were also filled with MnCl<sub>2</sub> solution, the concentration (1287.2  $\mu$ M) being adjusted to obtain a  $T_2$  value close to that of bread dough, i.e.  $T_2 = 20$  ms. SE images of the phantom were used to normalize dough MR images to correct for inherent magnet and coil spatial inhomogeneities.

#### 239 3.2. MRI device and parameters

240 A 0.2 T electromagnet scanner in open configuration (Magnetom Open, Siemens, 241 Erlangen, Germany) equipped with a head coil was used to image the bread dough. Double GE sequences (TE<sub>1</sub> = 4 ms, TE<sub>2</sub> = 12 ms) were performed to assess signal loss and  $T_{2\chi}$  during 242 dough proving. SE sequences (TE = 8 ms) were used to obtain local porosity. For all 243 244 sequences, slice thickness (ST), bandwidth (BW), field of view (FOV), matrix size (N), 245 number of accumulations (Acc) and repetition time (TR) were set as follows: ST = 5 mm, BW = 279 Hz/pixel,  $FOV = 128 \times 128 \text{ mm}^2$ ,  $N = 128 \times 128$ , Acc = 4, TR = 300 ms. Each sequence 246 247 lasted 2 min 33 s, and GE and SE sequences were alternated.

248 3.3. Porosity map

Assuming an initial uniform porosity at a reference time, the porosity of a voxel (i, j) can be estimated from its gray level during proving in SE imaging

251 [18]: 
$$\varepsilon(i, j) = 1 - (1 - \varepsilon_{init}) \cdot \frac{SE(i, j)}{MGL_{init}}$$
 (6)

with  $\varepsilon_{init}$  the initial overall porosity ( $\varepsilon_{init} = 10\%$ ) at t<sub>init</sub> = 9.5 min, i.e. the time to obtain the first SE image, *MGL<sub>init</sub>* the corresponding mean gray level computed on a reference ROI, and *SE(i,j)* the gray level in the SE image at the location (i,j). We verified that the mean value computed in the porosity map was concordant with the overall porosity computed from the total dough volume measurement.

257 *3.4. T*<sub>2</sub>χ map

258 With  $GE_1(i,j)$  and  $GE_2(i,j)$  as the GE images at  $TE_1$  and  $TE_2$ , the  $T_2\chi$  map can be defined 259 as follows:

260 
$$T_{2\chi}(i,j) = \frac{(TE_2 - TE_1)}{\ln(\frac{GE_1(i,j)}{GE_2(i,j)})}$$
(7)

In fact, the exponential signal decay constant was estimated from two points only. Prior simulations showed that  $TE_1 = 4 \text{ ms}$  and  $TE_2 = 12 \text{ ms}$  were suitable for a 0.2 T magnetic field.

#### 263 3.5. Bubble size estimation algorithm

The algorithm principle was to combine  $T_{2\chi}$  and porosity maps to estimate local bubble radii. Indeed, from Eq. (4) we can extract bubble radii (*r*) and compute the map of radii r(i,j)according to Eq. (8):

267 
$$r(i,j) = r_{break}(i,j) - \frac{1}{\lambda(i,j)} \cdot \ln(\frac{K(i,j)}{T_{2\chi}(i,j) - T_{2\chi_{init}}(i,j)} - 1)$$
(8)

268  $r_{break}$ , K,  $\lambda$  and  $T_{2\chi_{init}}$  were determined for each pixel from the porosity map using Eq. (5).

 $T_{2\chi}$  was simulated from a set of several bubbles, i.e. signal loss was characterized in the 269 270 center of the image to take into account the influence of the neighboring bubbles. This local 271 approach based on several voxels should result in a coherent estimation algorithm, and 272 therefore r(i,j) was not directly computed pixel-by-pixel but through a 3x3 averaging mask. 273 This allowed replacement of the pixel value by a local average around the pixel of interest. 274 Indeed, the non-linear property of the logarithm function made the subsequent use of an 275 averaging on the map of radii impossible. Prior to computation of Eq. (8), we therefore applied the mask on  $r_{break}$ , K,  $\lambda$ ,  $T_{2\chi}$  and  $T_{2\chi_{init}}$  maps. 276

277

#### 278 **4. Results**

#### 279 4.1. Non-yeasted dough

The porosity map of the non-yeasted dough was computed according to Eq. (6) and is represented in Fig. 4. The largest bubbles were easily detectable and the rest of the dough was quite uniform and dense. Mean porosity on a large reference ROI was found to be 12.6%. The corresponding mean value in the  $T_{2\chi}$  map computed on the same ROI was 12.9 ms. Compared

with simulation results ( $T_{2\chi} = 13$  ms and  $\varepsilon = 12.5$  % in Fig. 2.a), these two values were in 284 good agreement. The  $T_2$  value of bread dough used for simulation and the general hypothesis 285 286 of regularly spaced bubbles were justified to reach a good order of magnitude. The bubble 287 size estimation algorithm was then applied to obtain the map of radii depicted in Fig. 5. The 288 large bubbles were even more clearly revealed due to increased contrast compared to the 289 neighboring environment. Fig.5 also demonstrates localized regions of smaller bubbles which 290 were hardly visible in the porosity map (Fig. 4). The map of radii depicted different regions 291 corresponding to different bubble sizes. Smaller bubbles can be seen surrounded by larger 292 bubbles, themselves surrounded by larger bubbles and so on. This radius evolution was 293 spatially reproduced on the dough area, leading to different visible networks. The histogram representing the number of pixels according to their estimated radii extracted from the 294 295 reference ROI is presented in Fig. 6. A Gaussian probability density function was found to fit well the experimental distribution of bubble size, as shown in Fig. 6 ( $\sigma = 0.062$ ,  $\mu = 0.37$ ,  $R^2 =$ 296 297 0.991 with TableCurve2D® software, Jandel Scientific, version 5). The mean value computed 298 in the reference ROI of the r(i,j) map was r = 0.38 pix<sub>ima</sub>. As detailed in the discussion 299 section, estimated radius values were subjected to certain limits inherent to the method and 300 values were not systematically significant. However, they can provide relative information 301 about spatial differences (for a given protocol) and between different protocols of dough 302 production.

#### 303 *4.2.* Dough during proving

As in the previous experiment with the non-yeasted dough, r(i,j) was then computed for dough during proving. The three first maps at  $t_1 = 7 \text{ min}$ ,  $t_2 = 12.5 \text{ min}$  and  $t_3 = 18 \text{ min}$  after completing mixing are represented in Fig. 7 on the same scale as Fig. 5. Histograms were computed at  $t_1$ ,  $t_2$  and  $t_3$ , taking the same reference ROI occupying almost all the dough at  $t_1$ (Fig. 8). As in the non-yeasted dough, distributions were Gaussian and shifted to the higher 309 radius during proving, consistent with the growth of bubbles under 310 desolubilization/vaporization of CO<sub>2</sub>, ethanol and water vapor. The average radius estimated 311 from the algorithm as a function of the proving time is also presented in Fig. 9. Bubble size 312 distribution was slightly shifted compared to non-yeasted dough, which may be attributed to 313 the start of fermentation between the end of mixing and the first MRI acquisition (see 314 Materials and Methods section). At t<sub>2</sub>, modification in the dough structure was clearly 315 observed (Fig. 7.b). Groups of bubbles were growing and were identifiable from  $t_1$  to  $t_3$  and 316 also for longer proving times. The overall networks, i.e. virtual limits between regions of the 317 same radius, seemed to be almost unchanged in the horizontal direction during proving while 318 bubbles underwent overall growth. From Fig. 7 and a complete set of maps of radii, dough 319 evolution seemed to be related to the starting structure of the dough. Image processing 320 methods on maps of radii would be useful for tracking purposes.

### 321 **5. Discussion**

In this section, we first discuss the simulation of  $T_{2\chi}$  and the behavior of the signal. We then explain the limitations of the method. Finally, results on non-yeasted dough and dough during proving will be analyzed in the light of this previous discussion.

#### 325 5.1. Simulated evolution of $T_{2\chi}$

Simulation results concerning  $T_{2\chi}$  as a function of bubble radius (Fig. 2) were original and to our knowledge have never been reported in the literature. Several studies have been undertaken to study NMR signal dephasing due to the presence of mesoscopic field inhomogeneities in the static dephasing regime [25, 26]. The relaxation constant was thus found to be independent of cylinder or bubble radii. From the free induction decay due to randomly distributed spherical particles, Yablonskiy [33] derived an expression for the relaxation rate  $R_2^{\#}$  defined in section 2.1:

333 
$$\mathbf{R}_{2}^{\#} = \frac{4\pi}{9\sqrt{3}} \varepsilon.\gamma.\Delta_{\chi}.B_{0} \tag{9}$$

334

Eq. (9) explains  $T_{2\chi}$  shortening with increasing porosity (Fig. 3.a) or main magnetic 335 field (Fig. 2.b), and the values are compatible with our present study. Indeed, if  $\varepsilon = 12.6\%$ , Eq. (9) resulted in  $T_2^* = 1 / (R_2 + R_2^{\#}) = 12.2$  ms and  $T_2 \chi = 13$  ms in our simulation. 336

337 We demonstrated three levels for  $T_{2\chi}$  observation. First, the plateau in Fig. 2.a indicated 338 that, for a given porosity (or density),  $T_{2\chi}$  values are independent of bubble radii. This confirmed all the results encountered in mesoscopic scale studies. We showed that  $T_{2\chi}$  was 339 340 dependent on radius in a certain range (macroscopic scale) and thus constitutes a sensitive 341 indicator of the medium alveolar structure. The second plateau occurred for large bubbles 342 compared to voxel size.  $T_{2\chi}$  was higher for larger bubbles because the signal from an air voxel 343 could not decrease with TE due to the almost total absence of signal. In fact, heterogeneity 344 was high in the ROI corresponding to the second plateau with complete air voxels, whereas 345 for the first plateau each voxel contained small air bubbles. There were therefore more susceptibility interfaces for small bubbles, thus explaining the lower  $T_{2\chi}$  values. At a constant 346 347 radius, an increase in porosity means a reduction in the center to center distance (ccd) and 348 thus intravoxel dephasing is greater, as described in [27] and in Eq. (1), explaining the 349 decrease of  $T_{2\chi}$  with porosity.

350 Simulations at 1.5 T (Fig. 2.b) showed the potential of this method since there was a greater difference between  $T_{2\chi_{init}}$  values for low porosities and an augmentation of the K 351 352 factor. By using Eq. (8) to reveal bubbles, the estimated radius dynamic is thus enlarged and 353 the algorithm would be more stable. Furthermore, a major gain in signal at high field would 354 make better resolution possible. This is valuable for the study of the very small bubbles 355 present in non-yeasted dough. Indeed, Bellido et al. reported a mean value of 100 µm in wheat 356 flour dough [16].

#### 357 5.2. Limitations of the method

Fig 2.a illustrates some drawbacks of the method.  $T_{2\chi}$  is the same for small radii constituting the first plateau. Bubbles within this range cannot be discriminated. Fortunately, the process of dough proving with the growth of bubbles and the increase in local porosity was suited to this sensitive domain. Due to the shape of the  $T_{2\chi}$  curve (Fig. 2.a), computation of r(i,j) according to Eq. (8) worked well in the dynamic area between the two plateaux.

363 Although bread dough structure is highly heterogeneous (gaseous phase and dough 364 films included), the model was quite accurate in terms of signal loss. A more accurate virtual 365 description of the object will undoubtedly make possible the calculation of values of greater 366 accuracy. In fact, the expansion of bubbles rapidly becomes heterogeneous and anisotropic 367 (non-spherical and distorted bubbles) due to mechanisms of coalescence [23]. However, using 368 the estimation algorithm, it is possible to distinguish between bubbles according to their 369 radius. At this point, r(i,j) values have to be considered as variables related to length rather 370 than a physical variable. Babin et al. [22] found dough structures were heterogeneous and 371 dependent on the recipe, with bubble radii in the range 0.05-0.3 mm at the very beginning of 372 dough proving. Maps of the radii in the present study were thus thought to be overestimated.

#### 373 5.3. Non-yeasted dough and dough during proving

374 There is a real value in combining SE and GE sequences to assess local porosity and 375 signal loss and thus deduce bubble radii. A porosity map (Fig. 4) provides important 376 information (density is a criterion used by bakers to distinguish between bread recipes or to 377 evaluate defects originating from the flour or the different stages in bread making) but a map 378 of radii (Fig. 5) supplements the scientific understanding of evolution of the dough structure, 379 from small to large bubbles. Additionally, the field of view available with MRI offers the 380 possibility of assessing macroscopic heterogeneities (scale of a few millimeters to 381 centimeters) in samples of realistic size comparable with industrial practices. Several image 382 processing techniques based on this map could be developed to classify or segment bubbles.
383 The histogram presented in Fig. 6 shows a Gaussian distribution of bubbles. Bellido et al. [16]
384 used microcomputed tomography on dough and found a log-normal distribution. These first
385 results were encouraging for the study of non-yeasted bread dough.

Bikard et al. [23] analyzed the foaming phase during proving using 3D simulation. They characterized the influence of parameters such as dough viscosity, kinetics or initial number of bubbles on the evolution of average bubble radius over time. Fig. 9 obtained from maps of radii at different times of proving is in good agreement with their results while looking at curve shapes. From a quantitative point of view, it was hard to compare as values were highly dependent on the above parameters that we could not quantify in our experiments.

Using digital image analysis, Rouille et al. [34] showed that reduction in the number of small bubbles ( $\emptyset < 1$  mm) was proportionally balanced by an increase in the number of larger bubbles ( $1 \text{ mm} < \emptyset < 2 \text{ mm}$ ) during proving. This is also obvious from Figs. 7-8.

Image analysis-based methods [17, 20] showed limitations due to small sample size (magnetic resonance microscopy) and long imaging times. Even for small pixel size (115  $\mu$ m), resolution is not sufficient to detect small bubbles at the beginning of the proving time [34]. Susceptibility-based methods are able to reveal information at a lower scale due to the expansion of field inhomogeneities and associated signal loss.

### 400 **6.** Conclusion

Because they arise from object-dependent field inhomogeneities, magnetic susceptibility effects can be a source of quantitative data to characterize alveolar products such as bread dough. While classic image analysis techniques are limited by image or temporal resolution, this new method combines local porosity, local signal loss and simulation predictions to assess bubble radii which can be smaller or the same order of magnitude than the voxel. However, fully quantitative accuracy of the method would suppose comparisons between 407 different techniques or the use of a test object with known geometry and susceptibility. 408 Moreover, further simulation studies with more realistic object geometries to mimic bread 409 dough would make possible more quantitative results. Growth and distribution of bubbles 410 were observed during proving. Maps of radii gave information about dough structure and 411 evolution. The principle of the estimation algorithm was shown at 0.2 T and simulations 412 predicted greater distinction between bubble sizes at 1.5 T, especially for low porosities 413 which are also encountered at key stages in breadmaking.

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- 510 Figure legends
- 511

512 Figure 1: Overview of simulation framework. Field inhomogeneities map was dependent on object parameters (r, ccd,  $\Delta \chi$ , Nb) and was an input of MRI simulator SIMRI.  $T_{2\chi}$  was 513 514 determined for each object from ROI in GE simulated images at different TE from 4 to 12 ms. 515 **Figure 2**:  $T_{2\chi}$  value as a function of bubble radii at different porosities with sigmoid fitting 516 curves. a) At  $B_0 = 0.2$  T. b) At  $B_0 = 1.5$  T. Three stages were identified: an initial plateau where  $T_{2\chi}$  was only dependent on porosity (density), a rapid increase in  $T_{2\chi}$  value and a 517 518 second plateau corresponding to high air proportion in voxels. 519 **Figure 3**: a) *K* and  $T_{2\chi_{init}}$  as functions of porosity. b)  $r_{break}$  and  $\lambda$  as functions of porosity. 520 Second order polynomial relationship with porosity was found for each simulation parameter. 521 Figure 4: Porosity map of non-yeasted bread dough. Large bubbles can be distinguished as 522 high-value pixels. The squared reference ROI is represented. 523 Figure 5: Map of estimated radii of bubbles in non-yeasted bread dough. Contrast 524 enhancement between the different regions of bubble sizes highlighting dough structure with 525 small bubbles surrounded by larger bubbles, etc. Figure 6: Histogram with Gaussian distribution ( $\sigma = 0.062$ ,  $\mu = 0.37$ ,  $R^2 = 0.991$ ) of the non-526 527 yeasted dough computed on the reference ROI with a 0.02 pix<sub>ima</sub> interval.

528 **Figure 7**: Map of estimated radii of bubbles in bread dough during proving. **a**) At  $t = 7 \min b$ )

- 529 At t = 12.5 min c) At t = 18 min showing distribution and growth of bubbles.
- Figure 8: Histograms extracted from maps of radii showing distribution of bubble sizes atthree proving times.
- 532 **Figure 9:** Average estimated radius as a function of proving time.



534 Fig.1



 $T_2\chi$  function of bubbles radius at 0.2 T

535

536 Fig2a























548 Fig.6







552 Fig.8

