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Implicit time discretization of the mean curvature flow with a discontinuous forcing term

A. Chambolle ^{*}, M. Novaga [†]

Abstract

We consider an implicit time discretization for the motion of a hypersurface driven by its anisotropic mean curvature. We prove some convergence results of the scheme under very general assumptions on the forcing term, which include in particular the case of a typical path of the Brownian motion. We compare this limit with other available solutions, whenever they are defined. As a by-product of the analysis, we also provide a simple proof of the coincidence of the limit flow with the regular evolutions, defined for small times, in the case of a regular forcing term.

1 Introduction

Mean curvature flow has attracted a lot of attention in the past few years. Being one of the simplest evolution of hypersurfaces of \mathbb{R}^n , in its analysis arise many difficult issues mainly related to the formation of singularities, which sometimes lead to changes of the topology. To deal with this phenomenon, several notions of weak solutions have been proposed, such as (only to mention some) the varifold theory of Brakke [10], the level-set solution defined through the viscosity theory [18, 19, 15], the minimal barrier method of De Giorgi [16], the limit of a reaction-diffusion equations [14, 21] and the minimizing movements method [1, 25, 2], that corresponds to an implicit time-discrete scheme.

Each of these methods has different features and presents advantages and disadvantages. In particular, the level-set method always provides a unique solution, globally defined in time in the class of compact subsets of \mathbb{R}^n , but it is often very difficult to prove that such a solution is a regular hypersurface. There are even some singular situations in which this solution becomes a compact set with nonempty interior, showing the so-called *fattening phenomenon*. The minimal barrier method is a geometric counterpart of the level-set method and produces essentially the same solution [6].

^{*}CMAP, CNRS UMR 7641, Ecole Polytechnique, 91128, Palaiseau Cedex, France,
e-mail: antonin.chambolle@polytechnique.fr

[†]Dipartimento di Matematica, Università di Pisa, via Buonarroti 2, 56127 Pisa, Italy, e-mail: novaga@dm.unipi.it

On the contrary, the minimizing movements method produces a solution, called the *flat flow*, which can be nonunique but is always a (possibly nonsmooth) hypersurface. One of the difficulties in this approach is to show that the solution coincides with the classical smooth solution, whenever the latter exists, a property which is very easy to prove in the context of level-set viscosity solutions. One faces similar difficulties in proving that the flat flow is always contained in the level-set solution.

In this paper we study the (anisotropic) mean curvature flow with a possibly discontinuous driving force, by adapting the minimizing movements method, which has been originally developed without any forcing term. More precisely, we consider the evolution $E(t)$ of a set whose boundary is driven by the velocity

$$V(x, t) = -(\kappa_\phi(x, t) + g(x, t))n_\phi(x, t) \quad (1)$$

for any $x \in \partial E(t)$, where $\kappa_\phi(x, t)$ and $n_\phi(x, t)$ are respectively the ϕ -curvature and ϕ -normal to $\partial E(t)$ at x (see Section 2 for the precise definitions, and [9] for a general introduction on curvatures in Finsler geometry).

The purpose of this paper is twofold:

1. We extend the method of minimizing movements (and the proofs of consistency in [13]) to evolutions with a driving force, providing simple proofs of the coincidence with regular solutions and the inclusion in the level-set solution (see Cor. 4.6, Prop. 4.8, 4.9).
2. Our approach applies to the case where the forcing term is discontinuous. One important example is a forcing term which is the time derivative of an Hölder continuous function $G(t)$, *e.g.*, a typical path of the Brownian motion dW/dt . It also covers the case of spatially correlated Brownian motion, typically of the form $g(x, t) = dW/dt(t) + g^0(x, t)$ where g^0 is Lipschitz-continuous in x and continuous in t (see (18) for a precise formulation). A theory yielding existence and uniqueness for such evolutions, based on a level-set formulation in the framework of the viscosity theory, has been recently developed in [23, 24], and a corresponding theory in the framework of minimal barriers, valid only for x -independent forcing terms, has been proposed in [17]. We also refer to [26, 27] for a similar approach to a related problem, which still uses an implicit time discretization procedure.

We do not address in the present paper the issue of continuity in time (in a suitable topology) of the limit flat flow, even if we prove some weaker continuity results with respect to the Hausdorff distance (see Propositions 4.3, 4.4 and Remark 4.5).

The plan of this paper is as follows: in Section 2, we define the appropriate notion of sub- and superflow associated to the evolution equation (1), and we recall the definition of maximal and minimal barriers in the sense of De Giorgi. These allow to define generalized evolutions, which are essentially equivalent to the evolutions defined in terms of level sets of viscosity solutions, when the forcing term is regular enough.

Then, our implicit time-discretization scheme is defined in Section 3, and we show that it is consistent with our sub- and superflows (Thm. 3.3). As a corollary, we obtain a comparison results for sub- and superflows, which follows from the monotonicity of our scheme.

This consistency result is used in Section 4 to study the convergence of our time-discrete scheme, as the time step goes to zero. We define a notion of weak solution starting from an initial surface ∂E (given by $\Gamma(t) = E^*(t) \setminus E_*(t)$, where $E_*(t) \subset E^*(t)$ are two evolving sets with $E_*(0) = \text{int}(E)$, $E^*(0) = \overline{E}$), which coincides with the barrier solution as long as the latter is unique (Cor. 4.6). Under additional assumptions on the evolution law, we deduce that it is contained in the zero level-set of the viscosity solution (Prop. 4.8). This is also true for a particular class of stochastic evolutions (Prop. 4.9).

In Sections 4.2-4.3, under some further assumption on the forcing term (which still allows a stochastic forcing), we build a level-set evolution $u(t)$ starting from any bounded, uniformly continuous function u_0 . In particular, for all initial data $\{u_0 = s\}$ but a countable number, it shows that we can define a generalized flow $\{u(t) = s\}$ which remains a continuous hypersurface (with possible singularities). We do not prove that it is unique, though. If the forcing is regular, $u(t)$ is the same as the unique viscosity solution of the geometric equation associated to the flow. In general, we expect it to coincide with the solution defined by Lions and Souganidis in [23, 24].

2 Preliminary definitions and results

Let $\phi : \mathbb{R}^N \rightarrow \mathbb{R}$ be a norm on \mathbb{R}^N (that is, an even, convex, one-homogeneous function) such that $\phi \in C^2(\mathbb{R}^N \setminus \{0\})$ (we shall simply say that ϕ is smooth) and $\nabla^2(\phi^2) \geq c\text{Id}$ for some $c > 0$, so that ϕ is uniformly convex or *elliptic*. Most of our results could be extended to more general norms (or even possibly non-even convex one-homogeneous functions), but the proof of the consistency theorem 3.3, in the form presented here, needs such a regularity. Moreover, providing a clear and sound definition of a sub/superflow, as in Definition 2.1 below, is more difficult if the anisotropy is only Lipschitz-continuous, or non elliptic.

Let ϕ° be the polar norm, that is, $\phi^\circ(\xi) := \sup_{\phi(\eta) \leq 1} \xi \cdot \eta$, for all $\xi \in \mathbb{R}^N$. It turns out that also ϕ° is smooth and elliptic. In the sequel, the couple (ϕ, ϕ°) will be referred as *the anisotropy*. A ball of radius $r > 0$ centered in $x_0 \in \mathbb{R}^N$ for the norm ϕ , *i.e.*, the set $W_\phi(x_0, r) := \{\phi(x - x_0) \leq r\}$, will be called a *Wulff shape* (we set for simplicity $W_\phi := W_\phi(0, 1)$).

When $E, F \subset \mathbb{R}^N$, we denote by $\text{dist}_\phi(E, F)$ the distance between E and F with respect to ϕ :

$$\text{dist}_\phi(E, F) := \inf_{x \in E, y \in F} \phi(x - y).$$

Given a set $E \subset \mathbb{R}^N$, we also define $d_E(x)$, the signed distance function to ∂E (with

respect to the norm ϕ), by

$$d_E(x) := \inf_{y \in E} \phi(x - y) - \inf_{y \in \mathbb{R}^N \setminus E} \phi(y - x).$$

We let $n_\phi(x) := \nabla \phi^\circ(\nabla d_E(x))$ and $\kappa_\phi(x) := \operatorname{div} n_\phi(x)$ be respectively the ϕ -normal and the ϕ -curvature of ∂E at x . Notice that if ∂E is of class C^2 , then the functions n_ϕ and κ_ϕ are defined and continuous in an open neighbourhood of ∂E . We refer to [9] for a general introduction to the anisotropic curvature flow.

We say that E satisfies an interior (resp. exterior) εW_ϕ -condition, $\varepsilon > 0$, if $E = \{d_E < -\varepsilon\} + \varepsilon W_\phi$ (resp. $\mathbb{R}^N \setminus E = \{d_E > \varepsilon\} + \varepsilon W_\phi$), which is equivalent to require that at each point of ∂E , there is a Wulff shape of radius ε inside E (resp., outside E), that is tangent to ∂E at x .

2.1 Evolution law

Definition 2.1. Let $E(t) \subset \mathbb{R}^N$, $t \in [t_0, t_1]$ be closed sets. We say that $E(t)$ is a superflow of (1), if there exist a bounded open set $A \subset \mathbb{R}^N$, with $\bigcup_{t_0 \leq t \leq t_1} \partial E(t) \times \{t\} \subset A \times [t_0, t_1]$, and $\delta > 0$, such that $d(x, t) = d_{E(t)}(x) \in C^0([t_0, t_1]; C^2(A))$, and

$$d(x, s) - d(x, t) \geq \int_t^s \operatorname{div} \nabla \phi^\circ(\nabla d)(x, \tau) d\tau + G(x, s) - G(x, t) + \delta(s - t), \quad (2)$$

for a.e. $x \in A$ and any t, s with $t_0 \leq t \leq s \leq t_1$, where $G(x, t) := \int_0^t g(x, s) ds$.

We say that $E(t)$ is a subflow whenever there exist $A \subset \mathbb{R}^N$ as above and $\delta < 0$, such that the reverse inequality holds in (2).

We indicate with \mathcal{F}^+ (resp. \mathcal{F}^-) the family of all superflows (resp. subflows) of (1).

We observe that if g is continuous in (x, t) , and d is C^1 in t , condition (2) is equivalent to require

$$\frac{\partial d}{\partial t} > \operatorname{div} \nabla \phi^\circ(\nabla d) + g$$

in $A \times [t_0, t_1]$. On the other hand, Definition 2.1 still makes sense if the driving term is the “time-derivative” of a function $G \in C^0([t_0, t_1]; L^\infty(A))$, even when G is nondifferentiable with respect to t .

2.2 Barriers

We recall the definition of minimal and maximal barrier in the sense of De Giorgi. We refer to [8, 5, 7] for a more general introduction to this topic.

Definition 2.2. We say that a function $\Phi : [t_0, +\infty) \rightarrow \mathcal{P}(\mathbb{R}^N)$ ($\mathcal{P}(\mathbb{R}^N)$ is the set of all subsets of \mathbb{R}^N) is a barrier with respect to \mathcal{F}^+ if for any $\Sigma(t) \in \mathcal{F}^+$, $t \in [a, b] \subset [t_0, +\infty)$, $\Sigma(a) \subseteq \Phi(a)$ implies $\Sigma(b) \subseteq \Phi(b)$.

Similarly, we say that Φ is a barrier with respect to \mathcal{F}^- if for any $\Sigma(t) \in \mathcal{F}^-$, $t \in [a, b] \subset [t_0, +\infty)$, $\Sigma(a) \supseteq \Phi(a)$ implies $\Sigma(b) \supseteq \Phi(b)$.

In the following we denote by $\mathcal{B}_{t_0}^\pm$ the class of all barriers with respect to \mathcal{F}^\pm , defined on $[t_0, +\infty)$.

Definition 2.3. Let $E \subseteq \mathbb{R}^N$, $t_0 \in \mathbb{R}$. The minimal barrier $\mathcal{M}(E, t_0) : [t_0, +\infty) \rightarrow \mathcal{P}(\mathbb{R}^N)$ starting from E at time t_0 is defined as:

$$\mathcal{M}(E, t_0)(t) := \bigcap \left\{ \Phi(t) : \Phi \in \mathcal{B}_{t_0}^+, \Phi(t_0) \supseteq E \right\}.$$

We define the maximal barrier $\mathcal{N}(E, t_0) : [t_0, +\infty) \rightarrow \mathcal{P}(\mathbb{R}^N)$ starting from E at time t_0 as:

$$\mathcal{N}(E, t_0)(t) := \bigcup \left\{ \Phi(t) : \Phi \in \mathcal{B}_{t_0}^-, \Phi(t_0) \subseteq E \right\}.$$

We also define the upper and lower regularized barriers as

$$\begin{aligned} \mathcal{M}_*(E, t_0)(t) &:= \bigcup_{\rho > 0} \mathcal{M}(E_\rho^-, t_0)(t) & \mathcal{M}^*(E, t_0)(t) &:= \bigcap_{\rho > 0} \mathcal{M}(E_\rho^+, t_0)(t), \\ \mathcal{N}_*(E, t_0)(t) &:= \bigcup_{\rho > 0} \mathcal{N}(E_\rho^-, t_0)(t) & \mathcal{N}^*(E, t_0)(t) &:= \bigcap_{\rho > 0} \mathcal{N}(E_\rho^+, t_0)(t), \end{aligned}$$

where $E_\rho^\pm = \{d_E \leq \pm \rho\}$.

We recall the following result, proved in [6] (see also [22] for the case of the motion by mean curvature).

Theorem 2.4. Assume that $G(x, t) = \int_0^t g(x, s) ds$, with g continuous. Then, we have $\mathcal{M}^*(E, t_0)(t) = \mathcal{N}^*(E, t_0)(t)$ and $\mathcal{M}_*(E, t_0)(t) = \mathcal{N}_*(E, t_0)(t)$ for any $E \subset \mathbb{R}^N$ and $t \geq t_0$. Moreover, the set $\mathcal{M}^*(E, t_0)(t) \setminus \mathcal{M}_*(E, t_0)(t)$ coincides with the zero level-set of the viscosity solution of the parabolic equation corresponding to (1).

The parabolic equation which is mentioned here is equation (21) or (22). In the sequel, we shall omit the explicit dependence of barriers on t_0 whenever $t_0 = 0$.

2.3 Anisotropic total variation

The total variation of a function $w \in L^1(\Omega)$ is defined as

$$\sup \left\{ \int_\Omega u(x) \operatorname{div} \psi(x) dx : \psi \in C_0^1(\Omega; \mathbb{R}^N), |\psi(x)| \leq 1 \ \forall x \in \Omega \right\}.$$

It turns out that it is finite if and only if the distributional derivative Dw is a bounded Radon measure. In this case, the total variation is equal to the variation $|Dw|(\Omega) = \int_\Omega |Dw|$ of the measure Dw , and w belongs to the space $BV(\Omega)$ of functions with bounded variation.

Given (ϕ, ϕ°) a couple of mutually polar norms in \mathbb{R}^N (an anisotropy), one defines in the same way the anisotropic total variation

$$\int_\Omega \phi^\circ(Dw) = \sup \left\{ \int_\Omega u(x) \operatorname{div} \psi(x) dx : \psi \in C_0^1(\Omega; \mathbb{R}^N), \phi(\psi(x)) \leq 1 \ \forall x \in \Omega \right\}.$$

Clearly, it is finite if and only if $w \in BV(\Omega)$. In the case $w = \chi_E$, the characteristic function of a measurable set E , then $w \in BV(\Omega)$ if and only if E is a set of finite

perimeter in Ω (a *Caccioppoli set*). In this case, one can define a reduced boundary ∂^*E (which is \mathcal{H}^{N-1} -equivalent to the measure theoretical boundary, that is, the set of points where E has Lebesgue density neither 0 nor 1), on which is well defined a normal unit vector $\nu_E(x)$, and such that $D\chi_E = \nu_E \mathcal{H}^{N-1} \llcorner \partial^*E$. Then, one has

$$\int_{\Omega} |D\chi_E| = \mathcal{H}^{N-1}(\partial^*E) \quad \text{and} \quad \int_{\Omega} \phi^{\circ}(D\chi_E) = \int_{\partial^*E} \phi^{\circ}(\nu_E(x)) d\mathcal{H}^{N-1}(x).$$

See [20, 3] for more details.

3 The implicit time discretization

Let Ω be a bounded, convex, open subset of \mathbb{R}^N . Let $G \in C^0([0, +\infty), L^{\infty}(\Omega))$ and let $\omega_{G,T}$ its modulus of continuity in $[0, T]$. Let (ϕ, ϕ°) be the anisotropy, which we assume to be smooth and elliptic. Let $E \subseteq \mathbb{R}^N$. Given $s > t \geq 0$, let w denote the unique solution of

$$\min_{w \in L^2(\Omega)} \int_{\Omega} \phi^{\circ}(Dw) + \frac{1}{2(s-t)} \int_{\Omega} \left(w(x) - d_E(x) - G(x, s) + G(x, t) \right)^2 dx. \quad (3)$$

We let $T_{t,s}(E) = \{x \in \Omega : w(x) < 0\}$. The existence and uniqueness of w minimizing (3) does not raise any difficulty, since the energy which is minimized is trivially strictly convex, and lower-semicontinuous in $L^2(\Omega)$.

Notice that the set $T_{t,s}(E)$ is the minimizer of a prescribed curvature problem, with bounded mean curvature. Indeed, reasoning as in [12, 11, 4], one can check that this set is a solution of the variational problem

$$\min \left(\int_{\Omega \cap \partial^*F} \phi^{\circ}(\nu_F(x)) d\mathcal{H}^{N-1}(x) + \frac{1}{s-t} \int_F d_E(x) + G(x, s) - G(x, t) dx \right),$$

where the minimum is taken among the subsets F of Ω of finite perimeter. It follows that the set $T_{t,s}(E)$ has boundary of class $C^{1,\alpha}$ inside Ω , outside a compact singular set of zero \mathcal{H}^{N-1} -dimension [1] (when $N = 2$, the set $T_{t,s}(E)$ has boundary of class $C^{1,1}$). The variational problem above is the generalization of the approach proposed in [1, 25], for building mean curvature flows without driving terms, through an implicit time discretization.

For $s = t + h$, the Euler-Lagrange equation for w at a point $x \in \partial T_{t,t+h}(E)$ formally reads as

$$d_E(x) = -h \left(\kappa_{\phi}(x) + \frac{G(x, t+h) - G(x, t)}{h} \right),$$

with κ_{ϕ} being the ϕ -curvature at x of $\partial T_{t,t+h}(E)$, so that it corresponds to an implicit time-discretization of (1). Observe also that this approximation is monotone: indeed if $E \subseteq E'$ then $d_E \geq d_{E'}$, which yields $w \geq w'$, w and w' being the solutions of (3) for the distance functions d_E and $d_{E'}$ respectively. We deduce that $\{w < 0\} \subseteq \{w' < 0\}$, that is, $T_{t,s}(E) \subseteq T_{t,s}(E')$.

We will soon show (Theorem 3.3) that this scheme is also consistent, in some sense, with the evolution (1). Before this, let us prove that it is independent on Ω , in the sense that when $\partial E \subset \Omega$, then for $s - t$ is small enough the set $T_{t,s}(E)$ is also the zero sublevel-set of any function w' solving (3) in any larger open set $\Omega' \supseteq \Omega$. This justifies why we may ignore the dependency on Ω in our notation. Here and in the rest of the paper we shall assume that G is defined in the whole space: $G \in C^0([0, +\infty); L^\infty(\mathbb{R}^N))$.

Proposition 3.1. *For any $\delta > 0$ and $T > 0$, there exists $h_0 > 0$ such that if E is a closed set with compact boundary $\partial E \subset \Omega$, such that $\text{dist}_\phi(\partial\Omega, \partial E) \geq \delta$, then when $h \leq h_0$, for any $t \leq T$, the set $T_{t,t+h}(E)$ is the same if computed in Ω or in any larger open set $\Omega' \supseteq \Omega$.*

Before proving this proposition, we show a result that allows us to control in some uniform way the speed at which an initial Wulff shape $\{\phi(x - x_0) \leq \rho\}$ decreases in an iteration of the algorithm. The convexity of Ω is needed in the proof of this result.

Lemma 3.2. *Let $x_0 \in \Omega$ and $\rho > 0$, and let $t \geq 0$. Let w solve*

$$\min_{w \in L^2(\Omega)} \int_{\Omega} \phi^\circ(Dw) + \frac{1}{2h} \int_{\Omega} (w(x) - (\phi(x - x_0) - \rho) - G(x, t+h) + G(x, t))^2 dx. \quad (4)$$

Then

$$w(x) \leq \begin{cases} \phi(x - x_0) + h \frac{N-1}{\phi(x - x_0)} + \Delta_h(t) - \rho & \text{if } \phi(x - x_0) \geq \sqrt{h(N+1)} \\ \sqrt{h} \frac{2N}{\sqrt{N+1}} + \Delta_h(t) - \rho & \text{otherwise,} \end{cases} \quad (5)$$

where $\Delta_h(t) := \|G(\cdot, t+h) - G(\cdot, t)\|_{L^\infty(\Omega)}$.

Proof. The Euler-Lagrange equation for (4) can be written as follows: there exists a field $z \in L^\infty(\Omega; \mathbb{R}^N)$, with $z \in \partial\phi^\circ(\nabla w)$ a.e. and $z \cdot \nu_\Omega = 0$ on $\partial\Omega$, such that

$$w(x) - \phi(x - x_0) + \rho - G(x, t+h) + G(x, t) - h \operatorname{div} z(x) = 0, \quad (6)$$

see for instance [11, 4].

Let \overline{w} denote the function given in the right-hand side of equation (5). Let \overline{z} be the field given by

$$\overline{z}(x) = \begin{cases} \frac{x - x_0}{\phi(x - x_0)} & \text{if } \phi(x - x_0) \geq \sqrt{h(N+1)} \\ \left(1 - \left(\frac{\phi(x - x_0)}{\sqrt{h(N+1)}} - 1\right)^2\right) \frac{x - x_0}{\phi(x - x_0)} & \text{otherwise.} \end{cases}$$

One checks, as in [11, App. B], that $\overline{z} \in \partial\phi^\circ(\nabla \overline{w}(x))$ a.e., and

$$\frac{\overline{w}(x) - \phi(x - x_0) + \rho}{h} - \operatorname{div} \overline{z}(x) = \frac{\Delta_h(t)}{h}$$

a.e. in Ω . Moreover, if $x \in \partial\Omega$, $\bar{z}(x) \cdot \nu_\Omega(x)$ has the sign of $(x - x_0) \cdot \nu_\Omega(x)$, which is nonnegative since Ω is convex. By definition of $\Delta_h(t)$, we deduce that \bar{w} is a supersolution for (6). It follows that $\bar{w} \geq w$ a.e. in Ω : indeed, we have

$$\begin{aligned} \int_{\Omega} [(w - \bar{w})^+]^2 &= \int_{\{w > \bar{w}\}} (w - \bar{w})(\operatorname{div} z - \operatorname{div} \bar{z}) \\ &= \int_{\partial\Omega \cap \{w > \bar{w}\}} (w - \bar{w})(z - \bar{z}) \cdot \nu_\Omega - \int_{\{w > \bar{w}\}} (\nabla w - \nabla \bar{w}) \cdot (z - \bar{z}) \\ &\leq - \int_{\partial\Omega \cap \{w > \bar{w}\}} (w - \bar{w}) \bar{z} \cdot \nu_\Omega \leq 0, \end{aligned}$$

which shows the inequality. \square

Proof of Proposition 3.1. We assume $E \subset \Omega$, the proof being identical in the other case $\mathbb{R}^N \setminus E \subset \Omega$. Let w solve

$$\min_{w \in L^2(\Omega)} \int_{\Omega} \phi^\circ(Dw) + \frac{1}{2h} \int_{\Omega} (w(x) - d_E(x) - G(x, t+h) + G(x, t))^2 dx,$$

and let $x \in \Omega$ with $d_E(x) \geq \delta/2$. One has $d_E \geq \delta/2 - \phi(\cdot - x)$ in Ω . Invoking Lemma 3.2, we deduce that $w(x) \geq \delta/2 - \Delta_h(t) - 2N\sqrt{h}/\sqrt{N+1}$. Hence if $h_0 \leq 1$ is such that $\omega_{G,T+1}(h_0) + 2N\sqrt{h_0}/\sqrt{N+1} \leq \delta/4$, we find that when $h \leq h_0$ we have $w(x) \geq \delta/4$.

Let now $\Omega' \supseteq \Omega$. If $h \leq h_0$, we have in particular that $d_E(x) + G(x, t+h) - G(x, t) \geq \delta/4$ for any $x \in \Omega' \setminus \Omega$, $t \leq T$. We can hence reproduce the proof of Corollary A.2 in [12], that shows that if w' is the solution of the same problem as w , but in Ω' instead of Ω , then $w' \wedge (\delta/4)$ is the function equal to $w \wedge (\delta/4)$ in Ω and to $\delta/4$ in $\Omega' \setminus \Omega$. We deduce $\{w < 0\} = \{w' < 0\}$. Observe that in this proof, the larger domain Ω' does not need to be convex. \square

The previous proposition allows to define in a unique and intrinsic way the evolution $T_{t,t+h}(E)$ in \mathbb{R}^N for any $t \geq 0$ and $h > 0$, of a set E with compact boundary $\partial E \Subset \mathbb{R}^N$, by considering the corresponding set computed in a ball with radius large enough. Therefore, from now on we shall assume $\partial E \Subset \mathbb{R}^N$ and we shall omit the dependence on Ω in the construction of the limit flow. We now prove our main consistency result.

Theorem 3.3. *Let $E(t)$, $t \in [t_0, t_1]$ be a superflow of (1). Then there exists h_0 such that for any $h < h_0$ and any t with $t_0 \leq t < t+h \leq t_1$, $T_{t,t+h}(E(t)) \supseteq E(t+h)$. Respectively, if $E(t)$ is a subflow of (1), then $T_{t,t+h}(E(t)) \subseteq E(t+h)$ for h small enough.*

Proof. Let $A \subset \mathbb{R}^N$ be the open set associated to the superflow $E(t)$ (cf. Definition 2.1) and let Ω be a bounded, convex open set with $A \Subset \Omega$.

We first observe that there exists $\varepsilon > 0$ such that $C := \{(x, t) : t_0 \leq t \leq t_1, |d(x, t)| \leq \varepsilon\} \subset A \times [t_0, t_1]$. Since $d(\cdot, t)$ is uniformly bounded in $C^2(A)$, we can also assume, possibly reducing ε , that $E(t)$ satisfies for all t an interior and exterior

εW_ϕ -condition. Given t, h with $t_0 \leq t < t+h \leq t_1$, we build from $d(\cdot, t+h)$ a supersolution for problem (3). Consider $\psi : \mathbb{R} \rightarrow \mathbb{R}$ a smooth increasing function with $\psi(s) \geq s$ and $\psi(s) = s$ for $|s| \leq \varepsilon/2$. We set, for $x \in B = \{|d(\cdot, t)| < \varepsilon\}$, $v(x) := \psi(d(x, t+h))$. Then, for $x \in B$, from (2) it follows

$$\begin{aligned} \frac{v(x) - d_{E(t)}(x) - G(x, t+h) + G(x, t)}{h} \\ \geq \frac{d(x, t+h) - d(x, t) - G(x, t+h) + G(x, t)}{h} \\ \geq \frac{1}{h} \int_t^{t+h} (\operatorname{div} \nabla \phi^\circ(\nabla d)(x, s)) \, ds + \delta. \end{aligned}$$

Let now ω be a modulus of continuity for $\operatorname{div} \nabla \phi^\circ(\nabla d)$ in C : we find

$$\frac{v(x) - d_{E(t)}(x) - G(x, t+h) + G(x, t)}{h} \geq \operatorname{div} \nabla \phi^\circ(\nabla d)(x, t+h) + \delta - \omega(h).$$

Observe that for any $x \in B$ it holds $\nabla v(x) = \psi'(d(x, t+h)) \nabla d(x, t+h)$, so that (recall that $\nabla \phi^\circ$ 0-homogeneous), $\nabla \phi^\circ(\nabla v(x)) = \nabla \phi^\circ(\nabla d(x, t+h))$ hence $\operatorname{div} \nabla \phi^\circ(\nabla d)(x, t+h) = \operatorname{div} \nabla \phi^\circ(\nabla v)(x)$. Therefore, if h is small enough so that $\omega(h) \leq \delta$, we get

$$\frac{v(x) - d_{E(t)}(x) - G(x, t+h) + G(x, t)}{h} \geq \operatorname{div} \nabla \phi^\circ(\nabla v)(x).$$

Let w solve (3), with $E = E(t)$ and $s = t+h$. We will show that we may choose ψ in order to have $v \geq w$ on ∂B , so that v is a supersolution for the problem

$$\min \left\{ \int_B \phi^\circ(Du) + \frac{1}{2h} \int_B (u(x) - d_{E(t)}(x) - G(x, t+h) + G(x, t))^2 \, dx : \right. \\ \left. u = w \text{ on } \partial B \right\}$$

(which is solved by w). We will deduce that $v \geq w$ in B , hence $\{w < 0\} \supseteq \{v < 0\} = \{d(\cdot, t+h) < 0\}$, that is, $T_{t, t+h}(E(t)) \supseteq E(t+h)$.

First of all, d is uniformly continuous in time, so that if h is small enough, one has $d(x, t+h) \geq 3\varepsilon/4$ if $d(x, t) = \varepsilon$. If $M > \operatorname{diam} \Omega$, then $M \geq w$ in Ω . We may choose a function ψ with $\psi(3\varepsilon/4) \geq M$, so that $v(x) \geq M \geq w(x)$ if $d(x, t) = \varepsilon$.

On the other hand, since $E(t)$ satisfies the interior εW_ϕ -condition, one deduces from Lemma 3.2 that $w(x) \leq 2N\sqrt{h}/\sqrt{N+1} + \Delta_h(t) - \varepsilon$ whenever $d(x, t) = -\varepsilon$. We observe that $\Delta_h(t) \rightarrow 0$ as $h \rightarrow 0$ uniformly in $[t_0, t_1]$. Hence if h is small enough, we find that $w(x) \leq -3\varepsilon/4$. We can choose ψ such that $\psi(s) \geq -3\varepsilon/4$ for any s , so that $v(x) \geq w(x)$ if $d(x, t) = -\varepsilon$. We conclude that $v \geq w$ on ∂B . Hence v is a supersolution for (3), which implies $T_{t, t+h}(E(t)) \supseteq E(t+h)$.

If $E(t)$ is a subflow, we can reproduce the same proof to show that $T_{t, t+h}(E(t)) \subseteq E(t+h)$. \square

We deduce the following comparison result for sub/superflows.

Corollary 3.4. *Assume that $E_1(t)$, $E_2(t)$ are respectively a superflow and a subflow of (1) on $[t_0, t_1]$, such that $E_1(t_0) \subseteq E_2(t_0)$. Then $E_1(t) \subseteq E_2(t)$ for all $t \in [t_0, t_1]$.*

Proof. By the previous theorem, there exists h_0 such that $T_{t,t+h}(E_1(t)) \supseteq E_1(t+h)$ and $T_{t,t+h}(E_2(t)) \subseteq E_2(t+h)$ for any $t \in [t_0, t_1 - h]$, as soon as $h \leq h_0$. Hence, if $t \in [t_0, t_1]$, we just let $n \geq 1$ be such that $(t - t_0)/n = h \leq h_0$. Then, letting $t_k = t_0 + kh$, one can easily check by induction that $E_1(t_k) \subseteq T_{t_{k-1}, t_k}(E_1(t_{k-1})) \subseteq T_{t_{k-1}, t_k}(E_2(t_{k-1})) \subseteq E_2(t_k)$ for any $1 \leq k \leq n$, which implies the thesis since $t = t_n$. \square

Remark 3.5. It could be interesting, from a numerical analysis point of view, to modify slightly the algorithm presented in this paper by introducing a threshold $S > 0$ and replace in problem (3) the distance function d_E with a truncated distance function $(-S \wedge d^E) \vee S$. Almost all of the results presented in this paper would remain identical (in particular, the consistency still holds). Only the comparison results in section 4.2 are not completely clear, since they rely on comparisons of the true distance functions. However, it is reasonable to believe that in the limit $h \rightarrow 0$ also these results hold.

4 Convergence of the algorithm

In this section, we study the limit of the iterates of our variational algorithm, as the time-step goes to zero. First (Section 4.1), under some local boundedness assumption on the forcing term, we find a limit which has some continuity properties (Prop. 4.3 and 4.4), and we show that it is consistent with other notions of weak solutions (Cor. 4.6, and with further assumptions on the forcing term, Prop. 4.8 and 4.9).

Then in 4.2, assuming some spatial regularity (18) on the forcing term, we find a comparison principle for our limits (however, with some limitation), and it allows to build level-sets solutions starting from an arbitrary bounded uniformly continuous function (Section 4.3).

4.1 The discrete flow and its limit

Given $E \subset \mathbb{R}^N$, closed with compact boundary, and $h > 0$, we define the “tube” $E_h \subset \mathbb{R}^N \times [0, +\infty)$ as follows

$$E_h(t) := T_{[\frac{t}{h}]h - h, [\frac{t}{h}]h} \cdots T_{2h, 3h} T_{h, 2h} T_{0, h}(E), \quad (7)$$

where $[x]$ denotes the integer part of x . We then define $E_h := \bigcup_{t \geq 0} E_h(t) \times \{t\}$.

There exists a sequence $(h_n)_{n \geq 1}$ such that both E_{h_n} and $\mathbb{R}^N \times [0, +\infty) \setminus E_{h_n} = {}^c E_{h_n}$ converge in the Hausdorff distance (locally in time) to E^* and ${}^c E_*$ respectively. Such convergence is equivalent to the locally uniform convergence, in $\mathbb{R}^N \times [0, +\infty)$, of the distance functions $\text{dist}((x, t), E_{h_n})$ and $\text{dist}((x, t), {}^c E_{h_n})$ to the distance functions $\text{dist}((x, t), E^*)$ and $\text{dist}((x, t), {}^c E_*)$, see [11, App. A]. In particular,

for any $(x, t) \in E^*$ (resp., ${}^c E_*$), there exists $(x_n, t_n) \in E_{h_n}$ (resp., ${}^c E_{h_n}$) such that $(x_n, t_n) \rightarrow (x, t)$, and if $(x_n, t_n) \in E_{h_n}$ (resp., ${}^c E_{h_n}$) and converge to some point $(x, t) \in \Omega \times [0, +\infty)$, then $(x, t) \in E^*$ (resp., ${}^c E_*$). In the sequel we denote by $(h)_{h>0}$ the sequence $(h_n)_{n \geq 1}$.

Clearly, E_* is open while E^* is closed, and $E_* \subset E^*$. For any $t \geq 0$, we denote by $E^*(t)$ (resp. $E_*(t)$) the section $\{x : (x, t) \in E^*\}$ (resp. $\{x : (x, t) \in E_*\}$). For any $t \geq 0$, we let $\Gamma(t) = E^*(t) \setminus E_*(t)$, which in some sense is our generalized evolution starting from ∂E .

From the definition of E_* , E^* it follows

$$E_*(0) \subseteq \text{int}(E), \quad E \subseteq E^*(0), \quad (8)$$

in particular $\Gamma(0) \supseteq \partial E$. If $F(t)$ is a superflow on $[t_0, t_1]$ such that $F(t_0) \subset E_*(t_0)$, since $(F(t))$ being assumed to be closed for any time t , and E_* being open) $\text{dist}(F(t_0) \times \{t_0\}, {}^c E_*) > 0$, one sees that for h small enough, $F(t_0) \subset E_h([t_0/h]h) \cap E_h([t_0/h]h + h)$. It then follows from Theorem 3.3 that (if h is enough small) $F(t) \subset E_h(t)$ for any $t \in [t_0, t_1]$, and passing to the limit we get $F \subset E_* \cap (\mathbb{R}^N \times [t_0, t_1])$. Hence E_* satisfies a comparison principle for superflows that start inside and analogously E^* satisfies a comparison principle for subflows starting outside, so that we have shown the following:

Proposition 4.1. *The set-valued functions $E_*(\cdot)$ and $E^*(\cdot)$ are barriers on $[0, +\infty)$, with respect to \mathcal{F}^+ and \mathcal{F}^- respectively, that is: $E_* \in \mathcal{B}^+$ and $E^* \in \mathcal{B}^-$. In particular, we get from Definition 2.3 that for all $t \geq 0$,*

$$E_*(t) \supseteq \mathcal{M}_*(E_*(0))(t) \quad \text{and} \quad E^*(t) \subseteq \mathcal{M}^*(E^*(0))(t). \quad (9)$$

On the other hand, it is not clear if we also have $E_* \in \mathcal{B}^-$ and $E^* \in \mathcal{B}^+$.

Let us now show that $E_*(0) = \text{int}(E)$ and $E^*(0) = E$, so that we can substitute $E_*(0)$ and $E^*(0)$ with E in (9). In order to do so, we further require that the function G satisfies the following regularity assumption: for any $T > 0$ there exists a constant $C(T)$ such that

$$\left| \frac{G(x, s) - G(y, s) - G(x, t) + G(y, t)}{s - t} \right| \leq C(T), \quad (10)$$

for any $s, t \leq T$. Note that this is equivalent to require that G can be written as the sum $G(x, t) = G_1(t) + G_2(x, t)$, with $G_1 \in C^0([0, +\infty))$ and $G_2 \in \text{Lip}_{\text{loc}}([0, +\infty); L^\infty(\mathbb{R}^N))$.

We first construct explicit super/subflows starting from a Wulff shape $W_\phi(x_0, r)$ of radius $r > 0$ (or its complement), at time $t \geq 0$. More precisely, we construct superflows $W_{x_0, t, r}^+(s)$, with $s \in [t, +\infty)$, starting from $W_\phi(x_0, r)$ at time t , which are smooth on $[t, t + \tau]$ and vanish after the time $t + \tau$, where the duration τ depends only on r , and such that for all $h > 0$, $s \geq t$, we have $T_{s, s+h}(W_{x_0, t, r}^+(s)) \supseteq W_{x_0, t, r}^+(s + h)$.

Lemma 4.2. *Let $x_0 \in \mathbb{R}^N$ and $r > 0$. We consider the functions*

$$d_{\pm}(x, s) := \pm \left(\phi(x - x_0) - r + \frac{s - t}{2\tau} r \right) + G(x_0, s) - G(x_0, t), \quad (11)$$

for $(x, s) \in \mathbb{R}^n \times [t, t + \tau]$, where τ is such that

$$\omega_{G, t+\tau}(\tau) \leq \frac{r}{4}, \quad (12)$$

$$\tau \leq \frac{r^2}{2(C(t + \tau) + 4(N - 1))} \wedge \frac{r^2}{16(N + 1)}, \quad (13)$$

where $\omega_{G, t+\tau}$ is as before a modulus of continuity of G on $[0, t + \tau]$ and $C(\cdot)$ is the constant appearing in (10).

Let $W_{x_0, t, r}^+(s) := \{d_+(\cdot, s) \leq 0\}$ when $s \in [t, t + \tau]$, and $W_{x_0, t, r}^+(s) := \emptyset$ for $s > t + \tau$. Then, for any $s \geq t$, $h > 0$, we have $T_{s, s+h}(W_{x_0, t, r}^+(s)) \supseteq W_{x_0, t, r}^+(s + h)$. On the other hand, if $W_{x_0, t, r}^-(s) := \{d_-(\cdot, s) \leq 0\}$ when $s \in [t, t + \tau]$, and $W_{x_0, t, r}^-(s) := \mathbb{R}^N$ for $s > t + \tau$, then, for any $s \geq t$, $h > 0$, we have $T_{s, s+h}(W_{x_0, t, r}^-(s)) \subseteq W_{x_0, t, r}^-(s + h)$.

Notice that, letting $\tau(r)$ be the maximal time τ satisfying (12) and (13) for a given $r > 0$, we have $\tau(r) > 0$ and

$$\lim_{r \rightarrow \infty} \tau(r) = +\infty.$$

Notice also that the condition $\omega_{G, t+\tau}(\tau) \leq \frac{r}{4}$ ensures the inclusion $W_{\phi}(x_0, r/4) \subseteq W_{x_0, t, r}^+(s)$ for $s \in [t, t + \tau]$, hence in particular the set $W_{x_0, t, r}^+(s)$ is nonempty. In fact, one could check that $W_{x_0, t, r}^+$ is a superflow (in the sense of Definition 2.1) on $[t, t + \tau]$, while $W_{x_0, t, r}^-$ is a subflow. Then, the thesis would follow from Theorem 3.3, at least for h small enough. The statement of Lemma 4.2 is slightly more precise, as it holds without any restriction on h so that the superflow property which is shown is, in particular, uniform in x_0 .

Proof. Let $s \in [t, t + \tau]$ and $h > 0$. If $s + h > t + \tau$, $W_{x_0, t, r}^+(s) = \emptyset$ so that the statement is obvious, hence we may assume $s + h \leq t + \tau$. For any $x \in \mathbb{R}^N$, by (10) we have

$$\begin{aligned} d_+(x, s) &+ G(x, s + h) - G(x, s) \\ &= \phi(x - x_0) - r + G(x_0, s) - G(x_0, t) + G(x, s + h) - G(x, s) + \frac{s - t}{2\tau} r \\ &\leq \phi(x - x_0) - r + G(x_0, s + h) - G(x_0, t) + \frac{s - t}{2\tau} r + C(t + \tau)h. \end{aligned}$$

Let now Ω be an open and bounded subset of \mathbb{R}^N which is big enough to guarantee that the set $T_{s, s+h}(W_{x_0, t, r}^+(s))$ does not depend on Ω (Proposition 3.1). Hence the solution w of (3), with E replaced by $W_{x_0, t, r}^+(s)$ and (t, s) replaced by $(s, s + h)$, is less than the solution of

$$\begin{aligned} &\min_{v \in BV(\Omega)} \int_{\Omega} \phi^{\circ}(Dv) \\ &+ \frac{1}{2h} \int_{\Omega} \left(v - \phi(x - x_0) + r - G(x_0, s + h) + G(x_0, t) - \frac{s - t}{2\tau} r - C(t + \tau)h \right)^2 dx, \end{aligned}$$

which, in turn, is less (as shown in the proof of Lemma 3.2) than the function

$$\begin{cases} \phi(x - x_0) + h \frac{N-1}{\phi(x - x_0)} - r + G(x_0, s+h) - G(x_0, t) + \frac{s-t}{2\tau}r + C(t+\tau)h \\ \quad \text{if } \phi(x - x_0) \geq \sqrt{h(N+1)}, \\ \sqrt{h} \frac{2N}{\sqrt{N+1}} - r + G(x_0, s+h) - G(x_0, t) + \frac{s-t}{2\tau}r + C(t+\tau)h \\ \quad \text{otherwise.} \end{cases}$$

Hence, we see that

$$w(x) \leq d_+(x, s+h) + h \frac{N-1}{\phi(x - x_0)} - h \frac{r}{2\tau} + C(t+\tau)h$$

when $\phi(x - x_0) \geq \sqrt{h(N+1)}$. Now, since $\tau \leq r^2/(16(N+1))$ and $h \leq \tau$, we get $r/4 \geq \sqrt{h(N+1)}$ so that we can replace the last condition with the stronger condition $\phi(x - x_0) \geq r/4$. On the other hand, if both $\phi(x - x_0) \geq r/4$ and $\tau \leq r^2/(2(C(t+\tau) + 4(N-1)))$, then

$$\frac{r}{2\tau} \geq \frac{C(t+\tau) + 4(N-1)}{r} \geq C(t+\tau) + \frac{N-1}{\phi(x - x_0)},$$

so that $w(x) \leq d_+(x, s+h)$. This shows that $T_{s,s+h}(W_{x_0,t,r}^+(s)) \supseteq W_{x_0,t,r}^+(s+h)$.

The proof of the similar thesis for $W_{x_0,t,r}^-$ is analogous. \square

By the previous lemma, if $E_h(t) \supseteq W_\phi(x_0, r)$, then $E_h(t+nh) \supseteq W_{x_0,t,r}^+(t+nh)$ for any $n \geq 1$ and, in particular, $E_h(t+nh) \supseteq W_\phi(x_0, r/4)$ as long as $nh \leq \tau(r)$. In other words, for any $r > 0$, then $nh \leq \tau(r)$ yields

$$E_h(t+nh) \supseteq \{x : d_{E_h(t)}(x) \leq -r\}.$$

This shows that the (discrete) evolution will not vanish “suddenly”, in the sense that any subset in the interior of $E_h(t)$ remains in the interior of $E_h(s)$ for $s > t$ sufficiently close to t .

We can easily deduce the same “semicontinuity” property for E_* : indeed, when $d_{E_*(t)}(x) = -r$, then, for any $r' < r$, $W_\phi(x, r') \subset E_h(t)$ as soon as h is small enough, so that $W_\phi(x, r'/4) \subset E_h(t + [\tau/h]h)$ for all $\tau < \tau(r')$. Letting first $h \rightarrow 0$ and then $r' \rightarrow r$, we find that, if $\tau < \tau(r)$, it follows

$$E_*(t+\tau) \supseteq \{x : d_{E_*(t)}(x) \leq -r\}. \quad (14)$$

In the same way we obtain

$$E^*(t+\tau) \subseteq \{x : d_{E^*(t)}(x) < r\}. \quad (15)$$

Moreover, one can easily verify that the same properties hold at $t = 0$ with $E^*(t)$ replaced with E and $E_*(t)$ replaced with $\text{int}(E)$, where E is the initial set. From (14) and (15) we also get

$$\Gamma(t+\tau) \subseteq \{\text{dist}_\phi(\cdot, \Gamma(t)) < r\}. \quad (16)$$

As a consequence, we obtain the following semicontinuity property for the tubes E_*, E^* .

Proposition 4.3. *Assume that G satisfies (10). Let E be a closed subset of \mathbb{R}^N with compact boundary. Let O, F be an open and a closed subset of \mathbb{R}^N respectively. Let $t \geq 0$ and let $(\tau_n)_{n \geq 0}$ be a sequence of nonnegative numbers going to 0. Then*

- *If ${}^c E_*(t + \tau_n) \rightarrow {}^c O$ in the Hausdorff sense, $E_*(t) \subseteq O$,*
- *If $E^*(t + \tau_n) \rightarrow F$ in the Hausdorff sense, $F \subseteq E^*(t)$,*

In particular, we deduce that any Hausdorff limit of a sequence $\Gamma(t + \tau_n)$ is contained in $\Gamma(t)$. Moreover, if $t = 0$, we can replace $E_(t)$ with $\text{int}(E)$ in the first statement and $E^*(t)$ with E in the second. In particular, choosing $\tau_n \equiv 0$, we get*

$$\text{int}(E) \subseteq E_*(0) \subseteq E^*(0) \subseteq E,$$

which implies, recalling (8), that

$$E_*(0) = \text{int}(E) \text{ and } E^*(0) = E. \quad (17)$$

Notice that (17) shows that $\Gamma(0) = \partial E$.

Since $O \subset F$ in the above proposition, we see also that if $E^*(t) = \overline{E_*(t)}$, then $E^*(t + \tau) \rightarrow E^*(t)$ in the Hausdorff sense as $\tau \rightarrow 0$, whereas if $E_*(t) = \text{int}(E^*(t))$, then ${}^c E_*(t + \tau) \rightarrow {}^c E_*(t)$, and if both are true, then $\Gamma(t + \tau) = E^*(t + \tau) \setminus E_*(t + \tau) \rightarrow \partial E^*(t)$ as $\tau \rightarrow 0$. To show this one just needs to show that for any $x \in \partial E^*(t)$, there exists $x_\tau \in \Gamma(t + \tau)$ that converge to x as $\tau \rightarrow 0$. We know that there exists $y_\tau \in E^*(t + \tau)$ and $z_\tau \notin E_*(t + \tau)$ such that both y_τ and z_τ converge to x . Then, the segment $[y_\tau, z_\tau]$ must intersect the set $\Gamma(t + \tau)$ and any point x_τ in this intersection will satisfy the desired property.

Notice also that, if E is such that $E = \overline{\text{int}(E)}$, we deduce that $\Gamma(t) = E^*(t) \setminus E_*(t)$ converges to ∂E as $t \rightarrow 0$, in the Hausdorff sense.

The left continuity of the tubes E_*, E^* is given by the following proposition.

Proposition 4.4. *Assume that G satisfies (10). Let E be a closed subset of \mathbb{R}^N with compact boundary, and let $t > 0$. Then ${}^c E_*(t - \tau) \rightarrow {}^c E_*(t)$ in the Hausdorff sense as $\tau \rightarrow 0$, $\tau \geq 0$, while $E^*(t - \tau) \rightarrow E^*(t)$. Moreover, $\Gamma(t - \tau) \rightarrow \Gamma(t)$.*

Proof. We sketch the proof of this proposition. As for the previous proposition, one will deduce from (14) that if ${}^c E_*(t - \tau_n) \rightarrow {}^c O$ in the Hausdorff sense, along a subsequence τ_n going to 0, then $O \subseteq E_*(t)$. On the other hand, since ${}^c E_*$ is closed in $\Omega \times [0, +\infty)$, one must have ${}^c O \subseteq {}^c E_*(t)$. Thus $O = E_*(t)$ and the thesis follows. In the same way, (15) yields that if $E^*(t - \tau_n) \rightarrow F$ in the Hausdorff sense, then $E^*(t) \subseteq F$. From the closedness of F^* we conclude in the same way that $F = E^*(t)$.

The last assertion follows from (16): first of all, any Hausdorff limit F of a subsequence $\Gamma(t - \tau_n) = E^*(t - \tau_n) \setminus E_*(t - \tau_n)$ is inside $\Gamma(t) = E^*(t) \setminus E_*(t)$, by the

previous results. Now, since $\text{dist}_\phi(\cdot, \Gamma(t - \tau_n))$ converges uniformly to $\text{dist}_\phi(\cdot, F)$ as $n \rightarrow \infty$, $\Gamma(t) \subseteq \{\text{dist}_\phi(\cdot, F) \leq r\}$ for any $r > 0$, hence it lies in F . Thus $F = \Gamma(t)$. \square

Remark 4.5. Notice that in general we cannot expect the maps $t \mapsto E_*(t)$ and $t \mapsto E^*(t)$ to be continuous in the Hausdorff distance: indeed, this would prevent small disconnected parts from disappearing in finite time, a phenomenon which is known to happen even when $G \equiv 0$. On the other hand, these maps are likely to be continuous in the L^1 -topology, under suitable assumptions on G (when $G \equiv 0$ it is proved in [1, Theorem 4.4]).

From Propositions 4.1 and 4.3, and in particular (9) and (17), we get the following corollary.

Corollary 4.6. *If G satisfies condition (10), we have*

$$\Gamma(t) \subseteq \mathcal{N}^*(E, 0)(t) \setminus \mathcal{M}_*(E, 0)(t).$$

In particular, as long as $\mathcal{N}^(E, 0)(t) \setminus \mathcal{M}_*(E, 0)(t)$ has no interior (nonfattening condition), then the motions $E^*(t)$ and $E_*(t)$ are uniquely defined and do not depend on the sequence along which the limits are obtained.*

Remark 4.7. Notice that, as long as the set E has compact boundary, all the results of this section can be easily extended to functions G which are only locally bounded in x , i.e. $G(x, t) = G_1(t) + G_2(x, t)$, with $G_1 \in C^0([0, +\infty))$ and $G_2 \in \text{Lip}_{\text{loc}}([0, +\infty); L_{\text{loc}}^\infty(\mathbb{R}^N))$.

Proposition 4.8. *If $G(x, t) = \int_0^t g(x, s) ds$, with g continuous, then $\Gamma(t)$ is contained in the zero level-set of the corresponding viscosity solution.*

Proof. It follows immediately from Theorem 2.4 and Corollary 4.6. \square

From Corollary 4.6 and from [17, Section 3] we also have the following consistency result in the case of an x -independent forcing term.

Proposition 4.9. *Let $G(x, t) = G(t) \in C^0([0, +\infty))$ and let $\phi(x) = |x|$ (i.e. isotropic mean curvature flow). Then $\Gamma(t)$ is contained in the minimal barrier solution defined in [17]. In particular, if ∂E is of class $C^{2,\alpha}$, $E^*(t) = \overline{E_*(t)}$ and $\Gamma(t) = \partial E^*(t)$ coincides with the unique (local in time) solution of (1) given in [17].*

Remark 4.10. As already pointed out in the Introduction, the viscosity theory can be applied under more general assumptions on G than what is required in Proposition 4.8 (see [23, 24]). However, it is still not clear which is the relation between the limit set $\Gamma(t)$ and the zero level-set of such viscosity solutions, except for the particular case of an x -independent forcing term, where the equality holds for small times as a consequence of Proposition 4.9 (if ∂E is regular enough).

4.2 An inclusion principle

Let us now consider the case where the driving term is the “time-derivative” of a function $G(x, t)$ that satisfies

$$\left| \frac{G(x, s) - G(y, s) - G(x, t) + G(y, t)}{s - t} \right| \leq C(T)|x - y|. \quad (18)$$

This condition is stronger than condition (10) (see also Remark 4.7) and is for instance true whenever $G(x, t) = G_1(t) + G_2(x, t)$, with $G_1 \in C^0([0, +\infty))$ and $G_2 \in C^1([0, +\infty); \text{Lip}(\mathbb{R}^N))$. In particular, all the results of Section 4.1 still hold under assumption (18).

Given a closed set $E \subset \mathbb{R}^N$ with nonempty compact boundary ∂E , we define the maximal existence time $T_E^* \in [0, +\infty]$ for the flow E^* as the supremum of all times t such that $E^*(t) \neq \emptyset$ and $E_*(t) \neq \mathbb{R}^N$. The fact that $T_E^* > 0$ is ensured by Proposition 4.3, whenever $\text{int}(E) \neq \emptyset$.

Consider now two closed sets E^1 and E^2 , with nonempty compact boundary, and assume $E^1 \subset E^2$ and $D := \text{dist}_\phi(\partial E^1, \partial E^2) > 0$. Notice that, if G depends only on time, then for each z such that $\phi(z) \leq D$, we have $z + E^1 \subset E^2$, so that $T_{0,h}(z + E^1) \subset T_{0,h}(E^2)$ for any $h > 0$. Since G does not depend on x , we get $T_{0,h}(z + E^1) = z + T_{0,h}(E^1)$. It follows $W_\phi(0, D) + T_{0,h}(E^1) \subset T_{0,h}(E^2)$, which implies $\text{dist}_\phi(\partial T_{0,h}(E^1), \partial T_{0,h}(E^2)) \geq D$. By induction, we deduce that $\text{dist}_\phi(\partial E_h^1(t), \partial E_h^2(t)) \geq D$ for any $t \geq 0$ (where we set the distance equal to $+\infty$ if one of the two sets disappears).

For a general G the estimate is slightly trickier, even if it follows the same idea. Assume $T^* = \min\{T_{E^1}^*, T_{E^2}^*\} > 0$ and let $T < T^*$. By Proposition 3.1, we can find a “large” bounded open set $\Omega \subset \mathbb{R}^N$ such that the sets $E_h^1(t)$ and $E_h^2(t)$ defined in (7) do not depend on Ω , for $t \in [0, T]$ and h small enough. In particular, we can assume that $\partial E_h^1(t)$ and $\partial E_h^2(t)$ remain at a positive distance from $\partial\Omega$ for any $t \in [0, T]$. Let w_1, w_2 be the solutions of the variational problem (3), for $t = 0$ and $s = h$, with d_E replaced by d_{E^1}, d_{E^2} respectively. Notice that, for $z \in \mathbb{R}^N$, the set $z + T_{0,h}(E^1) = z + \{w_1 < 0\}$ coincides with the set $\{x \in z + \Omega : w_1(x - z) < 0\}$, and the function $\tilde{w}_1(x) = w_1(x - z)$, defined in $z + \Omega$, is the solution of

$$\min_{w \in L^2(z + \Omega)} \int_{z + \Omega} \phi^\circ(Dw) + \frac{1}{2h} \int_{z + \Omega} (w(x) - d_{E^1}(x - z) - G(x - z, h) + G(x - z, 0))^2 dx.$$

Possibly enlarging the set Ω , we can assume that both functions w_2 and \tilde{w}_1 are solutions of their respective variational problems in the same domain (for instance, $\Omega \cup (z + \Omega)$). Then, since $d_{E^1}(x - z) \geq d_{E^1}(x) - \phi(z) \geq d_{E^2}(x) + D - \phi(z)$ and, using (18), $-G(x - z, h) + G(x - z, 0) \geq -G(x, h) + G(x, 0) - C(T)|z|h$, one finds that $\tilde{w}_1 \geq w_2 + D - \phi(z) - C(T)|z|h$. In particular, if $\phi(z) \leq D/(1 + hC'(T))$ with $C'(T) = C(T) \sup_{z \neq 0} |z|/\phi(z)$, we get $\{\tilde{w}_1 \leq 0\} \subset \{w_2 \leq 0\}$, which in turn implies $\text{dist}_\phi(\partial T_{0,h}(E^1), \partial T_{0,h}(E^2)) \geq D/(1 + hC'(T))$. By an induction argument, we deduce that $\text{dist}_\phi(\partial E_h^1(t), \partial E_h^2(t)) \geq D(1 + hC'(T))^{-[t/h]}$ for any $t \in [0, T - h]$.

and $h > 0$ small enough. Observe that, as $h \rightarrow 0$, we have $D(1 + hC'(T))^{-[t/h]} \rightarrow De^{-C'(T)t}$. We will show that this estimate also holds in the limit, for the motions $(E^1)^*$ and $(E^2)_*$, obtained along the *same* subsequence $(h_k)_{k \geq 1}$ (which we will still denote by $(h)_{h>0}$).

Fix $\delta < De^{-C'(T)T}$. Then, if h is small enough, we have $\delta \leq \text{dist}_\phi(E_h^1(t), \Omega \setminus E_h^2(t))$ for any $t \in [0, T]$. Given a fixed $t < T$, choose a subsequence (h_k) such that both Hausdorff limits of $E_{h_k}^1(t)$ and $\Omega \setminus E_{h_k}^2(t)$ exist in $\overline{\Omega}$, and denote by K and L , respectively, these limits. Since $\text{dist}_\phi(E_h^1(t), \Omega \setminus E_h^2(t)) \geq \delta$, in the limit we find $\text{dist}_\phi(K, L) \geq \delta$. We also have $K \subseteq (E^1)^*(t)$ and $L \subseteq \overline{\Omega} \setminus (E^2)_*(t)$. Define now $K^{\delta/2} = K + W_\phi(0, \delta/2)$, which has its boundary between ∂K and ∂L and lies at distance at least $\delta/2$ from both boundaries. Let $\delta' < \delta$ and set $\delta'' = (\delta + \delta')/2$. If $x \in \partial K^{\delta/2}$, then $W_\phi(x, \delta''/2) \subseteq \overline{\Omega} \setminus (K \cup L)$ so that if h_k is small enough, $W_\phi(x, \delta''/2) \subset E_{h_k}^2(t)$ and $W_\phi(x, \delta''/2) \cap E_{h_k}^1(t) = \emptyset$. By Lemma 4.2 there exists $\tau > 0$, depending only on δ'' and $\delta' < \delta''$, such that $W_\phi(x, \delta'/2) \subset E_{h_k}^2(s)$ and $W_\phi(x, \delta'/2) \cap E_{h_k}^1(s) = \emptyset$ for all $t \leq s < t + \tau$. In the limit, this implies that for $t \leq s < t + \tau$, $W_\phi(x, \delta'/2) \subset (E^2)_*(s)$ and $(E^1)^*(s) \cap W_\phi(x, \delta'/2) = \emptyset$. Since x is an arbitrary point of $\partial K^{\delta/2}$, this implies that both $\text{dist}_\phi(\partial K^{\delta/2}, \partial(E^1)^*(s)) \geq \frac{\delta'}{2}$ and $\text{dist}_\phi(\partial K^{\delta/2}, \partial(E^2)_*(s)) \geq \frac{\delta'}{2}$. We deduce that for any s in $(t, t + \tau)$, $\text{dist}_\phi(\partial(E^1)^*(s), \partial(E^2)_*(s)) \geq \delta'$. Since t is arbitrary in $[0, T]$ and τ does not depend on t , we deduce that in fact for any $t \in [0, T]$, $\text{dist}_\phi(\partial(E^1)^*(t), \partial(E^2)_*(t)) \geq \delta'$. (The case $t = 0$ follows directly from Proposition 4.3.) We may send $\delta' \rightarrow \delta$ to see that the inequality holds with δ instead of δ' . In fact, we can deduce from the previous argument that the distance between the two sets decreases at most like $De^{-C'(T)t}$, $t \in [0, T]$.

We have obtained the following result.

Proposition 4.11. *Let $E^1 \subset E^2 \subset \mathbb{R}^N$ be two closed sets with nonempty compact boundary, and assume $\text{dist}_\phi(\partial E^1, \partial E^2) > 0$. Denote by E_h^1 and E_h^2 , $h > 0$, the corresponding discrete evolutions in $\mathbb{R}^N \times [0, +\infty)$. Let (h_k) be a subsequence such that $E_{h_k}^1 \rightarrow (E^1)^*$, and ${}^c E_{h_k}^2 \rightarrow {}^c (E^2)_*$ in the Hausdorff sense. Assume also that $E_{h_k}^2 \rightarrow (E^2)^*$. Then, for any $t \geq 0$ we have $(E^1)^*(t) \subseteq (E^2)_*(t)$.*

In particular, for any $t \in [0, T]$, with $T := \max\{T_{E^1}^, T_{E^2}^*\}$, we have, for any t' with $t < t' < T$,*

$$\text{dist}_\phi(\partial(E^1)^*(t), \partial(E^2)_*(t)) \geq \text{dist}_\phi(\partial E^1, \partial E^2)e^{-C'(t')t} > 0 \quad (19)$$

where C' is proportional to the constant in (18).

Remark 4.12. Notice that the inclusions $(E^1)^*(t) \subseteq (E^2)^*(t)$ and $(E^1)_*(t) \subseteq (E^2)_*(t)$ always hold (without any assumption on $\text{dist}_\phi(\partial E^1, \partial E^2)$).

Remark 4.13. We remark that we do not know if the thesis of Proposition 4.11 still holds if $(E^1)^*$ and $(E^2)^*$ are limits of $E_{h_k}^1$ and $E_{h_j}^2$ (respectively) along different subsequences. This would be an important result, yielding for instance the uniqueness of the level-set solution $u(x, t)$, defined in Section 4.3.

4.3 The level-set approach

Consider now a function $u_0 \in BUC(\mathbb{R}^N)$, such that for each $t \in \mathbb{R}$, the level-set $\partial\{u_0 > t\}$ is bounded. For all $q \in \mathbb{Q}$ consider the level-sets $E^q := \{u_0 \geq q\}$ and let $E_h^q \subset \mathbb{R}^N \times [0, +\infty)$ be the discrete evolutions of E^q . Then, a diagonal argument shows that, along a subsequence $(h_k)_{k \geq 1}$, we have $E_{h_k}^q \rightarrow (E^q)^*$ and ${}^c E_{h_k}^q \rightarrow {}^c(E^q)_*$ locally in the Hausdorff sense, i.e. the distance functions $\text{dist}(\cdot, E_{h_k}^q)$ and $\text{dist}(\cdot, \mathbb{R}^N \times [0, +\infty) \setminus E_{h_k}^q)$ converge to $\text{dist}(\cdot, (E^q)^*)$ and $\text{dist}(\cdot, \mathbb{R}^N \times [0, +\infty) \setminus (E^q)_*)$ respectively, uniformly in $\mathbb{R}^N \times [0, T]$ for any $T > 0$.

Observe that (Remark 4.12) for each $q, r \in \mathbb{Q}$ with $q \geq r$, we have $(E^q)^*(t) \subseteq (E^r)^*(t)$ and $(E^q)_*(t) \subseteq (E^r)_*(t)$, for any $t \geq 0$. Hence we can define two functions $u^*, u_* : \mathbb{R}^N \times [0, +\infty) \rightarrow \mathbb{R}$, by letting

$$u^*(x, t) := \sup\{q \in \mathbb{Q} : x \in (E^q)^*(t)\}, \quad u_*(x, t) := \sup\{q \in \mathbb{Q} : x \in (E^q)_*(t)\}.$$

By Proposition 4.11, we know that $(E^q)^*(t) \subset (E^r)_*(t)$ for any $t \geq 0$ whenever $q > r$, which implies $u^*(x, t) = u_*(x, t)$ for any $(x, t) \in \mathbb{R}^N \times [0, +\infty)$. Indeed, if $q > u^*(x, t)$, then $x \notin (E^q)^*(t)$, so that it is neither in $(E^q)_*(t)$, hence $u_*(x, t) \leq u^*(x, t)$; on the other hand, if $q > u_*(x, t)$, then if $q' \in (u_*(x, t), q) \cap \mathbb{Q}$, $x \notin (E^{q'})_*(t) \supset (E^q)^*(t)$, hence $x \notin (E^q)^*(t)$: we deduce $u^*(x, t) \leq u_*(x, t)$. We simply denote by $u(x, t)$ this common value.

Let us observe that, for each $t \geq 0$, from Proposition 4.11 (more exactly from the estimate (19)) it follows that $u(\cdot, t)$ is uniformly continuous on \mathbb{R}^N (with the same modulus of continuity as u_0 if $C(T) = 0$ in (18)). It also follows easily from Propositions 4.3 and 4.4 that if $(x_n, t_n) \rightarrow (x, t)$, then $u(x_n, t_n) \rightarrow u(x, t)$: indeed, for instance, one sees that if $u(x_n, t_n) < q$ for n large enough, then $x_n \notin (E^q)_*(t_n)$, hence in the limit $x \notin (E^q)_*(t)$ so that $u(x, t) \leq q$. This means that the function u is globally continuous on $\mathbb{R}^N \times [0, +\infty)$. In particular, we have the inclusions $(E^q)^*(t) \subset (E^s)_*(t)$ for any $t \geq 0$ whenever $q > s$, $q, s \in \mathbb{R}$. We deduce easily that (letting now $\Gamma^s(t) := (E^s)^*(t) \setminus (E^s)_*(t)$)

$$\bigcup_{t \geq 0} \Gamma^s(t) = (E^s)^* \setminus (E^s)_* \subseteq \{(x, t) : u(x, t) = s\} \quad (20)$$

Now, let $N := \{s \in \mathbb{R} : |\{(x, t) \in \mathbb{R}^N \times [0, +\infty) : u(x, t) = s\}| = 0\}$. The set N is at most countable. If $s \notin N$, then (20) is in fact an equality: one has $\{u > s\} = (E^s)_*$ and $\{u \geq s\} = (E^s)^*$. One can deduce that $\partial E_{h_k}^s$ converges to $\{u = s\}$ in the local Hausdorff sense. For these values of s , the flow defined by our algorithm is a “true” evolution of hypersurfaces. Indeed, at any time $t \geq 0$, we can show that $\{u(\cdot, t) = s\}$ has empty interior. Otherwise, there would exist $W_\phi(x, \rho) \subseteq \{u(\cdot, t) = s\}$. In particular, if $q > s > q'$, $q, q' \in \mathbb{Q}$, we would have $W_\phi(x, \rho) \subseteq (E^{q'})_*(t)$ while $W_\phi(x, \rho) \cap (E^q)^*(t) = \emptyset$. By (14) and (15), it would follow that, if $t \leq t' < t + \tau_{t+1}(\rho/2)$, then $W_\phi(x, \rho/2) \subset (E^{q'})_*(t)$ and $W_\phi(x, \rho/2) \cap (E^q)^*(t) = \emptyset$, which would imply that $\{u = s\}$ has nonempty interior, leading to a contradiction. We

can not prove in general the uniqueness of the flow $(E^s)^*$, since it could depend on the subsequence (h_k) along which the first limits have been taken. If on the contrary $s \in N$, then a *fattening* of the corresponding level-set happens, and we can only deduce the inclusion (20). As in the case of classical level-set solutions, we expect nonuniqueness of the limit flow in this situation (and in this situation *only*).

If $G(x, t) = \int_0^t g(x, s) ds$ with g continuous, then (by Proposition 4.8), one sees that u is the unique viscosity solution [15] of

$$\frac{\partial u}{\partial t} = \phi^\circ(\nabla u) (\operatorname{div} \nabla \phi^\circ(\nabla u) + g). \quad (21)$$

In particular, in this case, the limit u is the same along any subsequence. We can then deduce that $\partial E_h^s(t) \rightarrow \{u(\cdot, t) = s\}$ as $h \rightarrow 0$, for each level $s \notin N$, or each time t before the moment the level $s \in N$ fattens.

In case G is an arbitrary driving term satisfying (18), we conjecture that our u is still the viscosity solution of

$$\frac{\partial u}{\partial t} = \phi^\circ(\nabla u) \left(\operatorname{div} \nabla \phi^\circ(\nabla u) + \frac{\partial G}{\partial t} \right) \quad (22)$$

built by Lions and Souganidis [23, 24]. However, to show this, we would need either to show the stability of our construction under small perturbations of G (that would allow us to approximate G with smooth functions), or a comparison result like Theorem 2.4 between barriers and viscosity solutions (in the sense of [23, 24]) and then use Corollary 4.6.

Let us eventually make a few remarks. We first observe that our construction can still be performed if ϕ and ϕ° are nonsmooth: typically, in the *crystalline* case, where the Wulff shape $\{\phi \leq 1\}$ is a polyhedron. In this case, our proof of consistency does not hold (neither is clear how to extend the definition of a sub/superflow). On the other hand, most of the results are still valid, including the comparison principle in Proposition 4.11, and the construction of the level-set function u starting from u_0 still makes sense.

We also mention that in the convex case, if $G = G(t)$, by the same arguments as in [11] we can show that the evolution (defined in \mathbb{R}^N) remains convex for all time, including when the anisotropy is nonsmooth. We also expect that the results in [4] still hold with similar proofs, and that a unique “regular” evolution can be defined for small times as the unique limit of our algorithm, when the initial convex set satisfies an interior εW_ϕ -condition. This would in turn yield the uniqueness of the level-set function defined above.

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