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Existence and conservation laws for the Boltzmann-Fermi equation in a general domain

Thibaut Allemand

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Abstract

We give an existence theorem for the Boltzmann-Fermi-Dirac equation for integrable collision cross-sections in general domains with specular reflection at the boundaries, using the characteristic lines of the free transport. We then show that the solution satisfies the local conservations of mass, momentum and kinetic energy, which is the main novelty compared to previous results. This is done thanks to a dispersion technique.

1 Introduction

In this note, we are concerned with the Boltzmann equation for a gas of fermions. Our goal is double: first, proving that the existence result and the properties of the solution proved in [3] can be extended to more general domains than \mathbb{R}^3 , if we supplement the equation with a specular reflection condition at the boundaries; and then, establishing the local conservation laws. These two properties are needed in the study of hydrodynamic limits, especially in the incompressible inviscid regime, which is studied in a forthcoming paper [1].

Let us detail the problem. Let Ω be a subset of \mathbb{R}^3 having bounded curvature, and regular enough such that the function $x \mapsto n(x)$, where $n(x)$ is the outer unit normal at point x of $\partial\Omega$, can be extended into a continuously differentiable function on \mathbb{R}^3 .

The Boltzmann equation for a gas of fermions reads

$$\partial_t f + v \cdot \nabla_x f = Q(f) \quad (1)$$

where $f(t, x, v)$ is the density of particles which at time $t \in \mathbb{R}_+$ are at point $x \in \Omega$ with velocity $v \in \mathbb{R}^3$. It is supplemented with an initial condition

$$f(0, x, v) = f_0(x, v) \quad \forall (x, v) \in \Omega \times \mathbb{R}^3 \quad (2)$$

and a boundary condition

$$f(t, x, v) = f(t, x, R_x(v)) \quad \forall (x, v) \in \partial\Omega \times \mathbb{R}^3 \text{ such that } n(x) \cdot v < 0 \quad (3)$$

where $R_x(v)$ is the specular reflection law

$$R_x(v) = v - 2(v \cdot n(x))n(x).$$

The collision integral $Q(f)$ is given by

$$Q(f) = \int_{S^2} \int_{\mathbb{R}^3} b(v - v_*, \omega) (f' f'_* (1 - f)(1 - f_*) - f f_* (1 - f')(1 - f'_*)) dv_* d\omega \quad (4)$$

with the usual notations

$$f_* = f(t, x, v_*), \quad f' = f(t, x, v'), \quad f'_* = f(t, x, v'_*)$$

and where the precollisional velocities (v', v'_*) are deduced from the postcollisional ones by

$$\begin{cases} v' = v - (v - v_*) \cdot \omega \omega, \\ v'_* = v_* + (v - v_*) \cdot \omega \omega. \end{cases}$$

The collision integral differs from the classical one by the terms $(1 - f)$, which take into account the Pauli exclusion principle.

The function $b(w, \omega)$, known as the collision cross-section, is measurable, a.e. positive, and is assumed to be in L^1 as in [3]. For the study of the conservations, we will moreover assume that

$$b(w, \omega) = q(|w|, |w \cdot \omega|),$$

which is physically relevant, and endows (formally) $Q(f)$ with symmetry properties:

$$\int_{\mathbb{R}^3} (1, v, |v|^2) Q(f) dv = 0$$

for all f such that the integrals make sense. Then, integrating equation (1) against $1, v, |v|^2$, we obtain that the solution f satisfies formally the following conservation laws:

$$\begin{cases} \partial_t \int_{\mathbb{R}^3} f dv + \nabla_x \cdot \left(\int_{\mathbb{R}^3} v f dv \right) = 0 \\ \partial_t \int_{\mathbb{R}^3} v f dv + \nabla_x \cdot \left(\int_{\mathbb{R}^3} v \otimes v f dv \right) = 0 \\ \partial_t \int_{\mathbb{R}^3} |v|^2 f dv + \nabla_x \cdot \left(\int_{\mathbb{R}^3} v |v|^2 f dv \right) = 0 \end{cases} \quad (5)$$

This note aims at showing first an existence theorem for the initial-boundary value (IBV) problem (1)-(3), and then that the local conservation laws (5) are satisfied rigorously by the solution of the problem.

2 The IBV problem

Following [5] and [4], we define the characteristic lines of the free transport equation

$$\partial_t f + v \cdot \nabla_x f = 0$$

in the following way: for $(x, v) \in \Omega \times \mathbb{R}^3$, the characteristic line is given by $x + vt$, at least if t is small. Then, let $t_0(x, v)$ be the first value of t for which $x + t_0 v \in \partial\Omega$. Then, for $t > t_0$, the trajectory continues as $x + t_0 v + (t - t_0) R_{x + t_0 v}(v)$. If it intersects one more time the boundary of Ω at time $t_1(x, v)$, then the trajectory continues with velocity $R_{x + t_0 v + (t_1 - t_0) R_{x + t_0 v}(v)}(R_{x + t_0 v}(v))$, and so on. We then

define a family of maps $\{\Psi^t\}$ called the trajectory maps, $\Psi^t(x_0, v_0)$ being the point in phase space at which we arrive at time t following the trajectory line issuing from (x_0, v_0) . For $t = t_0, t_1, \dots$, we define Ψ^t to be continuous (in time) from the right. Since Ω has bounded curvature, the trajectory intersects the boundary finitely many times in finite times. Moreover, the energy is conserved at each reflection, Ψ^t maps $\Omega \times \mathbb{R}^3$ onto itself for every t , and the jacobian of Ψ^t is always unity [5].

With this construction, we now define

$$f^\#(t, x, v) = f(t, \Psi^t(x, v)),$$

that is, we conjugate f with the free transport semigroup. This notation allows to reformulate the boundary condition in a simpler way. Indeed, notice that if $\Psi^{t-0}(x_0, v_0) = (x, v)$, then $\Psi^t(x_0, v_0) = \Psi^{t+0}(x_0, v_0) = (x, R(v))$, so that the boundary condition can be written

$$f^\#(t-0, x_0, v_0) = f^\#(t+0, x_0, v_0),$$

that is, $t \mapsto f^\#(t, x, v)$ is continuous.

The existence result for the IBV problem (1)-(3) is a consequence of the following lemma:

Lemma 2.1. *Let $f, h \in L^1_{loc}(\mathbb{R}_+ \times \Omega \times \mathbb{R}^3)$. Then f is a solution of*

$$\begin{cases} \partial_t f + v \cdot \nabla_x f = h & \text{in } \mathcal{D}'(\mathbb{R}_+ \times \Omega \times \mathbb{R}^3) \\ f(t, x, v) = f(t, x, R_x(v)) & \forall (x, v) \in \partial\Omega \times \mathbb{R}^3 \text{ s.t. } n(x) \cdot v < 0 \end{cases} \quad (6)$$

if and only if for almost all $(x, v) \in \Omega \times \mathbb{R}^3$, $f^\#$ is absolutely continuous with respect to t , $h^\#(t, x, v) \in L^1_{loc}(\mathbb{R}_+)$ and

$$f^\#(t, x, v) = f_0^\#(t, x, v) + \int_0^t h^\#(s, x, v) ds.$$

Proof. It is very similar to what is done in [2] [4]. It consists in multiplying (6) by $\chi(\Psi^t(x, v))\phi(t)$ and integrating, with $\chi \in \mathcal{D}(\Omega \times \mathbb{R}^3)$ and $\phi \in \mathcal{D}(\mathbb{R}_+)$. Using then the change of variables $(X, V) = \Psi^{-t}(x, v)$ leads to

$$\int_{\Omega \times \mathbb{R}^3} \chi(x, v) \left(\int_{\mathbb{R}_+} [f^\# \phi'(t) + h^\# \phi(t)] dt \right) dx dv = 0$$

and then to the announced result, since this is true for every $\chi \in \mathcal{D}(\Omega \times \mathbb{R}^3)$ and $\phi \in \mathcal{D}(\mathbb{R}_+)$. \square

Equipped with this lemma, it is now possible to apply all the strategy developed in [3], and we obtain the following two theorems:

Theorem 2.2. *Let Ω be either \mathbb{R}^3 or a regular subset of \mathbb{R}^3 . Let the collision cross-section be such that*

$$0 \leq b \in L^1(\mathbb{R}^3 \times S^2), \quad (7)$$

and let

$$f_0 \in L^\infty(\Omega \times \mathbb{R}^3), \quad 0 \leq f_0 \leq 1. \quad (8)$$

Then, the problem (1)-(3) has a unique solution f satisfying

$$f \in L^\infty(\mathbb{R}_+ \times \Omega \times \mathbb{R}^3), \quad 0 \leq f \leq 1 \quad a.e..$$

Moreover, f is absolutely continuous with respect to t .

Sketch of proof. The main idea is to show that the function

$$T : f \mapsto f_0(\Psi^{-t}(x, v)) + \int_0^t Q(\bar{f})(s, \Psi^{s-t}(x, v)) ds$$

has a fixed point, with

$$\bar{f} = \begin{cases} 0 & \text{if } f \leq 0 \\ f & \text{if } 0 \leq f \leq 1 \\ 1 & \text{if } f \geq 1 \end{cases}$$

This is achieved by showing that T is contractive in $L^\infty([0, \theta] \times \Omega \times \mathbb{R}^3)$ if θ is small enough. Then, since

$$-B \max(f, 0) \leq -B\bar{f} \leq Q(\bar{f}) \leq B(1 - \bar{f}) \leq B(1 - \min(1, f))$$

where $B = \|b\|_{L^1(\mathbb{R}^3 \times S^2)}$, it comes

$$-B \max(f^\sharp, 0) \leq \partial_t f^\sharp \leq B(1 - \min(1, f^\sharp)),$$

which ensures that $f = \bar{f}$. It is then possible to reiterate the process on $[\theta, 2\theta]$, and so on, to construct a global solution. \square

Under the (physically relevant) assumption that the cross-section writes

$$b(w, \omega) = q(|w|, |w \cdot \omega|), \quad (9)$$

the following macroscopic conservation laws hold:

Theorem 2.3. *Let f_0 satisfy (8), and b satisfy (9) and (7). Then*

- if $f_0 \in L^1(\Omega \times \mathbb{R}^3)$, then the solution f to (1) belongs to $C^0(\mathbb{R}_+; L^1(\Omega \times \mathbb{R}^3))$ and

$$\iint_{\Omega \times \mathbb{R}^3} f(t, x, v) dx dv = \iint_{\Omega \times \mathbb{R}^3} f_0(x, v) dx dv, \quad \forall t \in \mathbb{R}_+;$$

- if $\iint_{\Omega \times \mathbb{R}^3} |v|^2 f_0 dx dv < +\infty$, then the solution f to (1) is such that the function $(t, x, v) \mapsto |v|^2 f(t, x, v)$ belongs to $C^0(\mathbb{R}_+; L^1(\Omega \times \mathbb{R}^3))$ and

$$\iint_{\Omega \times \mathbb{R}^3} |v|^2 f(t, x, v) dx dv = \iint_{\Omega \times \mathbb{R}^3} |v|^2 f_0(x, v) dx dv, \quad \forall t \in \mathbb{R}_+.$$

Moreover, the function $(t, x, v) \mapsto |v|^2 Q(f)(t, x, v)$ belongs to $L^\infty(\mathbb{R}_+; L^1(\Omega \times \mathbb{R}^3))$.

3 Local conservation laws

We first recall a theorem of [6]:

Theorem 3.1. *Let $f \in L^1([0, T]; L^1(\Omega \times \mathbb{R}^3))$ be solution of*

$$\partial_t f + v \cdot \nabla_x f = g, \quad f(t=0) = f_0$$

with

$$\iint_{\Omega \times \mathbb{R}^3} |v|^3 f_0(x, v) dx dv < +\infty$$

and

$$\int_0^T \iint_{\Omega \times \mathbb{R}^3} |v|^2 g(t, x, v) dx dv dt < C_0,$$

then, for any bounded subset K of Ω , we have

$$\int_0^T dt \int_K dx \int_{\mathbb{R}^3} |v|^3 f(t, x, v) dv \leq C_K.$$

In view of all the preceding results, we can state our theorem about the local conservation laws:

Theorem 3.2. *Let the cross-section b satisfy (9) and (7). Assume that the initial data satisfy (8), and*

$$\iint_{\Omega \times \mathbb{R}^3} (1 + |v|^3) f_0(x, v) dx dv < +\infty.$$

Then, the solution f to (1)-(3) satisfies, in distributional sense, the local conservation laws

$$\begin{cases} \partial_t \int_{\mathbb{R}^3} f dv + \nabla_x \cdot \left(\int_{\mathbb{R}^3} v f dv \right) = 0 \\ \partial_t \int_{\mathbb{R}^3} v f dv + \nabla_x \cdot \left(\int_{\mathbb{R}^3} v \otimes v f dv \right) = 0 \\ \partial_t \int_{\mathbb{R}^3} |v|^2 f dv + \nabla_x \cdot \left(\int_{\mathbb{R}^3} v |v|^2 f dv \right) = 0 \end{cases} \quad (10)$$

Proof. In the sense of distributions, f satisfies (1)

$$\partial_t f + v \cdot \nabla_x f = Q(f).$$

Let $\phi \in \mathcal{D}(\mathbb{R}_+ \times \Omega)$ and $\Psi_R \in \mathcal{D}(\mathbb{R}^3)$ satisfy:

$$\Psi_R(v) = \begin{cases} 1 & \text{if } |v| \leq R \\ 0 & \text{if } |v| > R. \end{cases}$$

Then equation (1) implies

$$\begin{aligned} & \int_{\mathbb{R}_+ \times \Omega} \partial_t \phi \left[\int_{\mathbb{R}^3} \Psi_R(v) \begin{pmatrix} 1 \\ v \\ |v|^2 \end{pmatrix} f dv \right] dx dt + \int_{\mathbb{R}_+ \times \Omega} \nabla_x \phi \left[\int_{\mathbb{R}^3} \Psi_R(v) \begin{pmatrix} v \\ v \otimes v \\ v |v|^2 \end{pmatrix} f dv \right] dx dt \\ &= - \int_{\mathbb{R}_+ \times \Omega} \phi \left[\int_{\mathbb{R}^3} \Psi_R(v) \begin{pmatrix} 1 \\ v \\ |v|^2 \end{pmatrix} Q(f) dv \right] dx dt \end{aligned}$$

We may pass to the limit $R \rightarrow +\infty$ thanks to 2.3 and theorem 3.1. At the end, we get the desired result. \square

References

- [1] T. Allemand, *The incompressible Euler limit of the Boltzmann equation for a gas of fermions*, in preparation.
- [2] R.J. DiPerna and P.-L. Lions, *On the Cauchy problem for Boltzmann equations: global existence and weak stability*, *Ann. of Math. (2)*, **130** (1989), 321–366.
- [3] J. Dolbeault, *Kinetic models and quantum effects: a modified Boltzmann equation for Fermi-Dirac particles*, *Arch. Rational Mech. Anal.*, **127** (1994), 101–131.
- [4] K. Hamdache, *Initial-boundary value problems for the Boltzmann equation: global existence of weak solutions* *Arch. Rational Mech. Anal.*, **119** (1992), 309–353.
- [5] S. Kaniel and M. Shinbrot, *The Boltzmann equation. I. Uniqueness and local existence*, *Comm. Math. Phys.*, **58** (1978), 65–84.
- [6] B. Perthame, *Global existence to the BGK model of Boltzmann equation*, *Journal of Differential equations*, **82** (1989), 191–205.