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Monte Carlo Study of the Spin Transport in Magnetic Materials

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Abstract

The resistivity in magnetic materials has been theoretically shown to depend on the spin-spin correlation function which in turn depends on the magnetic-field, the density of conduction electron, the magnetic ordering stability, etc. However, these theories involved a lot of approximations, so their validity remained to be confirmed. The purpose of this work is to show by newly improved extensive Monte Carlo (MC) simulation the resistivity of the spin resistivity from low- T ordered phase to high- T paramagnetic phase in ferromagnetic and antiferromagnetic films. We take into account the interaction between the itinerant spins and the localized lattice spins as well as the interaction between itinerant spins themselves. We show that in ferromagnets the resistivity shows a sharp peak at the magnetic phase transition in agreement with previous theories in spite of their numerous approximations. Resistivity in antiferromagnets on the other hand shows no peak for the SC, BCC and diamond lattices. Discussion on the origin of these resistivity behaviors is given.

Keywords: spin transport, Monte Carlo simulation, magnetic resistivity, magnetic materials

PACS: 72.25.-b, 75.47.-m

1. Introduction

The magnetic resistivity has been extensively studied by both theories and experiments in the last fifty years. Experiments have shown that the resistivity indeed depends on the itinerant spin orientation and the lattice spin ordering [1–6]. At low temperature (T), the main magnetic scattering is due to spin-wave excitations [7, 8]. The resistivity is proportional to T^2 . However at higher T the spin-wave theory is not valid, such a calculation of the resistivity is not possible, in particular in the critical region around the Curie temperature T_c in simple ferromagnets, let alone other complicated magnetic orderings. Experiments on various magnetic materials have found in particular an anomalous behavior of the resistivity at the critical temperature where the system undergoes the ferromagnetic-paramagnetic phase transition [2–6]. Very recent experiments such as those performed on ferromagnetic SrRuO₃ thin films[9], superconducting BaFe₂As₂ single crystals[10], La_{1-x}Sr_xMnO₃[11], Mn_{1-x}Cr_xTe[12] and other compounds[13–16] show different forms of anomaly of the magnetic resistivity at the transition temperature. de Gennes and Friedel's first explanation in 1958[17] for the resistivity behavior near T_c was based on the interaction between the spins of conduction electrons and the lattice spins. The resistivity was thus expected to depend

strongly on the spin ordering of the system. They have suggested that the magnetic resistivity is proportional to the spin-spin correlation, therefore it should behave as the magnetic susceptibility with a divergence at T_c due to "long-range" fluctuations of the magnetization. Other authors [18–20] subsequently suggested that the shape of the resistivity results mainly from "short-range" correlation at $T \geq T_c$. Fisher and Langer [19] have shown in particular that the form of the resistivity cusp depends on the correlation range. To see more details on the role of the spin-spin correlation, we quote a work by Haas[21] and a more recent work of Kataoka[22] where the spin-spin correlation function has been calculated. Recently, Zarand et al [23] have used the picture that the itinerant spin is mainly scattered by impurities which are characterized by a "localization length" in the sense of Anderson's localization. They found that the peak's height depends on this localization length. Note that since the giant magnetoresistance (GMR) was discovered experimentally twenty years ago in magnetic multilayers [24, 25], intensive investigations on the spin resistivity, both experimentally and theoretically, have been carried out[26, 27]. The "spintronics" was born with a spectacular rapid development in relation with industrial applications. For recent overviews, the reader is referred to Refs. [28] and [29]. In spite of these intensive investigations, except our works[30, 31], there have been no Monte Carlo (MC) simulations performed regarding the temperature dependence of the spin transport. In these works, we have investigated by MC simulations the

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effects of magnetic ordering on the spin current in magnetic multilayers.

In this paper we improve our previous MC simulations to study the transport of itinerant electrons in ferromagnetic and antiferromagnetic crystals. We use the Ising model and take into account interactions between lattice spins and itinerant spins. We show that in ferromagnets we obtain, with our new MC averaging method, much better results for the magnetic resistivity which shows a huge peak at the transition temperature while in antiferromagnets, the resistivity does not show such a peak.

The paper is organized as follows. Section 2 is devoted to the description of our model and the rules that govern its dynamics. In section 3, we describe our MC method and discuss the results we obtained for ferromagnets. Results on antiferromagnets are shown in Section 4. Concluding remarks are given in Section 5.

2. Model

2.1. Interactions

We consider a thin film of FCC structure with two symmetrical (001) surfaces. The total number of cells is $N_x \times N_y \times N_z$ where each cell has four spins. Spins localized at FCC lattice sites are called "lattice spins" hereafter. They interact with each other through the following Hamiltonian

$$\mathcal{H}_l = -J \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j, \quad (1)$$

where \mathbf{S}_i is the Ising spin at lattice site i , $\sum_{\langle i,j \rangle}$ indicates the sum over every nearest-neighbor (NN) spin pair $(\mathbf{S}_i, \mathbf{S}_j)$, J being the NN interaction. We consider in this paper both $J > 0$ (ferromagnets) and $J < 0$ (antiferromagnets).

We consider a flow of itinerant spins interacting with each other and with the lattice spins. The interaction between itinerant spins is defined by

$$\mathcal{H}_m = - \sum_{\langle i,j \rangle} K_{i,j} \mathbf{s}_i \cdot \mathbf{s}_j, \quad (2)$$

where \mathbf{s}_i is the itinerant Ising spin at position \vec{r}_i , and $\sum_{\langle i,j \rangle}$ denotes a sum over every spin pair $(\mathbf{s}_i, \mathbf{s}_j)$. The interaction $K_{i,j}$ depends on the distance between the two spins, i.e. $r_{ij} = |\mathbf{r}_i - \mathbf{r}_j|$. A specific form of $K_{i,j}$ will be chosen below. The interaction between itinerant spins and lattice spins is given by

$$\mathcal{H}_r = - \sum_{\langle i,j \rangle} I_{i,j} \mathbf{s}_i \cdot \mathbf{S}_j, \quad (3)$$

where the interaction $I_{i,j}$ depends on the distance between the itinerant spin \mathbf{s}_i and the lattice spin \mathbf{S}_i . For the sake

of simplicity, we assume the same form for $K_{i,j}$ and $I_{i,j}$, namely,

$$K_{i,j} = K_0 \exp(-r_{ij}) \quad (4)$$

$$I_{i,j} = I_0 \exp(-r_{ij}) \quad (5)$$

where K_0 and I_0 are constants.

2.2. Monte Carlo Method

Before calculating the resistivity, we determine the critical temperature T_c below which the system is in the ordered phase using Eq. (1). To this end, we perform standard Metropolis MC simulations to determine various physical quantities at different T [32].

Once the lattice is equilibrated at a given T , we inject N_0 itinerant spins into the system. The itinerant spins move in the x direction under the effect of an electric field. We use the periodic boundary conditions to ensure that the average density of itinerant spins remains constant with evolving time (stationary regime).

Note that unlike in the previous works[30, 31] where the lattice spin configuration is frozen while calculating the resistivity, we use here several thousands of configurations in each of which the resistivity is averaged with many thousands of passages. In our previous works, though the overall number of MC steps per spin was as high as in this work, the fact that we have used only a dozen lattice configurations for resistivity calculation has shown strong fluctuations in the result. In this work, we have made a new device in two steps:

i) For each lattice configuration all itinerant spins move through the system during typically one thousand MC steps. Then we thermalize again the lattice for several thousands of MC steps before continuing the averaging of the resistivity for another thousand MC steps per spin. We repeat this cycle for 200 times. In doing so, each itinerant spin was averaged over 10^6 MC step using "uncorrelated" 200 lattice configurations in all.

ii) The resistivity R is defined as $R = \frac{1}{n}$ where n is the number of itinerant spins crossing a unit area perpendicular to the x direction per unit of MC time. To know this number, we count them at three "detector" surfaces perpendicular to the x direction: the first at $N_x/4$, the second at $N_x/2$ and the third at $3N_x/4$. Averaging the resistivity at these three system positions helps to improve further the results (in our previous works[31] we counted them only at the end of the sample).

As will be shown below, these extensive configuration and space averages give much better results with respect to those in previous works. The dynamics of itinerant spins is governed by the following interactions:

i) an electric field \mathbf{E} which is applied in the x direction. Its energy is given by

$$\mathcal{H}_E = -e\mathbf{E} \cdot \mathbf{r}_i, \quad (6)$$

where \mathbf{r}_i is the distance traveled by the itinerant spin \mathbf{s}_i in a MC step, e its charge. The orientation of \mathbf{r}_i is taken

at random, its magnitude is taken from a uniform distribution between 0 and r_0 where r_0 is the nearest-neighbor distance;

ii) a chemical potential term ("concentration gradient" effect) given by

$$\mathcal{H}_c = Dn(\mathbf{r}), \quad (7)$$

where $n(\mathbf{r})$ is the concentration of itinerant spins in a sphere of radius D_2 centered at \mathbf{r} . D is a constant taken equal to K_0 for simplicity;

iii) interactions between a given itinerant spin and lattice spins inside a sphere of radius D_1 (Eq. 3);

iv) interactions between a given itinerant spin and other itinerant spins inside a sphere of radius D_2 (Eq. 2).

The simulation is carried out as follows: at a given T we calculate the energy of an itinerant spin \mathbf{s}_i by taking into account all the interactions described above. Then we tentatively move the spin under consideration to a new position with a step \mathbf{r}_i in an arbitrary direction. Note that this move is immediately rejected if the new position is inside a sphere of radius Δ_0 centered at a lattice spin or an itinerant spin. This excluded space represents the Pauli exclusion principle in the one hand, and the interaction with lattice phonons on the other hand. If the new position does not lie in a forbidden region of space, then the move is accepted with a probability given by the standard Metropolis algorithm[32].

3. Results on Ferromagnetic Thin Films

We let N_0 itinerant spins travel through the system several thousands of passages until a steady state is reached before averaging the spin resistivity.

The parameters we use in our calculations for ferromagnets are $s = S = 1$ and $N_x = N_y = 20$ and $N_z = 8$. Other parameters are $D_1 = D_2 = 1$ (all distances are in unit of the FCC cell length), $K_0 = 1$ ($I_0 = 0.5$ and $D = 0.5$ unless otherwise stated, $N_0 = 8 \times 20^2$ (namely one itinerant spin per FCC unit cell), $\Delta_0 = 0.05$ and $r_0 = \sqrt{2}/2$, the FCC nearest-neighbor distance. At each T the equilibration time for the lattice spins lies around 10^6 MC steps per spin and we compute statistical averages for the resistivity over 10^6 MC steps per spin. Note that, as described in subsection 2.2, in this work the averaging length 10^6 MC steps per spin has been divided into 200 segments of 5000 MC steps; between two consecutive segments we thermalize again our lattice over several thousands of MC steps per spin to explore a maximum of lattice configurations encountered by itinerant spins. In doing so we reduce statistical fluctuations observed in our old works[30, 31].

We show in Fig. 1 the lattice magnetization versus T for $N_z = 8$, $N_x = N_y = 20$, $J = 1$. We find $T_c \simeq 9.58$ for the critical temperature of the lattice spins.

We show in Fig. 2 the resistivity versus T without magnetic field for several values of I_0 , interaction between itinerant spins and lattice spins. Several remarks are in order:

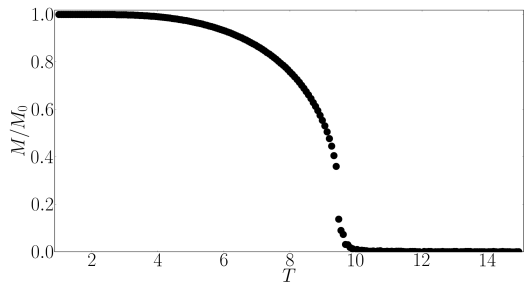


Figure 1: Lattice magnetization versus temperature T for $N_z = 8$, $N_x = N_y = 20$. T_c is $\simeq 9.58$ in unit of $J = 1$.

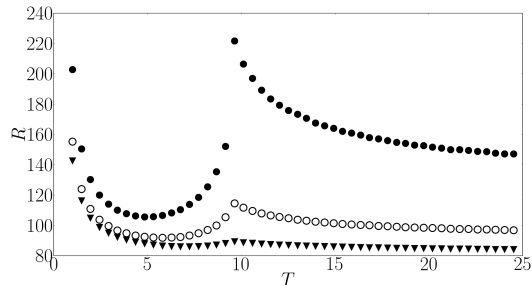


Figure 2: Resistivity R in arbitrary unit versus temperature T for several values of I_0 : 2 (black circles), 1 (void circles), 0.5 (black triangles). Other parameters: $N_x = N_y = 20$, $N_z = 8$, $E = 1$, $K_0 = 0.5$, $D = 0.5$, $D_1 = 1$.

i) As seen here, no significant fluctuations of the data are observed in the whole temperature range, thanks to our new averaging device (see subsection 2.2).

ii) At T_c , R exhibits a peak at the transition temperature. The height of the peak decreases with decreasing I_0 . We see thus that the peak is a consequence of the interaction between itinerant spins and lattice spins. The resistivity shows almost no peak for $I_0 = 0.5$. This case corresponds to a metal where the interaction between itinerant and lattice spins is very weak.

iii) We can explain the existence of the peak by the following argument: the peak is due to the coupling through I_0 of itinerant spins to the fluctuations of the lattice spins in the critical region around T_c . In our recent work, we found from our MC simulation[31] that the resistivity's peak is due to the scattering by antiparallel-spin clusters which exist when one enters the critical region. Below the transition temperature, there exists a single large cluster of lattice spins with some isolated "defects" (i. e. clusters of antiparallel spins), so that the resistance decreases with decreasing T just after T_c .

iv) However, at very low T , the resistivity increases with decreasing T . The origin of this behavior comes from the freezing of the itinerant spins due to their interaction with lattice spins and with themselves. This is very similar to the crystallization of interacting particles at low T . We have tested this interpretation by reducing the strength of the interactions K_0 and I_0 . As a matter of fact, R increases more slowly with decreasing T . This is seen in Fig. 2 at low

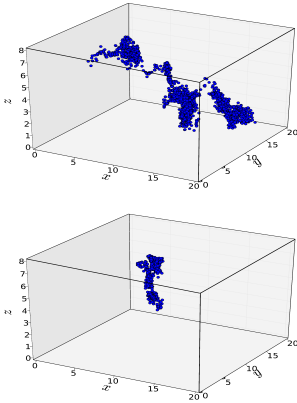


Figure 3: Travel path of an itinerant spin at $T = 5$ (upper) and at $T = 9.79$ (lower) for $I_0 = 2$. Other parameters are the same as in Fig. 2. See text for comments.

T where R is smaller for smaller I_0 . Note that the increase of R at very low T was observed in many experiments on various materials not limited to ferromagnets[10–12, 14].

v) In the paramagnetic phase, as T increases, small clusters will be broken more and more into single disordered spins, so that there is no more energy barrier between successive positions of itinerant spins on their trajectory. The resistance, though high, is thus decreasing with increasing T and saturated as $T \rightarrow \infty$.

During the simulation, we have followed one itinerant spin among 3200 and recorded its successive positions. We show in Fig. 3 its travel path at $T = 5$ where R is very low and at $T = 9.79$ where R is highest. As seen, even at $T = 5$, the spin spends a lot of time to overcome the scattering by lattice spins. Almost four thousands of trial moves are needed to get through the system. Meanwhile, at $T = 9.79$, the spin under consideration spends all this time in a small area. Note that at this peak's temperature, very few spins can go until the end unless we increase the electric field.

3.1. Effect of magnetic field.

Kataoka[22] has shown by a Boltzmann's equation formalism that the magnetic field reduces the peak's height. This is what we observed in simulations. We show results of R for several fields in Fig. 4. The peak reduction is stronger for stronger fields. This is easily understood: when a magnetic field is applied on a ferromagnet, the phase transition is suppressed because the magnetization is not zero at any T . The field reduces critical fluctuations, and hence the number of clusters of antiparallel spins. The peak of the resistivity is therefore reduced and disappears at high fields.

To close this section, let us emphasize that in ferromagnets, we have improved the previous results[31] by using a new averaging procedure. We found that the height of the peak is intimately related to the strength of the interaction between itinerant spins and lattice spins (Fig. 2).

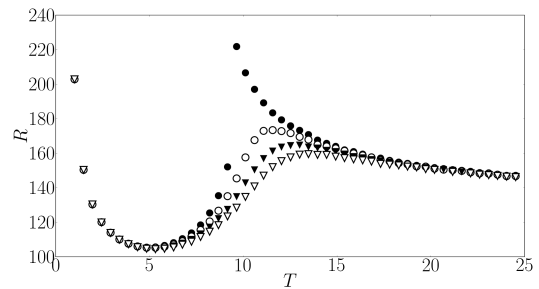


Figure 4: Resistivity R in arbitrary unit versus temperature T , for different magnetic fields B : 0 (black circles), 0.25 (void circles), 0.5 (black triangles), 0.75 (void triangles). $I_0 = 2$ and other parameters taken the same as in Fig. 2.

4. Results on Antiferromagnetic Thin Films

In the case of antiferromagnets, we study first a film with simple cubic (SC) lattice structure. This is because the FCC lattice used in the ferromagnetic case shown above becomes fully frustrated if we use an antiferromagnetic interaction. The frustrated case is very particular [33], it cannot be treated on the same footing as the non frustrated case.

Before showing the results on a SC antiferromagnet, let us emphasize the following point. The picture of defect clusters of down spins embedded in a up-spin sea that we used above to explain the behavior of the resistivity in ferromagnets should be modified in the case of antiferromagnets: in antiferromagnets defects are domain walls, clusters on the two sides of a wall both have antiferromagnetic ordering with opposite parity. An itinerant spin crossing a domain wall does not have the same scattering as in a ferromagnet. Its scattering depends on the numbers of up spins and down spins in the sphere of radius D_1 . In other words, the scattering depends on the energy landscape in the crystal: the itinerant spin will stay a longer time where its energy is low, and a shorter time where its energy is high.

The resistivity versus T in zero magnetic field is shown in Fig. 5 with $D_1 = D_2 = 1$. Several remarks are in order: i) one observes the absence of a peak of R ; ii) the variation of R with T has the same shape as the internal energy versus T shown in Fig. 6, therefore dR/dT shows a peak similar to the specific heat. The peak of dR/dT has been experimentally observed in many materials, in particular in MnSi [3, 4] among others [2].

The absence of a peak at T_N observed here certainly comes from the fact that the motion of an itinerant electron is not sharply slowed down at T_N by numerous clusters of opposite spins. Let us say it again in another manner: the absence of a peak at the transition is due to the fact that the motion of an itinerant spin depends on its immediate environment: in ferromagnets, the variation of its energy ΔE going from a "parallel" cluster to a nearby "defect" (or antiparallel) cluster is much larger than the energy variation going from a cluster of antiferromagnetic

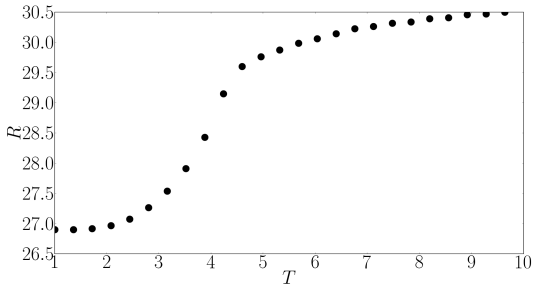


Figure 5: SC AF case. Resistivity R in arbitrary unit versus temperature T , in zero magnetic field, with electric field $E = 1$, $I_0 = K_0 = 0.5$.

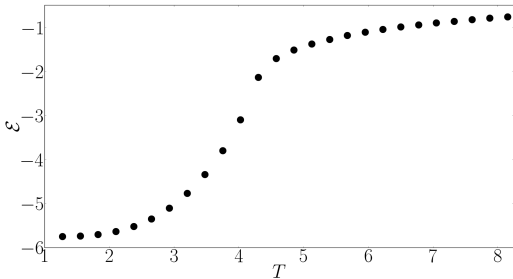


Figure 6: SC AF case. Internal energy \mathcal{E} in unit $|J| = 1$ versus temperature T , in zero magnetic field, with electric field $E = 1$, $I_0 = K_0 = 0.5$.

ordering to a cluster which is a defect but a defect with an antiferromagnetic structure in the SC antiferromagnetic case. The smaller ΔE gives rise to a larger spin mobility i. e. a smaller R . Note that experimental data show just a shoulder in antiferromagnetic LaFeAsO[16].

We show now in Fig. 7 the effect of D_1 on the resistivity at a given temperature. We observe here an oscillatory behavior of R . By analyzing the ratio of numbers of up spins and down spins in the sphere of radius D_1 , we found that this ratio oscillates with varying D_1 : the maxima (minima) of R correspond to the largest (smallest) numbers of parallel (antiparallel) spins in the sphere. This finding is consistent with what we said before, namely R is large when the energy of itinerant spin is low (i. e. large number of parallel spins).

We show in Fig. 8 the resistivity versus T for several D_1 . As seen here, the change of R at a given $T > T_N$ with varying D_1 is much smaller than that for $T < T_N$. It is interesting to note that R does not depend on D_1 at T_N . Further analysis should be carried out to understand this behavior at the transition point.

Finally, to compare with the SC AF case, let us show the results of the BCC and diamond-lattice antiferromagnets in Figs. 9 and 10. As seen, these cases show no peak for any value of D_1 . Note that R in both cases increases as T decreases to zero, unlike the SC AF case. The origin of this increase lies in the freezing of itinerant spins at low T . The degree of freezing depends on the lattice structure and on the strength of the interactions of the itinerant spins

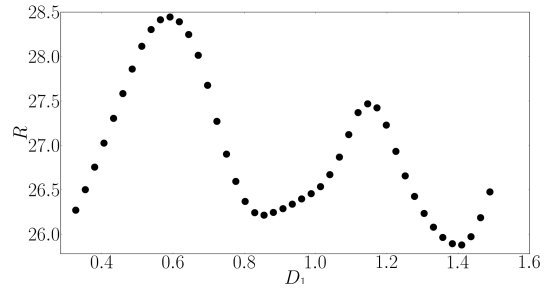


Figure 7: SC AF case. Resistivity R in arbitrary unit versus D_1 at $T = 1$, in zero magnetic field, with electric field $E = 1$, $I_0 = K_0 = 0.5$.

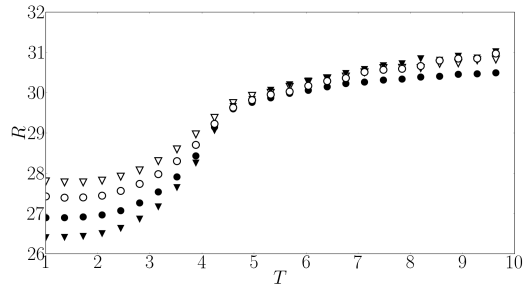


Figure 8: SC AF case. Resistivity R versus T for several values of D_1 : 1 (black circles), 1.2 (void triangles), 1.4 (black triangles), 1.6 (void circles), $E = 1$, $B = 0$, $I_0 = K_0 = 0.5$, $D = 0.35$.

with their environment as discussed in the ferromagnetic section.

4.1. Discussion

In the case of ferromagnets, the coupling of the motion of itinerant spins to the correlation of the lattice spins gives rise to the peak of the resistivity in the transition region. Depending on the strength of this coupling, the peak can be very sharp or rounded at T_c . The picture of scattering by clusters suggested by us is naturally consistent with the correlation interpretation.

In the case of antiferromagnets, the polarized itinerant spins are coupled to both parallel and antiparallel lattice spins. Due to the opposite correlation signs, their respec-

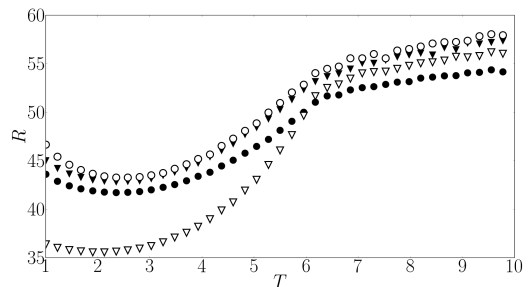


Figure 9: BCC AF case. Resistivity R versus T for several values of D_1 : 1 (black circles), 1.2 (void triangles), 1.4 (black triangles), 1.6 (void circles), $E = 1$, $B = 0$, $I_0 = K_0 = 1$, $D = 0.5$.

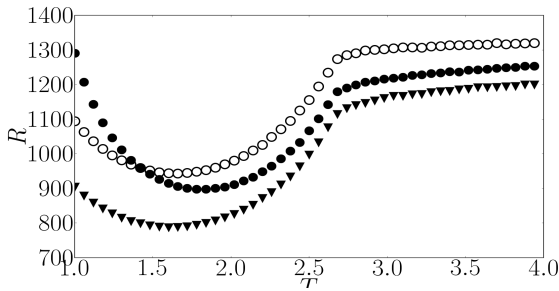


Figure 10: Diamond-lattice AF case. Resistivity R versus T for D_1 : 0.5 (black triangles), 0.75 (black circles), 1 (void circles). $E = 1$, $B = 0$, $I_0 = K_0 = 0.5$, $D = 0.35$.

tive effects are partially canceled out giving rise to an effective coupling weaker than that in ferromagnets. It is therefore not surprising that the peak is absent unlike in the ferromagnetic case.

5. Conclusion

In this work, we have improved our MC simulations by averaging the resistivity over a large number of lattice spin configurations. The results shown above are much better than those in our previous works [30, 31]. Though the physics in the ferromagnetic case are not qualitatively altered but the precision on the peak position of R is excellent and the statistical fluctuations in the paramagnetic phase are reduced. The spin resistivity is strongly dependent on the temperature. In ferromagnets, at very low T the itinerant spins are somewhat frozen. As T increases, their motion is thermally activated making a decrease of R . However, as T increases further, the system enters the transition region, R increases and undergoes a huge peak at the ferromagnetic transition temperature. At higher temperatures, the lattice spins are disordered, the resistivity is still large but it decreases with increasing T . The existence of the peak in ferromagnets is in agreement with theories, in particular those of Zarand et al [23] and Haas [21] where interaction between itinerant spins and lattice spins is dominant.

We have also shown here results of some antiferromagnets. The absence of a peak of R in the SC, BCC and diamond lattices confirms the prediction of Haas [21]. Let us emphasize that our results on frustrated FCC AF [33] and on Heisenberg BCC AF [34] show that the shape of the resistivity in antiferromagnets strongly depends on the spin model and the nature, i. e. first or second order, of the lattice phase transition. The extension of the Boltzmann's theory to the case of antiferromagnets is under way.

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[1] A. Fert and I. A. Campbell, Phys. Rev. Lett. **21**, 1190 (1968); I. A. Campbell, Phys. Rev. Lett. **24**, 269 (1970).

- [2] F. C. Shwerner and L. J. Cuddy, Phys. Rev. **2**, 1575 (1970).
[3] Alla E. Petrova, E. D. Bauer, Vladimir Krasnorussky, and Sergei M. Stishov, Phys. Rev. B **74**, 092401 (2006).
[4] S. M. Stishov, A.E. Petrova, S. Khasanov, G. Kh. Panova, A.A.Shikov, J. C. Lashley, D. Wu, and T. A. Lograsso, Phys. Rev. B **76**, 052405 (2007).
[5] F. Matsukura, H. Ohno, A. Shen and Y. Sugawara, Phys. Rev. B **57**, R2037 (1998).
[6] R. Brucas and M. Hanson; J. Magn. Magn. Mater. **310**, 2521 (2007).
[7] T. Kasuya, Prog. Theor. Phys. Vol. 16,P. 58 (1956), No. 1.
[8] E. A. Turov, Iza. Akad. Nauk. SSSR. Serb. Fiz. **19**, 426 (1955).
[9] Jing Xia, W. Siemons, G. Koster, M. R. Beasley and A. Kapitulnik, Phys. Rev. B **79**, 140407(R) (2009).
[10] X. F. Wang et al, Phys. Rev. Lett. **102**, 117005 (2009).
[11] Tiffany S. Santos, Steven J. May, J. L. Robertson and Anand Bhattacharya, Phys. Rev. B **80**, 155114 (2009).
[12] Y. B. Li, Y. Q. Zhang, N. K. Sun, Q. Zhang, D. Li, J. Li and Z. D. Zhang, Phys. Rev. B **72**, 193308 (2005).
[13] C. L. Lu, X. Chen, S. Dong, K. F. Wang, H. L. Cai, J.-M. Liu, D. Li and Z. D. Zhang, Phys. Rev. B **79**, 245105 (2009).
[14] J. Du, D. Li, Y. B. Li, N. K. Sun, J. Li and Z. D. Zhang, Phys. Rev. B **76**, 094401 (2007).
[15] Y. Q. Zhang, Z. D. Zhang and J. Aarts, Phys. Rev. B **79**, 224422 (2009).
[16] M. A. McGuire, A. D. Christianson, A. S. Sefat, B. C. Sales, M. D. Lumsden, R. Jin, E. A. Payzant, D. Mandrus, Y. Luan, V. Keppens, V. Varadarajan, J. W. Brill, R. P. Hermann, M. T. Sougrati, F. Grandjean and G. J. Long: Phys. Rev. B **78** (2008) 094517.
[17] P.-G. de Gennes and J. Friedel, J. Phys. Chem. Solids **4**, 71 (1958).
[18] Paul P. Craig, Walter I. Goldburg, T. A. Kitchens and J. I. Budnick, Phys. Rev. Lett. **19**, 1334 (1967).
[19] M. E. Fisher and J.S. Langer, Phys. Rev. Lett. **20**, 665 (1968).
[20] S. Alexander, J. S. Helman and I. Balberg, Phys. Rev. B. **13**, 304 (1975).
[21] C. Haas, Phys. Rev **168**, 531 (1968).
[22] Mitsuo Kataoka, Phys. Rev. B **63**, 134435-1 (2001).
[23] G. Zarand, C. P. Moca and B. Janko, Phys. Rev. Lett. **94**, 247202 (2005).
[24] M. N. Baibich, J. M. Broto, A. Fert, F. Nguyen Van Dau, F. Petroff, P. Etienne, G. Creuzet, A. Friederich and J. Chazelas, Phys. Rev. Lett. **61**, 2472 (1988).
[25] P. Grunberg, R. Schreiber, Y. Pang, M. B. Brodsky and H. Sowers, Phys. Rev. Lett. **57**, 2442 (1986); G. Binash, P. grunberg, F. Saurenbach and W. Zinn, Phys. Rev. B **39**, 4828 (1989).
[26] A. Barthél my, A. Fert, J.-P. Contour and M. Bowen, J. Mag. Mag. Mater. **242-245**, 68 (2002).
[27] See review by E. Y. Tsymbal and D. G. Pettifor, *Solid State Physics*, Academic Press (San Diego), Vol. 56, pp. 113-237 (2001).
[28] See review on Semiconductor Spintronics by T. Dietl, in *Lectures Notes*, vol. 712, Springer (Berlin), pp. 1-46 (2007).
[29] See review on Oxide Spintronics by M. Bibes and A. Barth l my, in a Special Issue of IEEE Transactions on Electron Devices on Spintronics, IEEE Trans. Electron. Devices **54**, 1003 (2007).
[30] K. Akabli, H. T. Diep and S. Reynal, J. Phys.: Condens. Matter **19**, 356204 (2007).
[31] K. Akabli and H. T. Diep, Phys. Rev. B **77**, 165433 (2008).
[32] K. Binder and D. W. Heermann, *Monte Carlo Simulation in Statistical Physics*, Springer, New York (2002).
[33] Y. Magnin, K. Akabli and H. T. Diep, in preparation.
[34] K. Akabli, Y. Magnin, Masataka, I. Harada and H. T. Diep, submitted to J. Phys. Soc. Jpn.