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WEIGHTED MAXIMAL REGULARITY ESTIMATES AND SOLVABILITY OF NON-SMOOTH ELLIPTIC SYSTEMS I

PASCAL AUSCHER AND ANDREAS AXELSSON

ABSTRACT. We develop new solvability methods for divergence form second order, real and complex, elliptic systems above Lipschitz graphs, with L_2 boundary data. The coefficients A may depend on all variables, but are assumed to be close to coefficients A_0 that are independent of the coordinate transversal to the boundary, in the Carleson sense $||A - A_0||_C$ defined by Dahlberg. We obtain a number of *a priori* estimates and boundary behaviour results under finiteness of $||A - A_0||_C$. Our methods yield full characterization of weak solutions, whose gradients have L_2 estimates of a non-tangential maximal function or of the square function, via an integral representation acting on the conormal gradient, with a singular operatorvalued kernel. Also, the non-tangential maximal function of a weak solution is controlled in L_2 by the square function of its gradient. This estimate is new for systems in such generality, and even for real non-symmetric equations in dimension 3 or higher. The existence of a proof *a priori* to well-posedness, is also a new fact.

As corollaries, we obtain well-posedness of the Dirichlet, Neumann and Dirichlet regularity problems under smallness of $||A - A_0||_C$ and well-posedness for A_0 , improving earlier results for real symmetric equations. Our methods build on an algebraic reduction to a first order system first made for coefficients A_0 by the two authors and A. McIntosh in order to use functional calculus related to the Kato conjecture solution, and the main analytic tool for coefficients A is an operational calculus to prove weighted maximal regularity estimates.

MSC classes: 35J55, 35J25, 42B25, 47N20

Keywords: elliptic systems, maximal regularity, Dirichlet and Neumann problems, square function, non-tangential maximal function, Carleson measure, functional and operational calculus

1. INTRODUCTION

In this article, we present and develop new representations and new solvability methods for boundary value problems (BVPs) for divergence form second order, real and complex, elliptic systems. We look here at BVPs in domains Lipschitz diffeomorphic to the upper half space $\mathbf{R}^{1+n}_{+} := \{(t, x) \in \mathbf{R} \times \mathbf{R}^n : t > 0\}, n \ge 1$. The same problems on bounded domains Lipschitz diffeomorphic to the unit ball, contain noticeable differences, such as use of Fredholm theory, which we address in a forthcoming paper [6]. Here, we focus on the fundamental scale-invariant estimates.

The system of equations is

(1)
$$Lu^{\alpha}(t,x) = \sum_{i,j=0}^{n} \sum_{\beta=1}^{m} \partial_i \left(A_{i,j}^{\alpha,\beta}(t,x) \partial_j u^{\beta}(t,x) \right) = 0, \qquad \alpha = 1, \dots, m$$

in \mathbf{R}^{1+n}_+ , where $\partial_0 = \frac{\partial}{\partial t}$ and $\partial_i = \frac{\partial}{\partial x_i}$, $1 \le i \le n$. We assume

(2)
$$A = (A_{i,j}^{\alpha,\beta}(t,x))_{i,j=0,\dots,n}^{\alpha,\beta=1,\dots,m} \in L_{\infty}(\mathbf{R}^{1+n};\mathcal{L}(\mathbf{C}^{(1+n)m})),$$

and that A is strictly accretive on \mathcal{H} , meaning that there exists $\kappa > 0$ such that

(3)
$$\sum_{i,j=0}^{n} \sum_{\alpha,\beta=1}^{m} \int_{\mathbf{R}^{n}} \operatorname{Re}(A_{i,j}^{\alpha,\beta}(t,x)f_{j}^{\beta}(x)\overline{f_{i}^{\alpha}(x)})dx \ge \kappa \sum_{i=0}^{n} \sum_{\alpha=1}^{m} \int_{\mathbf{R}^{n}} |f_{i}^{\alpha}(x)|^{2}dx,$$

for all $f \in \mathcal{H}$ and a.e. t > 0. The definition of \mathcal{H} , a subspace of $L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$, will be given in Section 2.

We seek to prove well-posedness for (1), *i.e.* unique solvability in appropriate spaces given Dirichlet data $u|_{t=0}$, Neumann data $\partial_{\nu_A} u|_{t=0}$ or Dirichlet regularity data $\nabla_x u|_{t=0}$, assumed to satisfy an L_2 condition. Note that the continuity estimate required for well-posedness in the sense of Hadamard is not included in our notion of well-posedness, but will be shown to hold. For the Neumann and Dirichlet regularity problems, we will work in the class of weak solutions whose gradient $\nabla_{t,x} u$ has L_2 modified non-tangential maximal function $\tilde{N}_*(\nabla_{t,x} u)$ in L_2 . (See Definition 2.1.) Under our assumptions, we shall describe the limiting behaviour of $\nabla_{t,x} u$ at t = 0and ∞ and obtain a perturbation result for well-posedness in this class. For the Dirichlet problem, it is more natural given our method to work in the class of weak solutions with square function estimate $\iint_{\mathbf{R}^{1+n}_+} |\nabla_{t,x} u|^2 t dt dx < \infty$. Under our assumptions, we shall describe the limiting behaviour of u at t = 0 and ∞ and prove a rigidity theorem that shows new a priori non-tangential maximal estimates and L_2 estimates, and obtain a perturbation result for well-posedness.

Let us begin by pointing out that the coefficients depend on t, which makes these problems not always solvable in such generality. In Caffarelli, Fabes and Kenig [15], the necessity of a square Dini condition is pointed out. There has been a wealth of results for real symmetric equations (*i.e.* m = 1 and $A_{ij} = A_{ji} \in \mathbf{R}$, $\mathcal{H} = L_2(\mathbf{R}^n; \mathbf{C}^{1+n}))$. In Fabes, Jerison and Kenig [26], the L_2 Dirichlet problem is solved under the square Dini condition and continuity. Dahlberg removed continuity and proved in [19] that if the discrepancy $A_1 - A_2$ of two matrices A_1 , A_2 satisfies a small Carleson condition, *i.e.* if $||A_1 - A_2||_C$ from Definition 2.1 is small enough, then L_{p_1} -solvability of the Dirichlet problem with coefficients A_1 implies L_{p_2} -solvability of the Dirichlet problem with coefficients A_2 with $p_2 = p_1$. The smallness condition was removed in Fefferman, Kenig and Pipher [28], but then the value of p_2 becomes unspecified. R. Fefferman obtained in [27] the same conclusions as Dahlberg with $p_2 = p_1$, under large perturbation conditions of different nature. See also Lim [41]. Kenig and Pipher [33] proved that the L_p -Neumann and regularity problems are uniquely solvable if the discrepancy A(t, x) - A(0, x) satisfies Dahlberg's small Carleson condition, depending on $p \in (1, 2 + \epsilon)$. Moreover, in [34] they proved small perturbation results for the Neumann and regularity problems analogous the result [19] for the Dirichlet problem, as well as large perturbation results for the regularity problem analogous to [28] for the Dirichlet problem.

Some related results of Kenig and Pipher [35] (going back to questions of Dahlberg [19]), Dindos, Petermichl and Pipher [24] and Dindos and Rule [25] are obtained under smallness of a Carleson condition on $t\nabla_{t,x}A(t,x)$. See also Rios' work [42]. Such an hypothesis does not compare to the one on A(t,x) - A(0,x).

We note that these results are obtained for L_p data, for appropriate p's, including p = 2. This is using all the available technology for *real scalar equations*, starting from the maximum principle, hence *L*-harmonic measure, and Green's functions. Moreover, as far as solvability is concerned, the main thrust of these works is to get

p = 2 with non-tangential maximal estimates, using for this the classical variational solutions, or those obtained via the maximum principle.

Of course, t-dependent coefficients incorporate the t-independent ones. We refer to the book by Kenig [30] and references therein, and to Alfonseca, Auscher, Axelsson, Hofmann and Kim [4] for more recent results on L_{∞} perturbation of real symmetric (or complex and constant) equations. See also below.

We mention a series of works for two dimensional equations on the upper halfplane with t-independent coefficients. Auscher and Tchamitchian [11] study complex coefficients equations with diagonal A (which we call here block form) and describe Dirichlet, regularity and Neumann problems for L_p data for p > 1 and even for data of Hardy type for $p \leq 1$. This is a precursor of the work for systems here, as it built upon new proofs relying on Calderón-Zygmund operators (which are no longer available here) of the one dimensional Kato conjecture proved earlier by Coifman, McIntosh and Meyer [16] and its extension by Kenig and Meyer [32]. For real equations of non block forms, Kenig, Koch, Pipher and Toro [31] show that the Dirichlet problem is well-posed for large enough p (and obtain counterexamples for any given specific p) by showing that L-harmonic measure is absolutely continuous with respect to Lebesgue measure on the boundary¹. Kenig and Rule [36] then obtained well-posedness for the Neumann and regularity problems with p-1 > 0small enough (and obtained counterexamples for any specific p > 1). The recent thesis of Barton [14] deals with complex, t-independent L_{∞} perturbations of the situation in [31, 36], and she obtains well-posedness of the Neumann and regularity problems in L_p for p-1 > 0 small and even at p = 1 with data in the classical Hardy space.

As the reader has observed, we consider complex systems and we wish to obtain L_2 solvability under conditions as general as possible (we mention that L_p solvability with our methods when $p \neq 2$ is under study at this time). For this, we need other tools than those mentioned above. In fact, the tools we develop and that we describe next would not have been conceivable prior to the full solution in all dimensions of the Kato problem and its extensions. In Auscher, Axelsson and McIntosh [9], a new method was presented for solving BVPs with *t*-independent coefficients, following an earlier setup designed in Auscher, Axelsson and Hofmann [7]. The main discovery in [9] is that the equation (1) becomes particularly simple when solving for the conormal gradient defined by

(4)
$$f = \nabla_A u := \begin{bmatrix} \partial_{\nu_A} u \\ \nabla_x u \end{bmatrix},$$

where $\partial_{\nu_A} u$ denotes the (inward for convenience) conormal derivative (see Section 4), instead of the potential u itself. It is a set of generalized Cauchy–Riemann equations expressed as an autonomous first order system

(5)
$$\partial_t f + DBf = 0,$$

where D is a self-adjoint (but not positive) first order differential operator with constant coefficients that is elliptic in some sense and B is multiplication with a bounded matrix B(x), which is strictly accretive on the space \mathcal{H} in (3) and related to A(x) = A(t, x), t > 0, by an explicit algebraic formula. The operator DB is a

¹A recent preprint by Dindos, Kenig and Pipher [23] posted during the revision of this article shows this is related to well-posedness with BMO data.

bisectorial operator and can be shown to have an L_2 -bounded holomorphic functional calculus for any (t-independent) matrix A satisfying (2) and (3). This fact was proved earlier by Axelsson, Keith and McIntosh [13, Theorem 3.1] elaborating on the technology for the solution of the Kato problem by Auscher, Hofmann, Lacey, McIntosh and Tchamitchian [10]; a more direct proof is proposed in Auscher, Axelsson and McIntosh [8]. As explained there, the main difficulty is the non-injectivity of D. The upshot is the possibility of solving (5) by a semi-group formula $f = e^{-t|DB|} f_0$ with f_0 in a suitable trace space, and such f has non-tangential and square function estimates. The BVP can then be solved in an appropriate class if and only if the map from the trace functions to boundary data is invertible. This is the scheme for the Neumann and regularity problems, for which the boundary data is simply the normal or tangential part of $\nabla_A u$. For the Dirichlet problem, it turns out that a "dual" scheme involving the operator BD can be used similarly. The one-to-one correspondence between trace functions f_0 and boundary data may fail, see Axelsson [12], and it is here that restrictions on A appear. It is known to hold if A is (complex) self-adjoint or block form (*i.e.* no cross derivatives $\partial_0 A \partial_i$ or $\partial_i A \partial_0$, $i \geq 1$, in (1)), or constant. Another consequence of this method, and this is why considering complex coefficients is useful, is that the set of t-independent A's for which solvability holds is open in L_{∞} . See [9].

Our work for t-dependent coefficients takes the algebraic reduction to (5) as a starting point, the conormal gradient becoming the central object. We shall state the main results in Section 2 and explain the strategy in Section 3. It involves in particular study of a highly singular integral operator S_A , with an operator-valued kernel. On a technical level, proper definition and handling of this operator is most efficiently done using operational calculus rather than the usual maximal regularity treatment originally due to de Simon [22] (see Kunstmann and Weis [38, Chapter 1] for an overview) and this avoids having to assume qualitatively that A is smooth in the calculations. We use the terminology operational calculus, following the thesis [1] of Albrecht, for the extension of functional calculus when not only scalar holomorphic functions are applied to the underlying operator (in our case DB_0 with $B_0(x) = B(0,x)$, but more general operator-valued holomorphic functions. The Hilbert space theory we use here to prove boundedness on appropriate functional spaces in Section 7, surveyed in Section 6.1, is a special case of the general theory developed in Albrecht, Franks and McIntosh [3, Section 4], Lancien, Lancien and LeMerdy [39], and Lancien and LeMerdy [40]. For further details and references, we refer to Kunstmann and Weis [38, Chapter 12].

The Carleson control on the discrepancy A(t,x) - A(0,x) from [18, 28, 33, 34] appears in a very natural way in the estimates of S_A , and well-posedness of the three BVPs with coefficients A(t,x) will follow under smallness of this control and wellposedness for coefficients A(0,x). We mention that the Dirichlet problem could be obtained by an abstract duality procedure from a regularity problem, in agreement with the results of [33, 34] for real symmetric equations. See also Kilty and Shen [37], and Shen [43]. We will formalize this abstract procedure in our subsequent work [6]. We remark however that although the hypotheses are the same for each BVP, the perturbation results can be proved independently of one another. For example, one does not need knowledge on well-posedness of regularity for A or of Dirichlet for A^* to obtain well-posedness of Neumann for A, in contrast with the results in [34]. We do not know how to prove well-posedness under the finiteness of $||A(t, x) - A(0, x)||_C$ only. (In light of [28, 34], this would first require to extend our methods to solvability for L_p data.) However, thanks to our representations, we do obtain under this hypothesis a number of *a priori* estimates and boundary behaviour on solutions to the equation (1) without any assumption on well-posedness. For example, and we concentrate on this to finish this introduction, we show that if $||A(t, x) - A(0, x)||_C < \infty$, all weak solutions to (1) with coefficients A satisfy, up to a constant, $\sup L_2$ estimates as well as non-tangential maximal functions estimates. More precisely, we obtain

$$\max(\sup_{t>0} \|u_t\|_2, \|\tilde{N}_*(u)\|_2) \lesssim \|\nabla_{t,x}u\|_{L_2(tdt;L_2)}$$

provided the right hand side is finite and the solution u vanishes at infinity in some sense (see Section 2 for precise meaning). Note in particular that this applies when A is t-independent and in that case, this is implicit from [9, Corollary 4.2] when restricted to the class of solutions considered there. Domination of the non-tangential maximal function $\|\widetilde{N}_*(u)\|_2$ by the square function $\|\nabla_{t,x}u\|_{L_2(tdt;L_2)} \approx \|\widetilde{S}(u)\|_2$, $S(u)(x) = (\int_{|y-x| < t} |\nabla_{t,y}u|^2 dt dy/t^{n-1})^{1/2}$, is reminiscent of the result of Dahlberg, Jerison and Kenig [20], and also of Dahlberg, Kenig, Pipher and Verchota [21]. But there is a difference. In [20] comparability of $N_*(u)$ and S(u) is obtained for solutions of the equation (1) under (2) and (3), A real and m = 1, in all $L_q(\mathbf{R}^n; d\mu)$ spaces, $0 < q < \infty$, with μ a doubling A_{∞} weight with respect to L-harmonic measure. If the Dirichlet problem in the class $\|N_*(u)\|_p < \infty$ is proved to be solvable for one $1 , then Lebesgue measure is <math>A_{\infty}$ of L-harmonic measure, hence $\|N_*(u)\|_q \approx \|S(u)\|_q$. This fact follows in particular from combining [29] and [26] under $||A(t,x) - A(0,x)||_C < \infty$ and A real symmetric. In [21], comparability $||N_*(u)||_q \approx ||S(u)||_q, 0 < q < \infty$, is obtained for real symmetric constant elliptic (in the sense of Legendre–Hadamard) second order systems (and also higher order but the formulation becomes different) on bounded Lipschitz domains owing to the fact that L_2 solvability of the Dirichlet problem was known (see the introduction of [21]). This comparability also follows for real non-symmetric scalar equations in two dimensions combining the results of [31] and again [20]. Here, although we obtain only one part of the comparison, it is essential to note that this is an *a priori* estimate valid independently of well-posedness. The existence of an *a priori* proof is new even for real symmetric scalar equations under the Carleson control (for example for all the *t*-independent ones), and is permitted by the solution of the Kato square root problem and its extensions.

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2. Statement of results

In this section we state our results concerning solvability of boundary value problems on the half space \mathbf{R}^{1+n} , and show how they extend to domains Ω which are Lipschitz diffeomorphic to \mathbf{R}^{1+n}_{+} .

Let us first fix notation. We write $\{e_0, e_1, \ldots, e_n\}$ for the standard basis for \mathbf{R}^{1+n} with e_0 "upward" pointing into \mathbf{R}^{1+n}_+ , and write $t = x_0$ for the vertical coordinate. For the vertical derivative, we write $\partial_0 = \partial_t$. For an *m*-tuple of vectors $v = (v_i^{\alpha})_{0 \le i \le n}^{1 \le \alpha \le m}$, we write v_{\perp} and v_{\parallel} for the normal and tangential parts of v, *i.e.* $(v_{\perp})_0^{\alpha} = v_0^{\alpha}$ and $(v_{\perp})_i^{\alpha} = 0$ when $1 \le i \le n$, whereas $(v_{\parallel})_i^{\alpha} = v_i^{\alpha}$ when $1 \le i \le n$ and $(v_{\parallel})_0^{\alpha} = 0$. We write $f_t(x) := f(t, x)$ for functions in \mathbf{R}^{1+n}_+ . As compared to [9], we here use subscript 0 to denote restriction to the boundary \mathbf{R}^n at t = 0, rather than the normal component of f. We also prefer to use small letters f, g, \ldots to denote functions in \mathbf{R}^{1+n}_+ , since this is where we work most of the time, not on the boundary as in [9].

For tuples of functions and vector fields, gradient and divergence act as $(\nabla_{t,x}u)_i^{\alpha} = \partial_i u^{\alpha}$ and $(\operatorname{div}_{t,x} f)^{\alpha} = \sum_{i=0}^n \partial_i f_i^{\alpha}$, with corresponding tangential versions $\nabla_x u = (\nabla_{t,x}u)_{\parallel}$ and $(\operatorname{div}_x f)^{\alpha} = \sum_{i=1}^n \partial_i f_i^{\alpha}$. With $\operatorname{curl}_{t,x} f = 0$ we understand that $\partial_j f_i^{\alpha} = \partial_i f_j^{\alpha}$, for all $i, j = 0, \ldots, n$. Similarly, write $\operatorname{curl}_x f_{\parallel} = 0$ if $\partial_j f_i^{\alpha} = \partial_i f_j^{\alpha}$, for all $i, j = 1, \ldots, n$.

We assume that A is strictly accretive on \mathcal{H} , i.e.

$$\int_{\mathbf{R}^n} \operatorname{Re}(A(t,x)f(x), f(x)) dx \ge \kappa \int_{\mathbf{R}^n} |f(x)|^2 dx,$$

holds for some constant $\kappa > 0$, uniformly for t > 0 and all f belonging to the closed subspace

(6)
$$\mathcal{H} := \mathsf{N}(\operatorname{curl}_x) = \{ g \in L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m}) \; ; \; \operatorname{curl}_x(g_{\parallel}) = 0 \}.$$

For scalar equations, *i.e.* m = 1, (3) amounts to the pointwise condition

$$\operatorname{Re}(A(t,x)\zeta,\zeta) \ge \kappa |\zeta|^2$$
, for all $\zeta \in \mathbf{C}^{1+n}$, a.e. $(t,x) \in \mathbf{R}^{1+n}_+$.

For systems, (3) is stronger than a strict Gårding inequality on \mathbf{R}^{1+n}_+ (*i.e.* integration would be on \mathbf{R}^{1+n}_+ and f such that $\operatorname{curl}_{t,x} f = 0$); still (3) is natural given the type of perturbation we consider here.

The boundary value problems we consider are to find $u \in \mathcal{D}'(\mathbf{R}^{1+n}_+; \mathbf{C}^m)$ solving the divergence form second order elliptic system

$$\operatorname{div}_{t,x} A \nabla_{t,x} u = 0 \text{ in } \mathbf{R}^{1+n}_+$$

in distributional sense, with appropriate interior estimates of $\nabla_{t,x} u$ and satisfying one of the following three natural boundary conditions.

- The Dirichlet condition $u = \varphi$ on \mathbb{R}^n , given $\varphi \in L_2(\mathbb{R}^n; \mathbb{C}^m)$.
- The Dirichlet regularity condition $\nabla_x u = \varphi$ on \mathbf{R}^n , given $\varphi \in L_2(\mathbf{R}^n; \mathbf{C}^{nm})$ satisfying $\operatorname{curl}_x \varphi = 0$.
- The Neumann condition $(e_0, A\nabla_{t,x}u) = \varphi$ on \mathbf{R}^n , given $\varphi \in L_2(\mathbf{R}^n; \mathbf{C}^m)$.

Note that since we shall impose distributional $\nabla_{t,x} u \in L_2^{\text{loc}}$, u can be identified with a function in $W_2^{1,\text{loc}}$, *i.e.* with a weak solution.

Definition 2.1. The modified non-tangential maximal function of a function f in \mathbf{R}^{1+n}_+ is

$$\widetilde{N}_*(f)(x) := \sup_{t>0} t^{-(1+n)/2} ||f||_{L_2(W(t,x))}, \qquad x \in \mathbf{R}^n,$$

where $W(t, x) := (c_0^{-1}t, c_0t) \times B(x; c_1t)$, for some fixed constants $c_0 > 1$, $c_1 > 0$. The modified Carleson norm of a function g in \mathbf{R}^{1+n}_+ is

$$||g||_C := \left(\sup_Q \frac{1}{|Q|} \iint_{(0,l(Q))\times Q} \sup_{W(t,x)} |g|^2 \frac{dtdx}{t}\right)^{1/2},$$

where the supremum is taken over all cubes Q in \mathbb{R}^n , with l(Q) denoting their side lengths.

Note that different choices for c_0, c_1 will give different, but equivalent norms $||N_*(f)||_2$, as well as equivalent norms $||g||_C$. Furthermore, this maximal function is really non-tangential since $\widetilde{N}_*(f)$ and the closely related maximal function $\sup_{|y-x| \leq t} t^{-(1+n)/2} ||f||_{L_2(W(t,y))}$ have equivalent L_2 norms. The latter was introduced in [33]. The modified Carleson norm originates from Dahlberg [19].

We will use the modified Carleson norm to measure the size of perturbations of t-independent coefficients A_0 . (In fact we shall use a possibly weaker quantity $\|\cdot\|_{*}$; see Section 5.) Intuitively, $\|A - A_0\|_C < \infty$ means that in a certain sense $A(t,x) = A_0(x)$ at t = 0, but also that A(t,x) is close to $A_0(x)$ at all scales since we are dealing with a scale invariant norm. Also, given A, such a A_0 is unique and has controlled bounds thanks to the following lemma. We remark here that the modified Carleson norm there can be weakened to the usual one. See Section 5 for proofs.

Lemma 2.2. Fix $A : \mathbf{R}^{1+n}_+ \to \mathcal{L}(\mathbf{C}^{(1+n)m})$ with $||A||_{\infty} < \infty$ and strictly accretive on \mathcal{H} , with constant of accretivity $\kappa > 0$. Assume that A_0 are t-independent coefficients such that $||A - A_0||_C < \infty$. Then A_0 is uniquely determined by A, i.e. if A'_0 are tindependent coefficients such that $||A - A'_0||_C < \infty$, then $A'_0 = A_0$ almost everywhere. Furthermore A_0 is strictly accretive on \mathcal{H} , with

$$\kappa \le \kappa_0 \le \|A_0\|_{\infty} \le \|A\|_{\infty},$$

where κ_0 denotes the constant of accretivity of A_0 .

For the Neumann and Dirichlet regularity problems, our result is the following.

Theorem 2.3. Consider coefficients $A \in L_{\infty}(\mathbf{R}^{1+n}; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ which are strictly accretive on \mathcal{H} .

(i) A priori estimates: Assume that $u \in W_2^{1,\text{loc}}(\mathbf{R}^{1+n}_+, \mathbf{C}^m)$ has gradient $\nabla_{t,x}u$ with estimate $\|\widetilde{N}_*(\nabla_{t,x}u)\|_2 < \infty$, and where u satisfies (1) in \mathbf{R}^{1+n}_+ distributional sense. If there exists t-independent $A_0 \in L_{\infty}(\mathbf{R}^n; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ such that $||A - A_0||_C < \infty$, then $\nabla_{t,x} u$ has limits

$$\lim_{t \to 0} t^{-1} \int_{t}^{2t} \|\nabla_{s,x} u_s - g_0\|_2^2 ds = 0 = \lim_{t \to \infty} t^{-1} \int_{t}^{2t} \|\nabla_{s,x} u_s\|_2^2 ds,$$

for some function $g_0 \in L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$, with estimate $||g_0||_2 \lesssim ||\widetilde{N}_*(\nabla_{t,x}u)||_2$. (ii) Well-posedness: By the Neumann problem with coefficients A (or A_0) being well-posed, we mean that given $\varphi \in L_2(\mathbf{R}^n; \mathbf{C}^m)$, there is a function $u \in$ $W_2^{1,\text{loc}}(\mathbf{R}^{1+n}_+, \mathbf{C}^m)$, unique modulo constants, solving (1), with coefficients A

(or A_0), and having estimates as in (i) and trace $g_0 = \lim_{t\to 0} \nabla_{t,x} u$ such that $(A_0g_0)_{\perp} = \varphi$.

The following perturbation result holds. If the Neumann problem for A_0 is well-posed, then there exists $\epsilon > 0$ such that if $||A - A_0||_C < \epsilon$, then the Neumann problem is well-posed for A.

The corresponding result holds when the Neumann problem is replaced by the regularity problem and the boundary condition $(A_0g_0)_{\perp} = \varphi$ is replaced by $(g_0)_{\parallel} = \varphi \in L_2(\mathbf{R}^n; \mathbf{C}^{nm})$, where φ satisfies $\operatorname{curl}_x \varphi = 0$. Moreover, for both BVPs the solutions u have estimates

$$||N_*(\nabla_{t,x}u)||_2 \approx ||g_0||_2 \approx ||\varphi||_2$$

(iii) Further regularity: Assume that A_0 is as in (i), with $||A - A_0||_C$ sufficiently small and consider solutions u as in (i).

If A satisfy the t-regularity condition $||t\partial_t A||_C < \infty$, then

$$\int_0^\infty \|\partial_t \nabla_{t,x} u\|_2^2 t dt \lesssim \sup_{t>0} \|\nabla_{t,x} u_t\|_2^2 \approx \|\widetilde{N}_*(\nabla_{t,x} u)\|_2^2,$$

 $t \mapsto \nabla_{t,x} u_t \in L_2 \text{ is continuous and } \lim_{t \to 0} \|\nabla_{t,x} u_t - g_0\|_2 = 0 = \lim_{t \to \infty} \|\nabla_{t,x} u_t\|_2.$ The converse estimate $\|\widetilde{N}_*(\nabla_{t,x} u)\|_2^2 \lesssim \int_0^\infty \|\partial_t \nabla_{t,x} u\|_2^2 t dt$ holds provided $\|t\partial_t A\|_C$ is sufficiently small.

If $\max(\|t\partial_i A\|_C, \|t\partial_t A\|_C) < \infty$ holds for some $i = 1, \ldots, n$, then

$$\int_0^\infty \|\partial_i \nabla_{t,x} u\|_2^2 t dt \lesssim \|\widetilde{N}_*(\nabla_{t,x} u)\|_2^2$$

The estimate $\|\widetilde{N}_*(\nabla_{t,x}u)\|_2^2 \lesssim \sum_{i=1}^n \int_0^\infty \|\partial_i \nabla_{t,x}u\|_2^2 t dt$ holds provided $\|t \nabla_{t,x}A\|_C$ is sufficiently small.

Under the hypothesis $||A - A_0||_C < \infty$, the implicit constant in (i) depends on n, m, $||A||_{\infty}$, κ . In (ii) and (iii), under the smallness hypothesis on $||A - A_0||_C$, which depends on n, m, $||A||_{\infty}$, κ (and in (ii) also on the "well-posedness constant" for A_0), the implicit constants depend on n, m, $||A||_{\infty}$, κ (and in (iii) also on the regularity assumptions on A). However, some inequalities are true without smallness and/or well-posedness as the reader can check on reading details in Section 8. Note that in (ii), the uniqueness holds in the class defined by $||\tilde{N}_*(\nabla_{t,x}u)||_2 < \infty$.

For the Dirichlet problem, our main result is the following, including a rigidity result of solutions with square function estimates.

Theorem 2.4. Consider coefficients $A \in L_{\infty}(\mathbf{R}^{1+n}_+; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ which are strictly accretive on \mathcal{H} .

(i) A priori representation and estimates: Assume that $u \in W_2^{1,\text{loc}}(\mathbf{R}_+^{1+n}, \mathbf{C}^m)$ has gradient bounds $\int_0^\infty \|\nabla_{t,x} u_t\|_2^2 t dt < \infty$ and satisfies (1) in \mathbf{R}_+^{1+n} distributional sense. If there exists t-independent $A_0 \in L_\infty(\mathbf{R}^n; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ such that $\|A - A_0\|_C < \infty$, then $u = \hat{u} + c$ almost everywhere, for a unique $\hat{u} \in C(\mathbf{R}_+; L_2(\mathbf{R}^n; \mathbf{C}^m))$ and constant $c \in \mathbf{C}^m$. Identifying the functions u and $\hat{u} + c$, it has L_2 limits

$$\lim_{t \to 0} \|u_t - \hat{u}_0 - c\|_2 = 0 = \lim_{t \to \infty} \|u_t - c\|_2,$$

for some $\hat{u}_0 \in L_2(\mathbf{R}^n; \mathbf{C}^m)$, and we have estimates

$$\max(\|\widetilde{N}_{*}(\hat{u})\|_{2}^{2}, \sup_{t>0} \|\hat{u}_{t}\|_{2}^{2}) \lesssim \int_{0}^{\infty} \|\nabla_{t,x}u\|_{2}^{2} t dt.$$

(ii) Well-posedness: By the Dirichlet problem with coefficients A (or A₀) being well-posed, we mean that given φ ∈ L₂(**R**ⁿ; **C**^m), there is a unique u ∈ W₂^{1,loc}(**R**¹⁺ⁿ₊, **C**^m) having gradient bounds and solving (1), with coefficients A (or A₀), as in (i), and trace u₀ = φ.

The following perturbation result holds. If the Dirichlet problem for A_0 is well-posed, then there exists $\epsilon > 0$ such that if $||A - A_0||_C < \epsilon$, then the Dirichlet problem is well-posed for A. Moreover, these solutions u have estimates

$$\|\widetilde{N}_{*}(u)\|_{2}^{2} \approx \sup_{t>0} \|u_{t}\|_{2}^{2} \approx \int_{0}^{\infty} \|\nabla_{t,x}u\|_{2}^{2} t dt \approx \|\varphi\|_{2}^{2}.$$

Note that in (ii), the boundary condition $u_0 \in L_2(\mathbf{R}^n; \mathbf{C}^m)$ forces c = 0 and $u = \hat{u}$. Under the hypothesis $||A - A_0||_C < \infty$, the implicit constants in (i) depend on n, m, $||A||_{\infty}$, κ . In (ii), under the smallness hypothesis on $||A - A_0||_C$, which depends on n, m, $||A||_{\infty}$, κ and the "well-posedness constant" for A_0 , the implicit constants depend on n, m, $||A||_{\infty}$, κ . However, some inequalities are true without smallness and/or well-posedness as the reader can check on reading details in Section 9. Note that in (ii), uniqueness holds in the class defined by $\int_0^\infty ||\nabla_{t,x}u||_2^2 t dt < \infty$.

As mentioned briefly in the introduction, the hypothesis on well-posedness with *t*independent coefficients A_0 is satisfied, for all three BVPs, for Hermitean coefficients, *i.e.* $A_0(x)^* = A_0(x)$, for block form coefficients, *i.e.* $(A_0)_{\perp\parallel} = 0 = (A_0)_{\parallel\perp}$, and for constant coefficients, *i.e.* $A_0(x) = A_0$, as well as for sufficiently small *t*-independent $L_{\infty}(\mathbf{R}^n; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ perturbations thereof. This was proved in [9, Theorem 2.2]. That the notions of well-posedness of these BVPs used in [9] coincide with the ones here, for *t*-independent coefficients, follows from Corollaries 8.4 and 9.4.

Note that we do not assume pointwise bounds on the solutions, hence we use N_* instead of the usual non-tangential maximal function.

When m = 1 and A, A_0 are real symmetric (and \mathbf{R}^{1+n}_+ replaced by the unit ball), Theorem 2.3(ii) is in [33], and Theorem 2.4(ii) is in [18] (and [20] for the square function estimate). The rest of Theorems 2.3 and 2.4 is mostly new. In Section 3, a more detailed road map to the proofs is given.

Proof of Theorems 2.3 and 2.4. For the Neumann and regularity problems in \mathbf{R}^{1+n}_+ , part (i) follows from Theorem 8.2, part (ii) follows from Corollary 8.5, and part (iii) is proved in Theorem 10.4.

For the Dirichlet problem in \mathbf{R}^{1+n}_+ , part (i) follows from Theorems 9.2 and 9.3, and part (ii) follows from Corollary 9.5, except for the estimate of the non-tangential maximal function, which is proved in Theorem 10.1.

We end this section with a remark on the Lipschitz invariance of the above results. Let $\Omega \subset \mathbf{R}^{1+n}$ be a domain which is Lipschitz diffeomorphic to \mathbf{R}^{1+n}_+ , and let ρ : $\mathbf{R}^{1+n}_+ \to \Omega$ be the Lipschitz diffeomorphism. Denote the boundary by $\Sigma := \partial \Omega$ and the restricted boundary Lipschitz diffeomorphism by $\rho_0 : \mathbf{R}^n \to \Sigma$.

Given a function $\tilde{u} : \Omega \to \mathbf{C}^m$, we pull it back to $u := \tilde{u} \circ \rho : \mathbf{R}^{1+n}_+ \to \mathbf{C}^m$. By the chain rule, we have $\nabla_{t,x} u = \rho^* (\nabla_{t,x} \tilde{u})$, where the pullback of an *m*-tuple of vector fields f, is defined as $\rho^*(f)(x)^{\alpha} := \underline{\rho}^t(x)f^{\alpha}(\rho(x))$, with $\underline{\rho}^t$ denoting the transpose of Jacobian matrix $\underline{\rho}$. If \tilde{u} satisfies $\operatorname{div}_{t,x} \tilde{A} \nabla_{t,x} \tilde{u} = 0$ in Ω , with coefficients $\tilde{A} \in L_{\infty}(\Omega; \mathcal{L}(\mathbf{C}^{(1+n)m}))$, then u will satisfy $\operatorname{div}_{t,x} A \nabla_{t,x} u = 0$ in \mathbf{R}^{1+n}_+ , where $A \in L_{\infty}(\mathbf{R}^{1+n}_+; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ is defined as

(7)
$$A(\mathbf{x}) := |J(\rho)(\mathbf{x})|(\underline{\rho}(\mathbf{x}))^{-1}\tilde{A}(\rho(\mathbf{x}))(\underline{\rho}^{t}(\mathbf{x}))^{-1}, \quad \mathbf{x} \in \mathbf{R}^{1+n}_{+}.$$

Here $J(\rho)$ is the Jacobian determinant of ρ .

The boundary conditions on \tilde{u} on Σ translate in the following way to boundary conditions on u on \mathbb{R}^n .

- The Dirichlet condition $\tilde{u} = \tilde{\varphi}$ on Σ is equivalent to the Dirichlet condition $u = \varphi$ on \mathbb{R}^n , where $\varphi := \tilde{\varphi} \circ \rho_0 \in L_2(\mathbb{R}^n; \mathbb{C}^m)$.
- The Dirichlet regularity condition $\nabla_{\Sigma} \tilde{u} = \tilde{\varphi}$ on Σ (∇_{Σ} denoting the tangential gradient on Σ), is equivalent to $\nabla_x u = \varphi$ on \mathbf{R}^n , where $\varphi := \rho_0^*(\tilde{\varphi}) \in L_2(\mathbf{R}^n; \mathbf{C}^{nm})$.
- The Neumann condition $(\nu, \tilde{A}\nabla_{\Omega}\tilde{u}) = \tilde{\varphi}$ on Σ (contrary to tradition, ν being the inward unit normal vector field on Σ) is equivalent to $(e_0, A\nabla_{t,x}u) = \varphi$ on \mathbb{R}^n , where $\varphi := |J(\rho_0)|\tilde{\varphi} \circ \rho_0 \in L_2(\mathbb{R}^n; \mathbb{C}^m)$.

In this way the Dirichlet/regularity/Neumann problem with coefficients \tilde{A} in the Lipschitz domain Ω is equivalent to the Dirichlet/regularity/Neumann problem with coefficients A in the half space \mathbf{R}^{1+n}_+ , and it is straightforward to extend Theorems 2.3 and 2.4 to the Lipschitz domain Ω .

3. Road map to the proofs

For the readers convenience, we give in this section an informal explanation of the main ideas behind the methods and the proofs of Theorems 2.3 and 2.4. In particular, the precise definitions of classes of solutions will be given later. Our basic idea for constructing solutions u to the divergence form equation (1) in \mathbf{R}^{1+n}_+ is to consider it as a first order system with the gradient $\nabla_{t,x}u$ as the unknown function. In fact, solving for the *t*-derivatives in the equation, the divergence form equation for u becomes a vector-valued ODE

$$\partial_t (\nabla_{t,x} u) + T_A (\nabla_{t,x} u) = 0,$$

where T_A is an operator only involving the first order derivatives along \mathbb{R}^n and multiplication by entries of A(t, x). The divergence form equation was first studied through this ODE in [7]. However, it turns out that if one instead of $\nabla_{t,x}u$ takes the conormal gradient $\nabla_A u$ defined by (4) as the unknown, then the corresponding operator T_A has a simpler structure; the ODE reads

(8)
$$\partial_t f + DBf = 0, \quad \text{with } f := \nabla_A u,$$

where $D := \begin{bmatrix} 0 & \operatorname{div}_x \\ -\nabla_x & 0 \end{bmatrix}$ and B is a second strictly accretive coefficient matrix determined by A. This was the key discovery in [9] when coefficients A do not depend on t. This carries over to t-dependent coefficients and as this result is central to us, we give full proof of it in Proposition 4.1.

3.1. The Neumann and regularity problems. The first order approach is most natural for solving the Neumann and regularity BVPs, since these boundary conditions are conditions on the conormal gradient f, not on the potential u. Indeed, the Neumann BVP means that the normal part $(f_0)_{\perp} = \partial_{\nu_A} u|_{\mathbf{R}^n}$ is given at the boundary t = 0, whereas the regularity condition is that the tangential part $(f_0)_{\parallel} = \nabla_x u|_{\mathbf{R}^n}$ is given. Note that both BVPs mean that "one half" of the function f_0 is prescribed. This is natural for a first order elliptic equation.

On the other hand, the set of all traces $f_0 = f|_{\mathbf{R}^n}$ of solutions f to the equation $\partial_t f + DBf = 0$ in \mathbf{R}^{1+n}_+ , with appropriate estimates, is a subspace of $L_2 = L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$ which we denote $E_A^+ \mathcal{H}$. Here the reader should have the classical situation in mind, where (8) is the Cauchy–Riemann equations and $E_A^+ \mathcal{H}$ is the upper Hardy subspace of L_2 . Just like in this classical situation, it turns out that for t-independent coefficients and small perturbations thereof, $E_A^+ \mathcal{H}$ is a closed proper subspace, being "one half" of L_2 , and there is a bounded Hardy type projection E_A^+ onto $E_A^+ \mathcal{H}$. Moreover, there is a Cauchy type reproducing formula for the solution fto the ODE, given $f_0 \in E_A^+ \mathcal{H}$. Hence there is a one-to-one correspondence between solutions f and their traces $f_0 \in E_A^+ \mathcal{H}$.

Proving these facts for small perturbations of t-independent coefficients is the main work in this paper. For t-independent coefficients, this result is in [9]. Before explaining the proofs in more detail, assume for a moment the stated properties of E_A^+ and $E_A^+\mathcal{H}$, in order to explain the implications for BVPs. The unique solvability of the Neumann BVP means that for each boundary data $\varphi \in L_2(\mathbb{R}^n; \mathbb{C}^m)$ there is a unique $f \in E_A^+\mathcal{H}$ such that $f_{\perp} = \varphi$. Hence well-posedness of the Neumann BVP is equivalent to

$$E_A^+ \mathcal{H} \to L_2(\mathbf{R}^n; \mathbf{C}^m) : f \mapsto f_\perp$$

being an isomorphism. Similarly well-posedness of the regularity BVP is equivalent to

$$E_A^+ \mathcal{H} \to \{g \in L_2(\mathbf{R}^n; \mathbf{C}^{nm}) ; \operatorname{curl}_x g = 0\} : f \mapsto f_{\parallel}$$

being an isomorphism. Even for t-independent coefficients, these maps are not always invertible. For counter examples based on [36] in this context, see [12]. However, we mentioned three important classes of t-independent coefficients where techniques are available to prove invertibility.

Concerning this approach to solving BVPs, it should be noted that the problem is divided into two parts: (i) understanding representation formulas and the trace space $E_A^+\mathcal{H}$ for solutions to the differential equation, and (ii) understanding the relation between the full trace space $E_A^+\mathcal{H}$ and the boundary condition (Neumann or regularity). Carefully note that (i) only involves the differential equation and not the boundary conditions. This is of great importance, since it means that it suffices to study the ODE (8), and explains why the trace spaces $E_A^+\mathcal{H}$ may be well behaved even when the BVPs fail to be well-posed. The main harmonic analysis work goes into proving that the projections E_A^+ are bounded for all complex *t*-independent coefficients, and small *t*-dependent perturbations thereof. From this it is deduced that the projections E_A^+ depend continuously on *A* in a certain Carleson sense, and as a consequence in Corollary 8.5 that well-posedness of the Neumann and regularity problems is stable under small perturbations of the coefficients.

We next explain our methods for solving the ODE, *i.e.* solving problem (i). For this we study (8), where we first consider t-independent coefficients $B = B_0$,

and we write $E_A^+ = E_0^+$. In this case, we view DB_0 as an unbounded operator in $L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$, and at a first glance the solution to (8) with initial/boundary data f_0 seems to be $f_t = e^{-tDB_0}f_0$. However, the problem is that DB_0 is not a sectorial operator, but instead bisectorial, *i.e.* its spectrum is contained in a double sector around the real axis. This indefiniteness means that the operators e^{-tDB_0} are not well defined on $L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$ for any $t \neq 0$. Another technical problem is that DB_0 has an infinite dimensional null space. The fact is that there are topological splittings

$$L_2 = \mathcal{H} \oplus \mathsf{N}(DB_0) = \left(E_0^+ \mathcal{H} \oplus E_0^- \mathcal{H}\right) \oplus \mathsf{N}(DB_0).$$

where $\mathcal{H} = \overline{\mathsf{R}(DB_0)} = \overline{\mathsf{R}(D)}$ is the closure of the range. The splitting of \mathcal{H} into the spectral subspace $E_0^+\mathcal{H}$ for the sector in the right half plane and the spectral subspace $E_0^-\mathcal{H}$ for the sector in the left half plane is a deep result, and builds on the Kato square root problem as discussed in the introduction. This proof also shows that DB_0 has square function estimates, which in particular shows that $-DB_0$ generates a bounded holomorphic semigroup in $E_0^+\mathcal{H}$, and that DB_0 generates a bounded holomorphic semigroup in $E_0^-\mathcal{H}$.

Given any $f_0 \in E_0^+ \mathcal{H}$, differentiation as seen in [9] shows that the generalized Cauchy reproducing formula

$$f = C_0^+ f_0,$$

with $C_0^+ f_0(t,x) = (e^{-tDB_0}E_0^+ f_0)(x)$, yields a solution to (8) with trace f_0 . Conversely, given a solution f to (8), it is a fact that f_t belongs to the range \mathcal{H} for any conormal gradient f and t > 0. (Note that it follows from (8) that $\partial_t f_t \in \mathcal{H}$.) We apply the projections E_0^{\pm} and suitable exponentials to the equation, giving

$$\begin{cases} \partial_s (e^{-(t-s)DB_0} E_0^+ f_s) = 0, & s \in (0,t), \\ \partial_s (e^{(s-t)DB_0} E_0^- f_s) = 0, & s \in (t,\infty). \end{cases}$$

Integration with limits $\lim_{s\to\infty} f_s = 0$ and $\lim_{s\to 0} f_s = f_0$ indicates that the trace f_0 belongs to the subspace $E_0^+ \mathcal{H}$ and that the generalized Cauchy reproducing formula $f = C_0^+ f_0$ holds. This converse and existence of the limits are shown in Corollary 8.4. For further details concerning the relation between C_0^+ and the classical Cauchy integral, see [7, Ex. 1.2] and [12, Thm. 2.1]. Note that $C_0^+ f_0 = 0$ for all $f_0 \in E_0^- \mathcal{H}$.

For t-dependent perturbations B of a given t-independent coefficient matrix B_0 , we write the ODE as $\partial_t f + DB_0 f_t = D\mathcal{E}_t f_t$, with $\mathcal{E} := B_0 - B$. The above argument now gives

$$\begin{cases} \partial_s (e^{-(t-s)DB_0} E_0^+ f_s) = e^{-(t-s)DB_0} E_0^+ D\mathcal{E}_s f_s, & s \in (0,t), \\ \partial_s (e^{(s-t)DB_0} E_0^- f_s) = e^{(s-t)DB_0} E_0^- D\mathcal{E}_s f_s, & s \in (t,\infty), \end{cases}$$

and integration and subtraction of the equations give the integral equation

(9)
$$f = C_0^+ h^+ + S_A f,$$

for some $h^+ \in E_0^+ \mathcal{H}$ (which will be shown to be $h^+ = E_0^+ f_0$) and where S_A is the integral operator given by

(10)
$$S_A f_t := \int_0^t e^{-(t-s)DB_0} E_0^+ D\mathcal{E}_s f_s ds - \int_t^\infty e^{(s-t)DB_0} E_0^- D\mathcal{E}_s f_s ds, \quad t > 0.$$

To construct solutions to (9), it is therefore natural to think of the Picard fixed point theorem. For this, we need an appropriate function space of functions in \mathbf{R}^{1+n}_+

that contains the free evolution $g = C_0^+ h^+$ and on which S_A is bounded. By [9], the non-tangential maximal function in Definition 2.1 of g belongs to $L_2(\mathbf{R}^n)$. Thus the space \mathcal{X} defined by $\|\widetilde{N}_*(f)\|_2 < \infty$ is a natural candidate (see Section 5).

We turn to a closer look at S_A , and that allows us to justify the equivalence between (9) and the ODE. This operator involves forward and backward maximal regularity operators. Usual treatment of maximal regularity is not sufficient for our needs. A convenient way to study S_A , in the spirit of functional calculus, is as follows. Define, for $\lambda \in \mathbf{C}$ with $\operatorname{Re} \lambda > 0$, the bounded integral operator $F(\lambda) :$ $L_2(\mathbf{R}_+; \mathbf{C}) \to L_2(\mathbf{R}_+; \mathbf{C}) : u = (u_s)_{s>0} \mapsto F(\lambda)u$ by

$$F(\lambda)u_t := \int_0^t \lambda e^{-(t-s)\lambda} u_s ds.$$

By letting $F(\lambda)$ act pointwise in $x \in \mathbf{R}^n$, it defines a bounded operator on $L_2(\mathbf{R}_+; \mathcal{H})$. In this space we also have the operator

$$|DB_0|f := \begin{cases} DB_0f, & f \in E_0^+\mathcal{H}, \\ -DB_0f, & f \in E_0^-\mathcal{H}, \end{cases}$$

by letting it act pointwise in $t \in \mathbf{R}_+$. The operator $|DB_0|$ is a sectorial operator in \mathcal{H} , hence in $L_2(\mathbf{R}_+; \mathcal{H})$, and on the sector containing its spectrum, $\lambda \mapsto F(\lambda)$ defines an operator-valued holomorphic function. Similar to Dunford functional calculus, we can apply $(F(\lambda))_{\lambda}$ to $|DB_0|$, since $|DB_0|$ commutes with each $F(\lambda)$, and we obtain an operator $F(|DB_0|)$ on $L_2(\mathbf{R}_+; \mathcal{H})$. An advantage of this method, referred to as *operational calculus*, is that boundedness of the singular integral $f = (f_s)_{s>0} \mapsto$ $F(|DB_0|)f$ defined by for all t > 0,

$$F(|DB_0|)f_t := \int_0^t |DB_0|e^{-(t-s)|DB_0|} f_s ds,$$

can be easily derived from the square function estimates for DB_0 , in exactly the same way as the boundedness of E_0^{\pm} , $e^{-t|DB_0|}$, or more general operators in the functional calculus of DB_0 is proved.

The upshot of all this is a representation

(11)
$$S_A = F(|DB_0|)\hat{E}_0^+ \mathcal{E} + F^*(|DB_0|)\hat{E}_0^- \mathcal{E},$$

where \widehat{E}_0^{\pm} are bounded operators such that $E_0^{\pm}D = (DB_0)\widehat{E}_0^{\pm}$ and $F^*(\lambda)u_t := \int_t^{\infty} \lambda e^{-(s-t)\lambda} u_s ds$. Since all operators on the right hand side in (11) are bounded on $L_2(\mathbf{R}_+; \mathcal{H})$, so is S_A . More generally, this boundedness proof goes through on weighted spaces $L_2(\mathbf{R}_+, t^{\alpha}dt; \mathcal{H})$ for $|\alpha| < 1$. Details are in Section 6. However, for $\alpha = -1$, which is the natural scaling for BVPs with L_2 data, boundedness of $F^*(|DB_0|)$ fails. Recall nonetheless that the free evolution belongs to \mathcal{X} , which contains $L_2(\mathbf{R}_+, dt/t; L_2)$ as a subspace. (As compared to the space $\mathcal{X}, L_2(\mathbf{R}_+, dt/t; L_2)$ consists of functions that vanish in a certain sense at t = 0.) Thus we can replace $L_2(\mathbf{R}_+, dt/t; L_2)$ by the larger space \mathcal{X} . To ensure boundedness of S_A on \mathcal{X} , we still need additional control on \mathcal{E} such as finiteness of $||\mathcal{E}||_C$ of Definition 2.1 (in fact a possibly smaller quantity suffices) and our fundamental estimate is

(12)
$$||N_*(S_A f)||_2 \lesssim ||A - A_0||_C ||N_*(f)||_2,$$

obtained from the chain

$$\mathcal{X} \xrightarrow{\mathcal{E}} L_2(\mathbf{R}_+, dt/t; L_2) \xrightarrow{\widehat{E}_0^-} L_2(\mathbf{R}_+, dt/t; \mathcal{H}) \xrightarrow{F^*(|DB_0|)} \mathcal{X}.$$

See Lemma 5.5 for this modified Carleson embedding theorem in the first arrow and Theorem 6.8 for remaining details. This allows us to prove that the trace of f at t = 0 exists in a certain sense and one sees that $E_0^+ f_0 = h^+$. Details of this representation of solutions f are in Theorem 8.2.

Now, smallness of this Carleson norm implies smallness of $||S_A||_{\mathcal{X}\to\mathcal{X}}$, in which case (9) rewrites as a Cauchy reproducing formula

$$f = (I - S_A)^{-1} C_0^+ h^+,$$

by inverting $I - S_A$ in \mathcal{X} . Conversely, for any $h^+ \in E_0^+ \mathcal{H}$, the Picard iteration scheme applies to produce a solution $f = (I - S_A)^{-1} C_0^+ h^+$ to (9) whose trace f_0 is determined by a linear relation $E_A^+ h^+ = f_0$.

This provides us with the Hardy type projections E_A^+ needed to solve the Neumann and regularity problems as described above. Details are in Corollary 8.5. Additional *a priori* square function estimates on solutions, under further regularity assumption on the coefficients, can be shown. They are proved in Section 10.2.

3.2. The Dirichlet problem. For the Dirichlet problem with L_2 boundary data, it is not obvious that the above first order approach applies. Nevertheless it is possible to adapt the arguments and we describe this now. Instead of (8) for the conormal gradient f, we want to work with the potential u at the L_2 level. The heuristic is that u solves the divergence form equation if and only if there is a vector-valued potential v such that

(13)
$$\partial_t v + BDv = 0$$

and $u = -v_{\perp}$. (The minus sign is just for convenience.) On one hand, we have seen that u is a divergence form equation if and only if $(\partial_t + DB)(\nabla_A u) = 0$. On the other hand, applying D to the ODE (13) gives $(\partial_t + DB)(Dv) = 0$. Thus the claim amounts to rewrite the conormal gradient $f = \nabla_A u$ as f = Dv. (Note that the latter equation is pointwise in t.) Evaluating the tangential part f_{\parallel} shows the relation $u = -v_{\perp}$. Once again the reader should compare this to the classical situation of Cauchy–Riemann's equations, in which case u would be a harmonic function, and v the analytic function having u and its harmonic conjugate function as imaginary and real parts.

We mention that the v we construct via the conormal gradient again does not quite satisfy (13) because of the null space of D, but that is not a problem as we only need its normal part which does not depend on this null space. Having this heuristic in mind we proceed as before in two steps: (i) understanding representation of solutions, and (ii) understanding its trace and the relation to solvability. We mostly concentrate on (i) as (ii) will follow right away.

For t-independent coefficients $B = B_0$, we note that B_0D is another bisectorial operator, just like DB_0 . It is not injective (the null space is \mathcal{H}^{\perp}) and the spectral projections \tilde{E}_0^{\pm} for the sectors in the right and left half planes split the range $B_0\mathcal{H}$ into two closed subspaces. Similar to the argument for DB_0 , we have that solutions obeying a certain square function estimate to the ODE (13) all are of the form $v = \tilde{C}_0^+ v_0 + c$ for a unique $v_0 \in \tilde{E}_0^+ L_2$ and some $c \in \mathbf{C}^{(1+n)m}$, where \tilde{C}_0^+ is defined by $(\tilde{C}_0^+ v_0)(t, x) = (e^{-tB_0D}\tilde{E}_0^+ v_0)(x)$. Thus, for t-independent coefficients, we have the representation

$$u = c - (\widetilde{C}_0^+ v_0)_{\perp}, \qquad v_0 \in \widetilde{E}_0^+ L_2, c \in \mathbf{C}^m,$$

for solutions u to (1) obeying a square function estimate. This is in Corollary 9.4 and improves [9] where this was shown for a smaller class of solutions. Note that the conormal gradient of u can be calculated as $f = D\tilde{C}_0^+ v_0$.

For t-dependent perturbations B of a given t-independent coefficient matrix B_0 , suppose we are given a solution u. Since we do not know v yet, we go via $f = \nabla_A u$ (which in the end will be Dv) to solve (13). Since f satisfies $(\partial_t + DB)f = 0$, as in Section 3.1, we look for a functional setting in which we obtain an equation of the form $f_t = e^{-tDB_0}h^+ + S_A f_t$ for some h^+ in a positive spectral space so as to again apply the Picard fixed point theorem to construct solutions. The main difference here is in the free evolution term $g_t = e^{-tDB_0}h^+$. Indeed, f_0 , which we should relate to h^+ , may only be defined in a space of Sobolev type with regularity index -1 (and in fact, we do not care about f_0).

Square function estimates (see [9]) lead us to the solution. Indeed, we have $\int_0^\infty \|e^{-tDB_0}h^+\|_2^2 t dt \approx \|D^{-1}h^+\|_2^2$ provided $D^{-1}h^+ \in L_2$. Here D^{-1} is defined as a closed operator as we assumed h^+ in a positive spectral space. Hence the natural (considering the method) space of conormal gradients for Dirichlet problems is $\mathcal{Y} := L_2(\mathbf{R}_+, t dt; L_2)$. Indeed, imposing the free evolution g to be in \mathcal{Y} allows us to obtain $g = D\widetilde{C}_0^+ \widetilde{h}^+$ for $\widetilde{h}^+ \in \widetilde{E}_0^+ L_2$ determined by $\widetilde{h}^+ = D^{-1}h^+$. (Note that D and B_0 are interchanged.) This is consistent with the t-independent case.

The next step is thus to bound S_A in \mathcal{Y} . We use again the operational calculus representation (11). In the space \mathcal{Y} , it is now the operator $F(|DB_0|)$ that fails to be bounded, but again the additional Carleson control ensures boundedness of $||S_A||_{\mathcal{Y}\to\mathcal{Y}}$ and our second fundamental estimate reads

(14)
$$\|S_A f\|_{L_2(tdt;L_2)} \lesssim \|A - A_0\|_C \|f\|_{L_2(tdt;L_2)}$$

Details are in Proposition 7.1. This allows us to justify the formal manipulations and to obtain an *a priori* representation of conormal gradients $f = \nabla_A u$ in \mathcal{Y} of solutions u by

(15)
$$f = D\widetilde{C}_0^+ \widetilde{h}^+ + S_A f, \qquad \widetilde{h}^+ \in \widetilde{E}_0^+ L_2.$$

We next want to exhibit the announced vector-valued potential v which must satisfy f = Dv. Remark that D being non-injective, there is some freedom in the choice but we basically want to factor out D in (15). This is granted for the free evolution and for S_A we obtain starting from (10) that $S_A f = D\tilde{S}_A f$, where

(16)
$$\widetilde{S}_A f_t := \int_0^t e^{-(t-s)B_0 D} \widetilde{E}_0^+ \mathcal{E}_s f_s ds - \int_t^\infty e^{(s-t)B_0 D} \widetilde{E}_0^- \mathcal{E}_s f_s ds$$

All this follows from the intertwining $f(DB_0)D = Df(B_0D)$ of the functional calculi of DB_0 and B_0D . Thus we can set

$$v := \widetilde{C}_0^+ \widetilde{h}^+ + \widetilde{S}_A f,$$

and then show that there exists a constant $c \in \mathbf{C}^m$ such that

$$u = c - (v)_{\perp},$$

for solutions u to (1) obeying an initial square function estimate with f being the conormal gradient of u. Again, this is not the only possible choice for v but any other choice has identical normal part and this is what we need to recover u. Moreover, this choice has good estimates. Details of this representation of solutions u to the divergence form equation are in Theorems 9.2 and 9.3. We note that this proves

that all solutions with gradient satisfying a square function estimates are, up to constants, continuous in t with values in $L_2(\mathbf{R}^n, \mathbf{C}^m)$. This representation also enables us to show existence of the trace $v_0 = \lim_{t\to 0} v_t$ in $L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$ and \tilde{h}^+ is determined by $\tilde{h}^+ = \tilde{E}_0^+ v_0$. Most importantly, the representation allows us to prove non-tangential maximal function estimates in Section 10.1. We remark they are *a priori* estimates.

Upon smallness of the Carleson control, these representations above rewrite

$$f = (I - S_A)^{-1} D \widetilde{C}_0^+ \widetilde{h}^+, \qquad \widetilde{h}^+ \in \widetilde{E}_0^+ L_2,$$
$$v = \widetilde{C}_0^+ \widetilde{h}^+ + \widetilde{S}_A f,$$
$$u = c - (v)_\perp,$$

and, conversely, this can serve (via Picard's iteration scheme to obtain the first equation) as an ansatz to produce a solution u starting from a given $\tilde{h}^+ \in \tilde{E}_0^+ L_2$ and constant c = 0, and $v_0 \in L_2$ is determined by a linear relation $v_0 = \tilde{E}_A^+ \tilde{h}^+$. With this ansatz, well-posedness of the Dirichlet problem is seen to be equivalent to

$$\widetilde{E}_0^+ L_2 \to L_2(\mathbf{R}^n; \mathbf{C}^m) : \widetilde{h}^+ \mapsto u_0 = -(v_0)_\perp$$

being an isomorphism. This allows us to prove stability of well-posedness of the Dirichlet problem under small perturbations in Corollary 9.5.

4. INTEGRATION OF THE DIFFERENTIAL EQUATION

Following [9], we construct solutions u to the divergence form system (1), by replacing u by its gradient g as the unknown function. Consequently (1) for u is replaced by (17) below for g. Proposition 4.1 reformulates this first order system (17) further, by solving for the *t*-derivatives, as the vector-valued ODE (18) for the conormal gradient

$$f = \nabla_A u = [\partial_{\nu_A} u, \nabla_x u]^t, \quad \text{where } [\alpha, v]^t := \begin{bmatrix} \alpha \\ v \end{bmatrix}$$

for $\alpha \in \mathbf{C}^m$ and $v \in \mathbf{C}^{nm}$, and $\partial_{\nu_A} u := (A \nabla_{t,x} u)_{\perp}$ denotes the (inward!) conormal derivative of u.

According to the decomposition of m-tuples into normal and tangential parts as introduced in Section 2, we split the matrix as

$$A(t,x) = \begin{bmatrix} A_{\perp\perp}(t,x) & A_{\perp\parallel}(t,x) \\ A_{\parallel\perp}(t,x) & A_{\parallel\parallel}(t,x) \end{bmatrix}$$

Note that with our assumption that A be strictly accretive on \mathcal{H} for a.e. t > 0, the matrix $A_{\perp\perp}$ is invertible.

Proposition 4.1. The pointwise transformation

$$A \mapsto \hat{A} := \begin{bmatrix} A_{\perp\perp}^{-1} & -A_{\perp\perp}^{-1}A_{\perp\parallel} \\ A_{\parallel\perp}A_{\perp\perp}^{-1} & A_{\parallel\parallel} - A_{\parallel\perp}A_{\perp\perp}^{-1}A_{\perp\parallel} \end{bmatrix}$$

is a self-inverse bijective transformation of the set of bounded matrices which are strictly accretive on \mathcal{H} .

For a pair of coefficient matrices $A = \hat{B}$ and $B = \hat{A}$, the pointwise map $g \mapsto f = [(Ag)_{\perp}, g_{\parallel}]^t$ gives a one-one correspondence, with inverse $g = [(Bf)_{\perp}, f_{\parallel}]^t$, between solutions $g \in L_2^{\text{loc}}(\mathbf{R}_+; L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m}))$ to the equations

(17)
$$\begin{cases} \operatorname{div}_{t,x}(Ag) = 0\\ \operatorname{curl}_{t,x}g = 0 \end{cases}$$

and solutions $f \in L_2^{\text{loc}}(\mathbf{R}_+; \mathcal{H})$ to the generalized Cauchy-Riemann equations

(18)
$$\partial_t f + DBf = 0,$$

where the derivatives are taken in \mathbf{R}^{1+n}_+ distributional sense, and $D := \begin{bmatrix} 0 & \operatorname{div}_x \\ -\nabla_x & 0 \end{bmatrix}$.

This was proved in [9, Section 3] for t-independent coefficients. The proof goes through without changes for t-dependent coefficients, but for completeness we give the proof of this important result. Note that $\overline{\mathsf{R}(D)} = \mathcal{H}$.

Proof. We first look at the correspondence $A \mapsto \hat{A} = B$. Fix an arbitrary t > 0and we write A for $A(t, \cdot)$. From the accretivity of A on \mathcal{H} , it follows that the component $A_{\perp\perp}$ is pointwise strictly accretive, hence invertible, and therefore so is $\overline{A} := \begin{bmatrix} A_{\perp\perp} & A_{\perp\parallel} \\ 0 & I \end{bmatrix}$. Thus, multiplication by \overline{A} is an isomorphism on \mathcal{H} and, letting $\underline{A} := \begin{bmatrix} I & 0 \\ A_{\parallel\perp} & A_{\parallel\parallel} \end{bmatrix}$, $B = \underline{A}\overline{A}^{-1}$ is bounded if A is so. We calculate, for any fixed $q \in \mathcal{H}$ and $f = \overline{A}q$,

$$\begin{aligned} \operatorname{Re}(Bf,f) &= \operatorname{Re}(\hat{A}\overline{A}g,\overline{A}g) = \operatorname{Re}\left(\begin{bmatrix} I & 0\\ A_{\parallel\perp} & A_{\parallel\parallel} \end{bmatrix} \begin{bmatrix} g_{\perp}\\ g_{\parallel} \end{bmatrix}, \begin{bmatrix} A_{\perp\perp} & A_{\perp\parallel}\\ 0 & I \end{bmatrix} \begin{bmatrix} g_{\perp}\\ g_{\parallel} \end{bmatrix} \right) \\ &= \operatorname{Re}\left(\begin{bmatrix} A_{\perp\perp} & A_{\perp\parallel}\\ A_{\parallel\perp} & A_{\parallel\parallel} \end{bmatrix} \begin{bmatrix} g_{\perp}\\ g_{\parallel} \end{bmatrix}, \begin{bmatrix} g_{\perp}\\ g_{\parallel} \end{bmatrix} \right) = \operatorname{Re}(Ag,g). \end{aligned}$$

This shows that $B = \hat{A}$ is strictly accretive if A is so. That $\hat{A} = A$ is straightforward to verify, and this shows that A and \hat{A} are in one-to-one correspondence.

Next consider a pair of functions g and f in $L_2^{\text{loc}}(\mathbf{R}_+; L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m}))$ such that $f = \overline{A}g$. Equations (17) for g are equivalent to

(19)
$$\begin{cases} \partial_t (Ag)_{\perp} + \operatorname{div}_x (A_{\parallel \perp} g_{\perp} + A_{\parallel \parallel} g_{\parallel}) = 0, \\ \partial_t g_{\parallel} - \nabla_x g_{\perp} = 0, \\ \operatorname{curl}_x g_{\parallel} = 0. \end{cases}$$

The last equation is equivalent to $f_t \in \mathcal{H}$. Moreover, using that $(Ag)_{\perp} = f_{\perp}$, $g_{\parallel} = f_{\parallel}$ and $g_{\perp} = (Bf)_{\perp} = A_{\perp \perp}^{-1}(f_{\perp} - A_{\perp \parallel}f_{\parallel})$, the first two equations are seen to be equivalent to the equation $\partial_t f + DBf = 0$. This proves the proposition.

Remark 4.2. In terms of the second order equation where g is a gradient and f is the corresponding conormal gradient, the identity $\operatorname{Re}(Bf, f) = \operatorname{Re}(Ag, g)$ rewrites $\operatorname{Re}(B\nabla_A u, \nabla_A u) = \operatorname{Re}(A\nabla_{t,x} u, \nabla_{t,x} u)$ for any appropriate u (not necessarily a solution).

We now want to construct solutions to (18). Let us first recall the situation when $B(t,x) = B_0(x)$ does not depend on the *t*-variable. In this case, we view B_0 as

a multiplication operator in the boundary function space $L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$. Define closed and open sectors and double sectors in the complex plane by

$$S_{\omega+} := \{\lambda \in \mathbf{C} : |\arg \lambda| \le \omega\} \cup \{0\}, \qquad S_{\omega} := S_{\omega+} \cup (-S_{\omega+}),$$
$$S_{\nu+}^o := \{\lambda \in \mathbf{C} : \lambda \ne 0, |\arg \lambda| < \nu\}, \qquad S_{\nu}^o := S_{\nu+}^o \cup (-S_{\nu+}^o),$$

and define the angle of accretivity of B_0 to be

$$\omega := \sup_{f \neq 0, f \in \mathcal{H}} |\arg(B_0 f, f)| < \pi/2.$$

The method for constructing solutions to the elliptic divergence form system, developed in [7, 9], uses holomorphic functional calculus of the *infinitesimal generator* DB_0 appearing in the ODE (18), and the following was proved.

- (i) The operator DB_0 is a closed and densely defined ω -bisectorial operator, *i.e.* $\sigma(DB_0) \subset S_\omega$, where ω is the angle of accretivity of B_0 . Moreover, there are resolvent bounds $\|(\lambda - DB_0)^{-1}\| \leq 1/\text{dist}(\lambda, S_\omega)$ when $\lambda \notin S_\omega$.
- (ii) The function space splits topologically as

$$L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m}) = \mathcal{H} \oplus \mathsf{N}(DB_0),$$

and the restriction of DB_0 to $\mathcal{H} = \overline{\mathsf{R}(D)}$ is a closed, densely defined and injective operator with dense range in \mathcal{H} , with same estimates on spectrum and resolvents as in (i).

(iii) The operator DB_0 has a bounded holomorphic functional calculus in \mathcal{H} , *i.e.* for each bounded holomorphic function $b(\lambda)$ on a double sector S_{ν}^{o} , $\omega < \nu < \pi/2$, the operator $b(DB_0)$ in \mathcal{H} is bounded with estimates

$$\|b(DB_0)\|_{\mathcal{H}\to\mathcal{H}} \lesssim \|b\|_{L_{\infty}(S_{\nu}^o)}.$$

For background material on sectorial operators (which is straightforward to adapt to bi-sectorial operators) and their holomorphic functional calculi, see [2]. The construction of the operators $b(DB_0)$ is explained in detail in Section 6.1, in the more general case of operational calculus. The two most important functions $b(\lambda)$ here are the following.

- The characteristic functions $\chi^+(\lambda)$ and $\chi^-(\lambda)$ for the right and left half planes, which give the generalized Hardy projections $E_0^{\pm} := \chi^{\pm}(DB_0)$.
- The exponential functions $e^{-t|\lambda|}$, t > 0, which give the operators $e^{-t|DB_0|}$. Here $|\lambda| := \lambda \operatorname{sgn}(\lambda)$ and $\operatorname{sgn}(\lambda) := \chi^+(\lambda) - \chi^-(\lambda)$.

A key result that we make use of frequently, is that the boundedness of the projections E_0^{\pm} shows that there is a topological splitting

(20)
$$\mathcal{H} = E_0^+ \mathcal{H} \oplus E_0^- \mathcal{H}$$

of $\mathcal{H} = \overline{\mathsf{R}(D)} = \overline{\mathsf{R}(DB_0)}$ into complementary closed subspaces $E_0^{\pm}\mathcal{H} := \mathsf{R}(E_0^{\pm})$.

We also recall the definition of the generalized Cauchy extension C_0^+ from Section 3.

Proposition 4.3. The generalized Cauchy extension

$$f_t = (C_0^+ f_0)(t, \cdot) := e^{-t|DB_0|} E_0^+ f_0$$

of $f_0 \in E_0^+\mathcal{H}$ gives a solution to $\partial_t f + DB_0 f = 0$, in the strong sense $f \in C^1(\mathbf{R}_+; L_2) \cap C^0(\mathbf{R}_+; D(DB_0))$, with L_2 bounds $\sup_{t>0} \|f_t\|_2 \approx \|f_0\|_2$ and L_2 limits $\lim_{t\to 0} f_t = f_0$ and $\lim_{t\to\infty} f_t = 0$.

Now consider more general t-dependent coefficients B(t, x). Fix some t-independent coefficients B_0 , strictly accretive on \mathcal{H} . (This B_0 should be thought of as the boundary trace of B, acting in \mathbf{R}^{1+n}_+ independently of t.) To construct solutions to the ODE, we rewrite it as

(21)
$$\partial_t f + DB_0 f = D\mathcal{E}f, \quad \text{where } \mathcal{E}_t := B_0 - B_t.$$

However, while $\partial_t f + DB_0 f = 0$ can be interpreted in the strong sense, (21) will be understood in the sense of distributions. The following proposition rewrites this equation in integral form. It uses operators \widehat{E}_0^{\pm} , defined as

(22)
$$\widehat{E}_0^{\pm} := E_0^{\pm} B_0^{-1} P_{B_0 \mathcal{H}},$$

where $P_{B_0\mathcal{H}}$ denotes the projection onto $B_0\mathcal{H}$ in the topological splitting $L_2 = B_0\mathcal{H} \oplus \mathcal{H}^{\perp}$ and B_0^{-1} is the inverse of $B_0 : \mathcal{H} \to B_0\mathcal{H}$. Beware that B_0^{-1} is not necessarily a multiplication operator and is only defined on the subspace $B_0\mathcal{H}$. Note also that unlike E_0^{\pm} , \widehat{E}_0^{\pm} are not projections.

Proposition 4.4. If $f \in L_2^{\text{loc}}(\mathbf{R}_+; \mathcal{H})$ satisfies $\partial_t f + DBf = 0$ in \mathbf{R}_+^{1+n} distributional sense, then

$$-\int_{0}^{t} \eta'_{+}(s)e^{-(t-s)|DB_{0}|}E_{0}^{+}f_{s}ds = \int_{0}^{t} \eta_{+}(s)DB_{0}e^{-(t-s)|DB_{0}|}\widehat{E}_{0}^{+}\mathcal{E}_{s}f_{s}ds,$$
$$-\int_{t}^{\infty} \eta'_{-}(s)e^{-(s-t)|DB_{0}|}E_{0}^{-}f_{s}ds = \int_{t}^{\infty} \eta_{-}(s)DB_{0}e^{-(s-t)|DB_{0}|}\widehat{E}_{0}^{-}\mathcal{E}_{s}f_{s}ds,$$

for all t > 0 and smooth bump functions $\eta_{\pm}(s) \ge 0$, where η_{+} is compactly supported in (0, t), and η_{-} is compactly supported in (t, ∞) .

Proof. In Section 3.1, we showed formally how to integrate the differential equation and arrived at (9). To make this rigorously, we proceed as follows. By assumption

(23)
$$\int_0^\infty \left(\left(-\partial_s \phi_s, f_s \right) + \left(D\phi_s, B_0 f_s \right) \right) ds = \int_0^\infty \left(D\phi_s, \mathcal{E}_s f_s \right) ds$$

for all $\phi \in C_0^{\infty}(\mathbf{R}^{1+n}; \mathbf{C}^{(1+n)m})$. To prove the identity on (0, t), let $\phi_0 \in \mathcal{H}$ be any boundary function and define $\phi_s := \eta_+(s)(e^{-(t-s)|DB_0|}E_0^+)^*\phi_0 \in C_0^{\infty}(\mathbf{R}_+; \mathsf{D}(D))$. To show that we can use this as test function, take $\eta \in C_0^{\infty}(\mathbf{R}^n)$ with $\eta = 1$ in a neighbourhood of x = 0 and $\int_{\mathbf{R}^n} \eta = 1$ and write $\eta_{\epsilon} := \epsilon^{-n}\eta(x/\epsilon)$ and

$$\phi_s^{R,\epsilon}(x) := \eta_+(s)\eta(x/R) \left(\eta_\epsilon * \left((e^{-|(t-s)DB_0|} E_0^{\pm})^* \phi_0 \right) \right)(x).$$

It is straightforward to verify that $\partial_s \phi^{R,\epsilon} \to \partial_s \phi$ and $D\phi^{R,\epsilon} \to D\phi$ in $L_2(\text{supp } \eta_+ \times \mathbf{R}^n; \mathbf{C}^{(1+n)m})$ when $R \to \infty, \epsilon \to 0$.

From (23) we obtain

$$\begin{split} \int_0^t (-\eta'_+(s)(e^{-(t-s)|DB_0|}E_0^+)^*\phi_0 - \eta_+(s)(DB_0e^{-(t-s)|DB_0|}E_0^+)^*\phi_0, f_s)ds \\ &+ \int_0^t (\eta_+(s)D(e^{-(t-s)|DB_0|}E_0^+)^*\phi_0, B_0f_s)ds \\ &= \int_0^t (\eta_+(s)D(e^{-(t-s)|DB_0|}E_0^+)^*\phi_0, \mathcal{E}_sf_s)ds. \end{split}$$

Since $B_0^* D(e^{-(t-s)|DB_0|}E_0^+)^* = (e^{-(t-s)|DB_0|}E_0^+ DB_0)^* = (DB_0 e^{-(t-s)|DB_0|}E_0^+)^*$, the last two terms on the left hand side cancel. Using that $E_0^+ D = E_0^+ DP_{B_0\mathcal{H}} = E_0^+ (DB_0)B_0^{-1}P_{B_0\mathcal{H}} = DB_0\widehat{E}_0^+$ on the right hand side, we have proved that

$$-\left(\phi_0, \int_0^t \eta'_+(s)e^{-(t-s)|DB_0|}E_0^+f_sds\right) = \left(\phi_0, \int_0^t \eta_+(s)e^{-(t-s)|DB_0|}E_0^+D\mathcal{E}_sf_sds\right).$$

Since this holds for all ϕ_0 , the (0, t) integral formula follows. The proof for the (t, ∞) integral formula is similar.

Our goal is to take limits to arrive at an integrated equation. Formally, if we let η_{\pm} approximate the characteristic functions for (0, t) and (t, ∞) respectively, we obtain in the limit from Proposition 4.4 that

$$E_0^+ f_t - e^{-t|DB_0|} E_0^+ f_0 = \int_0^t DB_0 e^{-(t-s)|DB_0|} \widehat{E}_0^+ \mathcal{E}_s f_s ds,$$
$$0 - E_0^- f_t = \int_t^\infty DB_0 e^{-(s-t)|DB_0|} \widehat{E}_0^- \mathcal{E}_s f_s ds,$$

if $\lim_{t\to 0} f_t = f_0$ and $\lim_{t\to\infty} f_t = 0$ in appropriate sense (and yet to be proved). Subtraction yields $f = C_0^+ f_0 + S_A f$, which we wish to solve as

(24)
$$f = (I - S_A)^{-1} C_0^+ f_0,$$

where the integral operator S_A is

(25)
$$S_A f_t = \int_0^t DB_0 e^{-(t-s)|DB_0|} \widehat{E}_0^+ \mathcal{E}_s f_s ds - \int_t^\infty DB_0 e^{-(s-t)|DB_0|} \widehat{E}_0^- \mathcal{E}_s f_s ds.$$

and C_0^+ is the generalized Cauchy integral defined via the semigroup $e^{-t|DB_0|}$.

The equation (24) can also be viewed as a generalized Cauchy integral formula, for t-dependent coefficients A, and we shall see that, given any $f_0 \in L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$, it constructs a solution f_t to the elliptic equation. However, for this one and also justification of the limiting arguments, one needs a suitable functional setting we now introduce.

5. NATURAL FUNCTION SPACES

It is well known that solutions g to (17) with L_2 boundary data typically satisfy certain square function estimates, as well as non-tangential maximal function estimates. In this section, we study the basic properties of some natural function spaces related to BVPs with L_2 boundary data.

Definition 5.1. In \mathbf{R}^{1+n}_+ , define the Banach/Hilbert spaces

$$\mathcal{X} := \{ f : \mathbf{R}_{+}^{1+n} \to \mathbf{C}^{(1+n)m} \; ; \; \widetilde{N}_{*}(f) \in L_{2}(\mathbf{R}^{n}) \}, \\ \mathcal{Y} := \{ f : \mathbf{R}_{+}^{1+n} \to \mathbf{C}^{(1+n)m} \; ; \; \int_{0}^{\infty} \|f_{t}\|_{2}^{2} t dt < \infty \},$$

with the obvious norms. Here \widetilde{N}_* denotes the modified non-tangential maximal function from Definition 2.1. By $\mathcal{Y}^* = L_2(\mathbf{R}^{1+n}_+, dt/t; \mathbf{C}^{(1+n)m})$ we denote the dual space of \mathcal{Y} , relative to $L_2(\mathbf{R}^{1+n}_+; \mathbf{C}^{(1+n)m})$.

In Sections 8 and 9 we demonstrate that the maximal function space \mathcal{X} is the natural space to solve the Neumann and regularity problems in, whereas \mathcal{Y} is natural for the Dirichlet problem. Natural is meant with respect to the method. That the spaces \mathcal{Y} and \mathcal{X} are relevant for L_2 BVPs with *t*-independent coefficients is clear from the following theorem. For proofs, we refer to [9, Proposition 2.3] and [7, Proposition 2.56].

Theorem 5.2. Let f_0 belong to the spectral subspace $E_0^+\mathcal{H}$. Then its generalized Cauchy extension $f = C_0^+ f_0$ as in Proposition 4.3 has estimates

$$\|\partial_t f\|_{\mathcal{Y}} \approx \|f\|_{\mathcal{X}} \approx \|f_0\|_2.$$

We will show in Corollary 8.4 that any distributional solution $f \in \mathcal{X}$ to $\partial_t f + DB_0 f = 0$ is the generalized Cauchy extension of some $f_0 \in E_0^+ \mathcal{H}$.

Clearly $\mathcal{Y} \subset L_2^{\text{loc}}(\mathbf{R}_+; L_2)$. The following lemma shows that \mathcal{X} is locally L_2 inside \mathbf{R}^{1+n}_+ as well, and is quite close to \mathcal{Y}^* .

Lemma 5.3. There are estimates

$$\sup_{t>0} \frac{1}{t} \int_{t}^{2t} \|f_s\|_2^2 ds \lesssim \|\widetilde{N}_*(f)\|_2^2 \lesssim \int_0^\infty \|f_s\|_2^2 \frac{ds}{s}$$

In particular $\mathcal{Y}^* \subset \mathcal{X}$.

Proof. The second inequality follows by integrating the pointwise estimate

$$\widetilde{N}_{*}(f)(x)^{2} \approx \sup_{t>0} \iint_{W(t,x)} |f(s,y)|^{2} \frac{dsdy}{s^{1+n}} \le \iint_{|y-x|< c_{0}c_{1}s} |f(s,y)|^{2} \frac{dsdy}{s^{1+n}}.$$

For the lower bound on $\|\tilde{N}_*(f)\|_2$, it suffices to estimate $t^{-1} \int_t^{c_0 t} \|f_s\|_2^2 ds$, uniformly for t > 0. To this end, split $\mathbf{R}^n = \bigcup_k Q_k$, where Q_k all are disjoint cubes with diagonal lengths $c_1 t$. Then

$$t^{-1} \int_{t}^{c_0 t} \int_{Q_k} |f(s,y)|^2 ds dy \lesssim |Q_k| \inf_{x \in Q_k} |\widetilde{N}_*(f)(x)|^2 \lesssim \int_{Q_k} |\widetilde{N}_*(f)(x)|^2 dx.$$

Summation over k gives the stated estimate.

The space \mathcal{Y}^* is a subspace of \mathcal{X} of functions with zero trace at the boundary \mathbf{R}^n , in the square L_2 -Dini sense $\lim_{t\to 0} t^{-1} \int_t^{2t} ||f_s||_2^2 ds = 0$. A fundamental quantity is the norm of multiplication operators mapping \mathcal{X} into \mathcal{Y}^* .

Definition 5.4. For functions $\mathcal{E} : \mathbf{R}^{1+n}_+ \to \mathcal{L}(\mathbf{C}^{(1+n)m})$, we denote

$$\|\mathcal{E}\|_* := \|\mathcal{E}\|_{\mathcal{X} \to \mathcal{Y}^*} = \sup_{\|f\|_{\mathcal{X}} = 1} \|\mathcal{E}f\|_{\mathcal{Y}^*}$$

the norm of pointwise multiplication by \mathcal{E} .

The following lemma gives a sufficient Carleson condition for a multiplication operator to map into this subspace.

Lemma 5.5. For functions $\mathcal{E} : \mathbf{R}^{1+n}_+ \to \mathcal{L}(\mathbf{C}^{(1+n)m})$, we have estimates $\|\mathcal{E}\|_{\infty} \lesssim \|\mathcal{E}\|_* \lesssim \|\mathcal{E}\|_C$,

where $\|\mathcal{E}\|_C$ denotes the modified Carleson norm from Definition 2.1.

Proof. For the first estimate, fix t and consider only f supported on (t, 2t) in the definition of $\|\mathcal{E}\|_{\mathcal{X}\to\mathcal{Y}^*}$. Lemma 5.3 shows that

$$\sup \|\mathcal{E}f\|_{\mathcal{Y}^*} / \|f\|_{\mathcal{X}} \approx \sup(t^{-1/2} \|\mathcal{E}f\|_2) / (t^{-1/2} \|f\|_2) = \operatorname{ess\,sup}_{t < s < 2t} \|\mathcal{E}_s\|_{\infty},$$

where the first two suprema are over all $0 \neq f \in L_2((t, 2t) \times \mathbf{R}^n; \mathbf{C}^{(1+n)m})$. Taking supremum over t shows the estimate $\|\mathcal{E}\|_{\infty} \lesssim \|\mathcal{E}\|_{*}$.

For the second estimate, we calculate

$$\begin{split} \|\mathcal{E}f\|_{\mathcal{Y}^*}^2 &\approx \iint_{\mathbf{R}^{1+n}_+} \left(\frac{1}{t^{1+n}} \iint_{W(t,x)} ds dy\right) |\mathcal{E}(t,x)f(t,x)|^2 \frac{dt dx}{t} \\ &\approx \iint_{\mathbf{R}^{1+n}_+} \left(\frac{1}{s^{1+n}} \iint_{W(s,y)} |\mathcal{E}(t,x)f(t,x)|^2 \frac{dt dx}{t}\right) ds dy \\ &\lesssim \iint_{\mathbf{R}^{1+n}_+} \left(\frac{1}{s} \sup_{W(s,y)} |\mathcal{E}|^2\right) \left(\frac{1}{s^{1+n}} \iint_{W(s,y)} |f(t,x)|^2 dt dx\right) ds dy \lesssim \|\mathcal{E}\|_C^2 \|f\|_{\mathcal{X}}^2, \end{split}$$
here the final estimate is by Carleson's theorem. \Box

where the final estimate is by Carleson's theorem.

We have not been able to prove that the $\|\cdot\|_*$ norm is equivalent to the modified Carleson norm, that is to prove the appropriate lower bound. It is however easy to see that the $\|\cdot\|_*$ norm dominates the standard Carleson norm $\|\cdot\|_c$. Indeed, choosing f as the characteristic function for the Carleson box $(0, l(Q)) \times Q$ (times a unit vector field) in the estimate $\|\mathcal{E}f\|_{\mathcal{Y}^*} \leq \|\mathcal{E}\|_* \|f\|_{\mathcal{X}}$, shows that

$$\|\mathcal{E}\|_c := \sup_Q \frac{1}{|Q|} \iint_{(0,l(Q))\times Q} |\mathcal{E}(t,x)|^2 \frac{dtdx}{t} \lesssim \|\mathcal{E}\|_*^2.$$

Furthermore, it is straightforward to see that the modified Carleson norm is dominated by the corresponding modified square Dini norm

$$\|\mathcal{E}\|_{C}^{2} \lesssim \int_{0}^{\infty} \sup_{c_{0}^{-1}t < s < c_{0}t} \|\mathcal{E}_{s}\|_{\infty}^{2} \frac{dt}{t}.$$

Proof of Lemma 2.2. We shall prove the lemma assuming only finiteness of the standard Carleson norm (hence the lemma also holds for the star norm). Thus we assume $||A - A_0||_c < \infty$ and do the rest of the proof replacing $|| \cdot ||_C$ by $|| \cdot ||_c$.

To prove uniqueness, we use $\|A'_0 - A_0\|_c \leq \|A - A'_0\|_c + \|A - A_0\|_c < \infty$ to obtain

$$\int_{0}^{l(Q)} \left(\frac{1}{|Q|} \int_{Q} |A'_{0}(x) - A_{0}(x)|^{2} dx \right) \frac{dt}{t} < \infty$$

for all cubes $Q \subset \mathbf{R}^n$, which only is possible if $A'_0 = A_0$ almost everywhere, by *t*-independence.

Similarly for $A - A_0$, we have $\int_0^{l(Q)} (|Q|^{-1} \int_Q |A(t,x) - A_0(x)|^2 dx) dt/t < \infty$ for any cube Q, and it follows that we have essential infimum

$$\operatorname{ess\,inf}_{0 < t < l(Q)} \frac{1}{|Q|} \int_{Q} |A(t, x) - A_0(x)|^2 dx = 0.$$

To prove $||A_0||_{\infty} \leq ||A||_{\infty}$, assume $\epsilon > 0$ and pick Q such that $(|Q|^{-1} \int_Q |A_0|^2 dx)^{1/2} \geq 1$ $||A_0||_{\infty} - \epsilon$. Then choose $t \in (0, l(Q))$ such that $|Q|^{-1} \int_Q |A(t, x) - A_0(x)|^2 dx \le \epsilon^2$. We assume that t is a Lebesgue point of $t \mapsto |Q|^{-1} \int_Q |A(t,x) - A_0(x)|^2 dx$ and of $t\mapsto |Q|^{-1}\int_Q |A(t,x)|^2 dx.$ This yields $||A_0||_{\infty} -\epsilon \leq (|Q|^{-1}\int_Q |A_0|^2 dx)^{1/2} \leq ||A||_{\infty} +\epsilon.$ Letting $\epsilon \to 0$ proves the claim.

To prove $\kappa_0 \geq \kappa$, assume $\epsilon > 0$ and pick $f \in \mathcal{H} \cap C_0^{\infty}(\mathbb{R}^n; \mathbb{C}^{(1+n)m})$. Pick Q such that supp $f \subset Q$. Then

$$\operatorname{Re}(A_0 f, f) = \operatorname{Re}(A_t f, f) + \operatorname{Re}((A_0 - A_t) f, f) \ge \\ \kappa \|f\|_2^2 - |Q| \|f\|_\infty^2 \left(|Q|^{-1} \int_Q |A(t, x) - A_0(x)|^2 dx \right)^{1/2},$$

for all $t \in (0, l(Q))$. Taking the essential supremum over such t gives $\operatorname{Re}(A_0 f, f) \geq \kappa \|f\|_2^2$. Since $\mathcal{H} \cap C_0^{\infty}(\mathbf{R}^n; \mathbf{C}^{(1+n)m})$ is dense in \mathcal{H} , taking infimum over f, this proves the claim. \Box

6. HOLOMORPHIC OPERATIONAL CALCULUS

Throughout this section Λ denotes a closed, densely defined ω -sectorial operator in an arbitrary Hilbert space \mathcal{H} , *i.e.* $\sigma(\Lambda) \subset S_{\omega+}$, and we assume resolvent bounds $\|(\lambda - \Lambda)^{-1}\|_{\mathcal{H}\to\mathcal{H}} \leq 1/\text{dist}(\lambda, S_{\omega+})$. For simplicity, we assume throughout that Λ is injective, and therefore has dense range. In our applications Λ will be $|DB_0|$, and \mathcal{H} will be the Hilbert space from (6). See Section 7.

The goal in this section is to develop the theory needed to make rigorous the limiting argument following Proposition 4.4. To this end, we study uniform boundedness and convergence of model operators

(26)
$$S_{\epsilon}^{+}f_{t} := \int_{0}^{t} \eta_{\epsilon}^{+}(t,s)\Lambda e^{-(t-s)\Lambda}f_{s}ds,$$

(27)
$$S_{\epsilon}^{-}f_{t} := \int_{t}^{\infty} \eta_{\epsilon}^{-}(t,s)\Lambda e^{-(s-t)\Lambda}f_{s}ds,$$

acting on functions $f_t(x) = f(t, x)$ in a Hilbert space $L_2(\mathbf{R}_+, d\mu(t); \mathcal{H})$. For uniform boundedness issues, it suffices that the bump functions $\eta_{\epsilon}^+(t, s)$ and $\eta_{\epsilon}^-(t, s)$ are uniformly bounded and compactly supported within $\{(s, t) ; 0 < s < t\}$ and $\{(s, t) ; 0 < t < s\}$ and $\{(s, t) ; 0 < t < s\}$ and $\{(s, t) ; 0 < t < s\}$ respectively. For convergence issues and to link to the ODE, they should approximate the characteristic functions of the above sets. A convenient choice which we shall use systematically is the following. Define $\eta^0(t)$ to be the piecewise linear continuous function with support $[1, \infty)$, which equals 1 on $(2, \infty)$ and is linear on (1, 2). Then let $\eta_{\epsilon}(t) := \eta^0(t/\epsilon)(1 - \eta^0(2\epsilon t))$ and

$$\eta_{\epsilon}^{\pm}(t,s) := \eta^{0}(\pm (t-s)/\epsilon)\eta_{\epsilon}(t)\eta_{\epsilon}(s).$$

We study the operators S_{ϵ}^{\pm} from the point of view of operational calculus. This means for example that we view $S_{\epsilon}^{+} = F(\Lambda)$ as obtained from the underlying operator Λ (acting horizontally, *i.e.* in the variable x) by applying the operator-valued function $\lambda \mapsto F(\lambda)$, where

$$(F(\lambda)f)_t := \int_0^t \eta_{\epsilon}^+(t,s)\lambda e^{-(t-s)\lambda} f_s ds,$$

which depends holomorphically on λ in a sector $S_{\nu+}^o$ containing the spectrum of Λ . Note that each of these vertically acting, *i.e.* acting in the *t*-variable, operators $F(\lambda)$ commute with Λ . 6.1. Operational calculus in Hilbert space. Consider Λ as above. Let $\mathcal{K} := L_2(\mathbf{R}_+, d\mu(t); \mathcal{H})$ for some Borel measure μ . We extend the resolvents $(\lambda - \Lambda)^{-1} \in \mathcal{L}(\mathcal{H}), \lambda \notin S_{\omega+}$, to bounded operators on \mathcal{K} (and we use the same notation, letting $((\lambda - \Lambda)^{-1}f)_t := (\lambda - \Lambda)^{-1}(f_t)$ for all $f \in \mathcal{K}$ and a.e. t > 0). These extensions of the resolvents to \mathcal{K} clearly inherit the bounds from \mathcal{H} . We may think of them as being the resolvents of an ω -sectorial operator $\Lambda = \Lambda_{\mathcal{K}}$, although this extended unbounded operator $\Lambda_{\mathcal{K}}$ is not needed below.

Define the commutant of Λ to be

$$\Lambda' := \{ T \in \mathcal{L}(\mathcal{K}) \ ; \ (\lambda - \Lambda)^{-1} T = T(\lambda - \Lambda)^{-1} \text{ for } \lambda \notin S_{\omega +} \}.$$

Fix $\omega < \nu < \pi/2$, and consider classes of operator-valued holomorphic functions

$$H(S_{\nu+}^{o}; \Lambda') := \{ \text{holomorphic } F : S_{\nu+}^{o} \to \Lambda' \},$$

$$\Psi(S_{\nu+}^{o}; \Lambda') := \{ F \in H(S_{\nu+}^{o}; \Lambda') \; ; \; \|F(\lambda)\| \lesssim \min(|\lambda|^{a}, |\lambda|^{-a}), \text{ some } a > 0 \},$$

$$H_{\infty}(S_{\nu+}^{o}; \Lambda') := \{ F \in H(S_{\nu+}^{o}; \Lambda') \; ; \; \sup_{\lambda \in S_{\nu+}^{o}} \|F(\lambda)\| < \infty \}.$$

Through Dunford calculus, we define for $F \in \Psi(S_{\nu+}^o; \Lambda')$ the operator

(28)
$$F(\Lambda) := \frac{1}{2\pi i} \int_{\gamma} F(\lambda) (\lambda - \Lambda)^{-1} d\lambda,$$

where γ is the unbounded contour $\{re^{\pm i\theta}; r > 0\}, \omega < \theta < \nu$, parametrized counter clockwise around $S_{\omega+}$. This yields a bounded operator $F(\Lambda)$, since the bounds on F and the resolvents guarantee that the integral converges absolutely.

Remark 6.1. Functional calculus of the operator Λ is a special case of this operational calculus (28). Applying a scalar holomorphic function $f(\lambda)$ to Λ with functional calculus is the same as applying the operator-valued holomorphic function $F(\lambda) = f(\lambda)I$ to Λ with operational calculus. For the functional calculus, we write $\Psi(S_{\nu+}^o)$ and $H_{\infty}(S_{\nu+}^o)$ for the corresponding classes of scalar symbol functions.

We also remark that a more general functional and operational calculi for bisectorial operators like DB_0 are developed entirely similar to those of sectorial operators Λ , replacing the sector $S_{\omega+}$ by the bisector S_{ω} .

The following three propositions contain all the theory of operational calculus that we need. To be self-contained and illustrate their simplicity, we give full proofs, although the propositions are proved in exactly the same way as for functional calculus, and can be found in [2].

Proposition 6.2. If $F, G \in \Psi(S_{\nu+}^o; \Lambda')$, then

$$F(\Lambda)G(\Lambda) = (FG)(\Lambda).$$

Note that we need not assume that $F(\lambda)$ and $G(\mu)$ commute for any $\lambda, \mu \in S^o_{\nu+}$.

Proof. We use contours γ_1 and γ_2 , with angles $\omega < \theta_1 < \theta_2 < \pi/2$, so that γ_2 encircles γ_1 . Cauchy's theorem now yields

$$(2\pi i)^{2} F(\Lambda) G(\Lambda) = \left(\int_{\gamma_{1}} \frac{F(\lambda)}{\lambda - \Lambda} d\lambda \right) \left(\int_{\gamma_{2}} \frac{G(\mu)}{\mu - \Lambda} d\mu \right)$$

$$= \int_{\gamma_{1}} \int_{\gamma_{2}} F(\lambda) G(\mu) \frac{1}{\mu - \lambda} \left(\frac{1}{\lambda - \Lambda} - \frac{1}{\mu - \Lambda} \right) d\lambda d\mu$$

$$= \int_{\gamma_{1}} \frac{F(\lambda)}{\lambda - \Lambda} \left(\int_{\gamma_{2}} \frac{G(\mu)}{\mu - \lambda} d\mu \right) d\lambda - \int_{\gamma_{2}} \left(\int_{\gamma_{1}} \frac{F(\lambda)}{\mu - \lambda} d\lambda \right) \frac{G(\mu)}{\mu - \Lambda} d\mu$$

$$= \int_{\gamma_{1}} \frac{F(\lambda)}{\lambda - \Lambda} 2\pi i G(\lambda) d\lambda - 0 = (2\pi i)^{2} (FG)(\Lambda),$$

sing the resolvent equation.

using the resolvent equation.

Proposition 6.3. Assume that Λ satisfies square function estimates, i.e. assume that

$$\int_0^\infty \|\psi(t\Lambda)u\|_{\mathcal{H}}^2 \frac{dt}{t} \approx \|u\|_{\mathcal{H}}^2, \qquad \text{for all } u \in \mathcal{H}$$

and some fixed $\psi \in \Psi(S^o_{\nu+})$. Then there exists $C < \infty$ such that

$$\|F(\Lambda)\| \le C \sup_{\lambda \in S_{\nu+}^o} \|F(\lambda)\|, \quad \text{for all } F \in \Psi(S_{\nu+}^o; \Lambda').$$

We remark that if square function estimates for Λ hold with one such ψ , then they hold for any non-zero $\psi \in \Psi(S_{\nu+}^o)$.

Proof. Note that the square function estimates extend to $u \in \mathcal{K}$, with $\|\cdot\|_{\mathcal{K}}$ instead of $\|\cdot\|_{\mathcal{H}}$. We drop \mathcal{K} in $\|\cdot\|_{\mathcal{K}}$. Using the resolution of identity $\int_0^\infty \psi^2(s\Lambda) u ds/s = cu$, where $0 < c < \infty$ is a constant, and the square function estimates, we calculate

$$\begin{split} \|F(\Lambda)u\|^2 &\approx \int_0^\infty \|\psi(t\Lambda)F(\Lambda)u\|^2 \frac{dt}{t} \\ &\approx \int_0^\infty \left\|\int_0^\infty (\psi(t\Lambda)F(\Lambda)\psi(s\Lambda))(\psi(s\Lambda)u)\frac{ds}{s}\right\|^2 \frac{dt}{t} \\ &\lesssim \sup_{S_{\nu+}^o} \|F(\lambda)\|^2 \int_0^\infty \left(\int_0^\infty \eta(t/s)\frac{ds}{s}\right) \left(\int_0^\infty \eta(t/s)\|\psi(s\Lambda)u\|^2 \frac{ds}{s}\right) \frac{dt}{t} \\ &\lesssim \sup_{S_{\nu+}^o} \|F(\lambda)\|^2 \int_0^\infty \|\psi(s\Lambda)u\|^2 \frac{ds}{s} \lesssim \sup_{S_{\nu+}^o} \|F(\lambda)\|^2 \|u\|^2. \end{split}$$

We have used the estimate

$$\|\psi(t\Lambda)F(\Lambda)\psi(s\Lambda)\| \lesssim \int_{\gamma} \|F(\lambda)\| |\psi(t\lambda)\psi(s\lambda)\lambda^{-1}d\lambda| \lesssim \sup_{\lambda \in S_{\nu+}^o} \|F(\lambda)\|\eta(t/s),$$

where $\eta(x) := \min\{x^a, x^{-a}\}(1 + |\log x|)$ for some a > 0.

Proposition 6.4. Assume that Λ satisfies square function estimates as in Proposition 6.3. Let $F_n \in \Psi(S_{\nu+}^o; \Lambda')$, $n = 1, 2, \ldots$, satisfy $\sup_{n,\lambda} \|F_n(\lambda)\| < \infty$, and let $F \in H_{\infty}(S_{\nu+}^{o}; \Lambda')$. Assume that for each fixed $v \in \mathcal{K}$ and $\lambda \in S_{\nu+}^{o}$, we have strong

convergence $\lim_{n\to\infty} ||F_n(\lambda)v - F(\lambda)v|| = 0$. Then the operators $F_n(\Lambda)$ converge strongly to a bounded operator $F(\Lambda)$, *i.e.*

$$F_n(\Lambda)u \to F(\Lambda)u, \quad \text{for all } u \in \mathcal{K}, \text{ as } n \to \infty.$$

Proof. Since $\sup_n ||F_n(\Lambda)|| < \infty$ by Proposition 6.3, it suffices to consider $u = \psi(\Lambda)v$ for some fixed $\psi \in \Psi(S_{\nu+}^o) \setminus \{0\}$, since $\mathsf{R}(\psi(\Lambda))$ is dense in \mathcal{K} . From (28), we get

$$||F_n(\Lambda)u - F_m(\Lambda)u|| \lesssim \int_{\gamma} ||(F_n(\lambda) - F_m(\lambda))v|| |\psi(\lambda)\lambda^{-1}d\lambda|,$$

where $||(F_n(\lambda) - F_m(\lambda))v|| \leq ||v||$ and $|\psi(\lambda)|/|\lambda|$ is integrable. The dominated convergence theorem applies and proves the proposition.

Propositions 6.2, 6.3 and 6.4 show that we have a continuous Banach algebra homomorphism

$$H_{\infty}(S^o_{\nu+};\Lambda') \to \mathcal{L}(\mathcal{K}): F \mapsto F(\Lambda),$$

provided that Λ satisfies square function estimates as in Proposition 6.3. This is the operational calculus that we need. Note that with $F(\Lambda)$ defined in this way for all $F \in H_{\infty}(S_{\nu+}^{o}; \Lambda')$, Proposition 6.4 continues to hold for any $F_n \in H_{\infty}(S_{\nu+}^{o}; \Lambda')$.

6.2. Maximal regularity estimates. Here, we apply the operational calculus from Section 6.1 to prove weighted bounds on the operators S_{ϵ}^{\pm} from (26) and (27).

Theorem 6.5. The operators S_{ϵ}^+ are uniformly bounded and converge strongly as $\epsilon \to 0$ on the weighted space $L_2(t^{\alpha}dt; \mathcal{H})$ if $\alpha < 1$. The operators S_{ϵ}^- are uniformly bounded and converge strongly as $\epsilon \to 0$ on the weighted space $L_2(t^{\alpha}dt; \mathcal{H})$ if $\alpha > -1$.

Note that the case $\alpha = 0$ is the usual maximal regularity result in $L_2(dt; \mathcal{H})$. The methods here provide a proof of it.

To establish boundedness of the integral operators $F(\lambda)$, we rely on the following version of Schur's lemma. The proof is straightforward using Cauchy–Schwarz' inequality.

Lemma 6.6. Consider the integral operator $f_t \mapsto \int_0^\infty k(t,s) f_s ds$, with C-valued kernel k(t,s). If the kernel has the bounds

$$\sup_{t} \frac{1}{t^{\beta_{2}-\alpha}} \int_{0}^{\infty} |k(t,s)| s^{\beta_{1}} ds = C_{1} < \infty, \qquad \sup_{s} \frac{1}{s^{\beta_{1}+\alpha}} \int_{0}^{\infty} |k(t,s)| t^{\beta_{2}} dt = C_{2} < \infty,$$

for some $\beta_1, \beta_2 \in \mathbf{R}$, then the integral operator is bounded on $L_2(t^{\alpha}dt; \mathcal{H})$ with norm at most $\sqrt{C_1C_2}$.

The second result that we need shows that when the integral operators $F(\lambda)$ define a holomorphic function in $\Psi(S_{\nu+}^o; \mathcal{L}(\mathcal{K}))$, then the resulting operator $F(\Lambda)$ can be represented as an integral operator with operator-valued kernel.

Lemma 6.7. Consider a family of integral operators $F(\lambda)f_t = \int_0^\infty k_\lambda(t,s)f_s ds$ such that the **C**-valued kernels have the bounds

$$\sup_{t} \frac{1}{t^{\beta_2 - \alpha}} \int_0^\infty |k_\lambda(t, s)| s^{\beta_1} ds \le \eta(\lambda), \qquad \sup_{s} \frac{1}{s^{\beta_1 + \alpha}} \int_0^\infty |k_\lambda(t, s)| t^{\beta_2} dt \le \eta(\lambda).$$

If $\sup_{\lambda \in S_{\nu+}^o} \eta(\lambda) < \infty$, if $\lambda \mapsto k_{\lambda}(t,s)$ is holomorphic in $S_{\nu+}^o$ for a.e. (t,s), and if $\iint_K |\partial_{\lambda}k_{\lambda}(t,s)| dtds$ is locally bounded in λ for each compact set K, then $F \in H_{\infty}(S_{\nu+}^o; \mathcal{L}(L_2(t^{\alpha}dt; \mathcal{H})))$. If furthermore $\eta(\lambda) \leq \min(|\lambda|^a, |\lambda|^{-a})$ for $\lambda \in S^o_{\nu+}$ and some a > 0, then $F \in \Psi(S^o_{\nu+}; \mathcal{L}(L_2(t^{\alpha}dt; \mathcal{H})))$, and

$$F(\Lambda)f_t = \int_0^\infty k_\Lambda(t,s)f_s ds, \qquad \text{for all } f \in L_2(t^\alpha dt;\mathcal{H}) \text{ and a.e. } t,$$

where the operator-valued kernels $k_{\Lambda}(t,s)$ are defined through (28) for a.e. (t,s).

Proof. Schur's lemma 6.6 provides the bounds on $F(\lambda)$. To show that the operatorvalued function F is holomorphic, by local boundedness it suffices to show that the scalar function

$$\lambda \mapsto \iint (h_t, k_\lambda(t, s) f_s) ds dt$$

is holomorphic, for all bounded and compactly supported f, h. The hypothesis on $\partial_{\lambda}k_{\lambda}(t,s)$ guarantees this.

To prove the representation formula for $F(\Lambda)$, it suffices to show that for each $f \in L_2(t^{\alpha}dt; \mathcal{H}), v \in \mathcal{H}$, and a.e. t, changing order of integration is possible in

$$\iint (v, k_{\lambda}(t, s)(\lambda - \Lambda)^{-1} f_s) ds d\lambda.$$

Thus, by Fubini, one needs to show

$$\iint |k_{\lambda}(t,s)| \|f_s\| ds \frac{|d\lambda|}{|\lambda|} < \infty, \quad \text{for a.e. } t.$$

The bounds on $k_{\lambda}(t,s)$ in the hypothesis guarantee this.

Proof of Theorem 6.5. Since S_{ϵ}^+ in $L_2(t^{\alpha}dt; \mathcal{H})$ and S_{ϵ}^- in $L_2(t^{-\alpha}dt; \mathcal{H})$, with Λ replaced by Λ^* , are adjoint operators, it suffices to consider S_{ϵ}^+ . Let

$$F_{\epsilon}(\lambda)f_t := \int_0^t \eta_{\epsilon}^+(t,s)\lambda e^{-(t-s)\lambda} f_s ds.$$

Uniform boundedness of the integral operators $F_{\epsilon}(\lambda)$ follows from Lemma 6.6 with $\beta_1 = -\alpha, \beta_2 = 0$, using the estimate $\int_0^y e^x x^{-\alpha} dx \leq e^y y^{-\alpha}$, which holds if and only if $\alpha \in (-\infty, 1)$. Indeed, since $\lambda \in S_{\nu+}^o$ with $\nu < \pi/2$, we have $\lambda_1 := \operatorname{Re} \lambda \approx |\lambda|$ and

$$\int_0^t |\lambda e^{-\lambda(t-s)}| s^{-\alpha} ds \approx \int_0^t \lambda_1 e^{-\lambda_1(t-s)} s^{-\alpha} ds = \lambda_1^{\alpha} e^{-\lambda_1 t} \int_0^{\lambda_1 t} e^x x^{-\alpha} dx \lesssim t^{-\alpha},$$

Similarly, $\int_{s}^{\infty} |\lambda e^{-\lambda(t-s)}| dt \lesssim e^{\lambda_{1}s} \int_{\lambda_{1}s}^{\infty} e^{-x} dx = 1.$

Again using Lemma 6.6, we note for fixed $\epsilon > 0$ the crude estimate $||F_{\epsilon}(\lambda)|| \lesssim |\lambda|e^{-\epsilon \operatorname{Re}\lambda}$, and with Lemma 6.7 we verify that $F_{\epsilon} \in \Psi(S_{\nu+}^{o}; \mathcal{L}(L_{2}(t^{\alpha}dt; \mathcal{H}))))$, and

$$F_{\epsilon}(\Lambda)f_t = \int_0^t \eta_{\epsilon}^+(t,s)\Lambda e^{-(t-s)\Lambda}f_s ds = S_{\epsilon}^+f_t, \quad \text{for a.e. } t.$$

To prove strong convergence, by Proposition 6.4 it suffices to show strong convergence of $F_{\epsilon}(\lambda)$ as $\epsilon \to 0$, for fixed $\lambda \in S_{\nu+}^{o}$. By uniform boundedness of $F_{\epsilon}(\lambda)$, it suffices to show that $F_{\epsilon}(\lambda)f$ converges in $L_2(t^{\alpha}dt;\mathcal{H})$ as $\epsilon \to 0$ for each f in the dense set $\bigcup_{\delta>0} L_2((\delta, \delta^{-1}), t^{\alpha}dt;\mathcal{H})$. This will follow from norm convergence of $F_{\epsilon}(\lambda)$ in $\mathcal{L}(L_2((\delta, \delta^{-1}), t^{\alpha}dt;\mathcal{H}), L_2(t^{\alpha}dt;\mathcal{H}))$ for each fixed $\delta > 0$. To see this, we

use Lemma 6.6 with $\beta_1 = -\alpha$ and $\beta_2 = 0$. As above C_1 is uniformly bounded. One verifies decay to 0 as $\epsilon \to 0$ of

$$\sup_{s \in (\delta, \delta^{-1})} \int_{(2\epsilon)^{-1}}^{\infty} \lambda_1 e^{-(t-s)\lambda_1} dt \quad \text{and} \quad \sup_{s \in (\delta, \delta^{-1})} \int_s^{s+2\epsilon} \lambda_1 e^{-(t-s)\lambda_1} dt.$$

This shows that $C_2 \to 0$ as $\epsilon \to 0$, which proves the strong convergence and the theorem.

6.3. Endpoint cases. The operators S_{ϵ}^{-} are not uniformly bounded on $L_{2}(t^{\alpha}dt; \mathcal{H})$ when $\alpha \leq -1$, and therefore no limit operator S^{-} exists in these spaces. Indeed, if $\eta(t)$ is a smooth approximation to the Dirac delta at t = 1 and $f \in \mathcal{H}$, then $S_{\epsilon}^{-}(\eta f)_{t}$ is independent of ϵ for $\epsilon < t/2$, with non-zero value $\approx \Lambda e^{-\Lambda} f \in \mathcal{H}$ for $t \approx 0$. Thus $\sup_{\epsilon>0} \int_{0}^{\infty} \|S_{\epsilon}^{-}(\eta f)_{t}\|_{\mathcal{H}}^{2} t^{\alpha} dt = \infty$ if $\alpha \leq -1$. By duality S_{ϵ}^{+} cannot be uniformly bounded when $\alpha \geq 1$.

In this section we study the endpoint cases $\alpha = \pm 1$. It is convenient here (and as we apply these abstract results in the subsequent paper [6]) to introduce the abstract spaces $Y := L_2(tdt; \mathcal{H})$ and $Y^* := L_2(dt/t; \mathcal{H})$. Note that they differ from $\mathcal{Y}, \mathcal{Y}^*$ by the target space \mathcal{H} being here an arbitrary Hilbert space. To obtain a uniform boundedness result for S_{ϵ}^- , assume there exists an auxiliary Banach space X with continuous embeddings

(29)
$$Y^* \subset X \subset L_2^{\text{loc}}(dt; \mathcal{H}),$$

i.e. $\int_a^b \|f_t\|_{\mathcal{H}}^2 dt \lesssim \|f\|_X^2 \lesssim \int_0^\infty \|f_t\|_{\mathcal{H}}^2 dt/t$ hold for each fixed $0 < a < b < \infty$, and such that the map $u \mapsto (e^{-t\Lambda}u)_{t>0}$ is bounded $\mathcal{H} \to X$, *i.e.*

(30)
$$||e^{-t\Lambda}u||_X \lesssim ||u||_{\mathcal{H}}, \quad \text{for all } u \in \mathcal{H}.$$

Theorem 6.8. Consider the model operators S_{ϵ}^+ and S_{ϵ}^- from (26-27) and Y, Y^{*} and X as above.

The operators S_{ϵ}^+ are uniformly bounded on Y^* and converge strongly to a limit operator $S^+ \in \mathcal{L}(Y^*, Y^*)$ as $\epsilon \to 0$.

The operators S_{ϵ}^{-} are uniformly bounded $Y^* \to X$, and there is a limit operator $S^{-} \in \mathcal{L}(Y^*, X)$ such that $\lim_{\epsilon \to 0} ||S_{\epsilon}^{-}f - S^{-}f||_{L_2(a,b;\mathcal{H})} = 0$ for any fixed $0 < a < b < \infty$ and $f \in Y^*$.

For the proof, we shall need the first part of the following lemma. The second part will be required in Propositions 7.1 and 7.2 below.

Lemma 6.9. The operators

$$\int_0^\infty \eta_\epsilon(s)\Lambda e^{-s\Lambda} f_s ds: Y^* \to \mathcal{H}$$

are bounded, uniformly in ϵ , and converge strongly as $\epsilon \to 0$. Let $U_s : \mathcal{H} \to \mathcal{H}$ be bounded operators such that $\mathcal{H} \to Y^*; h \mapsto (U_s^* e^{-s\Lambda^*} h)_{s>0}$ is bounded. Then the operators

$$\int_0^\infty \eta_\epsilon(s) e^{-s\Lambda} U_s f_s ds : Y \to \mathcal{H}$$

are bounded, uniformly in ϵ , and converge strongly as $\epsilon \to 0$.

Proof. For the first operator, square function estimates for Λ^* give

$$\left\|\int_0^\infty \eta_\epsilon(s)\Lambda e^{-s\Lambda}f_s ds\right\|_{\mathcal{H}} = \sup_{\|h\|_2=1} \left|\int_0^\infty (s\Lambda^* e^{-s\Lambda^*}h, f_s)\eta_\epsilon(s)\frac{ds}{s}\right| \lesssim \|\eta_\epsilon f\|_{Y^*} \lesssim \|f\|_{Y^*}.$$

For the second operator

$$\begin{split} \left\| \int_0^\infty \eta_\epsilon(s) e^{-s\Lambda} U_s f_s ds \right\|_{\mathcal{H}} &\lesssim \sup_{\|h\|_2 = 1} \left| \int_0^\infty (U_s^* e^{-s\Lambda^*} h, f_s) \eta_\epsilon(s) ds \right| \\ &\lesssim \sup_{\|h\|_2 = 1} \| U_s^* e^{-s\Lambda^*} h \|_{Y^*} \| \eta_\epsilon f \|_Y \lesssim \| \eta_\epsilon f \|_Y \lesssim \| f \|_Y, \end{split}$$

where in the second last estimate the hypothesis is used. (Note that the \mathcal{H} -bound on U_s is not used quantitatively.)

To see the strong convergence, replace η_{ϵ} by $\eta_{\epsilon} - \eta_{\epsilon'}$ and use the dominated convergence theorem.

Proof of Theorem 6.8. The result for S_{ϵ}^+ is contained in Theorem 6.5, so it suffices to consider S_{ϵ}^- . Write

$$(31) \quad S_{\epsilon}^{-}f_{t} = \int_{t}^{\infty} \eta_{\epsilon}^{-}(t,s)\Lambda e^{-(s-t)\Lambda}f_{s}ds = \int_{t}^{\infty} \eta_{\epsilon}^{-}(t,s)\Lambda (e^{-(s-t)\Lambda} - e^{-(s+t)\Lambda})f_{s}ds - \int_{0}^{t+2\epsilon} (\eta_{\epsilon}(t)\eta_{\epsilon}(s) - \eta_{\epsilon}^{-}(t,s))\Lambda e^{-(s+t)\Lambda}f_{s}ds + \eta_{\epsilon}(t)e^{-t\Lambda}\int_{0}^{\infty} \eta_{\epsilon}(s)\Lambda e^{-s\Lambda}f_{s}ds = I - II + III.$$

We show that it is only the last term which is singular in the sense that it is not uniformly bounded on Y^* . Consider the term I and the symbol $F_{\epsilon}^{I}(\lambda)u_t = \int_{t}^{\infty} \eta_{\epsilon}^{-}(t,s)k_{\lambda}(t,s)u_sds$, where $k_{\lambda}(t,s) := \lambda e^{-(s-t)\lambda}(1-e^{-2t\lambda})$. Boundedness of $F_{\epsilon}^{I}(\lambda)$ on Y^* , uniformly in λ and ϵ follows from Lemma 6.6 and the estimates $\int_{t}^{\infty} |k_{\lambda}(t,s)| sds \lesssim t$ and $\int_{0}^{s} |k_{\lambda}(t,s)| dt \lesssim 1$. For example

$$\int_{t}^{\infty} |k_{\lambda}(t,s)| s ds \lesssim \min(1,t\lambda_{1})e^{t\lambda_{1}} \int_{t}^{\infty} \lambda_{1}e^{-s\lambda_{1}} s ds = t\min(1,t\lambda_{1})(1+1/(t\lambda_{1})) \lesssim t,$$

with $\lambda_1 := \operatorname{Re} \lambda$. On the other hand, for fixed $\epsilon > 0$, it is straightforward to verify with Lemma 6.6 that $\|F_{\epsilon}^{I}(\lambda)\|_{Y^* \to Y^*} \lesssim |\lambda| e^{-\epsilon \operatorname{Re} \lambda}$, and with Lemma 6.7 that $F_{\epsilon}^{I} \in \Psi(S_{\nu+}^{o}; \mathcal{L}(Y^*))$ and

$$F_{\epsilon}^{I}(\Lambda)f_{t} = \int_{t}^{\infty} \eta_{\epsilon}^{-}(t,s)\Lambda(e^{-(s-t)\Lambda} - e^{-(s+t)\Lambda})f_{s}ds, \quad \text{for a.e. } t$$

To prove strong convergence, as in the proof of Theorem 6.5, by uniform boundedness it suffices to show norm convergence of $F_{\epsilon}^{I}(\lambda)$ in $\mathcal{L}(L_{2}((\delta, \delta^{-1}), t^{-1}dt; \mathcal{H}), Y^{*})$ for each fixed $\delta > 0$. This follows from Lemma 6.6, where one verifies decay to 0 as $\epsilon \to 0$ of $\sup_{s \in (\delta, \delta^{-1})} \int_{0}^{2\epsilon} |k_{\lambda}(t, s)| dt$ and $\sup_{s \in (\delta, \delta^{-1})} \int_{s-2\epsilon}^{s} |k_{\lambda}(t, s)| dt$, and hence $C_{2} \to 0$, for fixed $\lambda \in S_{\nu+}^{o}$. Together with the uniform bound $\sup_{t} t^{-1} \int_{t}^{\infty} |k_{\lambda}(t, s)| s ds < \infty$, this proves the strong convergence for the term I.

Consider next the term II and the symbol

$$F_{\epsilon}^{II}(\lambda)u_t = \int_0^{t+2\epsilon} (\eta_{\epsilon}(t)\eta_{\epsilon}(s) - \eta_{\epsilon}^-(t,s))\lambda e^{-(s+t)\lambda}u_s ds$$

Boundedness of $F_{\epsilon}^{II}(\lambda)$ on Y^* , uniformly in λ and ϵ follows from Lemma 6.6 and the estimates $\int_0^{3t} |\lambda e^{-(s+t)\lambda}| sds \lesssim t$ and $\int_{s/3}^{\infty} |\lambda e^{-(s+t)\lambda}| dt \lesssim 1$. On the other hand, for fixed $\epsilon > 0$, we verify with Lemma 6.6 that $\|F_{\epsilon}^{II}(\lambda)\|_{Y^* \to Y^*} \lesssim |\lambda| e^{-\epsilon \operatorname{Re} \lambda}$, and with Lemma 6.7 that $F_{\epsilon}^{II} \in \Psi(S_{\nu+}^o; \mathcal{L}(Y^*))$ and

$$F_{\epsilon}^{II}(\Lambda)f_t = \int_0^{t+2\epsilon} (\eta_{\epsilon}(t)\eta_{\epsilon}(s) - \eta_{\epsilon}^-(t,s))\Lambda e^{-(s+t)\Lambda}u_s ds, \quad \text{for a.e. } t.$$

With the same technique as for the term I, the strong convergence of the term II follows from the decay to 0 as $\epsilon \to 0$ of $\sup_{s \in (\delta, \delta^{-1})} \int_{s-2\epsilon}^{s} |\lambda e^{-(s+t)\lambda}| dt$.

It remains to estimate the principal term III. Since the variables t and s separate, we can factor this term through the boundary space \mathcal{H} as a composition $Y^* \to \mathcal{H} \to X$, where Lemma 6.9 and the assumed bounds $e^{-t\Lambda} : \mathcal{H} \to X$ prove boundedness, uniform in ϵ , as well as strong convergence as maps $Y^* \to \mathcal{H} \to L_2(a, b; \mathcal{H})$. This completes the proof.

7. Estimates of the integral operators S_A and S_A

Let us come back to our concrete situation. Consider the operator DB_0 from Section 4. We set $\Lambda = |DB_0| := DB_0 \operatorname{sgn}(DB_0)$ on $\mathcal{H} = \overline{\mathsf{R}}(D)$, and see that Λ satisfies the assumptions of Section 6. It is a closed, densely defined, injective operator with $\sigma(\Lambda) \subset S_{\omega+}$ and $||(\lambda - \Lambda)^{-1}||_{\mathcal{H} \to \mathcal{H}} \leq 1/\operatorname{dist}(\lambda, S_{\omega+})$ (this follows from the resolvent bounds on DB_0). We apply the abstract theory from Section 6.3 to this Λ and spaces $Y^* := \mathcal{Y}^* \cap L_2^{\operatorname{loc}}(\mathbf{R}_+; \mathcal{H}), X := \mathcal{X} \cap L_2^{\operatorname{loc}}(\mathbf{R}_+; \mathcal{H})$ and $Y := \mathcal{Y} \cap L_2^{\operatorname{loc}}(\mathbf{R}_+; \mathcal{H})$. Note that the continuous embeddings (29) follow from Lemma 5.3 and the boundedness hypothesis (30) on $\mathcal{H} \to X : h \mapsto (e^{-t|DB_0|}h)_{t>0}$ follows from Theorem 5.2 (and the analogous result for the lower half space \mathbf{R}_-^{1+n} , *i.e.* $f_0 \in E_0^-\mathcal{H}$ giving a solution of $\partial_t f + DB_0 f = 0$ for t < 0.). We shall use the operational calculus of Λ to rigorously define and estimate the operator S_A in (25).

The strategy for the Dirichlet problem described in Section 3.2 leads us to consider the functional calculus of B_0D and the integral operator \tilde{S}_A from (16). If B_0 were invertible on all L_2 , then DB_0 and B_0D would be similar operators, but this is not the case in general. Still, whenever B_0 is strictly accretive on \mathcal{H} , it is true that B_0D is an ω -bisectorial operator with resolvent bounds. Furthermore, the L_2 space splits as

$$L_2 = B_0 \mathcal{H} \oplus \mathcal{H}^\perp$$

(cf. (22)) and B_0D restricts to an injective closed operator with dense range in $B_0\mathcal{H}$. This operator has square function estimates, and therefore bounded functional and operational calculus in $B_0\mathcal{H}$, as in Section 6.1. For proofs and further details, see [8]. We set $\tilde{\Lambda} := |B_0D|$ and $\tilde{E}_0^{\pm} := \chi^{\pm}(B_0D)$. We extend an operator $b(B_0D)$ in the functional calculus to an operator on all L_2 by letting $b(B_0D) = 0$ on $\mathcal{H}^{\perp} = \mathcal{N}(B_0D)$. With this notation $\tilde{E}_0^{\pm}(B_0\mathcal{H}) = \tilde{E}_0^{\pm}L_2$, and we shall prefer the latter to ease the notation.

A important relation between the functional calculus of DB_0 and B_0D is

(32)
$$B_0 b(DB_0) = b(B_0 D) B_0,$$

where we also extend operators $b(DB_0)$ to all L_2 , letting $b(DB_0)|_{\mathsf{N}(DB_0)} := 0$. The equation (32) clearly holds for resolvents $b(z) = (\lambda - z)^{-1}$. The general case follows from Dunford integration (28) and taking strong limits as in Proposition 6.4 (adapted

to bisectorial operators). Note that (32) in particular shows that for appropriate b and u

$$b(DB_0)Du = Db(B_0D)u.$$

A final observation is that with $\Lambda = |DB_0|$ and $\tilde{\Lambda} = |B_0D|$, then $\tilde{\Lambda}^* = |DB_0^*|$ and $\Lambda^* = |B_0^*D|$. So Λ and $\tilde{\Lambda}^*$ are of the same type, and the same holds for $\tilde{\Lambda}$ and Λ^* .

The boundedness result for the operator S_A is as follows.

Proposition 7.1. Assume that $\mathcal{E} : \mathbf{R}^{1+n}_+ \to \mathcal{L}(\mathbf{C}^{(1+n)m})$ satisfies $\|\mathcal{E}\|_* < \infty$, and define operators

$$S_A^{\epsilon}f_t := \int_0^t \eta_{\epsilon}^+(t,s)\Lambda e^{-(t-s)\Lambda}\widehat{E}_0^+ \mathcal{E}_s f_s ds + \int_t^\infty \eta_{\epsilon}^-(t,s)\Lambda e^{-(s-t)\Lambda}\widehat{E}_0^- \mathcal{E}_s f_s ds.$$

Then $||S_A^{\epsilon}||_{\mathcal{X}\to\mathcal{X}} \lesssim ||\mathcal{E}||_*$ and $||S_A^{\epsilon}||_{\mathcal{Y}\to\mathcal{Y}} \lesssim ||\mathcal{E}||_*$, uniformly for $\epsilon > 0$. In the space \mathcal{X} there is a limit operator $S_A = S_A^{\mathcal{X}} \in \mathcal{L}(\mathcal{X}; \mathcal{X})$ such that

$$\lim_{\epsilon \to 0} \|S_A^{\epsilon} f - S_A^{\mathcal{X}} f\|_{L_2(a,b;L_2)} = 0, \qquad \text{for any } f \in \mathcal{X}, 0 < a < b < \infty.$$

In the space \mathcal{Y} , there is a limit operator $S_A = S_A^{\mathcal{Y}} \in \mathcal{L}(\mathcal{Y}; \mathcal{Y})$ such that

$$\lim_{\epsilon \to 0} \|S_A^{\epsilon} f - S_A^{\mathcal{Y}} f\|_{\mathcal{Y}} = 0, \qquad \text{for any } f \in \mathcal{Y}.$$

Proof. The result on \mathcal{X} is a direct consequence of Theorem 6.8, since $\|\mathcal{E}f\|_{\mathcal{Y}^*} \leq \|\mathcal{E}\|_* \|f\|_{\mathcal{X}}$ and $S_A^{\epsilon} = S_{\epsilon}^+ \widehat{E}_0^+ \mathcal{E} + S_{\epsilon}^- \widehat{E}_0^- \mathcal{E}$. Note that $\mathsf{R}(\widehat{E}_0^{\pm}) \subset \mathcal{H} \subset L_2$.

Consider now the space \mathcal{Y} . The second term $S_{\epsilon}^{-}\widehat{E}_{0}^{-}\mathcal{E}$ is bounded on \mathcal{Y} , uniformly in ϵ , and converges strongly on \mathcal{Y} . This follows from Theorem 6.5 and the boundedness $\|\widehat{E}_{0}^{-}\mathcal{E}\|_{\mathcal{Y}\to\mathcal{Y}} \lesssim \|\mathcal{E}\|_{\infty} \lesssim \|\mathcal{E}\|_{*} < \infty$. The term $S_{\epsilon}^{+}\widehat{E}_{0}^{+}\mathcal{E}$ we split as

$$\int_{0}^{t} \eta_{\epsilon}^{+}(t,s)\Lambda e^{-(t-s)\Lambda} \widehat{E}_{0}^{+} \mathcal{E}_{s} f_{s} ds = \int_{0}^{t} \eta_{\epsilon}^{+}(t,s)\Lambda (e^{-(t-s)\Lambda} - e^{-(t+s)\Lambda}) \widehat{E}_{0}^{+} \mathcal{E}_{s} f_{s} ds$$
$$- \int_{t-2\epsilon}^{\infty} (\eta_{\epsilon}(t)\eta_{\epsilon}(s) - \eta_{\epsilon}^{+}(t,s))\Lambda e^{-(t+s)\Lambda} \widehat{E}_{0}^{+} \mathcal{E}_{s} f_{s} ds$$
$$+ \eta_{\epsilon}(t)\Lambda e^{-t\Lambda} \int_{0}^{\infty} \eta_{\epsilon}(s) e^{-s\Lambda} \widehat{E}_{0}^{+} \mathcal{E}_{s} f_{s} ds.$$

The result for the first two terms follows from the proof of Theorem 6.8 by duality, only using the boundedness of \mathcal{E} on \mathcal{Y} . For the last term, as the variables t and s split, it suffices to show uniform boundedness and convergence of

$$L_2 \to \mathcal{Y} : h \mapsto \eta_{\epsilon}(t) \Lambda e^{-t\Lambda} h$$

and

$$\mathcal{Y} \to L_2 : f_t \mapsto \int_0^\infty \eta_\epsilon(s) e^{-s\Lambda} \widehat{E}_0^+ \mathcal{E}_s f_s ds$$

separately. For the first operator, this follows directly from the square function estimates for Λ . To handle the second, it suffices to estimate $B_0 \int_0^\infty \eta_\epsilon(s) e^{-s\Lambda} \widehat{E}_0^+ \mathcal{E}_s f_s ds = \int_0^\infty \eta_\epsilon(s) e^{-s\Lambda} \widetilde{E}_0^+ \mathcal{E}_s f_s ds$, since B_0 is strictly accretive on $\mathcal{H} \supset \mathsf{R}(e^{-s\Lambda} \widehat{E}_0^+)$. To this end, we apply Lemma 6.9 with $U_s := \widetilde{E}_0^+ \mathcal{E}_s P_{\mathcal{H}}$, where $P_{\mathcal{H}}$ is orthogonal projection onto \mathcal{H} , and Λ replaced by $\tilde{\Lambda}$. The hypothesis there on boundedness of

$$\mathcal{H} \to Y^* : h \mapsto U_s^* e^{-s\bar{\Lambda}^*} h = P_{\mathcal{H}} \mathcal{E}_s^* e^{-s|DB_0^*|} \chi^+ (DB_0^*) h,$$

follows from the maximal estimate in Theorem 5.2 (with B_0 replaced by B_0^*), the assumed boundedness of $\mathcal{E}^* : \mathcal{X} \to \mathcal{Y}^*$ and L_2 boundedness of $\chi^+(DB_0^*)$ and $P_{\mathcal{H}}$. This completes the proof.

By inspection of the proofs above, the limit operator S_A , both for $f \in \mathcal{X}$ and $f \in \mathcal{Y}$, is seen to be

$$S_A f_t = \lim_{\epsilon \to 0} \left(\int_{\epsilon}^{t-\epsilon} \Lambda e^{-(t-s)\Lambda} \widehat{E}_0^+ \mathcal{E}_s f_s ds + \int_{t+\epsilon}^{\epsilon^{-1}} \Lambda e^{-(s-t)\Lambda} \widehat{E}_0^- \mathcal{E}_s f_s ds \right),$$

with convergence in $L_2(a, b; L_2)$ for any $0 < a < b < \infty$. This holds since we may equally well choose to work with the characteristic function $\eta^0(t) = \chi_{(1,\infty)}(t)$ instead of the piecewise linear function η^0 defined below (26-27). The only places we need the continuity of η^0 are in Theorems 8.2 and 9.2 below.

We now turn to defining rigorously the second integral operator needed for the Dirichlet problem. Write $C_b(X, V)$ for the space of bounded and continuous functions on X with values in V.

Proposition 7.2. The operators

$$\widetilde{S}_A^{\epsilon} f_t := \int_0^t \eta_{\epsilon}^+(t,s) e^{-(t-s)\tilde{\Lambda}} \widetilde{E}_0^+ \mathcal{E}_s f_s ds - \int_t^\infty \eta_{\epsilon}^-(t,s) e^{-(s-t)\tilde{\Lambda}} \widetilde{E}_0^- \mathcal{E}_s f_s ds$$

are bounded $\mathcal{Y} \to C_b(\overline{\mathbf{R}}_+; L_2)$, with $\sup_{t>0} \|\widetilde{S}_A^{\epsilon} f_t\|_2 \lesssim \|\mathcal{E}\|_* \|f\|_{\mathcal{Y}}$, uniformly for $\epsilon > 0$, and there is a limit operator $\widetilde{S}_A \in \mathcal{L}(\mathcal{Y}, C_b(\overline{\mathbf{R}}_+; L_2))$ such that $\lim_{\epsilon \to 0} \|\widetilde{S}_A^{\epsilon} f_t - \widetilde{S}_A f_t\|_2 = 0$ locally uniformly for $t \in (0, \infty)$, for any $f \in \mathcal{Y}$. The limit operator satisfies $S_A f = D\widetilde{S}_A f$ in \mathbf{R}_+^{1+n} distributional sense, where $S_A = S_A^{\mathcal{Y}}$ is the operator from Proposition 7.1, and has limits

$$\lim_{t \to 0} \|\widetilde{S}_A f_t - \tilde{h}^-\|_2 = 0 = \lim_{t \to \infty} \|\widetilde{S}_A f_t\|_2,$$

where $\tilde{h}^- := -\int_0^\infty e^{-s\tilde{\Lambda}} \widetilde{E}_0^- \mathcal{E}_s f_s ds \in \widetilde{E}_0^- L_2$, for any $f \in \mathcal{Y}$.

Note that $\widetilde{S}_A^{\epsilon} f_t = 0$ when $t \notin (\epsilon, \epsilon^{-1})$, so convergence $\widetilde{S}_A^{\epsilon} f_t \to \widetilde{S}_A f_t$ is not uniform up to t = 0. By inspection of the proof below, the limit operator is seen to be

(33)
$$\widetilde{S}_A f_t = \int_0^t e^{-(t-s)\tilde{\Lambda}} \widetilde{E}_0^+ \mathcal{E}_s f_s ds - \int_t^\infty e^{-(s-t)\tilde{\Lambda}} \widetilde{E}_0^- \mathcal{E}_s f_s ds$$

where the integrals are weakly convergent in L_2 for all $f \in \mathcal{Y}$ and t > 0.

Proof. The estimates for $\widetilde{S}^{\epsilon}_{A}$ are more straightforward than those for S^{ϵ}_{A} since there is no singularity at s = t. For the (0, t)-integral, split it as

$$\int_0^t \eta_{\epsilon}^+(t,s) e^{-(t-s)\tilde{\Lambda}} (I - e^{-2s\tilde{\Lambda}}) \widetilde{E}_0^+ \mathcal{E}_s f_s ds + e^{-t\tilde{\Lambda}} \int_0^t \eta_{\epsilon}^+(t,s) e^{-s\tilde{\Lambda}} \widetilde{E}_0^+ \mathcal{E}_s f_s ds.$$

For the first term, we write $e^{-(t-s)\tilde{\Lambda}}(I-e^{-2s\tilde{\Lambda}}) = \frac{s}{t-s}((t-s)\tilde{\Lambda}e^{-(t-s)\tilde{\Lambda}})((I-e^{-2s\tilde{\Lambda}})/(s\tilde{\Lambda}))$ to obtain the estimate $||e^{-(t-s)\tilde{\Lambda}}(I-e^{-2s\tilde{\Lambda}})|| \leq s/t$. From this uniform boundedness and convergence, locally uniformly in t, as $\epsilon \to 0$ follows by Cauchy–Schwarz inequality. For the second term we use uniform boundedness of $e^{-t\tilde{\Lambda}}$ and duality to estimate it by

$$\sup_{\|h\|_{2}=1} \left| \int_0^t (\mathcal{E}_s^* e^{-s\tilde{\Lambda}^*} (\widetilde{E}_0^+)^* h, f_s) \eta_{\epsilon}^+(t, s) ds \right| \lesssim \|\mathcal{E}^*\|_* \|f\|_{\mathcal{Y}},$$

using Lemma 6.9 as in the proof of Proposition 7.1. Moreover, the L_2 difference between the integral at ϵ and ϵ' is bounded by $\int_0^t ||f_s||_2^2 |\eta_{\epsilon}^+(t,s) - \eta_{\epsilon'}^+(t,s)|^2 s ds \to 0$ as $\epsilon, \epsilon' \to 0$ for fixed t, which proves the convergence.

The proof for the (t, ∞) -integral in S_A^{ϵ} is similar, splitting it as

$$\int_{t}^{\infty} \eta_{\epsilon}^{-}(t,s) e^{-(s-t)\tilde{\Lambda}} (I - e^{-2t\tilde{\Lambda}}) \widetilde{E}_{0}^{-} \mathcal{E}_{s} f_{s} ds + e^{-t\tilde{\Lambda}} \int_{t}^{\infty} \eta_{\epsilon}^{-}(t,s) e^{-s\tilde{\Lambda}} \widetilde{E}_{0}^{-} \mathcal{E}_{s} f_{s} ds,$$

and using the estimate $||e^{-(s-t)\tilde{\Lambda}}(I-e^{-2t\tilde{\Lambda}})|| \lesssim t/s$ for the first term and Lemma 6.9 for the second.

Since clearly $\widetilde{S}_A^{\epsilon} f \in C_b(\mathbf{R}_+; L_2)$, its locally uniform limit $\widetilde{S}_A f$ also belongs to $C_b(\mathbf{R}_+; L_2)$. To find the limits of $\widetilde{S}_A f_t$ at 0 and ∞ , since $\widetilde{S}_A : \mathcal{Y} \to C_b(\mathbf{R}_+; L_2)$ is bounded it suffices to consider $f \in \mathcal{Y}$ such that $f_t = 0$ for $t \notin (a, b)$, with $0 < a < b < \infty$ fixed but arbitrary. In this case,

$$\widetilde{S}_A f_t = \int_{a < s < \min(t,b)} e^{-(t-s)\tilde{\Lambda}} \widetilde{E}_0^+ \mathcal{E}_s f_s ds - \int_{\max(t,a) < s < b} e^{-(s-t)\tilde{\Lambda}} \widetilde{E}_0^- \mathcal{E}_s f_s ds$$

satisfies $\widetilde{E}_0^+ \widetilde{S}_A f_t = 0$ for t < a and $\widetilde{E}_0^- \widetilde{S}_A f_t = 0$ when t > b, from which the two limits $\lim_{t\to 0} \widetilde{E}_0^+ \widetilde{S}_A f_t = 0 = \lim_{t\to\infty} \widetilde{E}_0^- \widetilde{S}_A f_t$ follow. For the remaining two limits $\lim_{t\to\infty} \widetilde{E}_0^+ \widetilde{S}_A f_t$ and $\lim_{t\to 0} \widetilde{E}_0^- \widetilde{S}_A f_t$, we use that

$$\lim_{t \to \infty} \int_a^b \|e^{-(t-s)\tilde{\Lambda}} \widetilde{E}_0^+ \mathcal{E}_s f_s\|_2 ds = 0 = \lim_{t \to 0} \int_a^b \|(e^{-(s-t)\tilde{\Lambda}} - e^{-s\tilde{\Lambda}}) \widetilde{E}_0^- \mathcal{E}_s f_s\|_2 ds$$

by dominated convergence.

To verify the identity $S_A = D\widetilde{S}_A$, note that $\int_0^\infty (\phi_t, S_A^\epsilon f_t) dt = \int_0^\infty (D\phi_t, \widetilde{S}_A^\epsilon f_t) dt$ for all $f \in \mathcal{Y}$ and $\phi \in C_0^\infty(\mathbf{R}_+^{1+n}; \mathbf{C}^{(1+n)m})$. This uses the relations $B_0 \widehat{E}_0^{\pm} = B_0 E_0^{\pm} B_0^{-1} P_{B_0 \mathcal{H}} = \widetilde{E}_0^{\pm} P_{B_0 \mathcal{H}} = \widetilde{E}_0^{\pm}$ by (32). Let $\epsilon \to 0$ and use S_A^ϵ and \widetilde{S}_A^ϵ convergence. This completes the proof.

8. The Neumann and regularity problems

Throughout this section, A denotes t-dependent coefficients satisfying (2) and (3), and $A_0 \in L_{\infty}(\mathbf{R}^n; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ denotes t-independent coefficients which are strictly accretive on \mathcal{H} . We let $B := \hat{A}$ and $B_0 := \hat{A}_0$ be the transformed strictly accretive coefficients from Proposition 4.1, and define $\mathcal{E} := B_0 - B$.

For the Neumann and regularity problems, one seeks estimates of the gradient $g = \nabla_{t,x} u$ rather than the potential u. With a slight abuse of notation, we say below that g solves the divergence form equation when u does so.

Definition 8.1. By an \mathcal{X} -solution to the divergence form equation, with coefficients A, we mean a function $g \in L_2^{\text{loc}}(\mathbf{R}_+; L_2(\mathbf{R}^n; \mathbf{C}^{(1+n)m}))$, with estimate $\|\widetilde{N}_*(g)\|_2 < \infty$, which satisfies

$$\begin{cases} \operatorname{div}_{t,x} A g = 0, \\ \operatorname{curl}_{t,x} g = 0, \end{cases}$$

in \mathbf{R}^{1+n}_+ distributional sense.

Note that the boundary behaviour of g is not specified in this definition; we show existence of a limit in appropriate sense (see also Section 11). This will allow us to formulate in what sense the boundary data is prescribed. We first prove the following representation and regularity result for \mathcal{X} -solutions. Recall that $\Lambda = |DB_0|$. **Theorem 8.2.** Assume that $\|\mathcal{E}\|_* < \infty$. Then g is an \mathcal{X} -solution to the divergence form equation with coefficients A if and only if the corresponding conormal gradient $f = [(Ag)_{\perp}, g_{\parallel}]^t \in \mathcal{X}$ satisfies the equation

 $f_t = e^{-t\Lambda}h^+ + S_A f_t, \qquad for some \ h^+ \in E_0^+ \mathcal{H}.$

In this case, f has limits

(34)
$$\lim_{t \to 0} t^{-1} \int_{t}^{2t} \|f_s - f_0\|_2^2 ds = 0 = \lim_{t \to \infty} t^{-1} \int_{t}^{2t} \|f_s\|_2^2 ds$$

where $f_0 := h^+ + h^-$ and $h^- := \int_0^\infty \Lambda e^{-s\Lambda} \widehat{E}_0^- \mathcal{E}_s f_s ds \in E_0^- \mathcal{H}$, with estimates

$$\max(\|h^+\|_2, \|h^-\|_2) \approx \|f_0\|_2 \approx \|g_0\|_2 \lesssim \|f\|_{\mathcal{X}} \approx \|g\|_{\mathcal{X}}$$

The limits analogous to (34), replacing f by g and f_0 by $g_0 := [(B_0 f_0)_{\perp}, (f_0)_{\parallel}]^t$, hold. If furthermore $\|\mathcal{E}\|_*$ is sufficiently small, then there are estimates

$$||h^{-}||_{2} \lesssim ||h^{+}||_{2} \approx ||f_{0}||_{2} \approx ||g_{0}||_{2} \approx ||f||_{\mathcal{X}} \approx ||g||_{\mathcal{X}}$$

Note that these limits for \mathcal{X} -solutions are stronger than L_2 convergence of Cesaro means $t^{-1} \int_t^{2t} g_s ds$, and that we do get limits for g and f, *i.e.* for the full gradient and conormal gradient. That $\|g_0\|_2 \approx \|f_0\|_2$ and $\|g\|_{\mathcal{X}} \approx \|f\|_{\mathcal{X}}$ are easy consequences of Proposition 4.1.

Proof. (i) Assume that g is an \mathcal{X} -solution, and consider f. To show that $f_t = e^{-t\Lambda}h^+ + S_A f_t$, we choose η_{ϵ}^{\pm} for η^{\pm} in Proposition 4.4 and subtract the equations to obtain

$$(35) \quad -\int_0^t (\partial_s \eta_\epsilon^+)(t,s) e^{-(t-s)\Lambda} E_0^+ f_s ds + \int_t^\infty (\partial_s \eta_\epsilon^-)(t,s) e^{-(s-t)\Lambda} E_0^- f_s ds$$
$$= \int_0^t \eta_\epsilon^+(t,s) \Lambda e^{-(t-s)\Lambda} \widehat{E}_0^+ \mathcal{E}_s f_s ds + \int_t^\infty \eta_\epsilon^-(t,s) \Lambda e^{-(s-t)\Lambda} \widehat{E}_0^- \mathcal{E}_s f_s ds.$$

Note that $DB_0 = \pm |DB_0| = \pm \Lambda$ on $E_0^{\pm} \mathcal{H}$. We fix $0 < a < b < \infty$ and consider the equation in $L_2(a, b; \mathcal{H})$. By Proposition 7.1, the right hand side converges to $S_A f$ in $L_2(a, b; \mathcal{H})$. When $t \in (a, b)$ and ϵ is small, the left hand side equals

(36)
$$\epsilon^{-1} \int_{\epsilon}^{2\epsilon} e^{-s\Lambda} (E_0^+ f_{t-s} + E_0^- f_{t+s}) ds$$

 $-\epsilon^{-1} \int_{\epsilon}^{2\epsilon} e^{-(t-s)\Lambda} E_0^+ f_s ds - 2\epsilon \int_{(2\epsilon)^{-1}}^{\epsilon^{-1}} e^{-(s-t)\Lambda} E_0^- f_s ds.$

To prove that the first term converges to f in $L_2(a, b; \mathcal{H})$, adding and subtracting the term $\epsilon^{-1} \int_{\epsilon}^{2\epsilon} e^{-s\Lambda} f_t ds = e^{-\epsilon\Lambda} (\epsilon\Lambda)^{-1} (I - e^{-\epsilon\Lambda}) f_t$ shows that the square of the $L_2(a, b; \mathcal{H})$ norm of the difference is bounded by

$$\int_{a}^{b} \left\| \left(I - e^{-\epsilon\Lambda} \frac{I - e^{-\epsilon\Lambda}}{\epsilon\Lambda} \right) f_t \right\|_{2}^{2} dt + \int_{a}^{b} \epsilon^{-1} \int_{\epsilon}^{2\epsilon} \|f_t - E_0^+ f_{t-s} - E_0^- f_{t+s}\|_{2}^{2} ds dt \to 0$$

as $\epsilon \to 0$, using Proposition 6.4 for the functional calculus, dominated convergence, and the identity $f_t = E_0^+ f_t + E_0^- f_t$.

Next consider the last term in (36). For any $\phi \in L_2(a, b; \mathcal{H})$, we have

$$\int_{a}^{b} \left(\epsilon \int_{(2\epsilon)^{-1}}^{\epsilon^{-1}} e^{-(s-t)\Lambda} E_{0}^{-} f_{s} ds, \phi_{t} \right) dt$$
$$= \epsilon \int_{(2\epsilon)^{-1}}^{\epsilon^{-1}} \left(f_{s}, \int_{a}^{b} (e^{-(s-t)\Lambda^{*}} - e^{-s\Lambda^{*}}) (E_{0}^{-})^{*} \phi_{t} dt + e^{-s\Lambda^{*}} (E_{0}^{-})^{*} \int_{a}^{b} \phi_{t} dt \right) ds.$$

From the sup $-L_2$ estimate in Lemma 5.3 for f, the estimate $||e^{-(s-t)\Lambda^*} - e^{-s\Lambda^*}|| \leq t/s$ and the strong limit $\lim_{s\to\infty} e^{-s\Lambda^*} (E_0^-)^* = 0$, it follows that the last term in (36) converges weakly to 0. Hence the middle term must converge weakly in $L_2(a, b; L_2)$ as well, and we may replace $e^{-(t-s)\Lambda}$ by $e^{-t\Lambda}$ since $||e^{-(t-s)\Lambda} - e^{-t\Lambda}|| \leq s/t$. We get that

$$\int_{a}^{b} \left(e^{-t\Lambda} (\epsilon^{-1} \int_{\epsilon}^{2\epsilon} E_{0}^{+} f_{s} ds), \phi_{t} \right) dt = \left(\epsilon^{-1} \int_{\epsilon}^{2\epsilon} E_{0}^{+} f_{s} ds, \int_{a}^{b} e^{-t\Lambda^{*}} \phi_{t} dt \right)$$

converges for all $\phi \in L_2(a, b; L_2)$. Since $\epsilon^{-1} \int_{\epsilon}^{2\epsilon} E_0^+ f_s ds$ are uniformly bounded in \mathcal{H} by Lemma 5.3, and since functions $\int_a^b e^{-t\Lambda^*} \phi_t dt$ are dense in $B_0\mathcal{H} \approx \mathcal{H}^*$ (for example $\int_{\epsilon}^{2\epsilon} e^{-t\Lambda^*} \epsilon^{-1} \phi dt \rightarrow P_{B_0\mathcal{H}} \phi$), it follows that $\epsilon^{-1} \int_{\epsilon}^{2\epsilon} E_0^+ f_s ds$ converges weakly to a function $h^+ \in E_0^+\mathcal{H}$, and that the weak limit of the middle term in (36) is $e^{-t\Lambda}h^+$. In total, this proves that $f_t - e^{-t\Lambda}h^+ = S_A f_t$.

(ii) Conversely, assume that $f \in \mathcal{X}$ solves $f_t = e^{-t\Lambda}h^+ + S_A f_t$. To verify that f satisfies the differential equation, note that $(\partial_t + DB_0)e^{-t\Lambda}h^+ = 0$. It suffices to show that for $\phi \in C_0^{\infty}(\mathbf{R}^{1+n}_+; \mathbf{C}^{(1+n)m})$ there is convergence

$$\int (-\partial_t \phi_t + B_0^* D\phi_t, f_t^{\epsilon}) dt \to \int (D\phi_s, \mathcal{E}_s f_s) ds, \qquad \epsilon \to 0.$$

where $f_t^{\epsilon} := S_A^{\epsilon} f_t$. For the term $S_{\epsilon}^+ \widehat{E}_0^+ \mathcal{E} f$, Fubini's theorem and integration by parts give

$$\begin{split} \int_0^\infty \int_0^t \eta_\epsilon^+(t,s)((-\partial_t + \Lambda^*)\phi_t, \Lambda e^{-(t-s)\Lambda}\widehat{E}_0^+\mathcal{E}_s f_s)dsdt \\ &= -\int_0^\infty \left(\int_s^\infty \eta_\epsilon^+(t,s)\partial_t (e^{-(t-s)\Lambda^*}\Lambda^*\phi_t)dt, \widehat{E}_0^+\mathcal{E}_s f_s\right)ds \\ &= \int_0^\infty \left(\int_s^\infty (\partial_t \eta_\epsilon^+)(t,s)e^{-(t-s)\Lambda^*}\Lambda^*\phi_tdt, \widehat{E}_0^+\mathcal{E}_s f_s\right)ds \\ &\to \int_0^\infty (\Lambda^*\phi_s, \widehat{E}_0^+\mathcal{E}_s f_s)ds = \int_0^\infty (D\phi_s, \widetilde{E}_0^+\mathcal{E}_s f_s)ds. \end{split}$$

Adding the corresponding limit for the term $S_{\epsilon}^{-} \widehat{E}_{0}^{-} \mathcal{E}f$ gives the stated result. Note that $\widetilde{E}_{0}^{+} + \widetilde{E}_{0}^{-} = P_{B_{0}\mathcal{H}}$ and $DP_{B_{0}\mathcal{H}} = D$.

(iii) To show the limits, note that $E_0^+ f - e^{-t\Lambda} h^+ = S^+ \widehat{E}_0^+ \mathcal{E} f \in \mathcal{Y}^*$, and by inspection of the proof of Theorem 6.8 we see that $E_0^- f - e^{-t\Lambda} \int_0^\infty \Lambda e^{-s\Lambda} \widehat{E}_0^- \mathcal{E}_s f_s ds \in \mathcal{Y}^*$. From this, the limits for f follow. To see the limit for g at t = 0, write $B_t f_t - B_0 f_0 = B_0 (f_t - f_0) + \mathcal{E}_t f_t$. Since $\mathcal{E}_t f_t \in \mathcal{Y}^*$, we have $\lim_{t\to 0} t^{-1} \int_t^{2t} \|\mathcal{E}_s f_s\|_2^2 ds = 0$. The limit of $B_0 (f_t - f_0)$ at t = 0, as well as that of g at $t = \infty$, is immediate from the limits of f. (iv) It remains to prove the estimates. Note that (20) and Lemma 5.3 show that

$$\max(\|h^+\|_2^2, \|h^-\|_2^2) \approx \|f_0\|_2^2 = \lim_{t \to 0} t^{-1} \int_t^{2t} \|f_s\|_2^2 ds \lesssim \|f\|_{\mathcal{X}}^2.$$

Proposition 7.1 shows that $||S_A||_{\mathcal{X}\to\mathcal{X}} \leq 1/2$ if $||\mathcal{E}||_*$ is sufficiently small. In this case $I - S_A$ is an isomorphism on \mathcal{X} with $||(I - S_A)^{-1}||_{\mathcal{X}\to\mathcal{X}} \leq 2$. Using this together with Theorem 5.2, we get estimates $||f||_{\mathcal{X}} = ||(I - S_A)^{-1}e^{-t\Lambda}h^+||_{\mathcal{X}} \approx ||h^+||_2$. This proves the stated estimates and completes the proof.

Remark 8.3. We may ask whether \mathcal{Y}^* could be used as a solution space for gradients of solutions. The answer is no because we have seen that functions in \mathcal{Y}^* vanish in some sense at the boundary. A second question is then how far gradients of solutions are from being in \mathcal{Y}^* . Inspection of the proofs of Theorem 6.8 and Theorem 8.2 reveals that $S_A f_t = e^{-t\Lambda}h^- + \hat{f}_t$ with $\hat{f} \in \mathcal{Y}^*$ and $h^- \in E_0^- \mathcal{H}$. Hence, if g is an \mathcal{X} -solution, the corresponding conormal gradient f can be represented (assuming $\|\mathcal{E}\|_* < \infty$) as $f_t = e^{-t\Lambda}f_0 + \hat{f}_t$, since $f_0 = h^+ + h^-$. Note in particular that $f - e^{-t\Lambda}f_0 \in \mathcal{Y}^* \subsetneq \mathcal{X}$, *i.e.* the free evolution $e^{-t\Lambda}f_0$ is the term responsible for f(hence g) to belong to \mathcal{X} and not to \mathcal{Y}^* .

We note the following immediate corollary to Theorem 8.2.

Corollary 8.4. Assume that coefficients $A = A_0$ are t-independent. Then g is an \mathcal{X} -solution to the divergence form equation if and only if the associated conormal gradient f is a generalized Cauchy extension $C_0^+h^+$ of some $h^+ \in E_0^+\mathcal{H}$, i.e.

$$f_t = e^{-t\Lambda} h^+ \qquad for \ all \ t > 0,$$

and $h^+ = \lim_{t\to 0} f_t$ in L_2 sense. In particular, the class of \mathcal{X} -solutions in Definition 8.1 coincides with the class of solutions in [9, Definition 2.1(i-ii)] for t-independent coefficients.

That the solutions in [9] are of this form was shown in the proof of [9, Theorem 2.3]. Note that the operator $T_A|_{\mathcal{H}}$ used in [9] is similar to our operator $DB_0|_{\mathcal{H}}$, as in [9, Definition 3.1].

We are now ready study BVPs. We recall that for the class of solutions used in [9], with *t*-independent coefficients A_0 , well-posedness of the Neumann and regularity problems was shown to be equivalent to the maps

$$E_0^+ \mathcal{H} \to L_2(\mathbf{R}^n; \mathbf{C}^m) : h^+ \mapsto (h^+)_{\perp},$$

$$E_0^+ \mathcal{H} \to \{ f \in L_2(\mathbf{R}^n; \mathbf{C}^{nm}) \; ; \; \operatorname{curl}_x f = 0 \} : h^+ \mapsto (h^+)_{\parallel},$$

being isomorphisms respectively. From Corollary 8.4, it is equivalent to well-posedness in the class of \mathcal{X} -solutions.

We now turn to *t*-dependent perturbations of the coefficients.

Corollary 8.5. Assume that the Neumann problem for A_0 is well-posed. Then there exists $\epsilon > 0$ such that for any t-dependent coefficient matrix A with $\|\mathcal{E}\|_* < \epsilon$, the Neumann problem is well-posed for A in the following sense.

Given any function $\varphi \in L_2(\mathbb{R}^n; \mathbb{C}^m)$, there is a unique \mathcal{X} -solution g to the divergence form equation with coefficients A, whose trace g_0 satisfies $(A_0g_0)_{\perp} = \varphi$. Moreover, this solution has estimates

$$\|N_*(g)\|_2 \approx \|g_0\|_2 \approx \|\varphi\|_2.$$

The same holds true when the Neumann problem is replaced by the regularity problem and the boundary condition $(A_0g_0)_{\perp} = \varphi$ is replaced by $(g_0)_{\parallel} = \varphi \in L_2(\mathbf{R}^n; \mathbf{C}^{nm})$ such that $\operatorname{curl}_x \varphi = 0$.

Proof. Throughout the proof, we assume that $\|\mathcal{E}\|_*$ is small enough, so that $I - S_A$ is invertible on \mathcal{X} by Proposition 7.1. To solve the Neumann problem, we compute f by making the ansatz

(37)
$$f = (I - S_A)^{-1} C_0^+ h^+$$

for some $h^+ \in E_0^+ \mathcal{H}$ to be determined, and calculate its full trace

$$f_0 = h^+ + \int_0^\infty \Lambda e^{-s\Lambda} \widehat{E}_0^- \mathcal{E}_s f_s ds$$

using Theorem 8.2. We see that f satisfies the Neumann boundary condition $(f_0)_{\perp} = \varphi$ if and only if h^+ solves the equation $\Gamma_A h^+ = \varphi$, where $\Gamma_A : E_0^+ \mathcal{H} \to L_2(\mathbf{R}^n; \mathbf{C}^m)$ is the operator

$$\Gamma_A: h^+ \mapsto \left(h^+ + \int_0^\infty \Lambda e^{-s\Lambda} \widehat{E}_0^- \mathcal{E}_s f_s ds\right)_\perp$$

with f given by (37). Since $\Gamma_{A_0}h^+ = (h^+)_{\perp}$, a computation using Lemma 6.9 and the boundedness of $(I - S_A)^{-1}$ on \mathcal{X} implies $\|\Gamma_A - \Gamma_{A_0}\|_{L_2 \to L_2} \leq \|A - A_0\|_*$. By assumption Γ_{A_0} is an invertible operator, and thus Γ_A remains an isomorphism whenever $\|A - A_0\|_*$ is sufficiently small. Thus, in this case we can, given φ , calculate $h^+ = \Gamma_A^{-1}\varphi$ with $\|h^+\|_2 \approx \|\varphi\|_2$ and find a unique solution g to the Neumann problem, with estimates $\|g\|_{\mathcal{X}} \approx \|g_0\|_2 \approx \|h^+\|_2 \approx \|\varphi\|_2$.

For the regularity problem, we proceed as for the Neumann problem, but instead solve for h^+ in the equation $\left(h^+ + \int_0^\infty \Lambda e^{-s\Lambda} \widehat{E}_0^- \mathcal{E}_s f_s ds\right)_{\parallel} = \varphi$.

9. The Dirichlet problem

Throughout this section, A denotes t-dependent coefficients satisfying (2) and (3), and $A_0 \in L_{\infty}(\mathbb{R}^n; \mathcal{L}(\mathbb{C}^{(1+n)m}))$ denotes t-independent coefficients which are strictly accretive on \mathcal{H} . We let $B := \hat{A}$ and $B_0 := \hat{A}_0$ be the transformed strictly accretive coefficients from Proposition 4.1, and define $\mathcal{E} := B_0 - B$.

Definition 9.1. By a \mathcal{Y} -solution to the divergence form equation, with coefficients A, we mean a function $u \in W_2^{1,\text{loc}}(\mathbf{R}^{1+n}_+, \mathbf{C}^m)$, with estimate $\int_0^\infty ||g_t||_2^2 t dt < \infty$ of its gradient $g := \nabla_{t,x} u$ which satisfies $\operatorname{div}_{t,x} Ag = 0$ in \mathbf{R}^{1+n}_+ , in distributional sense.

We will prove in Theorem 9.3, for certain coefficients A, that any \mathcal{Y} -solution belongs to $C(\mathbf{R}_+; L_2)$, modulo constants. Note also that we do not assume any limits of u at t = 0 or $t = \infty$, but will prove such below. This will allow us to formulate in what sense the boundary values are prescribed. When discussing $C(\mathbf{R}_+; L_2)$ limits and bounds of \mathcal{Y} -solutions u, we shall identify the function u with a $C(\mathbf{R}_+; L_2)$ function modulo constants.

Our representation result for gradients of \mathcal{Y} -solutions is the following. Recall that $\tilde{\Lambda} = |B_0 D|$.

Theorem 9.2. Assume that $\|\mathcal{E}\|_* < \infty$. Consider $u \in W_2^{1,\text{loc}}(\mathbf{R}^{1+n}_+, \mathbf{C}^m)$ with conormal gradient $f = [(A\nabla_{t,x}u)_{\perp}, \nabla_x u]^t \in \mathcal{Y}$. Then f solves $\partial_t f + DBf = 0$ in \mathbf{R}^{1+n}_+ distributional sense if and only if it satisfies the equation

$$f_t = De^{-t\tilde{\Lambda}}\tilde{h}^+ + S_A f_t, \quad \text{for some } \tilde{h}^+ \in \widetilde{E}_0^+ L_2.$$

In this case, let $v_t := e^{-t\tilde{\Lambda}}\tilde{h}^+ + \tilde{S}_A f_t$. Then f = Dv and v_t has L_2 limits $\lim_{t \to 0} \|v_t - v_0\|_2 = 0 = \lim_{t \to \infty} \|v_t\|_2,$ (38)

where $v_0 := \tilde{h}^+ + \tilde{h}^-$ and $\tilde{h}^- := -\int_0^\infty e^{-s\tilde{\Lambda}} \tilde{E}_0^- \mathcal{E}_s f_s ds \in \tilde{E}_0^- L_2$, with estimates $\max(\|\tilde{h}^+\|_2, \|\tilde{h}^-\|_2) \approx \|v_0\|_2 \lesssim \sup_{t>0} \|v_t\|_2 \lesssim \|f\|_{\mathcal{Y}}.$

If furthermore $\|\mathcal{E}\|_*$ is sufficiently small, then there are estimates

$$\|\tilde{h}^-\|_2 \lesssim \|\tilde{h}^+\|_2 \approx \sup_{t>0} \|v_t\|_2 \approx \|f\|_{\mathcal{Y}} \approx \|\nabla_{t,x}u\|_{\mathcal{Y}}.$$

Proof. (i) Assume that $f \in \mathcal{Y}$ satisfies the differential equation. As in the proof of Theorem 8.2, we aim to take limits $\epsilon \to 0$ in equation (35). By Proposition 7.1, the right hand side converges in \mathcal{Y} to $S_A f$. Fix $0 < a < b < \infty$. For $t \in (a, b)$ and small ϵ , the left hand side equals

(39)
$$\epsilon^{-1} \int_{\epsilon}^{2\epsilon} e^{-s\Lambda} (E_0^+ f_{t-s} + E_0^- f_{t+s}) ds$$

 $-\epsilon^{-1} \int_{\epsilon}^{2\epsilon} e^{-(t-s)\Lambda} E_0^+ f_s ds - 2\epsilon \int_{(2\epsilon)^{-1}}^{\epsilon^{-1}} e^{-(s-t)\Lambda} E_0^- f_s ds.$

As in the proof of Theorem 8.2, the first term converges to f in $L_2(a, b; L_2)$. The L_2 -norm of the last term is bounded by $\epsilon \int_{(2\epsilon)^{-1}}^{\epsilon^{-1}} \|f_s\|_2 ds \lesssim \epsilon (\int_{(2\epsilon)^{-1}}^{\epsilon^{-1}} \|f_s\|^2 s ds)^{1/2}$, and hence converges to 0, uniformly for $t \in (a, \overline{b})$. We conclude that $\tilde{f}_t^{\epsilon} := \epsilon^{-1} \int_{\epsilon}^{2\epsilon} e^{-(t-s)\Lambda} E_0^+ f_s ds$ converges in $L_2(a, b; L_2)$ as $\epsilon \to 0$.

In fact, since $\sup_{t>0} \|e^{-t\Lambda}\|_{L_2\to L_2} < \infty$ we have

$$\|\tilde{f}_{t_0}^{\epsilon} - \tilde{f}_{t_0}^{\epsilon'}\|_2 \le \frac{1}{b-a} \int_a^b \|e^{-(t_0-t)\Lambda} (\tilde{f}_t^{\epsilon} - \tilde{f}_t^{\epsilon'})\|_2 dt \lesssim \left(\int_a^b \|\tilde{f}_t^{\epsilon} - \tilde{f}_t^{\epsilon'}\|_2^2 dt\right)^{1/2},$$

when $t_0 > b$. Hence, since (a, b) is arbitrary, \tilde{f}_t^{ϵ} converges in L_2 , locally uniformly in t. Call the limit \tilde{f} and note that it coincides with $f - S_A f \in \mathcal{Y}$ for a.e. t > 0. Fix $t_0 > 0$ and note that $\tilde{f}_{t+t_0} = \lim_{\epsilon \to 0} e^{-t\Lambda} \tilde{f}_{t_0}^{\epsilon} = e^{-t\Lambda} \tilde{f}_{t_0}^{\epsilon}$ and that in fact $\tilde{f}_{t_0} \in E_0^+ \mathcal{H}$ by the definition of \tilde{f}_{t_0} . The estimate

$$\sup_{t_0>0} \int_0^\infty \|e^{-t\Lambda} \tilde{f}_{t_0}\|_2^2 t dt \le \|\tilde{f}\|_{\mathcal{Y}}^2 \lesssim \|f\|_{\mathcal{Y}}^2$$

follows. Consider the restriction Λ_+ of Λ to $E_0^+\mathcal{H}$, which is a closed and injective operator with dense domain and range. We claim that $\tilde{f}_{t_0} \in \mathsf{D}(\Lambda_+^{-1})$. To see this, by duality it suffices to show that

$$|((\Lambda_{+}^{-1})^*\phi, \tilde{f}_{t_0})| \lesssim ||\phi||_2, \quad \text{for all } \phi \in \mathsf{D}((\Lambda_{+}^{-1})^*).$$

As in the proof of Proposition 6.3, we use an identity $\int_0^\infty (t\Lambda_+ e^{-t\Lambda_+})^2 \tilde{f}_{t_0} dt/t = 4^{-1} \tilde{f}_{t_0}$ to estimate

$$|((\Lambda_{+}^{-1})^{*}\phi,\tilde{f}_{t_{0}})| \approx \left|\int_{0}^{\infty} (t\Lambda_{+}^{*}e^{-t\Lambda_{+}^{*}}\phi,te^{-t\Lambda_{+}}\tilde{f}_{t_{0}})\frac{dt}{t}\right| \lesssim \|\phi\|_{2}\|f\|_{\mathcal{Y}}$$

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Hence the claim. As $\mathsf{D}(\Lambda_{+}^{-1}) = \mathsf{R}(\Lambda_{+}) \subset \mathsf{R}(D)$, this shows that $\tilde{f}_{t_0} = D\tilde{h}_{t_0}^+$, where $\tilde{h}_{t_0}^+ \in B_0 E_0^+ \mathcal{H} = \tilde{E}_0^+ L_2$ has bounds $\|\tilde{h}_{t_0}^+\|_2 \lesssim \|f\|_{\mathcal{Y}}$, uniformly in t_0 . From the identity $\tilde{f}_{t+t_0} = e^{-t\Lambda} \tilde{f}_{t_0} = e^{-t\Lambda} D\tilde{h}_{t_0}^+$, we get

$$\int_{a}^{b} (\phi_t, \tilde{f}_{t+t_0}) dt = \left(\int_{a}^{b} De^{-t\Lambda^*} \phi_t dt, \tilde{h}_{t_0}^+ \right)$$

for any $\phi \in L_2(a, b; L_2)$. Here the left hand side converges as $t_0 \to 0$, and the functions $\int_a^b De^{-t\Lambda^*}\phi_t dt$ are dense in \mathcal{H} . (For example $\int_{\epsilon}^{2\epsilon} De^{-t\Lambda^*}\epsilon^{-1}\phi dt \to D\phi$.) Since $\|\tilde{h}_{t_0}^+\|_2$ is uniformly bounded, it follows that $\tilde{h}_{t_0}^+ \to \tilde{h}^+$ weakly in $\tilde{E}_0^+L_2$ as $t_0 \to 0$. Letting $t_0 \to 0$ in $\tilde{f}_{t+t_0} = e^{-t\Lambda}D\tilde{h}_{t_0}^+ = De^{-t\tilde{\Lambda}}\tilde{h}_{t_0}^+$, we obtain $f_t - S_A f_t = \tilde{f}_t = De^{-t\tilde{\Lambda}}\tilde{h}^+$ for a.e. t > 0.

(ii) Conversely, assume that $f \in \mathcal{Y}$ and $f_t = De^{-t\tilde{\Lambda}}\tilde{h}^+ + S_A f_t$. As in Theorem 8.2, we verify that f satisfies the differential equation, and we omit the details.

(iii) Introduce $v_t = e^{-t\tilde{\Lambda}}\tilde{h}^+ + \tilde{S}_A f_t$, so that f = Dv and the stated limits on v follow from Propositions 6.4 and 7.2.

To prove the estimates, note that the square function estimates for B_0D and the accretivity of B_0 on \mathcal{H} show that

$$\|\tilde{h}^+\|_2 \approx \|B_0 D e^{-t\tilde{\Lambda}} \tilde{h}^+\|_{\mathcal{Y}} \approx \|D e^{-t\tilde{\Lambda}} \tilde{h}^+\|_{\mathcal{Y}} \lesssim \|f\|_{\mathcal{Y}} + \|S_A f\|_{\mathcal{Y}} \lesssim \|f\|_{\mathcal{Y}}.$$

From Proposition 7.2, we also obtain the estimates $\max(\|\tilde{h}^+\|_2, \|\tilde{h}^-\|_2) \approx \|v_0\|_2 \leq \sup_{t>0} \|v_t\|_2 \lesssim \|\tilde{h}^+\|_2 + \|f\|_{\mathcal{Y}} \lesssim \|f\|_{\mathcal{Y}}$, where we have used the topological splitting $B_0\mathcal{H} = \widetilde{E}_0^+L_2 \oplus \widetilde{E}_0^-L_2$ in the first equivalence.

(iv) Finally, Proposition 7.1 shows that $||S_A||_{\mathcal{Y}\to\mathcal{Y}} \leq 1/2$ if $||\mathcal{E}||_*$ is sufficiently small. In this case $I - S_A$ is an isomorphism on \mathcal{Y} , giving the estimate

$$\|f\|_{\mathcal{Y}} \lesssim \|De^{-t\tilde{\Lambda}}\tilde{h}^+\|_{\mathcal{Y}}.$$

As $\|De^{-t\tilde{\Lambda}}\tilde{h}^+\|_{\mathcal{Y}} \approx \|\tilde{h}^+\|_2$, this proves the stated estimates and completes the proof.

We can now prove a rigidity theorem for \mathcal{Y} -solutions.

Theorem 9.3. Let u be a \mathcal{Y} -solution to the divergence form equation, with coefficients A. Assume that $\|\mathcal{E}\|_* < \infty$. Then there is a constant $c \in \mathbf{C}^m$ such that

$$u = c - v_{\perp}$$

almost everywhere, where $v \in C(\mathbf{R}_+; L_2)$ is the vector-valued potential from Theorem 9.2. Identifying the functions u and $c - v_{\perp}$, there are L_2 -limits $\lim_{t\to 0} ||u_t - u_0||_2 = 0$, $u_0 := c - (v_0)_{\perp}$, and $\lim_{t\to\infty} ||u_t - c||_2 = 0$, and there are bounds

$$||u_0 - c||_2 \le \sup_{t>0} ||u_t - c||_2 \le ||\nabla_{t,x}u||_{\mathcal{Y}}.$$

Proof. Let f, \tilde{h}^+ and v be as in Theorem 9.2. We claim that v satisfies the equation

$$(\partial_t + BD)v_t = -P\mathcal{E}_t f_t$$

in \mathbf{R}^{1+n}_+ distributional sense, where $P = I - \widetilde{E}^+_0 - \widetilde{E}^-_0$ is the projection onto \mathcal{H}^{\perp} along $B_0\mathcal{H}$. To prove this, compute

$$\partial_t (e^{-t\tilde{\Lambda}}\tilde{h}^+ + \widetilde{S}^{\epsilon}_A f_t) = -B_0 D e^{-t\tilde{\Lambda}}\tilde{h}^+ - B_0 D \widetilde{S}^{\epsilon}_A f_t + \epsilon^{-1} \int_{\epsilon}^{\infty} e^{-s\tilde{\Lambda}} (\widetilde{E}^+_0 \mathcal{E}_{t-s} f_{t-s} + \widetilde{E}^-_0 \mathcal{E}_{t+s} f_{t+s}) ds,$$

for $\epsilon \ll t$. This uses the result for the operator B_0D analogous to Proposition 4.3. The claim follows by letting $\epsilon \to 0$, using Proposition 7.1 for convergence of $D\widetilde{S}_A^{\epsilon} = S_A^{\epsilon}$, and a calculation as in part (i) of the proof of Theorem 8.2 for the last term.

It is clear that $\nabla_x u = f_{\parallel} = -\nabla_x v_{\perp}$. Using the claim above, we also obtain

$$\partial_t (v_t)_\perp + (BDv_t)_\perp = 0,$$

because normal parts of functions in \mathcal{H}^{\perp} are zero. Since $(BDv)_{\perp} = (Bf)_{\perp} = g_{\perp} = \partial_t u$, we conclude that $\nabla_{t,x} u = -\nabla_{t,x} v_{\perp}$ in \mathbf{R}^{1+n}_+ distributional sense. Hence $u = c - v_{\perp}$. The stated limits and bounds now follow from Theorem 9.2.

The constant c in Theorem 9.3 can also be calculated as the limit

$$c = \lim_{d \to \infty} (u, \tau_d \phi), \qquad \tau_d \phi(t, x) := \phi(t - d, x),$$

for any $\phi \in C_0^{\infty}(\mathbf{R}^{1+n}; \mathbf{C}^m)$ with $\int \phi dt dx = 1$. In particular this limit does not depend on ϕ . So if this limit is zero, we obtain a solution that vanishes at infinity in L_2 sense. This is akin to the classical pointwise limit at infinity required to eliminate constants for representing harmonic functions in the upper half-space. Here the averages replace the pointwise values.

We also note the following corollary to Theorem 9.3.

Corollary 9.4. Assume that coefficients $A = A_0$ are t-independent. Then u is a \mathcal{Y} -solution to the divergence form equation if and only if it is the normal part of a generalized Cauchy extension $\widetilde{C}_0^+ \widetilde{h}^+$ of some $\widetilde{h}^+ \in \widetilde{E}_0^+ L_2$, modulo constants, i.e.

$$u_t = (e^{-t\Lambda}\tilde{h}^+)_{\perp} + c$$
 for all $t > 0$ and some $c \in \mathbf{C}^m$.

In particular, the class of \mathcal{Y} -solutions in Definition 9.1 that vanish at infinity coincides with the class of solutions in [9, Definition 2.1(iii)] for t-independent coefficients.

That the solutions considered in [9] are of this form follows from [9, Lemma 4.2] and the proof of [9, Theorem 2.3]. Note that the operator $T_A|_{\mathcal{H}}$ used in [9] is similar to our operator $B_0D|_{B_0\mathcal{H}}$, as in [9, Definition 3.1].

We are now ready to study BVPs. We recall that for the class of solutions used in [9], with *t*-independent coefficients A_0 , well-posedness of the Dirichlet problem was shown to be equivalent to the map

$$\widetilde{E}_0^+ L_2 \to L_2(\mathbf{R}^n; \mathbf{C}^m) : \widetilde{h}^+ \mapsto (\widetilde{h}^+)_\perp$$

being an isomorphism. From Corollary 9.4, it is equivalent to well-posedness in the class of \mathcal{Y} -solutions. Remark that well-posedness implies that the map $u_0 \mapsto u_t$ is a C_0 -semigroup on $L_2(\mathbf{R}^n; \mathbf{C}^m)$ and this corollary also shows that the results in [5] concerning the domain of this semi-group obtained for solutions in the sense of [9] apply to \mathcal{Y} -solutions.

We now turn to t-dependent perturbations of the coefficients.

Corollary 9.5. Assume that the Dirichlet problem for A_0 is well-posed.

Then there exists $\epsilon > 0$ such that for any t-dependent coefficient matrix A with $\|\mathcal{E}\|_* < \epsilon$, the Dirichlet problem is well-posed for A in the following sense.

Given any function $\varphi \in L_2(\mathbf{R}^n; \mathbf{C}^m)$, there is a unique \mathcal{Y} -solution u to the divergence form equation with coefficients A, with boundary trace $u_0 = \varphi$. Moreover, this solution has estimates

$$\|\nabla_{t,x}u\|_{\mathcal{Y}} \approx \sup_{t>0} \|u_t\|_2 \approx \|\varphi\|_2.$$

Proof. Throughout the proof, we assume that $\|\mathcal{E}\|_*$ is small enough, so that $I - S_A$ is invertible on \mathcal{Y} by Proposition 7.1. To solve the Dirichlet problem, we make the ansatz

(40)
$$u = \left((I + \widetilde{S}_A (I - S_A)^{-1} D) \widetilde{C}_0^+ \widetilde{h}^+ \right)_\perp$$

for some $\tilde{h}^+ \in \tilde{E}_0^+ L_2$. Theorems 9.2 and 9.3 show that u is a \mathcal{Y} -solution to the divergence form equation with coefficients A and that all \mathcal{Y} -solutions with L_2 trace are of this form. Moreover, the Dirichlet boundary condition $u_0 = \varphi$ is satisfied if and only if \tilde{h}^+ solves the equation $\tilde{\Gamma}_A \tilde{h}^+ = \varphi$, where $\tilde{\Gamma}_A : \tilde{E}_0^+ L_2 \to L_2(\mathbf{R}^n; \mathbf{C}^m)$ is the operator

$$\widetilde{\Gamma}_A: \widetilde{h}^+ \mapsto \left(\widetilde{h}^+ - \int_0^\infty e^{-s\widetilde{\Lambda}} \widetilde{E}_0^- \mathcal{E}_s f_s ds\right)_\perp,$$

where $f := (I - S_A)^{-1} D \widetilde{C}_0^+ \widetilde{h}^+$. Since $\widetilde{\Gamma}_{A_0} \widetilde{h}^+ = (\widetilde{h}^+)_{\perp}$, Lemma 6.9 and the boundedness of $(I - S_A)^{-1}$ on \mathcal{Y} imply that $\|\widetilde{\Gamma}_A - \widetilde{\Gamma}_{A_0}\|_{L_2 \to L_2} \lesssim \|A - A_0\|_*$. By assumption $\widetilde{\Gamma}_{A_0}$ is an invertible operator, and thus $\widetilde{\Gamma}_A$ remains an isomorphism whenever $\|\mathcal{E}\|_*$ is sufficiently small. Thus, in this case we can, given φ , calculate $\widetilde{h}^+ = \widetilde{\Gamma}_A^{-1} \varphi$ with $\|\widetilde{h}^+\|_2 \approx \|\varphi\|_2$ and find a unique solution u to the Dirichlet problem. From Theorem 9.2, we get estimates

$$\|\varphi\|_2 \le \sup_{t>0} \|u_t\|_2 \lesssim \|\nabla_{t,x}u\|_{\mathcal{Y}} \approx \|\tilde{h}^+\|_2 \approx \|\varphi\|_2.$$

This proves the theorem.

Remark 9.6. The tangential part v_{\parallel} of the vector-valued potential $v = (I + \tilde{S}_A(I - S_A)^{-1}D)\tilde{C}_0^+\tilde{h}^+$ can be viewed as a set of generalized conjugate functions to the Dirichlet solution u. Our proof of Theorem 9.2 above eliminates the need of the technical condition at ∞ on these conjugate functions which was required in [9, Definition 3.1].

10. Further estimates

In Section 8, we constructed solutions, with estimates on the modified nontangential maximal function, to the Neumann and regularity problems with L_2 boundary data, and in Section 9 we constructed solutions, with estimates on the square function, to the Dirichlet problem with L_2 boundary data. In this section, we prove two theorems which give modified non-tangential maximal function estimates for the Dirichlet problem, and square function estimates for the Neumann/regularity problems.

10.1. Maximal function estimates for \mathcal{Y} -solutions.

Theorem 10.1. Let A_0 be t-independent coefficients which are strictly accretive on \mathcal{H} , and assume that $||A - A_0||_C < \infty$. Then any \mathcal{Y} -solution u to the divergence form equation with coefficients A, with boundary trace $u_0 \in L_2(\mathbb{R}^n, \mathbb{C}^m)$, has modified non-tangential maximal estimates

$$||u_0||_2 \lesssim ||N_*(u)||_2 \lesssim ||\nabla_{t,x}u||_{\mathcal{Y}}.$$

The core of the proof reduces to the following estimate of the operator S_A .

Lemma 10.2. For any fixed $p \in [1, 2)$, the operator \widetilde{S}_A has estimates

$$\|\widetilde{N}^p_*((\widetilde{S}_A f)_{\perp})\|_2 \lesssim \|\mathcal{E}\|_C \|f\|_{\mathcal{Y}}.$$

Here $\widetilde{N}^p_*(f)(x) := \sup_{t>0} t^{-(1+n)/p} ||f||_{L_p(W(t,x))}$ is an L_p modified non-tangential maximal function.

Proof of Theorem 10.1 modulo Lemma 10.2. As in Theorems 9.2 and 9.3, any \mathcal{Y} solution u that vanishes at infinity can be written

$$u_t = (e^{-t\tilde{\Lambda}}\tilde{h}^+ + \tilde{S}_A f_t)_{\perp}, \qquad \tilde{h}^+ \in \tilde{E}_0^+ L_2, f \in \mathcal{Y}$$

From Poincaré's inequality $||u - u_{W(t,x)}||_{L_2(W(t,x))} \lesssim t ||\nabla_{s,y}u||_{L_2(W(t,x))}$, where $u_{W(t,x)}$ denotes the average, we obtain the estimate $||\widetilde{N}_*(u)||_2 \lesssim ||\widetilde{N}_*^1(u)||_2 + ||\nabla_{t,x}u||_{\mathcal{Y}}$. Theorem 5.2, Lemma 10.2 and Theorem 9.2 now apply to give the estimate

$$\|N^1_*(u)\|_2 \lesssim \|\tilde{h}^+\|_2 + \|f\|_{\mathcal{Y}} \approx \|\nabla_{t,x}u\|_{\mathcal{Y}}.$$

To see the first estimate, write $\tilde{h}^+ = B_0 h^+$ with $h^+ \in E_0^+ \mathcal{H}$, and apply Theorem 5.2 to get $\|e^{-t\tilde{\Lambda}}B_0h^+\|_{\mathcal{X}} = \|B_0e^{-t\Lambda}h^+\|_{\mathcal{X}} \lesssim \|h^+\|_2 \approx \|\tilde{h}^+\|_2$. The lower estimate follows from Lemma 5.3 since

$$\|\widetilde{N}_*(u)\|_2^2 \gtrsim \lim_{t \to 0} t^{-1} \int_t^{2t} \|u_s\|_2^2 ds = \|u_0\|_2^2.$$

Proof of Lemma 10.2. Before we start, we remark that $p \mapsto \|\widetilde{N}^p_*((\widetilde{S}_A f)_{\perp})\|_2$ is increasing, so it suffices to consider p close to 2. We shall fix the value of p eventually in (iii) below, when we see how close to 2 it need to be. Next it suffices to prove the inequality for $t \mapsto f_t$ compactly supported in \mathbf{R}_+ . Indeed, combining Lemma 5.3 and Proposition 7.2, for all $\epsilon > 0$ and $f \in \mathcal{Y}$ we have (since $p \leq 2$)

$$\begin{split} \|\widetilde{N}_{*}^{p}(\chi_{(\epsilon,\epsilon^{-1})}(t)(\widetilde{S}_{A}f)_{\perp})\|_{2}^{2} &\leq \|\widetilde{N}_{*}(\chi_{(\epsilon,\epsilon^{-1})}(t)(\widetilde{S}_{A}f)_{\perp})\|_{2}^{2} \\ &\lesssim \int_{\epsilon}^{\epsilon^{-1}} \|(\widetilde{S}_{A}f)_{\perp}\|_{2}^{2} \frac{dt}{t} \lesssim \ln\epsilon \sup_{t>0} \|\widetilde{S}_{A}f\|_{2}^{2} \lesssim \ln\epsilon \|f\|_{\mathcal{Y}}^{2}. \end{split}$$

Thus, if $f_{\delta} := \chi_{(\delta, \delta^{-1})}(t) f$ for $f \in \mathcal{Y}$, we have for fixed $\epsilon > 0$

$$\|\widetilde{N}^p_*(\chi_{(\epsilon,\epsilon^{-1})}(t)(\widetilde{S}_A f)_{\perp})\|_2 \le \liminf_{\delta \to 0} \|\widetilde{N}^p_*(\chi_{(\epsilon,\epsilon^{-1})}(t)(\widetilde{S}_A f_{\delta})_{\perp})\|_2.$$

Now our assumption gives

 $\|\widetilde{N}^p_*(\chi_{(\epsilon,\epsilon^{-1})}(t)(\widetilde{S}_A f_{\delta})_{\perp})\|_2 \lesssim \|\mathcal{E}\|_C \|f_{\delta}\|_{\mathcal{Y}} \lesssim \|\mathcal{E}\|_C \|f\|_{\mathcal{Y}},$ uniformly in ϵ , so for all $f \in \mathcal{Y}$ and $\epsilon > 0$ we obtain

 $\|\widetilde{N}^p_*(\chi_{(\epsilon,\epsilon^{-1})}(t)(\widetilde{S}_A f)_{\perp})\|_2 \lesssim \|\mathcal{E}\|_C \|f\|_{\mathcal{Y}}.$

It remains to let $\epsilon \to 0$ and apply the monotone convergence theorem.

(i) We now fix $t \mapsto f_t$ compactly supported in \mathbf{R}_+ and write

$$\widetilde{S}_A f_t = \int_0^t e^{-(t-s)\tilde{\Lambda}} \widetilde{E}_0^+ \mathcal{E}_s f_s ds - \int_t^\infty e^{-(s-t)\tilde{\Lambda}} \widetilde{E}_0^- \mathcal{E}_s f_s ds =: I - II$$

Most of the time we use the pointwise inequality $\widetilde{N}_*^p \leq \widetilde{N}_*$. It is only for one term, estimated in (iii) below, that we require p < 2.

Split the integral I as

$$I = \int_0^t e^{-(t-s)\tilde{\Lambda}} (I - e^{-2s\tilde{\Lambda}}) \widetilde{E}_0^+ \mathcal{E}_s f_s ds + e^{-t\tilde{\Lambda}} \int_0^t e^{-s\tilde{\Lambda}} \widetilde{E}_0^+ \mathcal{E}_s f_s ds = I_1 + I_2.$$

As in the proof of Proposition 7.2, the kernel of I_1 has bounds s/t, giving the estimate

$$(41) \quad \|\widetilde{N}_*(I_1)\|_2^2 \lesssim \int_0^\infty \|I_1\|_2^2 \frac{dt}{t} \lesssim \int_0^\infty \left(\int_0^t \frac{s}{t} \frac{ds}{s}\right) \left(\int_0^t \frac{s}{t} \|\mathcal{E}_s f_s\|_2^2 s ds\right) \frac{dt}{t} \\ \lesssim \int_0^\infty \|\mathcal{E}_s f_s\|_2^2 s ds \le \|\mathcal{E}\|_\infty^2 \|f\|_{\mathcal{Y}}^2.$$

Similarly we split $II = II_1 + II_2$ by writing $e^{-(s-t)\tilde{\Lambda}} = e^{-(s-t)\tilde{\Lambda}}(I - e^{-2t\tilde{\Lambda}}) + e^{-t\tilde{\Lambda}}e^{-s\tilde{\Lambda}}$, and a Schur estimate similar to (41) give the bound for II_1 . Next we write

$$II_2 = e^{-t\tilde{\Lambda}} \int_0^\infty e^{-s\tilde{\Lambda}} \widetilde{E}_0^- \mathcal{E}_s f_s ds - e^{-t\tilde{\Lambda}} \int_0^t e^{-s\tilde{\Lambda}} \widetilde{E}_0^- \mathcal{E}_s f_s ds =: II_3 - II_4.$$

By Theorem 5.2, the term II_3 has bound

$$\begin{aligned} \left\| \widetilde{N}_{*} \left(B_{0} e^{-t\Lambda} B_{0}^{-1} P_{B_{0}\mathcal{H}} \int_{0}^{\infty} e^{-s\tilde{\Lambda}} \widetilde{E}_{0}^{-} \mathcal{E}_{s} f_{s} ds \right) \right\|_{2} &\lesssim \left\| \int_{0}^{\infty} e^{-s\tilde{\Lambda}} \widetilde{E}_{0}^{-} \mathcal{E}_{s} f_{s} ds \right\|_{2} \\ &= \sup_{\|h\|_{2}=1} \left| \int_{0}^{\infty} (\mathcal{E}_{s}^{*} e^{-s\tilde{\Lambda}^{*}} (\widetilde{E}_{0}^{-})^{*} h, f_{s}) ds \right| \lesssim \|\mathcal{E}\|_{*} \|f\|_{\mathcal{Y}}. \end{aligned}$$

(ii) It remains to consider $I_2 + II_4 = (\widetilde{E}_0^+ + \widetilde{E}_0^-)e^{-t\tilde{\Lambda}}\int_0^t e^{-s\tilde{\Lambda}}\mathcal{E}_s f_s ds$. Note that $(\widetilde{E}_0^+ + \widetilde{E}_0^-) = P_{B_0\mathcal{H}}$. Since we only consider the normal component of $I_2 + II_4$ and $(P_{B_0\mathcal{H}}h)_{\perp} = h_{\perp}$ for any h, it remains to estimate $e^{-t\tilde{\Lambda}}\int_0^t e^{-s\tilde{\Lambda}}\mathcal{E}_s f_s ds$. To make use of off-diagonal estimates (see Lemma 10.3), we need to replace $e^{-t\tilde{\Lambda}}$ by the resolvents $(I + itB_0D)^{-1}$. To this end, define $\psi_t(z) := e^{-t|z|} - (1 + itz)^{-1}$ and split the integral

$$e^{-t\tilde{\Lambda}} \int_0^t e^{-s\tilde{\Lambda}} \mathcal{E}_s f_s ds = \psi_t(B_0 D) \int_0^\infty e^{-s\tilde{\Lambda}} \mathcal{E}_s f_s ds - \int_t^\infty \psi_t(B_0 D) e^{-s\tilde{\Lambda}} \mathcal{E}_s f_s ds + \int_0^t (I + itB_0 D)^{-1} (e^{-s\tilde{\Lambda}} - I) \mathcal{E}_s f_s ds + (I + itB_0 D)^{-1} \int_0^t \mathcal{E}_s f_s ds.$$

For the first term, square function estimates show that $\psi_t(B_0D): L_2 \to \mathcal{Y}^* \subset \mathcal{X}$ is continuous, and a duality argument like for II_3 gives the bound. For the second and third terms, we note the operator estimates

$$\|\psi_t(B_0D)e^{-s\tilde{\Lambda}}\| = \left\|\frac{t}{s}\frac{e^{-t|B_0D|} - (I + itB_0D)^{-1}}{tB_0D}(sB_0D)e^{-s|B_0D|}\right\| \lesssim t/s,$$

and

$$\|(I+itB_0D)^{-1}(e^{-s\tilde{\Lambda}}-I)\| \lesssim \left\|\frac{s}{t}\frac{tB_0D}{I+itB_0D}\frac{e^{-s|B_0D|}-I}{sB_0D}\right\| \lesssim s/t.$$

Schur estimates similar to (41) give the N_* bounds.

(iii) It remains to prove the estimate

$$\left\|\widetilde{N}^p_*\left((I+itB_0D)^{-1}\int_0^t \mathcal{E}_s f_s ds\right)\right\|_2 \lesssim \|\mathcal{E}\|_C \|f\|_{\mathcal{Y}}.$$

To show this, fix a Whitney box $W(t_0, x_0)$, take $h \in L_q(W(t_0, x_0); \mathbb{C}^{(1+n)m})$, and let h = 0 outside $W(t_0, x_0)$. Here 1/p + 1/q = 1, p < 2 and q > 2. To bound the $L_p(W(t_0, x_0))$ norm, we do the duality argument

$$\begin{aligned} \frac{1}{t_0} \int_{c_0^{-1}t_0}^{c_0t_0} \left((I + itB_0D)^{-1} \int_0^t \mathcal{E}_s f_s ds, h_t \right) dt \\ &= \int_0^{c_0t_0} \left(\mathcal{E}_s f_s, \frac{1}{t_0} \int_{\max(c_0^{-1}t_0, s)}^{c_0t_0} (I - itDB_0^*)^{-1} h_t dt \right) ds \\ &\leq \int_{\mathbf{R}^n} \int_0^{c_0t_0} |\mathcal{E}(s, y)| |f(s, y)| H(y) ds dy, \end{aligned}$$

where

$$H(y) := \frac{1}{t_0} \int_{c_0^{-1}t_0}^{c_0t_0} |(I - itDB_0^*)^{-1}h_t(y)| dt.$$

To handle the tails of $(I - itDB_0^*)^{-1}h_t$, we split the space into annular regions $\mathbf{R}^n = \bigcup_{k=0}^{\infty} A_k$, where $A_0 := B(x_0; t_0)$ and $A_k := (2^k A_0) \setminus (2^{k-1}A_0)$ for $k \ge 1$. Define $f_k(s, y) := \chi_{(0,c_0t_0)}(s)\chi_{A_k}(y)f(s, y)$ and $H_k(y) := \chi_{A_k}(y)H(y)$. Then Whitney averaging as in the proof of Lemma 5.5 gives

$$\begin{split} \int_{\mathbf{R}^{n}} \int_{0}^{c_{0}t_{0}} |\mathcal{E}(s,y)| |f(s,y)| H(y) ds dy &\leq \sum_{k=0}^{\infty} \iint_{\mathbf{R}^{1+n}_{+}} |\mathcal{E}(s,y)| s |f_{k}(s,y)| H_{k}(y) \frac{ds dy}{s} \\ &\approx \sum_{k=0}^{\infty} \iint_{\mathbf{R}^{1+n}_{+}} \left(\frac{1}{t^{1+n}} \iint_{W(t,x)} |\mathcal{E}(s,y)| s |f_{k}(s,y)| H_{k}(y) ds dy \right) \frac{dt dx}{t} \\ &\lesssim \sum_{k=0}^{\infty} \iint_{\mathbf{R}^{1+n}_{+}} \sup_{W(t,x)} |\mathcal{E}| \left(\frac{1}{t^{1+n}} \iint_{W(t,x)} |sf_{k}|^{2} \right)^{1/2} \left(\frac{1}{t^{1+n}} \iint_{W(t,x)} |H_{k}|^{2} \right)^{1/2} \frac{dt dx}{t} \\ &\lesssim \sum_{k=0}^{\infty} \|\mathcal{E}\|_{C} \int_{\mathbf{R}^{n}} \mathcal{A} \left(\frac{1}{\sqrt{t^{1+n}}} \|sf_{k}\|_{L_{2}(W(t,x))} \frac{1}{\sqrt{t^{1+n}}} \|H_{k}\|_{L_{2}(W(t,x))} \right) (z) dz \\ &\lesssim \sum_{k=0}^{\infty} \|\mathcal{E}\|_{C} \int_{\mathbf{R}^{n}} \mathcal{A} \left(\frac{1}{\sqrt{t^{1+n}}} \|sf_{k}\|_{L_{2}(W(t,x))} \right) (z) N_{*} \left(\frac{1}{\sqrt{t^{1+n}}} \|H_{k}\|_{L_{2}(W(t,x))} \right) (z) dz \\ &\lesssim \sum_{k=0}^{\infty} \|\mathcal{E}\|_{C} \int_{\mathbf{R}^{n}} \mathcal{A} \left(\frac{1}{\sqrt{t^{1+n}}} \|sf_{k}\|_{L_{2}(W(t,x))} \right) (z) N_{*} \left(\frac{1}{\sqrt{t^{1+n}}} \|H_{k}\|_{L_{2}(W(t,x))} \right) (z) dz \\ &\lesssim \sum_{k=0}^{\infty} \|\mathcal{E}\|_{C} \|\mathcal{A}(sf_{k})\|_{L_{p}(\mathbf{R}^{n})} \|M(|H_{k}|^{2})^{1/2}\|_{L_{q}(\mathbf{R}^{n})}. \end{split}$$

Here \mathcal{A} denotes the area function $\mathcal{A}g(z) := (\iint_{|y-z| < cs} |g(s,y)|^2 s^{-(1+n)} ds dy)^{1/2}$ and $N_*g(z) := \sup_{|y-z| < cs} |g(s,x)|$ is the non-tangential maximal function, where $c \in (0,\infty)$ is some fixed constant, and M is the Hardy–Littlewood maximal function.

On the fourth line we used the tent space estimate by Coifman, Meyer and Stein in [17, Theorem 1(a)]. Since $M: L_{q/2} \to L_{q/2}$ is bounded, we have

$$\begin{split} \|M(|H_k|^2)^{1/2}\|_{L_q(\mathbf{R}^n)} &\lesssim \|H\|_{L_q(A_k)} \leq \frac{1}{t_0} \int_{c_0^{-1}t_0}^{c_0t_0} \|(I - itDB_0^*)^{-1}h_t\|_{L_q(A_k)} dt \\ &\lesssim 2^{-km} \frac{1}{t_0} \int_{c_0^{-1}t_0}^{c_0t_0} \|h_t\|_{L_q(B(x_0;c_0t_0))} dt \leq 2^{-km} t_0^{-q} \|h\|_{L_q(W(t_0,x_0))} dt \end{split}$$

The third estimate uses Lemma 10.3 below, and thus is where we choose p < 2 sufficiently close to 2 so that $2 < q < 2 + \delta$. We obtain the maximal function estimate

$$\begin{split} \widetilde{N}_{*}^{p} \left((I + itB_{0}D)^{-1} \int_{0}^{t} \mathcal{E}_{s}f_{s}ds \right) (x_{0}) &\lesssim \|\mathcal{E}\|_{C} \sup_{t_{0} > 0} \sum_{k=0}^{\infty} 2^{-km} t_{0}^{n/q-n} \|\mathcal{A}(sf_{k})\|_{L_{p}(\mathbf{R}^{n})} \\ &\lesssim \|\mathcal{E}\|_{C} \sum_{k=0}^{\infty} 2^{-k(m-n/p)} \sup_{t_{0} > 0} \left(\frac{1}{(2^{k}t_{0})^{n}} \int_{B(x_{0};(2^{k}+cc_{0})t_{0})} |\mathcal{A}(sf)|^{p} dx \right)^{1/p} \\ &\lesssim \|\mathcal{E}\|_{C} \mathcal{M}(\mathcal{A}(sf)^{p})^{1/p} (x_{0}), \end{split}$$

where c is the constant from the definition of \mathcal{A} and m > n/p. Since $M : L_{2/p} \to L_{2/p}$ is bounded, this yields

$$\left\|\widetilde{N}_{*}^{p}\left((I+itB_{0}D)^{-1}\int_{0}^{t}\mathcal{E}_{s}f_{s}ds\right)\right\|_{2}$$

$$\lesssim \|\mathcal{E}\|_{C}\|M(\mathcal{A}(sf)^{p})^{1/p}\|_{2} \lesssim \|\mathcal{E}\|_{C}\|\mathcal{A}(sf)\|_{2} \approx \|\mathcal{E}\|_{C}\|f\|_{\mathcal{Y}}.$$

This completes the proof of the maximal function estimate.

The following lemma, which we used above, is contained in [7, Lemma 2.57]. However, we give a more direct proof here, since the algebraic setup in [7] was quite different.

Lemma 10.3. Let B_0 be t-independent coefficients, strictly accretive on $\mathcal{H} = \overline{R(D)}$. Then for each positive integer m, there is $C_m < \infty$ and $\delta > 0$ such that

$$\|(1+itDB_0)^{-1}f\|_{L_q(E)} \le \frac{C_m}{(1+\operatorname{dist}(E,F)/t)^m} \|f\|_{L_q(F)}$$

for all t > 0 and sets $E, F \subset \mathbf{R}^n$ such that supp $f \subset F$, and all q such that $|q-2| < \delta$. Here dist $(E, F) := \inf\{|x - y| ; x \in E, y \in F\}.$

Proof. For q = 2, these off-diagonal estimates can be proved as in [8, Proposition 5.1], using estimates on commutators with bump functions (and replacing the operator B_0D there by DB_0). By interpolation, it suffices to estimate $||(1 + itDB_0)^{-1}f||_{L_q(\mathbf{R}^n)\to L_q(\mathbf{R}^n)}$, uniformly for t and q in a neighbourhood of 2. To this end, assume that $(I + itDB_0)\tilde{f} = f$. As in Proposition 4.1, but replacing ∂_t by $(it)^{-1}$, this equation is equivalent to

$$\begin{cases} (A_0 \tilde{g})_{\perp} + it \operatorname{div}_x (A_0 \tilde{g})_{\parallel} = (A_0 g)_{\perp}, \\ \tilde{g}_{\parallel} - it \nabla_x \tilde{g}_{\perp} = g_{\parallel}, \end{cases}$$

where A_0, g, \tilde{g} are related to B_0, f, \tilde{f} , respectively, as in Proposition 4.1. Using the second equation to eliminate \tilde{g}_{\parallel} in the first, shows that \tilde{g}_{\perp} satisfies the divergence form equation

$$L\tilde{g}_{\perp} := \begin{bmatrix} 1 & it \operatorname{div}_x \end{bmatrix} A_0(x) \begin{bmatrix} 1 \\ it \nabla_x \end{bmatrix} \tilde{g}_{\perp} = \begin{bmatrix} 1 & it \operatorname{div}_x \end{bmatrix} \begin{bmatrix} A_{\perp \perp} g_{\perp} \\ -A_{\parallel \parallel} g_{\parallel} \end{bmatrix}$$

By the stability result of Sneiberg [44] it follows that the divergence form operator L is an isomorphism $L: W_q^1(\mathbf{R}^n) \to W_q^{-1}(\mathbf{R}^n)$ for $|q-2| < \delta$, giving us the desired estimate

$$\|f\|_q \approx \|\tilde{g}\|_q \lesssim \|\tilde{g}_{\perp}\|_q + t \|\nabla_x \tilde{g}_{\perp}\|_q + \|g_{\parallel}\|_q \lesssim \|g\|_q \approx \|f\|_q.$$

10.2. Square function estimates for \mathcal{X} -solutions under t-regularity for the coefficients. Looking closely at the equation $\operatorname{div}_{t,x}Ag = 0$, it seems unlikely that \mathcal{X} solutions g would in general satisfy the square function estimate $\int_0^\infty \|\partial_t g_t\|_2^2 t dt < \infty$, *i.e.* $\partial_t g_t \in \mathcal{Y}$, when A is t-dependent. More precisely, it writes $\partial_t (A_{\perp \perp} g_{\perp} + A_{\perp \parallel} g_{\parallel}) + d_{\perp \parallel} g_{\parallel} = 0$ $\operatorname{div}_x(A_{\parallel\perp}g_{\perp}+A_{\parallel\parallel}g_{\parallel})=0$, and as ∂_t and multiplication by A do not commute the quantity $\partial_t g_t$ does not arise. We show in the next result that $\partial_t g_t \in \mathcal{Y}$ can be obtained upon a further t-regularity assumption on A. This also improves the regularity of g_t itself. We do not claim this assumption is sharp nor necessary (in particular, it could well be that this regularity on the components $A_{\perp\perp}, A_{\perp\parallel}$ suffices). This regularity assumption is akin to the one in [35], going back to [19], towards A_{∞} property of the L-harmonic measure with respect to surface measure on bounded Lipschitz domains for real elliptic operators. See also [24, 25] for results with smallness assumptions of the derivatives of A. The difference with these works is that we are imposing our coefficients A to be perturbation of "good" t-independent coefficients. So our next result neither contains or is contained in the above cited works. Besides, it is again an a priori estimate on solutions, so it is valid independently of solvability issues.

Theorem 10.4. Let A_0 be t-independent coefficients which are strictly accretive on \mathcal{H} and assume that $||A - A_0||_*$ is sufficiently small.

If A satisfies the t-regularity condition

$$\|t\partial_t A\|_* < \infty$$

then any \mathcal{X} -solution g to the divergence form equation, with coefficients A, with boundary trace g_0 has regularity $\partial_t g_t \in L_2^{\text{loc}}(\mathbf{R}_+; L_2)$ with estimates

$$\|\partial_t g_t\|_{\mathcal{Y}} \lesssim \|g\|_{\mathcal{X}}.$$

We also have estimates $\sup_{t>0} \|g_t\|_2 \approx \|g\|_{\mathcal{X}}$, and $t \mapsto g_t \in L_2$ is continuous with limits $\lim_{t\to 0} \|g_t - g_0\|_2 = 0 = \lim_{t\to\infty} \|g_t\|_2$. The converse estimate $\|g\|_{\mathcal{X}} \lesssim \|\partial_t g\|_{\mathcal{Y}}$ holds for all \mathcal{X} -solutions g, provided $\|t\partial_t A\|_*$ is sufficiently small.

If $\max(\|t\partial_i A\|_*, \|t\partial_t A\|_*) < \infty$ for some i = 1, ..., n, then $\partial_i g_t \in L_2^{\text{loc}}(\mathbf{R}_+; L_2)$ for any \mathcal{X} -solution g to the divergence form equation with coefficients A, with estimates $\|\partial_i g_t\|_{\mathcal{Y}} \leq \|g\|_{\mathcal{X}}$. The estimate $\|g\|_{\mathcal{X}} \leq \|\nabla_x g\|_{\mathcal{Y}}$ holds for all \mathcal{X} -solutions g, provided $\|t\nabla_{t,x}A\|_*$ is sufficiently small.

We do not know whether the smallness assumptions are needed for the converse estimates to hold. We also remark that the same conclusions hold for the conormal gradient f, as is clear from the proof below.

Lemma 10.5. If $h \in \mathcal{X}$ has distribution derivative $\partial_t h \in \mathcal{Y}$, then $\partial_t(S_A h) \in \mathcal{Y}$ with estimates

$$\|\partial_t(S_A h)\|_{\mathcal{Y}} \lesssim (\|\mathcal{E}\|_* + \|t\partial_t \mathcal{E}\|_*)\|h\|_{\mathcal{X}} + \|\mathcal{E}\|_{\infty}\|\partial_t h\|_{\mathcal{Y}}$$

Proof of Theorem 10.4 modulo Lemma 10.5. (i) As in the proof of Corollary 8.5, any \mathcal{X} -solution can be written $g = [(Bf)_{\perp}, f_{\parallel}]^t$, where

$$(I - S_A)f = e^{-t\Lambda}h^+, \quad \text{for some } h^+ \in E_0^+\mathcal{H}.$$

Introduce the auxiliary Banach space $Z := \{h \in \mathcal{X} : \partial_t h \in \mathcal{Y}\} \subset \mathcal{X}$, with norm $\|h\|_Z := \|h\|_{\mathcal{X}} + a\|\partial_t h\|_{\mathcal{Y}}$. By Proposition 7.1 and Lemma 10.5 we have estimates $\|S_A h\|_{\mathcal{X}} \leq C \|h\|_{\mathcal{X}}$ and $\|\partial_t (S_A h)\|_{\mathcal{Y}} \leq D \|h\|_{\mathcal{X}} + C \|\partial_t h\|_{\mathcal{Y}}$, where we assume C < 1, and we choose the parameter a > 0 small enough so that

$$\|S_A\|_{Z \to Z} \le C + aD < 1.$$

Hence $I - S_A$ is invertible on both \mathcal{X} and Z. Since $e^{-t\Lambda}h^+ \in Z$ by Theorem 5.2, we conclude that $f \in Z$ with estimates $\|\partial_t f\|_{\mathcal{Y}} \lesssim \|f\|_Z \approx \|e^{-t\Lambda}h^+\|_Z \approx \|h^+\|_2$. For the gradient g, this gives the bound $\|\partial_t g\|_{\mathcal{Y}} \lesssim \|t\partial_t B\|_* \|f\|_{\mathcal{X}} + (\|B\|_{\infty} + 1) \|\partial_t f\|_{\mathcal{Y}} \lesssim \|h^+\|_2 \approx \|f\|_{\mathcal{X}} \approx \|g\|_{\mathcal{X}}$.

(ii) To prove the sup $-L_2$ estimate and trace result for g_t , write $\int_0^\infty s\eta(s)\partial_s g_s ds = \int_0^\infty (\eta(s) + s\eta'(s))g_s ds$, for some $\eta \in C_0^\infty(\mathbf{R}_+)$. Take the limit as η approaches the characteristic function for (0, t) to get

$$g_t = \frac{1}{t} \int_0^t g_s ds + \frac{1}{t} \int_0^t \partial_s g_s s ds, \qquad \text{a.e. } t > 0.$$

The last term has bound $(\int_0^t \|\partial_s g_s\|^2 s ds)^{1/2}$, whereas the first term satisfies

$$\left\|\frac{1}{t}\int_{0}^{t}g_{s}ds - g_{0}\right\|_{2}^{2} \leq \sum_{k=1}^{\infty} 2^{-k} \left(\frac{1}{2^{-k}t}\int_{2^{-k}t}^{2^{1-k}t} \|g_{s} - g_{0}\|_{2}^{2}ds\right) \to 0$$

as $t \to 0$. Hence the trace claims follow from the square function estimates $\|\partial_t g_t\|_{\mathcal{Y}} < \infty$. Moreover, the estimate $\sup_{t>0} \|g_t\|_2 \lesssim \|g\|_{\mathcal{X}} + \|\partial_t g\|_{\mathcal{Y}} \lesssim \|g\|_{\mathcal{X}}$ follows. The converse estimate follows from Theorem 8.2.

An integration by parts, similar to above, shows that

$$2g_{2t} = g_t + \frac{1}{t} \int_t^{2t} g_s ds + \frac{1}{t} \int_t^{2t} \partial_s g_s s ds, \quad \text{a.e. } t > 0.$$

Taking $\limsup_{t\to\infty}$ of both sides, shows $2\limsup_{t\to\infty} \|g_t\|_2 = \limsup_{t\to\infty} \|g_t\|_2$. Since $\|g_t\|_2$ is bounded, we conclude that $\lim_{t\to\infty} \|g_t\|_2 = 0$.

(iii) To show $||g||_{\mathcal{X}} \lesssim ||\partial_t g||_{\mathcal{Y}}$, consider f satisfying $e^{-t\Lambda}h^+ = f_t - S_A f_t$. Theorem 5.2 and Lemma 10.5 give

$$\|h^+\|_2 \approx \|\partial_t e^{-t\Lambda}h^+\|_{\mathcal{Y}} \lesssim \|\partial_t f\|_{\mathcal{Y}} + (\|\mathcal{E}\|_* + \|t\partial_t A\|_*)\|f\|_{\mathcal{X}} + \|\mathcal{E}\|_{\infty}\|\partial_t f\|_{\mathcal{Y}},$$

where by Theorem 8.2 we have $||f||_{\mathcal{X}} \approx ||h^+||_2$ as $||\mathcal{E}||_*$ is assumed small enough. If in addition $||t\partial_t A||_*$ is sufficiently small, then we obtain $||f||_{\mathcal{X}} \lesssim ||\partial_t f||_{\mathcal{Y}}$. As in (i), again using smallness of $||t\partial_t A||_*$, this implies $||g||_{\mathcal{X}} \lesssim ||\partial_t g||_{\mathcal{Y}}$.

(iv) To prove the x-regularity result, consider the equation $\partial_t f + DBf = 0$, which implies

$$\|\partial_t f\|_{\mathcal{Y}} = \|DP_{\mathcal{H}}Bf\|_{\mathcal{Y}} \approx \sum_{i=1}^n \|(P_{\mathcal{H}}B)(\partial_i f) + P_{\mathcal{H}}(\partial_i B)f\|_{\mathcal{Y}}$$

since $D = DP_{\mathcal{H}}$ and the operator D has estimates $||Dh||_2 \approx \sum_{i=1}^n ||\partial_i h||_2$ for all $h \in \mathsf{D}(D) \cap \mathcal{H}$. (The latter is straightforward to verify with the Fourier transform.) Here $P_{\mathcal{H}}$ denotes orthogonal projection onto \mathcal{H} ; it commutes with ∂_i . This yields the bound

$$\|\partial_i f\|_{\mathcal{Y}} \approx \|(P_{\mathcal{H}}B)\partial_i f\|_{\mathcal{Y}} \lesssim \|\partial_t f\|_{\mathcal{Y}} + \|t(\partial_i B)f\|_{\mathcal{Y}^*} \lesssim (1 + \|t\partial_i B\|_*)\|f\|_{\mathcal{X}} \lesssim \|f\|_{\mathcal{X}}$$

if $\max(\|t\partial_i A\|_*, \|t\partial_t A\|_*) < \infty$, where we used that $P_{\mathcal{H}}B_t : \mathcal{H} \to \mathcal{H}$ is an isomorphism in the first comparison. Conversely, if $\|t\partial_t A\|_*$ is sufficiently small, then

$$\|f\|_{\mathcal{X}} \lesssim \|\partial_t f\|_{\mathcal{Y}} \lesssim \sum_{i=1}^n (\|\partial_i f\|_{\mathcal{Y}} + \|t\partial_i B\|_* \|f\|_{\mathcal{X}}),$$

where the first estimate is by (iii). Using next that $\sum_{i=1}^{n} ||t\partial_i B||_*$ is small enough, this implies $||f||_{\mathcal{X}} \lesssim ||\nabla_x f||_{\mathcal{Y}}$.

As in (i) above, these estimates translate to $\|\partial_i g\|_{\mathcal{Y}} \lesssim \|g\|_{\mathcal{X}}$ and $\|g\|_{\mathcal{X}} \lesssim \|\nabla_x g\|_{\mathcal{Y}}$ respectively.

Proof of Lemma 10.5. Assume that the coefficients A satisfy $||A-A_0||_* < \infty$ and has distribution derivative $\partial_t A \in L^{\text{loc}}_{\infty}(\mathbf{R}^{1+n}_+; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ such that $||t\partial_t A||_* < \infty$. Fix $h \in \mathcal{X}$ with distribution derivative $\partial_t h \in \mathcal{Y}$. By Theorem 7.1, $\int_a^b ||S_A h_t - S_A^\epsilon h_t||_2^2 dt \to 0$ as $\epsilon \to 0$, where

$$S_A^{\epsilon}h_t := \int_0^t \eta_{\epsilon}^+(t,s)\Lambda e^{-(t-s)\Lambda}\widehat{E}_0^+ \mathcal{E}_s h_s ds - \int_t^\infty \eta_{\epsilon}^-(t,s)\Lambda e^{-(s-t)\Lambda}\widehat{E}_0^- \mathcal{E}_s h_s ds = I - II.$$

Hence it suffices to bound $\|\partial_t(S_A^{\epsilon}h)\|_{\mathcal{Y}}$, uniformly for $\epsilon > 0$.

(i) Differentiate I and write

$$t\partial_t(I) = \int_0^t (t\partial_t\eta_\epsilon^+)\Lambda e^{-(t-s)\Lambda} \widehat{E}_0^+ \mathcal{E}_s h_s ds - \int_0^t \eta_\epsilon^+ (t-s)\Lambda^2 e^{-(t-s)\Lambda} \widehat{E}_0^+ \mathcal{E}_s h_s ds$$
$$- \int_0^t \eta_\epsilon^+ (\partial_s\Lambda e^{-(t-s)\Lambda}) \widehat{E}_0^+ (s\mathcal{E}_s h_s) ds = \int_0^t (t\partial_t\eta_\epsilon^+ + s\partial_s\eta_\epsilon^+)\Lambda e^{-(t-s)\Lambda} \widehat{E}_0^+ \mathcal{E}_s h_s ds$$
$$- \int_0^t \eta_\epsilon^+ (t-s)\Lambda^2 e^{-(t-s)\Lambda} \widehat{E}_0^+ \mathcal{E}_s h_s ds + \int_0^t \eta_\epsilon^+ \Lambda e^{-(t-s)\Lambda} \widehat{E}_0^+ \partial_s (s\mathcal{E}_s h_s) ds = I_1 - I_2 + I_3.$$

Note that in I_3 the distribution derivative $\partial_s(s\mathcal{E}_sh_s)$ extends its action to test functions $s \mapsto (\eta_{\epsilon}^+(t,s)\lambda e^{-(t-s)\Lambda}\widehat{E}_0^+)^*\phi$, for any $\phi \in \mathcal{H}$. Theorem 6.5 and Lemma 5.5 give the estimate

$$\|I_3\|_{\mathcal{Y}^*} \lesssim \|\partial_t (t\mathcal{E}_t h_t)\|_{\mathcal{Y}^*} \lesssim (\|\mathcal{E}\|_* + \|t\partial_t \mathcal{E}\|_*)\|h\|_{\mathcal{X}} + \|\mathcal{E}\|_{\infty}\|\partial_t h\|_{\mathcal{Y}}.$$

To bound I_2 , we apply Lemma 6.7, using the bounds

$$\int_0^t |(t-s)\lambda^2 e^{-(t-s)\lambda}| s ds \lesssim t \quad \text{and} \quad \int_s^\infty |(t-s)\lambda^2 e^{-(t-s)\lambda}| dt \lesssim 1,$$

which shows $||I_2||_{\mathcal{Y}^*} \lesssim ||\mathcal{E}h||_{\mathcal{Y}^*} \lesssim ||\mathcal{E}||_* ||h||_{\mathcal{X}}$. To estimate I_1 , we calculate

$$(t\partial_t + s\partial_s)\eta_{\epsilon}^+(t,s) = \frac{t-s}{\epsilon}(\eta^0)'(\frac{t-s}{\epsilon})\eta_{\epsilon}(t)\eta_{\epsilon}(s) + \eta^0(\frac{t-s}{\epsilon})(t\eta_{\epsilon}'(t)\eta_{\epsilon}(s) + s\eta_{\epsilon}(t)\eta_{\epsilon}'(s)).$$

From this, we verify that $|(t\partial_t + s\partial_s)\eta_{\epsilon}^+| \lesssim \chi_{\operatorname{supp} \nabla \eta_{\epsilon}^+} \leq 1$. Hence an estimate as in the proof of Theorem 6.5 shows that $||I_1||_{\mathcal{Y}^*} \lesssim ||\mathcal{E}||_* ||h||_{\mathcal{X}}$.

(ii) Next we differentiate II and write

$$\begin{split} t\partial_t(II) &= \int_t^\infty (t\partial_t\eta_\epsilon^-)\Lambda e^{-(s-t)\Lambda} \widehat{E}_0^- \mathcal{E}_s h_s ds - \int_t^\infty t\eta_\epsilon^- (\partial_s\Lambda e^{-(s-t)\Lambda}) \widehat{E}_0^- \mathcal{E}_s h_s ds \\ &= \int_t^\infty t(\partial_t\eta_\epsilon^- + \partial_s\eta_\epsilon^-)\Lambda e^{-(s-t)\Lambda} \widehat{E}_0^- \mathcal{E}_s h_s ds + \int_t^\infty \eta_\epsilon^- \frac{t}{s} \Lambda e^{-(s-t)\Lambda} \widehat{E}_0^- s \partial_s (\mathcal{E}_s h_s) ds \\ &= II_1 + II_2. \end{split}$$

To bound II_2 , we apply Lemma 6.7 using the bounds

$$\int_{t}^{\infty} |(t/s)\lambda e^{-(s-t)\lambda}| s ds \lesssim t \quad \text{and} \quad \int_{0}^{s} |(t/s)\lambda e^{-(s-t)\lambda}| dt \lesssim 1,$$

which shows $\|II_2\|_{\mathcal{Y}^*} \lesssim \|t\partial_t \mathcal{E}\|_* \|h\|_{\mathcal{X}} + \|\mathcal{E}\|_{\infty} \|\partial_t h\|_{\mathcal{Y}}$. To estimate II_1 , we calculate

$$t(\partial_t + \partial_s)\eta_{\epsilon}^+(t,s) = t\eta^0(\frac{t-s}{\epsilon})(\eta_{\epsilon}'(t)\eta_{\epsilon}(s) + \eta_{\epsilon}(t)\eta_{\epsilon}'(s)).$$

The last term is supported on $s \in (1/(2\epsilon), 1/\epsilon)$, $t \in (\epsilon, s - \epsilon)$, where it is bounded by $\epsilon t \leq t/s$. Thus estimates as for II_2 apply. The first term is supported on $t \in (\epsilon, 2\epsilon)$, $s \in (t + \epsilon, 1/\epsilon)$ (and another component which can be taken together with the last term) and is bounded by 1. Splitting this remaining term as in (31), it suffices to estimate

$$\begin{aligned} \left\| \chi_{(\epsilon,2\epsilon)}(t)t\eta_{\epsilon}'(t)e^{-t\Lambda} \int_{0}^{\infty} \eta_{\epsilon}(s)\Lambda e^{-s\Lambda}\widehat{E}_{0}^{-}\mathcal{E}_{s}h_{s}ds \right\|_{\mathcal{Y}^{*}} \\ \lesssim \left(\frac{1}{\epsilon} \int_{\epsilon}^{2\epsilon} \left\| e^{-t\Lambda} \int_{0}^{\infty} \eta_{\epsilon}(s)\Lambda e^{-s\Lambda}\widehat{E}_{0}^{-}\mathcal{E}_{s}h_{s}ds \right\|_{2}^{2} dt \right)^{1/2} \\ \lesssim \left\| \int_{0}^{\infty} \eta_{\epsilon}(s)\Lambda e^{-s\Lambda}\widehat{E}_{0}^{-}\mathcal{E}_{s}h_{s}ds \right\|_{2} \lesssim \|\mathcal{E}h\|_{\mathcal{Y}^{*}} \lesssim \|\mathcal{E}\|_{*}\|h\|_{\mathcal{X}} \end{aligned}$$

using the uniform boundedness of $e^{-t\Lambda}$ and Lemma 6.9. This completes the proof.

11. MISCELLANEOUS REMARKS AND OPEN QUESTIONS

(i) The condition $\widetilde{N}_*(\nabla_{t,x}u) \in L_2$ implies that Whitney averages $\frac{1}{|W(t,y)|} \iint_{W(t,y)} u$ converge non-tangentially for almost every x, *i.e.* with $|y - x| < \alpha t$ for some $\alpha < \infty$, to some $u_0(x)$ with u_0 belonging to the closure of $C_0^{\infty}(\mathbb{R}^n)$ with respect to $\|\nabla_x f\|_2 < \infty$. Furthermore, $t^{-1} \int_t^{2t} \nabla_x u_s ds$ converges weakly to $\nabla_x u_0$ in L_2 as $t \to 0$ (compare Theorem 2.3(i)). In particular $\|\nabla_x u_0\|_2 \leq \|\widetilde{N}_*(\nabla_{t,x}u)\|_2$. This is essentially in [33, p. 461-462], where it is done on the unit ball instead of the upper half space, and with pointwise values instead of averages, working with u's solving a real symmetric equation. However, the result has nothing to do with BVPs, but is a result on a function space.

(ii) Assume that $A \in L_{\infty}(\mathbf{R}^{1+n}_{+}; \mathcal{L}(\mathbf{C}^{(1+n)m}))$ and that $\widetilde{N}_{*}(\nabla_{t,x}u) \in L_{2}$ with u satisfying (1) in \mathbf{R}^{1+n}_{+} distributional sense. Then there exists $g \in \dot{H}^{-1/2}(\mathbf{R}^{n}; \mathbf{C}^{m})$ such that

(42)
$$\iint_{\mathbf{R}^{1+n}_+} (A\nabla_{t,x}u, \nabla_{t,x}\phi) dt dx = (g, \phi|_{\mathbf{R}^n}), \quad \text{for all } \phi \in C_0^\infty(\mathbf{R}^{1+n}; \mathbf{C}^m).$$

If $\partial_{\nu_A} u(s,x) := (A\nabla_{t,x} u(s,x))_{\perp}$ for all $s > 0, x \in \mathbf{R}^n$, then $t^{-1} \int_t^{2t} \partial_{\nu_A} u_s ds$ converges weakly to -g in L_2 as $t \to 0$. In particular $||g||_2 \leq ||\widetilde{N}_*(\nabla_{t,x} u)||_2$. This is again essentially [33] for the unit ball. See [4, Lemma 4.3(iii)] for an argument in \mathbf{R}^{1+n}_+ . The equality (42) justifies that g is called the Neumann data. This result has nothing to do with accretivity of A, boundedness suffices. Compare again Theorem 2.3(i).

(iii) Theorem 2.4(i) contains a priori estimates on \mathcal{Y} -solutions. A natural question is to reverse the *a priori* estimates for such systems. Does a weak solution to (1) with $||A - A_0||_C < \infty$ and $\widetilde{N}_*(u) \in L_2$ satisfy $||\nabla_{t,x}u||_{\mathcal{Y}} \leq ||\widetilde{N}_*(u)||_2$? Same question replacing $\widetilde{N}_*(u) \in L_2$ with $\sup_{t>0} ||u_t||_2 < \infty$. The smallness of $||A - A_0||_C$, which implies well-posedness of the Dirichlet problem for \mathcal{Y} -solutions, yields a posteriori such estimates. It would be interesting to have positive answers a priori (*i.e.* independently of well-posedness) when $||A - A_0||_C < \infty$.

(iv) Is there existence of \mathcal{X} -solutions to the Neumann and regularity problems with L_2 data under $||A - A_0||_C < \infty$ (or even under the stronger $\int_0^\infty \omega_A(t)^2 dt/t < \infty$, where $\omega_A(t) := \sup_{0 < s < t} ||A_s - A_0||_{\infty}$)? Is there uniqueness under the same constraint on A, provided existence holds? Recall that tools such as Green's functions are not available here.

(v) Same questions for \mathcal{Y} -solutions and the Dirichlet problem with L_2 data.

(vi) It is likely that \mathcal{Y} -solutions have the a.e. non-tangential convergence property for averages: $\frac{1}{|W(t,y)|} \iint_{W(t,y)} u \to u_0(x)$ for a.e. $x \in \mathbf{R}^n$ and $(t,y) \to (0,x)$ in $|y-x| < \alpha t$. This requires an argument which we leave open.

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