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# Energy decay for solutions of the wave equation with general memory boundary conditions.

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## Abstract

We consider the wave equation in a smooth domain subject to Dirichlet boundary conditions on one part of the boundary and dissipative boundary conditions of memory-delay type on the remainder part of the boundary, where a general borelian measure is involved. Under quite weak assumptions on this measure, using the multiplier method and a standard integral inequality we show the exponential stability of the system. Some examples of measures satisfying our hypotheses are given, recovering and extending some of the results from the literature.

## Introduction

We consider the wave equation subject to Dirichlet boundary conditions on one part of the boundary and dissipative boundary conditions of memory-delay type on the remainder part of the boundary. More precisely, let  $\Omega$  be a bounded open connected set of  $\mathbb{R}^n$  ( $n \geq 2$ ) such that, in the sense of Nečas ([8]), its boundary  $\partial\Omega$  is of class  $\mathcal{C}^2$ . Throughout the paper,  $I$  denotes the  $n \times n$  identity matrix, while  $A^s$  denotes the symmetric part of a matrix  $A$ . Let  $m$  be a  $\mathcal{C}^1$  vector field on  $\bar{\Omega}$  such that

$$\inf_{\bar{\Omega}} \operatorname{div}(m) > \sup_{\bar{\Omega}} (\operatorname{div}(m) - 2\lambda_m) \quad (1)$$

where  $\lambda_m(x)$  is the smallest eigenvalue function of the real symmetric matrix  $\nabla m(x)^s$ .

**Remark 1** The set of all  $\mathcal{C}^1$  vector fields on  $\bar{\Omega}$  such that (1) holds is an open cone. If  $m$  is in this set, we denote

$$c(m) = \frac{1}{2} \left( \inf_{\bar{\Omega}} \operatorname{div}(m) - \sup_{\bar{\Omega}} (\operatorname{div}(m) - 2\lambda_m) \right).$$

**Example 1** • An affine example is given by

$$m(x) = (A_1 + A_2)(x - x_0),$$

where  $A_1$  is a definite positive matrix,  $A_2$  a skew-symmetric matrix and  $x_0$  any point in  $\mathbb{R}^n$ .

- A non linear example is

$$m(x) = (dI + A)(x - x_0) + F(x)$$

where  $d > 0$ ,  $A$  is a skew-symmetric matrix,  $x_0$  any point in  $\mathbb{R}^n$  and  $F$  is a  $\mathcal{C}^1$  vector field on  $\bar{\Omega}$  such that

$$\sup_{x \in \bar{\Omega}} \|(\nabla F(x))^s\| < \frac{d}{n}$$

( $\|\cdot\|$  stands for the usual 2-norm of matrices).

We define a partition of  $\partial\Omega$  in the following way. Denoting by  $\nu(x)$  the normal unit vector pointing outward of  $\Omega$  at a point  $x \in \partial\Omega$ , we consider a partition  $(\partial\Omega_N, \partial\Omega_D)$  of the boundary such that the measure of  $\partial\Omega_D$  is positive and that

$$\partial\Omega_N \subset \{x \in \partial\Omega, m(x) \cdot \nu(x) \geq 0\}, \partial\Omega_D \subset \{x \in \partial\Omega, m(x) \cdot \nu(x) \leq 0\}. \quad (2)$$

Furthermore, we assume

$$\overline{\partial\Omega_D} \cap \overline{\partial\Omega_N} = \emptyset \text{ or } m \cdot n \leq 0 \text{ on } \overline{\partial\Omega_D} \cap \overline{\partial\Omega_N} \quad (3)$$

where  $n$  stands for the normal unit vector pointing outward of  $\partial\Omega_N$  when considering  $\partial\Omega_N$  as a sub-manifold of  $\partial\Omega$ .

On this domain, we consider the following delayed wave problem:

$$(S) \begin{cases} u'' - \Delta u = 0 & \text{in } \mathbb{R}_+^* \times \Omega, \\ u = 0 & \text{on } \mathbb{R}_+^* \times \partial\Omega_D, \\ \partial_\nu u + m \cdot \nu \left( \mu_0 u'(t) + \int_0^t u'(t-s) d\mu(s) \right) = 0 & \text{on } \mathbb{R}_+^* \times \partial\Omega_N, \\ u(0) = u_0 & \text{in } \Omega, \\ u'(0) = u_1 & \text{in } \Omega, \end{cases}$$

where  $u'$  (resp.  $u''$ ) is the first (resp. second) time-derivative of  $u$ ,  $\partial_\nu u = \nabla u \cdot \nu$  is the normal outward derivative of  $u$  on  $\partial\Omega$ . Moreover  $\mu_0$  is some positive constant and  $\mu$  is a borelian measure on  $\mathbb{R}^+$ .

The above problem covers the case of a problem with memory type as studied for instance in [1, 3, 5, 9], when the measure  $\mu$  is given by

$$d\mu(s) = k(s)ds, \quad (4)$$

where  $ds$  stands for the Lebesgue measure and  $k$  is non negative kernel. But it also covers the case of a problem with a delay as studied for instance in [10, 11, 12], when the measure  $\mu$  is given by

$$\mu = \mu_1 \delta_\tau, \quad (5)$$

where  $\mu_1$  is a non negative constant and  $\tau > 0$  represents the delay. An intermediate case treated in [11] is the case when

$$d\mu(s) = k(s) \chi_{[\tau_1, \tau_2]}(s) ds, \quad (6)$$

where  $0 \leq \tau_1 < \tau_2$ ,  $\chi_{[\tau_1, \tau_2]}$  is the characteristic equation of the interval  $[\tau_1, \tau_2]$  and  $k$  is a non negative function in  $L^\infty([\tau_1, \tau_2])$ .

A closer look at the decay results obtained in these references shows that there are different ways to quantify the energy of  $(S)$ . More precisely for the measure of the form (4), the exponential or polynomial decay of an appropriated energy is proved in [1, 3, 5, 9], by combining the multiplier method (or differential geometry arguments) with the use of suitable Lyapounov functionals (or integral inequalities) under the assumptions that the kernel  $k$  is sufficiently smooth and has a certain decay at infinity. On the other hand for a measure of delay type like (5) or (6), the exponential stability of the system was proved in [10, 11, 12] by proving an observability estimate obtained by assuming that the term  $\int_0^t u'(t-s) d\mu(s)$  is sufficiently small with respect to  $\mu_0 u'(t)$ . Consequently, our goal is here to obtain some uniform decay results in the general context described above with a similar assumptions than in [10, 11, 12]. More precisely, we will show in this paper that if there exists  $\alpha > 0$  such that

$$\mu_{\text{tot}} := \int_0^{+\infty} e^{\alpha s} d|\mu|(s) < \mu_0 \quad (7)$$

where  $|\mu|$  is the absolute value of the measure  $\mu$ , then the above problem  $(S)$  is exponentially stable.

The paper is organized as follows: in the first two sections, we explain how to define an energy using some basic measure theory. Using well-known results, we obtain the existence of energy solutions. In this setting we present and prove our stabilization result in the third section. Examples of measures  $\mu$  satisfying our hypotheses are given in the end of the paper, where we show that we recover and extend some of the results from the references cited above.

Finally in the whole paper we use the notation  $A \lesssim B$  for the estimate  $A \leq CB$  with some constant  $C$  that only depends on  $\Omega$ ,  $m$  or  $\mu$ .

## 1 First results

In this section we show that the assumption (7) implies the existence of some borelian finite measure  $\lambda$  such that

$$\lambda(\mathbb{R}^+) < \mu_0, \quad |\mu| \leq \lambda \quad (8)$$

(in the sense that, for every measurable set  $\mathcal{B}$ ,  $|\mu|(\mathcal{B}) \leq \lambda(\mathcal{B})$ ) and

$$\text{for all measurable set } \mathcal{B}, \int_{\mathcal{B}} \lambda([s, +\infty)) ds \leq \alpha^{-1} \lambda(\mathcal{B}) \quad (9)$$

Indeed we show the following equivalence:

**Proposition 1** *Let  $\mu$  be a borelian positive measure on  $\mathbb{R}^+$  and  $\mu_0$  some positive constant. The following properties are equivalent:*

- $\exists \alpha > 0$  such that

$$\int_0^{+\infty} e^{\alpha s} d\mu(s) < \mu_0.$$

- There exists a borelian measure  $\lambda$  on  $\mathbb{R}^+$  such that

$$\lambda(\mathbb{R}^+) < \mu_0, \quad \mu \leq \lambda$$

and, for some constant  $\beta > 0$ ,

$$\text{for all measurable set } \mathcal{B}, \int_{\mathcal{B}} \lambda([s, +\infty)) ds \leq \beta^{-1} \lambda(\mathcal{B}).$$

**Proof** We introduce the application  $T$  from the set of positive borelian measures into itself as follows: if  $\mu$  is some positive borelian measure, we define a positive borelian measure  $T(\mu)$  by

$$T(\mu)(\mathcal{B}) = \int_{\mathcal{B}} \mu([s, +\infty)) ds,$$

if  $\mathcal{B}$  is any measurable set.

( $\Leftarrow$ ) If  $\lambda$  fulfills the second property, then it immediately follows that

$$\forall n \in \mathbb{N}, \beta^n T^n(\mu) \leq \lambda,$$

where as usual  $T^n$  is the composition  $T \circ T \cdots \circ T$   $n$ -times. A summation consequently gives, for any  $r \in (0, 1)$ ,

$$\sum_{n=0}^{\infty} (r\beta)^n T^n(\mu) \leq \sum_{n=0}^{\infty} r^n \lambda = (1-r)^{-1} \lambda.$$

Using Fubini theorem, we can now compute

$$\begin{aligned} T^n(\mu)(\mathbb{R}^+) &= \int_0^{+\infty} \left( \int_{s_{n+1}}^{+\infty} \cdots \int_{s_3}^{+\infty} \left( \int_{s_2}^{+\infty} d\mu(s_1) \right) ds_2 \cdots ds_n \right) ds_{n+1} \\ &= \int_0^{+\infty} \left( \int_0^{s_1} \cdots \left( \int_0^{s_n} ds_{n+1} \right) \cdots ds_2 \right) d\mu(s_1) \\ &= \int_0^{+\infty} \frac{s_1^n}{n!} d\mu(s_1) \end{aligned}$$

so that, using monotone convergence theorem, one can obtain

$$\int_0^{+\infty} e^{r\beta s} d\mu(s) \leq (1-r)^{-1} \lambda(\mathbb{R}^+)$$

and our proof ends using that  $(1-r)^{-1} \lambda(\mathbb{R}^+) < \mu_0$  for sufficiently small  $r$ .

( $\Rightarrow$ ) For any measurable set  $\mathcal{B}$ , we define

$$\lambda(\mathcal{B}) = \sum_{n=0}^{\infty} \alpha^n T^n(\mu)(\mathcal{B}).$$

It is clear that  $\lambda$  is a borelian measure such that  $\mu \leq \lambda$ . Moreover, if  $\mathcal{B}$  is a measurable set, one has, thanks to monotone convergence theorem

$$\begin{aligned} T(\lambda)(\mathcal{B}) &= \int_{\mathcal{B}} \lambda([s, +\infty)) ds \\ &= \sum_{n=0}^{\infty} \alpha^n \int_{\mathcal{B}} T^n(\mu)([s, +\infty)) ds \\ &\leq \alpha^{-1} \sum_{n=0}^{\infty} \alpha^{n+1} T^{n+1}(\mu)(\mathcal{B}), \end{aligned}$$

that is,  $T(\lambda) \leq \alpha^{-1} \lambda$ .

Finally, another use of monotone convergence theorem gives

$$\lambda(\mathbb{R}^+) = \int_0^{+\infty} e^{\alpha s} d\mu(s) < \mu_0.$$

□

**Remark 2** • If  $\mu$  satisfies our first property, one can choose  $\beta = \alpha$  in our second property.

- If  $\mu$  is supported in  $(0, \tau]$ , it is straightforward to see that, for some small enough constant  $c$ ,

$$d\lambda(s) = d\mu(s) + c \chi_{[0, \tau]}(s) ds$$

fulfills (7). This observation allows us to recover the choices of energy in [10, 11].

In the sequel, we can thus consider the measure  $\lambda$  obtained by the application of Proposition 1 to  $|\mu|$ .

## 2 Well-posedness

### 2.1 General results

Defining

$$H_D^1(\Omega) := \{u \in H^1(\Omega); u = 0 \text{ on } \partial\Omega_D\} \text{ and } H_0^1(\Omega) := \{u \in H^1(\Omega); u = 0 \text{ on } \partial\Omega\},$$

we here present an application of Theorem 4.4 of Propst and Prüss paper (see [13]) in the framework of hypothesis (3).

**Theorem 1** *Suppose  $u_0 \in H_D^1(\Omega)$ ,  $u_1 \in L^2(\Omega)$ . Then (S) admits a unique solution  $u \in \mathcal{C}(\mathbb{R}^+, H^1(\Omega)) \cap \mathcal{C}^1(\mathbb{R}^+, L^2(\Omega))$  in the weak sense of Propst and Prüss. Moreover, if  $u_0 \in H^2(\Omega) \cap H_D^1(\Omega)$ ,  $u_1 \in H_0^1(\Omega)$ , then  $u \in \mathcal{C}^1(\mathbb{R}^+, H^1(\Omega)) \cap \mathcal{C}^2(\mathbb{R}^+, L^2(\Omega))$  and the additional results hold*

$$\forall t \geq 0, \Delta u(t) \in L^2(\Omega) \quad \partial_\nu u(t)|_{\partial\Omega_N} \in H^{1/2}(\partial\Omega_N).$$

**Proof** The proof is the one proposed in [13], Theorem 4.4 except that, for smoother data, we can not use elliptic result in the general context of (3) to get more regularity.  $\square$

Inspired by [10, 11], we now define the energy of the solution of (S) at any positive time  $t$  by the following formula:

$$\begin{aligned} E(t) &= \frac{1}{2} \int_{\Omega} (u'(t, x))^2 + |\nabla u(t, x)|^2 dx + \frac{1}{2} \int_{\partial\Omega_N} m \cdot \nu \int_0^t \left( \int_0^s (u'(t-r, x))^2 dr \right) d\lambda(s) d\sigma \\ &+ \frac{1}{2} \int_{\partial\Omega_N} m \cdot \nu \int_t^\infty \left( \int_0^s (u'(s-r, x))^2 dr \right) d\lambda(s) d\sigma. \end{aligned}$$

**Remark 3** • In the definition of energy, the measure  $\lambda$  can be replaced by any positive borelian measure  $\nu$  such that

$$\nu \leq \lambda$$

such as, for instance,  $|\mu|$ . In fact, we will see later that conditions (8) and (9) are only here to ensure that the corresponding energy  $E_\lambda$  is non increasing, but the decay of another energy  $E_\nu$  is implied by the decay of  $E_\lambda$ .

- If  $\mu$  is compactly supported in  $[0, \tau]$ , for times greater than  $\tau$ , one can recover the energies from [10, 11] by choosing the measure  $\lambda$  supported in  $[0, \tau]$  given by Remark 2. Indeed, the last term in the energy is null for  $t > \tau$ , and the second term is reduced to

$$\frac{1}{2} \int_{\partial\Omega_N} m \cdot \nu \int_0^\tau \left( \int_0^s (u'(t-r, x))^2 dr \right) d\lambda(s) d\sigma.$$

We now identify our energy space.

**Proposition 2** *If  $u_0 \in H^2(\Omega) \cap H_D^1(\Omega)$ ,  $u_1 \in H_0^1(\Omega)$ , then  $u' \in L^\infty(\mathbb{R}^+, H^1(\Omega))$ . Consequently, for such initial conditions, the energy  $E(t)$  is well defined for any  $t > 0$  and it uniformly depends continuously on the initial data.*

**Proof** Let us first pick some solution of (S) with  $u_0 \in H^2(\Omega) \cap H_D^1(\Omega)$ ,  $u_1 \in H_D^1(\Omega)$ . We define the standard energy as

$$E_0(t) = \frac{1}{2} \int_{\Omega} (u'(t, x))^2 + |\nabla u(t, x)|^2 dx.$$

As in [6], it is classical that

$$E_0(0) - E_0(T) = - \int_0^T \int_{\partial\Omega_N} \partial_\nu u u' d\sigma dt.$$

Using the form of our boundary condition and Young inequality, one gets, for any  $\epsilon > 0$ ,

$$\begin{aligned} E_0(0) - E_0(T) &= \int_0^T \int_{\partial\Omega_N} (m \cdot \nu) \left( \mu_0 u'(t)^2 + u'(t) \int_0^t u'(t-s) d\mu(s) \right) d\sigma dt \\ &\geq \int_0^T \int_{\partial\Omega_N} (m \cdot \nu) \left( \left( \mu_0 - \frac{\epsilon}{2} \right) u'(t)^2 - \frac{1}{2\epsilon} \left( \int_0^t u'(t-s) d\mu(s) \right)^2 \right) d\sigma \end{aligned}$$

Using that  $\mu \leq |\mu|$  and Cauchy-Schwarz inequality consequently give us

$$E_0(0) - E_0(T) \geq \int_{\partial\Omega_N} (m \cdot \nu) \left( \left( \mu_0 - \frac{\epsilon}{2} \right) \int_0^T u'(t)^2 dt - \frac{\mu_{\text{tot}}}{2\epsilon} \int_0^T \int_0^t (u'(t-s))^2 d|\mu|(s) \right) d\sigma.$$

Using now Fubini theorem two times, one can obtain the following identities

$$\begin{aligned}
\int_0^T \int_0^t (u'(t-s))^2 d|\mu|(s) dt &= \int_0^T \left( \int_s^T u'(t-s)^2 dt \right) d|\mu|(s) \\
&= \int_0^T \left( \int_0^{T-s} u'(t)^2 dt \right) d|\mu|(s) \\
&= \int_0^T \left( \int_0^{T-t} d|\mu|(s) \right) u'(t)^2 dt
\end{aligned}$$

so that, using  $|\mu|([0, T-t]) \leq \mu_0$ ,

$$E_0(0) - E_0(T) \geq \int_{\partial\Omega_N} (m \cdot \nu) \left( \left( \mu_0 - \frac{\epsilon}{2} - \frac{\mu_{\text{tot}}^2}{2\epsilon} \right) \int_0^T u'(t)^2 dt \right) d\sigma.$$

The choice of  $\epsilon = \mu_{\text{tot}}$  finally gives us that  $E_0(T)$  is bounded. Using the density of  $H^2(\Omega) \cap H_D^1(\Omega) \times H_D^1(\Omega)$  in  $H_D^1(\Omega) \times L^2(\Omega)$ , we get the boundedness of  $E_0$  for solutions with initial data  $u_0 \in H_D^1(\Omega), u_1 \in L^2(\Omega)$ . In particular, if  $u_0 \in H_D^1(\Omega), u_1 \in L^2(\Omega)$ , we obtain that  $u \in L^\infty(\mathbb{R}^+, H^1(\Omega))$ .

Let now  $u$  be a solution of  $(S)$  with  $u_0 \in H^2(\Omega) \cap H_D^1(\Omega), u_1 \in H_0^1(\Omega)$ . Using Theorem 1, one can define the limit in  $L^2(\Omega)$   $u_2$  of  $u''(t)$  as  $t \rightarrow 0$  and in this situation, as in [13], it is easy to see that  $u'$  is solution of  $(S)$  with initial data  $u_1 \in H_D^1(\Omega)$  and  $u_2 \in L^2(\Omega)$ .

Indeed, one can see that Fubini's theorem gives

$$\int_0^t u'(t-s) d\mu(s) = \int_0^t \left( \int_0^s u''(s-r) d\mu(r) \right) ds,$$

provided  $u_1 = 0$  on  $\partial\Omega_N$ , so that

$$\frac{d}{dt} \left( \int_0^t u'(t-s) d\mu(s) \right) = \int_0^t u''(t-s) d\mu(s).$$

Using the proof above, one concludes that  $u' \in L^\infty(\mathbb{R}^+, H^1(\Omega))$  which, thanks to a classical trace result, give that  $u' \in L^\infty(\mathbb{R}^+, L^2(\partial\Omega))$ .

The first three terms of the energy  $E(t)$  are consequently defined for any time  $t > 0$ . We only need to take a look at the last one to achieve our result. Using Fubini theorem again, one has

$$\int_t^{+\infty} \left( \int_0^s (u'(t-r))^2 dr \right) d\lambda(s) = \int_0^{+\infty} u'(r)^2 \left( \int_{\max(r,t)}^{+\infty} d\lambda(s) \right) dr$$

so that, using (9),

$$\begin{aligned}
\int_{\partial\Omega_N} m \cdot \nu \int_t^\infty \left( \int_0^s (u'(s-r, x))^2 dr \right) d\lambda(s) d\sigma &\leq \|m\|_\infty \|u'\|_{L^\infty(L^2(\partial\Omega))} \int_0^{+\infty} \lambda([r, +\infty)) dr \\
&\leq \alpha^{-1} \|m\|_\infty \lambda(\mathbb{R}^+) \|u'\|_{L^\infty(L^2(\partial\Omega))}.
\end{aligned}$$

□

## 2.2 Compactly supported measure and semigroup approach

In this second approach, we assume that  $\mu$  is supported in  $[0, \tau]$  and that  $\overline{\partial\Omega_D} \cap \overline{\partial\Omega_N} = \emptyset$ . We here simply follow the result obtained by Nicaise-Pignotti ([11]).

First, observe that, for  $t > \tau$ ,  $(S)$  is reduced to

$$\begin{cases} u'' - \Delta u = 0 & \text{in } (\tau, +\infty) \times \Omega, \\ u = 0 & \text{on } (\tau, +\infty) \times \partial\Omega_D, \\ \partial_\nu u + m \cdot \nu (\mu_0 u'(t) + \int_0^\tau u'(t-s) d\mu(s)) = 0 & \text{on } (\tau, +\infty) \times \partial\Omega_N, \\ u(0) = u_0 & \text{in } \Omega, \\ u'(0) = u_1 & \text{in } \Omega. \end{cases}$$

We define  $X_\tau = L^2(\partial\Omega_N \times (0, 1) \times (0, \tau), d\sigma d\rho s d\mu(s))$  and  $Y_\tau = L^2(\partial\Omega_N \times (0, \tau); H^1(0, 1), d\sigma s d\mu(s))$ . One can use the same strategy as in the proof of Theorem 2.1 in [11] to get

**Theorem 2** • If  $u(\tau) \in H_D^1(\Omega)$ ,  $u'(\tau) \in L^2(\Omega)$  and  $u'(\tau - \rho s, x) \in X_\tau$ , (S) has a unique solution  $u \in \mathcal{C}([\tau, +\infty), H_D^1(\Omega)) \cap \mathcal{C}^1([\tau, +\infty), L^2(\Omega))$ . Moreover, if  $u(\tau) \in H^2(\Omega) \cap H_D^1(\Omega)$ ,  $u'(\tau) \in H^1(\Omega)$  and  $u'(\tau - \rho s, x) \in Y_\tau$ , then

$$\begin{cases} u \in \mathcal{C}^1([\tau, +\infty), H_D^1(\Omega)) \cap \mathcal{C}([\tau, +\infty), H^2(\Omega)); \\ t \mapsto su''(t - \rho s, x) \in \mathcal{C}([\tau, +\infty), X_\tau). \end{cases}$$

• If  $(u_\tau^n(x), v_\tau^n(x), g^n(s, \rho, x)) \rightarrow (u(\tau, x), u'(\tau, x), u'(\tau - \rho s, x))$  in  $H_D^1(\Omega) \times L^2(\Omega) \times X_\tau$ , then the solution  $u^n$  of

$$\begin{cases} u'' - \Delta u = 0 & \text{in } (\tau, +\infty) \times \Omega, \\ u = 0 & \text{on } (\tau, +\infty) \times \partial\Omega_D, \\ \partial_\nu u + m \cdot \nu (\mu_0 u'(t) + \int_0^\tau u'(t-s) d\mu(s)) = 0 & \text{on } (\tau, +\infty) \times \partial\Omega_N, \\ u(\tau) = u_\tau^n & \text{in } \Omega, \\ u'(\tau) = u_\tau^n & \text{in } \Omega, \\ u'(x, \tau - \rho s) = g^n(x, s, \rho) & \text{in } \Omega_N \times (0, \tau) \times (0, 1) \end{cases}$$

is such that  $E(u^n)$  converges uniformly with respect to time towards  $E(u)$ .

**Proof** We define  $z(x, \rho, s, t) = u'(t - \rho s, x)$  for  $x \in \partial\Omega_N$ ,  $t > \tau$ ,  $s \in (0, \tau)$ ,  $\rho \in (0, 1)$ . Problem (S) is then equivalent to

$$\begin{aligned} u'' - \Delta u &= 0 && \text{in } (\tau, +\infty) \times \Omega, \\ sz_t(x, \rho, s, t) + z_\rho(x, \rho, s, t) &= 0 && \text{in } \partial\Omega_N \times (0, 1) \times (0, \tau) \times (\tau, +\infty), \\ u &= 0 && \text{on } (\tau, +\infty) \times \partial\Omega_D, \\ \partial_\nu u + m \cdot \nu (\mu_0 u'(t) + \int_0^\tau u'(t-s) d\mu(s)) &= 0 && \text{on } (\tau, +\infty) \times \partial\Omega_N, \\ u(\tau) &= u(\tau) && \text{in } \Omega, \\ u'(\tau) &= u(\tau) && \text{in } \Omega, \\ z(x, 0, t, s) &= u'(t, x) && \text{on } \partial\Omega_N \times (\tau, +\infty) \times (0, \tau), \\ z(x, \rho, \tau, s) &= f_0(x, \rho, s) && \text{on } \partial\Omega_N \times (0, 1) \times (0, \tau), \end{aligned}$$

where  $f_0(x, \rho, s) = u'(\tau - \rho s, x)$ .

Consequently, (S) can be rewritten as

$$\begin{cases} U' = \mathcal{A}U \\ U(\tau) = (u(\tau), u'(\tau), f_0)^T \end{cases}$$

where the operator is defined by

$$\mathcal{A} \begin{pmatrix} u \\ v \\ z \end{pmatrix} = \begin{pmatrix} v \\ \Delta u \\ -s^{-1}z_\rho \end{pmatrix}$$

with domain

$$\mathcal{D}(\mathcal{A}) = \left\{ (u, v, z)^T \in H_D^1(\Omega) \times L^2(\Omega) \times Y_\tau; \Delta u \in L^2(\Omega), \right. \\ \left. \partial_\nu u(x) = -(m \cdot \nu) \left( \mu_0 v(t) + \int_0^\tau z(x, 1, s) d\mu(s) \right) \text{ on } \partial\Omega_N, v(x) = z(x, 0, s) \text{ on } \partial\Omega_N \times (0, \tau) \right\}.$$



The proof of Theorem 2.1 in [11] shows us that  $\mathcal{A}$  is a maximal monotone operator on the Hilbert space  $\mathcal{H} := H_D^1(\Omega) \times L^2(\Omega) \times X_\tau$  endowed with the product topology. It consequently generates a contraction semigroup on  $\mathcal{H}$ . Moreover, if  $(u(\tau, x), u'(\tau, x), u'(\tau - \rho s, x)) \in \mathcal{D}(\mathcal{A})$ , one gets that

$$\begin{cases} u \in \mathcal{C}^1([\tau, +\infty), H_D^1(\Omega)) \cap \mathcal{C}([\tau, +\infty), H^2(\Omega)); \\ t \mapsto su''(t - \rho s, x) \in \mathcal{C}([\tau, +\infty), X_\tau). \end{cases}$$

This ends the proof.  $\square$

We can consequently deduce another way to obtain solutions:

**Corollary 1** *Suppose that  $u_0 \in H^2(\Omega) \cap H_D^1(\Omega)$ ,  $u_1 \in H_0^1(\Omega)$ , then  $(S)$  has a unique solution  $u \in \mathcal{C}([\tau, +\infty), H_D^1(\Omega)) \cap \mathcal{C}^1([\tau, +\infty), L^2(\Omega))$ .*

**Proof** Thanks to Theorem 1, one only needs to check that if  $u \in \mathcal{C}^1([0, \tau], H_D^1(\Omega))$  then  $u'(x, \tau - \rho s) \in X_\tau$ ; and this is straightforward using Fubini theorem.  $\square$

### 3 Linear stabilization

We begin with a classical elementary result due to Komornik [6]:

**Lemma 1** *Let  $E : [0, +\infty[ \rightarrow \mathbb{R}_+$  be a non-decreasing function that fulfils:*

$$\forall t \geq 0, \quad \int_t^\infty E(s) ds \leq TE(t),$$

*for some  $T > 0$ . Then, one has:*

$$\forall t \geq T, E(t) \leq E(0) \exp\left(1 - \frac{t}{T}\right).$$

We will now show the following stabilization result:

**Theorem 3** *Assume (1)-(7). Then, if  $u_0 \in H^2(\Omega) \cap H_D^1(\Omega)$ ,  $u_1 \in H_0^1(\Omega)$ , there exists  $T > 0$  such that the energy  $E(t)$  of the solution  $u$  of  $(S)$  satisfies:*

$$\forall t \geq T, \quad E(t) \leq E(0) \exp\left(1 - \frac{t}{T}\right).$$

**Proof** Our goal is to perform the multiplier method and to deal with the delay terms to show that one can apply Lemma 1 to the energy.

**Lemma 2** *There exists  $C > 0$ , such that, for any solution  $u$  of  $(S)$  and any  $S \leq T$ ,*

$$E(S) - E(T) \geq C \int_S^T \int_{\partial\Omega_N} (m \cdot \nu) \left( (u'(t))^2 + \int_0^t (u'(t-s))^2 d\lambda(s) \right) d\sigma dt.$$

*In particular, the energy is a non-increasing function of time.*

**Proof** We start from the classical result that

$$E_0(S) - E_0(T) = - \int_S^T \int_{\partial\Omega} \partial_\nu uu' d\sigma dt.$$

As above, one gets, for any  $\epsilon > 0$ ,

$$E_0(S) - E_0(T) \geq \int_S^T \int_{\partial\Omega_N} (m \cdot \nu) \left( \left( \mu_0 - \frac{\epsilon}{2} \right) u'(t)^2 - \frac{\mu_{\text{tot}}}{2\epsilon} \int_0^t \int_0^t (u'(t-s))^2 d|\mu|(s) \right) d\sigma dt$$

We will now split  $E - E_0$  in two terms:

$$[E - E_0]_S^T = -\frac{1}{2} \left( \int_{\partial\Omega_N} (m \cdot \nu) [f(t, x) - g(t, x)]_S^T d\sigma \right),$$

where

$$\begin{aligned} f(t, x) &= \int_0^t \left( \int_0^s (u'(t-r))^2 dr \right) d\lambda(s), \\ g(t, x) &= \int_0^t \left( \int_0^s u'(r)^2 dr \right) d\lambda(s). \end{aligned}$$

A change of variable allows us to get

$$f(t, x) = \int_0^t \int_0^t u'(r)^2 dr d\lambda(s) - \int_0^t \int_0^{t-s} u'(r)^2 dr d\lambda(s).$$

An application of Fubini theorem consequently gives us

$$f(t, x) - g(t, x) = \int_0^t u'(r)^2 \lambda([0, r]) dr - \int_0^t \int_0^{t-s} u'(r)^2 dr d\lambda(s)$$

and, as above, one can use Fubini theorem to deduce that

$$\int_S^T \int_0^t (u'(t-s))^2 d\lambda(s) dt = \left[ \int_0^t \int_0^{t-s} u'(r)^2 dr d\lambda(s) \right]_S^T.$$

One now uses  $\lambda([0, r]) \leq \lambda(\mathbb{R}^+)$  to conclude that

$$[E - E_0]_S^T \geq \frac{1}{2} \int_S^T \int_{\partial\Omega_N} m \cdot \nu \left( \int_0^t (u'(t-s))^2 d\lambda(s) - \lambda(\mathbb{R}^+) u'(t)^2 \right) d\sigma dt.$$

Summing up and using that  $|\mu| \leq \lambda$ , we have obtained that

$$E(S) - E(T) \geq \int_S^T \int_{\partial\Omega_N} (m \cdot \nu) \left( \left( \mu_0 - \frac{\lambda(\mathbb{R}^+) + \epsilon}{2} \right) u'(t)^2 + \frac{1}{2} \left( 1 - \frac{\mu_{\text{tot}}}{\epsilon} \right) \int_0^t (u'(t-s))^2 d\lambda(s) \right) d\sigma dt.$$

We finally chose  $\epsilon = \mu_0$  which gives us our result since  $\lambda(\mathbb{R}^+) < \mu_0$ .  $\square$

In the multiplier method, one may use Rellich's relation, especially in the context of singularities. In our framework (3), the following Rellich inequality (see the proof of Theorem 4 in [4] or Proposition 4 in [2]) is useful

**Proposition 3** *For any  $u \in H^1(\Omega)$  such that*

$$\Delta u \in L^2(\Omega), u|_{\partial\Omega_D} \in H^{\frac{3}{2}}(\partial\Omega_D) \text{ and } \partial_\nu u|_{\partial\Omega_N} \in H^{\frac{1}{2}}(\partial\Omega_N).$$

*Then it satisfies  $2\partial_\nu u(m \cdot \nabla u) - (m \cdot \nu)|\nabla u|^2 \in L^1(\partial\Omega)$  and we have the following inequality*

$$2 \int_\Omega \Delta u(m \cdot \nabla u) dx \leq \int_\Omega (\operatorname{div}(m)I - 2(\nabla m)^s)(\nabla u, \nabla u) dx + \int_{\partial\Omega} (2\partial_\nu u(m \cdot \nabla u) - (m \cdot \nu)|\nabla u|^2) d\sigma.$$

With this result, we can prove the following multiplier estimate:

**Lemma 3** *Let  $Mu = 2m \cdot \nabla u + a_0 u$ , where  $a_0 := \frac{1}{2} (\inf_{\bar{\Omega}} \operatorname{div}(m)) + \sup_{\bar{\Omega}} (\operatorname{div}(m) - 2\lambda_m)$ . Then under the assumptions of Theorem 3, the following inequality holds true:*

$$\begin{aligned} c(m) \int_S^T \int_\Omega (u')^2 + |\nabla u|^2 dx dt &\leq - \left[ \int_\Omega u' Mu \right]_S^T \\ &+ \int_S^T \int_{\partial\Omega_N} Mu \partial_\nu u + (m \cdot \nu)((u')^2 - |\nabla u|^2) d\sigma dt. \end{aligned}$$

**Proof** Firstly, we consider  $M = 2m \cdot \nabla u + au$  where  $a$  will be fixed later. Using the fact that  $u$  is a regular solution of  $(S)$  and noting that  $u''Mu = (u'Mu)' - u'Mu'$ , an integration by parts gives:

$$\begin{aligned} 0 &= \int_S^T \int_{\Omega} (u'' - \Delta u) M u dx dt \\ &= \left[ \int_{\Omega} u' M u dx \right]_S^T - \int_S^T \int_{\Omega} (u' M u' + \Delta u M u) dx dt. \end{aligned}$$

Now, thanks to Proposition 3, we have :

$$\begin{aligned} \int_{\Omega} \Delta u M u dx &\leq a \int_{\Omega} \Delta u u dx + \int_{\Omega} (\operatorname{div}(m)I - 2(\nabla m)^s)(\nabla u, \nabla u) dx \\ &\quad + \int_{\partial\Omega} (2\partial_{\nu} u(m \cdot \nabla u) - (m \cdot \nu)|\nabla u|^2) d\sigma. \end{aligned}$$

Consequently, Green-Riemann formula leads to:

$$\int_{\Omega} \Delta u M u dx = \int_{\Omega} ((\operatorname{div}(m) - a)I - 2(\nabla m)^s)(\nabla u, \nabla u) dx + \int_{\partial\Omega} (\partial_{\nu} u M u - (m \cdot \nu)|\nabla u|^2) d\sigma.$$

Using the fact that  $\nabla u = \partial_{\nu} u \nu$  on  $\partial\Omega_D$  and  $m \cdot \nu \leq 0$  on  $\partial\Omega_D$ , we have then:

$$\int_{\Omega} \Delta u M u dx \leq \int_{\Omega} ((\operatorname{div}(m) - a)I - 2(\nabla m)^s)(\nabla u, \nabla u) dx + \int_{\partial\Omega_N} (\partial_{\nu} u M u - (m \cdot \nu)|\nabla u|^2) d\sigma.$$

On the other hand, another use of Green formula gives us:

$$\int_{\Omega} u' M u' dx = \int_{\Omega} (a - \operatorname{div}(m))(u')^2 dx + \int_{\partial\Omega_N} (m \cdot \nu)|u'|^2 d\sigma.$$

Consequently

$$\begin{aligned} &\int_S^T \int_{\Omega} (\operatorname{div}(m) - a)(u')^2 + ((a - \operatorname{div}(m))I + 2(\nabla m)^s)(\nabla u, \nabla u) dx dt \\ &\leq - \left[ \int_{\Omega} u' M u dx \right]_S^T + \int_S^T \int_{\partial\Omega_N} \partial_{\nu} u M u + (m \cdot \nu)((u')^2 - |\nabla u|^2) d\sigma dt. \end{aligned}$$

Our goal is now to find  $a$  such that  $\operatorname{div}(m) - a$  and  $(a - \operatorname{div}(m))I + 2(\nabla m)^s$  are uniformly minorized on  $\Omega$ . One has to find  $a$  such that, uniformly on  $\Omega$ ,

$$\begin{cases} \operatorname{div}(m) - a \geq c \\ 2\lambda_m + (a - \operatorname{div}(m)) \geq c \end{cases} \quad (10)$$

for some positive constant  $c$ . The latter condition is then equivalent to find  $a$  which fulfills

$$\inf_{\Omega} \operatorname{div}(m) > a > \sup_{\Omega} (\operatorname{div}(m) - 2\lambda_m),$$

and its existence is now guaranteed by (1). Moreover, it is straightforward to see that the greatest value of  $c$  such that (10) holds is

$$c(m) = \frac{1}{2} \left( \inf_{\Omega} \operatorname{div}(m) - \sup_{\Omega} (\operatorname{div}(m) - 2\lambda_m) \right)$$

and is obtained for  $a = a_0$ . This ends the proof.  $\square$

Consequently, the following result holds

**Lemma 4** *For every  $\tau \leq S < T < \infty$ , the following inequality holds true:*

$$\int_S^T \int_{\Omega} (u')^2 + |\nabla u|^2 dx dt \lesssim E(S).$$

**Proof** We start from Lemma 3.

First of all, Young and Poincaré inequalities give

$$|\int_{\Omega} u' M u dx| \lesssim E(t),$$

so that

$$-\left[\int_{\Omega} u' M u dx\right]_S^T \lesssim E(S) + E(T) \leq CE(S).$$

Now, from the boundary condition, one has

$$Mu \partial_{\nu} u + (m \cdot \nu)((u')^2 - |\nabla u|^2) = (m \cdot \nu) \left( \left( \mu_0 u' + \int_0^t u'(t-s) d\mu(s) \right) Mu + (u')^2 - |\nabla u|^2 \right).$$

Using the definition of  $Mu$  and Young inequality, we get for any  $\epsilon > 0$

$$Mu \partial_{\nu} u + (m \cdot \nu)((u')^2 - |\nabla u|^2) \leq (m \cdot \nu) \left( \left( 1 + \|m\|_{\infty}^2 + \mu_0^2 \frac{a_0^2}{2\epsilon} \right) (u')^2 + \frac{a_0^2}{2\epsilon} \left( \int_0^t u'(t-s) d\mu(s) \right)^2 + \epsilon u^2 \right).$$

Another use of Poincaré inequality consequently allow us to choose  $\epsilon > 0$  such that

$$\epsilon \int_{\partial\Omega_N} (m \cdot \nu) u^2 d\sigma \leq \frac{c(m)}{2} \int_{\Omega} |\nabla u|^2 dx.$$

Cauchy-Schwarz inequality consequently leads to

$$\frac{c(m)}{2} \int_S^T \int_{\Omega} (u')^2 + |\nabla u|^2 dx dt \lesssim E(S) + \int_S^T \int_{\partial\Omega_N} (m \cdot \nu) \left( u'(t)^2 + \int_0^t (u'(t-s))^2 d|\mu|(s) \right) d\sigma dt$$

and, since  $|\mu| \leq \lambda$ , Lemma 2 gives us the desired result:

$$c(m) \int_S^T \int_{\Omega} (u')^2 + |\nabla u|^2 dx dt \lesssim E(S).$$

□

To conclude we need to absorb the two last integral terms for which we use the following result.

**Lemma 5** • For any solution  $u$  and any  $S < T$ ,

$$\int_S^T \int_{\partial\Omega_N} m \cdot \nu \int_0^t \left( \int_0^s (u'(t-r, x))^2 dr \right) d\lambda(s) d\sigma dt \lesssim \int_S^T \int_{\partial\Omega_N} m \cdot \nu \int_0^t (u'(t-s, x))^2 d\lambda(s) d\sigma dt.$$

• For any solution  $u$  and any  $S < T$ ,

$$\begin{aligned} \int_S^T \int_{\partial\Omega_N} m \cdot \nu \int_t^{+\infty} \left( \int_0^s (u'(s-r, x))^2 dr \right) d\lambda(s) d\sigma dt &\lesssim \int_S^T \int_{\partial\Omega_N} m \cdot \nu \int_0^t (u'(t-s, x))^2 d\lambda(s) d\sigma dt \\ &+ \int_S^{+\infty} \int_{\partial\Omega_N} m \cdot \nu u'^2 d\sigma dt. \end{aligned}$$

**Proof**

• We start from the left hand side term. We fix  $x \in \partial\Omega_N, t \in [S, T]$  and we use Fubini theorem to estimate integrals with respect to time:

$$\begin{aligned} \int_0^t \left( \int_0^s (u'(t-r, x))^2 dr \right) d\lambda(s) &= \int_0^t (u'(t-r, x))^2 \lambda([r, t]) dr \\ &\leq \int_0^t (u'(t-r, x))^2 \lambda([r, +\infty)) dr \\ &\leq \alpha^{-1} \int_0^t (u'(t-r, x))^2 d\lambda(r) \end{aligned}$$

which gives the required result after an integration with respect to  $t$  and  $x$ .

- As above, fixing  $x \in \partial\Omega_N$ , we obtain

$$\begin{aligned} \int_S^T \int_t^{+\infty} \left( \int_0^s (u'(s-r, x))^2 dr \right) d\lambda(s) dt &= \int_S^T \int_0^{+\infty} (u'(r, x))^2 \lambda([\max(r, t), +\infty]) dr dt \\ &= \int_S^T \int_0^t (u'(t-r, x))^2 dr \lambda([t, +\infty)) dt + \int_S^T \left( \int_t^{+\infty} (u'(r, x))^2 \lambda([r, +\infty)) dr \right) dt. \end{aligned}$$

Since for all  $r \leq t$ ,  $\lambda([t, +\infty)) \leq \lambda([r, +\infty))$ , we first have

$$\begin{aligned} \int_S^T \int_0^t (u'(t-r, x))^2 dr \lambda([t, +\infty)) dt &\leq \int_S^T \int_0^t (u'(t-r, x))^2 \lambda([r, +\infty)) dr dt \\ &\leq \alpha^{-1} \int_S^T \int_0^t (u'(t-r, x))^2 d\lambda(r) dt. \end{aligned}$$

On the other hand, Fubini theorem gives us

$$\int_S^T \left( \int_t^{+\infty} (u'(r, x))^2 \lambda([r, +\infty)) dr \right) dt = \int_S^{+\infty} u'(r)^2 \lambda([r, +\infty)) (\min(T, r) - S) dr.$$

We now note that

$$r\lambda([r, +\infty)) \leq \int_r^{+\infty} s d\lambda(s) \leq \int_0^{+\infty} s d\lambda(s)$$

and

$$\int_0^{+\infty} s d\lambda(s) = \int_0^{+\infty} \lambda([t, +\infty)) dt \leq \alpha^{-1} \lambda(\mathbb{R}^+).$$

We consequently obtain

$$\int_S^T \left( \int_t^{+\infty} (u'(r, x))^2 \lambda([r, +\infty)) dr \right) dt \lesssim \int_S^{+\infty} u'^2,$$

which give the required result after an integration over  $\partial\Omega_N$ .

□

Up to now, we have proven that

$$\int_S^T E(t) dt \lesssim E(S) + \int_S^T \int_{\partial\Omega_N} m \cdot \nu \int_0^t (u'(t-s, x))^2 d\lambda(s) d\sigma dt + \int_S^{+\infty} \int_{\partial\Omega_N} m \cdot \nu u'^2 d\sigma dt.$$

Lemma 2 allows us to conclude since it gives

$$\int_S^{+\infty} \int_{\partial\Omega_N} m \cdot \nu u'^2 d\sigma dt \lesssim E(S)$$

and

$$\int_S^T \int_{\partial\Omega_N} m \cdot \nu \int_0^t (u'(t-s, x))^2 d\lambda(s) d\sigma dt \lesssim E(S).$$

□

**Remark 4** In the case of some compactly supported measure  $\mu$ , one can also obtain exponential decay result for the following problem

$$\begin{cases} u'' - \Delta u = 0 & \text{in } \mathbb{R}_+^* \times \mathbb{R}_+^*, \\ u = 0 & \text{on } \mathbb{R}_+^* \times \partial\Omega_D, \\ \partial_\nu u + \mu_0 u'(t) + \int_0^t u'(t-s) d\mu(s) = 0 & \text{on } \mathbb{R}_+^* \times \partial\Omega_N, \\ u(0) = u_0 & \text{in } \Omega, \\ u'(0) = u_1 & \text{in } \Omega, \end{cases}$$

as it was done in [11] using the work of Lasiecka-Triggiani-Yao [7] and since the system is time invariant for  $t \gg 1$ .

Moreover, a careful attention shows that our proof allows us to obtain decay for this system without assumption on the support of  $\mu$  provided that

$$\inf_{\partial\Omega_N} m \cdot \nu > 0.$$

## 4 Examples

We start with two general results and then particularize them to recover results from the literature.

**Example 2** If  $\mu$  is some borelian measure such that

$$|\mu|(\mathbb{R}^+) < \mu_0 \text{ and } \int_0^{+\infty} e^{\beta s} d|\mu|(s) < +\infty$$

for some  $\beta > 0$ , then  $\mu$  fulfils the assumption (7) for an appropriate  $\alpha$ . Indeed for any  $0 \leq \alpha \leq \beta$ , the expression

$$\int_0^{+\infty} e^{\alpha s} d|\mu|(s)$$

is finite and by the dominated convergence Theorem of Lebesgue we have

$$\int_0^{+\infty} e^{\alpha s} d|\mu|(s) \rightarrow |\mu|(\mathbb{R}^+) \text{ as } \alpha \rightarrow 0.$$

Consequently by the assumption  $|\mu|(\mathbb{R}^+) < \mu_0$ , we get (7) for  $\alpha$  small enough.

**Example 3** One can choose

$$\mu = \sum_{i=1}^{\infty} \mu_i \delta_{\tau_i},$$

where  $(\tau_i)_{i=1}^{\infty}$ ,  $(\mu_i)_{i=1}^{\infty}$  are some families such that  $\tau_i > 0$  and are two by two disjoint, and

$$\sum_{i=1}^{\infty} |\mu_i| e^{\alpha \tau_i} < \mu_0$$

for some  $\alpha > 0$ .

**Example 4** If we choose  $d\mu(s) = k(s)ds$  where  $k$  is a kernel satisfying

$$\int_0^{+\infty} |k(s)|ds < \mu_0 \text{ and } \int_0^{+\infty} |k(s)|e^{\beta s}ds < \infty$$

for some  $\beta > 0$ . Then as a consequence of Example 2, we get an exponential decay rate for the system (S) under the (very weak) condition above, in particular we do not need any differentiability assumptions on  $k$ , nor uniform exponential decay of  $k$  at infinity as in [1, 3, 5, 9].

**Example 5** Choose

$$d\mu(s) = k(s)\chi_{[\tau_1, \tau_2]}(s)ds,$$

where  $k$  is an integrable function in  $[\tau_1, \tau_2]$  such that

$$\int_{\tau_1}^{\tau_2} |k(s)|ds < \mu_0,$$

then we get an exponential decay for the system (S) as a consequence of Example 2 because the second assumption trivially holds. In that case we extend the results of [11] to a larger class of kernels  $k$ , for instance in the class of bounded variations functions.

**Example 6** Take

$$\mu(s) = \mu_1 \delta_\tau(s),$$

where  $\mu_1$  is a constant and  $\tau > 0$  represents the delay satisfying

$$|\mu_1| < \mu_0,$$

then we recover the decay results from [10, 12].

## References

- [1] AASIL, M.; CAVALCANTI, M.M.; SORIANO, J.A., 2000, Asymptotic stability and energy decay rates for solutions of the wave equation with memory in a star-shaped domain. *SIAM J. Control. Optim.*, **38**, no 5, 1581-1602.
- [2] BEY, R., LOHÉAC, J.-P., MOUSSAOUI, M., 1999, Singularities of the solution of a mixed problem for a general second order elliptic equation and boundary stabilization of the wave equation. *J. Math. pures et appl.*, **78**, 1043-1067.
- [3] CAVALCANTI, M. M.; GUESMIA, A., 2005, General decay rates of solutions to a nonlinear wave equation with boundary condition of memory type. *Differential Integral Equations*, **18**, 583-600.
- [4] CORNILLEAU, P.; LOHÉAC, J.-P.; OSSES, A., Nonlinear Neumann boundary stabilization of the wave equation using rotated multipliers. *Preprint*. To appear in Journal of Dynamical and Control Systems, 2009.
- [5] GUESMIA, A. 1999, Stabilisation de l'équation des ondes avec conditions aux limites de type mémoire. *Afrika Math.*, **10**, 14-25.
- [6] KOMORNIK, V., 1994, Exact controllability and stabilization ; the multiplier method. *Masson-John Wiley, Paris*.
- [7] LASIECKA, I.; TRIGGIANI, R.; YAO, P.F., 1999, Inverse/observability estimates for second order hyperbolic equations with variable coefficients. *J. Math. Anal. Appl.*, **235**, 13-57.
- [8] NEČAS, J., 1967, Les méthodes directes en théorie des équations elliptiques. *Masson, Paris*.
- [9] NICAISE, S.; PIGNOTTI, C., 2006, Stabilization of the wave equation with variable coefficients and boundary condition of memory type, *Asymptotic Analysis*, **50**, p. 31-67.
- [10] NICAISE, S.; PIGNOTTI, C., 2006, Stability and unstability results of the wave equation with a delay term in the boundary or internal feedbacks. *SIAM J. Control. Optim.*, **45**, no 5, 1561-1585.
- [11] NICAISE, S.; PIGNOTTI, C., 2008, Stabilization of the wave equation with boundary or distributed delay. *Diff. Int. Equ.*, **21**, no 9-10, 935-958.
- [12] NICAISE, S.; VALEIN, J., 2007, Stabilization of the wave equation on 1-d networks with a delay term in the nodal feedbacks, *NHM*, **2**, p. 425-479.
- [13] PROPST, G.; PRÜSS, J., 1996, On wave equations with boundary dissipation of memory type. *J. Integral Equations Appl.*, **8**, no. 1, 99-123.
- [14] OSSES, A., 1998, Une nouvelle famille de multiplicateurs et applications à la contrôlabilité exacte de l'équation des ondes. *C. R. Acad. Sci. Paris*, **326**, Série I, 1099-1104.
- [15] OSSES, A., 2001, A rotated multiplier applied to the controllability of waves, elasticity and tangential Stokes control. *SIAM J. Control Optim.*, **40**, no 3, 777-800.
- [16] RELICH F., 1940, Darstellung der Eigenwerte von  $\Delta u + \lambda u$  durch ein Randintegral. *Math. Zeitschrift*, **46**, 635-636.