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Bifurcation delay - the case of the sequence: stable focus - unstable focus - unstable node

Eric Benoit*

January 19, 2009

Abstract

Let us give a two dimensional family of real vector fields. We suppose that there exists a stationary point where the linearized vector field has successively a stable focus, an unstable focus and an unstable node. When the parameter moves slowly, a bifurcation delay appears due to the Hopf bifurcation. The studied question in this article is the continuation of the delay after the focus-node bifurcation.

AMS classification: 34D15, 34E15, 34E18, 34E20, 34M60

Keywords: Hopf-bifurcation, bifurcation-delay, slow-fast, canard, Airy, relief.

1 Introduction

"Singular perturbations" is a studied domain from many years ago. Since 1980, many contributions were written because new tools were applied to the subject. The main studied objects are the *slow fast vector fields* also known as *systems with two time-scales*. We will give the problem here with a more particular point of view: the *bifurcation delay*, as in articles [8, 2, 9, 7]. We write the studied system: $\varepsilon \dot{X} = f(t, X, \varepsilon)$, where ε is a real positive parameter which tends to zero. For a better understanding of the expression *dynamic bifurcation* it is better to write the system after a rescaling of the variable:

$$\begin{cases} \dot{X} &= f(a, X, \varepsilon) \\ \dot{a} &= \varepsilon \end{cases}$$

where a is a "slowly varying" parameter.

The main objects in this study are the eigenvalues of the linear part of equation $\dot{X} = f(a, X, 0)$ near the quasi-stationary point. Indeed, they give a characterization of the stability of the equilibrium of the fast vector field at this point. The aim of this study is to understand what happens when the stability of a quasi-stationary point changes. A bifurcation occurs when at least one of the eigenvalues has a null real part.

In this article we restrict our study to two-dimensional real systems. In this situation, the generic bifurcations are: the saddle-node bifurcation, the Hopf bifurcation and the focus-node bifurcation.

The saddle-node bifurcation is solved by the *turning point* theory: when the real part of one of the eigenvalue becomes positive, there is no delay and a trajectory of the systems leaves the neighborhood of the quasi-stationary point when it reaches the bifurcation. For this study, the study of one-dimensional systems is sufficient: we have a decomposition of the phase space where only the one-dimensional factor is interesting. There exist many articles on this subject, we will be interested particularly by [3] where the method of *relief* is used. The article [6] introduces the geometrical methods of *Fenichel's manifold*.

The Hopf delayed bifurcation is well explained in [10], we will upgrade the results in paragraph 2 below.

In a focus-node bifurcation, the stability of the quasi stationary point does not change, then, locally, there is no problem of canards or bifurcation delays. Indeed, when there is a bifurcation delay at a Hopf-bifurcation point, it is possible to evaluate the value of the delay, and the main question is to understand the influence of the focus node bifurcation to this delay.

In paragraph 2, the Hopf bifurcation alone is studied, as well as the focus-node bifurcation following a Hopf bifurcation in paragraph 3.

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In paragraphs 2.1 and 3.1, we assume that there exists a solution of the system approximated by the quasi steady state in the whole domain, so this trajectory has an infinite delay. The used methods are real, and the system has to be smooth (actually only C^2). In paragraphs 2.2 et 3.2, we avoid this very special hypothesis. It is here supposed that the system is analytic, and we study the solutions on complex domains. Unfortunately, I have not a proof for the main result of this article. But it seems to me that the problem is interesting, and the results are argmented.

We use Nelson's nonstandard terminology (see for example [5]). Indeed, almost all sentences can be translated in classical terms, where ε is considered as a variable and not as a parameter. Often, the translation is given on footnotes.

2 The delayed Hopf bifurcation

The problem is studied and essentially resolved in [10]. We give here the proofs to improve the results and to fix the ideas for the main paragraph of the article. The main tool is the *relief's* theory of J.L. Callot, explained in [4].

The studied equation is

$$\varepsilon \dot{X} = f(t, X, \varepsilon) \quad (1)$$

where f is analytic on a domain \mathcal{D} of $\mathbb{C} \times \mathbb{C}^2 \times \mathbb{C}$.

Hypothesis and notations

- H1 The function f is analytic. It takes real values when the arguments are real.
- H2 The parameter ε is real, positive, infinitesimal¹.
- H3 There exists an analytic function ϕ , defined on a complex domain \mathcal{D}_t so that $f(t, \phi(t), 0) = 0$. The curve $X = \phi(t)$ is called the *slow curve* of equation (1). We assume that the intersection of \mathcal{D}_t with the real axis is an interval $]t_m, t_M[$.
- H4 Let us denote $\lambda(t)$ and $\mu(t)$ for the eigenvalues of the jacobian matrix $D_X f$, computed at point $(t, \phi(t), 0)$. We assume that, for t real, the signs of the real and imaginary parts are given by the table below :

t	t_m	a	t_M
$\Re(\lambda(t))$	-	0	+
$\Re(\mu(t))$	-	0	+
$\Im(\lambda(t))$	-	-	-
$\Im(\mu(t))$	+	+	+

Then, when t increases from t_m to t_M , the quasi-steady state is first an attractive focus, then a repulsive focus, with a Hopf bifurcation at $t = a$.

2.1 Input-output function when there exists a big canard

In this section, we assume that there exists a *big canard* $\tilde{X}(t)$ i.e. a solution of equation (1) such that² $\tilde{X}(t) \simeq \phi(t)$ for all t in the S -interior of $]t_m, t_M[$. We now want to study the others solutions of equation (1) by comparison with \tilde{X} .

The main tool for that is a sequence of change of unknown: first, we perform a translation on X , depending on t to put the big canard on the axis:

$$X = \tilde{X}(t) + Y$$

It gives the system

$$\varepsilon \dot{Y} = g(t, Y, \varepsilon) \quad \text{with} \quad g(t, Y, \varepsilon) = f(t, \tilde{X}(t) + Y, \varepsilon) - f(t, \tilde{X}(t), \varepsilon)$$

¹In classical terms, we assume that ε leaves in a small complex sector: $|\varepsilon|$ bounded and $\arg(\varepsilon) \in]-\delta, \delta[$.

²Without nonstandard terminology, a *big canard* is a solution of equation (1) depending on the parameter ε such that

$$\forall t \in]t_m, t_M[, \lim_{\varepsilon > 0, \varepsilon \rightarrow 0} \tilde{X}(t, \varepsilon) = \phi(t)$$

The matrix $D_X f(t, \phi(t), 0)$ has two complex conjugate distinct eigenvalues (see hypothesis H4), then there exists a linear transformation $P(t)$ which transforms the jacobian matrix in a canonical form. We define the change of unknown

$$Y = P(t)Z$$

The new system, has the following form (we wrote only the interesting terms):

$$\varepsilon \dot{Z} = h(t, Z, \varepsilon) \quad \text{with} \quad h(t, Z, \varepsilon) = \begin{pmatrix} \alpha(t) & -\omega(t) \\ \omega(t) & \alpha(t) \end{pmatrix} Z + O(\varepsilon)Z + O(Z^2) \quad , \quad \lambda(t) = \alpha(t) - i\omega(t)$$

The next change is given by the polar coordinates:

$$Z = \begin{pmatrix} r \cos \theta \\ r \sin \theta \end{pmatrix}$$

$$\begin{cases} \varepsilon \dot{r} &= r(\alpha(t) + O(\varepsilon) + O(r)) \\ \varepsilon \dot{\theta} &= \omega(t) + O(\varepsilon) + O(r) \end{cases}$$

The last one is an exponential microscope³:

$$r = \exp\left(\frac{\rho}{\varepsilon}\right)$$

$$\begin{cases} \dot{\rho} &= \alpha(t) + O(\varepsilon) + e^{\frac{\rho}{\varepsilon}} k_1(r, \theta, \varepsilon) \\ \varepsilon \dot{\theta} &= \omega(t) + O(\varepsilon) + e^{\frac{\rho}{\varepsilon}} k_2(r, \theta, \varepsilon) \end{cases} \quad (2)$$

While ρ is non positive and non infinitesimal, r is exponentially small and the equation (2) gives a good approximation of ρ with $\dot{\rho} = \alpha$. When ρ becomes infinitesimal, with a more subtle argument (see [1]) using differential inequations, we can prove that r becomes non infinitesimal. This gives the proposition below:

Proposition 1 *Let us assume hypothesis H1 to H4 (Hopf bifurcation) for equation (1). However, we assume that there exists a canard $\tilde{X}(t)$ going along⁴ the slow curve at least on $]t_m, t_M[$. Then if $X(t)$ goes along the slow curve exactly⁵ on $]t_e, t_s[$ with $[t_e, t_s] \subset]t_m, t_M[$, then*

$$\int_{t_e}^{t_s} \Re(\lambda(\tau)) d\tau = 0$$

The input-output relation (between t_e and t_s) is defined by $\int_{t_e}^{t_s} \Re(\lambda(\tau)) d\tau = 0$. It is described by its graph (see figure 1). In this case, this relation is a function.

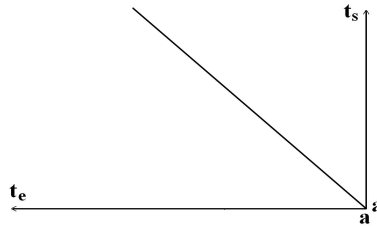


Figure 1: The input-output relation for equation (3) when there exists a big canard.

³All the preceding transformations were regular with respect to ε . This last one is singular at $\varepsilon = 0$.

⁴A solution $\tilde{X}(t, \varepsilon)$ goes along the slow curve at least on $]t_1, t_2[$ if

$$\forall t \in]t_1, t_2[, \lim_{\varepsilon > 0, \varepsilon \rightarrow 0} \tilde{X}(t, \varepsilon) = \phi(t)$$

⁵A solution $\tilde{X}(t, \varepsilon)$ goes along the slow curve exactly on $]t_e, t_s[$ if it goes along the slow curve at least on $]t_e, t_s[$, and if the interval $]t_e, t_s[$ is maximal for this property.

2.2 The bump and the anti-bump

In this paragraph, t becomes complex, in the domain \mathcal{D}_t . We assume that for all t in \mathcal{D}_t , the two eigenvalues $\lambda(t)$ and $\mu(t)$ are distinct. It is a necessary condition to apply Callot's theory of reliefs.

We define the reliefs R_λ and R_μ by:

$$\begin{aligned} F_\lambda(t) &= \int_a^t \lambda(t) dt & , & & R_\lambda(t) &= \Re(F_\lambda(t)) \\ F_\mu(t) &= \int_a^t \mu(t) dt & , & & R_\mu(t) &= \Re(F_\mu(t)) \end{aligned}$$

It is easy to see that $\lambda(\bar{t}) = \overline{\mu(t)}$, and $F_\lambda(\bar{t}) = \overline{F_\mu(t)}$, then $R_\lambda(\bar{t}) = R_\mu(t)$. The two functions R_λ and R_μ coincide on the real axis. We will denote $R(t)$.

Definition 1 We say that a path $\gamma : s \in [0, 1] \mapsto \mathcal{D}_t$ goes down the relief R_λ if and only if $\frac{d}{ds} R_\lambda(\gamma(s)) < 0$ for all s in $[0, 1]$.

Definition 2 Let us give a point t_e such that $(t_e, \phi(t_e), 0) \in \mathcal{D}$. We say that \mathcal{D}_t is a domain below t_e if and only if for all t in the S -interior of \mathcal{D}_t , there exist two paths in \mathcal{D}_t , from t_e to t , the first one goes down the relief R_λ and the second one down R_μ .

Theorem 2 (Callot) Let us assume that \mathcal{D}_t is a domain below t_e . A solution $X(t)$ of equation (1) with an initial condition $X(t_e)$ infinitesimally close to $\phi(t_e)$ is defined at least on the S -interior of \mathcal{D}_t where it is infinitesimally close to $\phi(t)$.

Let us apply this theorem to the following example, chosen as the typical example satisfying hypothesis H1 to H4 (Hopf bifurcation).

$$\begin{cases} \varepsilon \dot{x} &= tx + y + \varepsilon c_1 \\ \varepsilon \dot{y} &= -x + ty + \varepsilon c_2 \end{cases} \quad (3)$$

The eigenvalues are $\lambda = t - i$ et $\mu = t + i$. The level curves of the two reliefs $R_\lambda(t) = \frac{1}{2}(t - i)^2$ and $R_\mu(t) = \frac{1}{2}(t + i)^2$ are drawn on figure 2.

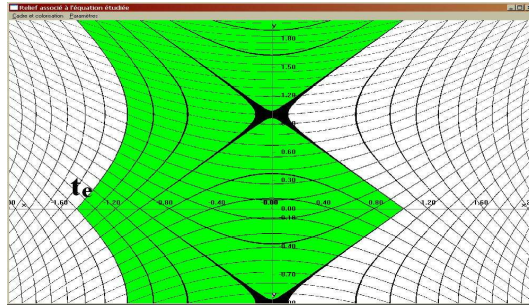


Figure 2: The level curves of the two reliefs of equation (3), and a domain below t_e

Generically¹ there is no surstability at point $t = i$ (see [3] for the definition of surstability). Consequently, we have the following results for all equations such that the reliefs are on the same type as those of figure 2:

Definition 3 Let us give t_c a point² where the eigenvalue vanishes: $\lambda(t_c) = 0$. The value of the relief at point t_c is a critical value of the relief R_λ . The bump³ is the real number t^* bigger than a , minimal such that $R_\lambda(t^*)$ is a critical value. The anti-bump is the real number t^{**} smaller than a , maximal such that $R_\lambda(t^{**})$ is a critical value.

¹I do not know the exact generic hypothesis. We have to combine the constraints given by the surstability theory of [3] and the fact that the equation (1) is real

²In some cases, it is possible that t_c is infinite. For $\varepsilon \dot{X} = \begin{pmatrix} \sin t & \cos t \\ -\cos t & \sin t \end{pmatrix} X + O(X^2) + O(\varepsilon)$ we have $t_c = +i\infty$.

³The name "bump" is a translation of the french name "butée"

For equation (3), the bump is $t^* = 1$ and the anti-bump $t^{**} = -1$.

Theorem 3 *A trajectory of equation (1) can go along the slow curve $X = \phi(t)$ exactly on $]t_e, t_s[$ if and only if one of the following is verified:*

$$\begin{aligned} t_e < t^{**} \quad \text{and} \quad t_s = t^* \\ t_e = t^{**} \quad \text{and} \quad t_s > t^* \\ t^{**} < t_e < a \quad \text{and} \quad a < t_s < t^* \quad \text{and} \quad R(t_e) = R(t_s) \end{aligned}$$

This theorem is illustrated by the graph of the input-output relation, drawn on figure 3.

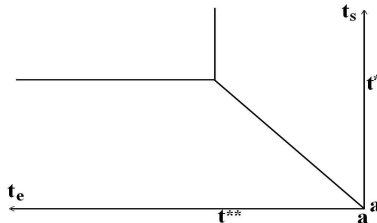


Figure 3: The input-output relation for equation (3)

3 Delayed Hopf bifurcation followed by a focus-node bifurcation

The studied equation is

$$\varepsilon \dot{X} = f(t, X, \varepsilon) \quad (4)$$

where f is analytic on a domain \mathcal{D} of $\mathbb{C} \times \mathbb{C}^2 \times \mathbb{C}$, and satisfies the following hypothesis:

Hypothesis and notations

HFN1 The analytic function f takes real values when the arguments are real.

HFN2 The parameter ε is real, positive, infinitesimal.

HFN3 There exists an analytic function ϕ , defined on a complex domain \mathcal{D}_t such that $f(t, \phi(t), 0) = 0$. The curve $X = \phi(t)$ is called the *slow curve* of equation 1. We assume that the intersection of \mathcal{D}_t with the real axis is an interval $]t_m, t_M[$.

HFN4 Let us denote $\lambda(t)$ and $\mu(t)$ for the eigenvalues of the jacobian matrix $D_X f$, computed at point $(t, \phi(t), 0)$. We assume that, for t real, the signs of the real and imaginary parts are given by the table below :

t	t_m	a	b	t_M
$\Re(\lambda(t))$	-	0	+	+
$\Re(\mu(t))$	-	0	+	+
$\Im(\lambda(t))$	-	-	0	0
$\Im(\mu(t))$	+	+	0	0

Then, when t increases on the real interval $]t_m, t_M[$, we have successively an attractive focus, a Hopf bifurcation at $t = a$, a repulsive focus, a focus-node bifurcation at $t = b$ and a repulsive node. At point $t = b$, the two eigenvalues coincide. We assume that $\lambda(t) = \mu(t)$ only at point b . Actually, in the complex plane, the two eigenvalues are the two determinations of a multiform function defined on a Riemann surface with a square root singularity at point b .

However, there is a symmetry: if the function $\sqrt{\cdot}$ is defined with a cut-off on the positive real axis, it satisfies $\overline{\sqrt{s}} = -\sqrt{\bar{s}}$ and we then have

$$\mu(\bar{t}) = \overline{\lambda(t)}$$

HFN5 For the same reason, the two reliefs

$$R_\lambda(t) = \Re \left(\int_a^t \lambda(t) dt \right) \quad \text{and} \quad R_\mu(t) = \Re \left(\int_a^t \mu(t) dt \right)$$

are the two determinations of a multiform function with a square root singularity at point $t = b$. However, there is a symmetry: if the function $\sqrt{\cdot}$ is defined with a cut-off on the positive real axis, it satisfies $\overline{\sqrt{s}} = -\sqrt{s}$ and we have then: $R_\mu(t) = R_\lambda(\bar{t})$ except on the cut-off half line $[b, +\infty[$. For real $t > b$, we choose determinations of square root such that $\lambda(t) < \mu(t)$. We assume that R_λ has a unique critical point with critical value R_c . We assume that $R_\lambda(b) < R_c$. An example is given and studied in paragraph 3.2.1.

3.1 Input-output function when there exists a big canard

We assume now that there exists a *big canard* $\tilde{X}(t)$ i.e. a solution of equation (1) such that $\tilde{X}(t) \simeq \phi(t)$ for all t in the S -interior of $]t_m, t_M[$. The study below is similar to paragraph 2.1. The added difficulty is the coincidence of the two eigenvalues at point b which do not allow to diagonalize the linear part.

The first change of unknown is $X = \tilde{X}(t) + Z$ which moves the big canard on the axis $X = 0$:

$$\varepsilon \dot{Z} = A(t)Z + O(\varepsilon)Z + O(Z^2) \quad , \quad A(t) = D_X f(t, \phi(t), 0)$$

Let us denote $\begin{pmatrix} \alpha(t) & \beta(t) \\ \gamma(t) & \delta(t) \end{pmatrix}$ the coefficients of the matrix $A(t)$. As in paragraph 2.1, the change of unknowns

$$Z = \begin{pmatrix} r \cos \theta \\ r \sin \theta \end{pmatrix} \quad , \quad r = \exp\left(\frac{\rho}{\varepsilon}\right)$$

gives the new system:

$$\begin{cases} \dot{\rho} &= \alpha(t) \cos^2 \theta + (\beta(t) + \gamma(t)) \cos \theta \sin \theta + \delta(t) \sin^2 \theta + O(\varepsilon) + e^{\frac{\rho}{\varepsilon}} k_1(r, \theta, \varepsilon) \\ \varepsilon \dot{\theta} &= \gamma(t) \cos^2 \theta + (\delta(t) - \alpha(t)) \cos \theta \sin \theta - \beta(t) \sin^2 \theta + O(\varepsilon) + e^{\frac{\rho}{\varepsilon}} k_2(r, \theta, \varepsilon) \end{cases} \quad (5)$$

For nonpositive ρ (more precisely, for infinitesimal r), the second equation is a slow-fast equation. Its slow curve is given by

$$\theta = \arctan \left(\frac{\delta(t) - \alpha(t) \pm \sqrt{\alpha(t)^2 - 2\alpha(t)\delta(t) + \delta(t)^2 + 4\beta(t)\gamma(t)}}{2\beta(t)} \right)$$

It has two branches when λ and μ are reals, one is attractive, the other is repulsive: see figure 4.

When θ goes along a branch of the slow curve, (and when r is infinitesimal), an easy computation shows that $\dot{\rho}$ is infinitely close to one of the eigenvalues λ or μ . The repulsive branch corresponds to the smallest eigenvalue (which is real positive). When $t < b$, the angle θ moves infinitely fast, and an averaging procedure is needed to evaluate the variation of ρ :

$$\langle \dot{\rho} \rangle = \frac{\int_{\theta_1}^{\theta_1+2\pi} \frac{\dot{\rho}}{\theta} d\theta}{\int_{\theta_1}^{\theta_1+2\pi} \frac{1}{\theta} d\theta}$$

An easy computation shows now that, in the S -interior of the domain $t < b$, $\rho < 0$, we have

$$\langle \dot{\rho} \rangle \simeq \frac{\alpha(t) + \delta(t)}{2} = \Re(\lambda(t)) = \Re(\mu(t))$$

Let us give an initial condition (t, θ) between the two branches of the slow curve and ρ negative non infinitesimal (in the example, we can take $t = 0.8$, $\theta = 0$, $\rho = -0.03$). For increasing t , the curve $(t, \theta(t))$ goes along the attractive branch of the slow curve, while ρ believes negative non infinitesimal. For decreasing t , the solution goes along the repulsive branch, then θ moves infinitely fast while ρ believes negative non infinitesimal. Consequently, we know the variation of $\rho(t)$ (see figure 4). As in paragraph 2.1, a more subtle argument is needed to prove that when ρ becomes infinitesimal, the variable r becomes non infinitesimal and the trajectory X leaves the neighborhood of the slow curve.

From this study, all the behaviours of $\rho(t)$ are known, depending on the initial condition. They are drawn on figure 5.

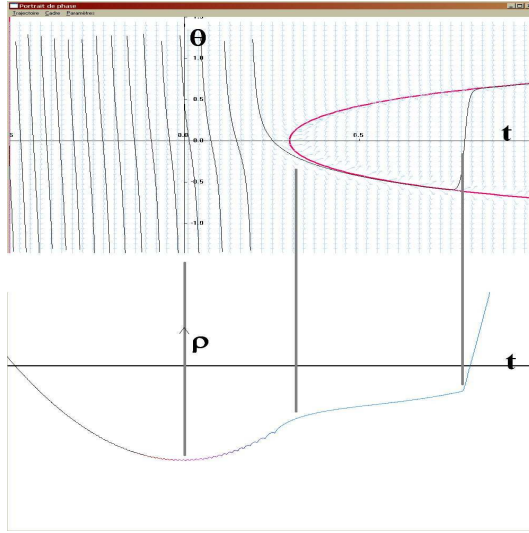


Figure 4: One of the trajectories of system $0.002\dot{X} = \begin{pmatrix} t & 1 \\ t-0.3 & t \end{pmatrix} X$ drawn with the variables (θ, ρ) . The slow curve is also drawn

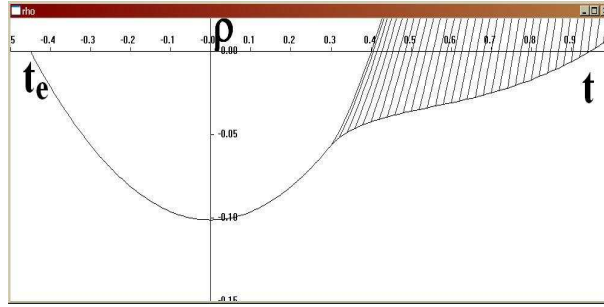


Figure 5: The possible behaviours of $\rho(t)$.

Proposition 4 *Let us give an equation of type (6) with hypothesis HFN1 to HFN5. Assume also that there exists a big canard $\tilde{X}(t)$ going along the slow curve on the whole interval $]t_m, t_M[$. If a trajectory $X(t)$ goes along the slow curve exactly on an interval $]t_e, t_s[$ with $[t_e, t_s] \subset]t_m, t_M[$, then*

$$\int_{t_e}^{t_s} \Re(\lambda(\tau)) d\tau \leq 0 \leq \int_{t_e}^{t_s} \Re(\mu(\tau)) d\tau$$

Conversely, if the inequalities above are satisfied, there exists a trajectory going along the slow curve exactly on $]t_e, t_s[$.

The input-output relation is described by its graph, drawn on figure 6.

We could give more precise results if we consider the two variables r and θ for the input-output relation. Indeed, when the point (t_e, t_s) is in the interior of the graph of the input-output relation, we know that, at time of output, θ is going along the attractive slow curve which corresponds to the unique fast trajectory tangent to the eigenspace of the biggest eigenvalue μ .

3.2 The focus-node bifurcation is a bump

Here is the main part of this article. Today, I am not able to prove the expecting results, but I have propositions in this direction. To explain the problem, I will give conjectures.

Let us define the anti-bump t^{**} and the two bumps t_λ^* and t_μ^* as in definition 3:

$$R_\lambda(t_c) = R_\mu(t_c) = R_\lambda(t^{**}) = R_\mu(t^{**}) = R_\lambda(t_\lambda^*) = R_\mu(t_\mu^*)$$

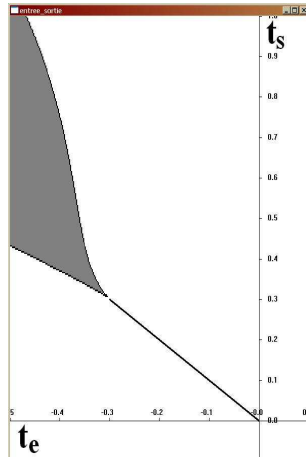


Figure 6: The input-output relation for equation (6) when there exists a big canard.

. We have $t^{**} < a < t_\mu^* = t_\lambda^* \leq b$ or $t^{**} < a < b < t_\mu^* < t_\lambda^*$. In the first case, the bump is before the focus node bifurcation, and the study of paragraph 2.2 is available. The interesting case is the second, where the computed bump is after the focus node bifurcation, this case is assumed with hypothesis HFN5.

Conjecture 5 *With hypothesis HFN1 to HFN5, the following proposition is generically wrong:*

*If a trajectory of (4) goes along the slow curve at least on $]t^{**}, a[$, then it goes until the slow curve at least on $[t^{**}, t_\mu^*]$.*

To work on this conjecture, we will study an example which is, in some sense, a normal form of the problem: the slow curve is moved on the t -axis and the fast vector field is linearized. The example is

$$\begin{cases} \varepsilon^3 \dot{x} &= tx + y + \varepsilon^3 c_1 \\ \varepsilon^3 \dot{y} &= (t - b)x + ty + \varepsilon^3 c_2 \end{cases} \quad (6)$$

Proposition 6 *A numerical simulation of equation (6) gave the figure 7. It confirms conjecture 5.*

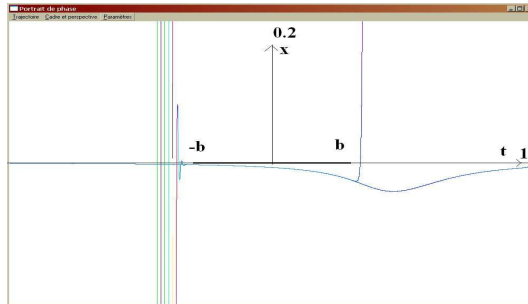


Figure 7: Trajectories X_- and X_+ : the first goes along the horizontal axis from $-\infty$ to b , where it jumps outside the neighborhood of the horizontal axis; the second one goes along the horizontal axis from $+\infty$ to $-b$ where it has big oscillations. The parameters are $b = 0.3$, $c_1 = 0$, $c_2 = -1$, $\varepsilon^3 = 0.002$, the trajectory is computed with a RK4 method, with step 0.0001. Other methods and other steps were tried, and the results are always very similar.

This proposition gives a good argument for the next conjecture, more precise than the first one:

Conjecture 7 *If a trajectory of system (4) goes along the slow curve in a neighborhood of a real t with $t < a$ and $R(t) > R(b)$, then it does not go along the slow curve after the focus-node bifurcation point b .*

So, generically, the input-output relation of equation (4) has a graph similar to the graph of figure 3; if $R(t^{**}) > R(b)$, we have to replace t^* by b et t^{**} par t_b^{**} where $R(t_b^{**}) = R(b)$. The delay of the Hopf bifurcation is stopped either by the bump (as in case of a Hopf bifurcation alone) either by the focus-node bifurcation.

Proposition 8 *If the conjecture 7 is true for one trajectory, then it is true for all of them.*

Proof Assume that equation (4) has a solution \tilde{X} which does not verify conjecture 7. Then, \tilde{X} goes along the slow curve on an interval $]t_1, t_2[$ with $t_1 < t_b^{**} < a < b < t_2$. If the problem is considered on a restricted interval $]t_1, t_2[$, the equation has a big canard, and we can apply the proposition 4. Then all trajectories going along the slow curve before t_b^{**} goes along the slow curve until b , and even a little more. ■

In this article, we will now study only equation (6). We changed ε into ε^3 only to avoid fractionnary exponents. The analytic structure with respect to ε is obviously modified, but does not matter for our purpose.

To study the phase portrait of equation (4) or (6), two trajectories are very important. They are called *distinguished trajectories* by JL.Callot and they are very classical. The first one, denoted X_+ goes along the slow curve for t near t_M . Similarly, X_- goes along the slow curve for t near t_m . These two trajectories are Fenichel's manifolds, they are unique when $t_m = -\infty$ and $t_M = +\infty$. For the particular equation (6), these two trajectories are drawn on figure 7. We have for this example a nice fact: X_- and X_+ have an explicit formula, using the Airy function (in an appendix (section 4) , we give classical needed results on Airy functions and Airy equation).

$$X_+(t) = \begin{pmatrix} x_+(t) \\ y_+(t) \end{pmatrix} = -e^{\frac{1}{2}\frac{t^2}{\varepsilon^3}} M(t) \int_t^{+\infty} e^{-\frac{1}{2}\frac{\tau^2}{\varepsilon^3}} M^{-1}(\tau) d\tau \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} \quad (7)$$

$$X_-(t) = \begin{pmatrix} x_-(t) \\ y_-(t) \end{pmatrix} = e^{\frac{1}{2}\frac{t^2}{\varepsilon^3}} M(t) \int_{-\infty}^t e^{-\frac{1}{2}\frac{\tau^2}{\varepsilon^3}} M^{-1}(\tau) d\tau \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} \quad (8)$$

$$\text{where } M(t) = \sqrt{\frac{\pi}{\varepsilon}} \begin{pmatrix} A(j\frac{t-b}{\varepsilon^2}) & A(j^2\frac{t-b}{\varepsilon^2}) \\ \varepsilon j A'(j\frac{t-b}{\varepsilon^2}) & \varepsilon j^2 A'(j^2\frac{t-b}{\varepsilon^2}) \end{pmatrix} \quad \text{with} \quad \det(M(t)) = \frac{i}{2} \quad (9)$$

All the integrals are convergent because the Airy function is bounded at infinity by $C|t|^{-\frac{3}{2}} e^{\frac{2}{3}|t|^{\frac{3}{2}}}$.

3.2.1 The relief

In this paragraph, we want to explore the methods used in paragraph 2.2 when there is a focus-node bifurcation. We also check the hypothesis HFN1 to HFN5.

Hypothesis HFN1 to HFN3 are obvious with the slow curve $\phi(t, X, 0) = 0$ and the domain $\mathcal{D} = \mathbb{C} \times \mathbb{C}^2 \times \mathbb{C}$.

The computation of the eigenvalues of the jacobian matrix $J(t) = \begin{pmatrix} t & 1 \\ t-b & t \end{pmatrix}$ gives

$$\lambda(t) = t - (t-b)^{\frac{1}{2}} \quad \mu(t) = t + (t-b)^{\frac{1}{2}}$$

The determination of the square root is needed to allow the formula above. In all this paragraph, we choose a cut-off on the positive real axis:

$$(re^{i\theta})^{\frac{1}{2}} = \sqrt{r} e^{\frac{i\theta}{2}} \quad \theta \in [0, 2\pi[$$

For the function $()^{\frac{3}{2}}$, we choose the same cut-off.

The relation $\overline{t^{\frac{1}{2}}} = -\overline{t}^{\frac{1}{2}}$ will be useful. Then, λ and μ are the two determinations of a multiform function. The cut-off is the semi-axis $[b, +\infty[$, and $\overline{\mu(t)} = \lambda(\overline{t})$.

For $a = 0$ and

$$b > \frac{1}{4}, \quad (10)$$

the hypothesis HFN4 is easy to check.

The two associated reliefs are given by

$$F_\lambda(t) = \frac{1}{2}t^2 - \frac{2}{3}(t-b)^{\frac{3}{2}} - \frac{2}{3}ib^{\frac{3}{2}} \quad F_\mu(t) = \frac{1}{2}t^2 + \frac{2}{3}(t-b)^{\frac{3}{2}} + \frac{2}{3}ib^{\frac{3}{2}}$$

$$R_\lambda(t) = \Re(F_\lambda(t)) \quad R_\mu(t) = \Re(F_\mu(t))$$

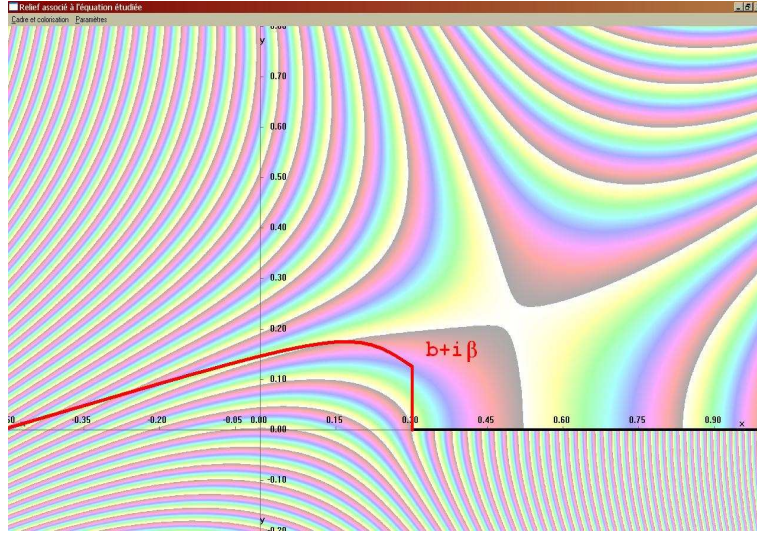


Figure 8: Level curves of relief R_λ for $b = 0.3$, and path used in paragraph 3.2.4.

Let us comment figure 8: the value of R_λ is $+\infty$ at both ends of the real axis. If a path goes from $t = -\infty$ to $t = +\infty$, it has to go down at least until the mountain pass, which is the unique critical point of the relief given by

$$t_c = \frac{1}{2} + i\sqrt{b - \frac{1}{4}} = 0.500 + 0.224 i$$

The value of the relief at this critical point is

$$R_c = R_\lambda(t_c) = \frac{1}{2}b - \frac{1}{12} = 0.067$$

We solve now on the real axis the equation $R_\lambda(t) = R_c$. The solution are t_e and t_s given by

$$\left\{ \begin{array}{l} \text{if } b > \frac{1}{2} + \frac{1}{6}\sqrt{3} \\ \text{if } b < \frac{1}{2} + \frac{1}{6}\sqrt{3} \end{array} \right. , \quad \left\{ \begin{array}{l} t_e = -\sqrt{b - \frac{1}{6}} \\ t_e = -\sqrt{b - \frac{1}{6}} \end{array} \right. \quad \left\{ \begin{array}{l} t_s = \sqrt{b - \frac{1}{6}} \\ t_{s1} = \dots \\ t_{s2} = \dots \end{array} \right.$$

The symbols ... in the formula above are the solutions of a polynomial in t of degree 4. The exact expression is not needed. For $b = 0.3$, we have

$$t_e = -0.365 \quad t_{s1} = 0.346 \quad t_{s2} = 0.525$$

The value t_{s1} is on the sheet right to the cut-off: $\arg(t_{s1}) = 2\pi$. Besides, t_{s2} is on the sheet left to the cut-off: $\arg(t_{s2}) = 0$. When we look on the polynomial which has t_{s1} and t_{s2} as roots, we can prove that the hypothesis HFN5 is satisfied for

$$\frac{1}{4} < b < \frac{1}{2} + \frac{1}{6}\sqrt{3} \quad (11)$$

3.2.2 Callot's domains

To study the canards of equation (6), we introduce two special solutions, called distinguished solutions by J.L. Callot: $X_+ = (x_+, y_+)$ has an asymptotic¹ condition $X_+(+\infty) = 0$ and $X_- = (x_-, y_-)$ has an asymptotic condition $X_-(-\infty) = 0$. They are unique. In this paragraph we build a domain \mathcal{D}_+ where X_+ is infinitesimal (it corresponds in the complex plane to the expression "going along a real interval"). In almost all situations, the builded domain is the maximal domain with this property.

¹Here the things are easier than in the general case because the domain \mathcal{D}_t contains the whole real axis. In general case, there is no unicity of the distinguished solution, but the difference remains exponentially smaller than the computed quantities.

For trajectory X_+ In this paragraph, it is better to change the cut-off, and we define (only in this paragraph)

$$(re^{i\theta})^{\frac{1}{2}} = \sqrt{r} e^{\frac{i\theta}{2}} \quad \theta \in [-\frac{1}{2}\pi, \frac{3}{2}\pi[$$

We are looking for a complex domain \mathcal{D}_+ such that the real point $+\infty$ is in \mathcal{D}_+ , the singularity b is not in \mathcal{D}_+ . We look for domains below $+\infty$ (see definition 2) for the relief R_λ and also below $+\infty$ for the relief R_μ .

On figure 9, such domain is drawn² in dark. Attention: at the left, the domain has a spike with a real part smaller than $-b$ and a nonzero imaginary part. The intersection of \mathcal{D}_+ with the real axis is $] -b, b[\cup] b, +\infty[$. The theorem of Callot (theorem 2) says that X^+ is infinitesimal on the whole S -interior of \mathcal{D}_+ .

Actually, a more precise study shows that the domain \mathcal{D}_+ is not the maximal domain where X^+ is infinitesimal: if we consider domains on the the Riemann surface (two sheets covering) we can add to \mathcal{D}_+ its conjugate (drawn in lightgray on the figure 9). Because the solution X^+ is analytic without singularity at point b , it is infinitesimal on the symetric domain.

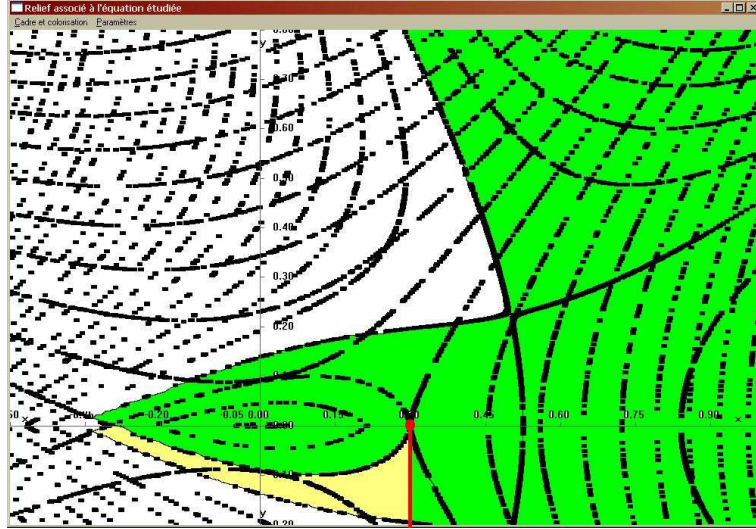


Figure 9: The domain \mathcal{D}_+ for $b = 0.3$

For trajectory X_- A similar method gives the domain \mathcal{D}_- such that X_- is infinitesimal on the S -interior of \mathcal{D}_- . It is easier because we do not need to consider a two sheets covering. The domain \mathcal{D}_- is drawn on figure 10.

3.2.3 Evaluation of $X_+(b)$

The slow curve $x = y = 0$ is repulsive for all positive t . Then the trajectory X_+ is infinitesimal at least for all t positive non infinitesimal (in fact it is infinitesimal on a larger interval). Its asymptotic expansion in power of ε^3 is given by formal identification in the equation: $X_+ = \sum_{n \geq 0} X_n(t) \varepsilon^{3n}$ has to verify the recurrence identities

$$\begin{cases} \dot{x}_{n-1} &= tx_n + y_n + \delta_{n-1}c_1 \\ \dot{y}_{n-1} &= (t-b)x_n + ty_n + \delta_{n-1}c_2 \end{cases}$$

where $\delta_{n-1} = 1$ if $n = 1$ and vanishes for all others n . The computation of the first terms is easy:

$$x(t) = \frac{-c_1 t + c_2}{t^2 - t + b} \varepsilon^3 + \frac{t(t^2 + t - 3b)c_1 + (-3t^2 + t + b)c_2}{(t^2 - t + b)^3} \varepsilon^6 + O(\varepsilon^9)$$

$$y(t) = \frac{(t-b)c_1 - t c_2}{t^2 - t + b} \varepsilon^3 + \frac{(-2t^3 + 3bt^2 + bt - b^2)c_1 + (t^3 + 2t^2 - 3bt - t + b)c_2}{(t^2 - t + b)^3} \varepsilon^6 + O(\varepsilon^9)$$

and we we have now proved the

²The picture is a little bit different when b is greater or smaller than $-\frac{3}{2} - \frac{1}{6}\sqrt{123}$. For this particular value, we have $R_\mu(b) = R_\mu(tc)$.

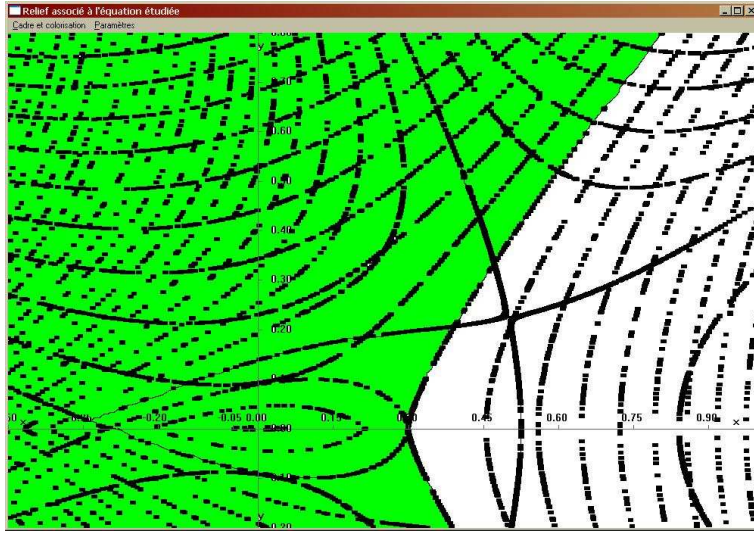


Figure 10: The domain \mathcal{D}_- for $b = 0.3$

Proposition 9

$$X_+(b) = \begin{pmatrix} \left(-\frac{1}{b}c_1 + \frac{1}{b^2}c_2 \right) \varepsilon^3 + \left(\left(\frac{1}{b^3} - \frac{2}{b^4} \right) c_1 + \left(-\frac{3}{b^4} + \frac{2}{b^5} \right) c_2 \right) \varepsilon^6 + O(\varepsilon^9) \\ -\frac{1}{b}c_2\varepsilon^3 + \left(\frac{1}{b^3}c_1 + \left(\frac{1}{b^3} - \frac{1}{b^4} \right) c_2 \right) \varepsilon^6 + O(\varepsilon^9) \end{pmatrix}$$

3.2.4 Evaluation of $X_-(b)$

The simple method above is not convenient to evaluate $X_-(b)$ because we expect that X_- does not go along the slow manifold in a neighborhood of b .

We will use the explicit formula (8) to evaluate $X_-(b)$. The computation is a little bit tedious. In all the formulae below, the symbol \mathcal{O} represent a quantity which goes to zero when $\varepsilon > 0$ goes to zero.

The inverse of the matrix M is easy to compute: we know the determinant of M (see the property 4 in the appendix on Airy's functions).

$$M^{-1}(\tau) = -2i\sqrt{\frac{\pi}{\varepsilon}} \begin{pmatrix} \varepsilon j^2 A' \left(j^2 \frac{\tau-b}{\varepsilon^2} \right) & -A \left(j^2 \frac{\tau-b}{\varepsilon^2} \right) \\ -\varepsilon j A' \left(j \frac{\tau-b}{\varepsilon^2} \right) & A \left(j \frac{\tau-b}{\varepsilon^2} \right) \end{pmatrix}$$

To compute the integrals in formula (8), we change the real path of integration $]-\infty, b]$. For some integrals we choose a path which goes down the relief R_λ from $-\infty$ to b , for other integrals, we choose the conjugate path which goes down the relief R_μ (the idea is the same as in Callot's proof of theorem 2). The path which goes down R_λ is drawn on figure 8. The end of the path is a vertical segment from $b + i\beta$ to b . At point b , it is tangent to the level curve of the relief, then, the path does not go down the relief with the precise definition 1. Thus, we have to be care with approximations at this point.

Let us denote

$$f(\tau) = e^{\frac{1}{2} \frac{b^2 - \tau^2}{\varepsilon^3}} A \left(j^2 \frac{\tau - b}{\varepsilon^2} \right)$$

It is one of the function we have to integrate to evaluate X_- .

Lemma 10 *Let us give τ such that $\tau - b$ is non infinitesimal and $\frac{1}{3}\pi < \arg(\tau - b) < \pi$. Then*

$$|f(\tau)| = \exp \left(\frac{-1}{\varepsilon^3} (R_\lambda(\tau) - R_\lambda(b) + \mathcal{O}) \right)$$

Proof Using the asymptotic expansion of A (see in appendix), we have:

$$A \left(j^2 \frac{\tau - b}{\varepsilon^2} \right) = \frac{1}{2\sqrt{\pi}} \exp \left(-\frac{2}{3} \left(j^2 \frac{\tau - b}{\varepsilon^2} \right)^{\frac{3}{2}} \right) \left(j^2 \frac{\tau - b}{\varepsilon^2} \right)^{-\frac{1}{4}} (1 + O(\varepsilon^3))$$

Substituting in the formula of f , we have:

$$\varepsilon^3 \ln |f(\tau)| = \Re \left(\frac{1}{2}b^2 - \frac{1}{2}\tau^2 - \frac{2}{3}(j^2(\tau - b))^{\frac{3}{2}} + \mathcal{O} \right)$$

We write $\tau - b$ in polar coordinates: $\tau - b = re^{i\theta}$, with $\theta \in]\frac{1}{3}\pi, \pi[$. Then $j^2(\tau - b) = re^{i(\theta - \frac{2}{3}\pi)}$. Because $\theta - \frac{2}{3}\pi$ has an argument between $-\frac{1}{2}\pi$ and $\frac{1}{2}\pi$, the power $\frac{3}{2}$ gives $(j^2(\tau - b))^{\frac{3}{2}} = r^{\frac{3}{2}}e^{i(\frac{3}{2}\theta - \pi)}$. This expression can be writed $-(\tau - b)^{\frac{3}{2}}$, with the same determination of $t^{\frac{3}{2}}$ as in R_λ . ■

The interesting consequence of this lemma is that along the considered path, the function f is increasing with a logarithmic derivative of type ε^{-3} . To precise, we need the following lemma:

Lemma 11 *There exist two constants k and δ standard³, positive such that*

$$\forall \sigma \in [0, \frac{\beta}{\varepsilon^2}] \quad , \quad |f(b + i\sigma\varepsilon^2)| < ke^{-\delta\sigma^{\frac{3}{2}}}$$

Proof By définition of f , we have

$$f(b + i\sigma\varepsilon^2) = e^{-\frac{b}{\varepsilon}\sigma i} e^{\frac{1}{2}\sigma^2\varepsilon} A(ij^2\sigma) \quad \text{then} \quad |f(b + i\sigma\varepsilon^2)| = e^{\frac{1}{2}\sigma^2\varepsilon} |A(ij^2\sigma)|$$

For real positive infinitely large σ , the asymptotic expansion of the Airy function give the estimation

$$|A(ij^2\sigma)| = \frac{1}{2\sqrt{\pi}} \left| e^{-\frac{2}{3}(ij^2\sigma)^{\frac{3}{2}}} \right| |ij^2\sigma|^{-\frac{1}{4}} (1 + \mathcal{O}) = \frac{1}{2\sqrt{\pi}} \sigma^{-\frac{1}{4}} e^{-\frac{\sqrt{2}}{3}\sigma^{\frac{3}{2}}} (1 + \mathcal{O})$$

(we know that $\Re((ij^2)^{\frac{3}{2}}) = \frac{\sqrt{2}}{2}$). Then if δ_1 is real standard, less than $\frac{\sqrt{2}}{3}$, we have the following inequality, true for all σ infinitely large:

$$|A(ij^2\sigma)| < e^{-\delta_1\sigma^{\frac{3}{2}}}$$

By permanence⁴, this inequality believes true for all real σ greater than some positive standard ω . We can deduce the following majoration:

$$\forall \sigma \in [\omega, \frac{\beta}{\varepsilon^2}] \quad , \quad |f(b + i\sigma\varepsilon^2)| < e^{\frac{1}{2}\sigma^2\varepsilon} e^{-\delta_1\sigma^{\frac{3}{2}}}$$

For $\sigma < \omega$, we have:

$$\forall \sigma \in [0, \omega] \quad , \quad |f(b + i\sigma\varepsilon^2)| < e^{\frac{1}{2}\omega^2\varepsilon} k_1 < 2k_1 \quad \text{with} \quad k_1 = \max_{\sigma \in [0, \omega]} |A(ij^2\sigma)|$$

Then we are looking for a constant δ such that

$$\forall \sigma \in [0, \frac{\beta}{\varepsilon^2}] \quad , \quad \frac{1}{2}\sigma^2\varepsilon - \delta_1\sigma^{\frac{3}{2}} < -\delta\sigma^{\frac{3}{2}}$$

The inequality is equivalent to $\sigma < \frac{4(\delta_1 - \delta)^2}{\varepsilon^2}$. A choice of δ less than $\delta_1 - \frac{1}{2}\sqrt{\beta}$ is convenient. This choice is possible only if $\delta_1 > \frac{1}{2}\sqrt{\beta}$ what is true as soon as $\beta < \frac{8}{9}$ and δ_1 near enough from $\frac{\sqrt{2}}{3}$.

To verify the majoration of the lemma for $\sigma < \omega$, we can choose

$$k = 2k_1 e^{\delta\omega^{\frac{3}{2}}}$$

.

The next lemma is the more technical part of the article. The purpose is to evaluate an oscillating integral with successive integrations by parts.

³here, it is the same to assume that k and δ are independent of ε

⁴The non standard arguments in these proofs can be replaced by classical arguments, but, for that, new quantified variables have to be added, and it seems to me that the idea of the proof is more understandable with nonstandard language.

Lemma 12 *We have the following expansion:*

$$\int_{b+i\beta}^b f(\tau)d\tau = -\frac{1}{b}A(0)\varepsilon^3 - \frac{j^2}{b^2}A'(0)\varepsilon^4 - \frac{j}{b^3}A''(0)\varepsilon^5 + \left(\frac{1}{b^3}A(0) - \frac{1}{b^4}A'''(0)\right)\varepsilon^6 + \left(\frac{3j^2}{b^4}A'(0) - \frac{j^2}{b^5}A''''(0)\right)\varepsilon^7 + \mathcal{O}\varepsilon^7$$

Proof Let us substitute τ by $b + i\varepsilon^2\sigma$ in the integral. We have

$$\int_{b+i\beta}^b f(\tau)d\tau = -i\varepsilon^2 \int_0^{\frac{\beta}{\varepsilon^2}} f(b + i\varepsilon^2\sigma)d\sigma = -i\varepsilon^2 \int_0^{\frac{\beta}{\varepsilon^2}} e^{-\frac{b}{\varepsilon}\sigma i} e^{\frac{1}{2}\sigma^2\varepsilon} A(ij^2\sigma)d\sigma$$

The exponential $e^{-\frac{b}{\varepsilon}\sigma i}$ is fast oscillating. The exponential $e^{\frac{1}{2}\sigma^2\varepsilon}$ is infinitely close to 1 for all non infinitely large σ and $A(ij^2\sigma)$ is decreasing. All properties are checked to apply the method of integrations by parts. But there is a difficulty: $e^{\frac{1}{2}\sigma^2\varepsilon}$ is increasing and does not believe close to 1. Now, let us explain the computations:

$$\begin{aligned} I &= \int_{b+i\beta}^b f(\tau)d\tau = I_1 + I_2 + I_3 \quad \text{with} \\ I_1 &= \frac{\varepsilon^3}{b} (f(b + \beta i) - f(b)) \quad I_2 = -\frac{\varepsilon^4}{b} \int_0^{\frac{\beta}{\varepsilon^2}} \sigma f(b + i\varepsilon^2\sigma)d\sigma \\ I_3 &= -\frac{ij^2\varepsilon^3}{b} \int_0^{\frac{\beta}{\varepsilon^2}} \hat{f}(b + i\varepsilon^2\sigma)d\sigma \quad \text{with} \quad \hat{f}(b + i\varepsilon^2\sigma) = e^{-\frac{b}{\varepsilon}\sigma i} e^{\frac{1}{2}\sigma^2\varepsilon} A'(ij^2\sigma) \end{aligned}$$

With lemma 11, we know that $f(b + \beta i)$ is exponentially smaller than $f(b) = A(0)$. Thus, we have

$$I_1 = -\frac{1}{b}A(0)\varepsilon^3 + \mathcal{O}\varepsilon^7$$

To estimate I_2 , we perform a new integration by parts:

$$\begin{aligned} I_2 &= J_1 + J_2 + J_3 + J_4 \quad \text{with} \\ J_1 &= -\frac{i\varepsilon^5}{b^2} \frac{\beta}{\varepsilon^2} f(b + \beta i) \quad J_2 = \frac{i\varepsilon^6}{b^2} \int_0^{\frac{\beta}{\varepsilon^2}} \sigma^2 f(b + i\varepsilon^2\sigma)d\sigma \\ J_3 &= -\frac{j^2\varepsilon^5}{b^2} \int_0^{\frac{\beta}{\varepsilon^2}} \sigma \hat{f}(b + i\varepsilon^2\sigma)d\sigma \quad J_4 = \frac{i\varepsilon^5}{b^2} \int_0^{\frac{\beta}{\varepsilon^2}} f(b + i\varepsilon^2\sigma)d\sigma \end{aligned}$$

Because $f(b + \beta i)$ is exponentially small, we have $J_1 = \mathcal{O}\varepsilon^7$. We have also $J_4 = -\frac{\varepsilon^3}{b^2}I$. If you substitute A' for A the expression I_3 is the same as $\frac{j^2\varepsilon}{b}I$. All the arguments are the same with function A' and function A . Let us denote \hat{I}_i, \hat{J}_i the expressions obtained from I_i and J_i when A' is substituted for A . Thus we have $I_3 = \frac{j^2\varepsilon}{b}\hat{I}$. To estimate J_2 we perform a new integration by parts exactly as for evaluation of I_2 : $J_2 = K_1 + K_2 + K_3 + K_4$. All the integrals are bounded by a non infinitely large real number because all the integrated functions are bounded (see lemma 11) by a integrable standard function. To summarize:

$$\begin{aligned} I &= I_1 + I_2 + I_3 \quad I_2 = J_1 + J_2 + J_3 + J_4 \quad J_2 = K_1 + K_2 + K_3 + K_4 \\ I_1 &= -\frac{1}{b}A(0)\varepsilon^3 + \mathcal{O}\varepsilon^7 \quad J_1 = \mathcal{O}\varepsilon^7 \quad K_1 = \mathcal{O}\varepsilon^7 \\ I_2 &= \mathcal{O}\varepsilon^3 \quad J_2 = \mathcal{O}\varepsilon^5 \quad K_2 = \frac{\varepsilon^8}{b^3} \int_0^{\frac{\beta}{\varepsilon^2}} \sigma^3 f(b + i\varepsilon^2\sigma)d\sigma = \mathcal{O}\varepsilon^7 \\ I_3 &= \frac{j^2\varepsilon}{b}\hat{I} \quad J_3 = \frac{j^2\varepsilon}{b}\hat{I}_2 \quad K_3 = \frac{j^2\varepsilon}{b}\hat{J}_2 \quad J_4 = -\frac{\varepsilon^3}{b^2}I \quad K_4 = -\frac{2\varepsilon^3}{b^2}I_2 \end{aligned} \quad (12)$$

Then, all the ingredients are given, and we can compute the asymptotic expansion of I in powers of ε . To start, we have $I = \mathcal{O}\varepsilon$. For similar reason, $\hat{I} = \mathcal{O}\varepsilon$. Then, using formulae 12, we have $I_3 = \mathcal{O}\varepsilon^2$, then $I = \mathcal{O}\varepsilon^2$. We iterate the process, inserting the known approximations in formulae 12, and we obtain a better approximation: $I_3 = \mathcal{O}\varepsilon^3$ then

$$I = -\frac{1}{b}A(0)\varepsilon^3 + \mathcal{O}\varepsilon^3$$

The next step:

$$I_3 = -\frac{j^2}{b^2}A'(0)\varepsilon^4 + \mathcal{O}\varepsilon^4 \quad J_3 = \mathcal{O}\varepsilon^4 \quad J_4 = \frac{1}{b^2}A(0)\varepsilon^6 + \mathcal{O}\varepsilon^6 \quad I_2 = \mathcal{O}\varepsilon^4$$

$$I = -\frac{1}{b}A(0)\varepsilon^3 - \frac{j^2}{b^2}A'(0)\varepsilon^4 + \mathcal{O}\varepsilon^4$$

The next step: (do not use the relation $A''(0) = 0$, because we have sometimes to substitute A' for A):

$$I_3 = -\frac{j^2}{b^2}A'(0)\varepsilon^4 - \frac{j}{b^3}A''(0)\varepsilon^5 + \mathcal{O}\varepsilon^5 \quad J_3 = \mathcal{O}\varepsilon^5 \quad I_2 = \mathcal{O}\varepsilon^5$$

$$I = -\frac{1}{b}A(0)\varepsilon^3 - \frac{j^2}{b^2}A'(0)\varepsilon^4 - \frac{j}{b^3}A''(0)\varepsilon^5 + \mathcal{O}\varepsilon^5$$

The next step:

$$I_3 = -\frac{j^2}{b^2}A'(0)\varepsilon^4 - \frac{j}{b^3}A''(0)\varepsilon^5 - \frac{1}{b^4}A'''(0)\varepsilon^6 + \mathcal{O}\varepsilon^6 \quad J_3 = \mathcal{O}\varepsilon^6 \quad K_3 = \mathcal{O}\varepsilon^6$$

$$J_4 = \frac{1}{b^3}A(0)\varepsilon^6 + \frac{j^2}{b^4}A'(0)\varepsilon^7 + \mathcal{O}\varepsilon^7 \quad K_4 = \mathcal{O}\varepsilon^8$$

$$J_2 = \mathcal{O}\varepsilon^6 \quad I_2 = \frac{1}{b^3}A(0)\varepsilon^6 + \mathcal{O}\varepsilon^6$$

$$I = -\frac{1}{b}A(0)\varepsilon^3 - \frac{j^2}{b^2}A'(0)\varepsilon^4 - \frac{j}{b^3}A''(0)\varepsilon^5 + \left(\frac{1}{b^3}A(0) - \frac{1}{b^4}A'''(0)\right)\varepsilon^6 + \mathcal{O}\varepsilon^6$$

The last step:

$$I_3 = -\frac{j^2}{b^2}A'(0)\varepsilon^4 - \frac{j}{b^3}A''(0)\varepsilon^5 - \frac{1}{b^4}A'''(0)\varepsilon^6 + \left(\frac{j^2}{b^4}A'(0) - \frac{j^2}{b^5}A''''(0)\right)\varepsilon^7 + \mathcal{O}\varepsilon^7 \quad J_3 = \frac{j^2}{b^4}A'(0)\varepsilon^7 + \mathcal{O}\varepsilon^7 \quad K_3 = \mathcal{O}\varepsilon^7$$

$$J_2 = \mathcal{O}\varepsilon^7 \quad I_2 = \frac{1}{b^3}A(0)\varepsilon^6 + \frac{2j^2}{b^4}A'(0)\varepsilon^7 + \mathcal{O}\varepsilon^7$$

$$I = -\frac{1}{b}A(0)\varepsilon^3 - \frac{j^2}{b^2}A'(0)\varepsilon^4 - \frac{j}{b^3}A''(0)\varepsilon^5 + \left(\frac{1}{b^3}A(0) - \frac{1}{b^4}A'''(0)\right)\varepsilon^6 + \left(\frac{3j^2}{b^4}A'(0) - \frac{j^2}{b^5}A''''(0)\right)\varepsilon^7 + \mathcal{O}\varepsilon^7$$

Lemma 13

$$\int_{-\infty}^{b+i\beta} f(\tau) d\tau = \exp\left(\frac{-1}{\varepsilon^3}(R_\lambda(b+i\beta) - R_\lambda(b) + \mathcal{O})\right)$$

Proof The chosen path goes down the relief R_λ , then the lemma is a corollary of the majoration of lemma 10. ■

Lemma 14

$$\int_{-\infty}^b e^{\frac{1}{2}\frac{b^2-\tau^2}{\varepsilon^3}} A\left(j^2\frac{\tau-b}{\varepsilon^2}\right) d\tau = -\frac{1}{b}A(0)\varepsilon^3 - \frac{j^2}{b^2}A'(0)\varepsilon^4 + \left(\frac{1}{b^3} - \frac{1}{b^4}\right)A(0)\varepsilon^6 + \left(\frac{3j^2}{b^4} - \frac{2j^2}{b^5}\right)A'(0)\varepsilon^7 + \mathcal{O}\varepsilon^7$$

$$\int_{-\infty}^b e^{\frac{1}{2}\frac{b^2-\tau^2}{\varepsilon^3}} A'\left(j^2\frac{\tau-b}{\varepsilon^2}\right) d\tau = -\frac{1}{b}A'(0)\varepsilon^3 - \frac{j}{b^3}A(0)\varepsilon^5 + \left(\frac{1}{b^3} - \frac{2}{b^4}\right)A'(0)\varepsilon^6 + \mathcal{O}\varepsilon^7$$

$$\int_{-\infty}^b e^{\frac{1}{2}\frac{b^2-\tau^2}{\varepsilon^3}} A\left(j\frac{\tau-b}{\varepsilon^2}\right) d\tau = -\frac{1}{b}A(0)\varepsilon^3 - \frac{j}{b^2}A'(0)\varepsilon^4 + \left(\frac{1}{b^3} - \frac{1}{b^4}\right)A(0)\varepsilon^6 + \left(\frac{3j}{b^4} - \frac{2j}{b^5}\right)A'(0)\varepsilon^7 + \mathcal{O}\varepsilon^7$$

$$\int_{-\infty}^b e^{\frac{1}{2}\frac{b^2-\tau^2}{\varepsilon^3}} A'\left(j\frac{\tau-b}{\varepsilon^2}\right) d\tau = -\frac{1}{b}A'(0)\varepsilon^3 - \frac{j^2}{b^3}A(0)\varepsilon^5 + \left(\frac{1}{b^3} - \frac{2}{b^4}\right)A'(0)\varepsilon^6 + \mathcal{O}\varepsilon^7$$

Proof With lemmas 12 and 13, the first part of the lemma is proved. A similar computation gives the second part, if we remember that $A''(0) = 0$ but $A'''(0) \neq 0$. So, the vanishing terms are not the same in the two formulas. The two last formulas are the complex conjugate of the two first one. ■

Proposition 15

$$X_-(b) = \begin{pmatrix} \left(-\frac{1}{b}c_1 + \frac{1}{b^2}c_2 \right) \varepsilon^3 + \left(\left(\frac{1}{b^3} - \frac{2}{b^4} \right) c_1 + \left(-\frac{3}{b^4} + \frac{2}{b^5} \right) c_2 \right) \varepsilon^6 + O(\varepsilon^9) \\ -\frac{1}{b}c_2\varepsilon^3 + \left(\frac{1}{b^3}c_1 + \left(\frac{1}{b^3} - \frac{1}{b^4} \right) c_2 \right) \varepsilon^6 + O(\varepsilon^9) \end{pmatrix}$$

Proof Insert the estimations of lemma 14 in the explicit formula (8), and, after tedious simplifications, the proposition is proved. ■

Conjecture 16 *The two values $X_-(b)$ and $X_+(b)$ have the same asymptotic expansion.*

With Maple, I checked that the two expansions coincide until terms in ε^9 .

4 Appendix: Airy's functions

The Airy's equation is linear, non autonomous of second order. It is

$$\frac{d^2x}{dt^2} = tx \tag{13}$$

The pair $(A(t), B(t))$ of Airy's functions is a fundamental system of solutions. The function satisfy the following properties (these results can be found in every book on special functions).

1. The value at the origin are:

$$A(0) = 3^{-\frac{2}{3}} \frac{1}{\Gamma(\frac{2}{3})} \quad A'(0) = -\frac{3^{\frac{1}{6}} \Gamma(\frac{2}{3})}{2 \pi} \quad B(0) = 3^{-\frac{1}{6}} \frac{1}{\Gamma(\frac{2}{3})} \quad B'(0) = \frac{3^{\frac{2}{3}} \Gamma(\frac{2}{3})}{2 \pi}$$

2. On a sector of angle less than $\frac{2}{3}\pi$, around the positive real axis⁵, the Airy's functions have an asymptotic expansion for t going to infinity:

$$\begin{aligned} A(t) &= \frac{1}{2\sqrt{\pi}} e^{-\frac{2}{3}t^{\frac{3}{2}}} t^{-\frac{1}{4}} (1 + O(t^{-\frac{3}{2}})) & A'(t) &= -\frac{1}{2\sqrt{\pi}} e^{-\frac{2}{3}t^{\frac{3}{2}}} t^{\frac{1}{4}} (1 + O(t^{-\frac{3}{2}})) \\ B(t) &= \frac{1}{\sqrt{\pi}} e^{\frac{2}{3}t^{\frac{3}{2}}} t^{-\frac{1}{4}} (1 + O(t^{-\frac{3}{2}})) & B'(t) &= \frac{1}{\sqrt{\pi}} e^{\frac{2}{3}t^{\frac{3}{2}}} t^{\frac{1}{4}} (1 + O(t^{-\frac{3}{2}})) \end{aligned}$$

The functions A et B are oscillating when t goes to $-\infty$.

3. Let us denote $j = e^{\frac{2}{3}i\pi} = -\frac{1}{2} + \frac{\sqrt{3}}{2}i$. The Airy's equation is invariant by the change of variable $t \mapsto jt$, then $A(jt)$ and $B(jt)$ are also solutions. So they can be written as a linear combination of $A(t)$ and $B(t)$. We perform an identification at point 0 to find the coefficients:

$$\begin{aligned} A(jt) &= -\frac{1}{2}j^2A(t) + \frac{1}{2}ij^2B(t) & B(jt) &= \frac{3}{2}ij^2A(t) - \frac{1}{2}j^2B(t) \\ A(j^2t) &= -\frac{1}{2}jA(t) - \frac{1}{2}ijB(t) & B(j^2t) &= -\frac{3}{2}ijA(t) - \frac{1}{2}jB(t) \end{aligned}$$

4. Classically, the couple $(A(t), B(t))$ is chosen for a base of the set of solutions. It could be better (in a study in the complex plane) to choose $(A(jt), A(j^2t))$ for base. With Liouville's theorem, we prove that the following determinant is constant, and we compute its value at the origin.

$$\det \begin{pmatrix} A(jt) & A(j^2t) \\ jA'(jt) & j^2A'(j^2t) \end{pmatrix} = \frac{i}{2\pi}$$

⁵Take care: the determination of $t^{\frac{3}{2}}$ is here the classical determination with a cut off on the negative real axis, not the determination choose along all this article

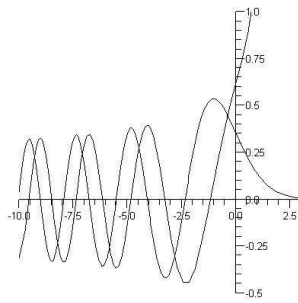


Figure 11: Graphs of Airy's real functions A and B

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