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# Derivation and Analysis of Piecewise Constant Conservative Approximation for Anisotropic Diffusion Problems 

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#### Abstract

A variational approach to derive a piecewise constant conservative approximation of anisotropic diffusion equations is presented. A priori error estimates are derived assuming usual mesh regularity constraints and a posteriori error indicator is proposed and analyzed for the model problem.


Key words:
1991 MSC: 65N30 Finite elements, conservative approximation, a posteriori error estimate

## 1 Introduction

Various phenomena in scientific fields such as geoscience, oil reservoir simulation, hydrogeology, biology ..., are generally modeled by anisotropic diffusion equations. The usual discretization schemes of this equations are finite difference, finite element or finite volume methods. The last are piecewise constant conservative approximation and are actually very popular in oil engineering, the reason probably being that complex coupled physical phenomena may be discretized on the same grids (see for instance [9] and references therein). But the well known five point on rectangles and four point schemes on triangles are not easily adapted to heterogeneous anisotropic diffusion operators, and so an enlarged stencil scheme which handles anisotropy on meshes satisfying an orthogonality property was proposed and analyzed in $[2,6,7]$. Let us recall that a huge literature exists in engineering study setting. However, even though these schemes perform well in the number of cases, their convergence

[^0]analysis often seems out of reach, unless some additional geometrical conditions are imposed. Moreover, actually in several applications the discretization meshes are imposed by engineering and computing considerations, therefore we have to deal with unstructured meshes.
A motivation for this work was to construct such a piecewise constant approximation for anisotropic diffusion problems which could satisfy the two assumptions : First, the resulting formulation is well-defined on general unstructured meshes, assuming usual finite element mesh regularity constraints. Secondly, the given scheme leads to standard algebraic system, for which we can use the existing efficient numerical solvers. This last point is of major importance in the coupling of physical models, from the implementation point of view and good adaptivity properties.

An outline of the paper is as follows. Among all the developments we briefly introduce the functional framework and some usual notations. In second section, we introduce the numerical scheme for the anisotropic diffusion problem and main approximation analysis results are given. A focus is made on the treatment of an additional reaction term. Finally, in section three we propose and analyze an a posteriori error indicator for the diffusion model problem.

## Functional Framework and some notations

Let $\omega$ be a bounded polygonal domain of $\mathbb{R}^{2}$. We denote by $H^{s}(\omega)$ the usual Sobolev space $W^{s, 2}(\omega)$ (see e.g [1]), endowed with the norm $\|.\|_{s, \omega}$ and $H_{0}^{s}(\omega)$ is the closure of $\mathcal{D}(\omega)$ in $H^{s}(\omega)$. For the semi-norm, we use the notation $|\cdot|_{s, \omega}$. We introduce the set $H$ (div, $\omega$ ) of vector fields $p \in\left(L^{2}(\omega)\right)^{d}$ and $\operatorname{div} p \in L^{2}(\omega)$. Equipped with the norm $\|\cdot\|_{H(\operatorname{div}, \omega)}^{2}=\|\cdot\|_{0, \omega}^{2}+\|\operatorname{div} \cdot\|_{0, \omega}^{2}$, $H(\operatorname{div}, \omega)$ is a Hilbert space. For any integer $k, P_{k}(\omega)$ is the set of polynomials of degree less than or equal to $k$.

## 2 Construction and analysis of the numerical scheme

Let $\Omega$ denote a bounded polygonal domain of $R^{2}$. We consider the anisotropic diffusion problem :

$$
\left\{\begin{align*}
-\operatorname{div}(K \nabla u)=f & \text { over } \Omega,  \tag{2.1}\\
u=0 & \text { on } \Gamma,
\end{align*}\right.
$$

with symmetric definite positive tensor K , assumed piecewise constant for simplicity, and $f \in L^{2}(\Omega)$. Let $\left(\mathcal{T}_{h}\right)$ be a family of triangulations of $\Omega$, by triangles, regular in the usual finite element sense [5]. For all $T$ in $\mathcal{T}_{h}$, there exist reals $d_{T, e}$ such that the bilinear form $a_{T}:\left(H^{1}(T)\right)^{2} \longrightarrow \mathbb{R}$ defined by

$$
\forall p, q \in\left(H^{1}(T)\right)^{2}, \quad a_{T}(p, q)=\sum_{e \in \partial T} d_{T, e}\left(\int_{e} p \cdot n_{T} d \gamma\right)\left(\int_{e} q \cdot n_{T} d \gamma\right) .
$$

verifies

$$
\forall p, q \in\left(P_{0}(T)\right)^{2}, \quad a_{T}(p, q)=\int_{T} K^{-1} p \cdot q d x
$$

where $n_{T}$ is the unit normal outward to $T$.

Remark 2.1. We could give explicit expressions of the parameters $d_{T, e}$. Let $T$ be a triangle with vertices $a, b, c$. The edge $a b$ is denoted $e$ and $\theta_{e}$ the opposite angle to $e$ then,

$$
d_{T, e}=\frac{\sqrt{\operatorname{det}(K)}<K^{-1} \overrightarrow{a c}, \overrightarrow{b c}>}{4 \operatorname{meas}(T)}
$$

Clearly if $K=\alpha I d$ then $d_{T, e}=\frac{\alpha}{2} \operatorname{cotan}\left(\theta_{e}\right)$.
In the sequel, we denote by $R_{T} x$ the unique element of $\left(P_{0}(T)\right)^{2}$ checking

$$
a_{T}\left(R_{T} x-x, q\right)=0, \quad \forall q \in\left(P_{0}(T)\right)^{2}
$$

Let $V_{h}$ be the nonconforming finite element space defined by

$$
\begin{aligned}
V_{h}= & \left\{\lambda_{h} \in L^{2}(\Omega) ; \lambda_{\left.h\right|_{T}} \in P_{1}(T),\right. \\
& \left.\forall T \in \mathcal{T}_{h} ; \forall e \text { interior edge, } \int_{e}\left[\lambda_{h}\right]_{e} d \gamma=0 \text { and } \forall e \subset \partial \Omega, \int_{e} \lambda_{h} d \gamma=0\right\}
\end{aligned}
$$

where $\left[\lambda_{h}\right]_{e}$ denotes the jump of the function $\lambda_{h}$ across the edge $e$.
The non-standard finite element approximation we propose for the model problem is the following :

$$
\left\{\begin{array}{l}
\text { Find } \lambda_{h} \in V_{h} \text { such that }  \tag{2.2}\\
\forall \mu_{h} \in V_{h}, \quad \sum_{T \in \mathcal{T}_{h}} \int_{T} K \nabla \lambda_{h} . \nabla \mu_{h}=\sum_{T \in \mathcal{T}_{h}} \int_{T} f_{T}\left(\mu_{h}+\frac{1}{2}\left(x_{G}-R_{T} x\right) . \nabla \mu_{h}\right) d x
\end{array}\right.
$$

where $\forall T \in \mathcal{T}_{h}, \quad f_{T}=\frac{1}{\operatorname{meas}(T)} \int_{T} f d x$ and $x_{G}$ is the barycenter of $T$.

First of all, by adapting standard arguments used in the analysis of nonconforming finite element approximation of elliptic problems [5], we can easily prove that the discrete problem has a unique solution; moreover, if the weak solution $u$ of the model problem belongs to the Sobolev space $H^{1+s}(\Omega)$ with $0<s \leq 1$, then:

$$
\begin{equation*}
\left(\sum_{T \in \mathcal{T}_{h}}\left|u-\lambda_{h}\right|_{1, T}^{2}\right)^{\frac{1}{2}} \leq C\left(h^{s}|u|_{1+s, \Omega}+\left(\sum_{T} h_{T}^{2}\|f\|_{0, T}^{2}+h_{T}^{2} \operatorname{dist}^{2}\left(x_{G}, R_{T} x\right)\right)^{\frac{1}{2}}\right) \tag{2.3}
\end{equation*}
$$

where $h_{T}$ is the diameter of the triangle $T$.

Let us set

$$
\forall T \in \mathcal{T}_{h}, \quad p_{h}=K . \nabla \lambda_{h}-f_{T} \frac{\left(x-R_{T} x\right)}{2} \quad \text { on } T .
$$

The key point of the construction of the scheme is that $p_{h}$ is an admissible field in the following sense :

Lemma 2.1 The vector field $p_{h}$ satisfies :

$$
p_{h} \in H(\operatorname{div}, \Omega) \quad \text { and } \quad \forall T \in \mathcal{T}_{h}-\operatorname{div} p_{h}=f_{T} \quad \text { on } T
$$

Proof. It is obvious that $-\operatorname{div} p_{h}=f_{T}, \forall T \in \mathcal{T}_{h}$. Moreover if $e$ is an interior edge of $\mathcal{T}_{h}$, $e=\partial T_{1} \cap \partial T_{2}$, with $T_{1}, T_{2} \in \mathcal{T}_{h}$ and $v_{h}^{e} \in V_{h}$ the associated basis function, i.e.,

$$
\text { for any edge } \sigma \text { of } \mathcal{T}_{h}, \quad \int_{\sigma} v_{h}^{e} d \gamma=\delta_{\sigma}^{e} \text {, the Kronecker delta. }
$$

Let us denote by $\left[p_{h} . n\right]_{e}$ the jump of the flux across the edge $e$. We have thus

$$
\begin{aligned}
{\left[p_{h} \cdot n\right]_{e}=\int_{e}\left[p_{h} \cdot n\right]_{e} v_{h}^{e} d \sigma } & =\sum_{i=1}^{2} \int_{T_{i}}\left(p_{h} \cdot \nabla v_{h}^{e}+v_{h}^{e} \operatorname{div} p_{h}\right) d x \\
& =\sum_{i=1}^{2} \int_{T_{i}}\left(\left(K . \nabla \lambda_{h}-f_{T} \frac{\left(x_{G}-R_{T} x\right)}{2}\right) \cdot \nabla v_{h}^{e}-f_{T} v_{h}^{e}\right) d x=0
\end{aligned}
$$

which yields $p_{h} \in H(\operatorname{div}, \Omega)$.

In order to define the numerical scheme, we need to introduce some notations. Let $\lambda_{h}$ be the solution of the discrete problem (2.2); For any $T \in \mathcal{T}_{h}$ and $\sigma$ edge of $T$, we set:

$$
F_{\sigma, T}=\int_{\sigma} p_{h} \cdot n_{T} d \gamma ; \quad u_{\sigma, T}=\frac{1}{\operatorname{meas}(\sigma)} \int_{\sigma} \lambda_{h} d \gamma
$$

and

$$
u_{T}=\frac{1}{\operatorname{meas}(T)} \int_{T}\left(\lambda_{h}+\frac{1}{2} \nabla \lambda_{h} \cdot\left(x_{G}-R_{T} x\right)\right) d x+\frac{\rho_{T, h}^{2}}{4} f_{T},
$$

where

$$
\rho_{T, h}^{2}=\frac{1}{\operatorname{meas}(T)} a_{T}\left(x-R_{T} x, x-R_{T} x\right) .
$$

Lemma 2.2 With the notations given above, one has the following scheme $\forall T, T_{1}, T_{2} \in \mathcal{T}_{h}$,

$$
\left\{\begin{align*}
-\sum_{\sigma \in \partial T} F_{\sigma, T} & =\operatorname{meas}(T) f_{T}  \tag{2.4}\\
F_{\sigma, T_{1}}+F_{\sigma, T_{2}} & =0, \quad \forall \sigma \in \partial T_{1} \cap \partial T_{2} \\
d_{T, \sigma} F_{\sigma, T}+u_{T} & =u_{\sigma, T}, \quad \forall \sigma \in \partial T \\
u_{\sigma, T} & =0, \quad i f \sigma \in \partial T \cap \partial \Omega
\end{align*}\right.
$$

Proof . First, we have

$$
-\sum_{\sigma \in \partial T} F_{\sigma, T}=\operatorname{meas}(T) f_{T}=\int_{T} f d x, \quad \forall T \in \mathcal{T}_{h}
$$

Indeed, we have :

$$
-\int_{T} \operatorname{div} p_{h} d x=-\sum_{e \in \partial T} \int_{\sigma} p_{h} \cdot n_{T} d x=\int_{T} f_{T} d x=\operatorname{meas}(T) f_{T}
$$

And for any interior edge $\sigma \in \partial T_{1} \cap \partial T_{2}, \quad F_{\sigma, T_{1}}+F_{\sigma, T_{2}}=0$ is obvious since $\left.p_{h} \in H(\operatorname{div}, \Omega)\right)$. Let $q_{h} \in R T_{0}(T)=\left(P_{0}(T)\right)^{2}+x P_{0}(T)$ such that $\left.q_{h} \cdot n_{T}\right|_{\sigma}=1$ and $q_{h} .\left.n_{T}\right|_{e}=0$, Vedge $e \neq \sigma$. By one hand, we have,

$$
B:=a_{T}\left(p_{h}, q_{h}\right)-\int_{\partial T} \lambda_{h} q_{h} \cdot n_{T}=d_{T, \sigma} \operatorname{meas}(\sigma) F_{\sigma, T}-u_{\sigma, T} \cdot \operatorname{meas}(\sigma)
$$

and on the other hand, if we set

$$
A=a_{T}\left(\frac{\operatorname{div} p_{h}}{2}\left(x-R_{T} x\right), q_{h}\right),
$$

and

$$
\bar{\lambda}_{h}=\frac{1}{\operatorname{meas}(T)} \int_{T} \lambda_{h} d x
$$

we get

$$
\begin{aligned}
& a_{T}\left(p_{h}, q_{h}\right)-\int_{\partial T} \lambda_{h} q_{h} \cdot n_{T}=a_{T}\left(p_{h}, q_{h}\right)-\int_{T} \nabla \lambda_{h} \cdot q_{h}-\int_{T} \overline{\lambda_{h}} \operatorname{div} q_{h} d x \\
& =a_{T}\left(p_{h}-\frac{\operatorname{div} p_{h}}{2}\left(x-R_{T} x\right), q_{h}\right)-\int_{T} \nabla \lambda_{h} \cdot q_{h}+A-\int_{T} \overline{\lambda_{h}} \operatorname{div} q_{h} d x \\
& =a_{T}\left(\nabla \lambda_{h}, q_{h}\right)-\int_{T} \nabla \lambda_{h} \cdot q_{h}+A-\int_{T} \overline{\lambda_{h}} \operatorname{div} q_{h} d x \\
& =a_{T}\left(\nabla \lambda_{h}, q_{h}-\frac{\operatorname{div} q_{h}}{2}\left(x-R_{T} x\right)\right)-\int_{T} \nabla \lambda_{h} \cdot q_{h}+A-\int_{T} \overline{\lambda_{h}} \operatorname{div} q_{h} d x \\
& \left.=\int_{T}\left(q_{h}-\frac{\operatorname{div} q_{h}}{2} x-R_{T} x\right)\right) \cdot \nabla \lambda_{h}-\int_{T} \nabla \lambda_{h} \cdot q_{h}+A-\int_{T} \overline{\lambda_{h}} \operatorname{div} q_{h} d x \\
& =-\int_{T}\left(\overline{\lambda_{h}}+\frac{1}{2} \nabla \lambda_{h}\left(x_{G}-R_{T} x\right)\right) \operatorname{div} q_{h}+A .
\end{aligned}
$$

However, since $\int_{T} \operatorname{div} q_{h}=\int_{\partial T} q_{h} \cdot n_{T} d \gamma=\operatorname{meas}(\sigma)$, we obtain

$$
B=-\operatorname{meas}(\sigma)\left(\overline{\lambda_{h}}+\frac{1}{2} \nabla \lambda_{h}\left(x_{G}-R_{T} x\right)\right)+A
$$

But we have also,

$$
\begin{aligned}
A & =a_{T}\left(\frac{\operatorname{div} p_{h}}{2}\left(x-R_{T} x\right), q_{h}\right) \\
& =a_{T}\left(\frac{\operatorname{div} p_{h}}{2}\left(x-R_{T} x\right), q_{h}-\frac{\operatorname{div} q_{h}}{2}\left(x-R_{T} x\right)\right) \\
& -\frac{\operatorname{div} q_{h}}{4} a_{T}\left(x-R_{T} x, x-R_{T} x\right) \cdot f_{T} \\
& =-\left(\int_{T} \operatorname{div} q_{h} d x\right) \cdot \frac{\rho_{T, h}}{4} \cdot f_{T}=-\operatorname{meas}(\sigma) \frac{\rho_{T, h}}{4} \cdot f_{T} .
\end{aligned}
$$

where

$$
\rho_{T, h}=\frac{1}{\operatorname{meas}(T)} a_{T}\left(x-R_{T} x, x-R_{T} x\right)
$$

which implies

$$
B=\left(\overline{\lambda_{h}}+\frac{1}{2} \nabla \lambda_{h}\left(x_{G}-R_{T} x\right)+\frac{\rho_{T, h}}{4} \cdot f_{T}\right)
$$

and thus scheme (2.4).
Using once more Lemma 2.1, we can derive the following a priori error estimate,
Lemma 2.3 If the weak solution $u$ of model problem (2.1) belongs to $H^{1+s}(\Omega), 0<s \leq 1$, one has:

$$
\begin{equation*}
\left(\sum_{T \in \mathcal{T}_{h}}\left\|u-u_{T}\right\|_{0, T}^{2}\right)^{\frac{1}{2}} \leq C\left(h^{s}|u|_{1+s, \Omega}+\left(\sum_{T} h_{T}^{2}\|f\|_{0, T}^{2}+h_{T}^{2} \operatorname{dist}^{2}\left(x_{G}, R_{T} x\right)\right)^{\frac{1}{2}}\right) \tag{2.5}
\end{equation*}
$$

### 2.1 A focus on the treatment of an additional reaction term

Let us consider the problem of diffusion-reaction equations:

$$
\left\{\begin{align*}
-\operatorname{div}(\nabla u)+c u=f & \text { over } \Omega  \tag{2.6}\\
u=0 & \text { on } \Gamma .
\end{align*}\right.
$$

where $f \in L^{2}(\Omega), c \in L^{\infty}(\Omega)$ with $c \geq 0$, and the following associated discrete problem :

$$
\left\{\begin{align*}
\text { Find } \lambda_{h} & \in V_{h} \text { such that }  \tag{2.7}\\
\forall \mu_{h} \in V_{h}, & \sum_{T \in \mathcal{T}_{h}} \int_{T} \nabla \lambda_{h} \cdot \nabla \mu_{h} d x \\
& +\sum_{T \in \mathcal{T}_{h}} \alpha_{T} c_{T} \int_{T}\left(\bar{\lambda}_{h}+\frac{1}{2}\left(x_{G}-R_{T} x\right) \cdot \nabla \lambda_{h}\right)\left(\bar{\mu}_{h}+\frac{1}{2}\left(x_{G}-R_{T} x\right) \cdot \nabla \mu_{h}\right) d x \\
& =\sum_{T \in \mathcal{T}_{h}} \alpha_{T} \int_{T} f_{T}\left(\mu_{h}+\frac{1}{2}\left(x_{G}-R_{T} x\right) \cdot \nabla \mu_{h}\right) d x
\end{align*}\right.
$$

where for all $T \in \mathcal{T}_{h}$,

$$
c_{T}=\frac{1}{\operatorname{meas}(T)} \int_{T} c d x, f_{T}=\frac{1}{\operatorname{meas}(T)} \int_{T} f d x, \alpha_{T}=\frac{4}{4+c_{T} \rho_{T}^{2}},
$$

$\bar{\lambda}_{h}=\frac{1}{\operatorname{meas}(T)} \int_{T} \lambda_{h} d x \quad$ (and analogously for $\bar{\mu}_{h}$ ), and $x_{G}$ is the barycenter of $T$.
Using the same arguments as before, we have in this case :
Lemma 2.4 Let $\lambda_{h}$ be the solution of the discrete problem. We introduce

$$
\begin{gathered}
p_{h}=\nabla \lambda_{h}-\frac{2}{4+c_{T} \rho_{T}^{2}}\left(f_{T}+c_{T}\left(\bar{\lambda}_{h}+\nabla \lambda_{h}\left(x_{G}-R_{T} x\right)\right)\right)\left(x-R_{T} x\right) . \\
F_{\sigma, T}=\int_{\sigma} p_{h} \cdot n_{T} d \gamma, \quad u_{\sigma}=\frac{1}{\operatorname{meas}(\sigma)} \int_{\sigma} \lambda_{h} d \gamma
\end{gathered}
$$

and

$$
u_{T}=\frac{4}{4+c_{T} \rho_{T}^{2}}\left(\frac{1}{\operatorname{meas}(T)} \int_{T}\left(\lambda_{h}+\frac{1}{2} \nabla \lambda_{h} \cdot\left(x_{G}-R_{T} x\right)\right)+\frac{\rho_{T, h}^{2}}{4} f_{T}\right),
$$

where

$$
\rho_{T, h}^{2}=\frac{1}{\operatorname{meas}(T)} a_{T}\left(x-R_{T} x, x-R_{T} x\right) .
$$

Then we have

$$
\begin{aligned}
-\sum_{\sigma \in \partial T} F_{\sigma, T}+c_{T} u_{T} & =\operatorname{meas}(T) f_{T} \\
F_{\sigma, T_{1}}+F_{\sigma, T_{2}} & =0, \quad \forall \sigma \in \partial T_{1} \cap \partial T_{2} \\
d_{T, \sigma} F_{\sigma, T}+u_{T} & =u_{\sigma, T}, \quad \forall \sigma \in \partial T \\
u_{\sigma, T} & =0, \quad i f \sigma \in \partial T \cap \partial \Omega .
\end{aligned}
$$

## 3 A posteriori error estimator for the diffusion model problem

Usually, error estimators for adaptive refinement require exact discrete solutions (see [10] and references therein), but in practical cases the exact solution is not available and so we are in the presence of solvers error. In this subsection, we introduce a posteriori error estimator for solutions obtained by black-box solver, in this case we are in the presence of many source of errors : approximation, error solvers, post processing error ...etc. We indicate the a posteriori error estimator for the diffusion model equation:

$$
\left\{\begin{align*}
-\Delta u=f & \text { over } \Omega  \tag{3.1}\\
u=0 & \text { on } \Gamma
\end{align*}\right.
$$

The solution is assumed to be obtained by any existing solver. The given estimator is valid also for equilibrium and mixed finite element approximations with or without numerical
integration.
Let $\mathcal{T}_{h}$ be a regular triangulation of $\Omega$ by triangles, $E$ is the set of all edges and $E_{I}$ the set of all interior edges. Given $T \in \mathcal{T}_{h}, \Delta(T)$ is the union of all elements of $\mathcal{T}_{h}$ sharing a vertex with $T, \omega_{T}$ is the union of all elements of $\mathcal{T}_{h}$ sharing an edge with $T$ and $E_{T}$ the set of all edges of $T$. We consider the finite dimensional space

$$
\begin{gather*}
V_{h}=\left\{v_{h} \in H^{1}(\Omega), \forall T \in \mathcal{T}_{h} \quad v_{h \mid T} \in \mathrm{P}_{1}(T)\right\}  \tag{3.2}\\
E_{h}=\left\{p_{h} \in\left(L^{2}(\Omega)\right)^{2}, \forall T \in \mathcal{T}_{h} \quad p_{h \mid T} \in \mathrm{RT}_{0}(T)=\left(P_{0}(T)\right)^{2}+x P_{0}(T)\right\} \tag{3.3}
\end{gather*}
$$

and

$$
\begin{equation*}
M_{h}=\left\{v_{h} \in L^{2}(\Omega), \quad \forall T \in \mathcal{T}_{h} \quad v_{h \mid T} \in \mathrm{P}_{0}(T)\right\} \tag{3.4}
\end{equation*}
$$

Let $p_{h} \in E_{h}$ and $u_{h} \in M_{h}$, for all $T \in \mathcal{T}_{h}$ we set

$$
\begin{gather*}
\varepsilon_{1, T}\left(p_{h}\right)=\sup _{v_{h} \in V_{h}(T)} \frac{\int_{\Delta(T)} p_{h} \nabla v_{h} d x-\int_{\Delta(T)} v_{h} f d x}{\left|v_{h}\right|_{1, \Delta(T)}},  \tag{3.5a}\\
\varepsilon_{2, T}\left(p_{h}\right)=\sup _{\phi_{h} \in V_{h}(T)} \frac{\int_{\Delta(T)} p_{h} \operatorname{curl} \phi_{h} d x}{\left|\phi_{h}\right|_{1, \Delta(T)}},  \tag{3.5b}\\
\varepsilon_{3, T}\left(p_{h}, u_{h}\right)=\sup _{q_{h} \in E_{h}(T)} \frac{\int_{\omega_{T}}\left(p_{h} q_{h}+u_{h} \operatorname{div} q_{h}\right) d x}{\left\|q_{h}\right\|_{H\left(\operatorname{div}, \omega_{T}\right)}}  \tag{3.5c}\\
\eta_{1, T}^{2}\left(p_{h}\right)=h_{T}^{2}\left\|f-f_{T}\right\|_{0, T}^{2}+\sum_{l \in E_{T}}\left(h_{l}\left\|\left[p_{h} \cdot t_{l}\right]_{l}\right\|_{0, l}^{2}\right) \tag{3.5d}
\end{gather*}
$$

and

$$
\begin{equation*}
\eta_{2, T}\left(p_{h}\right)=h_{T}\left\|p_{h}\right\|_{0, T} \tag{3.5e}
\end{equation*}
$$

where

$$
\begin{equation*}
V_{h}(T)=\left\{v_{h} \in H_{0}^{1}(\Delta(T)), \quad \forall T \in \Delta(T) \quad v_{h \mid T} \in P_{1}(T)\right\} \tag{3.6}
\end{equation*}
$$

and

$$
\begin{equation*}
E_{h}(T)=\left\{q_{h} \in E_{h} \cap H(\operatorname{div}, \Omega), \quad \forall T \notin \omega_{T} \quad q_{h \mid T}=0\right\} \tag{3.7}
\end{equation*}
$$

$h_{T}$ and $h_{l}$ are the diameters of $T$ and $l$ respectively. The outward normal to an edge $l$ of some $T \in \mathcal{T}_{h}$ is written as $n_{l}=\left(n_{1, l}, n_{2, l}\right)$ and we set $t_{l}=\left(n_{2, l},-n_{1, l}\right)$ for associated tangential direction. we denote by $\left[p_{h} . t_{l}\right]_{l}$ the jump of $p_{h} . t_{l}$ across the edge $l$.
In the sequel, $C, C_{1}, C_{2}$ are positive generic constants independent of $h$ (which may change from one line to other).

Remark 3.1. Let us notice that
(1) Since

$$
-\operatorname{div} p_{h}=\frac{1}{\operatorname{meas}(T)} \int_{T} f d x=f_{T} \text { on } T
$$

we have

$$
\varepsilon_{1, T}^{2}\left(p_{h}\right) \leq C \sum_{T \in \Delta(T)} h_{T}^{2}\left\|f-f_{T}\right\|_{0, T}^{2},
$$

which is higher order perturbation of the error.
(2) Let $u_{h}^{T}, \psi_{h}^{T} \in V_{h}(T)$ and $q_{h}^{T} \in E_{h}(T)$ be the unique solutions of the following respective problems
(P1) $\left\{\begin{array}{l}\text { Find } u_{h}^{T} \in V_{h}(T) \text { such that } \\ \forall v_{h} \in V_{h}(T), \int_{\Delta(T)} \nabla u_{h}^{T} \nabla v_{h} d x=\int_{\Delta(T)} p_{h} \nabla v_{h} d x-\int_{\Delta(T)} f v_{h} d x,\end{array}\right.$

$$
\left\{\begin{array}{l}
\text { Find } \psi_{h}^{T} \in V_{h}(T) \text { such that }  \tag{P2}\\
\forall \phi_{h} \in V_{h}(T), \int_{\Delta(T)} \operatorname{curl} \psi_{h}^{T} \operatorname{curl} \phi_{h} d x=\int_{\Delta(T)} p_{h} \operatorname{curl} \phi_{h} d x
\end{array}\right.
$$

$$
\left\{\begin{array}{l}
\text { Find } q_{h}^{T} \in E_{h}(T) \text { such that }  \tag{P3}\\
\forall s_{h} \in E_{h}(T), \quad \int_{\omega_{T}}\left(q_{h}^{T} s_{h}+\operatorname{div} q_{h}^{T} \operatorname{div} s_{h}\right) d x=\int_{\omega_{T}}\left(p_{h} s_{h}+u_{h} \operatorname{div} s_{h}\right) d x
\end{array}\right.
$$

It is easy to see that

$$
\left|u_{h}^{T}\right|_{1, T}=\varepsilon_{1, T}\left(p_{h}\right) \quad\left|\psi_{h}^{T}\right|_{1, T}=\varepsilon_{2, T}\left(p_{h}\right) \quad \text { and } \quad\left\|q_{h}^{T}\right\|_{H\left(\operatorname{div}, \omega_{T}\right)}=\varepsilon_{3, T}\left(p_{h}, u_{h}\right) .
$$

We have the following error estimates,
Theorem 3.1 Let $u \in H_{0}^{1}(\Omega)$ be the weak solution of the model problem (3.1), $p=\nabla u$, $p_{h} \in E_{h}$ and $u_{h} \in M_{h}$ the solution of the given numerical scheme. Then there exists a positive constant $C$ only depending on the minimum angle of $\mathcal{T}_{h}$ such that

$$
\begin{aligned}
\left\|u-u_{h}\right\|_{0, \Omega}+\left\|p-p_{h}\right\|_{0, \Omega} & \leq C\left\{\left(\sum_{T \in \mathcal{T}_{h}}\left(\eta_{1, T}^{2}\left(p_{h}\right)+\varepsilon_{1, T}^{2}\left(p_{h}\right)+\varepsilon_{2, T}^{2}\left(p_{h}\right)\right)\right)^{\frac{1}{2}}\right. \\
& \left.+\left(\sum_{T \in \mathcal{T}_{h}}\left(\eta_{2, T}^{2}\left(p_{h}\right)+\varepsilon_{3, T}^{2}\left(p_{h}, u_{h}\right)\right)\right)^{\frac{1}{2}}\right\} .
\end{aligned}
$$

Moreover, we have

$$
\begin{gathered}
\sum_{i=1}^{2} \eta_{i, T}\left(p_{h}\right)+\sum_{i=1}^{2} \varepsilon_{i, T}\left(p_{h}\right)+ \\
\varepsilon_{3, T}\left(p_{h}, u_{h}\right) \leq C_{1}\left(\left\|p-p_{h}\right\|_{0, \Delta(T)}+\left\|u-u_{h}\right\|_{0, \Delta(T)}\right) \\
+C_{2}\left(\sum_{T^{\prime} \in \omega_{T}} h_{T^{\prime}}^{2}\left\|f-f_{T^{\prime}}\right\|_{0, T^{\prime}}^{2}\right)^{\frac{1}{2}}
\end{gathered}
$$

Proof : First, Using Helmholtz-decomposition, we have $e_{h}=p-p_{h}=\nabla w+$ curl $\zeta$, with $w \in H_{0}^{1}(\Omega), \zeta \in H^{1}(\Omega)$ and $\int_{\Omega} \nabla w \operatorname{curl} \zeta d x=0$.
Let us remark that the orthogonality implies the following error decomposition :

$$
\begin{gather*}
\left\|e_{h}\right\|_{0, \Omega}^{2}=|w|_{1, \Omega}^{2}+\|\operatorname{curl} \zeta\|_{0, \Omega}^{2},  \tag{3.8}\\
|w|_{1, \Omega}^{2}=\int_{\Omega} e_{h} \nabla w d x \quad \text { and } \quad\|\operatorname{curl} \zeta\|_{0, \Omega}^{2}=\int_{\Omega} e_{h} \operatorname{curl} \zeta d x . \tag{3.9}
\end{gather*}
$$

Now, let $w^{I} \in V_{h}$ and $\zeta^{I} \in V_{h}$ be continuous approximations of $w$ and $\zeta$ respectively such that:

$$
\begin{gather*}
\forall T \in \mathcal{T}_{h}, \quad\left\|w-w^{I}\right\|_{0, T} \leq C \operatorname{meas}(T)^{\frac{1}{2}}|w|_{1, \Delta(T)}  \tag{3.10a}\\
\left|w^{I}\right|_{1, \Omega} \leq C|w|_{1, \Omega} \tag{3.10b}
\end{gather*}
$$

and

$$
\begin{equation*}
\forall l \in E_{T}, \quad\left\|w-w^{I}\right\|_{0, l} \leq C \operatorname{meas}(l)^{\frac{1}{2}}|w|_{1, \Delta(l)} \tag{3.10c}
\end{equation*}
$$

( and analogously for $\zeta$ ) where $\Delta(l)$ is the union of the elements $T$ sharing $l$. Moreover we assume that the interpolation preserves boundary conditions, that is, $w^{I} \in V_{h} \cap H_{0}^{1}(\Omega)$. It is well known that such approximations exist (see [5], [8]).
First, according to (3.9) we have by element-wise integration by parts, noting that $w-w^{I} \in H_{0}^{1}(\Omega)$,

$$
\begin{aligned}
& \|\nabla w\|_{0, \Omega}^{2}=\int_{\Omega} e_{h} \nabla\left(w-w^{I}\right) d x+\int_{\Omega} f w^{I} d x-\int_{\Omega} p_{h} \nabla w^{I} d x \\
= & \sum_{T \in \mathcal{T}} \int_{T}\left(f+\operatorname{div} p_{h}\right)\left(w^{I}-w\right) d x+\int_{\Omega} f w^{I} d x-\int_{\Omega} p_{h} \nabla w^{I} d x .
\end{aligned}
$$

From Cauchy's inequality and from (3.10a) and (3.10c),

$$
\sum_{T \in \mathcal{T}_{h}} \int_{T}\left(f+\operatorname{div} p_{h}\right)\left(w^{I}-w\right) d x \leq\left(\sum_{T \in \mathcal{T}_{h}} h_{T}^{2}\left\|f+\operatorname{div} p_{h}\right\|_{0, T}^{2}\right)^{\frac{1}{2}}|w|_{1, \Omega}
$$

Using (3.10b), we obtain

$$
\int_{\Omega} f w^{I} d x-\int_{\Omega} p_{h} \nabla w^{I} d x \leq C\left(\sum_{T \in \mathcal{T}_{h}} \varepsilon_{1, T}^{2}\left(p_{h}\right)\right)^{\frac{1}{2}}\left|w^{I}\right|_{1, \Omega} \leq C\left(\sum_{T \in \mathcal{T}_{h}} \varepsilon_{1, T}^{2}\left(p_{h}\right)\right)^{\frac{1}{2}}|w|_{1, \Omega}
$$

Then we have

$$
\begin{equation*}
\|\nabla w\|_{0, \Omega} \leq C\left(\sum_{T \in \mathcal{I}_{h}} h_{T}^{2}\left\|f+\operatorname{div} p_{h}\right\|_{0, T}^{2}+\sum_{T \in \mathcal{I}_{h}} \varepsilon_{1, T}^{2}\left(p_{h}\right)\right)^{\frac{1}{2}} \tag{3.11}
\end{equation*}
$$

Arguing as above, since $p=\nabla u$ and $p_{h} \in E_{h}$, by element-wise integration by parts we have

$$
\int_{\Omega} e_{h} \operatorname{curl}\left(\zeta-\zeta^{I}\right) d x=\sum_{l \in E} \int_{l}\left(\zeta-\zeta^{I}\right)\left[p_{h} . t_{l}\right]_{l} d \sigma
$$

By Cauchy's inequality (3.10c) we obtain

$$
\begin{equation*}
\int_{\Omega} e_{h} \operatorname{curl}\left(\zeta-\zeta^{I}\right) d x \leq C\left(\sum_{l \in E} h_{l}\left\|\left[p_{h} \cdot t_{l}\right]_{l}\right\|_{0, l}^{2}\right)^{\frac{1}{2}}\|\nabla \zeta\|_{0, \Omega} \tag{3.12}
\end{equation*}
$$

Finally, since $\|\operatorname{curl} \zeta\|_{0, \Omega}=\|\nabla \zeta\|_{0, \Omega}$,

$$
\begin{equation*}
\int_{\Omega} e_{h} \operatorname{curl} \zeta^{I} d x=-\int_{\Omega} p_{h} \operatorname{curl} \zeta^{I} d x \tag{3.13}
\end{equation*}
$$

and

$$
\|\operatorname{curl} \zeta\|_{0, \Omega}^{2}=\int_{\Omega} e_{h} \operatorname{curl} \zeta d x=\int_{\Omega} e_{h} \operatorname{curl}\left(\zeta-\zeta^{I}\right) d x+\int_{\Omega} e_{h} \operatorname{curl} \zeta^{I} d x
$$

and using (3.12) and (3.10b), we obtain

$$
\begin{equation*}
\|\operatorname{curl} \zeta\|_{0, \Omega} \leq C\left(\sum_{l \in E} h_{l}\left\|\left[p_{h} \cdot t_{l}\right]_{l}\right\|_{0, l}^{2}+\sum_{T \in \mathcal{T}_{h}} \varepsilon_{2, T}^{2}\left(p_{h}\right)\right)^{\frac{1}{2}} \tag{3.14}
\end{equation*}
$$

Using the Helmholtz decomposition (3.8) together with the estimates (3.11) and (3.14), we get

$$
\begin{equation*}
\left\|p-p_{h}\right\|_{0, \Omega}^{2} \leq C \sum_{T \in \mathcal{T}_{h}}\left(\eta_{1, T}^{2}\left(p_{h}\right)+\varepsilon_{1, T}^{2}\left(p_{h}\right)+\varepsilon_{2, T}^{2}\left(p_{h}\right)\right) \tag{3.15}
\end{equation*}
$$

Now, let $P_{h} u \in M_{h}$ defined by

$$
\forall T \in \mathcal{T}_{h}, \quad P_{h} u=\frac{1}{\operatorname{meas}(T)} \int_{T} u d x \text { on } T
$$

We have

$$
\begin{equation*}
\left\|u-P_{h} u\right\|_{0, T} \leq C h_{T}\|p\|_{0, T} \leq C\left(h_{T}\left\|p-p_{h}\right\|_{0, T}+h_{T}\left\|p_{h}\right\|_{0, T}\right) . \tag{3.16}
\end{equation*}
$$

On the other hand, using the inf-sup condition we have

$$
\left\|u_{h}-P_{h} u\right\|_{0, \Omega} \leq C \sup _{q_{h} \in E_{h}} \frac{\int_{\Omega}\left(u_{h}-P_{h} u\right) \operatorname{div} q_{h} d x}{\left\|q_{h}\right\|_{H(\operatorname{div}, \Omega)}}
$$

and since

$$
\int_{\Omega}\left(u_{h}-P_{h} u\right) \operatorname{div} q_{h} d x=\int_{\Omega}\left(u_{h} \operatorname{div} q_{h} d x+p_{h} q_{h}\right) d x+\int_{\Omega}\left(p-p_{h}\right) q_{h} d x
$$

we obtain easily that

$$
\begin{equation*}
\left\|u_{h}-P_{h} u\right\|_{0, \Omega} \leq C\left(\sum_{T \in \mathcal{T}_{h}} \varepsilon_{3, T}^{2}\left(p_{h}, u_{h}\right)\right)^{\frac{1}{2}}+\left\|p-p_{h}\right\|_{0, \Omega} \tag{3.17}
\end{equation*}
$$

By triangular inequality, and using (3.15) and (3.17), we have :

$$
\begin{aligned}
\left\|u-u_{h}\right\|_{0, \Omega}+\left\|p-p_{h}\right\|_{0, \Omega} \leq & C\left\{\left(\sum_{T \in \mathcal{T}_{h}}\left(\eta_{1, T}^{2}\left(p_{h}\right)+\varepsilon_{1, T}^{2}\left(p_{h}\right)+\varepsilon_{2, T}^{2}\left(p_{h}\right)\right)\right)^{\frac{1}{2}}\right. \\
& \left.+\left(\sum_{T \in \mathcal{T}_{h}}\left(\eta_{2, T}^{2}\left(p_{h}\right)+\varepsilon_{3, T}^{2}\left(p_{h}, u_{h}\right)\right)\right)^{\frac{1}{2}}\right\} .
\end{aligned}
$$

To indicate the efficiency of the a posteriori error estimator we follow Verfürth [10] and show a local reverse up to higher order perturbations.
For each $T \in \mathcal{T}_{h}$, we reset :

$$
\omega_{T}=\left\{T^{\prime} \in \mathcal{T}_{h} \text { such that } T \text { and } T^{\prime} \text { have a common edge }\right\},
$$

$$
\forall T \in \mathcal{T}_{h}, \quad f_{T}=\frac{1}{\operatorname{meas}(T)} \int_{T} f d x
$$

and for all $l \in E_{I}$, we denote by $T_{+}$and $T_{-}$the two elements of $\mathcal{T}_{h}$ sharing this edge. Let $b_{T}$ be the standard bubble function on $T$ with $\max _{T} b_{T}=1$, as defined in [10]. Then norms $\|\cdot\|_{0, T}$ and $\left\|b_{T} \cdot\right\|_{0, T}$ are equivalent on $P_{0}(T)$, and
$\left\|r_{T}\right\|_{0, T}^{2} \leq C \int_{T} b_{T} r_{T}\left(f_{T}-f+\operatorname{div}\left(p_{h}-p\right)\right) d x \leq C \int_{T} \nabla\left(b_{T} r_{T}\right) .\left(p_{h}-p\right) d x+C\left\|r_{T}\right\|_{0, T}\left\|f-f_{T}\right\|_{0, T}$ where $r_{T}:=f_{T}+\operatorname{div} p_{h}$ on $T$. Then we have

$$
\left\|r_{T}\right\|_{0, T}^{2} \leq C\left|b_{T} r_{T}\right|_{1, T}\left\|p-p_{h}\right\|_{0, T}+C\left\|r_{T}\right\|_{0, T}\left\|f-f_{T}\right\|_{0, T},
$$

Using the inverse estimate $\left|r_{T} b_{T}\right|_{1, T} \leq C h_{T}^{-1}\left\|r_{T}\right\|_{0, T}$, we obtain

$$
\begin{equation*}
h_{T}\left\|f_{T}+\operatorname{div} p_{h}\right\|_{0, T} \leq C_{1}\left\|p-p_{h}\right\|_{0, T}+C_{2} h_{T}\left\|f-f_{T}\right\|_{0, T} . \tag{3.18}
\end{equation*}
$$

Concerning the jump terms for $l \in E_{T} \cap E_{I}$, let $b_{l}$ be the standard bubble function on $T$ vanishing on $\partial T \backslash l$ such that $\max _{T} b_{l}=1$ (see [10]). Then again the norms $\|\cdot\|_{0, l}$ and $\left\|b_{l} \cdot\right\|_{0, l}$ are equivalent on $P_{1}(l)$. Let $l \in E_{I}$, then using the extension operator $P: \mathcal{C}^{0}(l) \longrightarrow$ $\mathcal{C}^{0}\left(T_{+} \cup T_{-}\right)$of [10], it follows that

$$
\left\|\left[p_{h} \cdot t_{l}\right]_{l}\right\|_{0, l}^{2} \leq C \int_{l} b_{l} P\left(\left[p_{h} \cdot t_{l}\right]_{l}\right)\left[p_{h} \cdot t_{l}\right]_{l} d \sigma=C \int_{T_{+} \cup T_{-}} p_{h} \operatorname{curl}\left(b_{l} P\left(\left[p_{h} \cdot t_{l}\right]_{l}\right)\right) d x
$$

Because of the inverse inequality

$$
\left\|b_{l} P\left(\left[p_{h} \cdot t_{l}\right]_{l}\right)\right\|_{1, T+\cup T_{-}} \leq C h_{l}^{-1}\left\|b_{l} P\left(\left[p_{h} \cdot t_{l}\right]_{l}\right)\right\|_{0, T+\cup T_{-}},
$$

the equality

$$
\int_{T_{+} \cup T_{-}} p \operatorname{curl}\left(b_{l} P\left(\left[p_{h} . t_{l}\right]_{l}\right)\right) d x=0
$$

using Cauchy's inequality and $\left\|b_{l} P\left(\left[p_{h} \cdot t_{l}\right]_{l}\right)\right\|_{0, T+\cup T_{-}} \leq C h_{l}^{\frac{1}{2}}\left\|\left[p_{h} \cdot t_{l}\right]_{l}\right\|_{0, l}$, we obtain

$$
\begin{equation*}
h_{l}^{\frac{1}{2}}\left\|\left[p_{h} \cdot t_{l}\right]_{l}\right\|_{0, l} \leq C\left\|p-p_{h}\right\|_{0, T_{+} \cup T_{-}} \tag{3.19}
\end{equation*}
$$

Finally, it is easy to see that

$$
\begin{equation*}
\forall i=1,2 \quad \varepsilon_{i, T}\left(p_{h}\right) \leq\left\|p-p_{h}\right\|_{0, \Delta(T)} . \tag{3.20}
\end{equation*}
$$

Combining (3.18)-(3.20) we obtain

$$
\begin{equation*}
\eta_{1, T}^{2}\left(p_{h}\right)+\varepsilon_{1, T}^{2}\left(p_{h}\right)+\varepsilon_{2, T}^{2}\left(p_{h}\right) \leq C_{1} \sum_{T^{\prime} \in \Delta(T)}\left\|p-p_{h}\right\|_{0, T^{\prime}}^{2}+C_{2} \sum_{T^{\prime} \in \omega_{T}} h_{T^{\prime}}^{2}\left\|f-f_{T^{\prime}}\right\|_{0, T^{\prime}}^{2} \tag{3.21}
\end{equation*}
$$

Introducing as above the bubble function $b_{T}$, we get

$$
\begin{equation*}
\left\|p_{h}\right\|_{0, T}^{2} \leq C \int p_{h} b_{T} p_{h} d x=C \int_{T} b_{T} p_{h}\left(p_{h}-p\right) d x+C \int_{T} \nabla u \cdot\left(b_{T} p_{h}\right) . \tag{3.22}
\end{equation*}
$$

Since $\int_{T} \operatorname{div}\left(b_{T} p_{h}\right) d x=0$ and $u_{h} \in M_{h}$, by element-wise integration by parts we have

$$
\begin{equation*}
\int_{T} \nabla u \cdot\left(b_{T} p_{h}\right) d x=\int_{T}\left(u_{h}-u\right) \operatorname{div}\left(b_{T} p_{h}\right) d x . \tag{3.23}
\end{equation*}
$$

Using the inverse inequality $\left\|\operatorname{div}\left(b_{T} p_{h}\right)\right\|_{0, T} \leq C h_{T}\left\|p_{h}\right\|_{0, T}$, and (3.22)-(3.23), we obtain

$$
\begin{equation*}
\eta_{2, T}\left(p_{h}\right) \leq C\left(\left\|p-p_{h}\right\|_{0, T}+\left\|u-u_{h}\right\|_{0, T}\right) . \tag{3.24}
\end{equation*}
$$

Finally, it is easy to check that

$$
\begin{equation*}
\varepsilon_{3, T}\left(p_{h}, u_{h}\right) \leq C\left(\left\|p-p_{h}\right\|_{0, \Delta(T)}+\left\|u-u_{h}\right\|_{0, \Delta(T)}\right) . \tag{3.25}
\end{equation*}
$$

Using the estimates (3.24)-(3.25), we get

$$
\begin{equation*}
\forall T \in \mathcal{T}_{h}, \quad \eta_{2, T}\left(p_{h}\right)+\varepsilon_{3, T}\left(p_{h}, u_{h}\right) \leq C\left(\left\|p-p_{h}\right\|_{0, \Delta(T)}+\left\|u-u_{h}\right\|_{0, \Delta(T)}\right) \tag{3.26}
\end{equation*}
$$

Finally, using (3.21) and (3.26) concludes the proof.

## References

[1] R.A. Adams. Sobolev Spaces. Academic Press, New York (1975).
[2] F. Brezzi, M. Fortin and L.D. Marini. Error analysis of piecewise constant pressure approximations of Darcy's law. Comput. Methods Appl. Mech. Engrg. 195, no. 13-16, pp. 1547-1559 (2006).
[3] C. Carstensen and R.H.W. Hoppe. Error reduction and convergence for an adaptive mixed finite element method. Math. Comp. Vol 75, No 255, pp. 1033-1042 (2006)
[4] C. Carstensen and R.H.W. Hoppe. Convergence analysis of an adaptive nonconforming finite element method. Numer. Math. Vol 103, No 2, pp. 251-266 (2006)
[5] P. G. Ciarlet. Handbook of numerical analysis. Vol. II. Eds. Ciarlet, P. G. and Lions, J.-L.. Finite element methods. Part 1. North-Holland, Amsterdam (1991).
[6] R. Eymard, T. Gallouët and R. Herbin. Finite volume methods. Handbook of numerical analysis, Vol. VII, pp 713-1020, North-Holland,Amsterdam 2000.
[7] R. Eymard, T. Gallouët and R. Herbin. A finite volume scheme for anisotropic diffusion problems. C.R. Acad. Sci. Paris, Ser I 339 (2004).
[8] Scott L. Ridgway and Zhang Shangyou . Finite element interpolation of non smooth functions satisfying boundary conditions . Math. Comp., Vol 54, No 190, pp 483-493(1990)
[9] William J. Layton, Friedhelm Schieweck and Ivan Yotov. Coupling fluid flow with porous media flow. SIAM J. Numer. Anal. Vol 40, No 6, pp. 2195-2218 (2003).
[10] R. Verfürth. A review of a posteriori error estimation and adaptive mesh-refinement techniques.Wiley Teubner, 3 (1996)


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