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# Optimal Force Generation in Parallel Manipulators for Passing through the Singular Positions 

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#### Abstract

It is known that a parallel manipulator at a singular configuration can gain one or more degrees of freedom and become uncontrollable. That is it might not reproduce a stable motion along a prescribed trajectory. However, it is proved experimentally that there is possible passing through the singular zones. This was simulated and shown on numerical examples and illustrated on several parallel structures. In this paper, we determine the optimal dynamic conditions generating a stable motion inside the singular zones. The obtained results show that the general condition for passing through a singularity can be defined as follows: the end-effector of the parallel manipulator can pass through the singular positions without perturbation of motion if the wrench applied on the end-effector by the legs and external efforts of the manipulator are orthogonal to the twist along the direction of the uncontrollable motion. This condition is obtained from the inverse dynamics and analytically demonstrated by the study of the Lagrangian of a general parallel manipulator. Numerical simulations are carried out using the software ADAMS and validated by experimental tests.

Index Terms - parallel manipulators, singularity, dynamics, force management, trajectory planning.


## 1. Introduction

Parallel manipulators have experienced an increase in popularity in recent years due to their higher rate of acceleration, payload to weight ratio, stiffness and low effective inertia compared to those of serial manipulators. However they have some drawbacks, like a small workspace and special singular zones in it. Thus, in the presence of singular positions, the workspace of the parallel manipulators, which is less than that of serial manipulators, becomes even smaller and limits their functional performance. The studies of singularity have had different stages of development. The previous works on this problem is reported in a great number of publications and can be classified by different criteria. They can be arranged, by example, in three major groups, which are distinguished by historical evolution and are characterized by study of the singularity from kinematic, kinetostatic and dynamic points of view.

The physical interpretation of a singularity in kinematics refers to those configurations of parallel manipulators in which the number of degrees of freedom of the mechanical structure changes instantaneously, either the manipulator gains some additional, uncontrollable degrees of freedom or loses some degrees of freedom. In this case the singularity analysis can be carried out on the base of the properties of the Jacobian matrices of the mechanical structure (i.e. when the Jacobian matrices relating the input speeds and the output speeds become rank deficient (Gosselin and Angeles 1990; Ma and Angeles 1992; Ottaviano, Gosselin and Ceccarelli 2001; Wen and Oapos Brien 2003)), by using Grassmann geometry (Merlet 1989) or screw theory (Bonev, Zlatanov and Gosselin 2003; Hunt 1987). However, it was observed that close to a singular configuration, a parallel manipulator loses its stiffness and its quality of motion transmission, and as a result, its payload capability. For this purpose, a kinetostatic approach has been applied for the evaluation of the quality of motion transmission in the singular zones of
parallel manipulators. The quality of motion transmission of parallel manipulators was successfully studied in (Kim and Choi 2001; Kim and Ryu 2004; Lee et al. 2002; Weiwei and Shuang 2006). The quality of motion of manipulators with three degrees of freedom has been evaluated by means of a kinetostatic indicator, which is similar to the pressure angle (AlbaGomez, Wenger and Pamanes 2005). In (Arakelian, Briot and Glazunov 2007), the pressure angle was used as an indicator of the quality of motion transmission and the nature of the inaccessibility of singular zones by parallel manipulators are shown.

The further study of singularity in parallel manipulators has revealed an interesting problem that concerns the path planning of parallel manipulators under the presence of singular positions, i.e. the motion feasibility in the neighborhood of singularities. In this case the dynamic conditions can be considered in the design process. One of the most evident solutions for the stable motion generation in the neighborhood of singularities is to use redundant sensors and actuators (Alvan and Slousch 2003; Collins and Long 1975; Dasgupta and Mruthyunjaya 1998a; Glazunov et al. 2004). However, it is an expensive solution to the problem because of the additional actuators and the complicated control of the manipulator caused by actuation redundancy. Another approach concerns with motion planning to pass through singularity (Bhattacharya, Hatwal and Ghosh 1998; Dasgupta and Mruthyunjaya 1998b; Hesselbach 2004; Kemal Ider 2005; Kevin Jui and Sun 2005; Nenchev, Bhattacharya and Uchiyama 1997; Perng and Hsiao 1999), i.e. a parallel manipulator may track a path through singular poses if its velocity and acceleration are properly constrained. This is a promising way for the solution of this problem. However only a few research papers on this approach have addressed the path planning for obtaining a good tracking performance but they have not adequately addressed the physical interpretation of dynamic aspects.

In this paper the dynamic condition for passing through the singular positions is defined in general. It allows the stable motion generation inside in the presence of singularity by means of the optimum force control. The disclosed condition can be formulated as follows: «In the presence of a type 2 singularity, the platform of the parallel manipulator can pass through the singular positions without perturbation of motion if the wrench applied on the platform by the legs and external forces is orthogonal to the direction of uncontrollable motion». In other terms, the condition is that the work of applied forces and moments on the platform along the uncontrollable motion is equal to zero. This condition is obtained from the inverse dynamics and analytically demonstrated by the study of the Lagrangian of a general parallel manipulator. The obtained results are illustrated by numerical simulations and validated by experimental tests.

The paper is organized as follows. The next section presents theoretical aspects of the examined problem. Based on the Lagrangian formulation, the condition of force distribution is defined, that allows the passing of any parallel manipulator through the type 2 singular positions. In Section 3, two applications illustrate the obtained theoretical results. In section 4, the numerical simulations carried out using the software ADAMS are validated by experimental tests.

## 2. Optimal dynamic conditions for passing through type 2 singularity

Let us consider a parallel manipulator of $m$ links, $n$ degrees of freedom and driven by $n$ actuators.

The Lagrangian dynamic formulation for a parallel manipulator can be expressed as:

$$
\begin{equation*}
\tau=\frac{\mathrm{d}}{\mathrm{dt}}\left(\frac{\partial L}{\partial \dot{\mathbf{q}}}\right)-\frac{\partial L}{\partial \mathbf{q}}+\mathbf{B}^{T} \boldsymbol{\lambda} \tag{1}
\end{equation*}
$$

where,
$-\tau$ is the vector of the input efforts;

- $L$ is the Lagrangian of the manipulator;
- $\mathbf{q}=\left[q_{1}, q_{2}, \ldots, q_{n}\right]^{T}$ and $\dot{\mathbf{q}}=\left[\dot{q}_{1}, \dot{q}_{2}, \ldots, \dot{q}_{n}\right]^{T}$ represent the vector of active joints variables and the active joints velocities, respectively;
- $\mathbf{x}=[x, y, z, \phi, \psi, \theta]^{T}$ and $\mathbf{v}=[\dot{x}, \dot{y}, \dot{z}, \dot{\phi}, \dot{\psi}, \dot{\theta}]^{T}$ are trajectory parameters and their derivatives, respectively; $x, y, z$ represent the position of the controlled point and $\phi, \psi$ and $\theta$ the rotation of the platform about three axes $\mathbf{a}_{\phi,} \mathbf{a}_{\psi}$ and $\mathbf{a}_{\theta} ;$
$-\boldsymbol{\lambda}$ is the Lagrange multipliers vector, which is related to the wrench applied on the platform by:

$$
\begin{equation*}
\boldsymbol{\lambda}=\mathbf{A}^{-T} \mathbf{W}_{\mathbf{p}} \tag{2}
\end{equation*}
$$

where,

- $\mathbf{A}$ and $\mathbf{B}$ are two matrices relating the vectors $\mathbf{v}$ and $\dot{\mathbf{q}}$ according to $\mathbf{A v}=\mathbf{B} \dot{\mathbf{q}}$. They can be found by the closure equations with respect to time.
- $\mathbf{W}_{\mathbf{p}}$ is the wrench applied on the platform by the legs and external forces (Khalil and Guégan 2002), which is defined as:

$$
\mathbf{W}_{\mathbf{p}}=\left(\frac{\mathrm{d}}{\mathrm{dt}}\left(\frac{\partial L}{\partial \mathbf{v}}\right)-\frac{\partial L}{\partial \mathbf{x}}\right)=\left[\begin{array}{l}
\mathbf{f}_{\mathbf{p}}  \tag{3}\\
\mathbf{n}_{\mathbf{p}}
\end{array}\right] .
$$

where $\mathbf{f}_{\mathbf{p}}$ is the force along the directions of the global frame and $\mathbf{n}_{\mathbf{p}}$ is the torque about the axes $\mathbf{a}_{\phi,} \mathbf{a}_{\psi}$ and $\mathbf{a}_{\theta}$.

The term $\mathbf{W}_{\mathbf{p}}$ can be rewritten in the base frame using a transformation matrix $\mathbf{D}$ (Merlet 2006):

$$
\begin{equation*}
\mathbf{W}_{\mathbf{p}}=\mathbf{D}\left({ }^{\mathbf{R}_{0}} \mathbf{W}_{\mathbf{p}}\right) \tag{4}
\end{equation*}
$$

with ${ }^{\mathbf{R}_{0}} \mathbf{W}_{\mathbf{p}}$ is the expression of the wrench $\mathbf{W}_{\mathbf{p}}$ in the base frame, and

$$
\mathbf{D}=\left[\begin{array}{ll}
\mathbf{I}_{3 \times 3} & \mathbf{0}_{3 \times 3}  \tag{5}\\
\mathbf{0}_{3 \times 3} & \mathbf{R}_{3 \times 3}
\end{array}\right]
$$

where $\mathbf{I}_{3 \times 3}, \mathbf{0}_{3 \times 3}$ and $\mathbf{R}_{3 \times 3}$ are respectively the identity matrix, the zero matrix and the transformation matrix between axes $\mathbf{a}_{\phi,} \mathbf{a}_{\psi}$ and $\mathbf{a}_{\theta}$ and the base frame, whose dimensions are $3 \times 3$.

By substituting (5) into (1), one can obtain:

$$
\begin{equation*}
\tau=\mathbf{W}_{\mathbf{b}}+\mathbf{J}^{T \mathbf{R}_{0}} \mathbf{W}_{\mathbf{p}}, \mathbf{W}_{\mathbf{b}}=\frac{\mathrm{d}}{\mathrm{dt}}\left(\frac{\partial L}{\partial \dot{\mathbf{q}}}\right)-\frac{\partial L}{\partial \mathbf{q}} \tag{6}
\end{equation*}
$$

where $\mathbf{J}=\left({ }^{\mathbf{R}_{0}} \mathbf{A}\right)^{-1} \mathbf{B}$ is the Jacobian matrix between twist $\mathbf{t}$ of the platform (expressed in the base frame) and $\dot{\mathbf{q}},{ }^{\mathbf{R}_{0}} \mathbf{A}=\mathbf{A D}$ is the expression of matrix $\mathbf{A}$ in the base frame.

For any prescribed trajectory $\mathbf{x}(t)$, the values of vectors $\ddot{\mathbf{q}}, \dot{\mathbf{q}}$ and $\mathbf{q}$ can be found using the inverse kinematics. Thus, taking into account that the manipulator is not in a type 1 singularity (Gosselin and Angeles 1990), the terms $\mathbf{W}_{\mathbf{b}}$ and ${ }^{\mathbf{R}_{0}} \mathbf{W}_{\mathbf{p}}$ can be computed. However, for a trajectory passing through a type 2 singularity, the determinant of matrix $\mathbf{J}$ tends to infinity. Numerically, the values of the efforts applied by the actuators become infinite. In practice, the manipulator either is locked in such a position of the end-effector or it generates an uncontrolled motion. That is the end-effector of the manipulator could produce a motion, different to the prescribed trajectory.

It is known that a type 2 singularity appears when the determinant of matrix ${ }^{\mathbf{R}_{0}} \mathbf{A}$ vanishes, in other words, when at least two of its columns are linearly dependant (Merlet 2006).

Let us rewrite the matrix ${ }^{\mathbf{R}_{0}} \mathbf{A}$ as:

$$
\mathbf{R}_{0} \mathbf{A}=\left[\begin{array}{cccc}
a_{11} & a_{12} & \ldots & a_{16}  \tag{7}\\
a_{21} & a_{22} & \ldots & a_{26} \\
\vdots & \vdots & \ddots & \vdots \\
a_{61} & a_{62} & \ldots & a_{66}
\end{array}\right]
$$

In the presence of type 2 singularity the columns of matrix ${ }^{\mathbf{R}_{0}} \mathbf{A}$ are linearly dependant, i.e.

$$
\begin{equation*}
\sum_{j=1}^{6} \alpha_{j} a_{i j}=0, i=1, \ldots, 6 \tag{8}
\end{equation*}
$$

where $\alpha_{j}$ are the coefficients, which in general can be functions of $q_{p}(p=1, \ldots, n)$. It should be noted that the vector $\mathbf{t}_{\mathrm{s}}=\left[\alpha_{1}, \alpha_{2}, \ldots, \alpha_{6}\right]^{T}$ represents the direction of the uncontrollable motion of the platform in a type 2 singularity.

Rewriting (8) in a vector form, we obtain:

$$
\begin{equation*}
\sum_{j=1}^{6} \alpha_{j} \mathbf{N}_{j}=\mathbf{0}, \mathbf{N}_{j}=\left[a_{1 j}, a_{2 j}, \ldots, a_{6 j}\right]^{T}, j=1, \ldots, 6 \tag{9}
\end{equation*}
$$

where $\mathbf{N}_{j}$ represents the $j$-th column of matrix ${ }^{\mathbf{R}_{0}} \mathbf{A}$.
By substituting (9) into (2), we obtain

$$
\begin{equation*}
\mathbf{N}_{j}^{T} \lambda=W_{j}, j=1, \ldots, 6 \tag{10}
\end{equation*}
$$

where $W_{j}$ is the $j$-th row of vector ${ }^{\mathbf{R}_{0}} \mathbf{W}_{\mathbf{p}}$.
Then, from (9) and (10) the following conditions are derived:

$$
\begin{equation*}
\sum_{j=1}^{6}\left(\alpha_{j} \mathbf{N}_{j}^{T} \lambda\right)=\sum_{j=1}^{6}\left(\alpha_{j} W_{j}\right)=0 \tag{11}
\end{equation*}
$$

The right term of eq.(11) corresponds to the scalar product of vectors $\mathbf{t}_{\mathrm{s}}$ and ${ }^{\mathbf{R}_{0}} \mathbf{W}_{\mathbf{p}}$.
Thus, in the presence of a type 2 singularity, it is possible to satisfy conditions (11) if the wrench applied on the platform by the legs and external efforts ${ }^{\mathbf{R}_{0}} \mathbf{W}_{\mathrm{p}}$ are orthogonal to the direction of the uncontrollable motion $\mathbf{t}_{\mathbf{s}}$. Otherwise, the dynamic model is not consistent. Obviously, in the presence of a type 2 singularity, the displacement of the end-effector of the manipulator has to be planned to satisfy (11).

Let us illustrate the considered problem by examples.

## 3. Illustrative examples

In this section, two examples are chosen to illustrate the obtained theoretical results discussed above. The first example presents a planar 5R parallel manipulator, which allows obtaining relatively simple mathematical models for demonstrating the expected results by numerical simulations. The second example presents a parallel manipulator which was developed in the I.N.S.A. of Rennes. This example was chosen for validation of numerical simulations carried out by the software ADAMS on the built prototype.

### 3.1. Example 1: Planar 5R parallel manipulator

In the planar 5R parallel manipulator, as shown in Fig. 1, the output point is connected to the base by two legs, each of which consists of three revolute joints and two links. In each of the two legs, the revolute joint connected to the base is actuated. Thus, such a manipulator is able to position its output point in a plane.


Fig. 1. Kinematic chain of the planar 5R parallel manipulator.


Fig. 2. Second kind of singularities of the planar 5R parallel manipulator.

As shown in Fig. 1, the actuated joints are denoted as $A$ and $E$ with input parameters $q_{1}$ and $q_{2}$. The common joint of the two legs is denoted as $C$, which is also the output point with controlled parameters $x$ and $y$. A fixed global reference system $\boldsymbol{x} \boldsymbol{O} \boldsymbol{y}$ is located at the center of $A E$ with the $\boldsymbol{y}$ axis normal to $\mathbf{A E}$ and the $\boldsymbol{x}$-axis directed along $\mathbf{A E}$. The lengths of the links $A B, B C, C D, D E$ are respectively denoted as $L_{1}, L_{2}, L_{3}$ and $L_{4}$. The positions of the centers of masses $S_{i}$ of links from joint centers $A, B, D$ and $E$ are respectively denoted by dimensionless lengths $r_{1}, r_{2}, r_{3}$ and $r_{4}$, i.e. $A S_{1}=r_{1} L_{1}, B S_{2}=r_{2} L_{2}, D S_{3}=r_{3} L_{3}$ and $E S_{4}=r_{4} L_{4}$.

The singularity analysis of this manipulator (Liu, Wang and Pritschow 2006) shows that the Type 2 singularities appear when legs 2 and 3 are parallel (Fig. 2).

In both cases, the gained degree of freedom is an infinitesimal translation perpendicular to the legs 2 and 3. However, if $L_{2}=L_{3}$, the gained degree of freedom in case (b) becomes a finite rotary motion about point $B$.

In order to simplify the analytic expressions, we consider that the gravity effects are along the $z$-axis and consequently the input torques are only due to inertia effects. To simplify the computation, it is also preferable to replace the masses of moving links by concentrated masses
(Seyferth 1974; Wu and Gosselin 2007). For a link $i$ with mass $m_{i}$ and its axial moment of inertia $I_{i}$, we have:

$$
\left[\begin{array}{ccc}
1 & 1 & 1  \tag{12}\\
r_{i} & 0 & 1-r_{i} \\
r_{i}^{2} L_{i}^{2} & 0 & \left(1-r_{i}\right)^{2} L_{i}^{2}
\end{array}\right]\left[\begin{array}{l}
m_{i 1} \\
m_{i 2} \\
m_{i 3}
\end{array}\right]=\left[\begin{array}{c}
m_{i} \\
0 \\
I_{i}
\end{array}\right],(i=1,2,3,4)
$$

where $m_{i j}(j=1,2,3)$ are the values of the three point masses placed at the centers of the revolute joints and at the center of masses of the link $i$.

In this case, the kinetic energy $T$ can be written as:

$$
\begin{equation*}
T=\frac{1}{2}\left(m_{S 1} \mathbf{V}_{S 1}^{2}+m_{S 2} \mathbf{V}_{S 2}^{2}+m_{S 3} \mathbf{V}_{S 3}^{2}+m_{S 4} \mathbf{V}_{S 4}^{2}+m_{B} \mathbf{V}_{B}^{2}+m_{C} \mathbf{V}_{C}^{2}+m_{D} \mathbf{V}_{D}^{2}\right) \tag{13}
\end{equation*}
$$

where, $\quad m_{S 1}=m_{12}, \quad m_{S 2}=m_{22}, \quad m_{S 3}=m_{32}, \quad m_{S 4}=m_{42}, \quad m_{B}=m_{13}+m_{21}, \quad m_{C}=m_{23}+m_{21}$, $m_{D}=m_{33}+m_{41}$. The terms $m_{i j}(i=1,2,3,4)$ are deduced from the relation (12), $\mathbf{V}_{S i}$ is the vector of the linear velocities of the center of masses $S_{i} ; \mathbf{V}_{B}, \mathbf{V}_{C}$ and $\mathbf{V}_{D}$ are the vectors of the linear velocities of the corresponding axes.

The input torques can be obtained from (6):

$$
\begin{equation*}
\boldsymbol{\tau}=\mathbf{W}_{\mathbf{b}}+\mathbf{J}_{5 R}^{T} \mathbf{W}_{\mathbf{p}} \tag{14}
\end{equation*}
$$

taking into account that for examined manipulator:

$$
\begin{equation*}
\mathbf{W}_{\mathbf{b}}=\mathbf{J}_{\mathbf{B}}^{T} \mathbf{F}_{\mathbf{B}}+\mathbf{J}_{\mathbf{D}}^{T} \mathbf{F}_{\mathbf{D}}, \tag{15}
\end{equation*}
$$

where,

$$
\begin{gather*}
\mathbf{J}_{\mathbf{B}}=\left[\begin{array}{cc}
-L_{1} \sin q_{1} & 0 \\
L_{1} \cos q_{1} & 0
\end{array}\right], \mathbf{J}_{\mathbf{D}}=\left[\begin{array}{cc}
0 & -L_{4} \sin q_{2} \\
0 & L_{4} \cos q_{2}
\end{array}\right],  \tag{16}\\
\mathbf{F}_{\mathbf{B}}=m_{B 1} \Gamma_{\mathbf{B}}+m_{C 1} \Gamma_{\mathbf{C}}, \mathbf{F}_{\mathbf{D}}=m_{D 2} \Gamma_{\mathbf{D}}+m_{C 3} \Gamma_{\mathbf{C}},  \tag{17}\\
\Gamma_{\mathbf{B}}=L_{1}\left(\ddot{q}_{1}\left[\begin{array}{c}
-\sin q_{1} \\
\cos q_{1}
\end{array}\right]-\dot{q}_{1}^{2}\left[\begin{array}{c}
\cos q_{1} \\
\sin q_{1}
\end{array}\right]\right), \Gamma_{\mathbf{D}}=L_{4}\left(\ddot{q}_{2}\left[\begin{array}{c}
-\sin q_{2} \\
\cos q_{2}
\end{array}\right]-\dot{q}_{2}^{2}\left[\begin{array}{c}
\cos q_{2} \\
\sin q_{2}
\end{array}\right]\right), \Gamma_{\mathbf{C}}=\left[\begin{array}{c}
\ddot{x} \\
\ddot{y}
\end{array}\right] \tag{18}
\end{gather*}
$$

$$
\begin{align*}
& m_{B 1}=m_{S 1} r_{1}^{2}+m_{B}+m_{S 2}\left(1-r_{2}\right)^{2}, m_{C 1}=m_{S 2} r_{2}\left(1-r_{2}\right),  \tag{19}\\
& m_{C 3}=m_{S 3} r_{3}\left(1-r_{3}\right), m_{D 2}=m_{S 4} r_{4}^{2}+m_{D}+m_{S 3}\left(1-r_{3}\right)^{2} . \tag{20}
\end{align*}
$$

The term $\mathbf{W}_{\mathbf{p}}$ is given by:

$$
\begin{gather*}
\mathbf{W}_{\mathbf{p}}=m_{C 1} \boldsymbol{\Gamma}_{\mathbf{B}}+m_{C 2} \boldsymbol{\Gamma}_{\mathbf{C}}+m_{C 3} \boldsymbol{\Gamma}_{\mathbf{D}}  \tag{21}\\
m_{C 2}=m_{s 2} r_{2}^{2}+m_{C}+m_{S 3} r_{3}^{2} \tag{22}
\end{gather*}
$$

and the Jacobian matrix $\mathbf{J}_{5 R}$ by:

$$
\begin{equation*}
\mathbf{J}_{5 R}=\mathbf{A}_{5 R}^{-1} \mathbf{B}_{5 R} \tag{23}
\end{equation*}
$$

where,

$$
\begin{gather*}
\mathbf{A}_{5 R}=\left[\begin{array}{ll}
a_{11} & a_{12} \\
a_{21} & a_{22}
\end{array}\right]=2\left[\begin{array}{cc}
x-L_{1} \cos q_{1}+a & y-L_{1} \sin q_{1} \\
x-L_{4} \cos q_{2}-a & y-L_{4} \sin q_{2}
\end{array}\right],  \tag{24}\\
\mathbf{B}_{5 R}=-\left[\begin{array}{cc}
L_{1}\left(a_{11} \sin q_{1}-a_{12} \cos q_{1}\right) & 0 \\
0 & L_{4}\left(a_{21} \sin q_{2}-a_{22} \cos q_{2}\right)
\end{array}\right] . \tag{25}
\end{gather*}
$$

and we determine $\mathbf{t}_{s}$ in according with (8):

$$
\begin{equation*}
\mathbf{t}_{\mathbf{s}}=\left[-\sin \psi_{1}, \cos \psi_{1}\right]^{T} \tag{26}
\end{equation*}
$$

Thus, the examined manipulator can pass through the given singular positions if the wrench $\mathbf{W}_{\mathbf{p}}$ determined by (21) is orthogonal to the direction of the uncontrollable motion $\mathbf{t}_{\mathbf{s}}$ described by (26).

Let us now consider the motion planning, which makes it possible to satisfy this condition. For this purpose the following parameters of manipulator's links are specified: $L_{1}=L_{2}=L_{3}=L_{4}$ $=0.25 \mathrm{~m} ; r_{1}=r_{2}=r_{3}=r_{4}=0.5 ; a=0.2 \mathrm{~m} ; m_{1}=m_{4}=2.81 \mathrm{~kg} ; I_{1}=I_{4}=0.02 \mathrm{~kg} / \mathrm{m}^{2} ; m_{2}=m_{3}=$ $1.41 \mathrm{~kg} ; I_{2}=I_{3}=0.01 \mathrm{~kg} / \mathrm{m}^{2}$.


Fig. 3. Initial, singular and final positions of the planar 5R parallel manipulator.

With regard to the prescribed trajectory generation, the point $C$ should reproduce a motion along a straight line between the initial position $C_{0}\left(x_{0}, y_{0}\right)=C_{0}(0.1,0.345)$ and the final point $C_{f}$ $\left(x_{f}, y_{f}\right)=C_{f}(-0.1,0.145)$ in $t_{f}=2 \mathrm{~s}$.

Thus, the given trajectory can be expressed as follows:

$$
\mathbf{x}=\left[\begin{array}{l}
x(t)  \tag{27}\\
y(t)
\end{array}\right]=\left[\begin{array}{l}
x_{0}+s(t)\left(x_{f}-x_{0}\right) \\
y_{0}+s(t)\left(y_{f}-y_{0}\right)
\end{array}\right]
$$

However, the manipulator will pass by a type 2 singular position at point $C_{s}\left(x_{s}, y_{s}\right)=C_{s}(0$, 0.245 (Fig. 3).

Developing the condition for passing through the singular position (11) for the planar 5R parallel manipulator at point $C_{s}$, we obtain:

$$
\begin{equation*}
m_{C 1} L_{1}\left(248 \dot{x}^{2}-48 \dot{y}^{2}\right)-3 \sqrt{6} m_{C 2} \ddot{y}=0 \tag{28}
\end{equation*}
$$

Then, taking into account that the velocity and the acceleration of the end-effector in initial and final positions are equal to zero, the following nine boundary conditions are found:

$$
\begin{gather*}
s\left(t_{0}\right)=0,  \tag{29}\\
s\left(t_{f}\right)=1,  \tag{30}\\
s\left(t_{s}=1 \mathrm{~s}\right)=0.5, \tag{31}
\end{gather*}
$$

$$
\begin{equation*}
\dot{s}\left(t_{0}\right)=0, \tag{32}
\end{equation*}
$$

$$
\begin{equation*}
\dot{s}\left(t_{f}\right)=0, \tag{33}
\end{equation*}
$$

$$
\begin{equation*}
\dot{s}\left(t_{s}\right)=\dot{y}_{s} /\left(y_{f}-y_{0}\right)=\dot{x}_{s} /\left(x_{f}-x_{0}\right)=1, \tag{34}
\end{equation*}
$$

$$
\begin{align*}
& \ddot{s}\left(t_{0}\right)=\ddot{s}_{0}=0 .  \tag{35}\\
& \ddot{s}\left(t_{f}\right)=\ddot{s}_{f}=0 . \tag{36}
\end{align*}
$$

$$
\begin{equation*}
\ddot{s}\left(t_{s}\right)=\ddot{s}_{s}=m_{C 1} L_{1}\left(248 \ddot{x}_{s}^{2}-48 \dot{y}_{s}^{2}\right) /\left(3\left(x_{f}-x_{0}\right) \sqrt{6} m_{C 2}\right) . \tag{37}
\end{equation*}
$$



Fig. 4. Input torques of the planar 5R parallel manipulator in the case of the eight order polynomial trajectory planning, obtained by the ADAMS software.


Fig. 5. Input torques of the planar 5R parallel manipulator in the case of the fifth order polynomial trajectory planning, obtained by the ADAMS software.

From (28) - (37), the following eighth order polynomial trajectory planning is found:

$$
\begin{equation*}
s(t)=-0.25851 t^{3}+3.84228 t^{4}-5.72792 t^{5}+3.58909 t^{6}-1.07101 t^{7}+0.12606 t^{8} \tag{38}
\end{equation*}
$$

Thus the generation of the motion by the obtained eighth order polynomial makes it possible to pass through the singularity without perturbation and the input torques remain in the limits of finite values, which are validated by numerical simulations carried out by the ADAMS software (Fig. 4).

Thus, we can assert that the obtained optimal dynamic conditions assume the passing of the manipulator's end-effector through the singular position.

Now, we would like to show that, in the case of the generation of the motion by any trajectory planning without meeting the adopted boundary conditions, the end-effector is not able to pass through the singular position. For this purpose the generation of motion between initial and final positions, let us generate by a fifth order polynomial trajectory planning:

$$
\begin{equation*}
s(t)=1.25 t^{3}-0.9375 t^{4}+0.1875 t^{5} \tag{39}
\end{equation*}
$$

The obtained numerical simulations carried out by the software ADAMS are given in Fig. 5. We can see that, when the manipulator is close to the singular configuration (for $t_{s}=1 \mathrm{~s}$ ), the values of the input torques tend to infinity.

### 3.2. Example 2: PAMINSA (Parallel Manipulator of the I.N.S.A.)

The second example presents a parallel manipulator, which was invented and developed at the I.N.S.A of Rennes (Arakelian et al. 2006). The particularity of this architecture is in decoupling of the displacements of the platform in the horizontal plane from the translations along the vertical axis. The advantages of such an approach was disclosed in (Arakelian et al. 2005; Briot et al. 2007b) and the singularity analysis is discussed in (Briot et al. 2007a; Briot and Arakelian 2007; Arakelian, Briot and Glazunov 2006).

The previous studies have revealed that there are type 2 singularities in the workspace of the symmetrical architecture of PAMINSA. Let us illustrate the proposed approach for the PAMINSA with 4 degrees of freedom (Fig. 6).

Each leg of this manipulator is realized by a pantograph mechanism (Fig. 7) with two input points $3_{i}$ and $8_{i}$, and an output point $5_{i}(i=1,2,3)$. Each input point $8_{i}$ is connected to the rotating drive $M_{i}$ by means of a prismatic guide mounted on a rotating link. This kind of architecture allows for generation of motion in the horizontal plane by the use of rotating actuators $M_{1}, M_{2}$, $M_{3}$, and the vertical translations by means of the linear actuator $M_{v}$. Thus, the displacements ( $x, y$, $\phi)$ of the platform in the horizontal plane $\boldsymbol{x} \boldsymbol{O} \boldsymbol{y}$, that are translations along the $\boldsymbol{x}$ and $\boldsymbol{y}$-axes and rotations about the $z$-axis, are independent of vertical translations $z$.


Fig. 6. PAMINSA with 4 DOF.


Fig. 7. Kinematic chain of each leg.

This implies that the kinematic models controlling the displacement of the manipulator can be divided into two parts:

- a model for the displacements in the horizontal plane, which is equivalent to a 3-RPR manipulator;
- a model for the translations along the vertical axis equivalent to the model for the vertical translations of a pantograph linkage.

The type 2 singularities of such a manipulator appear when (Arakelian, Briot and Glazunov 2006; Briot et al. 2007a):
a) the three legs of the manipulators are parallel, which is impossible for the developed PAMINSA manipulator.
b) the orientation of the platform is equal to $\cos ^{-1}\left(R_{p l} / R_{b}\right)$, where $R_{p l}$ and $R_{b}$ correspond respectively to the lengths $P C_{i}$ and $O M^{\prime}$ (Fig. 8). In this case, the manipulator gains one infinitesimal rotation around one vertical axis.
c) the platform is located in a circle defined by

$$
\begin{equation*}
x^{2}+y^{2}=R_{p l}^{2}+R_{b}^{2}-2 R_{p l} R_{b} \cos \phi \tag{40}
\end{equation*}
$$

In this case, the manipulator gains one finite rotation about one vertical axis (Cardanic self motion) (Briot et al. 2007a).


Fig. 8. Example of Type 2 singular configuration
(horizontal projection of the examined structure).

For both cases (b) and (c), the direction of the unconstrained motion can be represented by the twist $\mathbf{t}_{\mathbf{s}}=\left[0,0,1, x_{W}, y_{W}, 0\right]^{T}$, where $x_{W}$ and $y_{W}$ corresponds to the planar coordinates of the intersection point of the wrenches $\mathbf{R}_{i}$ applied on the platform by the three legs of the manipulator (Fig. 8).

Let us now study the inverse dynamics of the PAMINSA. The potential energy $V$ can be written as:

$$
\begin{equation*}
V=V_{p l}+\sum_{i=1}^{3} V_{l e g_{i}} \tag{41}
\end{equation*}
$$

where $V_{p l}$ is the potential energy of the platform and $V_{\text {leg }_{i}}$ is the potential energy of the leg $i(i=1$, $2,3)$.

By further considering that the coordinates of the all points of the pantograph linkages can be found as a linear combination of the coordinates of points $3_{i}, 5_{i}$ and $9_{i}$, one can express the terms $V_{p l}$ and $V_{\text {leg }_{i}}$ as follows:

$$
\begin{gather*}
V_{p l}=m_{p l} g z  \tag{42}\\
V_{l \operatorname{leg}_{i}}=C_{v 1} z_{5 i}+C_{v 2} z_{9 i}+C_{v 3} q_{v}+C_{v 4} \tag{43}
\end{gather*}
$$

Here, $C_{v j}(j=1,2,3)$ are constant terms whose dimension is equivalent to a mass multiplied by the gravitational acceleration $g, m_{p l}$ is the mass of the platform with a payload, and $z_{5 i}$ and $z_{9 i}$ are the altitude of joints $5_{i}$ and $9_{i}$. The expressions of the coordinates of joints $5_{i}$ and $9_{i}$ are given in appendix A. The expressions for $C_{v j}(j=1, \ldots, 4)$ are given in appendix B.

We consider that the links are perfect tubes. Therefore the tensor of inertia $\mathbf{I}_{j}$ of the link $B_{j i}$ at the center of masses will be written as:

$$
\mathbf{I}_{j}=\left[\begin{array}{ccc}
I_{X X}^{(j)} & 0 & 0  \tag{44}\\
0 & I_{Y Y}^{(j)} & 0 \\
0 & 0 & I_{Z Z}^{(j)}
\end{array}\right] \text {, with } I_{Y Y}^{(j)}=I_{Z Z}^{(j)}
$$

Thus, the kinetic energy $T$ of the manipulator can be represented as:

$$
\begin{equation*}
T=T_{p l}+\sum_{i=1}^{3} T_{l g_{i}}, \tag{45}
\end{equation*}
$$

where $T_{p l}$ is the kinetic energy of the platform, $T_{\text {legi }}$ is the kinetic energy of the leg $i$, as:

$$
\begin{equation*}
T_{p l}=\frac{1}{2}\left(m_{p l}\left(\dot{x}^{2}+\dot{y}^{2}+\dot{z}^{2}\right)+I_{p l} \dot{\phi}^{2}\right) \tag{46}
\end{equation*}
$$

where $I_{p l}$ is the axial moment of inertia of the platform about the vertical axis.

$$
\begin{gather*}
T_{\text {leg }_{i}}=T_{\text {transi }_{i}}+T_{\text {rot }_{i}}  \tag{47}\\
T_{\text {trans }}=C_{c 1}\left(\dot{x}_{5 i}^{2}+\dot{y}_{5 i}^{2}\right)+C_{c 2} \dot{z}_{5 i}^{2}+C_{c 3}\left(\dot{x}_{9 i}^{2}+\dot{y}_{9 i}^{2}+\dot{z}_{9 i}^{2}\right)+C_{c 4}\left(\dot{x}_{5 i} \dot{x}_{9 i}+\dot{y}_{5 i} \dot{y}_{9 i}\right)+C_{c 5} \dot{z}_{5 i} \dot{z}_{9 i}  \tag{48}\\
+C_{c 6} \dot{q}_{v}+C_{c 7} \dot{z}_{5 i} \dot{q}_{v}+C_{c 8} \dot{q}_{i}^{2}
\end{gather*}
$$

$T_{\text {rot }_{i}}$ is the kinetic energy of the rotating links.
Note that there are two types of rotations (see, Fig. 7):

- rotation due to the actuators $M_{i}(i=1,2,3)$ (angle $\left.q_{i}\right)$, which is about the vertical axis,
- rotations due to the displacement of the pantograph in the linkage plane (angles $\zeta_{i}$ and $\varepsilon_{i}$ denoted as the angles between the direction of the passive slider and the links $B_{4 i}$ and $B_{3 i}$ respectively).

Thus, the kinetic energy of the rotating links can be written as:

$$
\begin{equation*}
T_{\text {rot }}^{i} \text { }=C_{c 9} \dot{\zeta}_{i}^{2}+C_{c 10} \dot{\varepsilon}_{i}^{2}+\dot{q}_{i}^{2}\left(C_{c 13}+C_{c 10} \sin ^{2} \varsigma_{i}+C_{c 9} \cos ^{2} \varsigma_{i}+C_{c 12} \sin ^{2} \varepsilon_{i}+C_{c 11} \cos ^{2} \varepsilon_{i}\right) \tag{49}
\end{equation*}
$$

The expressions for $C_{c j}(j=1, \ldots, 13)$ are given in appendix C .
The input torques can be obtained from (6):

$$
\begin{equation*}
\boldsymbol{\tau}=\mathbf{W}_{\mathbf{b}}+\mathbf{J}_{P A M}^{T} \mathbf{W}_{\mathbf{p}} \tag{50}
\end{equation*}
$$

where the terms $\mathbf{J}_{\text {PAM }}, \mathbf{W}_{\mathbf{b}}$ and $\mathbf{W}_{\mathbf{p}}$ are presented in appendix D.
The following parameters of manipulator's links are specified for the trajectory generation:

- the radii of the circles circumscribed to the base and platform triangles are respectively equal to $R_{b}=0.35 \mathrm{~m}$ and $R_{p l}=0.1 \mathrm{~m}$;
- magnification factor of the pantograph: $k=3$;
- gravitational acceleration $g$ is equal to $9.81 \mathrm{~m} / \mathrm{s}^{2}$.
- lengths of the links of the pantograph linkages: $L_{B 1}=0.308 \mathrm{~m}, L_{B 2}=0.442 \mathrm{~m}, L_{B 3}=L_{B 8}=0.42$ $\mathrm{m}, L_{B 4}=k L_{B 7}=0.63 \mathrm{~m}, L_{B 5}=0.0275 \mathrm{~m}, L_{B 10}=0.3635 \mathrm{~m} ;$
- masses of the joints of the pantograph linkages: $m_{2}=0.214 \mathrm{~kg}, m_{3}=0.338 \mathrm{~kg}, m_{4}=0.262 \mathrm{~kg}$, $m_{5}=0.233 \mathrm{~kg}, m_{7}=3.08 \mathrm{~kg}, m_{8}=0.305 \mathrm{~kg}, m_{9}=0.259 \mathrm{~kg} ;$
- mass of the platform: $m_{p l}=2.301 \mathrm{~kg}$;
- masses of the links of the pantograph linkages: $m_{B 1}=1.221 \mathrm{~kg}, m_{B 2}=0.921 \mathrm{~kg}, m_{B 3}=0.406 \mathrm{~kg}$, $m_{B 4}=0.672 \mathrm{~kg}, m_{B 7}=0.107 \mathrm{~kg}, m_{B 8}=0.403 \mathrm{~kg}, m_{B 10}=0.436 \mathrm{~kg} ;$
- term of the inertia matrix of the platform: $I_{p l}=0.015 \mathrm{~kg} / \mathrm{m}^{2}$.
- terms of the inertia matrices of the links of the pantograph linkages:

$$
\begin{aligned}
& I_{X X}^{(B 3)}=0.0038 \mathrm{~kg} / \mathrm{m}^{2}, I_{Y Y}^{(B 3)}=0.02 \mathrm{~kg} / \mathrm{m}^{2}, I_{X X}^{(B 4)}=0.0012 \mathrm{~kg} / \mathrm{m}^{2}, I_{Y Y}^{(B 4)}=0.048 \mathrm{~kg} / \mathrm{m}^{2}, \\
& I_{X X}^{(B 7)}=8 \cdot 10^{-4} \mathrm{~kg} / \mathrm{m}^{2}, I_{Y Y}^{(B 7)}=0.003 \mathrm{~kg} / \mathrm{m}^{2}, I_{X X}^{(B 8)}=0.0024 \mathrm{~kg} / \mathrm{m}^{2}, I_{Y Y}^{(B 8)}=0.02 \mathrm{~kg} / \mathrm{m}^{2}, \\
& I_{B 2}=0.003 \mathrm{~kg} / \mathrm{m}^{2}, I_{B 10}=0.02 \mathrm{~kg} / \mathrm{m}^{2} .
\end{aligned}
$$

The point $P$ is desired to make a motion $x(t)$ along a straight line between points $P_{0}\left(x_{0}, y_{0}\right)=P_{0}$ $(0,0)$ and point $P_{f}\left(x_{f}, y_{f}\right)=P_{f}(0.3,0)$ in $t_{f}=2.4 \mathrm{~s}$. However, the manipulator will pass through a type 2 singular position at point $P_{s}\left(x_{s}, y_{s}\right)=(0.25,0)$ (Fig. 9).


Fig. 9. Displacement of the PAMINSA along the prescribed straight line (planar projection).

In order to carry out a comparative analysis for the optimized and not optimized dynamic conditions for passing through type 2 singularity, it has been considered two cases. The first is such a movement on the given trajectory, which is calculated from condition (11), and the second is an arbitrary motion.

At first let us consider an optimized trajectory which allows satisfying the condition (11), i.e. the force $\mathbf{W}_{\mathbf{p}}$ should be perpendicular to the to the twist $\mathbf{t}_{s}[0,0,1,0,0.1,0]^{T}$ defining the direction of the unconstrained motion. Developing the expression (11) for the PAMINSA at point $P_{s}$, we obtain:

$$
\begin{align*}
0 & =0.06441 \ddot{x}+1.2115 \ddot{y}-0.14649 \ddot{z}+0.04425 \ddot{\phi}+0.06827+6.85084 \dot{x}^{2}+0.11720 \dot{y}^{2} \\
& -0.18482 \dot{z}^{2}+0.02947 \dot{\phi}^{2}-0.85175 \dot{\phi} \dot{x}+0.05643 \dot{\phi} \dot{y}+0.19423 \dot{\phi} \dot{z}-5.17625 \dot{x} \dot{y}  \tag{51}\\
& +0.46477 \dot{x} \dot{z}+2.94694 \dot{y} \dot{z}
\end{align*}
$$

Now considering that the end-effector of the manipulator moves along a straight line directed along the $\boldsymbol{x}$-axis, we can note that $\dot{y}\left(t_{s}\right)=\dot{z}\left(t_{s}\right)=\ddot{y}\left(t_{s}\right)=\ddot{z}\left(t_{s}\right)=\dot{\phi}\left(t_{s}\right)=\ddot{\phi}\left(t_{s}\right)=0$. Thus, the relationships, which satisfy the passing through of the singular positions, taking into account that
the velocity and the acceleration of the platform in the initial and final positions are equal to zero, can be expressed by the following boundary conditions:

$$
\begin{gather*}
x\left(t_{0}\right)=x_{0},  \tag{52}\\
x\left(t_{f}\right)=x_{f},  \tag{53}\\
x\left(t_{s}=2 \mathrm{~s}\right)=x_{s},  \tag{54}\\
\dot{x}\left(t_{0}\right)=0,  \tag{55}\\
\dot{x}\left(t_{f}\right)=0,  \tag{56}\\
\ddot{x}\left(t_{0}\right)=0,  \tag{57}\\
\ddot{x}\left(t_{f}\right)=0,  \tag{58}\\
\dot{x}\left(t_{s}\right)=\dot{x}_{s}=0.05 \mathrm{~m} / \mathrm{s},  \tag{59}\\
\ddot{x}\left(t_{s}\right)=\ddot{x}  \tag{60}\\
s
\end{gather*}=-1.32583 \mathrm{~m} / \mathrm{s}^{2} . ~ \$
$$

In this case, a motion for passing of the platform through the singular position can be found from the following eighth order polynomial form:

$$
\begin{equation*}
x(t)=3.41 t^{8}-37.65 t^{7}+166.05 t^{6}-365.23 t^{5}+400.63 t^{4}-175.27 t^{3} \tag{61}
\end{equation*}
$$

However, a trajectory obtained by (61) cannot be reproduced by the prototype because of the limited capability of drivers' deceleration. Therefore, the trajectory was divided into two parts, i.e., the first sixth order polynomial trajectory assumes the motion from an initial to the singular position $\left(P_{0} P_{s}\right)$ and the second sixth order polynomial trajectory from singular to the final position $\left(P_{s} P_{f}\right)$. The core of the problem is same but it allows for generating motions for the prototype.

Thus, the trajectory planning equations can be written as:

$$
\begin{equation*}
x(t)=x_{0}+\left(x_{s}-x_{0}\right)\left(b_{3} t^{3}+b_{4} t^{4}+b_{5} t^{5}+b_{6} t^{6}\right) \text { for } t \leq t_{s} ; \tag{62}
\end{equation*}
$$

$x(t)=x_{s}+\left(x_{f}-x_{s}\right)\left(c_{1}\left(t-t_{s}\right)+c_{2}\left(t-t_{s}\right)^{2}+c_{4}\left(t-t_{s}\right)^{4}+c_{5}\left(t-t_{s}\right)^{5}+c_{6}\left(t-t_{s}\right)^{6}\right)$ for $t>t_{s}$.
with $b_{3}=-3.3033, b_{4}=5.10456, b_{5}=-2.45207, b_{6}=0.37844, c_{1}=1, c_{2}=-13.25829, c_{4}=$ $2365.3672, c_{5}=-11953.07236$ and $c_{6}=16158.76157$.

Thus, the motion obtained from the following sixth order polynomial equations
$x(t)=-0.826 t^{3}+1.276 t^{4}-0.613 t^{5}+0.095 t^{6}$ for $t \leq 2 \mathrm{~s} ;$
$x(t)=727227-2067183 t+244555.2 t^{2}-1541224 t^{3}+54571.1 t^{4}-102929 t^{5}+807.9 t^{6}$
for $t>2 \mathrm{~s}$;
allows for passing through the singularity without perturbation, and the input efforts take on finite values (Fig. 10).

It can be seen that the input torques remain in the limits of finite values, but, by the end of the motion there is an increase of the input efforts, caused by a quick deceleration to stop the manipulator before it reaches the workspace boundary. It will be shown further that in the case of the motion generated by any trajectory planning without meeting the adopted boundary conditions $(52)-(60)$, the manipulator platform is not able to pass through the singular position. For this purpose, the generation of motion between initial and final positions is carried out by a fifth order polynomial trajectory planning.

In this case, for $y(t)=0 \mathrm{~m}, z(t)=-0.45 \mathrm{~m}$ and $\phi(t)=0$, the fifth order polynomial trajectory planning is the following:

$$
\begin{equation*}
x(t)=0.217 t^{3}-0.137 t^{4}+0.023 t^{5} \tag{66}
\end{equation*}
$$

The obtained input efforts computed by the software ADAMS are represented in Fig. 11.
It can be noted that, while the manipulator passes through the singular configuration (for $t_{s} \approx$ 1.8 s ), the value of the input torques tend to infinity.

Let us now validate the obtained results by experimental tests.


Fig. 10. Input efforts of the PAMINSA in the case of the sixth order polynomial trajectory planning, computed with ADAMS software.


Fig. 11. Input efforts of the PAMINSA in the case of the fifth order polynomial trajectory planning, computed with ADAMS software.

## 4. Experimental validation of obtained results

For validating the results of the previous section, we have carried out experimental tests on the prototype of the PAMINSA developed in the I.N.S.A. of Rennes (Fig. 12).


Fig. 12. The prototype of PAMINSA developed in the I.N.S.A. of Rennes.


Fig. 13. Trajectory reproduction on the PAMINSA during the displacement of the platform with the fifth order polynomial law (view from below).


Fig. 14. Trajectory reproduction on the PAMINSA during the displacement of the platform with the sixth order polynomial law (view from below).

At first, we have applied an arbitrary fifth order control law and observed the reproduction of motion during the displacement of the platform. The obtained trajectory is shown in Fig. 13 (dotted line).

The different positions are classified by time. For positions from (a) to (d), the platform moves towards the singular zone but yet it is outside of it. In this case, the reproduction of the real trajectory is similar to the desirable. At position (e), the manipulator enters the singular zone, which is close to the circle of the theoretical singular loci, and starts an uncontrollable motion. Thus, since the motion generation is carried out by non optimized dynamic parameters, the platform moves along an unplanned trajectory (see positions (e) - (h) in Fig. 13).

Next, we have implemented the sixth order control laws as it was shown in the previous section and observed the behavior of the platform during the displacement (Fig. 14). The different positions are classified by time. During all these displacements, the manipulator retains its orientation and passes through the singular configuration without any perturbation.

Thus, we can note that the obtained optimum dynamic conditions allow the passing of the manipulator through the singular position.

## 5. Conclusion

At a singular configuration, a manipulator can gain one or more degrees of freedom, and at such a configuration it may becomes uncontrollable, i.e. it may not reproduce stable motion with prescribed trajectory. Nevertheless it is approved that there are several motion planning techniques, which allow passing through these singular zones. These approaches are simulated by numerical examples and illustrated on several parallel structures. It is a promising tendency for the solution of this problem. However, the attention was focused only on control aspects of this problem and very little attention has been paid to the dynamic interpretation, which is a crucial factor for governing the behavior of parallel manipulators at the singular zones.

In this paper we have found the optimal dynamic conditions, for making the pass through the type 2 singular configurations possible. The general definition of the condition for passing through the singular position is formulated as follows: in the presence of type 2 singular configuration, the platform of a parallel manipulator can pass through the singular positions without perturbation of motion if the wrench applied on the platform by the legs and external efforts are orthogonal to the direction of the uncontrollable motion, or in other words, if the work of applied forces and moments on the platform along the uncontrollable motion is equal to zero. This condition has been verified by numerical simulations carried out with the software ADAMS and validated by experimental tests on the prototype of four degrees of freedom parallel manipulator PAMINSA.

It should be noted that the formulated general conditions ensure any given trajectory generation in the manipulator workspace. We would like to point out that the trajectory is not imposed and only the conditions of force generation must be satisfied.

Thus, the passing of any parallel manipulator through the singular positions by the proposed technique is carried out by optimal generation of inertia forces. Hence, it is impossible to stop the manipulator in the singular locus and to start again from fixed position.

We would like to mention that we studied the optimal redistribution of forces only in singular positions of the manipulator but it should be noted that there are zones close to these positions, in which the manipulator loses the quality of motion. For more reliable generation of motion, it is desirable to ensure the given condition of force generation not only in the singular positions of the manipulator but also in the zones near to these positions. It should be also mentioned that a future development of our work is the study of the difficulties of controlling parallel robots in the neighbourhood of singular configurations.

Finally, it should be noted that for the case of non controllable external forces applied on the platform the proposed technique cannot be used. Therefore, the most prominent field of the industrial application is a "fast pick and place" manipulation, when the generation of motion is determined by input, gravitational and inertia forces.

## Appendix A

Coordinates of points $3_{i}, 5_{i}$ and $9_{i}(i=1,2,3)$ :

$$
\left[\begin{array}{l}
x_{3 i} \\
y_{3 i} \\
z_{3 i}
\end{array}\right]=\left[\begin{array}{c}
R_{b} \cos \gamma_{i} \\
R_{b} \sin \gamma_{i} \\
z_{5 i} / k
\end{array}\right],\left[\begin{array}{l}
x_{5 i} \\
y_{5 i} \\
z_{5 i}
\end{array}\right]=\left[\begin{array}{c}
x \\
y \\
z
\end{array}\right]+\left[\begin{array}{c}
R_{p l} \cos \left(\phi+\gamma_{i}\right) \\
R_{p l} \sin \left(\phi+\gamma_{i}\right) \\
L_{c}
\end{array}\right],
$$

where $\gamma_{i}=(-5 \pi / 6,-\pi / 6, \pi / 2)$,

$$
\left[\begin{array}{l}
x_{9 i} \\
y_{9 i} \\
z_{9 i}
\end{array}\right]=\left[\begin{array}{c}
x_{3 i}+X_{9 i} \cos q_{i} \\
y_{3 i}+X_{9 i} \sin q_{i} \\
F
\end{array}\right]
$$

with:

$$
\begin{aligned}
& X_{9 i}=A+B F, F=-(D-K) /(2 E), K=\sqrt{D^{2}-4 E C}, E=-\left(B^{2}+1\right), D=2 B\left(X_{8 i}-A\right), \\
& C=L_{B 3}^{2}-X_{8 i}^{2}+2 A X_{8 i}-A^{2}, B=z_{5 i} /\left(k X_{8 i}\right), A=\left(L_{B 4}^{2}-L_{B 3}^{2}-X_{5 i}^{2}+X_{8 i}^{2}-z_{5 i}^{2}\right) /\left(2 k X_{8 i}\right), \\
& X_{8 i}=-X_{5 i} /(k-1), X_{5 i}=\sqrt{\left(x_{5 i}-x_{3 i}\right)^{2}+\left(y_{5 i}-y_{3 i}\right)^{2}} .
\end{aligned}
$$

## Appendix B

Expressions of terms $C_{v j}(j=1, \ldots, 4)$ :
$C_{v 1}=g\left(m_{5}+\sum_{j=2,3,4,7}\left(\frac{m_{j}}{k}\right)+\frac{m_{B 4}}{2}+\sum_{j=1,7}\left(\frac{m_{B j}}{2 k}\right)+\sum_{j=2}^{3}\left(\frac{m_{B j}}{k}\right)\right)$,
$C_{v 2}=g\left(\frac{(k-1) m_{4}-m_{7}+k m_{9}}{k}+\frac{(k-2) m_{B 3}-m_{B 7}}{2 k}+\frac{m_{B 4}+m_{B 8}}{2}\right)$,
$C_{\nu 3}=g \frac{m_{B 1}}{2}, C_{v 4}=g L_{B 2}\left(m_{2}+\frac{m_{B 1}+m_{B 2}}{2}\right)$.

## Appendix C

Expressions of terms $C_{c j}(j=1, \ldots, 13)$ :
$C_{c 1}=\frac{1}{2}\left(\frac{m_{4}}{k^{2}}+m_{5}+\frac{m_{7}}{(k(k-1))^{2}}+\frac{m_{8}}{(k-1)^{2}}+\frac{m_{B 3}(k-2)^{2}}{(2 k(k-1))^{2}}+\frac{m_{B 4}}{4}+\frac{m_{B 7}(k+1)^{2}}{(2 k(k-1))^{2}}+\frac{m_{B 8}}{(2(k-1))^{2}}\right)$
$C_{c 2}=\frac{1}{2}\left(\sum_{j=2,3,4,7}\left(\frac{m_{j}}{k^{2}}\right)+m_{5}+\sum_{j=1,7}\left(\frac{m_{B j}}{4 k^{2}}\right)+\sum_{j=2}^{3}\left(\frac{m_{B j}}{k^{2}}\right)+\frac{m_{B 4}}{4}\right)$
$C_{c 3}=\frac{1}{2}\left(\frac{(k-1)^{2} m_{4}}{k^{2}}+\frac{m_{7}}{k^{2}}+m_{9}+\frac{m_{B 3}(k-2)^{2}}{4 k^{2}}+\frac{m_{B 4}}{4}+\frac{m_{B 7}}{4 k^{2}}+\frac{m_{B 8}}{4}\right)$
$C_{c 4}=\frac{1}{2}\left(\frac{2(k-1) m_{4}}{k^{2}}+\frac{m_{7}}{k^{2}(k-1)}+\frac{m_{B 3}(k-2)^{2}}{2 k^{2}(k-1)}+\frac{m_{B 4}}{2}+\frac{m_{B 7}(k+1)}{2 k^{2}(k-1)}-\frac{m_{B 8}}{2(k-1)}\right)$
$C_{c 5}=\frac{1}{2}\left(\frac{4\left((k-1) m_{4}-m_{7}\right)-m_{B 7}+2 m_{B 3}(k-2)}{2 k^{2}}+\frac{m_{B 4}}{2}\right)$
$C_{c 6}=\frac{m_{B 1}}{8}, C_{c 7}=\frac{m_{B 1}}{4 k}, C_{c 8}=\frac{m_{B 10} L_{B 10}^{2}}{8}, C_{c 9}=\frac{I_{Y Y}^{(B 4)}+I_{Y Y}^{(B 7)}}{2}, C_{c 10}=\frac{I_{X X}^{(B 4)}+I_{X X}^{(B 7)}}{2}$,
$C_{c 11}=\frac{I_{Y Y}^{(B 3)}+I_{Y Y}^{(B 8)}}{2}, C_{c 12}=\frac{I_{X X}^{(B 3)}+I_{X X}^{(B 8)}}{2}, C_{c 13}=\frac{I_{B 2}+I_{B 10}}{2}$.

## APPENDIX D

$\mathbf{J}_{P A M}=\mathbf{A}^{-\mathbf{1}} \mathbf{B}$ is the global Jacobian matrix where matrice $\mathbf{A}$ and $\mathbf{B}$ are:
$\mathbf{A}=-\left[\begin{array}{cccccc}\sin q_{1} & -\cos q_{1} & 0 & -z_{P C 1} \cos q_{1} & -z_{P C 1} \sin q_{1} & -\mathbf{P C}_{1}^{T} \mathbf{d}_{1} \\ \sin q_{2} & -\cos q_{2} & 0 & -z_{P C 2} \cos q_{2} & -z_{P C 2} \sin q_{2} & -\mathbf{P C}_{2}^{T} \mathbf{d}_{\mathbf{2}} \\ \sin q_{3} & -\cos q_{3} & 0 & -z_{P C 3} \cos q_{3} & -z_{P C 3} \sin q_{3} & -\mathbf{P C}_{\mathbf{3}}^{T} \mathbf{d}_{3} \\ 0 & 0 & -1 & y_{P C 1} & -x_{P C 1} & 0 \\ 0 & 0 & -1 & y_{P C 2} & -x_{P C 2} & 0 \\ 0 & 0 & -1 & y_{P C 3} & -x_{P C 3} & 0\end{array}\right]$,
$\mathbf{B}=\left[\begin{array}{cccc}\rho_{1} & 0 & 0 & 0 \\ 0 & \rho_{2} & 0 & 0 \\ 0 & 0 & \rho_{3} & 0 \\ 0 & 0 & 0 & k \\ 0 & 0 & 0 & k \\ 0 & 0 & 0 & k\end{array}\right]$ with $\rho_{i}=\sqrt{\left(x_{5 i}-x_{3 i}\right)^{2}+\left(y_{5 i}-y_{3 i}\right)^{2}}$,
$\mathbf{P C}_{\mathbf{i}}=\left[x_{P C i}, y_{P C i}, z_{P C i}\right]^{T}=\left[x_{5 i}-x, y_{5 i}-y, z_{5 i}-z\right]^{T}$ and $\mathbf{d}_{i}=\left[\begin{array}{lll}\cos q_{i} & \sin q_{i} & 0\end{array}\right]^{T}($ for $i=1,2,3)$.
The expressions of the terms $\mathbf{W}_{\mathbf{b}}$ and $\mathbf{W}_{\mathbf{p}}$ are:

$$
\begin{aligned}
\mathbf{W}_{\mathbf{b}}= & \sum_{i=1}^{3}\left(\mathbf{J}_{\mathbf{Q} 3 i}^{T} \mathbf{F}_{3 i}+\mathbf{J}_{\mathbf{Q} 4 i}^{T} \mathbf{F}_{4 i}+\mathbf{J}_{\mathbf{Q} 9 i}^{T} \mathbf{F}_{9 i}+\mathbf{J}_{\mathbf{Q} 7 i}^{T} \mathbf{F}_{7 i}+\mathbf{J}_{\mathbf{Q} 2 i}^{T} \mathbf{F}_{2 i}+\mathbf{J}_{\mathbf{Q} S 4 i}^{T} \mathbf{F}_{S 4 i}+\mathbf{J}_{\mathbf{Q} S 3 i}^{T} \mathbf{F}_{S 3 i}+\mathbf{J}_{\mathbf{Q} S 8 i}^{T} \mathbf{F}_{S 8 i}\right. \\
& \left.+\mathbf{J}_{\mathbf{Q} S 7 i}^{T} \mathbf{F}_{S 7 i}+\mathbf{J}_{\mathbf{Q} S 10 i}^{T} \mathbf{F}_{S 10 i}+\mathbf{J}_{\mathbf{Q} S 2 i}^{T} \mathbf{F}_{S 2 i}\right) \\
\mathbf{W}_{\mathbf{p}}= & \mathbf{F}_{\mathbf{P}}+\sum_{i=1}^{3}\left(\mathbf{J}_{\mathbf{X} 8 i}^{T} \mathbf{F}_{8 i}+\mathbf{J}_{\mathbf{X} 5 i}^{T} \mathbf{F}_{5 i}+\mathbf{J}_{\mathbf{X} 9 i}^{T} \mathbf{F}_{9 i}+\mathbf{J}_{\mathbf{X} 4 i}^{T} \mathbf{F}_{4 i}+\mathbf{J}_{\mathbf{X} 7 i}^{T} \mathbf{F}_{7 i}+\mathbf{J}_{\mathbf{X} S 4 i}^{T} \mathbf{F}_{S 4 i}+\mathbf{J}_{\mathbf{X} S 3 i}^{T} \mathbf{F}_{S 3 i}\right. \\
& \left.+\mathbf{J}_{\mathbf{X} S 8 i}^{T} \mathbf{F}_{S 8 i}+\mathbf{J}_{\mathbf{X} 57 i}^{T} \mathbf{F}_{S 7 i}\right)
\end{aligned}
$$



$$
\mathbf{J}_{\mathbf{Q} 3 i}=\left[\begin{array}{cc}
\mathbf{0}_{2 \times 3} & \mathbf{0}_{2 \times 1} \\
\mathbf{0}_{1 \times 3} & 1 \\
\mathbf{0}_{3 \times 3} & \mathbf{0}_{3 \times 1}
\end{array}\right], \mathbf{J}_{\mathbf{Q} 2 i}=\mathbf{J}_{\mathbf{Q} 3 i},
$$

$$
\left.\mathbf{J}_{\mathbf{Q} 9 i}=\left[\begin{array}{ll}
\mathbf{0}_{\mathbf{1 \times 3}} & \delta_{1 i}\left(\mathbf{R}_{\mathbf{P} q i} \operatorname{Rot}\left(\varepsilon_{i}, \boldsymbol{y}\right)\left[\begin{array}{c}
L_{B 3} \\
\mathbf{0}_{2 \times 1}
\end{array}\right]\right)^{T} \\
\mathbf{0}_{\mathbf{1} \times 3} & \delta_{2 i}\left(\mathbf{R}_{\mathbf{P} q i} \operatorname{Rot}\left(\varepsilon_{i}, \boldsymbol{y}\right)\left[\begin{array}{c}
L_{B 3} \\
\mathbf{0}_{2 \times 1}
\end{array}\right]\right)^{T} \\
\mathbf{0}_{\mathbf{1} \times 3} & \delta_{3 i}\left(\mathbf{R}_{\mathbf{P} q i} \boldsymbol{\operatorname { R o t } ( \varepsilon _ { i } , \boldsymbol { y } )}\right. \\
\mathbf{0}_{\mathbf{1} \times 3} & \mathbf{0}_{\mathbf{1 \times 3}}
\end{array}\right]_{\left.\left[\begin{array}{c}
L_{B 3} \\
\mathbf{0}_{2 \times 1}
\end{array}\right]\right)^{T}}\right]^{T}, \mathbf{R}_{\mathbf{P q i}}=\left[\begin{array}{ccc}
-\sin q_{i} & -\cos q_{i} & 0 \\
\cos q_{i} & -\sin q_{i} & 0 \\
0 & 0 & 0
\end{array}\right],
$$

$$
\mathbf{J}_{\mathbf{x} 9 i}=\left[\begin{array}{cc}
\boldsymbol{R o t}\left(q_{i}, \boldsymbol{z}\right) \mathbf{R}_{\mathbf{P} \varepsilon i}\left[\begin{array}{c}
L_{B 3} \\
\mathbf{0}_{\mathbf{3 \times 1}}
\end{array}\right. & \mathbf{0}_{\mathbf{3 \times 1}} \\
\mathbf{0}_{2 \times 1}
\end{array}\right]{\mathbf{\mathbf { 0 } _ { 3 \times 1 }}}_{\mathbf{x}\langle i \varepsilon i}+\mathbf{J}_{\mathbf{x} 8 i}, \mathbf{R}_{\mathbf{P} \varepsilon i}=\left[\begin{array}{ccc}
-\sin \varepsilon_{i} & 0 & \cos \varepsilon_{i} \\
0 & 0 & 0 \\
-\cos \varepsilon_{i} & 0 & -\sin \varepsilon_{i}
\end{array}\right],
$$

$$
\begin{aligned}
& \mathbf{J}_{\mathbf{X} \overline{G i d i}}=-\left[\begin{array}{ll}
L_{B 3} \sin \varepsilon_{i} & L_{B 4} \sin \zeta_{i} \\
L_{B 3} \cos \varepsilon_{i} & L_{B 4} \cos \zeta_{i}
\end{array}\right]^{-1}\left[\begin{array}{cc}
\frac{k}{k-1} & \mathbf{J}_{\mathbf{X} \rho i} \\
{\left[\begin{array}{lll}
\mathbf{0}_{\mathbf{1} 2} & 1 & \mathbf{0}_{\mathbf{1 \times 3}}
\end{array}\right] \mathbf{J}_{\mathbf{X} 5 i}}
\end{array}\right],
\end{aligned}
$$

$$
\begin{aligned}
& \mathbf{J}_{\mathbf{Q} \Omega 1 i}=\left[\begin{array}{cccc}
\mathbf{0}_{2 \times 1} & \mathbf{0}_{2 \times 1} & \mathbf{0}_{2 \times 1} & \mathbf{0}_{2 \times 1} \\
\delta_{1 i} & \delta_{2 i} & \delta_{3 i} & 0
\end{array}\right], \mathbf{J}_{\mathbf{X} 54 i}=0.5\left(\mathbf{J}_{\mathbf{X} 5 i}+\mathbf{J}_{\mathbf{X} 9 i}\right)+\left[\begin{array}{l}
\mathbf{0}_{\mathbf{3 \times 6}} \\
\mathbf{J}_{\mathbf{X} \Omega i}
\end{array}\right], \mathbf{J}_{\mathbf{X} \Omega 1 i}=\left[\begin{array}{cc}
0 & -\sin q_{i} \\
0 & \cos q_{i} \\
0 & 0
\end{array}\right] \mathbf{J}_{\mathbf{X} \leftrightarrows i i}, \\
& \mathbf{J}_{\mathbf{Q} 33 i}=0.5\left(\mathbf{J}_{\mathbf{Q} 4 i}+\mathbf{J}_{\mathbf{Q} 27}\right)+\left[\begin{array}{c}
\mathbf{0}_{3 \times 4} \\
\mathbf{J}_{\mathbf{Q} 22 i}
\end{array}\right], \mathbf{J}_{\mathbf{Q} \Omega 2 i}=\left[\begin{array}{cccc}
\mathbf{0}_{2 \times 1} & \mathbf{0}_{2 \times 1} & \mathbf{0}_{2 \times 1} & \mathbf{0}_{2 \times 1} \\
\delta_{1 i} & \delta_{2 i} & \delta_{3 i} & 0
\end{array}\right],
\end{aligned}
$$

$$
\begin{aligned}
& \mathbf{J}_{\mathbf{X} S 3 i}=0.5\left(\mathbf{J}_{\mathbf{X} 4 i}+\mathbf{J}_{\mathbf{X} 7 i}\right)+\left[\begin{array}{c}
\mathbf{0}_{3 \times 6} \\
\mathbf{J}_{\mathbf{X} \Omega 2 i}
\end{array}\right], \mathbf{J}_{\mathbf{X} \Omega 2 i}=\left[\begin{array}{cc}
-\sin q_{i} & 0 \\
\cos q_{i} & 0 \\
0 & 0
\end{array}\right] \mathbf{J}_{\mathbf{X}\langle i t i}, \mathbf{J}_{\mathbf{Q} 88 i}=0.5\left(\mathbf{J}_{\mathbf{Q} 8 i}+\mathbf{J}_{\mathbf{Q} 9 i}\right)+\left[\begin{array}{c}
\mathbf{0}_{\mathbf{3} \times \mathbf{4}} \\
\mathbf{J}_{\mathbf{Q} \Omega 2 i}
\end{array}\right], \\
& \mathbf{J}_{\mathbf{X} S 8 i}=0.5\left(\mathbf{J}_{\mathbf{x} 8 i}+\mathbf{J}_{\mathbf{x} 9 i}\right)+\left[\begin{array}{c}
\mathbf{0}_{3 \times 6} \\
\mathbf{J}_{\mathbf{x} \Omega 2 i}
\end{array}\right], \mathbf{J}_{\mathbf{Q} S 7 i}=0.5\left(\mathbf{J}_{\mathbf{Q} 8 i}+\mathbf{J}_{\mathbf{Q} 7 i}\right)+\left[\begin{array}{c}
\mathbf{0}_{\mathbf{3 \times 4}} \\
\mathbf{J}_{\mathbf{Q} \Omega 1 i}
\end{array}\right], \\
& \left.\mathbf{J}_{\mathbf{X} \ 7 i}=0.5\left(\mathbf{J}_{\mathbf{X} 8 i}+\mathbf{J}_{\mathbf{X} 7 i}\right)+\left[\begin{array}{c}
\mathbf{0}_{\mathbf{3 \times 6}} \\
\mathbf{J}_{\mathbf{X} \Omega 1 i}
\end{array}\right], \mathbf{J}_{\mathbf{Q} S 10 i}=\left[\frac{\partial\left[x_{S 10 i}, y_{S 10 i}, z_{S 10 i}\right]^{T}}{\partial \mathbf{q}} \mathbf{J}_{\mathbf{Q} \Omega 1 i}\right], \mathbf{J}_{\mathbf{Q} S 2 i}=\left[\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0
\end{array}\right] \mathbf{J}_{\mathbf{Q} 3 i}\right], \\
& \mathbf{F}_{\mathbf{P}}=\left[\begin{array}{llllll}
m_{p l} \ddot{x} & m_{p l} \ddot{y} & m_{p l} & \ddot{z} & 0 & 0
\end{array} I_{p l} \ddot{\phi}\right]^{T}, \\
& \mathbf{F}_{j i}=m_{j}\left[\begin{array}{llllll}
\ddot{x}_{j i} & \ddot{y}_{j i} & \ddot{z}_{j i} & 0 & 0 & 0
\end{array}\right]^{T}, \text { for } j=2,3,4,5,7,8,9 \\
& \mathbf{F}_{B 1 i}=\left[\begin{array}{llllll}
0 & 0 & m_{B 1} \ddot{q}_{v} & 0 & 0 & 0
\end{array}\right]^{T}, \\
& \mathbf{F}_{B j i}=\left[\begin{array}{llllll}
m_{B j} \ddot{x}_{S j i} & m_{B j} \ddot{y}_{S j i} & m_{B 2} \ddot{z}_{S j i} & 0 & 0 & I_{B j} \ddot{q}_{i}
\end{array}\right]^{T}, \text { for } j=2,10 \\
& \mathbf{F}_{B j i}=\left[\begin{array}{llll}
m_{B j} \ddot{x}_{S j i} & m_{B j} \ddot{y}_{S j i} & m_{B 2} \ddot{z}_{S j i} & \mathbf{C}_{S j i}^{T}
\end{array}\right]^{T}, \text { for } j=3,4,7,8 \\
& \mathbf{C}_{S j i}=\left(\mathbf{R}_{\mathbf{P} q i} \operatorname{Rot}\left(\alpha_{i}, \mathbf{y}\right)+\operatorname{Rot}\left(q_{i}, \mathbf{z}\right) \mathbf{R}_{\mathbf{P} \alpha i}\right) \mathbf{I}_{B j}\left(\boldsymbol{\operatorname { R o t }}\left(q_{i}, \mathbf{z}\right) \operatorname{Rot}\left(\alpha_{i}, \mathbf{y}\right)\right)^{T} \boldsymbol{\Omega}_{B j i} \\
& +\boldsymbol{\operatorname { R o t }}\left(q_{i}, \mathbf{z}\right) \boldsymbol{\operatorname { R o t }}\left(\alpha_{i}, \mathbf{y}\right) \mathbf{I}_{B j}\left(\mathbf{R}_{\mathbf{P}_{q i}} \boldsymbol{\operatorname { R o t }}\left(\alpha_{i}, \mathbf{y}\right)+\boldsymbol{\operatorname { R o t }}\left(q_{i}, \mathbf{z}\right) \mathbf{R}_{\mathbf{P} \alpha i}\right)^{T} \boldsymbol{\Omega}_{B j i} \\
& +\boldsymbol{\operatorname { R o t }}\left(q_{i}, \mathbf{z}\right) \boldsymbol{\operatorname { R o t }}\left(\alpha_{i}, \mathbf{y}\right) \mathbf{I}_{B j}\left(\boldsymbol{\operatorname { R o t }}\left(q_{i}, \mathbf{z}\right) \boldsymbol{\operatorname { R o t }}\left(\alpha_{i}, \mathbf{y}\right)\right)^{T} \dot{\boldsymbol{\Omega}}_{B j i}
\end{aligned}
$$

with $\alpha_{i}=\zeta_{i}$ if $j=4,7, \alpha_{i}=\varepsilon_{i}$ if $j=3,8$ and $\boldsymbol{\Omega}_{B j i}=\left[\begin{array}{lll}-\dot{\alpha}_{i} \sin q_{i} & \dot{\alpha}_{i} \cos q_{i} & \dot{q}_{i}\end{array}\right]^{T}$.
In these expressions, $\delta_{i j}$ represents the Krönecker symbol ( $\delta_{i j}=1$ if $j=i$ and $\delta_{i j}=0$ if $j \neq i$ ).

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## Figure Captions

Fig. 1. Kinematic chain of the planar 5R parallel manipulator.
Fig. 2. Second kind of singularities of the planar 5R parallel manipulator.
Fig. 3. Initial, singular and final positions of the planar 5R parallel manipulator.
Fig. 4. Input torques of the planar 5R parallel manipulator in the case of the sixth order polynomial trajectory planning, obtained by the ADAMS software.

Fig. 5. Input torques of the planar 5R parallel manipulator in the case of the fifth order polynomial trajectory planning, obtained by the ADAMS software.

Fig. 6. PAMINSA with 4 DOF.
Fig. 7. Kinematic chain of each leg.
Fig. 8. Example of Type 2 singular configuration (horizontal projection of the examined structure).

Fig. 9. Displacement of the PAMINSA along the prescribed straight line (planar projection).
Fig. 10. Input efforts of the PAMINSA in the case of the sixth order polynomial trajectory planning, computed with ADAMS software.

Fig. 11. Input efforts of the PAMINSA in the case of the fifth order polynomial trajectory planning, computed with ADAMS software.

Fig. 12. The prototype of PAMINSA developed in the I.N.S.A. of Rennes.
Fig. 13. Trajectory reproduction on the PAMINSA during the displacement of the platform with the fifth order polynomial law (view from below).

Fig. 14. Trajectory reproduction on the PAMINSA during the displacement of the platform with the sixth order polynomial law (view from below).

