



HAL
open science

Boundary blow-up solutions in the unit ball: asymptotics, uniqueness and symmetry

Ovidiu Costin, Louis Dupaigne

► **To cite this version:**

Ovidiu Costin, Louis Dupaigne. Boundary blow-up solutions in the unit ball : asymptotics, uniqueness and symmetry. 2009. hal-00356579v1

HAL Id: hal-00356579

<https://hal.science/hal-00356579v1>

Preprint submitted on 27 Jan 2009 (v1), last revised 10 Mar 2010 (v3)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Boundary blow-up solutions in the unit ball : asymptotics, uniqueness and symmetry

O. Costin¹ and L. Dupaigne²

January 27, 2009

¹Department of Mathematics, The Ohio State University,
100 Math Tower, 231 West 18th Avenue, Columbus, OH 43210-1174, USA
costin@math.ohio-state.edu

²LAMFA, UMR CNRS 6140, Université Picardie Jules Verne
33, rue St Leu, 80039 Amiens, France
louis.dupaigne@math.cnrs.fr

Abstract

We calculate the full asymptotic expansion of boundary blow-up solutions (see equation (1) below), for any nonlinearity f . Our approach enables us to state sharp qualitative results regarding uniqueness and radial symmetry of solutions, as well as a characterization of nonlinearities for which the blow-up rate is universal. At last, we study in more detail the standard nonlinearities $f(u) = u^p$, $p > 1$.

1 Introduction

Let B denote the unit ball of \mathbb{R}^N , $N \geq 1$ and let $f \in C(\mathbb{R})$, $f > 0$. We study the equation

$$(1) \quad \begin{cases} \Delta u = f(u) & \text{in } B, \\ u = +\infty & \text{on } \partial B, \end{cases}$$

where the boundary condition is understood in the sense that

$$\lim_{x \rightarrow x_0, x \in B} u(x) = +\infty \quad \text{for all } x_0 \in \partial B.$$

A function u satisfying (1) is called a boundary blow-up solution or simply a large solution. Existence of a solution of (1) is equivalent to the so-called Keller-Osserman condition :

$$(2) \quad \int^{+\infty} \frac{dt}{\sqrt{F(t)}} < +\infty, \quad \text{where } F' = f.$$

For a proof of this fact, see the seminal works of J.B. Keller [55, 5] and R. Osserman [77, 7] for the case of monotone f , as well as [33, 3] for the general case. From here on, we always assume that (2) holds.

Our goal here is to study asymptotics, uniqueness and symmetry properties of solutions. Our approach improves known results in at least two directions : firstly, aside the necessary condition (2), we need not make *any* additional assumption on f to obtain the sharp asymptotics of solutions. Secondly, we obtain the complete asymptotic expansion of solutions, to *all* orders. Here is a summary of our findings.

Theorem 1.1 *Let $f \in C(\mathbb{R})$, $f > 0$ and assume (2) holds. Let u_1, u_2 denote two solutions of (1). Then,*

$$\lim_{x \rightarrow x_0, x \in B} u_1(x) - u_2(x) = 0, \quad \text{for all } x_0 \in \partial B.$$

More precisely, there exists a constant $C = C(u_1, u_2, N, F) > 0$, such that for all $x \in B$,

$$(3) \quad |u_1(x) - u_2(x)| \leq C \int_{u_2(x)}^{+\infty} \frac{dt}{F(t)} dt.$$

In addition,

$$(4) \quad |F(u_1) - F(u_2)| \in L^\infty(B).$$

Estimates on the gradient of solutions can be obtained for a restricted class of nonlinearities, namely

Theorem 1.2 *Let $f \in C(\mathbb{R})$, $f > 0$ and assume (2) holds. Assume in addition that f is increasing up to a linear perturbation i.e. there exists an increasing function \tilde{f} and a constant K such that*

$$(5) \quad f(t) = \tilde{f}(t) - Kt,$$

for all $t \in \mathbb{R}$. Let u_1, u_2 denote two solutions of (1). Then,

$$(6) \quad |\nabla(u_1 - u_2)| \in L^\infty(B).$$

As a direct consequence of Theorem 1.1, we obtain the following uniqueness result.

Corollary 1.3 *Let $f \in C(\mathbb{R})$, $f > 0$ and assume (2) holds. Assume in addition that f is nondecreasing. Then, there exists a unique large solution of (1).*

Remark 1.4 *Many uniqueness theorems have been established in the literature (see e.g. the survey [11, 1]), and they hold for a general class of bounded domains Ω . However, in all of these results, additional assumptions on f are needed, such as convexity.*

Proof of Corollary 1.3 . Let u_1, u_2 denote two large solutions. It suffices to prove that $u_1 \leq u_2$. Assume this is not the case and let $\omega = \{x \in B : w(x) > 0\} \neq \emptyset$, where $w = u_1 - u_2$. Working if necessary on a connected component of ω , we may always assume that ω is connected. Using Theorem 1.1, we see that w solves the equation

$$\begin{cases} \Delta w = f(u_1) - f(u_2) \geq 0 & \text{in } \omega, \\ w = 0 & \text{on } \partial\omega. \end{cases}$$

By the Maximum Principle, $w \leq 0$ in ω , a contradiction. \square

When f is not increasing, uniqueness fails in general. One may ask however whether all solutions of (1) are radial. A. Porretta and L. Véron proved that this is indeed the case under the additional assumption that f is asymptotically convex (see [88,8]). We generalize their result as follows.

Corollary 1.5 *Let $f \in C(\mathbb{R})$, $f > 0$ and assume (2) holds. Let u denote a solution of (1). Assume in addition that f is increasing up to a linear perturbation (i.e. (5) holds for some nondecreasing function \tilde{f} and some constant K). Then, u is radially symmetric. Furthermore, $\frac{\partial u}{\partial r} > 0$ in $B \setminus \{0\}$.*

Remark 1.6 *We recall that in the setting of the classical symmetry result of B. Gidas, W. M. Ni and L. Nirenberg (see [44,4]), (5) is also assumed in order to prove symmetry. In the same article, the authors give an example of a nonlinearity f failing (5) for which there do exist nonradial solutions of the equation. In the context of large solutions, we are not able to produce such a counter-example (and expect none exists). In particular, we do not know whether radial symmetry remains true for simple nonlinearities such as $f(u) = u^2(1 + \sin u)$.*

Corollary 1.5 is a direct consequence of the moving plane method and Theorem 1.2 :

Proof of Corollary 1.5 . Let U denote a radial solution of (1). Since $f \geq 0$, U is a nondecreasing function of $r = |x|$ and we must have $\frac{dU}{dr}(r) \rightarrow +\infty$ as $r \rightarrow 1^-$. By (6), we conclude that any solution u of (1) satisfies $\frac{\partial u}{\partial r}(x) \rightarrow +\infty$ as $x \rightarrow \partial B$, while the tangential part of the gradient of u remains bounded. We then apply Theorem 2.1 in [88,8]. \square

In addition to the *relative* asymptotic information given by (3), (4) and (6), the *exact* asymptotic expansion of a solution can be calculated to all orders, as follows.

Theorem 1.7 *Let $f \in C(\mathbb{R})$, $f > 0$ and assume (2) holds. Let $U_0 \in \mathbb{R}$, $I = [U_0, +\infty)$ and let v_0 be the function defined for $u \in I$ by*

$$(7) \quad v_0(u) = \sqrt{2F(u)}.$$

Consider the Banach space

$$\mathcal{X} = \{v \in C(I; \mathbb{R}) : \exists M > 0 \text{ such that } |v| \leq Mv_0\},$$

endowed with the norm $\|v\| = \sup_I |v/v_0|$. If the constant U_0 is chosen sufficiently large, there exists $\rho \in (0, 1)$ such that there exists a unique solution v belonging to the ball $\mathcal{B}(v_0, \rho) \subset \mathcal{X}$ of the integral equation

$$(8) \quad v(u) = \sqrt{2 \left(F(u) - (N-1) \int_{U_0}^u \frac{v}{r} dt \right)}, \quad u \in I,$$

where $r = r(u, v)$ is given for $u \in I$, $v \in \mathcal{B}(v_0, \rho)$ by

$$(9) \quad r(u, v) = 1 - \int_u^{+\infty} \frac{1}{v} dt.$$

In addition, v is the limit in X of (v_k) defined for $k = 0$ by (7) and for $k \geq 1$, by

$$(10) \quad v_k(u) = \sqrt{2 \left(F(u) - (N-1) \int_{U_0}^u \frac{v_{k-1}}{1 - \int_t^{+\infty} \frac{1}{v_{k-1}} ds} dt \right)}$$

and the sequence v_k is asymptotic to v i.e. as $u \rightarrow +\infty$,

$$v_{k+1}(u) = v_k(u) + o(v_k(u))$$

and given any $k \in \mathbb{N}$,

$$v(u) = v_k(u) + O(v_{k+1}(u) - v_k(u)).$$

Let now u denote any solution of (1) and fix $r_0 \in (0, 1)$ such that $u(x) \geq U_0$ for $|x| \geq r_0$. For $k \geq 0$, define u_k for $r \geq r_0$ as the unique solution of

$$(11) \quad \begin{cases} \frac{du_k}{dr} = v_k(u_k) \\ \lim_{r \rightarrow 1^-} u_k(r) = +\infty, \end{cases}$$

where v_k is given by (10). Then, as $r \rightarrow 1^-$,

$$\int_{u_k(r)}^{u_{k+1}(r)} \frac{du}{v_0} = o \left(\int_{u_k(r)}^{+\infty} \frac{du}{v_0} \right)$$

and given any $k \in \mathbb{N}$, we have as $x \rightarrow \partial B$,

$$(12) \quad \int_{u_k(|x|)}^{u(x)} \frac{du}{v_0} = o \left(\int_{u_k(|x|)}^{+\infty} \frac{du}{v_0} \right)$$

Theorem 1.7 enables one to calculate (implicitly) the asymptotic expansion of a solution term by term. But how many terms in this expansion are singular? This is what we discuss in our last set of results.

We begin with the simplest class of nonlinearities f , those for which only one term in the expansion is singular, namely the function u_0 defined by (7) and (11). It turns out, as A.C. Lazer and P.J. Mc Kenna first demonstrated (see [66, 6]), that in this case $u_0(1-d(x))$ is the only singular term in the asymptotics of any blow-up solution on any smoothly bounded domain $\Omega \subset \subset \mathbb{R}^N$ and for any dimension $N \geq 1$, where $d(x)$ denotes the distance of a point $x \in \Omega$ to the boundary of Ω . In other words, the blow-up rate is universal. We characterize these nonlinearities as follows:

Theorem 1.8 *Let $\Omega \subset \mathbb{R}^N$ denote a bounded domain satisfying an inner and an outer sphere condition at each point of its boundary. Let $f \in C(\mathbb{R})$, $f \geq 0$, assume (2) holds and consider the equation*

$$(13) \quad \begin{cases} \Delta u = f(u) & \text{in } \Omega, \\ u = +\infty & \text{on } \partial\Omega. \end{cases}$$

Assume

$$(14) \quad \lim_{u \rightarrow +\infty} \sqrt{2F(u)} \int_u^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt = 0.$$

Then, any solution of (13) satisfies

$$(15) \quad \lim_{x \rightarrow \partial\Omega} u(x) - u_0(1 - d(x)) = 0,$$

where $d(x) = \text{dist}(x, \partial\Omega)$ and u_0 is defined by (7), (11).

We also have the following partial converse statement : if

$$(16) \quad \liminf_{u \rightarrow +\infty} \sqrt{2F(u)} \int_u^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt > 0,$$

then (15) always fails.

Remark 1.9 To our knowledge, (14) improves upon all known conditions for (15) to hold (see in particular [66, 6] and [22, 2]). Despite its unappealing technical looks, (14) uses information on the asymptotics of F only (in particular, no direct information on f is required). Nonlinearities such that $F(u) \sim e^u$ or $F(u) \sim u^p$, $p > 4$ as $u \rightarrow +\infty$ qualify. For $F(u) \sim u^4$, (16) holds and so the conclusion (15) fails.

Remark 1.10 Condition (14) can be weakened to :

$$\lim_{r \rightarrow 1^-} \sqrt{2F(u_0)} \int_{u_0}^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt = 0,$$

where $u_0 = u_0(r)$ is defined by (11). Similarly, (16) can be weakened to

$$\liminf_{r \rightarrow 1^-} \sqrt{2F(u_1)} \int_{u_1}^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt > 0.$$

As an immediate corollary, we obtain uniqueness on general domains, whenever only one singular term appears:

Corollary 1.11 Assume (14). If in addition, f is nondecreasing, then the solution of (13) is unique.

Proof. Simply repeat the proof of Corollary 1.3. \square

More than one term can be present in the asymptotic expansion of u . Finding all the (singular) terms in this expansion is of staggering algebraic complexity. To illustrate this, we provide the first three terms (in implicit form).

Proposition 1.12 Let u_2 be defined by (11) for $k = 2$. Let also R_1, R_2, R_3 denote three real-valued functions defined for $U \in \mathbb{R}$ sufficiently large

by

$$\begin{aligned}
R_0(U) &= \int_U^{+\infty} \frac{du}{\sqrt{2F}}, \\
R_1(U) &= (N-1) \int_U^{+\infty} \frac{\int^u \sqrt{2F} dt}{(2F)^{3/2}} du, \\
R_2(U) &= (N-1) \int_U^{+\infty} \left(- \int^u \left((N-1) \frac{\int^t \sqrt{2F} ds}{\sqrt{2F}} + \sqrt{2F} \int_u^{+\infty} \frac{ds}{\sqrt{2F}} \right) dt + \right. \\
&\quad \left. + \frac{5(N-1)}{4} \frac{\left(\int^u \sqrt{2F} dt \right)^2}{2F} \right) \frac{du}{(2F)^{3/2}}.
\end{aligned}$$

Then, for all $r \in (0, 1)$, r close to 1, there holds

$$1 - r = R_0(u_2(r)) + R_1(u_2(r)) + R_2(u_2(r))(1 + o(1)).$$

For specific nonlinearities, it is possible to invert the above identity. This is what we do for $f(u) = u^p$, $p > 1$:

Proposition 1.13 *Let $p > 1$ and $f(u) = u^p$, for $u > 0$. Then, the unique positive solution of (1) satisfies as $r \rightarrow 1^-$,*

$$u = d^{-\frac{2}{p-1}} \sum_{k=0}^{[2/(p-1)]} a_k d^k + o(1),$$

where $d(r) = 1 - r$ for $r \in (0, 1)$, and where each $a_k \in \mathbb{R}$ depends on N and p only.

Remark 1.14 *Note that the above result remains true for any nonlinearity f such that, for some positive constant c , $F(u) \sim cu^{p+1}$ as $u \rightarrow +\infty$ (and any solution of the equation).*

Outline of the paper

1. In the next section, we show that any solution u of (1) can be squeezed between two radial solutions U and V i.e. the inequality $U \leq u \leq V$ holds throughout B .
2. Thanks to this result, we need only find the asymptotics of *radial* solutions to prove Theorem 1.1. This is what we do in Section 3.
3. To obtain gradient estimates, the squeezing technique is insufficient and more work is needed. In Section 4, we estimate tangential derivatives via a standard comparison argument, while we gain control over the radial component through a more delicate potential theoretic argument.
4. Section 5 is dedicated to the proof of Theorem 1.7, that is we establish an algorithm for computing the asymptotics of solutions to all orders.
5. In Section 6, we characterize nonlinearities for which the blow-up rate is universal.

6. At last, Sections 7 and 8 contain the tedious calculations, which enable us to calculate the first three terms of the asymptotic expansion of u in implicit form for general f , and all terms explicitly for $f(u) = u^p$.

2 Ordering solutions

In this section, we prove that any solution of the equation is bounded above and below by radial blow-up solutions. To do so, we impose the following additional condition

$$(17) \quad \int^{+\infty} \frac{dt}{\sqrt{G(t)}} = +\infty, \quad \text{where } G'(t) = f(-t).$$

Remark 2.1 *Note that (17) is not restrictive. Indeed, if u denotes a solution of (1) and $m = \min_B u$, then u also solves (1) with nonlinearity \tilde{f} defined for $u \in \mathbb{R}$ by*

$$\tilde{f}(u) = \begin{cases} f(m) & \text{if } u < m \\ f(u) & \text{if } u \geq m. \end{cases}$$

Then, \tilde{f} clearly satisfies (17).

We now proceed through a series of three lemmas.

Lemma 2.2 *Assume (17) holds. For every $M \in \mathbb{R}$, there exists a radial function $\underline{v} \in C^2(B) \cap C(\overline{B})$ satisfying*

$$\Delta \underline{v} \geq f(\underline{v}) \quad \text{in } B$$

and such that

$$\underline{v} \leq -M \quad \text{in } B.$$

Proof. Let $g(t) = f(-t)$ for $t \in \mathbb{R}$ and let $a > M$ a parameter to be fixed later on. Let w denote a solution of

$$(18) \quad \begin{cases} -w'' = g(w) \\ w(0) = a \\ w'(0) = 0 \end{cases}$$

We claim that for $a > M$ sufficiently large, $w(1) \geq M$. Indeed w is clearly nonincreasing. So, either $w > M$. In particular, w is defined on all \mathbb{R}^+ and $w(1) \geq M$. Either there exists $R > 0$ such that $w(R) = M$ and we just need to prove that $R \geq 1$. To do so, multiply (18) by $-w'$ and integrate between 0 and $r \in (0, R)$:

$$-w' = \sqrt{2(G(a) - G(w))},$$

where G is an antiderivative of g . Integrate again between 0 and R :

$$\int_M^a \frac{dt}{\sqrt{2(G(a) - G(t))}} = \int_0^R \frac{-w'}{\sqrt{2(G(a) - G(w))}} dr = R.$$

Now,

$$\begin{aligned} R &= \int_M^a \frac{dt}{\sqrt{2(G(a) - G(t))}} \geq \int_M^{G^{-1}(G(a)/2)} \frac{dt}{\sqrt{2(G(a) - G(t))}} \geq \\ &\geq \int_M^{G^{-1}(G(a)/2)} \frac{dt}{\sqrt{2G(t)}}. \end{aligned}$$

By (17), we deduce that $R \geq 1$ for a sufficiently large. We have just proved that $w|_{(0,1)} \geq w(1) \geq M$. It follows that the function \underline{v} defined for $x \in B$ by $\underline{v}(x) = -w(|x|)$, is the desired subsolution. \square

Lemma 2.3 *Let $f \in C(\mathbb{R})$, $f > 0$ and assume (2) holds. Assume $\underline{v} \in C(\bar{B})$ satisfies*

$$\Delta \underline{v} \geq f(\underline{v}) \quad \text{in } B.$$

Then, there exists a radial large solution V of (1) such that $V \geq \underline{v}$.

Proof. Let $\bar{v} := N$. Then, \underline{v} and \bar{v} are respectively a sub and supersolution of

$$(19) \quad \begin{cases} \Delta v = f(v) & \text{in } B, \\ v = N & \text{on } \partial B, \end{cases}$$

provided N is chosen so large that $N > \|\underline{v}\|_{L^\infty(B)}$. Furthermore, $\underline{v} < \bar{v}$ in B for such values of N . By the method of sub and supersolutions (see e.g. Proposition 2.1 in [33,3]), there exists a minimal solution V_N of (19) such that $N \geq V_N \geq \underline{v}$. Note that V_N is radial, as follows from the classical symmetry result of Gidas, Ni and Nirenberg (see [44,4]) when f is locally Lipschitz, or simply from the fact that V_N is minimal, hence radial, for f merely continuous. Also, since V_N is minimal, we have that the sequence (V_N) is nondecreasing with respect to N (apply e.g. the Minimality Principle, Corollary 2.2, in [33,3]).

It turns out that the sequence (V_N) is uniformly bounded on compact sets of B . Indeed, fix $R_1 < 1$. There exists a solution \tilde{U} blowing up on the boundary of the ball of radius 1 and satisfying $\tilde{U} \geq \underline{v}$ in B_{R_1} , see Remark 2.9 in [33,3]. By minimality, $\underline{v} \leq V_N \leq \tilde{U}$ in B_{R_1} , whence (V_N) is uniformly bounded on B_{R_2} , for any given $R_2 < R_1$.

We have just proved that each V_N is radial and that the sequence (V_N) is nondecreasing and bounded on compact subsets of B . By standard elliptic regularity, it follows that (V_N) converges to a radial solution V of (1), such that $V \geq \underline{v}$ in B . \square

Lemma 2.4 *Assume (2) and (17) hold. Let u denote a solution of (1). Then, there exist two radial functions U, V solving (1) such that*

$$U \leq u \leq V \quad \text{in } B.$$

Proof. Let $-M$ denote the minimum value of u and let \underline{v} denote the subsolution given by Lemma 2.2. In particular, $\underline{v} \leq u$. By Lemma 2.3, there exists a solution $U \geq \underline{v}$ of (1) and we may assume that U is the

minimal solution relative to \underline{v} i.e. given any other solution $\tilde{u} \geq \underline{v}$ of (1), $U \leq \tilde{u}$. In particular, $U \leq u$. It remains to construct a radial solution V of (1) such that $u \leq V$. To do so, we fix $R < 1$. By Lemma 2.3, letting $\underline{v} = u|_{B_R}$, there exists a radial solution $v = V_R$ of

$$(20) \quad \begin{cases} \Delta v = f(v) & \text{in } B_R, \\ v = +\infty & \text{on } \partial B_R, \end{cases}$$

such that $V_R \geq u$ in B_R . Since V_R is constructed as the monotone limit of minimal solutions V_N (see the proof of the previous lemma), one can easily check that the mapping $R \mapsto V_R$ is nonincreasing (hence automatically bounded on compact sets of B). Hence, as $R \rightarrow 1$, V_R converges to a solution V of (1), which is radial and satisfies $V \geq u$ in B , as desired. \square

3 Asymptotics of radial solutions

Our next result establishes that the asymptotic expansion of a radial blow-up solution is unique. More precisely, consider the one-dimensional problem

$$(21) \quad \frac{d^2 \phi}{dr^2} = f(\phi), \quad r < 1, \quad \phi(r) \rightarrow +\infty \text{ as } r \rightarrow 1^-.$$

All solutions are given implicitly by

$$\int_{\phi}^{+\infty} \frac{ds}{\sqrt{2F(s)}} = 1 - r, \quad \text{where } F' = f.$$

We recall the following fact, first observed by C. Bandle and M. Marcus in [22, 2] :

Remark 3.1 *Let ϕ and ϕ_c denote two solutions of (21) corresponding to the antiderivatives F and $F + c$, respectively. Then $\phi(r) - \phi_c(r) \rightarrow 0$ as $r \rightarrow 1^-$.*

We improve this result in the following way :

Theorem 3.2 *Let $N \geq 1$ and let u_1, u_2 denote two strictly increasing functions solving*

$$(22) \quad \begin{cases} \frac{d^2 u}{dr^2} + \frac{N-1}{r} \frac{du}{dr} = f(u), & r < 1, \\ \lim_{r \rightarrow 1^-} u(r) = +\infty. \end{cases}$$

Then,

$$|u_1(r) - u_2(r)| \leq C \int_{u_2(r)}^{+\infty} \frac{dt}{F(t)} dt.$$

In addition, the quantity $|F(u_1) - F(u_2)|$ is bounded.

Remark 3.3 *Clearly, Theorem 1.1 follows as a direct consequence of Remark 2.1, Lemma 2.4 and Theorem 3.2.*

Proof. We want to think of the second term on the left-hand side of equation (22) as a lower order perturbation as $r \rightarrow 1$. So, we integrate (22) in the same way we would solve (21), namely we let $v = du/dr$ and multiply the equation by v . We get

$$\frac{d}{dr} \left(\frac{v^2}{2} \right) + \frac{N-1}{r} v^2 = \frac{d}{dr} (F(u)).$$

We define the resulting error term by

$$(23) \quad g := -\frac{v^2}{2} + F(u),$$

which, seen as function of r , satisfies the differential equation

$$(24) \quad \frac{dg}{dr} = \frac{N-1}{r} v^2.$$

Since u is a strictly increasing function, the change of independent variable $u = u(r)$ is valid. Thinking of g as a function of the variable u , we have $\frac{dg}{du} = \frac{dg}{dr} \frac{dr}{du} = \frac{1}{v} \frac{dg}{dr}$ and so

$$(25) \quad \frac{dg}{du} = \frac{N-1}{r} v.$$

Since (22) holds for r close to 1, the above equation holds for u in a neighborhood of $+\infty$. Solving (23) for v , we finally obtain

$$(26) \quad \begin{cases} \frac{dg}{du} = \frac{N-1}{r} v = \frac{N-1}{r} \sqrt{2(F(u) - g)}, \\ \frac{dr}{du} = 1/v = (2(F(u) - g))^{-1/2}. \end{cases}$$

We start by calculating the leading asymptotic behaviour of g at $+\infty$:

Lemma 3.4

$$\lim_{u \rightarrow +\infty} \frac{g(u)}{F(u)} = 0.$$

In addition,

$$(27) \quad \lim_{u \rightarrow +\infty} \frac{g(u)}{(N-1)G(u)} = 1, \quad \text{where } G \text{ is any antiderivative of } \sqrt{2F}.$$

Proof. First, we claim that

$$(28) \quad \lim_{u \rightarrow +\infty} \frac{G(u)}{F(u)} = 0.$$

Indeed, fix $\varepsilon > 0$ and recalling that (2) holds, choose $M > 0$ so large that $\int_M^{+\infty} \frac{dt}{\sqrt{2F(t)}} < \varepsilon$. By the definition of G , there exists a constant C_M such that

$$G(u) = C_M + \int_M^u \sqrt{2F(t)} dt.$$

Since F is nondecreasing it follows that

$$G(u) \leq C_M + 2F(u) \int_M^u \frac{dt}{\sqrt{2F(t)}} \leq C_M + 2\varepsilon F(u).$$

Dividing by $F(u)$ and letting $u \rightarrow +\infty$, (28) follows. Next, we claim that

$$(29) \quad \lim_{u \rightarrow +\infty} \frac{g(u)}{F(u)} = 0.$$

Note that by (25), $g(u)$ is increasing, thus bounded below by a constant c as $u \rightarrow +\infty$. Hence, by (26),

$$\frac{dg}{du} \leq \frac{N-1}{r} \sqrt{2(F(u)-c)} \leq 2(N-1) \sqrt{2(F(u)-c)},$$

where the last inequality holds if $r > 1/2$ i.e. if u is sufficiently large. Integrating on a given interval (u_0, u) , we obtain

$$c \leq g(u) \leq g(u_0) + 2(N-1) \int_{u_0}^u \sqrt{2(F(t)-c)} dt.$$

Using (28) and the fact that $\lim_{t \rightarrow +\infty} \frac{\sqrt{2F(t)}}{\sqrt{2(F(t)-c)}} = 1$, we deduce (29).

Now that (29) has been established, we return to (26) and infer that given $\varepsilon > 0$, we have for sufficiently large u ,

$$\frac{dg}{du} \geq \frac{N-1}{r} \sqrt{2(1-\varepsilon)F(u)} \geq (N-1) \sqrt{2(1-\varepsilon)F(u)}$$

and

$$\frac{dg}{du} \leq \frac{N-1}{r} \sqrt{2(1+\varepsilon)F(u)} \leq \frac{N-1}{1-\varepsilon} \sqrt{2(1+\varepsilon)F(u)}.$$

Integrating the above, we finally obtain for large u ,

$$(1-\varepsilon)(N-1) \int_{u_0}^u \sqrt{2F(t)} dt \leq g(u) - g(u_0) \leq (1+\varepsilon)^{3/2}(N-1) \int_{u_0}^u \sqrt{2F(t)} dt$$

and (27) follows. The fact that $g(u) \rightarrow +\infty$ as $u \rightarrow +\infty$ follows automatically. \square

Next, we prove that given two solutions u_1, u_2 , the corresponding error terms g_1, g_2 given by (23) differ by a bounded quantity.

Lemma 3.5 *Let u_1, u_2 denote two solutions of (22). Introduce $v_i = \frac{du_i}{dr}$ and*

$$g_i = -\frac{v_i^2}{2} + F(u_i), \quad \text{for } i = 1, 2.$$

Then, $g_1 - g_2$ is bounded.

Proof. We have seen in the course of the proof of Lemma 3.4 that each g_i (seen as a function of a variable u lying in some interval $(M_i, +\infty)$) solves the differential equation (26). Letting $w = g_1 - g_2$, it follows that

$$\begin{aligned} \frac{dw}{du} &= \frac{N-1}{r} \left(\sqrt{2(F(u)-g_1)} - \sqrt{2(F(u)-g_2)} \right) \\ &= -\frac{2(N-1)}{r} \frac{1}{\sqrt{2(F(u)-g_1)} + \sqrt{2(F(u)-g_2)}} w. \end{aligned}$$

So, for some constants C, u_0 ,

$$w = C \exp \left(- \int_{u_0}^u \frac{2(N-1)}{r} \frac{1}{\sqrt{2(F(t)-g_1)} + \sqrt{2(F(t)-g_2)}} dt \right),$$

whence $|w| \leq |C|$ for $u > u_0$. \square

Proof of Theorem 3.2 completed. Let u_1, u_2 denote two solutions of (22). By (26), each $u_i, i = 1, 2$, solves

$$\frac{du_i/dr}{\sqrt{2(F(u_i)-g_i)}} = 1.$$

Integrating, we obtain

$$\int_{u_1}^{+\infty} \frac{1}{\sqrt{2(F(t)-g_1)}} dt = 1 - r = \int_{u_2}^{+\infty} \frac{1}{\sqrt{2(F(t)-g_2)}} dt.$$

Without loss of generality, for a given r we may assume $u_2(r) \geq u_1(r)$. Split the left-hand side integral: $\int_{u_1}^{+\infty} = \int_{u_1}^{u_2} + \int_{u_2}^{+\infty}$. It follows that

$$\begin{aligned} \int_{u_1}^{u_2} \frac{1}{\sqrt{2(F(t)-g_1)}} dt &= \int_{u_2}^{+\infty} \left(\frac{1}{\sqrt{2(F(t)-g_2)}} - \frac{1}{\sqrt{2(F(t)-g_1)}} \right) dt \\ &= \int_{u_2}^{+\infty} \frac{\sqrt{2(F(t)-g_1)} - \sqrt{2(F(t)-g_2)}}{\sqrt{2(F(t)-g_1)}\sqrt{2(F(t)-g_2)}} dt \\ &= \int_{u_2}^{+\infty} \frac{g_2 - g_1}{\sqrt{2(F(t)-g_1)}\sqrt{2(F(t)-g_2)}(\sqrt{2(F(t)-g_1)} + \sqrt{2(F(t)-g_2)})} dt \end{aligned}$$

Recall that by Lemma 3.4, $g_i = o(F)$ as $t \rightarrow +\infty$. Recall also that $g_2 - g_1$ is bounded. So, for sufficiently large values of u_2 , we deduce

$$(30) \quad \int_{u_1}^{u_2} \frac{1}{\sqrt{2F(t)}} dt \leq C \int_{u_2}^{+\infty} \frac{dt}{F(t)^{3/2}}.$$

Since F is increasing, it follows that

$$\begin{aligned} 0 \leq \frac{u_2 - u_1}{\sqrt{F(u_2)}} &\leq C \int_{u_1}^{u_2} \frac{1}{\sqrt{2F(t)}} dt \leq C \int_{u_2}^{+\infty} \frac{dt}{F(t)^{3/2}} \\ &\leq \frac{C}{\sqrt{F(u_2)}} \int_{u_2}^{+\infty} \frac{dt}{F(t)}. \end{aligned}$$

Hence,

$$0 \leq u_2 - u_1 \leq C \int_{u_2}^{+\infty} \frac{dt}{F(t)},$$

as stated in Theorem 3.2. It remains to prove (4). Without loss of generality, we assume $u_1(r) \leq u_2(r)$ so

$$\begin{aligned}
\int_{u_1}^{u_2} \frac{dt}{\sqrt{F(t)}} &= \int_{u_1}^{+\infty} \frac{dt}{\sqrt{F(t)}} - \int_{u_2}^{+\infty} \frac{dt}{\sqrt{F(t)}} \\
&= \int_{u_2}^{+\infty} \frac{dt}{\sqrt{F(t - (u_2 - u_1))}} - \int_{u_2}^{+\infty} \frac{dt}{\sqrt{F(t)}} \\
&= \int_{u_2}^{+\infty} \frac{\sqrt{F(t)} - \sqrt{F(t - (u_2 - u_1))}}{\sqrt{F(t)}F(t - (u_2 - u_1))} dt \\
&= \int_{u_2}^{+\infty} \frac{F(t) - F(t - (u_2 - u_1))}{\sqrt{F(t)}F(t - (u_2 - u_1))(\sqrt{F(t)} + \sqrt{F(t - (u_2 - u_1))})} dt \\
&\geq (F(u_2) - F(u_1)) \int_{u_2}^{+\infty} \frac{dt}{F(t)^{3/2}}.
\end{aligned}$$

Recalling (30), (4) follows. \square

4 Gradient estimates

Proof of Theorem 1.2. Let $w = u_1 - u_2$ denote the difference of two solutions. Without loss of generality, we may assume that u_2 is the minimal solution of (1), so that $u_1 \geq u_2$ and u_2 is radial.

Step 1 : estimate of tangential derivatives

We begin by proving that any tangential derivative of w is bounded. Since the problem is invariant under rotation and since u_2 is radial, we need only show that $\frac{\partial u_1}{\partial x_2}(r, 0, \dots, 0)$ remains bounded as $r \rightarrow 1^-$. Given $x = (x_1, x_2, x') \in B$ and $\theta > 0$ small, we denote by $x_\theta = (x_1 \cos \theta - x_2 \sin \theta, x_1 \sin \theta + x_2 \cos \theta, x')$ the image of x under the rotation of angle θ above the x_1 -axis in the (x_1, x_2) plane. By the rotation invariance of the Laplace operator, the function u_θ defined for $x \in B$ by $u_\theta(x) = u_1(x_\theta)$, solves (1). Using (3) and assumption (5), we deduce that $w_\theta = u_1 - u_\theta$ solves

$$(31) \quad \begin{cases} \Delta w_\theta + K w_\theta = \tilde{f}(u_1) - \tilde{f}(u_\theta) & \text{in } B, \\ w_\theta = 0 & \text{on } \partial B. \end{cases}$$

By the Maximum Principle on small domains, there exists $R_0 \in (0, 1)$ such that the operator $L = \Delta + K$ is coercive on $B \setminus B_{R_0}$. As a consequence, we claim that there exists a constant $C > 0$ such that for all $x \in B \setminus B_{R_0}$,

$$(32) \quad |w_\theta(x)| \leq C \sup_{\partial B_{R_0}} |w_\theta|.$$

Let indeed $\zeta > 0$ denote the solution of

$$\begin{cases} \Delta \zeta + K \zeta = 0 & \text{in } B \setminus B_{R_0} \\ \zeta = 1 & \text{on } \partial B_{R_0} \\ \zeta = 0 & \text{on } \partial B. \end{cases}$$

We shall prove that $z^\pm := w_\theta - \pm \sup_{\partial B_{R_0}} |w_\theta| \zeta$ are respectively non-positive and nonnegative, which implies that (32) holds for the constant $C = \|\zeta\|_\infty$. We work say with z^+ and assume by contradiction that the open set $\omega = \{x \in B \setminus B_{R_0} : z^+(x) > 0\}$ is non-empty. Restricting the analysis to a connected component, we have

$$\begin{cases} \Delta z^+ + Kz^+ = \tilde{f}(u_1) - \tilde{f}(u_\theta) \geq 0 & \text{in } \omega \\ z^+ \leq 0 & \text{on } \partial\omega. \end{cases}$$

By the Maximum Principle, we conclude that $z^+ \leq 0$ in ω , a contradiction. We have just proved (32). Since $u_1 \in C^1(\overline{B_{R_0}})$, we deduce that for some constant $C > 0$ and all $x \in B \setminus B_{R_0}$,

$$|w_\theta(x)| \leq C\theta.$$

Applying the above inequality at the point $x = (r, 0, \dots, 0)$, $r \in (R_0, 1)$ and letting $\theta \rightarrow 0$, we finally deduce that

$$\left| \frac{\partial u_1}{\partial x_2}(r, 0, \dots, 0) \right| \leq C \quad \text{for all } r \in (R_0, 1),$$

as desired.

Step 2 : estimate of the radial derivative

It remains to control $\partial w / \partial r$. Fix $R \in (0, 1)$. Let $G_R(x, y)$ denote Green's function in the ball of radius R . Then, for $x \in B_R$,

$$(33) \quad \begin{aligned} w(x) &= \int_{\partial B_R} \frac{\partial G_R}{\partial \nu_y}(x, \cdot) w \, d\sigma + \int_{B_R} G_R(x, \cdot) (f(u_1) - f(u_2)) \, dy \\ &=: w_1(x) + w_2(x). \end{aligned}$$

We want to let $R \rightarrow 1$ in the above identity. To do so, we first observe that w_1 is harmonic. By the Maximum Principle, $|w_1| \leq \|w\|_{L^\infty(\partial B_R)}$. By estimate (3), we conclude that $w_1 \rightarrow 0$ as $R \rightarrow 1$. To estimate w_2 , we need the following crucial estimate :

Lemma 4.1 *Assume (5). Then,*

$$\sup_{\theta \in S^{N-1}} \int_0^1 |f(u_1) - f(u_2)|(r, \theta) \, dr < +\infty.$$

We shall also need the following elementary estimates.

Lemma 4.2 *There exists a constant $C > 0$ such that for all $1/2 < r, R < 1$ and all $x, y \in B_R$,*

$$(34) \quad G_R(x, y) = R^{2-N} G_1\left(\frac{x}{R}, \frac{y}{R}\right)$$

$$(35) \quad \begin{cases} \int_{\partial B_r} G_1(x, \cdot) \, d\sigma \leq 1 \\ \int_{\partial B_r} \left| \frac{\partial G_1}{\partial |x|}(x, \cdot) \right| \, d\sigma \leq C \end{cases}$$

We postpone the proofs of the above two lemmas and return to (33). Using polar coordinates,

$$\begin{aligned} w_2(x) &= \int_{B_R} G_R(x, \cdot) (f(u_1) - f(u_2)) dy \\ &= \int_0^R \left(\int_{\partial B_r} G_R(x, \cdot) (f(u_1) - f(u_2)) d\sigma \right) dr \end{aligned}$$

By Lemmas 4.1 and 4.2, we may easily pass to the limit in the above expression as $R \rightarrow 1$, so

$$w(x) = \int_B G_1(x, \cdot) (f(u_1) - f(u_2)) dy$$

Using again Lemmas 4.1 and 4.2, we also have that w is differentiable in the $r = |x|$ variable and

$$\frac{\partial w}{\partial r}(x) = \int_B \frac{\partial G_1}{\partial |x|}(x, \cdot) (f(u_1) - f(u_2)) dy$$

Using polar coordinates again and Lemmas 4.1 and 4.2, we finally obtain

$$\begin{aligned} \left| \frac{\partial w}{\partial r} \right| &\leq C + \\ + \sup_{r \in (1/2, 1)} \left(\int_{\partial B_r} \left| \frac{\partial G_1}{\partial |x|}(x, \cdot) \right| d\sigma \right) \sup_{\theta \in S^{N-1}} \left(\int_0^1 |(f(u_1) - f(u_2))(r, \theta)| dr \right) \\ &\leq C. \end{aligned}$$

It only remains to prove Lemmas 4.1 and 4.2.

Proof of Lemma 4.1. We first deal with the case where u_1, u_2 are radial and $u_1 \geq u_2$. By assumption (5), we have

$$\int_0^1 |f(u_1) - f(u_2)| dr \leq \int_0^1 (\tilde{f}(u_1) - \tilde{f}(u_2)) dr + K \|u_1 - u_2\|_{L^\infty(B)}.$$

Using (3), we see that $u_1 - u_2$ is bounded and so it remains to estimate $\tilde{f}(u_1) - \tilde{f}(u_2)$. By (26), each u_i , $i = 1, 2$, solves

$$\frac{du_i/dr}{\sqrt{2(F(u_i) - g_i)}} = 1.$$

We also know by Lemma 3.4 that $g_i = o(F(u_i))$. So,

$$\lim_{r \rightarrow 1} \frac{du_i/dr}{\sqrt{2F(u_i)}} = 1.$$

Using this fact, as well as Lemma 3.5 and (4), we obtain for $R \in (1/2, 1)$,

$$\begin{aligned}
\int_0^R (\tilde{f}(u_1) - \tilde{f}(u_2)) dr &\leq \int_0^R (f(u_1) - f(u_2)) dr + K \|u_1 - u_2\|_{L^\infty(B)} \\
&\leq C \int_0^R (f(u_1) - f(u_2)) \frac{du_1/dr}{\sqrt{2F(u_1)}} dr + C \\
&\leq C \int_0^R \left(f(u_1) \frac{du_1/dr}{\sqrt{2F(u_1)}} - f(u_2) \frac{du_2/dr}{\sqrt{2F(u_2)}} \right) dr + \\
&\quad + C \int_0^R f(u_2) \left(\frac{du_2/dr}{\sqrt{2F(u_2)}} - \frac{du_1/dr}{\sqrt{2F(u_1)}} \right) dr + C \\
&\leq C \left(\sqrt{F(u_1)} - \sqrt{F(u_2)} \right) (R) + C + \\
&\quad + C \int_0^R f(u_2) \left(\frac{\sqrt{2(F(u_2) - g_2)}}{\sqrt{2F(u_2)}} - \frac{\sqrt{2(F(u_1) - g_1)}}{\sqrt{2F(u_1)}} \right) dr \\
&\leq C \left(\frac{F(u_1) - F(u_2)}{\sqrt{F(u_1)} + \sqrt{F(u_2)}} \right) (R) + C + \\
&\quad + C \int_0^R f(u_2) \frac{\sqrt{F(u_2)F(u_1) - g_2F(u_1)} - \sqrt{F(u_1)F(u_2) - g_1F(u_2)}}{\sqrt{F(u_2)F(u_1)}} dr \\
&\leq \frac{C}{\sqrt{F(u_1(R))}} \|F(u_1) - F(u_2)\|_{L^\infty(B)} + C + \\
&\quad + C \int_0^R \frac{f(u_2)}{\sqrt{F(u_2)F(u_1)}} \frac{g_1F(u_2) - g_2F(u_1)}{\sqrt{F(u_2)F(u_1)}} dr \\
&\leq C + C \int_0^R \frac{f(u_2)}{\sqrt{F(u_2)F(u_1)}} \frac{(g_1 - g_2)F(u_2) + g_2(F(u_2) - F(u_1))}{\sqrt{F(u_2)F(u_1)}} dr \\
&\leq C + C \|g_1 - g_2\|_{L^\infty(B)} \int_0^R \frac{f(u_2)}{F(u_2)} dr + \\
&\quad + C \|F(u_1) - F(u_2)\|_{L^\infty(B)} \int_0^R \frac{g_2}{F(u_2)} dr \\
&\leq C + C \int_{1/2}^R \frac{f(u_2)}{F(u_2)} \frac{du_2/dr}{\sqrt{2F(u_2)}} dr + C \\
&\leq C + C \left(F^{-1/2}(1/2) - F^{-1/2}(R) \right) \leq C.
\end{aligned}$$

This proves the lemma for radial solutions. To obtain the estimate in the general case, we may always assume that u_2 is the minimal solution of (1), so that $u_2 \leq u_1$ and u_2 is radial. By Lemma 2.4, up to replacing f by \tilde{f} given by Remark 2.1, there exists another radial solution V such that

$V \geq u_1 \geq u_2$. Using assumption (5), we have

$$\begin{aligned} \int_0^1 |f(u_1) - f(u_2)| \, dr &\leq \int_0^1 \left(\tilde{f}(u_1) - \tilde{f}(u_2) \right) \, dr + K \|u_1 - u_2\|_{L^\infty(B)} \\ &\leq \int_0^1 \left(\tilde{f}(V) - \tilde{f}(u_2) \right) \, dr + K \|u_1 - u_2\|_{L^\infty(B)} \\ &\leq \int_0^1 (f(V) - f(u_2)) \, dr + 2K \|u_1 - u_2\|_{L^\infty(B)} \end{aligned}$$

By estimate (3), $u_1 - u_2$ is bounded and the result follows from the radial case. \square

Proof of Lemma 4.2. (34) is standard : write the representation formula (33) both in B_R and in B_1 , change variables in the B_1 integral and identify the kernels. Next, we prove that given any $r \in (0, 1)$, $\int_{\partial B_r} G_1(x, \cdot) d\sigma \leq 1$. It suffices to show that for any $\phi \in C_c(0, 1)$,

$$(36) \quad \int_0^1 \phi(r) \left(\int_{\partial B_r} G_1(x, \cdot) d\sigma \right) \, dr \leq \|\phi\|_{L^1(0,1)}.$$

By definition of Green's function, the left-hand side of the above inequality is the function v solving

$$\begin{cases} -\Delta v = \phi & \text{in } B, \\ v = 0 & \text{on } \partial B. \end{cases}$$

The above equation can also be integrated directly :

$$v'(r) = r^{1-N} \int_0^r \phi(t) t^{N-1} \, dt,$$

whence $|v'| \leq \|\phi\|_{L^1(0,1)}$ and $|v| \leq \|\phi\|_{L^1(0,1)}$ i.e. (36) holds. This proves that $\int_{\partial B_r} G_1(x, \cdot) d\sigma \leq 1$.

We turn to the second estimate in (35). Recall that the Green's function in the unit ball is expressed for $x, y \in B$, $x \neq y$, by

$$(37) \quad G_1(x, y) = \Gamma \left((R^2 + r^2 - 2Rr \cos \varphi)^{1/2} \right) - \Gamma \left((1 + R^2 r^2 - 2Rr \cos \varphi)^{1/2} \right),$$

where $R = |x|$, $r = |y|$, φ is the angle formed by the vectors x and y and Γ is the fundamental solution of the Laplace operator. Differentiating with respect to R , we obtain for some $C_N > 0$,

$$(38) \quad C_N \frac{\partial G_1}{\partial |x|}(x, y) = \frac{R - r \cos \varphi}{(R^2 + r^2 - 2Rr \cos \varphi)^{N/2}} - \frac{Rr^2 - r \cos \varphi}{(1 + R^2 r^2 - 2Rr \cos \varphi)^{N/2}} = \frac{R - r + r(1 - \cos \varphi)}{((R - r)^2 + 2Rr(1 - \cos \varphi))^{N/2}} - \frac{Rr^2 - r + r(1 - \cos \varphi)}{((1 - Rr)^2 + 2Rr(1 - \cos \varphi))^{N/2}} = A - B$$

We estimate A and leave the reader perform similar calculations for B . Clearly, given $\varepsilon > 0$, the expression (38) remains uniformly bounded in the range $1/2 < R, r < 1, \varepsilon < \varphi < 2\pi - \varepsilon$. Hence,

$$\int_{\partial B_r} |A| d\sigma \leq C_\varepsilon + C \int_{\partial B_r \cap [0 < \varphi < \varepsilon]} |A| d\sigma.$$

For $y \in \partial B_r \cap [0 < \varphi < \varepsilon]$, let $z = z(y)$ denote the intersection of the line (Oy) and the hyperplane P passing through x and tangent to the hypersphere ∂B_R . Then, there exists constants $c_1, c_2 > 0$ such that for all $y \in \partial B_r \cap [0 < \varphi < \varepsilon]$,

$$c_1(1 - \cos \phi) \leq |z - x|^2 \leq c_2(1 - \cos \phi).$$

Hence, letting $B^{N-1}(x, \rho) \subset P$ denote the $N-1$ -dimensional ball of radius $\rho > 0$ centered at x , we obtain

$$\begin{aligned} \int_{\partial B_r} |A| d\sigma &\leq C \left(1 + \int_{B^{N-1}(x, R \sin \varepsilon)} \frac{|R-r| + Cr|z-x|^2}{(|R-r|^2 + c|z-x|^2)^{N/2}} dz \right) \\ &\leq C \left(1 + \int_{B^{N-1}(O, R\varepsilon)} \frac{|R-r| + C|z|^2}{(|R-r|^2 + c|z|^2)^{N/2}} dz \right) \\ &\leq C \left(1 + \int_{B^{N-1}(O, \frac{R\varepsilon}{|R-r|})} \frac{|R-r| + C|R-r|^2|z|^2}{|R-r|^N (1+c|z|^2)^{N/2}} |R-r|^{N-1} dz \right) \\ &\leq C \left(1 + \int_{\mathbb{R}^{N-1}} \frac{1}{(1+c|z|^2)^{N/2}} dz + |R-r| \int_{B^{N-1}(O, \frac{R\varepsilon}{|R-r|})} |z|^{2-N} dz \right) \\ &\leq C. \end{aligned}$$

Working similarly with the B term in (38), we finally obtain the desired estimate (35). \square

5 Asymptotics to all orders

This section is devoted to the

Proof of Theorem 1.7. Our first task consists in applying the Fixed Point Theorem to the functional \mathcal{N} defined for $v \in \mathcal{B}(v_0, \rho)$, $u \in I$ by

$$(39) \quad [\mathcal{N}(v)](u) = \sqrt{2 \left(F(u) - (N-1) \int_{U_0}^u \frac{v}{r} dt \right)},$$

where r is given by (9). Let us check first that $\mathcal{N}(\mathcal{B}(v_0, \rho)) \subset \mathcal{B}(v_0, \rho)$. Take $v \in \mathcal{B}(v_0, \rho)$. Then,

$$(40) \quad 1 \geq r \geq 1 - \frac{1}{1-\rho} \int_{U_0}^{+\infty} \frac{1}{v_0} dt = 1 - \frac{1}{1-\rho} \int_{U_0}^{+\infty} \frac{1}{\sqrt{2F}} dt.$$

By (2), it follows that for $\rho < 1/4$ and U_0 sufficiently large, $1 \geq r \geq 1/2$. Hence,

$$\left| \int_{U_0}^u \frac{v}{r} dt \right| \leq C \int_{U_0}^u \sqrt{2F} dt = o(F(u)),$$

where we used Lemma 3.4. So for U_0 large and $u \geq U_0$,

$$\left| \frac{N-1}{F(u)} \int_{U_0}^u \frac{v}{r} dt \right| \leq \rho.$$

We deduce that

$$(41) \quad \left| \frac{\mathcal{N}(v) - v_0}{v_0} \right| = 1 - \sqrt{1 - \frac{N-1}{F(u)} \int_{U_0}^u \frac{v}{r} dt} \leq \frac{1}{2} \left| \frac{N-1}{F(u)} \int_{U_0}^u \frac{v}{r} dt \right| < \rho.$$

Next, we prove that \mathcal{N} is contractive. Given $v_1, v_2 \in \mathcal{B}(v_0, \rho)$, let $r_1 = r(u, v_1), r_2 = r(u, v_2)$ (where r is given by (9)). Then, by estimate (40), $1/2 \leq r_1, r_2 \leq 1$ and

$$\begin{aligned} \left| \frac{\mathcal{N}(v_1) - \mathcal{N}(v_2)}{v_0} \right| &= \left| \sqrt{1 - \frac{N-1}{F(u)} \int_{U_0}^u \frac{v_1}{r_1} dt} - \sqrt{1 - \frac{N-1}{F(u)} \int_{U_0}^u \frac{v_2}{r_2} dt} \right| \\ &\leq C \frac{N-1}{F(u)} \int_{U_0}^u \left| \frac{v_1}{r_1} - \frac{v_2}{r_2} \right| dt \\ &\leq \frac{C}{F(u)} \left(\int_{U_0}^u |v_1 - v_2| dt + \int_{U_0}^u v_0 \left| \frac{1}{r_1} - \frac{1}{r_2} \right| dt \right) \\ &\leq \frac{C}{F(u)} \left(\rho \int_{U_0}^u \sqrt{2F} dt + \int_{U_0}^u \sqrt{2F} |r_1 - r_2| dt \right) \\ &\leq \frac{C}{F(u)} \left(\rho \int_{U_0}^u \sqrt{2F} dt + \int_{U_0}^u \sqrt{2F} \left| \int_t^{+\infty} \left(\frac{1}{v_1} - \frac{1}{v_2} \right) ds \right| dt \right) \\ &\leq \frac{C\rho}{F(u)} \left(\int_{U_0}^u \sqrt{2F} dt + \int_{U_0}^u \sqrt{2F} \left| \int_t^{+\infty} \frac{1}{\sqrt{2F}} ds \right| dt \right) \\ &\leq \frac{C\rho}{F(u)} \int_{U_0}^u \sqrt{2F} dt. \end{aligned}$$

Using Lemma 3.4, we conclude that \mathcal{N} is contractive in $\mathcal{B}(v_0, \rho)$ if U_0 was chosen large enough in the first place. We may thus apply the Fixed Point Theorem.

So, it only remains to prove (12). We first observe that the sequence (v_k) defined by (10) is asymptotic i.e. $v_{k+1}(u) = v_k(u)(1 + o(1))$, as $u \rightarrow +\infty$. Since $v_{k+1} = \mathcal{N}(v_k)$, it suffices to prove that $\mathcal{N}(v_0) - v_0 = o(v_0)$ and iterate. By (41),

$$\left| \frac{\mathcal{N}(v_0) - v_0}{v_0} \right| \leq \frac{C}{F(u)} \int_{U_0}^u \sqrt{2F} dt$$

and the claim follows by Lemma 3.4. So, the sequence (v_k) is asymptotic and so must be the sequence (u_k) defined by (11). We are now in a position to prove (12). By Theorem 1.1, we may restrict to the case where u is radially symmetric. Let $v = du/dr$. By (22), v solves

$$\frac{dv}{dr} + \frac{N-1}{r} v = f(u)$$

Use the change of variable $u = u(r)$ to get

$$v \frac{dv}{du} + \frac{N-1}{r} v = f(u).$$

Integrating, it follows that for some constant C

$$(42) \quad \frac{v^2}{2} = F(u) + C - \int_{U_0}^u \frac{N-1}{r} v \, dt.$$

Up to replacing $F(u)$ by $\tilde{F}(u) = F(u) + C$ (which is harmless from the point of view of asymptotics), we may assume $C = 0$. So it suffices to prove that $v \in \mathcal{B}(v_0, \rho)$ to conclude that v coincides with the unique fixed point of \mathcal{N} , whence (12) will follow. By (42) (with $C = 0$), $v \leq v_0$ and so

$$\begin{aligned} 0 \leq v_0 - v &\leq \sqrt{2F(u)} - \sqrt{2\left(F(u) - \int_{U_0}^u \frac{N-1}{r} v_0 \, dt\right)} \\ &\leq C \frac{\int_{U_0}^u \sqrt{2F} \, dt}{\sqrt{2F(u)}}. \end{aligned}$$

By Lemma 3.4, it follows that

$$0 \leq \frac{v_0 - v}{v_0} \leq v_0 - v < \rho$$

and $v \in B(v_0, \rho)$ as desired. \square

6 Universal blow-up rate

In this section, we prove Theorem 1.8, that is we characterize nonlinearities for which the blow-up rate is universal.

Proof of Theorem 1.8 .

Step 1. We begin by establishing the theorem when $\Omega = B$ is the unit ball. In light of Theorem 1.1, it suffices to prove (15) for one given solution u of (1), which we may therefore assume to be radial. By (26), we have after integration that

$$(43) \quad \int_u^{+\infty} \frac{1}{\sqrt{2(F(t) - g)}} \, dt = 1 - r.$$

By definition of u_0 , we also have

$$(44) \quad \int_{u_0}^{+\infty} \frac{1}{\sqrt{2F(t)}} \, dt = 1 - r.$$

Observe that $u \geq u_0$, split the integral in (44) as $\int_{u_0}^{+\infty} = \int_{u_0}^u + \int_u^{+\infty}$ and equate (43) and (44). It follows that

$$\begin{aligned} \int_{u_0}^u \frac{1}{\sqrt{2F(t)}} \, dt &= \int_u^{+\infty} \left(\frac{1}{\sqrt{2(F(t) - g)}} - \frac{1}{\sqrt{2F(t)}} \right) \, dt \\ &= \int_u^{+\infty} \frac{\sqrt{2F(t)} - \sqrt{2(F(t) - g)}}{\sqrt{2F(t)}\sqrt{2(F(t) - g)}} \, dt \\ &= \int_u^{+\infty} \frac{g}{\sqrt{2F(t)}\sqrt{2(F(t) - g)}(\sqrt{2F(t)} + \sqrt{2(F(t) - g)})} \, dt \end{aligned}$$

Recall that by Lemma 3.4, $g = o(F)$ as $t \rightarrow +\infty$ and $g(u) \sim (N-1)G(u) = (N-1) \int_0^u \sqrt{2F} dt$. So, for sufficiently large values of u , we deduce

$$(45) \quad \int_{u_0}^u \frac{1}{\sqrt{2F(t)}} dt \leq C \int_u^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F(t))^{3/2}} dt.$$

Since F is nondecreasing, it follows that

$$0 \leq \frac{u - u_0}{\sqrt{2F(u)}} \leq \int_{u_0}^u \frac{1}{\sqrt{2F(t)}} dt \leq C \int_u^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F(t))^{3/2}} dt.$$

Hence,

$$0 \leq u - u_0 \leq C \sqrt{2F(u)} \int_u^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F(t))^{3/2}} dt,$$

and (15) follows from (14).

Step 2. Next, we prove that (15) holds for general domains Ω . To this end, we combine a standard approximation argument by inner and outer spheres (see e.g. [66, 6]) and the comparison technique of [33, 3]. Let u denote a solution of (13) and take a point $x_0 \in \partial\Omega$. Let $B \subset \Omega$ denote a ball which is tangent to $\partial\Omega$ at x_0 . Shrink B somewhat by letting $B_\varepsilon = (1 - \varepsilon)B$, $\varepsilon > 0$. Observe that $u \in C(\overline{B_\varepsilon})$ is a subsolution of

$$(46) \quad \begin{cases} \Delta U = f(U) & \text{in } B_\varepsilon, \\ U = +\infty & \text{on } \partial B_\varepsilon, \end{cases}$$

By Lemma 2.3, there exists a solution V_ε of (46), such that $V_\varepsilon \geq u$ in B_ε . Furthermore, V_ε can be chosen to be the minimal solution of (46) such that $V_\varepsilon \geq u$ in B_ε . In particular, V_ε is radial and $\varepsilon \rightarrow V_\varepsilon$ is nondecreasing. In addition, $\varepsilon \rightarrow V_\varepsilon$ is uniformly bounded on compact subsets of B (working as in the proof of Lemma 2.3), so V_ε converges as $\varepsilon \rightarrow 0$, to a solution V of (1) such that $V \geq u$ in B . By Step 1,

$$\lim_{\substack{x \rightarrow x_0 \\ x \in B}} V(x) - u_0(1 - d_B(x)) = 0,$$

where d_B denotes the distance to ∂B . Since $V \geq u$ and since the above discussion is valid for any point $x_0 \in \partial\Omega$, we finally obtain

$$(47) \quad \limsup_{x \rightarrow \partial\Omega} [u(x) - u_0(1 - d(x))] \leq 0,$$

where $d(x)$ is the distance to $\partial\Omega$. Choose now an exterior ball $B \subset \mathbb{R}^N \setminus \overline{\Omega}$ which is tangent to $\partial\Omega$ at x_0 . For $\varepsilon > 0$ small and $R > 0$ large, the annulus $A_\varepsilon = RB \setminus (1 - \varepsilon)B$ contains Ω . Let U_ε denote a large solution on A_ε , which we may assume to be minimal, radial and bounded above on Ω by u . Again $U_\varepsilon \rightarrow U$ as $\varepsilon \rightarrow 0$ where U is a radial large solution in $A = RB \setminus B \supset \Omega$. Repeating the analysis of Step 1. (which was purely local) for the case of a radial solution defined on an annulus rather than a ball, we easily deduce that

$$\lim_{\substack{x \rightarrow x_0 \\ x \in B}} U(x) - u_0(1 - d_B(x)) = 0.$$

Since $u \geq U$ and since the above discussion is valid for any point $x_0 \in \partial\Omega$, we obtain

$$(48) \quad \liminf_{x \rightarrow \partial\Omega} [u(x) - u_0(1 - d(x))] \geq 0.$$

So, by (48) and (47), we have that (15) holds in any smoothly bounded domain Ω .

Step 3. It only remains to prove that (15) fails when (16) holds. We use Theorem 1.7 to compute the second term in the asymptotic expansion of a solution. By (10),

$$\begin{aligned} v_1(u) &= \sqrt{2 \left(F(u) - (N-1) \int_0^u \sqrt{2F} dt (1 + o(1)) \right)} \\ &= \sqrt{2F(u)} \left(1 - (N-1) \frac{\int_0^u \sqrt{2F} dt}{2F(u)} (1 + o(1)) \right), \end{aligned}$$

whence

$$\begin{aligned} \frac{1}{v_1} &= \frac{1}{\sqrt{2F(u)}} \left(1 + (N-1) \frac{\int_0^u \sqrt{2F} dt}{2F(u)} (1 + o(1)) \right) \\ &= \frac{1}{\sqrt{2F(u)}} + (N-1) \frac{\int_0^u \sqrt{2F} dt}{(2F(u))^{3/2}} (1 + o(1)). \end{aligned}$$

Integrating (11) for $k = 1$, it follows that for r close enough to 1,

$$(49) \quad \int_{u_1}^{+\infty} \frac{dt}{\sqrt{2F}} + (N-1)(1 + o(1)) \int_{u_1}^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt = 1 - r.$$

Recall (44), split the integral in (44) as $\int_{u_0}^{+\infty} = \int_{u_0}^{u_1} + \int_{u_1}^{+\infty}$ and equate (49) and (44) to get

$$\int_{u_0}^{u_1} \frac{dt}{\sqrt{2F}} = (N-1)(1 + o(1)) \int_{u_1}^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt.$$

Since F is nondecreasing, we deduce that

$$(50) \quad \frac{u_1 - u_0}{\sqrt{2F(u_0)}} \geq (N-1)(1 + o(1)) \int_{u_1}^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt.$$

Note also that

$$(51) \quad \int_{u_0}^{u_1} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt \leq \int_{u_0}^{u_1} \frac{t}{2F} dt \leq \frac{(u_1 - u_0)^2}{4F(u_0)}.$$

Assume by contradiction that $\lim_{r \rightarrow 1^-} (u_1 - u_0)(r) = 0$. Then, (51) implies that

$$\int_{u_0}^{u_1} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt = o\left(\frac{u_1 - u_0}{\sqrt{2F(u_0)}}\right).$$

Using this information in (50), we obtain that

$$\frac{u_1 - u_0}{\sqrt{2F(u_0)}} \geq (N-1)(1+o(1)) \int_{u_0}^{+\infty} \frac{\int_0^t \sqrt{2F} ds}{(2F)^{3/2}} dt.$$

But (16) would then lead us to a contradiction with the assumption $\lim_{r \rightarrow 1^-} (u_1 - u_0)(r) = 0$. So we must have

$$\liminf_{r \rightarrow 1^-} (u_1 - u_0) > 0.$$

and so (15) fails. \square

7 The first three singular terms

In the previous section, we characterized nonlinearities for which only one term in the expansion is singular. In the present section, we calculate implicitly the next two terms in the expansion. We have not tried to characterize those f for which all remaining terms are nonsingular, but this can certainly be achieved. We leave the tenacious reader try her/his hand at this computational problem.

We begin by calculating the leading asymptotics of v_1, v_2 . By (10), we have

$$\frac{v_1^2}{2} = F - (N-1) \int^u \sqrt{2F} (1+o(1)) dt.$$

So,

$$\begin{aligned} \frac{v_1}{\sqrt{2}} &= \sqrt{F} \left(1 - (N-1) \frac{\int^u \sqrt{2F} dt}{F} (1+o(1)) \right)^{1/2} = \\ &= \sqrt{F} - \frac{N-1}{2} \frac{\int^u \sqrt{2F} dt}{\sqrt{F}} (1+o(1)). \end{aligned}$$

In other words,

$$v_1 = \sqrt{2F} - (N-1) \frac{\int^u \sqrt{2F} dt}{\sqrt{2F}} (1+o(1)).$$

To calculate v_2 , we introduce some notation. Given a positive measurable function v , set

$$Pv = \int^u v dt, \quad Qv = \frac{Pv}{v}, \quad Rv = \int_u^{+\infty} \frac{dt}{v} \quad \text{and} \quad Tv = (N-1)PQv + P(vRv).$$

v_1 is then expressed by

$$v_1 = v_0 - (N-1)(1+o(1))Qv_0,$$

while v_2 is given by

$$\begin{aligned} \frac{v_2^2}{2} &= F - (N-1) \int^u \frac{v_1}{1 - \int_t^{+\infty} \frac{ds}{v_0} (1+o(1))} dt = \\ &= F - (N-1) \int^u (v_0 - (N-1)Qv_0 + o(Q(v_0)))(1 - Rv_0 + o(Rv_0)) dt = \\ &= F - (N-1)Pv_0 + (N-1)Tv_0(1+o(1)). \end{aligned}$$

So,

$$\begin{aligned} v_2 &= (2F - 2(N-1)Pv_0 + 2(N-1)Tv_0(1+o(1)))^{1/2} = \\ &= v_0 \left(1 - (N-1)\frac{Pv_0}{F} + (N-1)\frac{Tv_0}{F}(1+o(1)) \right)^{1/2} = \\ &= v_0 \left(1 - \frac{N-1}{2}\frac{Pv_0}{F} + \frac{N-1}{2}\frac{Tv_0}{F} - \frac{3}{8}(N-1)^2 \left(\frac{Pv_0}{F} \right)^2 + o(Tv_0/F + (Pv_0/F)^2) \right). \end{aligned}$$

And so,

$$\begin{aligned} \frac{1}{v_2} &= \frac{1}{v_0} \left(1 + \frac{N-1}{2}\frac{Pv_0}{F} - \frac{N-1}{2}\frac{Tv_0}{F} + \frac{5}{8}(N-1)^2 \left(\frac{Pv_0}{F} \right)^2 + o(Tv_0/F + (Pv_0/F)^2) \right) = \\ &= \frac{1}{v_0} + (N-1)\frac{Pv_0}{v_0^3} + (N-1) \left(-\frac{Tv_0}{v_0^3} + \frac{5}{4}(N-1)\frac{(Pv_0)^2}{v_0^5} \right) (1+o(1)) = \\ &= \frac{1}{\sqrt{2F}} + (N-1)\frac{\int^u \sqrt{2F} dt}{(2F)^{3/2}} + \\ &+ \frac{(N-1)}{(2F)^{3/2}} \left(-\int^u \left((N-1)\frac{\int^t \sqrt{2F} ds}{\sqrt{2F}} + \sqrt{2F} \int_u^{+\infty} \frac{ds}{\sqrt{2F}} \right) dt + \frac{5(N-1)}{4} \frac{\left(\int^u \sqrt{2F} dt \right)^2}{2F} \right) (1+o(1)) \end{aligned}$$

Integrating once more, we finally obtain

$$\begin{aligned} 1-r &= \int_{u_2(r)}^{+\infty} \frac{du}{\sqrt{2F}} + (N-1) \int_{u_2(r)}^{+\infty} \frac{\int^u \sqrt{2F} dt}{(2F)^{3/2}} du + (1+o(1)) \times \\ &\times (N-1) \int_{u_2(r)}^{+\infty} \left(-\int^u \left((N-1)\frac{\int^t \sqrt{2F} ds}{\sqrt{2F}} + \sqrt{2F} \int_u^{+\infty} \frac{ds}{\sqrt{2F}} \right) dt + \frac{5(N-1)}{4} \frac{\left(\int^u \sqrt{2F} dt \right)^2}{2F} \right) \frac{du}{(2F)^{3/2}}. \end{aligned}$$

This proves Proposition 1.12.

8 An example: $f(u) = u^p$, $p > 1$

Finding the n -th term in the expansion for arbitrary $n \in \mathbb{N}$ is out of reach for general f , simply because of the algorithmic complexity of calculations. However, when additional information on f is available, one can guess the general form of the expansion and then try to establish it. This is precisely what we do in this section, with the nonlinearity $f(u) = u^p$, $p > 1$.

For notational convenience, we shall work with $F(u) = \frac{1}{2}u^{2q}$, where $2q-1 = p$, which simply amounts to working with a constant multiple of the original solution.

Recall (10) and (11). We want to prove inductively that there exists

numbers a_k, b_k depending on k, p, N only such that

$$(52) \quad v_n = u^q \sum_{k=0}^n b_k u^{-k(q-1)} + o(u^{q-n(q-1)}),$$

$$(53) \quad u_n = d^{-\frac{1}{q-1}} \sum_{k=0}^n a_k d^k + E_n(d^{-\frac{1}{q-1}+n+1}),$$

where $E_n(d^{-\frac{1}{q-1}+n+1}) \sim e_n d^{-\frac{1}{q-1}+n+1}$ for some $e_n \in \mathbb{R}$, as $d \rightarrow 0^+$. We have $v_0 = \sqrt{2F} = u^q$. Solving for u_0 in (11) yields $u_0 = c d^{-\frac{1}{q-1}}$. So, (53) and (52) hold for $n = 0$. Suppose now the result is true for a given $n \in \mathbb{N}$. In the computations below, the letter c_k denotes a number depending on k, p, N only, which value may change from line to line. By (52), we have

$$\begin{aligned} \frac{1}{v_n} &= u^{-q} \left(1 + \sum_{k=1}^n b_k u^{-k(q-1)} + o(u^{-n(q-1)}) \right)^{-1} = \\ &= u^{-q} \left(\sum_{k=0}^n c_k u^{-k(q-1)} + o(u^{-n(q-1)}) \right). \end{aligned}$$

So,

$$(54) \quad \int_t^{+\infty} \frac{ds}{v_n} = t^{1-q} \left(\sum_{k=0}^n c_k t^{-k(q-1)} \right) + o(t^{-(n+1)(q-1)}) =$$

$$(55) \quad = \sum_{k=1}^{n+1} c_k t^{-k(q-1)} + o(t^{-(n+1)(q-1)}).$$

It follows that

$$\frac{1}{1 - \int_t^{+\infty} \frac{ds}{v_n}} = \sum_{k=0}^{n+1} c_k t^{-k(q-1)} + o(t^{-(n+1)(q-1)}).$$

Whence,

$$\frac{v_n(t)}{1 - \int_t^{+\infty} \frac{ds}{v_n}} = t^q \sum_{k=0}^n c_k t^{-k(q-1)} + o(t^{q-n(q-1)}).$$

And so,

$$\begin{aligned} v_{n+1} &= \sqrt{2F - (N-1) \int^u \frac{v_n}{1 - \int_t^{+\infty} \frac{ds}{v_n}} dt} = \\ &= \sqrt{u^{2q} + u^{q+1} \sum_{k=0}^n c_k u^{-k(q-1)} + o(u^{1+q-n(q-1)})} = \\ &= u^q \left(1 + \sum_{k=1}^{n+1} c_k u^{-k(q-1)} + o(u^{-(n+1)(q-1)}) \right)^{1/2} = \\ &= u^q \sum_{k=0}^{n+1} c_k u^{-k(q-1)} + o(u^{q-(n+1)(q-1)}). \end{aligned}$$

This proves (52). Integrating (11), we obtain

$$(56) \quad \int_{u_n}^{+\infty} \frac{du}{v_n} = d.$$

Now, $v_{n+1} = v_n + c_{n+1}u^{q-(n+1)(q-1)}(1 + o(1))$. So,

$$\frac{1}{v_{n+1}} = \frac{1}{v_n} + c_{n+1}u^{-q-(n+1)(q-1)}(1 + o(1)).$$

It follows that

$$d = \int_{u_{n+1}}^{+\infty} \frac{du}{v_{n+1}} = \int_{u_{n+1}}^{+\infty} \frac{du}{v_n} + c_{n+1}u_{n+1}^{-(n+2)(q-1)}(1 + o(1)).$$

In addition, $v_n \sim v_0$, so $u_n \sim u_0$, and so $u_{n+1}^{-(q-1)} \sim d$. Using this in the above equation, we get

$$d + c_{n+1}d^{n+2}(1 + o(1)) = \int_{u_{n+1}}^{+\infty} \frac{du}{v_n}.$$

Recalling that v_n is defined by (56) and satisfies (53) by induction hypothesis, we conclude that

$$v_{n+1} = (d + c_{n+1}d^{n+2}(1 + o(1)))^{-\frac{1}{q-1}} \sum_{k=0}^n a_k (d + c_{n+1}d^{n+2}(1 + o(1)))^k + E_n(d^{-\frac{1}{q-1}+n+1}).$$

Expanding again the above expression, we finally obtain

$$v_{n+1} = d^{-\frac{1}{q-1}} \sum_{k=0}^{n+1} a_k d^k + E_{n+1}(d^{-\frac{1}{q-1}+n+2}),$$

which proves (53). Proposition 1.13 follows.

References

- [1]]cite.bundle1Catherine Bandle, *Asymptotic behavior of large solutions of elliptic equations*, An. Univ. Craiova Ser. Mat. Inform. **32** (2005), 1–8. MR **2215890** (2006j:35066)
- [2]]cite.bundle-marcus2Catherine Bandle and Moshe Marcus, *On second-order effects in the boundary behaviour of large solutions of semilinear elliptic problems*, Differential Integral Equations **11** (1998), no. 1, 23–34. MR **1607972** (98m:35049)
- [3]]cite.ddgr3Serge Dumont, Louis Dupaigne, Olivier Goubet, and Vicentiu Rădulescu, *Back to the Keller-Osserman condition for boundary blow-up solutions*, Adv. Nonlinear Stud. **7** (2007), no. 2, 271–298. MR **2308040** (2008e:35062)
- [4]]cite.gnn4B. Gidas, Wei Ming Ni, and L. Nirenberg, *Symmetry and related properties via the maximum principle*, Comm. Math. Phys. **68** (1979), no. 3, 209–243. MR **544879** (80h:35043)
- [5]]cite.keller5J. B. Keller, *On solutions of $\Delta u = f(u)$* , Comm. Pure Appl. Math. **10** (1957), 503–510. MR 0091407 (19,964c)

- [6]]cite.lm6A. C. Lazer and P. J. McKenna, *Asymptotic behavior of solutions of boundary blowup problems*, *Differential Integral Equations* **7** (1994), no. 3-4, 1001–1019. MR **1270115** (**95c**:35084)
- [7]]cite.osserman7Robert Osserman, *On the inequality $\Delta u \geq f(u)$* , *Pacific J. Math.* **7** (1957), 1641–1647. MR 0098239 (20 #4701)
- [8]]cite.pv8Alessio Porretta and Laurent Véron, *Symmetry of large solutions of nonlinear elliptic equations in a ball*, *J. Funct. Anal.* **236** (2006), no. 2, 581–591. MR **2240175** (**2007b**:35131)