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HITTING DENSITIES FOR SPECTRALLY POSITIVE STABLE PROCESSES

THOMAS SIMON

ABSTRACT. A multiplicative identity in law connecting the hitting times of completely asymmetric α -stable Lévy processes in duality is established. In the spectrally positive case, this identity allows with an elementary argument to compute fractional moments and to get series representations for the density. We also prove that the hitting times are unimodal as soon as $\alpha \leq 3/2$. Analogous results are obtained, in a much simplified manner, for the first passage time across a positive level.

1. A MULTIPLICATIVE IDENTITY IN LAW

Let $\{X_t, t \geq 0\}$ be a spectrally positive Lévy α -stable process ($1 < \alpha < 2$), starting from zero and normalized such that

$$(1.1) \quad \mathbb{E} [e^{-\lambda X_t}] = e^{t\lambda^\alpha}, \quad t, \lambda \geq 0.$$

For every $x > 0$, consider the first exit times

$$T_x = \inf\{t > 0, X_t > x\}, \quad \hat{T}_x = \inf\{t > 0, X_t < -x\}$$

and the first hitting times

$$\tau_x = \inf\{t > 0, X_t = x\}, \quad \hat{\tau}_x = \inf\{t > 0, X_t = -x\}.$$

Notice that since X has no negative jumps, one has $\hat{\tau}_x = \hat{T}_x$ a.s. and $\hat{\tau}_x$ is a positive stable random variable with index $1/\alpha$ - see e.g. Theorem 46.3 in [14]. Besides, (1.1) and the optional sampling theorem give the normalization

$$\mathbb{E} [e^{-\lambda \hat{\tau}_x}] = e^{-x\lambda^{1/\alpha}}, \quad x, \lambda \geq 0.$$

On the other hand, the process X crosses the level $x > 0$ by a jump so that $X_{T_x} > x$ and $\tau_x > T_x$ a.s. Introducing the overshoot $K_x = X_{T_x} - x$, the strong Markov property at time T_x entails that $\{X_t^x = X_{T_x+t} - X_{T_x}, t \geq 0\}$ is independent of (T_x, K_x) . In particular, setting $\hat{\tau}_y^x = \inf\{t > 0, X_t^x = -y\}$ for $y > 0$, we have

$$\tau_x = T_x + \hat{\tau}_{K_x}^x \stackrel{d}{=} T_x + K_x^\alpha \hat{\tau}_1 \stackrel{d}{=} x^\alpha (T_1 + K_1^\alpha \hat{\tau}_1)$$

where $\hat{\tau}_1$ is an independent copy of $\hat{\tau}_1$ and the two identities in law follow at once from the self-similarity relationships

$$(T_x, \tau_x, \hat{\tau}_x) \stackrel{d}{=} x^\alpha (T_1, \tau_1, \hat{\tau}_1) \quad \text{and} \quad K_x \stackrel{d}{=} xK_1,$$

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which are themselves simple consequences of the self-similarity of X with index $1/\alpha$. Specifying the above to $x = 1$ yields the basic identity

$$(1.2) \quad \tau_1 \stackrel{d}{=} T_1 + K_1^\alpha \hat{\tau}_1,$$

which has been known for a long time [6]. Notice that the laws of the three random variables appearing in the right-hand side of (1.2) are more or less explicit: the density of $\hat{\tau}_1$ is that of a positive $1/\alpha$ -stable variable, the density of K_1^α comes from a classical computation by Port involving some Beta integral - see e.g. Exercise VIII.3 in [2] - and the density of T_1 can be obtained in changing the variable in the main result of [1] thanks to the identity $T_1 = S_1^{-\alpha}$, with the notation $S_t = \sup\{X_s, s \leq t\}$ for the supremum process. However, lack of explicit information on the law of the bivariate random variable (T_1, K_1) prevents from applying (1.2) directly to derive an expression of the law of τ_1 . Recently Peskir [12] circumvented this difficulty with an identity of the Chapman-Kolmogorov type linking the laws of X_1, S_1 and τ_1 , and providing an explicit series representation at $+\infty$ for the density of τ_1 . In this paper we will follow the method of our previous article [15] in order to derive, firstly, a simple identity in law for τ_1 . Our main result reads:

Theorem. *One has*

$$(1.3) \quad \tau_1 \stackrel{d}{=} U_\alpha \times \hat{\tau}_1,$$

where U is an independent, positive random variable with density

$$f_{U_\alpha}(t) = \frac{-(\sin \pi \alpha)t^{1/\alpha}}{\pi(t^2 - 2t \cos \pi \alpha + 1)}.$$

Proof: From (1.2) we have for any $\lambda \geq 0$

$$\begin{aligned} \mathbb{E} [e^{-\lambda^\alpha \tau_1}] &= \mathbb{E} [e^{-\lambda^\alpha (T_1 + K_1^\alpha \hat{\tau}_1)}] = \mathbb{E} [e^{-\lambda^\alpha T_1} \mathbb{E} [e^{-\lambda^\alpha K_1^\alpha \hat{\tau}_1} | (T_1, K_1)]] \\ &= \mathbb{E} [e^{-(\lambda^\alpha T_1 + \lambda K_1)}] = \mathbb{E} [e^{-(T_\lambda + K_\lambda)}], \end{aligned}$$

where in the third equality we recalled that $\hat{\tau}_1$ is a $(1/\alpha)$ -stable positive variable with given normalization. On the other hand, the double Laplace transform of (T_1, K_1) had been evaluated long ago by Fristedt, and yields the simple expression

$$\int_0^\infty e^{-s\lambda} \mathbb{E} [e^{-(T_\lambda + K_\lambda)}] d\lambda = \left(\frac{1}{s-1} \right) - \left(\frac{\alpha}{s^\alpha - 1} \right)$$

for every $s > 1$ - see (2.16) in [6]. It is possible to invert the right-hand side in noticing that

$$\frac{1}{s-1} = \int_0^\infty e^{-s\lambda} e^{-\lambda} d\lambda \quad \text{and} \quad \frac{1}{s^\alpha - 1} = \int_0^\infty e^{-s\lambda} (\alpha \lambda^{\alpha-1} E'_\alpha(\lambda^\alpha)) d\lambda$$

for every $s > 1$, where E'_α is the derivative of the Mittag-Leffler function

$$E_\alpha(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(1 + \alpha n)}, \quad z \in \mathbb{C}.$$

Above, the computation involving E_α had been made originally by Humbert and we refer to the discussion after (5) in [15] for more details. Putting everything together gives an expression of the Laplace transform of τ_1 :

$$(1.4) \quad \mathbb{E} [e^{-\lambda^\alpha \tau_1}] = F_1(\lambda) - \alpha F'_\alpha(\lambda), \quad \lambda \geq 0,$$

where we used the notation $F_\alpha(x) = E_\alpha(x^\alpha)$ (notice that $F_1(x) = e^x$) for every $x \geq 0$. We will now prove that the function $G_\alpha : \lambda \mapsto F_1(\lambda) - \alpha F'_\alpha(\lambda)$ is completely monotone, reasoning as in Theorem 1 in [15]. The well-known curvilinear representation of the analytic function $1/\Gamma$ yields

$$F_1(\lambda) = \frac{1}{2\pi i} \int_{H_\lambda} \frac{e^t}{t - \lambda} dt \quad \text{and} \quad \alpha F'_\alpha(\lambda) = \frac{1}{2\pi i} \int_{H_\lambda} \frac{\alpha \lambda^{\alpha-1} e^t}{t^\alpha - \lambda^\alpha} dt$$

for every $\lambda > 0$, where H_λ is a Hankel path encircling the disk centered at the origin with radius λ - see e.g. Chapter IX.4 in [4]. We deduce

$$G_\alpha(\lambda) = \frac{1}{2\pi i} \int_{H_\lambda} e^t g_\lambda(t) dt, \quad \lambda > 0$$

where

$$g_\lambda(t) = \left(\frac{1}{t - \lambda} \right) - \left(\frac{\alpha \lambda^{\alpha-1}}{t^\alpha - \lambda^\alpha} \right)$$

is analytic on $\mathbb{C}/\{\lambda\} \cup (-\infty, 0]$. Besides, since $g_\lambda(t) \rightarrow (\alpha - 1)/2\lambda$ as $t \rightarrow \lambda$, it can be continued on the whole complex plane cut on the negative real axis, which entails

$$G_\alpha(\lambda) = \frac{1}{2\pi i} \int_H e^t g_\lambda(t) dt, \quad \lambda > 0$$

where H is a Hankel path independent of λ . We can then compute, as in Theorem 1 in [15],

$$\begin{aligned} G_\alpha(\lambda) &= \lim_{\rho \rightarrow 0} \frac{1}{2\pi i} \left(\int_\rho^\infty e^{-s} \left(\frac{1}{s + \lambda} + \frac{\alpha \lambda^{\alpha-1}}{e^{i\pi\alpha} s^\alpha - \lambda^\alpha} \right) - e^{-s} \left(\frac{1}{s + \lambda} + \frac{\alpha \lambda^{\alpha-1}}{e^{-i\pi\alpha} s^\alpha - \lambda^\alpha} \right) ds \right) \\ &= \frac{-\alpha \sin \pi\alpha}{\pi} \int_0^\infty e^{-s} \left(\frac{\lambda^{\alpha-1} s^\alpha}{s^{2\alpha} - 2s^\alpha \lambda^\alpha \cos \pi\alpha + \lambda^{2\alpha}} \right) ds \\ &= \frac{-\alpha \sin \pi\alpha}{\pi} \int_0^\infty e^{-\lambda s} \left(\frac{s^\alpha}{s^{2\alpha} - 2s^\alpha \cos \pi\alpha + 1} \right) ds, \end{aligned}$$

so that from (1.4) and the expression of the density f_{U_α} ,

$$\mathbb{E} [e^{-\lambda^\alpha \tau_1}] = \mathbb{E} [e^{-\lambda U_\alpha^{1/\alpha}}]$$

for all $\lambda \geq 0$. Reasoning now exactly as in Theorem 3 in [15] entails finally the desired identity in law $\tau_1 \stackrel{d}{=} U_\alpha \times \hat{\tau}_1$. \square

Remark. In Theorem 1 of [6], the double Laplace transform of the function $(t, x) \mapsto \mathbb{P}[\tau_x \geq t]$ was computed for general Lévy processes with no negative jumps. Our result can be viewed as an inversion of this Laplace transform in the particular stable case.

2. FRACTIONAL MOMENTS AND SERIES REPRESENTATIONS

In this section we will obtain two expressions for the density of τ_1 (a density which is readily known to exist from our main result, see also [10] for the general context) in terms of convergent series, the first one being useful in the neighbourhood of 0 and the second in the neighbourhood of $+\infty$. Finding different asymptotic expansions for densities of first passage times is a classical issue in probability, see Chapter 1 in [8]. Recall also that this work has been done long ago for $\hat{\tau}_1$, respectively by Pollard and Linnik - see (14.31) and (14.35) in [14]. Our approach is different from the above references and relies on Mellin inversion. We will hence first compute the fractional moments of τ_1 , as a simple consequence of our main result:

Corollary 1. *For any $s \in (-1 - 1/\alpha, 1 - 1/\alpha)$, one has*

$$\mathbb{E}[\tau_1^s] = \frac{\Gamma(1 - \alpha s) \sin(\pi(\alpha - 1)(s + 1/\alpha))}{\Gamma(1 - s) \sin(\pi(s + 1/\alpha))}.$$

Proof: From (1.3) we have $\mathbb{E}[\tau_1^s] = \mathbb{E}[\hat{\tau}_1^s] \mathbb{E}[U_\alpha^s]$ for every $s \in \mathbb{R}$. The computation of $\mathbb{E}[\hat{\tau}_1^s]$ is easy and classical:

$$\mathbb{E}[\hat{\tau}_1^s] = \frac{1}{\Gamma(-s)} \int_0^\infty \lambda^{-(s+1)} \mathbb{E}[e^{-\lambda \hat{\tau}_1}] d\lambda = \frac{1}{\Gamma(-s)} \int_0^\infty \lambda^{-(s+1)} e^{-\lambda^{1/\alpha}} d\lambda = \frac{\Gamma(1 - \alpha s)}{\Gamma(1 - s)}$$

for every $s < 1/\alpha$. Similarly, one can show that

$$\mathbb{E}[U_\alpha^s] = \frac{1}{\Gamma(-\alpha s)} \int_0^\infty \lambda^{-(\alpha s+1)} G_\alpha(\lambda) d\lambda = \frac{1}{\Gamma(-\alpha s)} \int_0^\infty \lambda^{-(\alpha s+1)} (e^\lambda - \alpha \lambda^{\alpha-1} E'_\alpha(\lambda^\alpha)) d\lambda$$

but I could not carry the computations further with the sole help of Mittag-Leffler functions. Instead, one can rest upon the residue theorem in order to evaluate directly

$$\mathbb{E}[U_\alpha^s] = \frac{-\sin \pi \alpha}{\pi} \int_0^\infty \frac{x^{s+1/\alpha}}{x^2 - 2x \cos \pi \alpha + 1} dx.$$

Setting $\lambda = s + 1/\alpha + 1 \in (0, 2)$ and

$$f(z) = \frac{(-z)^{\lambda-1}}{z^2 - 2z \cos \pi \alpha + 1}$$

one sees from Formula VIII.(10.2.3) in [4] that

$$\mathbb{E}[U_\alpha^s] = - \left(\frac{\sin \pi \alpha}{\sin \pi \lambda} \right) (\text{Res}_{x_\alpha} f + \text{Res}_{\bar{x}_\alpha} f)$$

where $x_\alpha = e^{i(2-\alpha)\pi}$ and $\bar{x}_\alpha = e^{i(\alpha-2)\pi}$ are the two conjugate roots of $x^2 - 2x \cos \pi \alpha + 1$. After some simple computations, one gets

$$\mathbb{E}[U_\alpha^s] = - \left(\frac{\sin(\pi(\lambda - 1)(\alpha - 1))}{\sin \pi \lambda} \right) = \frac{\sin(\pi(\lambda - 1)(\alpha - 1))}{\sin \pi(\lambda - 1)} = \frac{\sin(\pi(\alpha - 1)(s + 1/\alpha))}{\sin(\pi(s + 1/\alpha))}$$

for every $s \in (-1 - 1/\alpha, 1 - 1/\alpha)$, which together with the previous computation for $\hat{\tau}_1$ completes the proof (notice that $1 - 1/\alpha < 1/\alpha$). \square

Remark. The fractional moments of U_α and τ_1 are finite over the same strip $(-1 - 1/\alpha, 1 - 1/\alpha)$, whereas those of $\hat{\tau}_1$ are finite over the larger strip $(-\infty, 1/\alpha)$.

We will now derive series representations for the density f_{τ_1} of τ_1 , inverting the fractional moments which were computed just before. Setting $M(z) = \mathbb{E}[\tau_1^z]$ for every $z \in \mathbb{C}$ such that $\operatorname{Re}(z) \in (-1 - 1/\alpha, 1 - 1/\alpha)$, one has

$$M(i\lambda) = \mathbb{E}[e^{i\lambda\sigma_1}], \quad \lambda \in \mathbb{R},$$

where $\sigma_1 = \log \tau_1$ is a real random variable with finite polynomial moments by Corollary 1, and which is absolutely continuous with continuous density $f_{\sigma_1}(x) = e^x f_{\tau_1}(e^x)$ over \mathbb{R} . The Fourier inversion formula for f_{σ_1} entails then

$$(2.1) \quad f_{\tau_1}(x) = \frac{1}{2\pi x} \int_{\mathbb{R}} M(it) x^{-it} dt, \quad x \geq 0,$$

and we will start from this formula to obtain our series representation. The one at zero is particularly simple:

Proposition 2. *For every $x \geq 0$ one has*

$$f_{\tau_1}(x) = \sum_{n \geq 1} \frac{\alpha x^{1/\alpha+n-1}}{\Gamma(-\alpha n) \Gamma(1/\alpha + n)}.$$

Proof: Suppose first $x < 1$. We compute the integral in (2.1) with the help of the usual contour Γ_R joining $-R$ to R along the real axis and R to $-R$ along a half-circle in the trigonometric orientation, centered at the origin. Because $x < 1$, the integral along the half-circle is easily seen to converge to 0 when $R \rightarrow \infty$, so that it remains to consider the singularities of $L(t) = M(it)x^{-it}$ inside the contour. By Corollary 1, the latter are located at $i(n + 1/\alpha)$, $n \geq 0$. After standard computations, one finds that

$$\operatorname{Res}_{i(n+1/\alpha)} L = -i \left(\frac{\Gamma(2 + \alpha n) \sin(n\pi\alpha) x^{n+1/\alpha}}{\pi \Gamma(1 + 1/\alpha + n)} \right) = -i \left(\frac{\alpha x^{n+1/\alpha}}{\Gamma(-\alpha n) \Gamma(1/\alpha + n)} \right)$$

for any $n \geq 0$. By the residue theorem, this finishes the proof for $x \in (0, 1)$ (notice that $\operatorname{Res}_{i/\alpha} L = 0$) and by analyticity, one can then extend the formula over the whole \mathbb{R}^+ . \square

The series representation at infinity is slightly more complicated than the one at zero, involving actually two series. Up to some painless normalization and after some simplifications on the Gamma function, it had already been computed by Peskir [12] with an entirely different, probabilistic argument involving a previous computation made for f_{S_1} in [1].

Proposition 3. *For every $x > 0$ one has*

$$f_{\tau_1}(x) = \sum_{n \geq 1} \frac{-\alpha x^{1/\alpha-n-1}}{\Gamma(\alpha n) \Gamma(1/\alpha - n)} + \sum_{n \geq 1} \frac{x^{-n/\alpha-1}}{\Gamma(-n/\alpha) n!}.$$

Proof: By analyticity, it suffices to consider the case $x > 1$. We compute the integral in (2.1) with a contour $\bar{\Gamma}_R$ which is the reflection of Γ_R with respect to the real axis. Because $x > 1$, the integral along the half-circle converges to 0 when $R \rightarrow \infty$, so that one needs again to

consider the singularities of L inside $\bar{\Gamma}_R$, which are located at $i(1/\alpha - n)$ and $-in/\alpha$, $n \geq 0$. The term

$$\sum_{n \geq 1} \frac{-\alpha x^{1/\alpha - n - 1}}{\Gamma(\alpha n) \Gamma(1/\alpha - n)}$$

follows then from the computations of Proposition 2 in replacing n by $-n$ and taking into account the clockwise orientation of $\bar{\Gamma}_R$. The last term follows from the simple computation

$$\text{Res}_{i(n+1/\alpha)} L = \frac{ix^{-n/\alpha}}{\Gamma(-n/\alpha)n!}.$$

□

Remark. The two above representations are integrable term by term, which yields two series representations for the distribution function of τ_1 :

$$\mathbb{P}[\tau_1 \leq x] = \sum_{n \geq 1} \frac{\alpha x^{1/\alpha + n}}{\Gamma(-\alpha n) \Gamma(1 + 1/\alpha + n)}$$

and

$$\mathbb{P}[\tau_1 \geq x] = \sum_{n \geq 1} \frac{\alpha x^{1/\alpha - n}}{\Gamma(\alpha n) \Gamma(1 + 1/\alpha - n)} - \sum_{n \geq 1} \frac{x^{-n/\alpha}}{\Gamma(1 - n/\alpha)n!}.$$

One can also differentiate term by term in order to get expansions for the successive derivatives of f_{τ_1} : for any $p \geq 1$ one has

$$f_{\tau_1}^{(p)}(x) = \sum_{n \geq 1} \frac{\alpha x^{1/\alpha + n - 1 - p}}{\Gamma(-\alpha n) \Gamma(1/\alpha + n - p)}$$

and

$$f_{\tau_1}^{(p)}(x) = \sum_{n \geq 1} \frac{-\alpha x^{1/\alpha - n - 1 - p}}{\Gamma(\alpha n) \Gamma(1/\alpha - n - p)} + \sum_{n \geq 1} \frac{x^{-n/\alpha - 1 - p}}{\Gamma(-n/\alpha - p)n!}.$$

The first term of all these expansions give the behaviour of the function at zero respectively at infinity. For instance, one has

$$f_{\tau_1}(x) \sim \frac{\alpha x^{1/\alpha}}{\Gamma(-\alpha) \Gamma(1 + 1/\alpha)} \quad \text{and} \quad f'_{\tau_1}(x) \sim \frac{\alpha x^{1/\alpha - 1}}{\Gamma(-\alpha) \Gamma(1/\alpha)}$$

as $x \rightarrow 0^+$, whereas

$$f_{\tau_1}(x) \sim \frac{x^{1/\alpha - 2}}{\Gamma(\alpha - 1) \Gamma(1/\alpha)} \quad \text{and} \quad f'_{\tau_1}(x) \sim \frac{-\alpha x^{1/\alpha - 3}}{\Gamma(\alpha) \Gamma(1/\alpha - 2)}$$

as $x \rightarrow +\infty$. Notice finally that since $\alpha \in (1, 2)$, the function f_{τ_1} is ultimately completely monotone at infinity and f'_{τ_1} ultimately completely monotone at zero.

3. COMPLEMENTS ON THE FIRST EXIT TIME AND THE RUNNING SUPREMUM

In this section, we will derive similar series representations for the density f_{T_1} of the first exit time T_1 . As mentioned before, this random variable is connected to the supremum process $\{S_t, t \geq 0\}$, so that we will also get series representations for S_1 . Our basic tool is an identity in law analogous to our main result, which was proved in our previous article [15], Theorem 3. This identity reads

$$(3.1) \quad T_1 \stackrel{d}{=} T_\alpha \times \hat{\tau}_1$$

where is an independent random variable with density given by

$$f_{T_\alpha}(t) = \frac{-(\sin \pi \alpha)(1 + t^{1/\alpha})}{\pi \alpha (t^2 - 2t \cos \pi \alpha + 1)}$$

over \mathbb{R}^+ . As before, we begin in evaluating the fractional moments of T_1 :

Proposition 4. *For every $s \in (-1, 1 - 1/\alpha)$ one has*

$$\mathbb{E}[T_1^s] = \frac{\Gamma(1 + s) \sin(\pi/\alpha)}{\Gamma(1 + s\alpha) \sin(\pi(s + 1/\alpha))}.$$

Proof. For every $s \in (-1, 1 - 1/\alpha)$ we have

$$\begin{aligned} \mathbb{E}[T_\alpha^s] &= \frac{1}{\alpha} (\mathbb{E}[U_\alpha^{s-1/\alpha}] + \mathbb{E}[U_\alpha^s]) \\ &= \frac{1}{\alpha} \left(\frac{\sin(\pi s(\alpha - 1))}{\sin \pi s} - \frac{\sin(\pi s(\alpha - 1) - 1/\alpha)}{\sin(\pi(s + 1/\alpha))} \right) \\ &= \frac{\sin(\pi s \alpha) \sin(\pi/\alpha)}{\alpha \sin(\pi s) \sin(\pi(s + 1/\alpha))} \end{aligned}$$

where the last equality follows from tedious trigonometric transformations. From (3.1), the previous computation made for $\mathbb{E}[\hat{\tau}_1^s]$, and the complement formula for the Gamma function, we get

$$\mathbb{E}[T_1^s] = \frac{\sin(\pi/\alpha)}{\sin(\pi(s + 1/\alpha))} \left(\frac{\Gamma(1 - s\alpha) \sin(\pi s \alpha)}{\alpha \Gamma(1 - s) \sin(\pi s)} \right) = \frac{\Gamma(1 + s) \sin(\pi/\alpha)}{\Gamma(1 + s\alpha) \sin(\pi(s + 1/\alpha))}$$

as desired. □

Remark. Similarly as above, the fractional moments of T_α and τ_1 are finite over the same strip $(-1, 1 - 1/\alpha)$, whereas those of $\hat{\tau}_1$ are finite over the larger strip $(-\infty, 1/\alpha)$.

We notice in passing that this computation allows to show that the law of the independent quotient T_1/\hat{T}_1 is Pareto($1 - 1/\alpha$), a fact which had been proved in [5] in the general context of stable processes in duality:

Corollary 5. *The density of the independent quotient T_1/\hat{T}_1 is given by*

$$\frac{\sin(\pi/\alpha)}{\pi t^{1-1/\alpha}(1+t)}$$

over \mathbb{R}^+ .

Proof: By Mellin inversion, it is enough to identify the fractional moments. Proposition 4 and the classical computation for the moments of \hat{T}_1 entail

$$\mathbb{E} \left[\frac{T_1^s}{\hat{T}_1^s} \right] = \frac{\sin(\pi/\alpha)}{\sin(\pi(s + 1/\alpha))} = \int_0^\infty \frac{\sin(\pi/\alpha)t^s}{\pi t^{1-1/\alpha}(1+t)} dt$$

for every $s \in (-1/\alpha, 1 - 1/\alpha)$, where the second equality follows from Formula VIII.(10.2.4) in [4]. \square

We now come back to series representations for f_{T_1} , which are obtained from Proposition 4 and Mellin inversion similarly as in Propositions 2 and 3. Since we choose exactly the same contours, we will leave all the details to the reader. Notice that contrary to τ_1 , the two-series representation for T_1 is the one at zero and that the first series therein defines an entire function.

Corollary 6. *One has both representations*

$$f_{T_1}(x) = \sum_{n \geq 1} \frac{x^{1/\alpha - n - 1}}{\alpha \Gamma(\alpha n - 1) \Gamma(1 + 1/\alpha - n)}, \quad x > 0$$

and

$$f_{T_1}(x) = \sum_{n \geq 1} \frac{x^{n-1}}{\alpha \Gamma(-\alpha n)} + \sum_{n \geq 1} \frac{x^{1/\alpha + n - 1}}{\Gamma(-\alpha n) \Gamma(1/\alpha + n)}, \quad x \geq 0.$$

Thanks to the aforementioned identity $T_1 \stackrel{d}{=} S_1^{-\alpha}$, changing the variable gives corresponding series representations for the density $f_{S_1} = \alpha x^{-(\alpha+1)} f_{T_1}(x^{-\alpha})$. The first one had been obtained originally in [1] after solving some fractional integral equation of the Abel type. See also [11] for a new proof involving Wright's hypergeometric function. We stress that up to the Wiener-Hopf factorisation which is a fundamental tool for this type of questions, overall our argument to get this series only makes use of undergraduate mathematics. Besides, this simple method allows to get the series representation at infinity.

Corollary 7. *One has both representations*

$$f_{S_1}(x) = \sum_{n \geq 1} \frac{x^{\alpha n - 2}}{\Gamma(\alpha n - 1) \Gamma(1 + 1/\alpha - n)}, \quad x \geq 0$$

and

$$f_{S_1}(x) = \sum_{n \geq 1} \frac{x^{-(\alpha n + 1)}}{\Gamma(-\alpha n)} + \sum_{n \geq 1} \frac{\alpha x^{-(\alpha n + 2)}}{\Gamma(-\alpha n) \Gamma(1/\alpha + n)}, \quad x > 0.$$

Remark. Notice that one really needs the two series to get the behaviour of f_{T_1} and its derivatives at zero. For instance one has

$$f_{T_1}(x) \sim \frac{1}{\alpha \Gamma(-\alpha)} \quad \text{and} \quad f'_{T_1}(x) \sim \frac{x^{1/\alpha - 1}}{\Gamma(-\alpha) \Gamma(1/\alpha)} \sim \frac{1}{\alpha} f'_{\tau_1}(x)$$

as $x \rightarrow 0^+$. On the other hand, only one series is important to get the behaviour of f_{S_1} and its derivatives at infinity: one has

$$f_{S_1}(x) \sim \frac{x^{-(\alpha+1)}}{\Gamma(-\alpha)} \quad \text{and} \quad f_{S_1}^{(p)}(x) \sim \frac{x^{-(\alpha+2+p)}}{\Gamma(-(\alpha+1+p))}$$

as $x \rightarrow +\infty$ for every $p \geq 1$ (this is in accordance with the general asymptotics derived in [7], see the final remark therein).

4. SOME REMARKS ON UNIMODALITY

Recall that a real random variable X is said to be unimodal if there exists $a \in \mathbb{R}$ such that the functions $\mathbb{P}[X \leq x]$ and $\mathbb{P}[X > x]$ are convex respectively in $(-\infty, a)$ and $(a, +\infty)$. If X has a density f_X , this means that f_X increases on $(-\infty, a]$ and decreases on $[a, +\infty)$. We refer e.g. to Section 52 in [14] for more on this topic. Differentiating the density of U_α :

$$f'_{U_\alpha}(t) = \frac{t^{1/\alpha-1}}{\alpha(t^2 - 2t \cos \pi\alpha + 1)^2} ((1 - 2\alpha)t^2 + 2t(\alpha - 1) \cos \pi\alpha + 1),$$

we find that there is a unique positive root, so that the random variable U_α is unimodal. Since the same property is known to hold for $\hat{\tau}_1$ - see e.g. Theorem 53.1 in [14] for a much stronger result, in view of our main result it is natural to ask whether the variable τ_1 is unimodal as well. This seems also plausible from the behaviour of f_{τ_1} and f'_{τ_1} at zero. In general, the product of two positive independent unimodal random variables is not necessarily unimodal. However the notion of multiplicative strong unimodality which had been introduced in [3], will allow us to give a quick positive answer in half of the cases:

Proposition 8. *For every $\alpha \leq 3/2$, the random variable τ_1 is unimodal.*

Proof: Recalling that $\hat{\tau}_1$ is unimodal for every $\alpha \in (1, 2)$, from (1.3) and Theorem 3.6 in [3] it is enough to show that the function $t \mapsto f_{U_\alpha}(e^t)$ is log-concave over \mathbb{R} , in other words that the function $g_\alpha(t) = \log(e^t - 2 \cos \pi\alpha + e^{-t})$ is convex over \mathbb{R} . Differentiating twice yields

$$g''_\alpha(t) = \frac{4(1 - \cos \pi\alpha \cosh t)}{(e^t - 2 \cos \pi\alpha + e^{-t})^2}$$

and we see that g_α is convex over \mathbb{R} as soon as $\alpha \leq 3/2$. □

When $\alpha \in (3/2, 2)$ the above proof shows that U_α is not multiplicative strongly unimodal (MSU) in the terminology of [3]. In order to get the unimodality of τ_1 when $\alpha \in (3/2, 2)$, one could of course be tempted to obtain the MSU property for positive stable laws with index $\beta > 2/3$. Unfortunately, we showed in [16] that this property does not hold true as soon as $\beta > 1/2$. Nevertheless, in view of the unimodality of $\hat{\tau}_1$ and that of τ_1 for $\alpha = 2$ - see also Rösler's result for general diffusions on the line [13], we conjecture that τ_1 is unimodal for every value of α . At first sight the problem does not seem easy because of the unstability of the unimodality property under multiplicative convolution. See however our final remark. Differentiating now the density of T_α :

$$f'_{T_\alpha}(t) = \frac{t^{1/\alpha-1}}{\alpha(t^2 - 2t \cos \pi\alpha + 1)^2} ((1 - 2\alpha)t^2 - 2\alpha t^{2-1/\alpha} + 2t(\alpha - 1) \cos \pi\alpha + 2\alpha t^{1-1/\alpha} \cos \pi\alpha + 1),$$

we find after some simple analysis that there is also a unique positive root, so that the random variable T_α is unimodal. But contrary to U_α , one can check that it is never MSU, so that the above simple argument cannot be applied. Nevertheless we have the

Proposition 9. *For every $\alpha \leq 3/2$, the random variable T_1 is unimodal.*

Proof: From (3.1) and the so-called Kanter representation for positive stable laws - see Corollary 4.1 in [9], we have the identity

$$T_1 \stackrel{d}{=} T_\alpha \times L^{1-\alpha} \times b_\alpha(U)$$

where L is a standard exponential variable, U an independent uniform variable over $[0, \pi]$, and $b_\alpha(u) = (\sin((1 - 1/\alpha)u)/\sin(u))^{\alpha-1} \sin(u/\alpha)/\sin(u)$ is a strictly increasing function from $(0, \pi)$ to \mathbb{R}^+ . Besides, from the beginning of the proof of Theorem 4.1 in [9], we know that the logarithmic derivative of b_α is positive and strictly increasing, so that the same holds for b'_α itself, because b_α is also positive and strictly increasing. By Lemma 4.2 in [9] we conclude that $b_\alpha(U)$ is unimodal: again, from (3.1) and Theorem 3.6 in [3], it suffices to show that the density $t \mapsto f_{T_\alpha \times L^{1-\alpha}}(e^t)$ is log-concave over \mathbb{R} . We compute

$$f_{T_\alpha \times L^{1-\alpha}}(x) = \frac{-(\sin \pi \alpha) x^{\frac{\alpha}{1-\alpha}}}{\pi \alpha (\alpha - 1)} \int_0^\infty \exp - (x/u)^{\frac{1}{1-\alpha}} \frac{u^{\frac{1}{\alpha-1}} (1 + u^{1/\alpha})}{u^2 - 2u \cos \pi \alpha + 1} du, \quad x \geq 0.$$

It is easily seen that $u \mapsto 1 + u^{1/\alpha}$ is log-concave over \mathbb{R}^+ and $(t, u) \mapsto e^t/u$ log-convex, hence convex, over $\mathbb{R} \times \mathbb{R}^+$. In particular, the function

$$(t, u) \mapsto \exp - (e^t/u)^{\frac{1}{1-\alpha}} (1 + u^{1/\alpha})$$

is log-concave over $\mathbb{R} \times \mathbb{R}^+$ and by Prékopa's theorem, it is enough to show that the function $u \mapsto u^{\frac{1}{\alpha-1}}/(u^2 - 2u \cos \pi \alpha + 1)$ is log-concave over \mathbb{R}^+ . We compute its logarithmic derivative, which equals

$$\frac{(2\alpha - 3)u^4 + 4u^3(2 - \alpha) \cos \pi \alpha - (4(2 - \alpha) \cos^2 \pi \alpha + 2\alpha)u^2 + 4u \cos \pi \alpha - 1}{(\alpha - 1)u^2(u^2 - 2u \cos \pi \alpha + 1)^2}$$

and we see that it is negative over \mathbb{R}^+ as soon as $\alpha \leq 3/2$. □

Final Remark. The above proof does not work to obtain the unimodality of T_1 when $\alpha > 3/2$, and we do not know as yet how to tackle this situation. However, we can use the same argument to obtain the unimodality of τ_1 when $3/2 < \alpha \leq 1 + 1/\sqrt{2} < 2$. Details, which are technical, can be obtained upon request. All these questions will be the matter of future research.

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