



On the time schedule of Brownian Flights

Athanasios Batakis, Michel Zinsmeister

► To cite this version:

Athanasios Batakis, Michel Zinsmeister. On the time schedule of Brownian Flights. 2009. hal-00398627v3

HAL Id: hal-00398627

<https://hal.science/hal-00398627v3>

Preprint submitted on 8 Mar 2010

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

On the time schedule of Brownian Flights

Athanasios BATAKIS and Michel ZINSMEISTER

March 8, 2010

Abstract: *We are interested in the statistics of the duration of Brownian diffusions started at distance ϵ from the boundary of a given domain and stopped when they hit back this boundary.*

1 Introduction

The motivation of the following work has its origin in experimental physics. Some long molecules are solvable in a liquid (for instance imogolite in water or DNA in lithium) and the molecules forming the liquid show an intermittent dynamics, alternating diffusion in the bulb and adsorption on the long molecules. For the physicist's point of view, it is very important to have as precise as possible knowledge of the statistics of these brownian flights.

In [GKL⁺06] a connection is established between the statistics of the long flight lengths and the geometry of the long molecules (more precisely their Minkowski dimension). This connection has been made rigorous in [BLZ09]. These two papers concern almost exclusively lengths. How does one check experimentally the results? A very powerful tool for that is relaxation methods in nuclear magnetic resonance (see [DPP⁺08]): but this method only allows to compute (the statistics of the) duration of long flights. Some heuristic link between time and length was derived in [GKL⁺06], [DPP⁺08]. The aim of this paper is to make this heuristics rigorous.

2 Geometric Background

In the sequel, Ω will always denote a domain in \mathbb{R}^d with compact boundary. The crucial tool we need to use is the notion of Whitney cubes. We thus recall the

Proposition 2.1 (*cf. [Gra08], p. 463*) *Given any non-empty open proper subset Ω of \mathbb{R}^d , there exists a family of closed dyadic cubes $\{Q_j\}_j$ such that*

- $\bigcup_j Q_j = \Omega$ and the cubes Q_j 's have disjoint interiors
- $\sqrt{d}\ell(Q_j) \leq \text{dist}(Q_j, \partial\Omega) \leq 4\sqrt{d}\ell(Q_j)$

- if Q_j and Q_k touch then $\ell(Q_j) \leq 4\ell(Q_k)$
- for a given Whitney cube Q_j there are at most 12^d Whitney cubes Q_k 's that touch Q_j .

In this statement, $\ell(Q)$ stands for the side-length of the cube Q and, for $\lambda > 0$, λQ is the cube of the same center and of sidelength $\lambda\ell(Q)$. For $k \in \mathbb{Z}$, we denote by \mathcal{Q}_k , the collection of Whitney cubes Q_j with $\ell(Q_j) = 2^k$. We also recall the definition of the Minkowski sausage: for $r > 0$,

$$M_r = \{x \in \Omega ; \text{dist}(x, \partial\Omega) \leq r\}$$

and

$$\Gamma_r = \{x \in \Omega ; \text{dist}(x, \partial\Omega) = r\}$$

We then define \mathcal{S}_r as the collection of Whitney cubes intersecting Γ_r . Notice that \mathcal{S}_r is a finite set.

Definition 2.1 *Let $\varepsilon > 0$. We will call Brownian flight the random process $F_t, t \geq 0$ consisting in picking at random with equiprobability one of the dyadic Whitney cubes of \mathcal{S}_ε and starting from the center of the cube a Brownian motion B_t killed once it reaches $\partial\Omega$. We denote by $\tau_\Omega = \inf\{t ; F_t \notin \Omega\}$ the lifetime of this process.*

We are interested in the asymptotics of $\mathbb{P}(\tau_\Omega > t)$ as t grows, but this needs some explanation:

It is well known that, if Ω is bounded, this quantity decreases exponentially, as $t \rightarrow \infty$, as $e^{\lambda t}$, where λ is the first eigenvalue of the Laplacian. We define

$$R_\Omega = \min \left\{ 1, \sup_{x \in \Omega} \text{dist}(x, \partial\Omega) \right\}.$$

Our aim is to evaluate $\mathbb{P}(\tau_\Omega > t)$ in the interval $\varepsilon^2 \leq t \leq R_\Omega^2$, and this independently of ε . In fact, the estimate we are looking for is thus an estimate with respect to ε rather than for "pure" t .

We study first a simple example, since we will need the partial result anyhow. Let $[0, a]$ be a real segment and take $x \in (0, a)$. The probability that brownian motion started at x has not exit the interval $(0, a)$ by time t , $\mathbb{P}(\tau_x > t)$, is given by the following equivalent formulas (see [Fel71], pg. 342)

$$\mathbb{P}(\tau_x > t) = \frac{1}{\sqrt{2\pi}} \int_{-x/\sqrt{t}}^{-x/\sqrt{t}} \exp\left(-\frac{1}{2}y^2\right) dy + \sqrt{\frac{2}{\pi}} \sum_{k=1}^{\infty} (-1)^k \int_{(ka-x)\sqrt{t}}^{(ka+x)\sqrt{t}} \exp\left(-\frac{1}{2}y^2\right) dy$$

and

$$\mathbb{P}(\tau_x > t) = \frac{4}{\pi} \sum_{k=1}^{\infty} \exp\left(-\frac{(2n+1)^2\pi^2}{2a^2}t\right) \sin \frac{(2n+1)\pi x}{a}.$$

By symmetry we can assume $x \leq \frac{a}{2}$. If $\frac{a}{\sqrt{t}}$ is not too big, say $\frac{a}{\sqrt{t}} < 1/2$, we have an easy estimate of $\mathbb{P}(\tau_x > t)$ using the first formula : $\mathbb{P}(\tau_x > t) \sim \frac{x}{\sqrt{t}}$. On the other hand, for $\frac{a}{\sqrt{t}}$ not

too small we get, by the second formula, that $\mathbb{P}(\tau_x > t) \leq \exp\left(-\frac{\pi^2}{2a^2}t\right)$. Hence, that there is a change of regime, the decay of $\mathbb{P}(\tau_x > t)$ with t going from polynomial to exponential and the quantity $\frac{x}{\sqrt{t}}$ is relevant for small times.

In a higher dimensional context, let Q be the cube centered at 0 and of side r . By the preceding remark it follows that, T denoting the exit time from Q of the Brownian motion starting at 0, we have

$$\mathbb{P}(T > t) \leq c \left(\frac{r}{\sqrt{t}} \right)^d \quad (1)$$

where c depends only on d . Here, $\varepsilon = r \simeq R_\Omega$ and the opposite inequality thus holds for $\varepsilon^2 \leq t \leq R_\Omega^2$.

If the Brownian motion is started at distance ε from the boundary then the exit time is essentially the exit time from a half space and we thus get

$$\mathbb{P}(T > t) \leq c \left(\frac{\varepsilon}{\sqrt{t}} \right) = c \left(\frac{\varepsilon}{\sqrt{t}} \right)^{d-1+2-d}, \quad (2)$$

and, as we will see, the opposite inequality is valid for t not too big.

Our goal is to extend (2) to general domains with rough boundary. In order to describe the domain of validity of our result let us recall a few definitions. Let K be a compact subset of \mathbb{R}^d . For $j \geq 0$ let N_j be the number of dyadic cubes of the j -th generation (i.e of size 2^{-j}) that intersect K .

Definition 2.2 *The Minkowski dimension of K is*

$$d_M(K) = \limsup_{j \rightarrow \infty} \frac{\log_2(N_j)}{j}$$

Returning to our situation, we can define similarly the Whitney dimension of $\partial\Omega$ as

$$d_W = d_W(\partial\Omega) = \limsup_{j \rightarrow \infty} \frac{\log_2(W_j)}{j}, \quad (3)$$

where W_j is the number of elements of \mathcal{Q}_j .

Under very mild conditions (see [Bis96], [JK82], [BLZ09]) these two dimensions coincide. If the boundary of Ω has some self similarity we can moreover say that there is a constant $c > 0$ such that

$$\frac{1}{c} \varepsilon^{d_M} \leq \#\mathcal{S}_\varepsilon \leq c \varepsilon^{d_M}, \quad (4)$$

for all $\varepsilon \leq R_\Omega$, where $d_M = d_M(\partial\Omega)$.

In our main theorem we will assume that our domain Ω satisfies (4).

We also suppose that the domain Ω satisfies so-called Δ -regularity condition (see also [JW88], [Anc86], [HKM93]): there exists $L > 0$ such that for all $x \in \Omega$, if $d_x = \text{dist}(x, \partial\Omega) < R_\Omega$ then

$$\omega_{\mathbb{B}(x, 2d_x) \cap \Omega}^x(\partial\Omega) \geq L, \quad (5)$$

where $\omega_{\mathbb{B}(x, 2d_x) \cap \Omega}^x$ is the distribution law of the hitting point of Brownian motion starting at x and killed when reaching the boundary of $\mathbb{B}(x, 2d_x) \cap \Omega$. This is a very mild condition (satisfied, for instance, by all domains in \mathbb{R}^2 with non-trivial connected boundary) that appears frequently in related literature in various forms (for instance “uniform capacity condition” or Hardy inequality).

3 Time and length estimates for Brownian flights

We can now state the main result. Let Ω be a bounded domain in \mathbb{R}^d satisfying (4) and (5). If τ_Ω denotes the life-time of a Brownian flight F_t with parameter ε we have

Theorem 3.1 *There exists $c > 0$ depending only on constants in (4), (5) (and in particular not on ε) such that*

$$\frac{1}{c} \left(\frac{\varepsilon}{\sqrt{t}} \right)^{d_M+2-d} \leq \mathbb{P}(\tau_\Omega > t) \quad (6)$$

and

$$\mathbb{P}(\tau_\Omega > t) \leq c \left(\frac{\varepsilon}{\sqrt{t}} \right)^{d_M+2-d} \left| \log \left(\frac{\varepsilon}{\sqrt{t}} \right) \right|^{2d}, \quad (7)$$

for all $\varepsilon^2 < t < R_\Omega^2$.

This theorem has a “cousin” theorem, which was proved in [BLZ09].

Theorem 3.2 *Let $\varepsilon < r < R_\Omega$. The probability that the hitting point of F is at distance greater than r from the starting point x is comparable to*

$$\left(\frac{\#\mathcal{S}_r}{\#\mathcal{S}_\varepsilon} \right)^{d_M} \left(\frac{r}{\varepsilon} \right)^{d-2} \quad (8)$$

Notice that we do not assume (4) for this theorem. If we do, we have

$$\left(\frac{\#\mathcal{S}_r}{\#\mathcal{S}_\varepsilon} \right)^{d_M} \left(\frac{r}{\varepsilon} \right)^{d-2} \sim \left(\frac{r}{\varepsilon} \right)^{d_M-(d-2)} \quad (9)$$

Notice that the quantity on the left of (9) is the same as the one in (6) where we have replaced r by \sqrt{t} which is coherent with standard behaviour of Brownian motion.

4 Proof of the theorem 3.1

For $s > 0$ we denote by β_s the total time spent by Brownian flight F_t in the Minkowski sausage $\{x \in \Omega ; \text{dist}(x, \partial\Omega) \leq s\}$ and $\delta\beta_s = \beta_s - \beta_{s/2}$.

We define analogue quantities more adapted to the Whitney decomposition ; namely, $\delta\tilde{\beta}_{2^k}$ will denote the time spent by F_t inside $\tilde{M}_k = \bigcup\{Q ; Q \in \mathcal{Q}_k\}$.

If Q is a Whitney cube we define the "vicinity" of Q as

$$\tilde{Q} = Q \cup \bigcup\{Q' ; Q' \in \mathcal{E}\},$$

where

$$\mathcal{E} = \{Q' \in \bigcup_k \mathcal{Q}_k ; \lambda Q \cap Q' \neq \emptyset \text{ and } \lambda Q' \cap Q \neq \emptyset\},$$

and $\lambda = 8\sqrt{d}$ satisfy that for all Whitney cubes

$$\lambda Q \supset \mathbb{B}(x_Q, 2 \text{dist}(x_Q, \partial\Omega)),$$

x_Q being the center of Q .

We may now start the proof and we begin with

4.1 The upper bound

We separate the event $\{\tau_\Omega > t\}$ by the partition $\{\tau_\Omega = \beta_{\sqrt{t}}\}$ and $\{\tau_\Omega > \beta_{\sqrt{t}}\}$, that is whether the process goes or does not go at distance \sqrt{t} from the boundary. From theorem 3.2 we get

$$\mathbb{P}(\tau_\Omega > t \text{ and } \tau_\Omega > \beta_{\sqrt{t}}) \leq \mathbb{P}(\tau_\Omega > \beta_{\sqrt{t}}) \leq c \left(\frac{\varepsilon}{\sqrt{t}} \right)^{d_M+2-d}.$$

In order to estimate the term $\mathbb{P}(\tau_\Omega = \beta_{\sqrt{t}} > t)$ we begin by writing

$$\tau_\Omega = \sum_{k=-\infty}^{\log_2 \sqrt{t}} \delta\beta_{2^k},$$

and thus, putting $k_0 = \log_2 \epsilon$ and $k_1 = \log_2 \sqrt{t}$,

$$\mathbb{P} \left(\sum_{k=-\infty}^{\log_2 \sqrt{t}} \delta\beta_{2^k} > t \right) \leq \sum_{k=k_0+1}^{k_1} \mathbb{P} \left(\delta\beta_{2^k} > \frac{t}{2(k_1 - k_0)} \right) + \mathbb{P} \left(\beta_{2^{k_0}} > \frac{t}{2} \right) \quad (10)$$

We now invoke the following lemma whose proof is postponed to the next section.

Lemma 4.1 *There exists a number k^* depending only on d and constants C , $0 < p < 1$ depending only on d and L such that for all $t > 0$, $k \in \mathbb{Z}$ and $N \in \mathbb{N}$ we have*

$$\mathbb{P}(\delta\beta_{2^k} > t) \leq Cp^N +$$

$$\mathbb{P}\left(\exists Q \in \mathcal{Q}_k \cup \dots \cup \mathcal{Q}_{k-k^*}; \exists 0 < s_1 < s_2 < t \text{ with } F_{[s_1, s_2]} \subset \tilde{Q} \text{ and } s_2 - s_1 > t/N\right)$$

Using this lemma, we get

$$\mathbb{P}\left(\delta\beta_{2^j} > \frac{t}{2(k_1 - k_0)}\right) \leq cp^N +$$

$$\mathbb{P}\left(\exists Q \in \mathcal{Q}_j \cup \dots \cup \mathcal{Q}_{j-k^*} \exists s_1 < s_2 < \tau_\Omega; F_{[s_1, s_2]} \subset \tilde{Q}, s_2 - s_1 > \frac{t}{2(k_1 - k_0)N}\right)$$

By (1), (9) and the strong Markov property of Brownian motion we then have :

$$\mathbb{P}\left(\delta\beta_{2^j} > \frac{t}{2(k_1 - k_0)}\right) \leq cp^N +$$

$$\mathbb{P}\left(\exists Q \in \mathcal{Q}_j \cup \dots \cup \mathcal{Q}_{j-k^*} \exists s_1 \tau_\Omega; F_{s_1} \in \tilde{Q}\right) \times$$

$$\mathbb{P}\left(\exists s_2 > s_1, F_{[s_1, s_2]} \subset \tilde{Q} \text{ and } s_2 - s_1 > \frac{t}{2(k_1 - k_0)N}\right) \leq$$

$$cp^N + c \left(\frac{2^j}{\sqrt{\frac{t}{2(k_1 - k_0)N}}}\right)^d \left(\frac{\varepsilon}{2^j}\right)^{d_M + 2 - d}$$

Suming up the first term in (10) we get

$$\sum_{k=k_0+1}^{k_1} \mathbb{P}\left(\delta\beta_{2^k} > \frac{t}{2(k_1 - k_0)}\right) \leq c(k_1 - k_0)p^N + c \sum_{k=k_0+1}^{k_1} \left(\frac{2^k}{\sqrt{\frac{t}{2(k_1 - k_0)N}}}\right)^d \left(\frac{\varepsilon}{2^k}\right)^{d_M + 2 - d}$$

$$\leq c(k_1 - k_0)p^N + c(2(k_1 - k_0)N)^d \varepsilon^{d_M + d - 2} \left(\frac{1}{\sqrt{t}}\right)^d \sum_{k=k_0+1}^{k_1} 2^{2d - d_M + 2}$$

$$\leq c(k_1 - k_0)p^N + c(2(k_1 - k_0)N)^d \varepsilon^{d_M + d - 2} \left(\frac{1}{\sqrt{t}}\right)^d 2^{k_1(2d - d_M + 2)}$$

$$\leq cp^N \log_2\left(\frac{\sqrt{t}}{\varepsilon}\right) + c \left(\log_2\left(\frac{\sqrt{t}}{\varepsilon}\right) N\right)^d \left(\frac{\varepsilon}{\sqrt{t}}\right)^{d_M + 2 - d}$$

Take $N \simeq (d_M + 2 - d) \log_p \left(\frac{\varepsilon}{\sqrt{t}}\right)$ to obtain

$$\sum_{k=k_0+1}^{k_1} \mathbb{P}\left(\delta\beta_{2^k} > \frac{t}{2(k_1 - k_0)}\right) \leq c \left(\log\left(\frac{\varepsilon}{\sqrt{t}}\right)\right)^{2d} \left(\frac{\varepsilon}{\sqrt{t}}\right)^{d_M + 2 - d} \quad (11)$$

To bound the second term $\mathbb{P}(\beta_{2^{k_0}} > \frac{t}{2})$ of the sum (10) we need a lemma of the same nature as lemma 4.1. The proof of this lemma is also post poned to the next section.

Lemma 4.2 *Let \mathcal{R}_{k_0} be the collection of all dyadic cubes of sidelength 2^{k_0} intersecting $\partial\Omega$. There exist constants C , $0 < p < 1$ depending only on d and L such that for all $t > 0$, $k_0 \in \mathbb{Z}$ and $N \in \mathbb{N}$*

$$\begin{aligned} \mathbb{P}(\delta\beta_{2^{k_0}} > t) &\leq Cp^N + \\ \mathbb{P}(\exists Q \in \mathcal{R}_{k_0}; \exists 0 < s_1 < s_2 < t \text{ with } F_{[s_1, s_2]} \subset NQ \text{ and } s_2 - s_1 > t/N) \end{aligned}$$

We therefore deduce

$$\mathbb{P}(\delta\beta_{2^{k_0}} > \frac{t}{2}) \leq cp^N + \left(\frac{2^{k_0}N}{\sqrt{t}} \right)^d.$$

We minimize on $N \simeq \log_p(\varepsilon)$ to get

$$\mathbb{P}(\delta\beta_{2^{k_0}} > \frac{t}{2}) \leq c \left(\log(\varepsilon) \frac{\varepsilon}{\sqrt{t}} \right)^d.$$

Combining this last inequality and (11) we get the upper bound, since $d_M + d - 2 \leq d$.

4.2 The Lower Bound

Following the same reasoning for $k_1 = \lceil \log_2 \sqrt{t} \rceil + 1$ we get

$$\mathbb{P}(\tau_\Omega > t) \geq \mathbb{P}(\exists s_1 > 0 \text{ s.t. } F_{s_1} \in \bigcup_{Q \in \mathcal{Q}_{k_1}} Q \text{ and } \exists s_2 > s_1 + t \text{ s.t. } F_{[s_1, s_2]} \in \bigcup_{Q \in \mathcal{Q}_{k_1}} 2Q)$$

Using strong Markov property the this probability can be written as the product of $\mathbb{P}(\exists s_1 > 0 \text{ s.t. } F_{s_1} \in \mathcal{Q}_{k_1})$ with $\mathbb{P}(\exists s_2 > s_1 + t \text{ s.t. } F_{[s_1, s_2]} \in \bigcup_{Q \in \mathcal{Q}_{k_1}} 2Q)$. The second term of the product is greater than the probability that Brownian motion exits a cube of size $2^{k_1+1} \simeq \sqrt{t}$ at time greater than t which is bounded below by a positive constant depending only on d . The first one is simply the probability that Brownian flight gets to \mathcal{Q}_{k_1} which is equivalent to $\left(\frac{\varepsilon}{\sqrt{t}} \right)^{d_M+2-d}$ and the proof is complete.

5 Proofs of lemmas

Let un first deal with lemma 4.1. The proof of 4.2 is quite similar and will hence be abridged.

5.1 Proof of lemma 4.1

We need the following

Lemma 5.1 *Under the Δ -regularity hypothesis the probability that BM touches more than N Whitney cubes of a given size decreases as Cp^N , with $0 < p < 1$, C a positive constant.*

The proof of the lemma relies on an annuli reasoning.

Proof Let $(B_t)_{t>0}$ be Brownian motion started at any point $x \in \Omega$ and choose $k \in \mathbb{Z}$. Choose any $Q \in \mathcal{Q}_k$ and let λQ be the cube of the same center but λ times the side-length $\ell(Q)$ of Q . By the definition of Whitney cubes, there is a $\lambda = 8\sqrt{d}$ depending only on d such that

$$\frac{\lambda}{2}\ell(Q) \leq \text{dist}(Q, \partial\Omega) \leq 2\lambda\ell(Q).$$

Suppose that there exists $t_0 > 0$ such that $B_{t_0} \in Q$. By the Δ -regularity condition (5), the probability that there exists $t_1 > t_0$ with $B_{[t_0, t_1]} \subset \Omega$ and $B_{t_1} \notin \lambda Q$ is bounded above by $p < 1$ depending only on L, λ :

$$\mathbb{P}\left(\exists t_1 > t_0 ; B_{[t_0, t_1]} \subset \Omega \text{ and } B_{t_1} \notin \lambda Q \mid \exists t_0 > 0 ; B_{t_0} \in Q\right) < p \quad (12)$$

On the other hand, the number of Whitney cubes of \mathcal{Q}_k lying inside λQ is bounded by a constant $c_1 = c_1(d)$. The probability that there exists a Whitney cube $Q_1 \in \mathcal{Q}_k$ outside λQ that is visited by Brownian motion is hence bounded above by $p < 1$.

We study probability that there exist Whitney cubes $Q_1, \dots, Q_m \in \mathcal{Q}_k$ such that $Q_1 \cap \lambda Q = Q_2 \cap \lambda Q_1 = \dots = Q_m \cap \lambda Q_{m-1} = \emptyset$ all visited by Brownian motion. It is sufficient to prove that this probability decays exponentially with m .

By the strong Markov property the probability that there exists $t_m > t_{m-1} > \dots > t_0$ such that $B_{t_0} \in Q$, $B_{t_1} \in Q_1$..., $B_{t_m} \in Q_m$ is given by

$$\begin{aligned} & \mathbb{P}(\exists t_m > t_{m-1} > \dots > t_0 ; \text{ and } Q, \dots, Q_m \text{ as above such that } B_{t_m} \in Q_m, \dots, B_{t_0} \in Q) \\ &= \mathbb{P}(\exists t_m > t_{m-1} ; B_{t_m} \in Q_m \mid \exists t_{m-1} > \dots > t_0 B_{t_{m-1}} \in Q_{m-1}, \dots, B_{t_0} \in Q) \cdot \\ & \mathbb{P}(\exists t_{m-1} > \dots > t_0 B_{t_{m-1}} \in Q_{m-1}, \dots, B_{t_0} \in Q) \\ &= \mathbb{P}(\exists t_m > t_{m-1} ; B_{t_m} \in Q_m \text{ with } Q_m \cap \lambda Q_{m-1} = \emptyset \mid \exists t_{m-1} ; B_{t_{m-1}} \in Q_{m-1}) \cdot \\ & \mathbb{P}(\exists t_{m-1} > \dots > t_0 B_{t_{m-1}} \in Q_{m-1}, \dots, B_{t_0} \in Q) \end{aligned}$$

Now, by (12),

$$\mathbb{P}(\exists t_m > t_{m-1} ; B_{t_m} \in Q_m \text{ with } Q_m \cap \lambda Q_{m-1} = \emptyset \mid \exists t_{m-1} ; B_{t_{m-1}} \in Q_{m-1}) < p.$$

By induction we get that

$$\mathbb{P}(\exists t_m > t_{m-1} > \dots > t_0 ; \text{ and } Q, \dots, Q_m \text{ as above such that } B_{t_m} \in Q_m, \dots, B_{t_0} \in Q) < p^m$$

and hence the lemma. •

Recall that for a given dyadic Whitney cube Q we have defined the vicinity \tilde{Q} of Q as the union of all Whitney cubes Q' verifying

$$Q' \cap \lambda Q \neq \emptyset \text{ and } Q \cap \lambda Q' \neq \emptyset.$$

We can easily check that there are less than $(100\sqrt{d})^d$ such cubes Q' of size at most $\ell(Q)/12$ (the constants are not optimal). We say that the k -level layers are visited at least n times if there exist $t_0 < s_1 < t_1 < \dots < s_n < t_n$ satisfying

$$B_{t_j} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q \text{ and } B_{s_j} \notin \bigcup_{Q \in \mathcal{S}_{2^k}} \tilde{Q},$$

for all $j = 1, \dots, n$. For any $k \in \mathbb{Z}$ note

$$\nu_k = \sup\{n \in \mathbb{N} ; \text{ the } k\text{-level layers are visited at least } n \text{ times}\}$$

Lemma 5.2 *There exists $0 < p < 1$ and a positive constant C such that, given $k \in \mathbb{Z}$, for all $n \in \mathbb{N}$*

$$\mathbb{P}(\nu_k > n) \leq p^n \mathbb{P}(\exists t_0 > 0 \text{ and } Q \in \mathcal{S}_{2^k} ; B_{t_0} \in Q).$$

Proof The arguments are similar as in lemma 5.1. We only need to prove that $\mathbb{P}(\nu_k > 1) < p$ and apply strong Markov property. We have

$$\begin{aligned} \mathbb{P}(\nu_k > 1) &\leq \mathbb{P}\left(\exists 0 < t_0 < s_1 < t_1, Q \in \mathcal{S}_{2^k} ; B_{t_0} \in Q, B_{s_1} \notin \bigcup_{Q \in \mathcal{S}_{2^k}} \tilde{Q}, B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q\right) \\ &= \mathbb{P}\left(\exists t_1 > s_1 > t_0 ; B_{s_1} \notin \bigcup_{Q \in \mathcal{S}_{2^k}} \tilde{Q}, B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q \mid \exists t_0 > 0 ; B_{t_0} \in Q \in \mathcal{S}_{2^k}\right) \\ &\times \mathbb{P}(\exists t_0 > 0 \text{ and } Q \in \mathcal{S}_{2^k} ; B_{t_0} \in Q) \end{aligned}$$

To abbreviate formulas we note $\mathbb{P}_c(\cdot) = \mathbb{P}(\cdot \mid \exists t_0 > 0 ; B_{t_0} \in Q \in \mathcal{S}_{2^k})$. With this notation,

$$\begin{aligned} \mathbb{P}_c\left(\exists t_1 > s_1 > t_0 ; B_{s_1} \notin \bigcup_{Q \in \mathcal{S}_{2^k}} \tilde{Q}, B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q\right) &= \\ \mathbb{P}_c\left(\exists t_1 > s_1 ; B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q \mid A\right) \mathbb{P}_c(A) &+ \mathbb{P}_c\left(\exists t_1 > s_1 ; B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q \mid B\right) \mathbb{P}_c(B) \end{aligned}$$

where

$$A = \left\{ \exists s_1 > t_0 ; B_{s_1} \notin \bigcup_{Q \in \mathcal{S}_{2^k}} \lambda Q \right\} \text{ and}$$

$$B = \left\{ \exists s_1 > t_0 ; B_{s_1} \in Q' , \lambda Q' \cap \bigcup_{Q \in \mathcal{S}_{2^k}} Q = \emptyset , B_{s_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} \lambda Q \right\}$$

form a partition of the event $\left\{ \exists s_1 > t_0 ; B_{s_1} \notin \bigcup_{Q \in \mathcal{S}_{2^k}} \tilde{Q} \right\}$.

By (12), $\mathbb{P}_c(A) \leq p$. Similarly, by the strong Markov property of Brownian motion and by (12),

$$\mathbb{P}_c \left(\exists t_1 > s_1 ; B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q \mid B \right) = \mathbb{P} \left(\exists t_1 > s_1 ; B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q \mid B_{s_1} \in Q' \right) \leq p.$$

We deduce that

$$\mathbb{P}_c \left(\exists t_1 > s_1 > t_0 ; B_{s_1} \notin \bigcup_{Q \in \mathcal{S}_{2^k}} \tilde{Q} , B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q \right) \leq \mathbb{P}_c(A) + p(1 - \mathbb{P}_c(A))$$

The function $t \mapsto t + p(1 - t)$ being increasing on $[0, p]$ we get

$$\mathbb{P}_c \left(\exists t_1 > s_1 > t_0 ; B_{s_1} \notin \bigcup_{Q \in \mathcal{S}_{2^k}} \tilde{Q} , B_{t_1} \in \bigcup_{Q \in \mathcal{S}_{2^k}} Q \right) \leq 2p - p^2 < 1$$

and the lemma is proven. \bullet

Remark that by definition of dyadic Whitney cubes there exist k^* depending only on the dimension of the space ($k^* = \lceil \log_2(8\sqrt{d}) \rceil + 3$ will do) such that for all $k \in \mathbb{Z}$,

$$\{x \in \Omega ; 2^{k-1} \leq \text{dist}(x, \partial\Omega) \leq 2^k\} \subset \bigcup_{j=k-k^*}^k \bigcup_{Q \in \mathcal{Q}_j} Q. \quad (13)$$

Proof of lemma 4.1 We clearly have

$$\mathbb{P}(\delta\beta_{2^k} > t) \leq \mathbb{P}(\delta\beta_{2^k} > t , \nu_k + \dots + \nu_{k-k^*} > N) + \mathbb{P}(\delta\beta_{2^k} > t , \nu_k + \dots + \nu_{k-k^*} \leq N) \quad (14)$$

Given $t > 0$, $k \in \mathbb{Z}$, by lemma 5.2 we have, for all N ,

$$\mathbb{P}(\delta\beta_{2^k} > t , \nu_k + \dots + \nu_{k-k^*} > N) < k^* p^{\frac{1}{k^*}N} \mathbb{P}(\exists t_0 > 0 \text{ and } Q \in \mathcal{S}_{2^{k-k^*}} ; B_{t_0} \in Q) < c\tilde{p}^N$$

Let us estimate the second term of the sum (14). By (13) and the definition of ν_k we get

$$\begin{aligned} & \mathbb{P}(\delta\beta_{2^k} > t , \nu_k + \dots + \nu_{k-k^*} \leq N) \\ & \leq \mathbb{P}(\exists l \leq N , Q_1, \dots, Q_l \in \bigcup_{j=k-k^*}^k \mathcal{Q}_j ; Q_s \cap \tilde{Q}_{s-1} = \emptyset , \forall s = 2, \dots, l \text{ and} \\ & \quad \exists t_1 < s_1 \leq t_2 < s_2 < \dots \leq t_l < s_l ; B_{[t_i, s_i]} \subset \tilde{Q}_s \forall s = 1, \dots, l \text{ and } \sum_{i=1}^l s_i - t_i > t). \end{aligned}$$

Since $l \leq N$ we get that $\mathbb{P}(\delta\beta_{2^k} > t, \nu_k + \dots + \nu_{k-k^*} \leq N)$ is bounded above by

$$\mathbb{P}\left(\exists Q \in \mathcal{Q}_k \cup \dots \cup \mathcal{Q}_{k-k^*}; \exists 0 < t_{i_0} < s_{i_0} \text{ with } B_{[t_{i_0}, s_{i_0}]} \subset \tilde{Q} \text{ and } s_{i_0} - t_{i_0} > t/N\right),$$

which completes the proof. •

5.2 Proof of lemma 4.2

The ideas are the same but we will work with cubes touching the boundary instead of Whitney cubes.

Lemma 5.3 *Under the Δ -regularity hypothesis, the probability that BM started at distance $\varepsilon < r$ from the boundary gets at distance greater than R from the starting point without leaving the Minkowski sausage $M_r = \{x \in \Omega ; \text{dist}(x, \partial\Omega) \leq r\}$ is bounded above by $cp^{R/r}$ where $c > 0$ and $0 < p < 1$ are constants (depending only on L of the Δ -regularity hypothesis and on d).*

The proof, similar to the one of lemma 5.1, is therefore abridged.

Proof Let $x \in M_r$ and consider the annuli centered at x of inner radii $4\ell r$ and outer radii $4(\ell+1)r$ where $\ell = 0, \dots, \lfloor \frac{R}{4r} \rfloor$. Brownian motion started at x and moving at distance R from x before exiting Ω must go through all these annuli. The probability of going through such an annulus while staying at distance at most r from the boundary is bounded by a $p_0 \in (0, 1)$ by the Δ -regularity hypothesis. To see this take any point y in the middle of the annulus (i.e. at distance $\frac{6\ell+6}{r}$ from x) and consider the ball of center y and radius $2r$. If $d_y < r$, the probability to exit the ball without touching $\partial\Omega$ is uniformly bounded away from 1 by the same hypothesis. This probability being greater than the probability of going through the annulus we have the statement. By the independence of the “crossing annuli” events we get that the probability that Brownian motion goes through all the annuli is smaller than $cp_0^{\lfloor \frac{R}{4r} \rfloor} \sim c\tilde{p}^{\lfloor \frac{R}{r} \rfloor}$. •

We say that the Minkowski sausage M_r is visited by the Brownian motion at least k times if there exist $t_0 < s_1 < t_1 < \dots < s_n < t_n < \tau_\Omega$ satisfying $B_{t_i} \in M_r$ for all $i = 0, \dots, n$ and $B_{s_i} \notin M_{4r}$.

In a similar way with ν_k we define ξ_k as the

$$\xi_r = \sup\{n \in \mathbb{N} ; M_r \text{ has been visited at least } k \text{ times}\}.$$

Lemma 5.4 *There exists $0 < p < 1$ depending only on the Δ -regularity's L such that, given $r > 0$, for all $n \in \mathbb{N}$*

$$\mathbb{P}(\xi_r > n) \leq p^n.$$

Proof The proof of this lemma is a straightforward application of the Δ -regularity condition. It suffices to show that there exists $0 < p < 1$ such that

$$\mathbb{P}((\exists t_1 > s_1 > t_0 ; B_{t_0} \in M_r, B_{s_1} \notin M_{4r}, B_{t_1} \in M_r) < p,$$

and then apply the Markov property. Remark that, the probability

$$\mathbb{P}((\exists t_1 > s_1 > t_0 ; B_{t_0} \in M_r, B_{s_1} \notin M_{4r})$$

is smaller than the probability that brownian motion started at t_0 exits a ball of radius $2r > 2 \text{ dist}(B_{t_0}, \partial\Omega)$ without hitting $\partial\Omega$. By the Δ -regularity this last probability is bounded by a constant $p < 1$. •

Proof of lemma 4.2 As before we get :

$$\mathbb{P}(\delta\beta_{2^{k_0}} > t) = \mathbb{P}(\delta\beta_{2^{k_0}} > t, \xi_{2^{k_0}} > N) + \mathbb{P}(\delta\beta_{2^{k_0}} > t, \xi_{2^{k_0}} \leq N) \quad (15)$$

By lemma 5.4 we get $\mathbb{P}(\delta\beta_{2^{k_0}} > t, \xi_{2^{k_0}} > N) \leq p^N$, where $0 < p < 1$, for all $N \in \mathbb{N}$. Let us now deal with the second term of the sum.

$$\mathbb{P}(\delta\beta_{2^{k_0}} > t, \xi_{2^{k_0}} \leq N) \leq \mathbb{P}(\exists s_1 < s_2 < \tau_\Omega ; s_2 - s_1 > t/N, B_{[s_1, s_2]} \subset M_{4r})$$

Using lemma 5.3 we get that, for $R > 4r$, this probability is bounded by

$$cp^{\frac{R}{r}} + \mathbb{P}(\exists s_1 < s_2 < \tau_\Omega ; s_2 - s_1 > t/N, B_{[s_1, s_2]} \subset M_{4r} \cap \mathbb{B}(B_{s_1}, R)),$$

and the statement of the lemma 4.2 follows on taking $N = \lceil \frac{R}{r} \rceil$ •

6 Further Comments

We should point out that hypothesis (4) in theorem 3.1 can be dropped; in this case, the same reasoning as in subsection 4.2 gives the lower bound

$$\left(\frac{\#\mathcal{S}_{\sqrt{t}}}{\#\mathcal{S}_\epsilon} \right)^{d_M} \left(\frac{\sqrt{t}}{\epsilon} \right)^{d-2}.$$

The best upper bound is less evident; nevertheless a slight improvement of the above proofs gives

$$\mathbb{P}(\tau_\Omega > t) \leq c \left(\log \frac{\sqrt{t}}{\epsilon} \right)^{cd} \sup_N \left(2^{-Nd} + \sum_{k=0}^N \left(\frac{\#\mathcal{S}_{\sqrt{t}/2^k}}{\#\mathcal{S}_\epsilon} \right)^{d_M} \left(\frac{2^{-k}\sqrt{t}}{\epsilon} \right)^{d-2} \right).$$

References

- [Anc86] A. Ancona. On strong barriers and an inequality of Hardy for domains in \mathbb{R}^n . *Journal of the London Mathematical Society*, **34** (2): 274–290, 1986.
- [Bis96] C. J. Bishop. Minkowski dimension and the Poincaré exponent. *Michigan Mathematical Journal*, **43** : 231–246, 1996.
- [BLZ09] A. Batakis, P. Levitz, and M. Zinsmeister. Brownian flights. *To appear in the PAMQ*, 2009.
- [DPP⁺08] D.Constantin, P.Davidson, P.Levitz, O.Poncelet, and M.Zinsmeister. Intermittent brownian dynamics over a rigid strand: Heavily tailed relocation statistics in a simple geometry. *Physical Review E*, **78** (3), 2008.
- [Fel71] W. Feller. *An introduction to probability theory and its applications*, volume II. John Wiley and sons, second edition, 1971.
- [GKL⁺06] D. Grebenkov, K. Kolwankar, P. Levitz, B. Sapoval, and M. Zinsmeister. Brownian flights over a fractal nest and first-passage statistics on irregular surfaces. *Phys. Rev. Lett.*, **96** , 2006.
- [Gra08] L. Grafakos. *Classical Fourier Analysis*. Springer, New York, 2008.
- [HKM93] J. Heinonen, T. Kilpeläinen, and O. Martio. *Nonlinear Potential Theory of Degenerate Elliptic Equations*. Clarendon Press, 1993.
- [JK82] D. Jerison and C. Kenig. Boundary behaviour of harmonic functions in non-tangentially accessible domains. *Advances in Mathematics*, **46** : 80–147, 1982.
- [JW88] P. Jones and T. Wolff. Hausdorff dimension of harmonic measures in the plane. *Acta Mathematica*, **161** : 131–144, 1988.