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On certain time-space fractional evolution equations

Ahmad Z. FINO^{†*} and Mokhtar KIRANE[†]

Abstract

In this article, we present first a new technique to prove, in a general case, the recent result of Cazenave, Dickstein and Weissler (Nonlinear Analysis, 2008) on the blowing-up solutions to a temporally nonlocal nonlinear parabolic equation. Then, we study the blow-up rate, and the existence of global solutions. Furthermore, we establish necessary conditions for local or global existence.

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1 Introduction

In this paper, we investigate the nonlinear parabolic equation with a nonlocal in time nonlinearity

$$\begin{cases} u_t + (-\Delta)^{\beta/2}u = \frac{1}{\Gamma(1-\gamma)} \int_0^t (t-s)^{-\gamma} |u|^{p-1}u(s) ds & x \in \mathbb{R}^N, t > 0, \\ u(x, 0) = u_0(x) & x \in \mathbb{R}^N, \end{cases} \quad (1.1)$$

where $u_0 \in C_0(\mathbb{R}^N)$, $N \geq 1$, $0 < \beta \leq 2$, $0 < \gamma < 1$, $p > 1$ and the nonlocal operator $(-\Delta)^{\beta/2}$ is defined by

$$(-\Delta)^{\beta/2}v(x) := \mathcal{F}^{-1} \left(|\xi|^\beta \mathcal{F}(v)(\xi) \right) (x)$$

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for every $v \in D((-\Delta)^{\beta/2}) = H^\beta(\mathbb{R}^N)$, where $H^\beta(\mathbb{R}^N)$ is the homogeneous Sobolev space of order β defined by

$$H^\beta(\mathbb{R}^N) = \left\{ u \in \mathcal{S}' ; (-\Delta)^{\beta/2} u \in L^2(\mathbb{R}^N) \right\} \quad \text{if } \beta \notin \mathbb{N},$$

$$H^\beta(\mathbb{R}^N) = \left\{ u \in L^2(\mathbb{R}^N) ; (-\Delta)^{\beta/2} u \in L^2(\mathbb{R}^N) \right\} \quad \text{if } \beta \in \mathbb{N},$$

where \mathcal{S}' is the space of Schwartz distributions. \mathcal{F} stands for the Fourier transform and \mathcal{F}^{-1} for its inverse, Γ is the Euler gamma function and $C_0(\mathbb{R}^N)$ denotes the space of all continuous functions decaying to zero at infinity.

When Eq. (1.1) is considered with a nonlinearity of the form u^p , it reads

$$u_t + (-\Delta)^{\beta/2} u = u^p.$$

This equation has been considered by Nagasawa and Sirao [27], Kobayashi [21], Fino and Karch [13], Birkner, Lopez-Mimbela and Wakolbinger [6], Guedda and Kirane [15], and by Kirane and Qafsaoui [20].

The fractional Laplacian $(-\Delta)^{\beta/2}$ is related to Lévy flights in physics. Many observations and experiments related to Lévy flights (super-diffusion), e.g., collective slip diffusion on solid surfaces, quantum optics or Richardson turbulent diffusion, have been performed in recent years. The symmetric β -stable processes ($\beta \in (0, 2)$) are the basic characteristics for a class of jumping Lévy's processes. Compared with the continuous Brownian motion ($\beta = 2$), symmetric β -stable processes have infinite jumps in an arbitrary time interval. The large jumps of these processes make their variances and expectations infinite according to $\beta \in (0, 2)$ or $\beta \in (0, 1]$, respectively (see [18]). Let us also mention that when $\beta = 3/2$, the symmetric β -stable processes appear in the study of stellar dynamics (see [11]).

As a physical motivation, the problem (1.1) suggested to us the possibility of a physical model in which a super-diffusive medium is coupled to a classically diffusive medium. The right-side of (1.1) might be interpreted as the effect of a classically diffusive medium that is nonlinearly linked to a super-diffusive medium. Such a link might come in the form of a porous material with reactive properties that is partially insulated by contact with a classically diffusive material. For more informations see the recent result of Roberts and Olmstead [28].

Our article is motivated mathematically by the recent and very interesting paper by Cazenave, Dickstein and Weissler [10] which deals with the global existence and blow-up for the parabolic equation with nonlocal in time nonlinearity

$$u_t - \Delta u = \int_0^t (t-s)^{-\gamma} |u|^{p-1} u(s) ds \quad x \in \mathbb{R}^N, t > 0, \quad (1.2)$$

where $0 \leq \gamma < 1$, $p > 1$ and $u_0 \in C_0(\mathbb{R}^N)$, which is a particular case of (1.1); it corresponds to $\beta = 2$.

If we set

$$p_\gamma = 1 + \frac{2(2-\gamma)}{(N-2+2\gamma)_+}$$

and

$$p_* = \max \left\{ \frac{1}{\gamma}, p_\gamma \right\} \in (0, +\infty],$$

where $(\cdot)_+$ is the positive part, they proved that

- (i) If $\gamma \neq 0$, $p \leq p_*$, and $u_0 \geq 0$, $u_0 \not\equiv 0$, then u blows up in finite time.
- (ii) If $\gamma \neq 0$, $p > p_*$, and $u_0 \in L^{q_{sc}}(\mathbb{R}^N)$ (where $q_{sc} = N(p-1)/(4-2\gamma)$) with $\|u_0\|_{L^{q_{sc}}}$ sufficiently small, then u exists globally.

If $\gamma = 0$ then all nontrivial positive solutions blow-up as proved by Souplet in [31]. The study of Cazenave, Dickstein and Weissler reveals the surprising fact that for equation (1.2) the critical exponent in Fujita's sense p_* is not the one predicted by scaling. This can be explained by the fact that their equation can be formally converted into

$$D_{0|t}^\alpha u_t - D_{0|t}^\alpha \Delta u = |u|^{p-1}u, \quad (1.3)$$

where $D_{0|t}^\alpha$ is the left-sided Riemann-Liouville fractional derivative of order $\alpha \in (0, 1)$ defined in (2.8) below (we have set in (1.3), $\alpha = 1 - \gamma \in (0, 1)$).

Eq. (1.3) is pseudo-parabolic equation and as is well known scaling is efficient for detecting the Fujita exponent only for equations of parabolic type.

Needless to say that the equation considered by Cazenave, Dickstein and Weissler [10] is a genuine extension of the one considered by Fujita in his pioneering work [14].

In this article, concerning blowing-up solutions, we present a different proof from the one presented in [10], and for the more general equation (1.1). Our proof is more versatile and can be applied to more nonlinear equations (see the Remarks in Section 4).

Our analysis is based on the observation that the nonlinear differential equation (1.1) can be written in the form:

$$u_t + (-\Delta)^{\beta/2}u = J_{0|t}^\alpha (|u|^{p-1}u), \quad (1.4)$$

where $\alpha := 1 - \gamma \in (0, 1)$ and $J_{0|t}^\alpha$ is the Riemann-Liouville fractional integral defined in (2.10).

We will show that:

- (1) For $u_0 \geq 0$, $u_0 \not\equiv 0$, and $u_0 \in C_0(\mathbb{R}^N)$, if

$$p \leq 1 + \frac{\beta(2-\gamma)}{(N-\beta+\beta\gamma)_+} \quad \text{or} \quad p < \frac{1}{\gamma},$$

then all solutions of problem (1.1) blow-up in finite time.

- (2) For $u_0 \in C_0(\mathbb{R}^N) \cap L^{p_{sc}}(\mathbb{R}^N)$, where $p_{sc} := N(p-1)/\beta(2-\gamma)$, if

$$p > \max \left\{ 1 + \frac{\beta(2-\gamma)}{(N-\beta+\beta\gamma)_+}; \frac{1}{\gamma} \right\},$$

and $\|u_0\|_{L^{p_{sc}}}$ is sufficiently small, then u exists globally.

The method used to prove the blow-up theorem is the test function method introduced by Baras and Pierre [5], Mitidieri and Pohozaev [24, 25], Kirane et al. [15, 19]; it was also used by Baras and Kersner [4] for the study of the necessary conditions for the local existence.

Furthermore, in the case $\beta = 2$, we derive the blow-up rate estimates for the parabolic equation (1.1). We shall prove that, if $u_0 \in C_0(\mathbb{R}^N) \cap L^2(\mathbb{R}^N)$, $u_0 \geq 0$, $u_0 \not\equiv 0$ and if u is the blowing-up solution of (1.1) at the finite time $T^* > 0$, then there are constants $c, C > 0$ such that $c(T^* - t)^{-\alpha_1} \leq \sup_{\mathbb{R}^N} u(\cdot, t) \leq C(T^* - t)^{-\alpha_1}$ for $1 < p \leq 1 + 2(2 - \gamma)/(N - 2 + 2\gamma)_+$ or $1 < p < 1/\gamma$ and all $t \in (0, T^*)$, where $\alpha_1 := (2 - \gamma)/(p - 1)$. We use a scaling argument to reduce the problems of blow-up rate to Fujita-type theorems (it is similar to blow-up analysis in elliptic problems to reduce the problems of a priori bounds to Liouville-type theorems). As far as we know, this method was first applied to parabolic problems by Hu [16], and then was used in various parabolic equations and systems (see [9, 12]). We notice that in the limiting case when $\gamma \rightarrow 0$, we obtain the constant rate $2/(p - 1)$ found by P. Souplet [32].

For more informations, we refer the reader to the excellent paper of Andreucci and Tedeev [1] for the blow-up rate by an alternative method

The organization of this paper is as follows. In Section 2, we present some definitions and properties. In Section 3, we derive the local existence of solutions for the parabolic equation (1.1). Section 4 contains the blow-up result of solutions for (1.1). Section 5 is dedicated to the blow-up rate of solutions. Global existence is studied in Section 6. Finally, a necessary condition for local existence and a necessary condition for global existence are given in Section 7.

2 Preliminaries

In this section, we present some definitions and results concerning fractional Laplacians, fractional integrals and fractional derivatives that will be used hereafter.

First, if we take the usual linear fractional diffusion equation

$$u_t + (-\Delta)^{\beta/2} u = 0, \quad \beta \in (0, 2], \quad x \in \mathbb{R}^N, \quad t > 0, \quad (2.1)$$

then, its fundamental solution S_β can be represented via the Fourier transform by

$$S_\beta(t)(x) := S_\beta(x, t) = \frac{1}{(2\pi)^{N/2}} \int_{\mathbb{R}^N} e^{ix \cdot \xi - t|\xi|^\beta} d\xi. \quad (2.2)$$

It is well-known that this function satisfies

$$S_\beta(1) \in L^\infty(\mathbb{R}^N) \cap L^1(\mathbb{R}^N), \quad S_\beta(x, t) \geq 0, \quad \int_{\mathbb{R}^N} S_\beta(x, t) dx = 1, \quad (2.3)$$

for all $x \in \mathbb{R}^N$ and $t > 0$. Hence, using Young's inequality for the convolution and the following self-similar form $S_\beta(x, t) = t^{-N/\beta} S_\beta(xt^{-1/\beta}, 1)$, we have

$$\|S_\beta(t) * v\|_q \leq Ct^{-\frac{N}{\beta}(\frac{1}{r} - \frac{1}{q})} \|v\|_r, \quad (2.4)$$

for all $v \in L^r(\mathbb{R}^N)$ and all $1 \leq r \leq q \leq \infty$, $t > 0$.

Moreover, as $(-\Delta)^{\beta/2}$ is a self-adjoint operator with $D((-\Delta)^{\beta/2}) = H^\beta(\mathbb{R}^N)$, we have

$$\int_{\mathbb{R}^N} u(x)(-\Delta)^{\beta/2}v(x) dx = \int_{\mathbb{R}^N} v(x)(-\Delta)^{\beta/2}u(x) dx, \quad (2.5)$$

for all $u, v \in H^\beta(\mathbb{R}^N)$.

In an open bounded domain Ω , we denote by $\Delta_D^{\beta/2}$ the fractional Laplacian in Ω with homogeneous Dirichlet boundary condition. We have the following facts:

Let λ_k ($k = 1, \dots, +\infty$) be the eigenvalues for the Laplacian operator in $L^2(\Omega)$ with homogeneous Dirichlet boundary condition and let φ_k be the eigenfunction corresponding to λ_k . Then

$$\begin{cases} \Delta_D^{\beta/2} \varphi_k = \lambda_k^{\beta/2} \varphi_k & \text{in } \Omega \\ \varphi_k = 0 & \text{on } \partial\Omega \end{cases} \quad (2.6)$$

and

$$D(\Delta_D^{\beta/2}) = \left\{ u \in L^2(\Omega) \text{ s.t. } u|_{\partial\Omega} = 0; \|\Delta_D^{\beta/2} u\|_{L^2(\Omega)} := \sum_{k=1}^{k=+\infty} |\lambda_k^{\beta/2} \langle u, \varphi_k \rangle|^2 < +\infty \right\}.$$

So, for $u \in D(\Delta_D^{\beta/2})$ we have

$$\Delta_D^{\beta/2} u = \sum_{k=1}^{k=+\infty} \lambda_k^{\beta/2} \langle u, \varphi_k \rangle \varphi_k,$$

and then we get the following integration by parts

$$\int_{\Omega} u(x)\Delta_D^{\beta/2}v(x) dx = \int_{\Omega} v(x)\Delta_D^{\beta/2}u(x) dx, \quad (2.7)$$

for all $u, v \in D(\Delta_D^{\beta/2})$.

Next, if $AC[0, T]$ is the space of all functions which are absolutely continuous on $[0, T]$ with $0 < T < \infty$, then, for $f \in AC[0, T]$, the left-handed and right-handed Riemann-Liouville fractional derivatives $D_{0|t}^\alpha f(t)$ and $D_{t|T}^\alpha f(t)$ of order $\alpha \in (0, 1)$ are defined by (see [18])

$$D_{0|t}^\alpha f(t) := D J_{0|t}^{1-\alpha} f(t), \quad (2.8)$$

$$D_{t|T}^\alpha f(t) := -\frac{1}{\Gamma(1-\alpha)} D \int_t^T (s-t)^{-\alpha} f(s) ds, \quad (2.9)$$

for all $t \in [0, T]$, where $D := d/dt$ is the usual derivative, and

$$J_{0|t}^\alpha f(t) := \frac{1}{\Gamma(\alpha)} \int_0^t (t-s)^{\alpha-1} f(s) ds \quad (2.10)$$

is the Riemann-Liouville fractional integral defined in [18], for all $f \in L^q(0, T)$ ($1 \leq q \leq \infty$). Now, for every $f, g \in C([0, T])$, such that $D_{0|t}^\alpha f(t), D_{t|T}^\alpha g(t)$ exist and are continuous, for

all $t \in [0, T]$, $0 < \alpha < 1$, we have the formula of integration by parts (see (2.64) p. 46 in [30])

$$\int_0^T \left(D_{0|t}^\alpha f \right) (t) g(t) dt = \int_0^T f(t) \left(D_{t|T}^\alpha g \right) (t) dt. \quad (2.11)$$

Note also that, for any $f \in AC^2[0, T]$, we have (see (2.2.30) in [18])

$$-D.D_{t|T}^\alpha f = D_{t|T}^{1+\alpha} f, \quad (2.12)$$

where

$$AC^2[0, T] := \{f : [0, T] \rightarrow \mathbb{R} \text{ such that } Df \in AC[0, T]\}.$$

Moreover, for all $1 \leq q \leq \infty$, the following equality (see [18, Lemma 2.4 p.74])

$$D_{0|t}^\alpha J_{0|t}^\alpha = Id_{L^q(0, T)} \quad (2.13)$$

holds almost everywhere on $[0, T]$.

Later on, we will use the following results.

- If $w_1(t) = (1 - t/T)_+^\sigma$, $t \geq 0$, $T > 0$, $\sigma \gg 1$, then

$$D_{t|T}^\alpha w_1(t) = \frac{(1 - \alpha + \sigma)\Gamma(\sigma + 1)}{\Gamma(2 - \alpha + \sigma)} T^{-\alpha} \left(1 - \frac{t}{T}\right)_+^{\sigma - \alpha}, \quad (2.14)$$

$$D_{t|T}^{\alpha+1} w_1(t) = \frac{(1 - \alpha + \sigma)(\sigma - \alpha)\Gamma(\sigma + 1)}{\Gamma(2 - \alpha + \sigma)} T^{-(\alpha+1)} \left(1 - \frac{t}{T}\right)_+^{\sigma - \alpha - 1}, \quad (2.15)$$

for all $\alpha \in (0, 1)$; so

$$\left(D_{t|T}^\alpha w_1 \right) (T) = 0 \quad ; \quad \left(D_{t|T}^\alpha w_1 \right) (0) = C T^{-\alpha}, \quad (2.16)$$

where $C = (1 - \alpha + \sigma)\Gamma(\sigma + 1)/\Gamma(2 - \alpha + \sigma)$; indeed, using the Euler change of variable $y = (s - t)/(T - t)$, we get

$$\begin{aligned} D_{t|T}^\alpha w_1(t) &:= -\frac{1}{\Gamma(1 - \alpha)} D \left[\int_t^T (s - t)^{-\alpha} \left(1 - \frac{s}{T}\right)^\sigma ds \right] \\ &= -\frac{T^{-\sigma}}{\Gamma(1 - \alpha)} D \left[(T - t)^{1 - \alpha + \sigma} \int_0^1 (y)^{-\alpha} (1 - y)^\sigma ds \right] \\ &= \frac{(1 - \alpha + \sigma)B(1 - \alpha; \sigma + 1)}{\Gamma(1 - \alpha)} T^{-\sigma} (T - t)^{\sigma - \alpha}, \end{aligned}$$

where $B(\cdot; \cdot)$ stands for the beta function. Then, (2.14) follows using the relation

$$B(1 - \alpha; \sigma + 1) = \frac{\Gamma(1 - \alpha)\Gamma(\sigma + 1)}{\Gamma(2 - \alpha + \sigma)}.$$

Moreover, using (2.12) and (2.14), we obtain (2.15).

- If $w_2(t) = (1 - t^2/T^2)_+^\ell$, $T > 0$, $\ell \gg 1$, then, using the change of variable $y = (s -$

$t)/(T - t)$, we have

$$D_{t|T}^\alpha w_2(t) = \frac{T^{-\alpha}}{\Gamma(1-\alpha)} \sum_{k=0}^{\ell} C_1(\ell, k, \alpha) \left(1 - \frac{t}{T}\right)^{\ell+k-\alpha}, \quad (2.17)$$

$$D_{t|T}^{1+\alpha} w_2(t) = \frac{T^{-\alpha-1}}{\Gamma(1-\alpha)} \sum_{k=0}^{\ell} C_2(\ell, k, \alpha) \left(1 - \frac{t}{T}\right)^{\ell+k-\alpha-1}, \quad (2.18)$$

for all $-T \leq t \leq T$, $\alpha \in (0, 1)$, where

$$\begin{cases} C_1(\ell, k, \alpha) := c_\ell^k (1 - \alpha + \ell + k) 2^{\ell-k} (-1)^k \frac{\Gamma(k+\ell+1)\Gamma(1-\alpha)}{\Gamma(k+\ell+2-\alpha)}, \\ C_2(\ell, k, \alpha) := (\ell + k - \alpha) C_1(\ell, k, \alpha), \\ c_\ell^k := \frac{\ell!}{(\ell-k)!k!}; \end{cases}$$

so

$$\left(D_{t|T}^\alpha w_2\right)(T) = 0 \quad ; \quad \left(D_{t|T}^\alpha w_2\right)(-T) = C_3(\ell, k, \alpha) T^{-\alpha}, \quad (2.19)$$

where

$$C_3(\ell, k, \alpha) := \frac{2^{2\ell-\alpha}(-1)^\ell}{\Gamma(1-\alpha)} \sum_{k=0}^{\ell} c_\ell^k (1 - \alpha + \ell + k) \frac{\Gamma(k+\ell+1)\Gamma(1-\alpha)}{\Gamma(k+\ell+2-\alpha)}.$$

3 Local existence

This section is dedicated to proving the local existence and uniqueness of mild solutions to the problem (1.1). Let $T(t) := e^{-t(-\Delta)^{\beta/2}}$. As $(-\Delta)^{\beta/2}$ is a positive definite self-adjoint operator in $L^2(\mathbb{R}^N)$, $T(t)$ is a strongly continuous semigroup on $L^2(\mathbb{R}^N)$ generated by the fractional power $(-\Delta)^{\beta/2}$ (see Yosida [33]). It holds $T(t)v = S_\beta(t) * v$, where S_β is given by (2.2) and $u * v$ is the convolution of u and v . We start by giving the

Definition 1 (Mild solution)

Let $u_0 \in C_0(\mathbb{R}^N)$, $0 < \beta \leq 2$, $p > 1$ and $T > 0$. We say that $u \in C([0, T], C_0(\mathbb{R}^N))$ is a mild solution of the problem (1.1) if u satisfies the following integral equation

$$u(t) = T(t)u_0 + \int_0^t T(t-s) J_{0|s}^\alpha (|u|^{p-1}u) ds, \quad t \in [0, T]. \quad (3.1)$$

Theorem 1 (Local existence)

Given $u_0 \in C_0(\mathbb{R}^N)$ and $p > 1$, there exist a maximal time $T_{\max} > 0$ and a unique mild solution $u \in C([0, T_{\max}], C_0(\mathbb{R}^N))$ to the problem (1.1). Furthermore, either $T_{\max} = \infty$ or else $T_{\max} < \infty$ and $\|u\|_{L^\infty((0, t) \times \mathbb{R}^N)} \rightarrow \infty$ as $t \rightarrow T_{\max}$. In addition, if $u_0 \geq 0$, $u_0 \not\equiv 0$, then $u(t) > 0$ for all $0 < t < T_{\max}$. Moreover, if $u_0 \in L^r(\mathbb{R}^N)$, for $1 \leq r < \infty$, then $u \in C([0, T_{\max}], L^r(\mathbb{R}^N))$.

Proof. For arbitrary $T > 0$, we define the Banach space

$$E_T := \{u \in L^\infty((0, T), C_0(\mathbb{R}^N)); \|u\|_1 \leq 2\|u_0\|_{L^\infty}\},$$

where $\|\cdot\|_1 := \|\cdot\|_{L^\infty((0,T),L^\infty(\mathbb{R}^N))}$. Next, for every $u \in E_T$, we define

$$\Psi(u) := T(t)u_0 + \int_0^t T(t-s)J_{0|s}^\alpha (|u|^{p-1}u) ds.$$

As usual, we prove the local existence by the Banach fixed point theorem.

• $\Psi : \mathbf{E}_T \rightarrow \mathbf{E}_T$: Let $u \in E_T$, using (2.4), we obtain with $\|\cdot\|_\infty := \|\cdot\|_{L^\infty(\mathbb{R}^N)}$

$$\begin{aligned} \|\Psi(u)\|_1 &\leq \|u_0\|_\infty + \frac{1}{\Gamma(1-\gamma)} \left\| \int_0^t \int_0^s (s-\sigma)^{-\gamma} \|u(\sigma)\|_\infty^p d\sigma ds \right\|_{L^\infty(0,T)} \\ &= \|u_0\|_\infty + \frac{1}{\Gamma(1-\gamma)} \left\| \int_0^t \int_\sigma^t (s-\sigma)^{-\gamma} \|u(\sigma)\|_\infty^p ds d\sigma \right\|_{L^\infty(0,T)} \\ &\leq \|u_0\|_\infty + \frac{T^{2-\gamma}}{(1-\gamma)(2-\gamma)\Gamma(1-\gamma)} \|u\|_1^p \\ &\leq \|u_0\|_\infty + \frac{T^{2-\gamma} 2^p \|u_0\|_\infty^{p-1}}{\Gamma(3-\gamma)} \|u_0\|_\infty, \end{aligned}$$

thanks to the following formula $\Gamma(x+1) = x\Gamma(x)$ for all $x > 0$. Now, if we choose T small enough such that

$$\frac{T^{2-\gamma} 2^p \|u_0\|_\infty^{p-1}}{\Gamma(3-\gamma)} \leq 1, \quad (3.2)$$

we conclude that $\|\Psi(u)\|_1 \leq 2\|u_0\|_\infty$, and then $\Psi(u) \in E_T$.

• Ψ is a contracting map: For $u, v \in E_T$, taking account of (2.4), we have

$$\begin{aligned} \|\Psi(u) - \Psi(v)\|_1 &\leq \frac{1}{\Gamma(1-\gamma)} \left\| \int_0^t \int_0^s (s-\sigma)^{-\gamma} \left\| |u|^{p-1}u(\sigma) - |v|^{p-1}v(\sigma) \right\|_\infty d\sigma ds \right\|_{L^\infty(0,T)} \\ &= \frac{1}{\Gamma(1-\gamma)} \left\| \int_0^t \int_\sigma^t (s-\sigma)^{-\gamma} \left\| |u|^{p-1}u(\sigma) - |v|^{p-1}v(\sigma) \right\|_\infty ds d\sigma \right\|_{L^\infty(0,T)} \\ &\leq \frac{T^{2-\gamma}}{\Gamma(3-\gamma)} \left\| |u|^{p-1}u - |v|^{p-1}v \right\|_1 \\ &\leq \frac{C(p) 2^p \|u_0\|_\infty^{p-1} T^{2-\gamma}}{\Gamma(3-\gamma)} \|u - v\|_1 \\ &\leq \frac{1}{2} \|u - v\|_1, \end{aligned}$$

thanks to the following inequality

$$\left| |u|^{p-1}u - |v|^{p-1}v \right| \leq C(p) |u - v| (|u|^{p-1} + |v|^{p-1}); \quad (3.3)$$

T is chosen such that

$$\frac{T^{2-\gamma} 2^p \|u_0\|_\infty^{p-1} \max(2C(p), 1)}{\Gamma(3-\gamma)} \leq 1. \quad (3.4)$$

Then, by the Banach fixed point theorem, there exists a unique mild solution $u \in \Pi_T$, where $\Pi_T := L^\infty((0, T), C_0(\mathbb{R}^N))$, to the problem (1.1).

• **Uniqueness of solution:** If u, v are two mild solutions in E_T for some $T > 0$, using (2.4) and (3.3), we obtain

$$\begin{aligned} \|u(t) - v(t)\|_\infty &\leq \frac{C(p)2^p \|u_0\|_\infty^{p-1}}{\Gamma(1-\gamma)} \int_0^t \int_0^s (s-\sigma)^{-\gamma} \|u(\sigma) - v(\sigma)\|_\infty d\sigma ds \\ &= \frac{C(p)2^p \|u_0\|_\infty^{p-1}}{\Gamma(1-\gamma)} \int_0^t \int_\sigma^t (s-\sigma)^{-\gamma} \|u(\sigma) - v(\sigma)\|_\infty ds d\sigma \\ &= \frac{C(p)2^p \|u_0\|_\infty^{p-1}}{\Gamma(2-\gamma)} \int_0^t (t-\sigma)^{1-\gamma} \|u(\sigma) - v(\sigma)\|_\infty d\sigma. \end{aligned}$$

So the uniqueness follows from Gronwall's inequality for this type of nonlinearities (cf. [8]).

Next, using the uniqueness of solutions, we conclude the existence of a solution on a maximal interval $[0, T_{\max})$ where

$$T_{\max} := \sup \{T > 0 ; \text{there exist a mild solution } u \in \Pi_T \text{ to (1.1)}\} \leq +\infty.$$

Note that, using the continuity of the semigroup $T(t)$, we can easily conclude that

$$u \in C([0, T_{\max}), C_0(\mathbb{R}^N)).$$

Moreover, if $0 \leq t \leq t + \tau < T_{\max}$, using (3.1), we obtain

$$\begin{aligned} u(t + \tau) &= T(\tau)u(t) + \frac{1}{\Gamma(1-\gamma)} \int_0^\tau T(\tau-s) \int_0^s (s-\sigma)^{-\gamma} |u|^{p-1} u(t+\sigma) d\sigma ds \\ &\quad + \frac{1}{\Gamma(1-\gamma)} \int_0^\tau T(\tau-s) \int_0^t (t+s-\sigma)^{-\gamma} |u|^{p-1} u(\sigma) d\sigma ds. \end{aligned} \quad (3.5)$$

To prove that if $T_{\max} < \infty$, then $\|u(t)\|_{L^\infty(\mathbb{R}^N)} \rightarrow \infty$ as $t \rightarrow T_{\max}$, we proceed by contradiction. Suppose that u is a solution of (3.1) on some interval $[0, T]$ with $\|u\|_{L^\infty((0, T) \times \mathbb{R}^N)} < \infty$ and $T_{\max} < \infty$. So, using the fact that the last term in (3.5) depends only on the values of u in the interval $(0, t)$ and using again a fixed-point argument, we conclude that u can be extended to a solution on some interval $[0, T']$ with $T' > T$. If we repeat this iteration, we obtain a contradiction with the fact that the maximal time T_{\max} is finite.

• **Positivity of solutions:** If $u_0 \geq 0$ and $u_0 \not\equiv 0$, then we can construct a nonnegative solution on some interval $[0, T]$ by applying the fixed point argument in the set $E_T^+ = \{u \in E_T; u \geq 0\}$. In particular, it follows from (3.1) that $u(t) \geq T(t)u_0 > 0$ on $(0, T]$. It is not difficult by uniqueness to deduce that u stays positive on $(0, T_{\max})$.

• **Regularity of solutions:** If $u_0 \in L^r(\mathbb{R}^N) \cap C_0(\mathbb{R}^N)$, for $1 \leq r < \infty$, then by repeating the fixed point argument in the space

$$E_{T,r} := \{u \in L^\infty((0, T), C_0(\mathbb{R}^N) \cap L^r(\mathbb{R}^N)); \|u\|_1 \leq 2\|u_0\|_{L^\infty}, \|u\|_{\infty,r} \leq 2\|u_0\|_{L^r}\},$$

instead of E_T , where $\|\cdot\|_{\infty,r} := \|\cdot\|_{L^\infty((0,T),L^r(\mathbb{R}^N))}$, and by estimating $\|u^p\|_{L^r(\mathbb{R}^N)}$ by $\|u\|_{L^\infty(\mathbb{R}^N)}^{p-1}\|u\|_{L^r(\mathbb{R}^N)}$ in the contraction mapping argument, using (2.4), we obtain a unique solution in $E_{T,r}$; we conclude then that

$$u \in C([0, T_{\max}), C_0(\mathbb{R}^N) \cap L^r(\mathbb{R}^N)).$$

□

We say that u is a global solution if $T_{\max} = \infty$; when $T_{\max} < \infty$, u is said to blow up in a finite time and in this case we have $\|u(\cdot, t)\|_{L^\infty(\mathbb{R}^N)} \rightarrow \infty$ as $t \rightarrow T_{\max}$.

Remark. We note that classical or strong solution do not exist due to the singularity in time in the nonlinear term.

4 Blow-up of solutions

Now, we want to derive a blow-up result for Eq. (1.1). To do this, we show for later use that a mild solution is a weak solution.

Definition 2 (*Weak solution*)

Let $u_0 \in L^\infty_{Loc}(\mathbb{R}^N)$, $0 < \beta \leq 2$ and $T > 0$. We say that u is a weak solution of the problem (1.1) if $u \in L^p((0, T), L^\infty_{Loc}(\mathbb{R}^N))$ and verifies the equation

$$\begin{aligned} \int_{\Omega} u_0(x)\varphi(x, 0) + \int_0^T \int_{\Omega} J_{0|t}^\alpha(|u|^{p-1}u)(x, t)\varphi(x, t) &= \int_0^T \int_{\Omega} u(x, t)(-\Delta)^{\beta/2}\varphi(x, t) \\ &- \int_0^T \int_{\Omega} u(x, t)\varphi_t(x, t), \end{aligned} \quad (4.1)$$

for all $\varphi \in C^1([0, T], H^\beta(\mathbb{R}^N))$ such that $\Omega := \text{supp}\varphi$ is compact with $\varphi(\cdot, T) = 0$, where $\alpha := 1 - \gamma \in (0, 1)$.

Lemma 1

Consider $u_0 \in C_0(\mathbb{R}^N)$ and let $u \in C([0, T], C_0(\mathbb{R}^N))$ be a mild solution of (1.1), then u is a weak solution of (1.1), for all $0 < \beta \leq 2$ and all $T > 0$.

Proof. Let $T > 0$, $0 < \beta \leq 2$, $u_0 \in C_0(\mathbb{R}^N)$ and let $u \in C([0, T], C_0(\mathbb{R}^N))$ be a solution of (3.1). Given $\varphi \in C^1([0, T], H^\beta(\mathbb{R}^N))$ such that $\text{supp}\varphi =: \Omega$ is compact with $\varphi(\cdot, T) = 0$. Then after multiplying (3.1) by φ and integrating over \mathbb{R}^N , we have

$$\int_{\Omega} u(x, t)\varphi(x, t) = \int_{\Omega} T(t)u_0(x)\varphi(x, t) + \int_{\Omega} \left(\int_0^t T(t-s)J_{0|s}^\alpha(|u|^{p-1}u)(x, s) ds \right) \varphi(x, t).$$

So after differentiating in time, we obtain

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} u(x, t)\varphi(x, t) &= \int_{\Omega} \frac{d}{dt} (T(t)u_0(x)\varphi(x, t)) \\ &+ \int_{\Omega} \frac{d}{dt} \int_0^t T(t-s)J_{0|s}^\alpha(|u|^{p-1}u)(x, s) ds \varphi(x, t). \end{aligned} \quad (4.2)$$

Now, using (2.5) and a property of the semigroup $T(t)$ ([8, Chapter 3]), we have:

$$\begin{aligned} \int_{\Omega} \frac{d}{dt} (T(t)u_0(x)\varphi(x, t)) &= \int_{\Omega} A(T(t)u_0(x))\varphi(x, t) + \int_{\Omega} T(t)u_0(x)\varphi_t(x, t) \\ &= \int_{\Omega} T(t)u_0(x)A\varphi(x, t) + \int_{\Omega} T(t)u_0(x)\varphi_t(x, t), \end{aligned} \quad (4.3)$$

and

$$\begin{aligned} \int_{\Omega} \frac{d}{dt} \int_0^t T(t-s)f(x, s) ds \varphi(x, t) &= \int_{\Omega} f(x, t)\varphi(x, t) + \int_{\Omega} \int_0^t A(T(t-s)f(x, s)) ds \varphi(x, t) \\ &+ \int_{\Omega} \int_0^t T(t-s)f(x, s) ds \varphi_t(x, t) \\ &= \int_{\Omega} f(x, t)\varphi(x, t) + \int_{\Omega} \int_0^t T(t-s)f(x, s) ds A\varphi(x, t) \\ &+ \int_{\Omega} \int_0^t T(t-s)f(x, s) ds \varphi_t(x, t), \end{aligned} \quad (4.4)$$

where $f := J_{0|t}^{\alpha}(|u|^{p-1}u) \in C([0, T]; L^2(\Omega))$.

Thus, using (3.1), (4.3) and (4.4), we conclude that (4.2) imply

$$\frac{d}{dt} \int_{\Omega} u(x, t)\varphi(x, t) = \int_{\Omega} u(x, t)A\varphi(x, t) + \int_{\Omega} u(x, t)\varphi_t(x, t) + \int_{\Omega} f(x, t)\varphi(x, t).$$

Finally, we conclude by integrating in time over $[0, T]$ and use the fact that $\varphi(\cdot, T) = 0$. \square

Theorem 2 *Let $u_0 \in C_0(\mathbb{R}^N)$ be such that $u_0 \geq 0$ and $u_0 \not\equiv 0$. If*

$$p \leq 1 + \frac{\beta(2-\gamma)}{(N-\beta+\beta\gamma)_+} := p^* \quad \text{or} \quad p < \frac{1}{\gamma}, \quad (4.5)$$

for all $\beta \in (0, 2]$, then any mild solution to (1.1) blows-up in a finite time.

Note that in the case where $p = p^$ and $\beta \in (0, 2)$ we take $p > N/(N-\beta)$ with $N > \beta$.*

Proof. The proof is by contradiction. Suppose that u is a global mild solution to (1.1), then u is a solution of (1.1) in $C([0, T], C_0(\mathbb{R}^N))$ for all $T \gg 1$ such that $u(t) > 0$ for all $t \in [0, T]$.

Then, using Lemma 1, we have

$$\begin{aligned} \int_{\text{supp}\varphi} u_0(x)\varphi(x, 0) + \int_0^T \int_{\text{supp}\varphi} J_{0|t}^{\alpha}(u^p)(x, t)\varphi(x, t) &= \int_0^T \int_{\text{supp}\varphi} u(x, t)(-\Delta)^{\beta/2}\varphi(x, t) \\ &- \int_0^T \int_{\text{supp}\varphi} u(x, t)\varphi_t(x, t), \end{aligned}$$

for all test function $\varphi \in C^1([0, T], H^{\beta}(\mathbb{R}^N))$ such that $\text{supp}\varphi$ is compact with $\varphi(\cdot, T) = 0$, where $\alpha := 1 - \gamma \in (0, 1)$.

Now we take $\varphi(x, t) = D_{t|T}^\alpha(\tilde{\varphi}(x, t)) := D_{t|T}^\alpha\left((\varphi_1(x))^\ell \varphi_2(t)\right)$ with $\varphi_1(x) := \Phi(|x|/T^{1/\beta})$, $\varphi_2(t) := (1 - t/T)_+^\eta$, where $\ell \geq p/(p-1)$, $\eta \geq \max\{(\alpha p + 1)/(p-1); \alpha + 1\}$ and Φ a smooth nonnegative non-increasing function such that

$$\Phi(r) = \begin{cases} 1 & \text{if } 0 \leq r \leq 1, \\ 0 & \text{if } r \geq 2, \end{cases}$$

$0 \leq \Phi \leq 1$, $|\Phi'(r)| \leq C_1/r$, for all $r > 0$. Using (2.16), we then obtain

$$\begin{aligned} & \int_{\Omega} u_0(x) D_{t|T}^\alpha \tilde{\varphi}(x, 0) + \int_{\Omega_T} J_{0|t}^\alpha(u^p)(x, t) D_{t|T}^\alpha \tilde{\varphi}(x, t) \\ &= \int_{\Omega_T} u(x, t) (-\Delta)^{\beta/2} D_{t|T}^\alpha \tilde{\varphi}(x, t) - \int_{\Omega_T} u(x, t) D D_{t|T}^\alpha \tilde{\varphi}(x, t), \end{aligned} \quad (4.6)$$

where

$$\Omega_T := [0, T] \times \Omega \text{ for } \Omega = \left\{x \in \mathbb{R}^N ; |x| \leq 2T^{1/\beta}\right\}, \quad \int_{\Omega} = \int_{\Omega} dx \text{ and } \int_{\Omega_T} = \int_{\Omega_T} dx dt.$$

Furthermore, using (2.11) and (2.16) in the left hand side of (4.6), and (2.12) in the right hand side, we conclude that

$$\begin{aligned} & C T^{-\alpha} \int_{\Omega} u_0(x) \varphi_1^\ell(x) + \int_{\Omega_T} D_{0|t}^\alpha J_{0|t}^\alpha(u^p)(x, t) \tilde{\varphi}(x, t) \\ &= \int_{\Omega_T} u(x, t) (-\Delta)^{\beta/2} D_{t|T}^\alpha \tilde{\varphi}(x, t) + \int_{\Omega_T} u(x, t) D_{t|T}^{1+\alpha} \tilde{\varphi}(x, t). \end{aligned} \quad (4.7)$$

Moreover, using (2.13), we may write

$$\begin{aligned} & \int_{\Omega_T} u^p(x, t) \tilde{\varphi}(x, t) + C T^{-\alpha} \int_{\Omega} u_0(x) \varphi_1^\ell(x) \\ &= \int_{\Omega_T} u(x, t) (-\Delta)^{\beta/2} \varphi_1^\ell(x) D_{t|T}^\alpha \varphi_2(t) + \int_{\Omega_T} u(x, t) D_{t|T}^{1+\alpha} \tilde{\varphi}(x, t). \end{aligned} \quad (4.8)$$

So, Ju's inequality $(-\Delta)^{\beta/2}(\varphi_1^\ell) \leq \ell \varphi_1^{\ell-1} (-\Delta)^{\beta/2}(\varphi_1)$ (see the Appendix) allows us to write:

$$\begin{aligned} & \int_{\Omega_T} u^p(x, t) \tilde{\varphi}(x, t) + C T^{-\alpha} \int_{\Omega} u_0(x) \varphi_1^\ell(x) \\ & \leq C \int_{\Omega_T} u(x, t) \varphi_1^{\ell-1}(x) \left| (-\Delta)^{\beta/2} \varphi_1(x) D_{t|T}^\alpha \varphi_2(t) \right| \\ & \quad + \int_{\Omega_T} u(x, t) \varphi_1^\ell(x) \left| D_{t|T}^{1+\alpha} \varphi_2(t) \right| \\ & = C \int_{\Omega_T} u(x, t) \tilde{\varphi}^{1/p} \tilde{\varphi}^{-1/p} \varphi_1^{\ell-1}(x) \left| (-\Delta)^{\beta/2} \varphi_1(x) D_{t|T}^\alpha \varphi_2(t) \right| \\ & \quad + \int_{\Omega_T} u(x, t) \tilde{\varphi}^{1/p} \tilde{\varphi}^{-1/p} \varphi_1^\ell(x) \left| D_{t|T}^{1+\alpha} \varphi_2(t) \right| \end{aligned} \quad (4.9)$$

Therefore, using Young's inequality

$$ab \leq \frac{1}{2p} a^p + \frac{2^{\tilde{p}-1}}{\tilde{p}} b^{\tilde{p}} \quad \text{where } p\tilde{p} = p + \tilde{p}, \quad a > 0, b > 0, \quad p > 1, \tilde{p} > 1, \quad (4.10)$$

with

$$\begin{cases} a = u(x, t) \tilde{\varphi}^{1/p}, \\ b = \tilde{\varphi}^{-1/p} \varphi_1^{\ell-1}(x) \left| (-\Delta)^{\beta/2} \varphi_1(x) D_{t|T}^\alpha \varphi_2(t) \right|, \end{cases}$$

in the first integral of the right hand side of (4.9), and with

$$\begin{cases} a = u(x, t) \tilde{\varphi}^{1/p}, \\ b = \tilde{\varphi}^{-1/p} \varphi_1^\ell(x) \left| D_{t|T}^{1+\alpha} \varphi_2(t) \right|, \end{cases}$$

in the second integral of the right hand side of (4.9), we obtain

$$\begin{aligned} & \left(1 - \frac{1}{p}\right) \int_{\Omega_T} u^p(x, t) \tilde{\varphi}(x, t) \\ & \leq C \int_{\Omega_T} (\varphi_1(x))^{\ell-\tilde{p}} (\varphi_2(t))^{-\frac{1}{p-1}} \left| (-\Delta_x)^{\beta/2} \varphi_1(x) D_{t|T}^\alpha \varphi_2(t) \right|^{\tilde{p}} \\ & \quad + C \int_{\Omega_T} (\varphi_1(x))^\ell (\varphi_2(t))^{-\frac{1}{p-1}} \left| D_{t|T}^{1+\alpha} \varphi_2(t) \right|^{\tilde{p}}, \end{aligned} \quad (4.11)$$

as $u_0 \geq 0$. At this stage, we introduce the scaled variables: $\tau = T^{-1}t$, $\xi = T^{-1/\beta}x$, use formula (2.14) and (2.15) in the right hand-side of (4.11), to obtain:

$$\int_{\Omega_T} u^p(x, t) \tilde{\varphi}(x, t) \leq C T^{-\delta}, \quad (4.12)$$

where $\delta := (1 + \alpha)\tilde{p} - 1 - (N/\beta)$, $C = C(|\Omega_1|, |\Omega_2|)$, ($|\Omega_i|$ stands for the measure of Ω_i , for $i = 1, 2$), with

$$\Omega_1 := \{\xi \in \mathbb{R}^N ; |\xi| \leq 2\} \quad , \quad \Omega_2 := \{\tau \geq 0 ; \tau \leq 1\}.$$

Now, noting that, as

$$p \leq p^* \quad \text{or} \quad p < \frac{1}{\gamma} \quad \iff \quad \delta \geq 0 \quad \text{or} \quad p < \frac{1}{\gamma}, \quad (4.13)$$

we have to distinguish three cases:

- The case $p < p^*$ ($\delta > 0$): we pass to the limit in (4.12), as T goes to ∞ ; we get

$$\lim_{T \rightarrow \infty} \int_0^T \int_{|x| \leq 2T^{1/\beta}} u^p(x, t) \tilde{\varphi}(x, t) dx dt = 0.$$

Using the Lebesgue dominated convergence theorem, the continuity in time and space of u and the fact that $\tilde{\varphi}(x, t) \rightarrow 1$ as $T \rightarrow \infty$, we infer that

$$\int_0^\infty \int_{\mathbb{R}^N} u^p(x, t) dx dt = 0 \quad \implies \quad u \equiv 0.$$

Contradiction.

• The case $p = p^*$ ($\delta = 0$): using inequality (4.12) with $T \rightarrow \infty$, taking into account the fact that $p = p^*$, we have on one hand

$$u \in L^p((0, \infty), L^p(\mathbb{R}^N)); \quad (4.14)$$

on the other hand, we repeat the same calculation as above by taking this time $\varphi_1(x) := \Phi(|x|/(B^{-1/\beta}T^{1/\beta}))$, where $1 \leq B < T$ is large enough such that when $T \rightarrow \infty$ we don't have $B \rightarrow \infty$ at the same time, so we arrive at

$$\int_{\Sigma_T} u^p(x, t) \tilde{\varphi}(x, t) \leq C B^{-N/\beta} + C B^{-N/\beta+\tilde{p}}, \quad (4.15)$$

thanks to the following rescaling: $\tau = T^{-1}t$, $\xi = (T/B)^{-1/\beta}x$, where

$$\Sigma_T := [0, T] \times \left\{ x \in \mathbb{R}^N ; |x| \leq 2B^{-1/\beta}T^{1/\beta} \right\} \quad \text{and} \quad \int_{\Sigma_T} = \int_{\Sigma_T} dx dt.$$

Thus, using $p > N/(N - \beta)$ and taking the limit when $T \rightarrow \infty$ and then $B \rightarrow \infty$, we get:

$$\int_0^\infty \int_{\mathbb{R}^N} u^p(x, t) dx dt = 0 \quad \implies \quad u \equiv 0,$$

which is a contradiction.

Note that, in the case $\beta = 2$ it is not necessary to take the condition $p > N/(N - \beta)$ with $N > \beta$. Indeed, from (4.9) with the new function φ_1 , we may write

$$\begin{aligned} & \int_{\Sigma_T} u^p(x, t) \tilde{\varphi}(x, t) \\ & \leq C \int_{\Sigma_T} u(x, t) \tilde{\varphi}^{1/p} \tilde{\varphi}^{-1/p} (\varphi_1(x))^\ell \left| D_{t|T}^{1+\alpha} \varphi_2(t) \right| \\ & + C \int_{\Delta_B} u(x, t) \tilde{\varphi}^{1/p} \tilde{\varphi}^{-1/p} (\varphi_1(x))^{\ell-1} \left| \Delta_x \varphi_1(x) D_{t|T}^\alpha \varphi_2(t) \right|, \end{aligned} \quad (4.16)$$

where

$$\Delta_B = [0, T] \times \left\{ x \in \mathbb{R}^N ; B^{-1/2}T^{1/2} \leq |x| \leq 2B^{-1/2}T^{1/2} \right\} \subset \Sigma_T \quad \text{and} \quad \int_{\Delta_B} = \int_{\Delta_B} dx dt.$$

Moreover, using the following Young's inequality

$$ab \leq \frac{1}{p} a^p + \frac{1}{\tilde{p}} b^{\tilde{p}} \quad \text{where } p\tilde{p} = p + \tilde{p}, \quad a > 0, b > 0, \quad p > 1, \tilde{p} > 1, \quad (4.17)$$

with

$$\begin{cases} a = u(x, t) \tilde{\varphi}^{1/p}, \\ b = \tilde{\varphi}^{-1/p} (\varphi_1(x))^\ell \left| D_{t|T}^{1+\alpha} \varphi_2(t) \right|, \end{cases}$$

in the first integral of the right hand side of (4.16), and using Hölder's inequality

$$\int_{\Delta_B} ab \leq \left(\int_{\Delta_B} a^p \right)^{1/p} \left(\int_{\Delta_B} b^{\tilde{p}} \right)^{1/\tilde{p}}, \quad \text{where } p\tilde{p} = p + \tilde{p}, \quad a > 0, b > 0, \quad p > 1, \tilde{p} > 1,$$

with

$$\begin{cases} a = u(x, t) \tilde{\varphi}^{1/p}, \\ b = \tilde{\varphi}^{-1/p} (\varphi_1(x))^{\ell-1} \left| \Delta_x \varphi_1(x) D_{t|T}^\alpha \varphi_2(t) \right|, \end{cases}$$

in the second integral of the right hand side of (4.16), we obtain

$$\begin{aligned} & \left(1 - \frac{1}{p}\right) \int_{\Sigma_T} u^p(x, t) \tilde{\varphi}(x, t) \\ & \leq C \int_{\Sigma_T} (\varphi_1(x))^\ell (\varphi_2(t))^{-\frac{1}{p-1}} \left| D_{t|T}^{1+\alpha} \varphi_2(t) \right|^{\tilde{p}} \\ & \quad + C \left(\int_{\Delta_B} u^p \tilde{\varphi} \right)^{1/p} \left(\int_{\Delta_B} (\varphi_1(x))^{\ell-\tilde{p}} (\varphi_2(t))^{-\frac{1}{p-1}} \left| \Delta_x \varphi_1(x) D_{t|T}^\alpha \varphi_2(t) \right|^{\tilde{p}} \right)^{1/\tilde{p}}. \end{aligned} \quad (4.18)$$

Taking account of the scaled variables: $\tau = T^{-1}t$, $\xi = (T/B)^{-1/2}x$, and the fact that $\delta = 0$, we get

$$\int_{\Sigma_T} u^p(x, t) \tilde{\varphi}(x, t) \leq C B^{-N/2} + C B^{-\frac{N}{2\tilde{p}}+1} \left(\int_{\Delta_B} u^p \tilde{\varphi} \right)^{1/p}. \quad (4.19)$$

Now, as

$$\lim_{T \rightarrow \infty} \left(\int_{\Delta_B} u^p \tilde{\varphi} \right)^{1/p} = 0 \quad (\text{from (4.14)}),$$

then, passing to the limit in (4.19), as $T \rightarrow \infty$, we get

$$\int_0^\infty \int_{\mathbb{R}^N} u^p(x, t) dx dt \leq C B^{-N/2}.$$

We conclude that $u \equiv 0$ by taking the limit when B goes to infinity, contradiction.

• For the case $p < (1/\gamma)$, we repeat the same argument as in the case $p < p^*$ by choosing the test function as follows: $\varphi(x, t) = D_{t|T}^\alpha \bar{\varphi}(x, t) := D_{t|T}^\alpha (\varphi_3^\ell(x) \varphi_4(t))$ where $\varphi_3(x) = \Phi(|x|/R)$, $\varphi_4(t) = (1 - t/T)_+^\eta$ and $R \in (0, T)$ large enough such that in the case when $T \rightarrow \infty$ we don't have $R \rightarrow \infty$ at the same time, with the same functions Φ as above. We then obtain

$$\int_{\mathcal{C}_T} u^p(x, t) \bar{\varphi}(x, t) + C T^{-\alpha} \int_{\mathcal{C}} (\varphi_3(x))^\ell u_0(x)$$

$$\begin{aligned}
&\leq C \int_{\mathcal{C}_T} u(x, t) \bar{\varphi}^{1/p} \bar{\varphi}^{-1/p} (\varphi_3(x))^\ell \left| D_{t|T}^{1+\alpha} \varphi_4(t) \right| \\
&+ C \int_{\mathcal{C}_T} u(x, t) \bar{\varphi}^{1/p} \bar{\varphi}^{-1/p} (\varphi_3(x))^{\ell-1} \left| (-\Delta_x)^{\beta/2} \varphi_3(x) D_{t|T}^\alpha \varphi_4(t) \right|, \quad (4.20)
\end{aligned}$$

where

$$\mathcal{C}_T := [0, T] \times \mathcal{C} \quad \text{for } \mathcal{C} := \{x \in \mathbb{R}^N ; |x| \leq 2R\}, \quad \int_{\mathcal{C}} = \int_{\mathcal{C}} dx \quad \text{and} \quad \int_{\mathcal{C}_T} = \int_{\mathcal{C}_T} dx dt.$$

Now, by Young's inequality (4.10), with

$$\begin{cases} a = u(x, t) \bar{\varphi}^{1/p}, \\ b = \bar{\varphi}^{-1/p} (\varphi_3(x))^\ell \left| D_{t|T}^{1+\alpha} \varphi_4(t) \right|, \end{cases}$$

in the first integral of the right hand side of (4.20) and with

$$\begin{cases} a = u(x, t) \bar{\varphi}^{1/p}, \\ b = \bar{\varphi}^{-1/p} (\varphi_3(x))^{\ell-1} \left| (-\Delta_x)^{\beta/2} \varphi_3(x) D_{t|T}^\alpha \varphi_4(t) \right|, \end{cases}$$

in the second integral of the right hand side of (4.20) and using the positivity of u_0 , we get

$$\begin{aligned}
\left(1 - \frac{1}{p}\right) \int_{\mathcal{C}_T} u^p(x, t) \bar{\varphi}(x, t) &\leq C \int_{\mathcal{C}_T} (\varphi_3(x))^\ell (\varphi_4(t))^{-\frac{1}{p-1}} \left| D_{t|T}^{1+\alpha} \varphi_4(t) \right|^{\tilde{p}} \\
&+ C \int_{\mathcal{C}_T} (\varphi_3(x))^{\ell-\tilde{p}} (\varphi_4(t))^{-\frac{1}{p-1}} \left| (-\Delta_x)^{\beta/2} \varphi_3 D_{t|T}^\alpha \varphi_4 \right|^{\tilde{p}}.
\end{aligned}$$

Then, the new variables $\xi = R^{-1}x$, $\tau = T^{-1}t$ and (2.14) – (2.15) allow us to write the estimate

$$\int_{\mathcal{C}_T} u^p(x, t) \bar{\varphi}(x, t) dx dt \leq C T^{1-(1+\alpha)\tilde{p}} R^N + C T^{1-\alpha\tilde{p}} R^{N-\beta\tilde{p}}.$$

Taking the limit as $T \rightarrow \infty$, we infer, as $p < 1/\gamma \iff 1 - \alpha\tilde{p} < 0$, that

$$\int_0^\infty \int_{\mathcal{C}} u(x, t)^p (\varphi_3(x))^\ell dx dt = 0.$$

Finally, by taking $R \rightarrow \infty$, we get a contradiction as $u(x, t) > 0$ for all $x \in \mathbb{R}^N$, $t > 0$. \square

Remarks.

(1) We can extend our analysis to the equation

$$u_t = -(-\Delta)^{\beta/2} u + \frac{1}{\Gamma(1-\gamma)} \int_0^t \frac{\psi(x, s) |u(s)|^{p-1} u(s)}{(t-s)^\gamma} ds, \quad x \in \mathbb{R}^N, \quad (4.21)$$

where $p > 1$, $\beta \in (0, 2]$, $0 < \gamma < 1$ and $\psi \in L_{Loc}^1(\mathbb{R}^N \times (0, \infty))$, $\psi(\cdot, t) > 0$ for all $t > 0$,

$$\begin{cases} \psi(B^{-1/\beta} T^{1/\beta} \xi, T\tau) \geq C > 0 & \text{if } p \leq p^* \\ \psi(R\xi, T\tau) \geq C > 0 & \text{if } p < 1/\gamma, \end{cases}$$

for any $0 < R, B < T$, $\tau \in [0, 1]$ and $\xi \in [0, 2]$.

(2) If we take $\beta = 2$ and $v(x, t) = (\Gamma(1 - \gamma))^{(1-\gamma)/(p-1)} u(\Gamma(1 - \gamma)^{1/2}x, \Gamma(1 - \gamma)t)$ where u is a solution of (1.1), we recover the result in [10] as a particular case.

(3) In Theorem 2, we use precisely the weak solution, but in this case we say that we have a nonexistence of global weak solutions. Therefore, to obtain blow-up results, we use the mild solution and the fact that: either $T_{\max} = \infty$ or else $T_{\max} < \infty$ and $\|u\|_{L^\infty((0,t) \times \mathbb{R}^N)} \rightarrow \infty$ as $t \rightarrow T_{\max}$ [3].

(4) We can take the nonlocal porous-medium spatio-fractional problem which is our real motivation to extend the results of [10] :

$$\begin{cases} u_t + (-\Delta)^{\beta/2} |u|^{m-1} u = \frac{1}{\Gamma(1-\gamma)} \int_0^t (t-s)^{-\gamma} |u|^{p-1} u(s) ds & x \in \mathbb{R}^N, t > 0, \\ u(x, 0) = u_0(x) & x \in \mathbb{R}^N, \end{cases}$$

where $\beta \in (0, 2]$, $0 < \gamma < 1$, $1 \leq m < p$, $u_0 \geq 0$ and $u_0 \not\equiv 0$.
The threshold on p will be

$$p \leq 1 + \frac{(2-\gamma)(N(m-1) + \beta)}{(N - \beta + \beta\gamma)_+} \quad \text{or} \quad p < \frac{m}{\gamma}.$$

5 Blow-up Rate

In this section, we present the blow-up rate for the blowing-up solutions to the parabolic problem (1.1) in the case $\beta = 2$.

We take the solution of (1.1) with an initial condition satisfying

$$u_0 \in C_0(\mathbb{R}^N) \cap L^2(\mathbb{R}^N), \quad u(\cdot, 0) = u_0 \geq 0, \quad u_0 \not\equiv 0. \quad (5.1)$$

The following lemma will be used in the proof of Theorem 3 below.

Lemma 2 *Let φ be a nonnegative classical solution of*

$$\varphi_t = \Delta \varphi + J_{-\infty|t}^{1-\gamma}(\varphi^p) \quad \text{in } \mathbb{R}^N \times \mathbb{R}, \quad (5.2)$$

where $\gamma \in (0, 1)$, $p > 1$ and

$$J_{-\infty|t}^{1-\gamma}(\varphi^p)(t) := \frac{1}{\Gamma(1-\gamma)} \int_{-\infty}^t (t-s)^{-\gamma} \varphi^p(s) ds.$$

Then $\varphi \equiv 0$ whenever

$$p \leq p^* \quad \text{or} \quad p < \frac{1}{\gamma}. \quad (5.3)$$

Proof. We repeat the same computations as in Theorem 2 with $(1 - t^2/T^2)_+^\eta$ instead of $(1 - t/T)_+^\eta$ for $\eta \gg 1$, using (2.17) – (2.19) and taking account of the following inequality

$$J_{-\infty|t}^{1-\gamma}(\varphi^p) \geq J_{-T|t}^{1-\gamma}(\varphi^p).$$

Moreover, we take $\varphi_1^{\ell/p} \varphi_1^{-\ell/p}$ instead of $\tilde{\varphi}^{1/p} \tilde{\varphi}^{-1/p}$ (resp. $\bar{\varphi}^{1/p} \bar{\varphi}^{-1/p}$) in (4.9), (4.16) (resp. in (4.20)) for $\ell \gg 1$ to use the Young and Hölder's inequality.

Note that here, we use rather the ε -Young inequality

$$ab \leq \frac{\varepsilon}{2} a^p + C(\varepsilon) b^{\tilde{p}},$$

for $0 < \varepsilon < 1$. □

Theorem 3 *Let u_0 satisfies (5.1). For $p \leq p^*$ or $p < (1/\gamma)$, let $\alpha_1 := (2 - \gamma)/(p - 1)$ and let u be the blowing-up mild solution of (1.1) in a finite time $T_{\max} := T^*$. Then there exist two constants $c, C > 0$ such that*

$$c(T^* - t)^{-\alpha_1} \leq \sup_{\mathbb{R}^N} u(\cdot, t) \leq C(T^* - t)^{-\alpha_1}, \quad t \in (0, T^*). \quad (5.4)$$

Proof. The proof is split into two parts:

- The upper blow-up rate estimate. Let

$$M(t) := \sup_{\mathbb{R}^N \times (0, t]} u, \quad t \in (0, T^*).$$

Clearly, M is positive, continuous and nondecreasing in $(0, T^*)$. As $\lim_{t \rightarrow T^*} M(t) = \infty$, then for all $t_0 \in (0, T^*)$, we can define

$$t_0^+ := t^+(t_0) := \max\{t \in (t_0, T^*) : M(t) = 2M(t_0)\}.$$

Choose $A \geq 1$ and let

$$\lambda(t_0) := \left(\frac{1}{2A} M(t_0) \right)^{-1/(2\alpha_1)}. \quad (5.5)$$

we claim that

$$\lambda^{-2}(t_0)(t_0^+ - t_0) \leq D, \quad t_0 \in \left(\frac{T^*}{2}, T^* \right), \quad (5.6)$$

where $D > 0$ is a positive constant which does not depend on t_0 .

We proceed by contradiction. If (5.6) were false, then there would exist a sequence $t_n \rightarrow T^*$ such that

$$\lambda_n^{-2}(t_n^+ - t_n) \rightarrow \infty,$$

where $\lambda_n = \lambda(t_n)$ and $t_n^+ = t^+(t_n)$. For each t_n choose

$$(\hat{x}_n, \hat{t}_n) \in \mathbb{R}^N \times (0, t_n] \quad \text{such that} \quad u(\hat{x}_n, \hat{t}_n) \geq \frac{1}{2} M(t_n). \quad (5.7)$$

Obviously, $M(t_n) \rightarrow \infty$; hence, $\widehat{t}_n \rightarrow T^*$. Next, rescale the function u as

$$\varphi^{\lambda_n}(y, s) := \lambda_n^{2\alpha_1} u(\lambda_n y + \widehat{x}_n, \lambda_n^2 s + \widehat{t}_n), \quad (y, s) \in \mathbb{R}^N \times I_n(T^*), \quad (5.8)$$

where $I_n(t) := (-\lambda_n^{-2}\widehat{t}_n, \lambda_n^{-2}(t - \widehat{t}_n))$ for all $t > 0$. Then φ^{λ_n} is a mild solution of

$$\varphi_s = \Delta\varphi + J_{-\lambda_n^{-2}\widehat{t}_n|s}^\alpha(\varphi^p) \quad \text{in } \mathbb{R}^N \times I_n(T^*), \quad (5.9)$$

i.e., for $G(t) := G(x, t) := (4\pi t)^{-N/2} e^{-|x|^2/4t}$ and $*$ being the space convolution, we have

$$\varphi^{\lambda_n}(s) = G(s + \lambda_n^{-2}\widehat{t}_n) * \varphi^{\lambda_n}(-\lambda_n^{-2}\widehat{t}_n) + \int_{-\lambda_n^{-2}\widehat{t}_n}^s G(s - \sigma) * J_{-\lambda_n^{-2}\widehat{t}_n|\sigma}^\alpha((\varphi^{\lambda_n})^p) d\sigma \quad (5.10)$$

in $\mathbb{R}^N \times I_n(T^*)$; whereupon, as $\varphi^{\lambda_n}(0, 0) \geq A$,

$$0 \leq \varphi^{\lambda_n} \leq \lambda_n^{2\alpha_1} M(t_n^+) = \lambda_n^{2\alpha_1} 2M(t_n) = 4A \quad \text{in } \mathbb{R}^N \times I_n(t_n^+),$$

thanks to (5.5) and the definition of t_n^+ .

Moreover, as

$$\varphi^{\lambda_n} \in C([-\lambda_n^{-2}\widehat{t}_n, T], C_0(\mathbb{R}^N) \cap L^2(\mathbb{R}^N)) \quad \text{for all } T \in I_n(T^*),$$

so, as in Lemma 1, φ^{λ_n} is a weak solution of (5.9).

On the other hand, if we write φ^{λ_n} as $\varphi^{\lambda_n}(s) = v(s) + w(s)$ for all $s \in I_n(T^*)$, where

$$v(s) := G(s + \lambda_n^{-2}\widehat{t}_n) * \varphi^{\lambda_n}(-\lambda_n^{-2}\widehat{t}_n) \quad \text{and} \quad w(s) := \int_{-\lambda_n^{-2}\widehat{t}_n}^s G(s - \sigma) * J_{-\lambda_n^{-2}\widehat{t}_n|\sigma}^\alpha((\varphi^{\lambda_n})^p) d\sigma,$$

we have, see [8, Chapter 3], for $T \in I_n(T^*)$

$$v \in C((-\lambda_n^{-2}\widehat{t}_n, T); H^2(\mathbb{R}^N)) \cap C^1((-\lambda_n^{-2}\widehat{t}_n, T); L^2(\mathbb{R}^N)) \subset L^2((-\lambda_n^{-2}\widehat{t}_n, T); H^1(\mathbb{R}^N))$$

and, using the fact that $f(s) := J_{-\lambda_n^{-2}\widehat{t}_n|s}^\alpha((\varphi^{\lambda_n})^p) \in L^2((-\lambda_n^{-2}\widehat{t}_n, T); L^2(\mathbb{R}^N))$ and the maximal regularity theory, we have

$$w \in W^{1,2}((-\lambda_n^{-2}\widehat{t}_n, T); L^2(\mathbb{R}^N)) \cap L^2((-\lambda_n^{-2}\widehat{t}_n, T); H^2(\mathbb{R}^N)) \subset L^2((-\lambda_n^{-2}\widehat{t}_n, T); H^1(\mathbb{R}^N)).$$

It follows that

$$\varphi^{\lambda_n} \in C([-\lambda_n^{-2}\widehat{t}_n, T], L^2(\mathbb{R}^N)) \cap L^2((-\lambda_n^{-2}\widehat{t}_n, T), W^{1,2}(\mathbb{R}^N));$$

so from the parabolic interior regularity theory (cf. [23, Theorem 10.1 p. 204]) there is $\mu \in (0, 1)$ such that the sequence φ^{λ_n} is bounded in the $C_{\text{loc}}^{\mu, \mu/2}(\mathbb{R}^N \times \mathbb{R})$ -norm by a constant that does not depend on n , where $C_{\text{loc}}^{\mu, \mu/2}(\mathbb{R}^N \times \mathbb{R})$ is the locally Hölder space defined in [23]. Similar uniform estimates for $J_{-\lambda_n^{-2}\widehat{t}_n|s}^\alpha(\varphi^p)$ follow if μ is sufficiently small. The parabolic interior Schauder's estimates (see [22, Th. 8.11.1 p. 130]), using the existence theorems in

Hölder's space, imply now that the $C_{loc}^{2+\mu, 1+\mu/2}(\mathbb{R}^N \times \mathbb{R})$ -norm of φ^{λ_n} is uniformly bounded. Hence, we obtain a subsequence converging in $C_{loc}^{2+\mu, 1+\mu/2}(\mathbb{R}^N \times \mathbb{R})$ to a solution φ of

$$\varphi_s = \Delta\varphi + J_{-\infty|s}^\alpha(\varphi^p) \quad \text{in } \mathbb{R}^N \times (-\infty, +\infty),$$

such that $\varphi(0, 0) \geq A$ and $0 \leq \varphi \leq 4A$ in $\mathbb{R}^N \times \mathbb{R}$. Wherefrom, using Lemma 2, we infer that $\varphi \equiv 0$ in $\mathbb{R}^N \times (-\infty, +\infty)$. Contradiction with the fact that $\varphi(0, 0) \geq A > 1$. This proves (5.6).

Next we use an idea from Hu [16]. From (5.5) and (5.6) it follows that

$$(t_0^+ - t_0) \leq D(2A)^{1/\alpha_1} M(t_0)^{-1/\alpha_1} \quad \text{for any } t_0 \in \left(\frac{T^*}{2}, T^*\right).$$

Fix $t_0 \in (T^*/2, T^*)$ and denote $t_1 = t_0^+, t_2 = t_1^+, t_3 = t_2^+, \dots$. Then

$$\begin{aligned} t_{j+1} - t_j &\leq D(2A)^{1/\alpha_1} M(t_j)^{-1/\alpha_1}, \\ M(t_{j+1}) &= 2M(t_j), \end{aligned}$$

$j = 0, 1, 2, \dots$. Consequently,

$$\begin{aligned} T^* - t_0 &= \sum_{j=0}^{\infty} (t_{j+1} - t_j) \leq D(2A)^{1/\alpha_1} \sum_{j=0}^{\infty} M(t_j)^{-1/\alpha_1} \\ &= D(2A)^{1/\alpha_1} M(t_0)^{-1/\alpha_1} \sum_{j=0}^{\infty} 2^{-j/\alpha_1}. \end{aligned}$$

Finally, we conclude that

$$u(x, t_0) \leq M(t_0) \leq C(T^* - t_0)^{-\alpha_1}, \quad \forall t_0 \in (0, T^*)$$

where

$$C = 2A \left(D \sum_{j=0}^{\infty} 2^{-j/\alpha_1} \right)^{\alpha_1};$$

consequently

$$\sup_{\mathbb{R}^N} u(\cdot, t) \leq C(T^* - t)^{-\alpha_1}, \quad \forall t \in (0, T^*).$$

• The lower blow-up rate estimate. If we repeat the proof of the local existence of Theorem 1, by taking $\|u\|_1 \leq \theta$ instead of $\|u\|_1 \leq 2\|u_0\|_\infty$ in the space E_T for all positive constant $\theta > 0$ and all $0 < t < T$, then the condition (3.2) of T will be:

$$\|u_0\|_\infty + CT^{2-\gamma}\theta^p \leq \theta, \quad (5.11)$$

and then, like before, we infer that $\|u(t)\|_\infty \leq \theta$ for (almost) all $0 < t < T$. Consequently, if $\|u_0\|_\infty + Ct^{2-\gamma}\theta^p \leq \theta$, then $\|u(t)\|_\infty \leq \theta$. Applying this to any point in the trajectory, we see that if $0 \leq s < t$ and

$$(t - s)^{2-\gamma} \leq \frac{\theta - \|u(s)\|_\infty}{C\theta^p}, \quad (5.12)$$

then $\|u(t)\|_\infty \leq \theta$, for all $0 < t < T$.

Moreover, if $0 \leq s < T^*$ and $\|u(s)\|_\infty < \theta$, then:

$$(T^* - s)^{2-\gamma} > \frac{\theta - \|u(s)\|_\infty}{C\theta^p}. \quad (5.13)$$

Indeed, arguing by contradiction and assume that for some $\theta > \|u(s)\|_\infty$ and all $t \in (s, T^*)$ we have

$$(t - s)^{2-\gamma} \leq \frac{\theta - \|u(s)\|_\infty}{C\theta^p},$$

then, using (5.12), we infer that $\|u(t)\|_\infty \leq \theta$ for all $t \in (s, T^*)$; this contradicts the fact that $\|u(t)\|_\infty \rightarrow \infty$ as $t \rightarrow T^*$.

Next, for example, by setting $\theta = 2\|u(s)\|_\infty$ in (5.13), we see that for $0 < s < T^*$ we have:

$$(T^* - s)^{2-\gamma} > C'\|u(s)\|_\infty^{1-p},$$

and by the positivity and the continuity of u we get

$$c(T^* - s)^{-\alpha_1} < \sup_{x \in \mathbb{R}^N} u(x, s), \quad \forall s \in (0, T^*). \quad (5.14)$$

□

6 Global existence

In this section, we prove the existence of global solutions of (1.1) with initial data small enough. We give a similar proof as that in [10] just for the seek of completeness. In the following, we use the notation $p_{sc} := N(p-1)/\beta(2-\gamma)$. As $p^* > 1 + \beta(2-\gamma)/N$, we note that $p > p^* \Rightarrow p_{sc} > 1$.

Theorem 4 *Let $u_0 \in C_0(\mathbb{R}^N) \cap L^{p_{sc}}(\mathbb{R}^N)$ and $0 < \beta \leq 2$. If*

$$p > \max\left\{\frac{1}{\gamma}; p^*\right\}, \quad (6.1)$$

and $\|u_0\|_{L^{p_{sc}}}$ is sufficiently small, then u exists globally.

Note that we can take $|u_0(x)| \leq C|x|^{-\beta(2-\gamma)/(p-1)}$ instead of $u_0 \in L^{p_{sc}}(\mathbb{R}^N)$.

Proof. As $p > (1/\gamma)$, then we have the possibility to take a positive constant $q > 0$ so that:

$$\frac{2-\gamma}{p-1} - \frac{1}{p} < \frac{N}{\beta q} < \frac{1}{p-1}, \quad q \geq p. \quad (6.2)$$

It follows, using (6.1), that

$$q > \frac{N(p-1)}{\beta} > p_{sc} > 1. \quad (6.3)$$

Let

$$b := \frac{N}{\beta p_{sc}} - \frac{N}{\beta q} = \frac{2-\gamma}{p-1} - \frac{N}{\beta q}. \quad (6.4)$$

Then, using (6.2) – (6.4), we conclude that

$$b > \frac{1-\gamma}{p-1} > 0, \quad pb < 1, \quad \frac{N(p-1)}{\beta q} + (p-1)b + \gamma = 2. \quad (6.5)$$

As $u_0 \in L^{p_{sc}}$, using (2.4) and (6.4), we get, for all $t > 0$,

$$\sup_{t>0} t^b \|e^{-t(-\Delta)^{\beta/2}} u_0\|_{L^q} \leq C \|u_0\|_{L^{p_{sc}}} = \eta < \infty. \quad (6.6)$$

Set

$$\Xi := \left\{ u \in L^\infty((0, \infty), L^q(\mathbb{R}^N)); \sup_{t>0} t^b \|u(t)\|_{L^q} \leq \delta \right\}, \quad (6.7)$$

where $\delta > 0$ is to be chosen sufficiently small. If we define

$$d_\Xi(u, v) := \sup_{t>0} t^b \|u(t) - v(t)\|_{L^q}, \quad \forall u, v \in \Xi, \quad (6.8)$$

then (Ξ, d) is a complete metric space. Given $u \in \Xi$, let's set:

$$\Phi(u)(t) := e^{-t(-\Delta)^{\beta/2}} u_0 + \frac{1}{\Gamma(1-\gamma)} \int_0^t e^{-(t-s)(-\Delta)^{\beta/2}} \int_0^s (s-\sigma)^{-\gamma} |u|^{p-1} u(\sigma) \, d\sigma \, ds, \quad (6.9)$$

for all $t \geq 0$. We have by (2.4), (6.6) and (6.7)

$$\begin{aligned} t^b \|\Phi(u)(t)\|_{L^q} &\leq \eta + Ct^b \int_0^t (t-s)^{-\frac{N}{\beta}(\frac{p}{q}-\frac{1}{q})} \int_0^s (s-\sigma)^{-\gamma} \|u^p(\sigma)\|_{L^{\frac{q}{p}}} \, d\sigma \, ds \\ &\leq \eta + C\delta^p t^b \int_0^t \int_0^s (t-s)^{-\frac{N(p-1)}{\beta q}} (s-\sigma)^{-\gamma} \sigma^{-bp} \, d\sigma \, ds. \end{aligned} \quad (6.10)$$

Next, using (6.2) and $pb < 1$, we get

$$\begin{aligned} \int_0^t \int_0^s \frac{(t-s)^{-\frac{N}{\beta q}(p-1)}}{(s-\sigma)^\gamma} \sigma^{-bp} \, d\sigma \, ds &= \left(\int_0^1 (1-\sigma)^{-\gamma} \sigma^{-bp} \, d\sigma \right) \int_0^t \frac{(t-s)^{-\frac{N(p-1)}{\beta q}}}{s^{bp+\gamma-1}} \, ds \\ &= Ct^{-\frac{N(p-1)}{\beta q}-bp-\gamma+2} = Ct^{-b}, \end{aligned} \quad (6.11)$$

for all $t \geq 0$. So, we deduce from (6.10) – (6.11) that

$$t^b \|\Phi(u)(t)\|_{L^q} \leq \eta + C\delta^p. \quad (6.12)$$

Therefore, if η and δ are chosen small enough so that $\eta + C\delta^p \leq \delta$, we see that $\Phi : \Xi \rightarrow \Xi$. Similar calculations show that (assuming η and δ small enough) Φ is a strict contraction, so it has a unique fixed point $u \in \Xi$ which is a solution of (1.1). Now, we show that $u \in C([0, \infty), C_0(\mathbb{R}^N))$.

First, we show that $u \in C([0, T], C_0(\mathbb{R}^N))$ if $T > 0$ is sufficiently small. Indeed, note that the above argument shows uniqueness in Ξ_T , where, for any $T > 0$,

$$\Xi_T := \left\{ u \in L^\infty((0, T), L^q(\mathbb{R}^N)); \sup_{0<t<T} t^b \|u(t)\|_{L^q} \leq \delta \right\}.$$

Let \tilde{u} be the local solution of (1.1) constructed in Theorem 1. Since $u_0 \in C_0(\mathbb{R}^N) \cap L^{p_{sc}}(\mathbb{R}^N)$, then, using the fact that $u_0 \in L^q(\mathbb{R}^N)$ and (6.3), we have $\tilde{u} \in C([0, T_{\max}), L^q(\mathbb{R}^N))$ by Theorem 1. It follows that $\sup_{0 < t < T} t^b \|\tilde{u}(t)\|_{L^q} \leq \delta$ if $T > 0$ is sufficiently small. Therefore, by uniqueness, $u = \tilde{u}$ on $[0, T]$, so that $u \in C([0, T], C_0(\mathbb{R}^N))$.

Next, we show that $u \in C([T, \infty), C_0(\mathbb{R}^N))$ by a bootstrap argument. Indeed, for $t > T$, we write

$$\begin{aligned} u(t) - e^{-t(-\Delta)^{\beta/2}} u_0 &= \int_0^t e^{-(t-s)(-\Delta)^{\beta/2}} \int_0^T (s-\sigma)^{-\gamma} |u|^{p-1} u(\sigma) \, d\sigma \, ds \\ &+ \int_0^t e^{-(t-s)(-\Delta)^{\beta/2}} \int_T^s (s-\sigma)^{-\gamma} |u|^{p-1} u(\sigma) \, d\sigma \, ds \\ &\equiv I_1(t) + I_2(t). \end{aligned}$$

Since $u \in C([0, T], C_0(\mathbb{R}^N))$, it follows that $I_1 \in C([T, \infty), C_0(\mathbb{R}^N))$. Also, by the calculations used to construct the fixed point, using the fact that $t^{-b} \leq T^{-b} < \infty$ and $pq > q$, $I_1 \in C([T, \infty), L^q(\mathbb{R}^N))$. Next, note that $N(p/q - 1/q)/\beta < 1$ by (6.3). Therefore, there exists $r \in (q, \infty]$ such that

$$\frac{N}{\beta} \left(\frac{p}{q} - \frac{1}{r} \right) < 1. \quad (6.13)$$

Let $T < s < t$ (the case of $s \leq T \leq t$ is obvious). Since $u \in L^\infty((0, \infty), L^q(\mathbb{R}^N))$, we have $|u|^{p-1} u \in L^\infty((T, s), L^{q/p}(\mathbb{R}^N))$, and it easily follows, using (2.4) and (6.13), that $I_2 \in C([T, \infty), L^r(\mathbb{R}^N))$. As the terms $e^{-(-\Delta)^{\beta/2} t} u_0$ and I_1 both belong to $C([T, \infty), C_0(\mathbb{R}^N)) \cap C([T, \infty), L^q(\mathbb{R}^N))$, we see that $u \in C([T, \infty), L^r(\mathbb{R}^N))$. Iterating this procedure a finite number of times, we deduce that $u \in C([T, \infty), C_0(\mathbb{R}^N))$. This completes the proof.

7 Necessary conditions for local and global existence

In this section, we establish necessary conditions for the existence of local or global weak solutions to the problem (1.1); these conditions depend on the behavior of the initial condition for large x .

Theorem 5 (*Necessary conditions for global existence*)

Let $u_0 \in L_{Loc}^\infty(\mathbb{R}^N)$, $u_0 \geq 0$, $0 < \beta \leq 2$ and $p > 1$. If u is a global weak solution to problem (1.1), then there is a positive constant $C > 0$ such that

$$\liminf_{|x| \rightarrow \infty} (u_0(x) |x|^{\frac{\beta(2-\gamma)}{p-1}}) \leq C. \quad (7.1)$$

Proof. Let u be a global weak solution to (1.1), then $u \in L^p((0, R^\beta), L^\infty(B_{2R}))$ for all $R > 0$, where B_{2R} stands for the closed ball of center 0 and radius $2R$. So, we repeat the same calculation as in the proof of Theorem 1 (here in bounded domain) by taking $\varphi(x, t) := D_{t|T}^\alpha \tilde{\varphi}(x, t) := D_{t|T}^\alpha (\varphi_1(x/R) \varphi_2(t))$ instead of the one chosen in Theorem 1, where $0 \leq \varphi_1 \in D(\Delta_D^{\beta/2})$ is the first eigenfunction of the fractional Laplacian operator $\Delta_D^{\beta/2}$ in B_2 , with homogeneous Dirichlet boundary condition (2.6), associated to the first

eigenvalue $\lambda := \lambda_1^{\beta/2}$, and $\varphi_2(t) := (1 - t/R^\beta)_+^\ell$ for $\ell \gg 1$ large enough. Then, as for the estimate (4.11), we obtain, with $\Sigma := [0, R^\beta] \times B_{2R}$,

$$\begin{aligned} & \int_{\Sigma} u^p \tilde{\varphi} dx dt + C R^{-\alpha\beta} \int_{|x| \leq 2R} u_0(x) \varphi_1(x/R) dx \\ & \leq C \int_{\Sigma} \varphi_1(x/R) (\varphi_2(t))^{-\frac{1}{p-1}} \left| D_{t|R^\beta}^{1+\alpha} \varphi_2(t) \right|^{\tilde{p}} dx dt \\ & \quad + C \int_{\Sigma} (\varphi_1(x/R))^{-\frac{1}{p-1}} (\varphi_2(t))^{-\frac{1}{p-1}} \left| \Delta_D^{\beta/2} \varphi_1(x/R) D_{t|R^\beta}^\alpha \varphi_2(t) \right|^{\tilde{p}} dx dt, \end{aligned} \quad (7.2)$$

where $\alpha := 1 - \gamma$ and $\tilde{p} := p/(p-1)$. If we take the scaled variables $\tau = t/R^\beta$, $\xi = x/R$ and use the fact that $\Delta_D^{\beta/2} \varphi_1(x/R) = R^{-\beta} \lambda \varphi_1(x/R)$ in the right-hand side of (7.2), take into account the positivity of u , we infer that

$$\begin{aligned} & C R^{-\alpha\beta} \int_{|\xi| \leq 2} u_0(R\xi) \varphi_1(\xi) d\xi \\ & \leq C(R) \int_{|\xi| \leq 2} \varphi_1(\xi) d\xi \\ & = C(R) \int_{|\xi| \leq 2} |R\xi|^{\beta(1+\alpha)(\tilde{p}-1)} |R\xi|^{\beta(1+\alpha)(1-\tilde{p})} \varphi_1(\xi) d\xi \\ & \leq C(R) (2R)^{\beta(1+\alpha)(\tilde{p}-1)} \int_{|\xi| \leq 2} |R\xi|^{\beta(1+\alpha)(1-\tilde{p})} \varphi_1(\xi) d\xi \end{aligned}$$

where $\bar{C}(R) = R^{\beta-(1+\alpha)\beta\tilde{p}}(C + C\lambda)$, and so

$$\int_{|\xi| \leq 2} u_0(R\xi) \varphi_1(\xi) d\xi \leq C \int_{|\xi| \leq 2} |R\xi|^{\beta(1+\alpha)(1-\tilde{p})} \varphi_1(\xi) d\xi. \quad (7.3)$$

Using the estimate

$$\begin{aligned} \inf_{|\xi| > 1} (u_0(R\xi) |R\xi|^{\beta(1+\alpha)(\tilde{p}-1)}) \int_{|\xi| \leq 2} |R\xi|^{\beta(1+\alpha)(1-\tilde{p})} \varphi_1(\xi) d\xi & \leq \int_{1 < |\xi| \leq 2} u_0(R\xi) \varphi_1(\xi) d\xi \\ & \leq \int_{|\xi| \leq 2} u_0(R\xi) \varphi_1(\xi) d\xi \end{aligned}$$

in the left-hand side of (7.3), we conclude, after dividing by $\int_{|\xi| \leq 2} |R\xi|^{\beta(1+\alpha)(1-\tilde{p})} \varphi_1(\xi) d\xi$ that

$$\inf_{|\xi| > 1} (u_0(R\xi) |R\xi|^{\beta(1+\alpha)(\tilde{p}-1)}) \leq C. \quad (7.4)$$

Passing to the limit in (7.4), as $R \rightarrow \infty$, we obtain

$$\liminf_{|x| \rightarrow \infty} (u_0(x) |x|^{\beta(1+\alpha)(\tilde{p}-1)}) \leq C.$$

□

Corollary 1 (*sufficient conditions for the nonexistence of global solutions*)

Let $u_0 \in L_{Loc}^\infty(\mathbb{R}^N)$, $u_0 \geq 0$, $0 < \beta \leq 2$ and $p > 1$. If

$$\liminf_{|x| \rightarrow \infty} (u_0(x) |x|^{\frac{\beta(2-\gamma)}{p-1}}) = +\infty,$$

then the problem (1.1) cannot admit a global weak solution. \square

Next, we give a necessary condition for local existence where we obtain a similar estimate of T founded in the proof of Theorem 1, as $|x|$ goes to infinity.

Theorem 6 (*Necessary conditions for local existence*)

Let $u_0 \in L_{Loc}^\infty(\mathbb{R}^N)$, $u_0 \geq 0$, $\beta \in (0, 2]$ and $p > 1$. If u is a local weak solution to problem (1.1) on $[0, T]$ where $0 < T < +\infty$, then we have

$$\liminf_{|x| \rightarrow \infty} u_0(x) \leq C T^{-\frac{2-\gamma}{p-1}}, \quad (7.5)$$

for some positive constant $C > 0$.

Note that, if $A := \liminf_{|x| \rightarrow \infty} u_0(x)$, then we obtain a similar estimate as that found in (3.4),

$$\frac{T^{2-\gamma} A^{p-1}}{C^{p-1}} \leq 1.$$

Proof. We take here, for $R > 0$ sufficiently large, $\varphi(x, t) := D_{t|T}^\alpha \tilde{\varphi}(x, t) := D_{t|T}^\alpha (\varphi_1(x/R) \varphi_2(t))$ where $\varphi_2(t) := (1 - t/T)_+^\ell$ instead of the one chosen in Theorem 5. Then, as (7.2), we obtain

$$\begin{aligned} & \int_{\Sigma_1} u^p \tilde{\varphi} \, dx \, dt + C T^{-\alpha} \int_{|x| \leq 2R} u_0(x) \varphi_1(x/R) \, dx \\ & \leq C \int_{\Sigma_1} \varphi_1(x/R) (\varphi_2(t))^{-\frac{1}{p-1}} \left| D_{t|T}^{1+\alpha} \varphi_2(t) \right|^{\tilde{p}} \, dx \, dt \\ & \quad + C \int_{\Sigma_1} (\varphi_1(x/R))^{-\frac{2}{p-1}} (\varphi_2(t))^{-\frac{1}{p-1}} \left| \Delta_D^{\beta/2} \varphi_1(x/R) D_{t|T}^\alpha \varphi_2(t) \right|^{\tilde{p}} \, dx \, dt, \end{aligned} \quad (7.6)$$

where $\Sigma_1 := [0, T] \times \{x \in \mathbb{R}^N; |x| \leq 2R\}$, $\alpha := 1 - \gamma$ and $\tilde{p} := p/(p-1)$. Now, in the right-hand side of (7.6), we take the following scale of variables $\tau = T^{-1}t$, $\xi = R^{-1}x$ and we use the fact that $\Delta_D^{\beta/2} \varphi_1(x/R) = R^{-\beta} \lambda \varphi_1(x/R)$, while in the left-side we use the positivity of u , then we get

$$C T^{-\alpha} \int_{|\xi| \leq 2} u_0(R\xi) \varphi_1(\xi) \, d\xi \leq \left(C T^{1-(1+\alpha)\tilde{p}} + C \lambda T^{1-\alpha\tilde{p}} R^{-\beta\tilde{p}} \right) \int_{|\xi| \leq 2} \varphi_1(\xi) \, d\xi;$$

and so

$$\int_{|\xi| \leq 2} u_0(R\xi) \varphi_1(\xi) \, d\xi \leq C(R, T) \int_{|\xi| \leq 2} \varphi_1(\xi) \, d\xi \quad (7.7)$$

where $C(R, T) = C T^{(1+\alpha)(1-\tilde{p})} + C T^{1+\alpha(1-\tilde{p})} R^{-\beta\tilde{p}}$.

Using the estimate

$$\begin{aligned} \inf_{|\xi|>1} (u_0(R\xi)) \int_{|\xi|\leq 2} \varphi_1(\xi) d\xi &\leq \int_{1<|\xi|\leq 2} u_0(R\xi) \varphi_1(\xi) d\xi \\ &\leq \int_{|\xi|\leq 2} u_0(R\xi) \varphi_1(\xi) d\xi \end{aligned}$$

in the left-hand side of (7.7), we conclude, after dividing by $\int_{|\xi|\leq 2} \varphi_1(\xi) d\xi$, that

$$\inf_{|\xi|>1} u_0(R\xi) \leq C(R, T). \quad (7.8)$$

Passing to the limit in (7.8), as $R \rightarrow \infty$, we obtain

$$\liminf_{|x|\rightarrow\infty} u_0(x) \leq C T^{(1+\alpha)(1-\tilde{p})} = C T^{-\frac{2-\gamma}{p-1}}.$$

□

Appendix

In this appendix, we give a proof of Ju's inequality (see Proposition 3.3 in [17]), in dimension $N \geq 1$ where $\delta \in [0, 2]$ and $q \geq 1$, for all nonnegative Schwartz function ψ (in the general case)

$$(-\Delta)^{\delta/2} \psi^q \leq q \psi^{q-1} (-\Delta)^{\delta/2} \psi. \quad (7.9)$$

The cases $\delta = 0$ and $\delta = 2$ are obvious, as well as $q = 1$. If $\delta \in (0, 2)$ and $q > 1$, using [7, Definition 3.2], we have

$$(-\Delta)^{\delta/2} \psi(x) = -c_N(\delta) \int_{\mathbb{R}^N} \frac{\psi(x+z) - \psi(x)}{|z|^{N+\delta}} dz, \quad \forall x \in \mathbb{R}^N, \quad (7.10)$$

where $c_N(\delta) = 2^\delta \Gamma((N+\delta)/2) / (\pi^{N/2} \Gamma(1-\delta/2))$. Then

$$(\psi(x))^{q-1} (-\Delta)^{\delta/2} \psi(x) = -c_N(\delta) \int_{\mathbb{R}^N} \frac{(\psi(x))^{q-1} \psi(x+z) - (\psi(x))^q}{|z|^{N+\delta}} dz.$$

By Young's inequality we have

$$(\psi(x))^{q-1} \psi(x+z) \leq \frac{q-1}{q} (\psi(x))^q + \frac{1}{q} (\psi(x+z))^q.$$

Therefore,

$$(\psi(x))^{q-1} (-\Delta)^{\delta/2} \psi(x) \geq \frac{-c_N(\delta)}{q} \int_{\mathbb{R}^N} \frac{(\psi(x+z))^q - (\psi(x))^q}{|z|^{N+\delta}} dz = \frac{1}{q} (-\Delta)^{\delta/2} (\psi(x))^q.$$

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References

- [1] D. ANDREUCCI & A. F. TEDEEV, Universal bounds at the blow-up time for nonlinear parabolic equations, *Adv. Differential Equations* **10** (2005), no. 1, 89 – 120.
- [2] J. BALL, Strongly Continuous Semigroups, Weak Solutions, And The variation of Constants Formula, *Proc. Amer. Math. Soc.* **63** (1977) no. 2, 370 – 373.
- [3] J. BALL, Remarks on blow-up and nonexistence theorem for nonlinear evolution equations, *Quart. J. Math. Oxford* **28** (1977) no. 2, 473 – 486.
- [4] P. BARAS & R. KERSNER, Local and global solvability of a class of semilinear parabolic equations. *J. Differential Equations* **68** (1987), no. 2, 238 – 252.
- [5] P. BARAS & M. PIERRE, Critère d’existence de solutions positives pour des équations semi-linéaires non monotones, *Ann. Inst. H. Poincaré Anal. Non Linéaire* **2** (1985) 185212.
- [6] M. BIRKNER, J. A. LOPEZ-MIMBELA & A. WAKOLBINGER, Blow-up of semilinear PDE’s at the critical dimension. A probabilistic approach (English summary), *Proc. Amer. Math. Soc.* **130** (2002), no. 8, 2431 – 2442 (electronic).
- [7] K. BOGDAN & T. BYCZKOWSKI, Potential theory for the α -stable Schrödinger operator on bounded Lipschitz domains, *Studia Mathematica* **133** (1999), no. 1, 53 – 92.
- [8] T. CAZENAVE & A. HARAUX, *Introduction aux problèmes d’évolution semi-linéaires*, Ellipses, Paris, (1990).
- [9] M. CHLEBIK & M. FILA, From critical exponents to blow-up rates for parabolic problems, *Rend. Mat. Appl.* (7) 19 (4) (1999), 449 – 470.
- [10] T. CAZENAVE, F. DICKSTEIN & F. D. WEISSLER, An equation whose Fujita critical exponent is not given by scaling, *Nonlinear Analysis* **68** (2008), 862 – 874.
- [11] S. CHANDRASEKHAR, Stochastic problems in physics and astronomy. *Rev. Mod. Phys.* vol. **15** (1943), 1 – 89.
- [12] M. FILA & P. QUITTNER, The Blow-Up Rate for a Semilinear Parabolic System, *J. of Mathematical Analysis and Applications* **238** (1999), 468 – 476.
- [13] A. FINO & G. KARCH, Decay of mass for nonlinear equation with fractional Laplacian, *J. Monatsh. Math.* DOI 10.1007/s00605 – 009 – 0093 – 3.
- [14] H. FUJITA, On the blowing up of solutions of the problem for $u_t = \Delta u + u^{1+\alpha}$, *J. Fac. Sci. Univ. Tokyo* **13** (1966), 109 – 124.
- [15] M. GUEDDA & M. KIRANE, Criticality for some evolution equations, *Differential Equations* **37** (2001), 511 – 520.

- [16] B. HU, Remarks on the blowup estimate for solutions of the heat equation with a nonlinear boundary condition, *Differential Integral Equations* **9** (1996), 891 – 901.
- [17] N. JU, The Maximum Principle and the Global Attractor for the Dissipative 2-D Quasi-Geostrophic Equations, *Comm. Pure. Appl. Ana.* (2005), 161 – 181.
- [18] A. A. KILBAS, H. M. SRIVASTAVA & J. J. TRUJILLO, *Theory and Applications of Fractional Differential Equations*, 2006.
- [19] M. KIRANE, Y. LASKRI & N.-E. TATAR, Critical exponents of Fujita type for certain evolution equations and systems with Spatio-Temporal Fractional derivatives, *J. Math. Anal. Appl.* **312** (2005), 488 – 501.
- [20] M. KIRANE & M. QAFSAOUI, Global nonexistence for the cauchy problem of some nonlinear Reaction-Diffusion systems, *J. Math. Analysis and Appl.* **268** (2002), 217 – 243.
- [21] K. KOBAYASHI, On some semilinear evolution equations with time-lag, *Hiroshima Math. J.* **10** (1980), 189 – 227.
- [22] N. V. KRYLOV, *Lectures on Elliptic and Parabolic Equations in Hölder Spaces*, Amer. Math. Soc. (Graduate Studies in Mathematics) v. **12**, 1996.
- [23] O. A. LADYŽENSKAJA, V. A. SOLONNIKOV & N. N. URAL’CEVA, *Linear and Quasilinear Equations of Parabolic Type*, Amer. Math. Soc., Providence, RI, 1968.
- [24] E. MITIDIERI & S. I. POHOZAEV, A priori estimates and blow-up of solutions to nonlinear partial differential equations and inequalities, *Proc. Steklov. Inst. Math.* **234** (2001), 1 – 383.
- [25] E. MITIDIERI & S. I. POHOZAEV, Nonexistence of weak solutions for some degenerate elliptic and parabolic problems on \mathbb{R}^N , *J. Evol. Equ.* **1** (2001), 189 – 220.
- [26] S. MIYAJIMA & H. SHINDOH, Gaussian Estimates of Order α and L^p -Spectral Independence of Generators of C_0 -Semigroups, *Positivity* **11** (2007), 15 – 39.
- [27] M. NAGASAWA & T. SIRAO, Probabilistic treatment of the blowing up of solutions for a nonlinear integral equation, *Trans. Amer. Math. Soc.* **139** (1969), 301 – 310.
- [28] C. A. ROBERTS & W. E. OLMSTEAD, Blow-up in a subdiffusive medium of infinite extent. *Fract. Calc. Appl. Anal.* **12** (2009), no. 2, 179 – 194.
- [29] I. M. SOKOLOV & J. KLAFTER, From diffusion to anomalous diffusion: a century after Einstein’s Brownian Motion, *Chaos* **15** (2005), 6103 – 6109.
- [30] S. G. SAMKO, A. A. KILBAS & O. I. MARICHEV, *Fractional integrals and derivatives, Theory and Applications*, Gordon and Breach Science Publishers, 1987.
- [31] P. SOUPLET, Blow-up in nonlocal reaction-diffusion equations, *SIAM J. Math. Anal.* **29** (1998), 1301 – 1334.

- [32] P. SOUPLET, Monotonicity of solutions and blow-up for semilinear parabolic equations with nonlinear memory, *Z. angew. Math. Phys.* **55** (2004), 28 – 31.
- [33] K. YOSIDA, *Functional Analysis*, sixth Edition, Springer-Verlag, Berlin Heidelberg, New York 1980.