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## To cite this version:

Bernard Roynette, Marc Yor. An interesting family of Black-Scholes perpetuities. 2008. hal00344713

HAL Id: hal-00344713
https://hal.science/hal-00344713
Preprint submitted on 5 Dec 2008

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# An interesting family of Black-Scholes perpetuities 

B. Roynette ${ }^{(1)}$, M. Yor ${ }^{(2),(3)}$

12/11/2008

${ }^{(1)}$ Institut Elie Cartan, Université Henri Poincaré, B.P. 239, 54506 Vandoeuvre les Nancy Cedex email : bernard.roynette@iecn.u-nancy.fr<br>${ }^{(2)}$ Laboratoire de Probabilités et Modèles Aléatoires, Université Paris VI et VII, 4 place Jussieu - Case 188<br>F - 75252 Paris Cedex 05<br>email : deaproba@proba.jussieu.fr<br>${ }^{(3)}$ Institut Universitaire de France

Abstract We obtain the Laplace transform and integrability properties of the integral over $\mathbb{R}_{+}$of the call quantity associated with geometric Brownian motion with negative drift, thus adding a new element to the list of already studied Brownian perpetuities.

Key words Bessel-McDonald functions, Sturm-Liouville equation, perpetuities.
AMS Classification 60 J 25, 60 J 65, 60 G 44.

## 1 Motivation, introduction

1.1 Let $\left(B_{t}, t \geq 0\right)$ denote 1-dimensional Brownian motion starting from 0 ; associated to $B$, one considers the geometric Brownian motion :

$$
\left(\mathcal{E}_{t}:=\exp \left(B_{t}-\frac{t}{2}\right), t \geq 0\right)
$$

which is a positive martingale converging a.s. to 0 as $t \rightarrow \infty$. In recent years, the following questions have been asked to the second author, in connection with European option pricing :
i) to express as simply as possible the quantity :

$$
\int_{0}^{\infty} \lambda e^{-\lambda t} E\left(\left(\left(\mathcal{E}_{t}-k\right)^{+}\right) d t \quad(k \geq 0)\right.
$$

ii) to find the law of $\int_{0}^{s}\left(\mathcal{E}_{t}-k\right)^{+} d t$, for fixed $s$
iii) to find the law of $\int_{0}^{\infty} \lambda e^{-\lambda t}\left(\mathcal{E}_{t}-k\right)^{+} d t$
i) may be solved quite explicitly, and indeed this has led the authors, jointly with D. Madan, to write a series of papers relating European option prices and laws of last passage times of continuous positive martingales converging to 0 as $t \rightarrow \infty$ (see, e.g. [MRY1], [MRY2], [MRY3]).
Questions ii) and iii) are harder to solve explicitly, as, indeed, one may look for double Laplace transforms of either quantities, e.g. :

$$
\int_{0}^{\infty} e^{-\mu t} E\left(e^{-\lambda \int_{0}^{t}\left(\mathcal{E}_{s}-k\right)^{+} d s}\right) d t
$$

In the present paper, we are studying thoroughly the law of $\int_{0}^{\infty}\left(\mathcal{E}_{t}-k\right)^{+} d t$ which seems to be slightly less difficult than either question $i i$ ) or $i i i$ ). However, the results we obtain are not particulary simple.
1.2 These motivations having been presented, we now concentrate exclusively on the law of $\int_{0}^{\infty}\left(\mathcal{E}_{t}-k\right)^{+} d t$. It will be convenient to consider the two-parameter process :

$$
\begin{equation*}
\left(\mathcal{E}_{t}^{(x)}:=x \exp \left(B_{t}-\frac{t}{2}\right)=\exp \left(\log x+B_{t}-\frac{t}{2}\right), t \geq 0\right) \tag{1.1}
\end{equation*}
$$

$\left(\mathcal{E}_{t}^{(x)}, t \geq 0, x>0\right)$ is a Markov process taking values in $\mathbb{R}_{+}$which, most often, we shall denote as : $\left(\left(\mathcal{E}_{t}, t \geq 0 ; P_{x}, x>0\right)\right.$ since $\mathcal{E}_{0}^{(x)}=x$.
To simplify notation, we shall write $P$ for $P_{1}$ and $\mathcal{E}_{t}$ for $\mathcal{E}_{t}^{(1)}$.
For any $k>0$, let us define :

$$
\begin{equation*}
\Pi_{k}^{(x)}:=\int_{0}^{\infty}\left(\mathcal{E}_{t}^{(x)}-k\right)^{+} d t \tag{1.2}
\end{equation*}
$$

Again, to simplify, we shall write $\Pi_{k}$ for $\Pi_{k}^{(1)}$. Since $\mathcal{E}_{t}^{(x)} \underset{t \rightarrow \infty}{\longrightarrow} 0$ a.s., the integral which defines (1.2) is a.s. convergent, as $\left(\mathcal{E}_{t}^{(x)}-k\right)^{+}=0$ for $t \geq \mathcal{G}_{k}:=\sup \left\{t ; \mathcal{E}_{t}^{(x)}=k\right\}$, and $\mathcal{G}_{k}<\infty$ a.s.
1.3 We now explain about some reductions of the study of the laws of $\Pi_{k}^{(x)}$, to closely related problems.
i) From Itô's formula, we deduce :

$$
\begin{equation*}
\mathcal{E}_{t}^{(x)}=x+\int_{0}^{t} \mathcal{E}_{s}^{(x)} d B_{s}=\beta_{A_{t}^{(x)}} \tag{1.3}
\end{equation*}
$$

where $\left(\beta_{u}, u \geq 0\right)$ denotes the Dubins-Schwarz Brownian motion associated with $\left(\mathcal{E}_{t}^{(x)}, t \geq 0\right)$ (and $\beta_{0}=x$ ) and :

$$
\begin{equation*}
A_{t}^{(x)}=\left\langle\mathcal{E}^{(x)}\right\rangle_{t}=\int_{0}^{t}\left(\mathcal{E}_{s}^{(x)}\right)^{2} d s \tag{1.4}
\end{equation*}
$$

Hence :

$$
\begin{equation*}
\Pi_{k}^{(x)}=\int_{0}^{\infty}\left(\beta_{A_{s}^{(x)}}-k\right)^{+} d s=\int_{0}^{T_{0}(\beta)} \frac{\left(\beta_{v}-k\right)^{+}}{\beta_{v}^{2}} d v \tag{1.5}
\end{equation*}
$$

(after the change of variable $A_{s}^{(x)}=v$ and where $T_{0}(\beta)=\inf \left\{u \geq 0 ; \beta_{u}=0\right\}$ )

$$
\begin{equation*}
=\int_{k}^{\infty} \frac{(y-k)}{y^{2}} L_{T_{0}(\beta)}^{y} d y \quad\left(=\int_{k}^{\infty} d z \int_{z}^{\infty} \frac{d y}{y^{2}} L_{T_{0}(\beta)}^{y}\right) \tag{1.6}
\end{equation*}
$$

from the density of occupation formula, and where $L_{T_{0}}^{y}$ denotes the local time at time $T_{0}$ and level $y$ of Brownian motion ( $\beta_{u}, u \geq 0$ ).
ii) When $x=k$, the first Ray-Knight Theorem allows us to write, from (1.6) :

$$
\begin{equation*}
\Pi_{k}^{(k)}=\int_{k}^{\infty} \frac{(y-k)}{y^{2}} \lambda_{y-k} d y \tag{1.7}
\end{equation*}
$$

where ( $\lambda_{z}, z \geq 0$ ), conditionally on $\lambda_{0}=l$, is a 0 -dimensional squared Bessel process starting at $l$, and where $\lambda_{0}$ is an exponential variable with parameter $\frac{1}{2 k}$, i.e. with expectation $2 k$.
1.4 From the elementary relations:

$$
\begin{equation*}
\left.\left(x e^{B_{t}-\frac{t}{2}}-k\right)^{+}=x\left(e^{B_{t}-\frac{t}{2}}-\frac{k}{x}\right)^{+}=k\left(\frac{x}{k}\left(e^{B_{t}-\frac{t}{2}}\right)-1\right)\right)^{+} \tag{1.8}
\end{equation*}
$$

valid for every $x, k>0$, we deduce that the law of $\Pi_{k}$ under $P_{x}$ is that of $x \Pi_{\frac{k}{x}}$ under $P$, or that of $k \Pi_{1}$ under $P_{\frac{x}{k}}$. In other words, for every Borel positive function $\varphi$, we have :

$$
\begin{equation*}
E_{x}\left[\varphi\left(\Pi_{k}\right)\right]=E\left[\varphi\left(x \Pi_{\frac{k}{x}}\right)\right]=E_{\frac{x}{k}}\left(\varphi\left(k \Pi_{1}\right)\right) \tag{1.9}
\end{equation*}
$$

These relations allow us to reduce our study of the law of $\Pi_{k}$ under $P_{x}$ to that of $\Pi_{1}$ under $P_{\frac{x}{k}}$. We might as well limit ourselves to the study of $\Pi_{k}$ under $P$.
1.5 It is proven in Dufresne [Duf] and Yor [Y] (see also P. Salminen and M. Yor [SY1]) that, for every $a \neq 0$ and $\nu>0$ :

$$
\begin{equation*}
\int_{0}^{\infty} \exp \left(a B_{t}-\nu t\right) d t \stackrel{(\text { law })}{=} \frac{2}{a^{2} \gamma_{\frac{2 \nu}{a^{2}}}} \tag{1.10}
\end{equation*}
$$

where $\gamma_{b}$ denotes a gamma variable with parameter $b$, i.e. :

$$
\begin{equation*}
P\left(\gamma_{b} \in d t\right)=\frac{1}{\Gamma(b)} e^{-t} t^{b-1} d t \quad(t \geq 0) \tag{1.11}
\end{equation*}
$$

In particular, for $a=1$ and $\nu=\frac{1}{2}$

$$
\begin{equation*}
\Pi_{0} \stackrel{\text { (law) }}{=} \frac{2}{\gamma_{1}}=\frac{2}{\mathfrak{e}} \tag{1.12}
\end{equation*}
$$

where $\mathfrak{e}$ is a standard exponential variable.
1.6 Let $\nu \neq 0$ and:

$$
\mathcal{E}_{t}^{(x, \nu)}:=\exp \left(\nu\left(\log x+B_{t}\right)-\frac{\nu^{2} t}{2}\right)=x^{\nu} \exp \left(\nu B_{t}-\frac{\nu^{2} t}{2}\right)
$$

and define, for $k>0$ :

$$
\begin{equation*}
\Pi_{k}^{(x, \nu)}:=\int_{0}^{\infty}\left(\mathcal{E}_{s}^{(x, \nu)}-k\right)^{+} d s \tag{1.13}
\end{equation*}
$$

Since, by scaling, $\left(\mathcal{E}_{t}^{(x, \nu)}, t \geq 0\right) \stackrel{(\text { law })}{=}\left(x^{\nu} \mathcal{E}_{\nu^{2} t}, t \geq 0\right)$, we have :

$$
\begin{align*}
\Pi_{k}^{(x, \nu)} & \stackrel{\text { (law) }}{=} \int_{0}^{\infty}\left(x^{\nu} \mathcal{E}_{\nu^{2} s}-k\right)^{+} d s=\frac{x^{\nu}}{\nu^{2}} \int_{0}^{\infty}\left(\mathcal{E}_{v}-\frac{k}{x^{\nu}}\right)^{+} d v \\
& \stackrel{\text { (law) }}{=} \frac{x^{\nu}}{\nu^{2}} \Pi_{\frac{k}{x^{\nu}}} \tag{1.14}
\end{align*}
$$

Thus, the study of the law of $\Pi_{k}^{(x, \nu)}$ may be reduced very simply to that of $\Pi_{\frac{k}{r^{\nu}}}$. This is the reason why we have chosen, in this paper, to limit ourselves to $\nu=1$.
1.7 Here are our results :

Theorem 1.1 Let $\alpha \geq 0$. Then, for every $x>0, E_{x}\left(\left(\Pi_{1}\right)^{\alpha}\right)<\infty$ if and only if $\alpha<1$
Theorem 1.2 Let $\alpha<0$ and $x>0$. Then :
i) For every $x>1, E_{x}\left(\left(\Pi_{1}\right)^{\alpha}\right)<\infty$.
ii) For every $x<1, P_{x}\left(\Pi_{1}=0\right)=1-x$; hence $E_{x}\left(\left(\Pi_{1}\right)^{\alpha}\right)=+\infty$.
iii) For $x=1, E_{1}\left(\left(\Pi_{1}\right)^{\alpha}\right)<\infty$ if and only if $|\alpha|<\frac{1}{3}$.

Theorem 1.3 (Laplace transform of $\Pi_{1}$ ). For every $\theta \geq 0$ :

$$
E_{x}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right)= \begin{cases}\frac{\sqrt{x} K_{\gamma}(\sqrt{4 \theta x})}{\frac{1}{2} K_{\gamma}(\sqrt{4 \theta})-\sqrt{\theta} K_{\gamma}^{\prime}(\sqrt{4 \theta})} & \text { if } x \geq 1  \tag{1.15}\\ 1+x \frac{\frac{1}{2} K_{\gamma}(\sqrt{4 \theta})+\sqrt{\theta} K_{\gamma}^{\prime}(\sqrt{4 \theta})}{\frac{1}{2} K_{\gamma}(\sqrt{4 \theta})-\sqrt{\theta} K_{\gamma}^{\prime}(\sqrt{4 \theta})} & \text { if } 0<x \leq 1\end{cases}
$$

where $K_{\gamma}$ denotes the Bessel-McDonald function with index $\gamma$ (see [Leb], p. 108) and where $\gamma=\sqrt{1-4 \theta}$ if $4 \theta \leq 1$ and $\gamma=i \sqrt{4 \theta-1}$ if $4 \theta \geq 1$.
We shall also prove, as a consequence of Theorem 1.3 :
Theorem 1.4 Let $\left(\lambda_{x}, x \geq 0\right)$ denote a squared Bessel process with dimension 0 started at $l$ and denote its law by $Q_{l}^{(0)}$. Then :

- If $4 \theta \geq 1$ :

$$
\begin{align*}
Q_{l}^{(0)} & \left(\exp -\frac{\theta}{2} \int_{1}^{\infty} \frac{(x-1)}{x^{2}} \lambda_{x-1} d x\right)\left(=Q_{l}^{(0)}\left(\exp -\frac{\theta}{2} \int_{0}^{\infty} \frac{x}{(x+1)^{2}} \lambda_{x} d x\right)\right. \\
& =\exp \frac{l}{2}\left(\frac{1}{2}+\frac{\sqrt{\theta} K_{i \nu}^{\prime}(\sqrt{4 \theta})}{K_{i \nu}(\sqrt{4 \theta})}\right) \tag{1.16}
\end{align*}
$$

with $\nu=\sqrt{4 \theta-1}$

- If $4 \theta \leq 1$ :

$$
Q_{l}^{(0)}\left(\exp -\frac{\theta}{2} \int_{1}^{\infty} \frac{x-1}{x^{2}} \lambda_{x-1} d x\right)=\exp \frac{l}{2}\left(\frac{1}{2}+\frac{\sqrt{\theta} K_{\nu}^{\prime}(\sqrt{4 \theta})}{K_{\nu}(\sqrt{4 \theta})}\right)
$$

with $\nu=\sqrt{1-4 \theta}$.
In Section 3 of this work, we shall study the asymptotic behavior, as $\theta \rightarrow \infty$, of $E\left(e^{-\frac{\theta}{2} \Pi_{1}}\right)$ and we shall obtain :

$$
\begin{equation*}
E\left(e^{-\frac{\theta}{2} \Pi_{1}}\right) \underset{\theta \rightarrow \infty}{\sim} \frac{C}{\theta^{\frac{1}{3}}} \quad(C>0) \tag{1.17}
\end{equation*}
$$

Finally, in a short Section 4, we shall indicate how Theorems 1.1, 1.2 and 1.3 extend when we replace $\Pi_{k}$ by $\Pi_{k}^{(\rho)}$, with :

$$
\begin{array}{rll}
\Pi_{k}^{(\rho)} & :=\int_{0}^{\infty}\left(e^{\rho\left(B_{t}-\frac{t}{2}\right)}-k\right)^{+} d t & (\rho, k>0) \\
& \stackrel{(\text { law })}{=} \frac{1}{\rho^{2}} \int_{0}^{\infty}\left(e^{\left(B_{u}-\frac{u}{2 \rho}\right)}-k\right)^{+} d u & \text { (by scaling) } \tag{1.18}
\end{array}
$$

which, from (1.18), corresponds to consider an extension of our previous perpetuities relative to $\left(B_{t}-\frac{t}{2}, t \geq 0\right)$ to Brownian motion with drift $-\frac{1}{2 \rho}$, i.e. $\left(B_{t}-\frac{t}{2 \rho}, t \geq 0\right)$.
1.8 For studies of other perpetuities related to Brownian motion with drift, we refer the reader to Salminen-Yor ([SY1], [SY2], [SY3]).

## 2 Proofs of Theorems 1.1, 1.2 and 1.3

### 2.1 A first proof of Theorem 1.1

2.1 i) We now prove that, for every $x>0, E_{x}\left(\Pi_{1}\right)=\infty$ Indeed, from (1.6) :

$$
\begin{align*}
E_{1}\left(\Pi_{1}\right) & =\int_{1}^{\infty} \frac{(y-1)}{y^{2}} E_{1}\left(L_{T_{0}}^{y}\right) d y \\
& =\int_{1}^{\infty} \frac{(y-1)}{y^{2}} E_{1}\left(L_{T_{0}}^{1}\right) d y \\
& =2 \int_{1}^{\infty} \frac{(y-1)}{y^{2}} d y=+\infty \tag{2.1}
\end{align*}
$$

(since, for $y \geq 1, E_{1}\left(L_{T_{0}}^{y}\right)=E_{1}\left(L_{T_{0}}^{1}\right)=2$, as $L_{T_{0}}^{1}$ is an exponential variable with parameter $\frac{1}{2}$ ).
2.1 ii) We now prove that, for every $x>0, E_{x}\left(\Pi_{1}\right)=\infty$

- For $x \geq 1$, this is clear since $\Pi_{k}^{(x)}$ is an increasing function of $x$ (and a decreasing function of $k$ ). Hence :

$$
E_{x}\left(\Pi_{1}\right) \geq E_{1}\left(\Pi_{1}\right)=+\infty \quad(\text { from }(2.1))
$$

- For $x<1$, from the Markov property :

$$
\begin{equation*}
\Pi_{1}^{(x)} \stackrel{(\operatorname{law})}{=} 1_{\left\{T_{\log \left(\frac{1}{x}\right)}^{\left(-\frac{1}{2}\right)}<\infty\right\}} \cdot \Pi_{1} \tag{2.2}
\end{equation*}
$$

with $T_{\log \left(\frac{1}{x}\right)}^{\left(-\frac{1}{2}\right)}:=\inf \left\{t \geq 0 ; B_{t}-\frac{t}{2}=\log \left(\frac{1}{x}\right)\right\}$, and in $(2.2), \Pi_{1}$ and $T_{\log \left(\frac{1}{x}\right)}^{\left(-\frac{1}{x}\right)}$ are assumed independent. Hence :

$$
E\left(\Pi_{1}^{(x)}\right)=P\left(T_{\log \left(\frac{1}{x}\right)}^{\left(-\frac{1}{x}\right)}<\infty\right) E\left(\Pi_{1}\right)=+\infty
$$

from (2.1) and since $P\left(T_{\log \left(\frac{1}{x}\right)}^{\left(-\frac{1}{2}\right)}<\infty\right)=x$ (see Lemma 2.1 below).
2.1 iii) We prove that, for $0 \leq \alpha<1, E_{x}\left(\left(\Pi_{1}\right)^{\alpha}\right)<\infty$

We have :

$$
\begin{aligned}
E_{x}\left(\left(\Pi_{1}\right)^{\alpha}\right) & =x^{\alpha} E\left(\left(\Pi_{\frac{1}{x}}\right)^{\alpha}\right) \quad(\text { from } \\
& \leq x^{\alpha} E\left(\left(\Pi_{0}\right)^{\alpha}\right)
\end{aligned}
$$

(since $\Pi_{k}$ is a decreasing function of $k$ )

$$
\leq x^{\alpha} \int_{0}^{\infty} \frac{2}{y^{\alpha}} e^{-y} d y<\infty \quad(\text { from }(1.12))
$$

### 2.2 Second proof of Theorem 1.1

## It hinges upon:

## Lemma 2.1

Let $G_{a}^{(\nu)}:=\sup \left\{u \geq 0 ; B_{u}+\nu u=a\right\}(=0 \quad$ if this set is empty) and
$T_{a}^{(\nu)}:=\inf \left\{u \geq 0 ; B_{u}+\nu u=a\right\}(=+\infty$ if this set is empty).
i) If $\nu$ and a have the same sign :

- $E\left(e^{\mu G_{a}^{(\nu)}}\right)<\infty \quad$ if and only if $\mu<\frac{\nu^{2}}{2}$
- for every real $\alpha, E\left(\left(G_{a}^{(\nu)}\right)^{\alpha}\right)<\infty$
ii) If $\nu$ and a have opposite signs:
- $E\left(e^{\mu G_{a}^{(\nu)}}\right)<\infty \quad$ if and only if $\mu<\frac{\nu^{2}}{2}$
- for every real $\alpha>0, E\left(\left(G_{a}^{(\nu)}\right)^{\alpha}\right)<\infty$
- $P\left(T_{a}^{(\nu)}=\infty\right)=P\left(G_{a}^{(\nu)}=0\right)=1-e^{2 \nu a}$

Hence, for $\alpha<0, E\left(\left(G_{a}^{(\nu)}\right)^{\alpha}\right)=+\infty$
The proof of this Lemma is obvious. It hinges on the well-known formulae:

- If $\nu, a>0$ :

$$
P\left(G_{a}^{(\nu)} \in d t\right)=\frac{\nu}{\sqrt{2 \pi t}} \exp \left(-\frac{1}{2 t}(a-\nu t)^{2}\right) d t
$$

- If $\nu, a>0$ :

$$
P\left(G_{a}^{(-\nu)} \in d t\right)=\frac{\nu}{\sqrt{2 \pi t}} \exp \left(-\frac{1}{2 t}(a+\nu t)^{2}\right) d t
$$

- If $\nu$ and $a$ have opposites signs :

$$
P\left(G_{a}^{(\nu)}>0\right)=P\left(T_{a}^{(\nu)}<\infty\right)=e^{2 \nu a}
$$

- $G_{a}^{(\nu)} \stackrel{(\text { law })}{=} G_{(-a)}^{(-\nu)} \stackrel{(\text { law })}{=} \frac{1}{T_{\nu}^{(a)}}$

We now give a second proof of Theorem 1.1.

- We first show that, for $0 \leq \alpha<1, E_{x}\left(\left(\Pi_{1}\right)^{\alpha}\right)<\infty$. From the relation :

$$
\begin{aligned}
\Pi_{1}^{(x)} & =x \int_{0}^{\infty}\left(\mathcal{E}_{t}-\frac{1}{x}\right)^{+} d t=x \int_{0}^{G_{-\log x}^{\left(-\frac{1}{2}\right)}}\left(\mathcal{E}_{t}-\frac{1}{x}\right)^{+} d t \\
& \leq x\left(\sup _{t \geq 0} \mathcal{E}_{t}\right) \cdot G_{-\log x}^{\left(-\frac{1}{2}\right)}
\end{aligned}
$$

we deduce that, for $0 \leq \alpha<1, \alpha p<1$ and $\frac{1}{p}+\frac{1}{q}=1(p, q>1)$ :

$$
\begin{align*}
E_{x}\left(\left(\Pi_{1}\right)^{\alpha}\right) & \leq x^{\alpha} E\left[\left(\sup _{t \geq 0} \mathcal{E}_{t}\right)^{\alpha}\left(G_{-\log x}^{\left(-\frac{1}{2}\right)}\right)^{\alpha}\right] \\
& \leq x^{\alpha} E\left(\left(\sup _{t \geq 0} \mathcal{E}_{t}\right)^{\alpha p}\right)^{\frac{1}{p}} E\left(\left(G_{(-\log x)}^{\left(-\frac{1}{2}\right)}\right)^{\alpha q}\right)^{\frac{1}{q}} \\
& \leq C x^{\alpha} E\left(\left(\sup _{t \geq 0} \mathcal{E}_{t}\right)^{\alpha p}\right)^{\frac{1}{p}} \quad \text { (from Lemma } \tag{fromLemma2.1}
\end{align*}
$$

As $\left(\mathcal{E}_{t}, t \geq 0\right)$ is a positive martingale, starting from 1 , and converging a.s. to 0 as $t \rightarrow \infty$, we get (see [NY] or [RY], Chap. II, Ex. 3.12, p. 73) :

$$
\begin{equation*}
\sup _{s \geq 0} \mathcal{E}_{s} \stackrel{(\text { law })}{=} \frac{1}{U} \tag{2.8}
\end{equation*}
$$

with $U$ uniform on $[0,1]$. Thus :

$$
\begin{equation*}
E\left(\left(\sup _{t \geq 0} \mathcal{E}_{t}\right)^{\alpha p}\right)=\int_{0}^{1} \frac{1}{u^{\alpha p}} d u<\infty \tag{2.9}
\end{equation*}
$$

since $\alpha p<1$.

- We then show that $E_{x}\left(\Pi_{1}\right)=\infty$. First of all, it follows from [MRY1] that :

$$
\begin{equation*}
E_{1}\left(\left(\mathcal{E}_{t}-k\right)^{+}\right)=P\left(G_{\log k}^{\left(\frac{1}{2}\right)} \leq t\right) \tag{2.10}
\end{equation*}
$$

Hence :

$$
E_{1}\left(\Pi_{1}\right)=\int_{0}^{\infty} P\left(G_{0}^{\left(\frac{1}{2}\right)} \leq t\right) d t=E\left(\int_{G_{0}^{\left(\frac{1}{2}\right)}}^{\infty} d t\right)=+\infty
$$

since $G_{0}^{\left(\frac{1}{2}\right)}<+\infty$ a.s.
Likewise :

$$
\begin{aligned}
E_{x}\left(\Pi_{1}\right) & =x E_{1}\left(\Pi_{\frac{1}{x}}\right) \quad(\text { from }(1.9)) \\
& =x E\left(\int_{G_{\log \frac{1}{x}}^{\left(\frac{1}{2}\right)}}^{\infty} d t\right)=\infty
\end{aligned}
$$

from (2.10) and since $G_{\log \frac{1}{x}}^{\left(\frac{1}{2}\right)}<\infty$ a.s.

### 2.3 Proof of Theorem 1.2

i) We already prove point $i$ ):

Let $x>1$ and $x^{\prime}$ such that $1<x^{\prime}<x$. It is then obvious that:

$$
\begin{equation*}
\Pi_{1}^{(x)} \geq\left(x^{\prime}-1\right) T_{\log \left(\frac{x^{\prime}}{x}\right)}^{\left(-\frac{1}{x}\right)} \tag{2.11}
\end{equation*}
$$

(with $T_{a}^{\left(-\frac{1}{2}\right)}:=\inf \left\{t \geq 0 ; B_{t}-\frac{t}{2}=a\right\}$, since, if $t<T_{\log \frac{x^{\prime}}{x}}^{\left(-\frac{1}{2}\right)}$ then $e^{\log x+B_{t}-\frac{t}{2}}-1 \geq x^{\prime}-1$ ). Hence, with $\gamma>0$ :

$$
\begin{align*}
E_{x}\left(\frac{1}{\left(\Pi_{1}\right)^{\gamma}}\right) & =\int_{0}^{\infty} P_{x}\left(\frac{1}{\left(\Pi_{1}\right)^{\gamma}} \geq t\right) d t=\int_{0}^{\infty} P_{x}\left(\Pi_{1}<v\right) \frac{\gamma}{v^{1+\gamma}} d v \\
& \leq \int_{0}^{\infty} \frac{\gamma}{v^{1+\gamma}} d v P\left(T_{\log \frac{x^{\prime}}{x}}^{\left(-\frac{1}{x}\right)} \leq \frac{v}{x^{\prime}-1}\right) d v \quad(\text { from }(2.11)) \\
& =\int_{0}^{\infty} \frac{\gamma}{v^{1+\gamma}} d v\left(P\left(G_{\frac{1}{2}}^{\left(\log \left(\frac{x}{x^{\prime}}\right)\right)} \geq \frac{x^{\prime}-1}{v}\right) d v\right. \\
(\text { since } & \left.T_{a}^{(\nu)} \stackrel{(\operatorname{law})}{=} \frac{1}{G_{\nu}^{(a)}} \text { and } G_{\nu}^{(a)} \stackrel{(\operatorname{law})}{=} G_{-\nu}^{(-a)}\right) \\
& =\int_{0}^{\infty} \frac{\gamma}{\left(x^{\prime}-1\right)^{\gamma}} u^{\gamma-1} d u P\left(G_{\frac{1}{2}}^{\left(\log \left(\frac{x}{x^{\prime}}\right)\right)} \geq u\right) d u \\
& =\frac{1}{\left(x^{\prime}-1\right)^{\gamma}} E\left(\left(G_{\frac{1}{2}}^{\left(\log \frac{x}{x^{\prime}}\right)}\right)^{\gamma}\right)<\infty \tag{2.12}
\end{align*}
$$

(from point $i$ ) of Lemma 2.1).
ii) We now prove point ii) :

It is clear that, for $x<1$ :

$$
\left\{\Pi_{1}^{(x)}=0\right\}=\left\{T_{\log \frac{1}{x}}^{\left(-\frac{1}{2}\right)}=\infty\right\}
$$

Thus, $P\left\{\Pi_{1}^{(x)}=0\right\}=P\left\{T_{\log \frac{1}{x}}^{\left(-\frac{1}{2}\right)}=\infty\right\}=P\left\{G_{\frac{1}{2}}^{(\log x)}=0\right\}=1-x>0$ from Lemma 2.1 and (1.26).
iii) We now prove point iii):

For this purpose, we write, for $\gamma>0$ :

$$
\begin{equation*}
E_{1}\left(\frac{1}{\left(\Pi_{1}\right)^{\gamma}}\right)=\frac{1}{\Gamma(\gamma)} \int_{0}^{\infty} E\left(e^{-\theta \Pi_{1}}\right) \theta^{\gamma-1} d \theta \tag{2.13}
\end{equation*}
$$

and we show, in the next Section 3, that :

$$
E_{1}\left(e^{-\theta \Pi_{1}}\right) \underset{\theta \rightarrow \infty}{\sim} \frac{C}{\theta^{\frac{1}{3}}} \quad(C>0)
$$

Thus, $E\left(\frac{1}{\left(\Pi_{1}\right)^{\gamma}}\right)<\infty$ if and only if $\int_{1}^{\infty} \theta^{\gamma-1-\frac{1}{3}} d \theta<\infty$, that is, if and only if $\gamma<\frac{1}{3}$.

### 2.4 Proof of Theorem 1.3

a) A useful Lemma

Lemma 2.2 Let $\theta \geq 0$ and $K_{\gamma}$ the Bessel McDonald function with index $\gamma$, such that $\gamma=\sqrt{1-4 \theta}$ if $\theta \leq \frac{1}{4}$ and $\gamma=i \sqrt{4 \theta-1}$ if $\theta \geq \frac{1}{4}$.

1) Define the function $\varphi_{\theta}: \mathbb{R}_{+} \rightarrow \mathbb{R}$ by :

$$
\begin{equation*}
\varphi_{\theta}(y)=\sqrt{y} K_{\gamma}(\sqrt{4 \theta y}) \quad(y \geq 0) \tag{2.14}
\end{equation*}
$$

Then :
i) $\varphi_{\theta}$ is a real valued function which satisfies :

$$
\begin{equation*}
\varphi_{\theta}^{\prime \prime}(y)+\left(-\frac{\theta}{y}+\frac{\theta}{y^{2}}\right) \varphi_{\theta}(y)=0 \tag{2.15}
\end{equation*}
$$

ii) $\varphi_{\theta}$, restricted to $[1, \infty[$ is positive, convex, bounded and decreasing.
2) We define the function $\widetilde{\varphi}_{\theta}: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$by :

$$
\widetilde{\varphi}_{\theta}(y)= \begin{cases}\varphi_{\theta}(y) & \text { if } y \geq 1  \tag{2.16}\\ \left(\varphi_{\theta}(1)-\varphi_{\theta}^{\prime}(1)\right)+y \varphi_{\theta}^{\prime}(1) & \text { if } 0 \leq y \leq 1\end{cases}
$$

Then $\widetilde{\varphi}_{\theta}$ is a bounded, positive, convex, decreasing function which satisfies :

$$
\begin{equation*}
\widetilde{\varphi}_{\theta}^{\prime \prime}(y)+\left(-\frac{\theta}{y}+\frac{\theta}{y^{2}}\right) 1_{y \geq 1} \cdot \widetilde{\varphi}_{\theta}(y)=0 \tag{2.17}
\end{equation*}
$$

## b) Proof of Lemma 2.2

i) Relation (2.15) (as well as relation (2.17)) follows from a direct computation, using the equation $K_{\gamma}^{\prime \prime}(z)+\frac{1}{z} K_{\gamma}^{\prime}(z)-\left(1+\frac{\gamma^{2}}{z^{2}}\right) K_{\gamma}(z)=0$ (see [Leb], p. 110) and the fact that $\gamma^{2}=1-4 \theta$ (see Petiau [Pet], p. 306, formula (8), which needs to be corrected by replacing $a$ by $-a$, or Kamke [Kam], p. 440).
We distinguish two cases :
Case 1: $4 \theta \leq 1, \gamma=\sqrt{1-4 \theta}$. The function $K_{\gamma}$, hence also $\varphi_{\theta}$, is positive. Furthermore $\varphi_{\theta}$ is bounded on $\mathbb{R}_{+}$since :

$$
\varphi_{\theta}(y) \underset{y \rightarrow 0}{\sim} C y^{\frac{1-\gamma}{2}}, \varphi_{\theta}(y) \underset{y \rightarrow \infty}{\sim} C^{\prime} y^{\frac{1}{4}} e^{-\sqrt{4 \theta y}} \quad(\text { see [Leb], p. } 123 \text { and 136) }
$$

On the other hand, from (2.15), the function $\varphi_{\theta}$ is convex on the interval $[1, \infty[$. As it is convex, positive, and bounded, it is decreasing. Lemma 2.2 is thus proven in this case.


Figure 1: Graphs of $\varphi_{\theta}$ and $\widetilde{\varphi}_{\theta}$
Case 2: $4 \theta \geq 1, \gamma=i \sqrt{4 \theta-1}$ and here $\varphi_{\theta}(y)=\sqrt{y} K_{i \nu}(\sqrt{4 \theta y})$ with $\nu=\sqrt{4 \theta-1}$. From the integral representation formula (see [Leb], p. 119) :

$$
\begin{equation*}
K_{i \nu}(y)=\int_{0}^{\infty} e^{-y \cosh u} \cos (\nu u) d u \quad(y \geq 0) \tag{2.18}
\end{equation*}
$$

we deduce that $K_{i \nu}(y)$ is real valued, hence so is $\varphi_{\theta}(y)$.
On the other hand, $\varphi_{\theta}$ is bounded on $[0, \infty[$. Indeed, for $y \geq 0$, from (2.18) :

$$
\left|\sqrt{4 \theta} \varphi_{\theta}\left(\frac{y^{2}}{4 \theta}\right)\right|=y\left|K_{i \nu}(y)\right| \leq y \int_{0}^{\infty} e^{-y \frac{e^{u}}{2}} d u \leq y \int_{\frac{y}{2}}^{\infty} e^{-v} \frac{d v}{v}
$$

Thus:

$$
\left\{\begin{array}{l}
\left.\sqrt{4 \theta} \varphi_{\theta}\left(\frac{y^{2}}{4 \theta}\right) \right\rvert\, \leq e^{-\frac{y}{4}} \quad \text { for } y \text { large enough, hence } \\
\left.\sqrt{4 \theta} \varphi_{\theta}\left(\frac{y^{2}}{4 \theta}\right) \right\rvert\, \underset{y \rightarrow \infty}{\longrightarrow} \quad \text { and } \\
\left|\sqrt{4 \theta} \varphi_{\theta}\left(\frac{y^{2}}{4 \theta}\right)\right| \leq y(C+|\log (y)|) \underset{y \rightarrow 0}{\longrightarrow} 0
\end{array}\right.
$$

On the other hand, it is clear, from (2.18), that $K_{i \nu}(y)>0$ hence $\varphi_{\theta}(y)>0$, for $y$ large enough.

- We now show that $K_{i \nu}$ is decreasing on $\left[\nu, \infty\left[\right.\right.$. If not, there would exist a point $y_{0}>\nu$ such that:

$$
K_{i \nu}\left(y_{0}\right)>0, K_{i \nu}^{\prime}\left(y_{0}\right)=0, K_{i \nu}^{\prime \prime}\left(y_{0}\right) \leq 0
$$

However,

$$
K_{i \nu}^{\prime \prime}\left(y_{0}\right)+\frac{1}{y_{0}} K_{i \nu}^{\prime}\left(y_{0}\right)=\left(1-\frac{\nu^{2}}{y_{0}^{2}}\right) K_{i \nu}\left(y_{0}\right)
$$

Thus $K_{i \nu}^{\prime \prime}\left(y_{0}\right)>0$, which is absurd. Since $K_{i \nu}$ is decreasing on $[\nu, \infty[$, and positive near $\infty$, $K_{i \nu}$ is positive on $\left[\nu, \infty\left[\right.\right.$. Thus, $\varphi_{\theta}$ is positive on $[1, \infty[$ (since, if $y \geq 1, \sqrt{4 \theta y} \geq \sqrt{4 \theta} \geq$ $\sqrt{4 \theta-1}=\nu$ ). From the relation (2.15), we then deduce that $\varphi_{\theta}$ is convex on $[1, \infty[$. Since it is bounded, convex and positive, it is decreasing. Lemma 2.2 is proven.

## c) End of the proof of Theorem 1.3.

Let $\left(M_{t}^{\theta}:=\widetilde{\varphi}_{\theta}\left(B_{t}\right) \exp -\frac{\theta}{2} \int_{0}^{t} \frac{\widetilde{\varphi}_{\theta}^{\prime \prime}}{\varphi_{\theta}}\left(B_{s}\right) d s, t \geq 0\right)$. Then $\left(M_{t}^{\theta}, t \geq 0\right)$ is a local martingale. It is equal, from (2.17) to :

$$
\begin{equation*}
M_{t}^{\theta}=\widetilde{\varphi}_{\theta}\left(B_{t}\right) \exp -\frac{\theta}{2} \int_{0}^{t} \frac{\left(B_{s}-1\right)^{+}}{B_{s}^{2}} d s \tag{2.19}
\end{equation*}
$$

and, from Lemma 2.2, for every $x \geq 0,\left(M_{t \wedge T_{0}}^{\theta}, t \geq 0\right)$ is bounded. Thus, from Doob's optional stopping Theorem :

$$
\begin{aligned}
\widetilde{\varphi}_{\theta}(x) & =E_{x}\left(M_{0}^{\theta}\right)=E_{x}\left[\widetilde{\varphi}_{\theta}\left(B_{T_{0}}\right) \exp -\frac{\theta}{2} \int_{0}^{T_{0}(B)} \frac{\left(B_{s}-1\right)^{+}}{B_{s}^{2}} d s\right] \\
& =\widetilde{\varphi}_{\theta}(0) E_{x}\left[\exp -\frac{\theta}{2} \Pi_{1}\right] \quad(\text { from }(1.5))
\end{aligned}
$$

Thus :

$$
E_{x}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right)=\frac{\widetilde{\varphi}_{\theta}(x)}{\widetilde{\varphi}_{\theta}(0)}
$$

This is precisely Theorem 1.3, owing to formula (2.16) which yields $\widetilde{\varphi}_{\theta}$ explicitly.
Remark 2.3 Theorem 1.3 allows to recover formula (1.9) : $\Pi_{0} \stackrel{(\text { law })}{=} \frac{2}{\mathfrak{e}}$. Indeed, on one hand :

$$
E_{1}\left(e^{-\frac{\theta}{2} \Pi_{0}}\right)=E\left(e^{-\frac{\theta}{2} \frac{2}{c}}\right)=\int_{0}^{\infty} e^{-\frac{\theta}{z}-z} d z=2 \sqrt{\theta} K_{1}(\sqrt{4 \theta}) \quad(\text { see [Leb], p. 119) }
$$

and, on the other hand :

$$
\begin{aligned}
E_{1}\left(e^{-\frac{\theta}{2} \Pi_{0}}\right)= & \lim _{\epsilon \downarrow 0} E_{1}\left(e^{-\frac{\theta}{2} \Pi_{\epsilon}}\right)=\lim _{\epsilon \downarrow 0} E\left(\exp -\frac{\theta}{2} \int_{0}^{\infty}\left(e^{B_{t}-\frac{t}{2}}-\epsilon\right)^{+} d t\right) \\
= & \lim _{\epsilon\rfloor 0} E\left(e^{-\frac{\theta \epsilon}{2} \int_{0}^{\infty}\left(\frac{1}{\epsilon}\left(e^{\left.B_{t}-\frac{t}{2}-1\right)^{+} d t}\right)=\lim _{\epsilon] 0} E_{\frac{1}{\epsilon}}\left(e^{-\frac{\theta \epsilon}{2} \Pi_{1}}\right)\right.}=\right. \\
= & \lim _{\epsilon \downarrow 0} \frac{\frac{1}{\sqrt{\epsilon}} K_{\sqrt{1-4 \theta \epsilon}}(\sqrt{4 \theta})}{\frac{1}{2} K_{\sqrt{1-4 \theta \epsilon}}(\sqrt{4 \theta \epsilon})-\frac{1}{2} \sqrt{4 \theta \epsilon} K_{\sqrt{1-4 \theta \epsilon}}^{\prime}(\sqrt{4 \theta \epsilon})} \\
& \left(\text { from }(1.5), \text { replacing } \theta \text { by } \theta \epsilon \text { and } x \text { by } \frac{1}{\epsilon}\right) \\
= & \lim _{\epsilon\rfloor 0} \frac{\frac{1}{\sqrt{\epsilon}} K_{1}(\sqrt{4 \theta})+O(1)}{\frac{1}{2} K_{\sqrt{1-4 \theta \epsilon}}(\sqrt{4 \theta \epsilon})[1-\sqrt{1-4 \theta \epsilon}]+\sqrt{\theta \epsilon} K_{\sqrt{1-4 \theta \epsilon}+1}(\sqrt{4 \theta \epsilon})} \\
& \left(\operatorname{since} z K_{\nu}^{\prime}(z)=\nu K_{\nu}(z)-z K_{\nu+1}(z) ; \text { see }[\text { Leb], p. 110) }\right. \\
= & \lim _{\epsilon \downarrow 0} \frac{\frac{1}{\sqrt{\epsilon}} K_{1}(\sqrt{4 \theta})+O(1)}{O(1)+\sqrt{\theta \epsilon}\left(\frac{1}{2 \theta \epsilon}+O(1)\right)}=2 \sqrt{\theta} K_{1}(\sqrt{4 \theta})
\end{aligned}
$$

### 2.5 Proof of Theorem 1.4

## From Theorem 1.3, we know that, for $\theta \geq \frac{1}{4}$ :

$$
\begin{align*}
E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right) & =\frac{2 K_{i \nu}(\sqrt{4 \theta})}{K_{i \nu}(\sqrt{4 \theta})-\sqrt{4 \theta} K_{i \nu}^{\prime}(\sqrt{4 \theta})} \\
& =\frac{1}{\frac{1}{2}-\sqrt{\theta} \frac{K_{i \nu}^{\prime}(\sqrt{4 \theta})}{K_{i \nu}(\sqrt{4 \theta})}}=\frac{1}{2} \int_{0}^{\infty} e^{-\frac{l}{2}+\frac{l}{2}\left(\frac{1}{2}+\sqrt{\theta} \frac{K_{i \nu}^{\prime}(\sqrt{4 ब})}{K_{i \nu}(\sqrt{4 \theta})}\right)} d l \tag{2.20}
\end{align*}
$$

On the other hand, from (1.7) :

$$
\begin{align*}
E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right) & =E_{1}\left(\exp -\frac{\theta}{2} \int_{1}^{\infty} \frac{(y-1)^{+}}{y^{2}} \lambda_{y-1} d y\right) \\
& =\frac{1}{2} \int_{0}^{\infty} e^{-\frac{l}{2}} d l E\left(\left.\exp -\frac{\theta}{2} \int_{1}^{\infty} \frac{(y-1)^{+}}{y^{2}} \lambda_{y-1} d y \right\rvert\, \lambda_{0}=l\right) \\
& =\frac{1}{2} \int_{0}^{\infty} e^{-\frac{l}{2}} d l Q_{l}^{(0)}\left(\exp -\frac{\theta}{2} \int_{0}^{\infty} \frac{y}{(1+y)^{2}} \lambda_{y} d y\right) \tag{2.21}
\end{align*}
$$

where $Q_{l}^{(0)}$ denotes the expectation relative to a squared Bessel process of dimension 0 , starting from $l$. The comparison of (2.21) and (2.20) implies Theorem 1.4 in the case
$\theta \geq \frac{1}{4}$, since the Laplace transform is one-to-one. The proof of Theorem 1.4 in the case $\theta \leq \frac{1}{4}$ is the same.
Remark 2.4 Let $\mu$ denote a positive $\sigma$-finite measure on $\mathbb{R}_{+}$. Let $\Phi$ denote the unique decreasing, positive solution on $\mathbb{R}_{+}$of the Sturm-Liouville equation $\Phi^{\prime \prime}=\mu \Phi$ (and such that $\Phi(0)=1$ ). It is well known (see [RY], Chap. IX, p. 444) that :

$$
\begin{equation*}
Q_{l}^{(0)}\left(\exp \left(-\frac{1}{2} \int_{0}^{\infty} \lambda_{y} \mu(d y)\right)\right)=\exp \left(\frac{l}{2} \Phi^{\prime}\left(0_{+}\right)\right) \tag{2.22}
\end{equation*}
$$

(Observe that, as $\left(\lambda_{y}, y \geq 0\right)$ has compact support a.s., $\int_{0}^{\infty} \lambda_{y} \mu(d y)<\infty$ if $\mu$ is $\sigma$-finite on $\mathbb{R}_{+}$).
Theorem 1.4 may be recovered easily by applying formula (2.22) with $\mu(d y)=\frac{y}{1+y^{2}} d y$. On the other hand, if $\mu(d y)=a(y) d y$, with

$$
c_{1} 1_{\left[0, \gamma_{1}\right]}(y) \leq a(y) \leq c_{2} 1_{\left[0, \gamma_{2}\right]}(y) \quad\left(0<c_{1} \leq c_{2}, 0<\gamma_{1} \leq \gamma_{2}\right),
$$

we deduce from [RY], Chap. XI, Corollary 1.8, that :

$$
\begin{equation*}
\exp -\frac{l}{2} c_{2}^{\prime} \sqrt{\theta} \leq E_{l}^{(0)}\left(e^{-\frac{\theta}{2} \int_{0}^{\infty} \lambda_{y} \mu(d y)}\right) \leq \exp -\frac{l}{2} c_{1}^{\prime} \sqrt{\theta} \tag{2.23}
\end{equation*}
$$

for $\theta$ large enough, whereas, as we shall show in Section 3 :

$$
\begin{equation*}
Q_{l}^{(0)}\left(\exp -\frac{\theta}{2} \int_{0}^{\infty} \frac{y}{1+y^{2}} \lambda_{y} d y\right) \underset{\theta \rightarrow \infty}{\sim} C e^{-\frac{l}{2} \theta^{\frac{1}{3}}} \tag{2.24}
\end{equation*}
$$

## 3 Asymptotic behavior of $E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right)$ as $\theta \rightarrow \infty$

We shall now end the proof of Theorem 1.2 by showing :
Theorem 3.1 There is the equivalence result:

$$
\begin{equation*}
E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right) \underset{\theta \rightarrow \infty}{\sim} \frac{C}{\theta^{\frac{1}{3}}} \tag{3.1}
\end{equation*}
$$

## Proof of Theorem 3.1

We recall that, from Theorem 1.3 :

$$
\begin{equation*}
E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right)=\frac{2 K_{i \nu}(\sqrt{4 \theta})}{K_{i \nu}(\sqrt{4 \theta})-\sqrt{4 \theta} K_{i \nu}^{\prime}(\sqrt{4 \theta})} \tag{3.2}
\end{equation*}
$$

with $\nu=\sqrt{4 \theta-1}\left(\theta \geq \frac{1}{4}\right)$. We shall successively find an equivalent of the numerator and the denominator of (3.2), the difficulty arising from the fact that, in $K_{i \nu}(\sqrt{4 \theta})$ (and
$\left.K_{i \nu}^{\prime}(\sqrt{4 \theta})\right)$ the argument $\sqrt{4 \theta}$ and index $i \nu=i \sqrt{4 \theta-1}$ tend both to infinity as $\theta \rightarrow \infty$. To overcome this difficulty, we shall use some results about Bessel functions found in Watson ([Wat], p. 245-248), which we now recall.
i) Let $H_{i \nu}^{(1)}$ the first Hankel function (see [Leb], p. 120) ; it is related to $K_{i \nu}$ via the formula :

$$
\begin{equation*}
K_{i \nu}(z)=\frac{i \pi}{2} e^{-\frac{\nu \pi}{2}} H_{i \nu}^{(1)}\left(z e^{\frac{i \pi}{2}}\right) \tag{3.3}
\end{equation*}
$$

We define $\epsilon$ by the formula : $i \nu=i z(1-\epsilon)$ and assume that, as $z \rightarrow \infty, \epsilon$ remains bounded (with, of course, $\nu$ depending on $z$ ). Then, there is the second order asymptotic expansion :

$$
\begin{equation*}
H_{\nu}^{(1)}(z)=-\frac{2}{3 \pi}\left\{e^{\frac{2}{3} \pi i}\left(\sin \frac{\pi}{3}\right) \frac{\Gamma\left(\frac{1}{3}\right)}{\left(\frac{1}{6} z\right)^{\frac{1}{3}}}+e^{\frac{4}{3} \pi i}(\epsilon z)\left(\sin \frac{2 \pi}{3}\right) \frac{\Gamma\left(\frac{2}{3}\right)}{\left(\frac{1}{6} z\right)^{\frac{2}{3}}}+o\left(\frac{1}{z^{\frac{2}{3}}}\right)\right\} \tag{3.4}
\end{equation*}
$$

ii) Let us study the numerator of (3.2) :

$$
\begin{align*}
& N= 2 K_{i \nu}(\sqrt{4 \theta})=2 \frac{i \pi}{2} e^{-\frac{\nu \pi}{2}} H_{i \nu}^{(1)}\left(e^{\frac{i \pi}{2}} \sqrt{4 \theta}\right) \\
&=(i \pi) e^{-\frac{\nu \pi}{2}}\left(-\frac{2}{3 \pi}\right)\left\{C_{1} e^{\frac{2}{3} \pi i} \frac{1}{\left(e^{\frac{i \pi}{2}} \sqrt{4 \theta}\right)^{\frac{1}{3}}}+o\left(\frac{1}{\theta^{\frac{1}{6}}}\right)\right\}  \tag{3.5}\\
& \quad\left(\text { with } C_{1}=\left(\sin \frac{\pi}{3}\right) \Gamma\left(\frac{1}{3}\right) 6^{\frac{1}{3}}\right) \\
& \underset{\theta \rightarrow \infty}{\sim} \frac{2^{\frac{2}{3}}}{3} C_{1} e^{-\frac{\nu \pi}{2}} \frac{1}{\theta^{\frac{1}{6}}} \tag{3.6}
\end{align*}
$$

Here, we have used the first order expansion (3.4) with : $e^{\frac{2}{3} i \pi} \cdot e^{-\frac{i \pi}{6}}=i\left(\frac{4 \pi}{6}-\frac{\pi}{6}=\frac{3 \pi}{6}=\frac{\pi}{2}\right)$ and the fact that here

$$
\begin{aligned}
\epsilon z & =z-\nu \\
& =i \sqrt{4 \theta}-i \sqrt{4 \theta-1}=i \sqrt{4 \theta}\left(1-\sqrt{1-\frac{1}{4 \theta}}\right) \underset{\theta \rightarrow \infty}{\longrightarrow} 0
\end{aligned}
$$

iii) We now study the denominator of (3.2)

$$
\begin{align*}
D & =K_{i \nu}(\sqrt{4 \theta})-\sqrt{4 \theta} K_{i \nu}^{\prime}(\sqrt{4 \theta}) \\
& =K_{i \nu}(\sqrt{4 \theta})+\sqrt{4 \theta} K_{i \nu-1}(\sqrt{4 \theta})+i \sqrt{4 \theta-1} K_{i \nu}(\sqrt{4 \theta}) \tag{3.7}
\end{align*}
$$

(after using $\frac{d}{d z}\left(z^{\mu} K_{\mu}(z)=-z^{\mu} K_{\mu-1}(z)\right) ;$ Leb], p. 110). Since we have already studied the asymptotic behavior of $K_{i \nu}(\sqrt{4 \theta})$, it remains to study that of

$$
\begin{equation*}
\Delta(\theta):=\sqrt{4 \theta} K_{i \nu-1}(\sqrt{4 \theta})+i \sqrt{4 \theta-1} K_{i \nu}(\sqrt{4 \theta}) \tag{3.8}
\end{equation*}
$$

Now, since $\quad \sqrt{4 \theta} K_{i \nu-1}(\sqrt{4 \theta})+i \sqrt{4 \theta-1} K_{i \nu}(\sqrt{4 \theta}) \quad$ is real (from (2.18)) and since $i \sqrt{4 \theta-1} K_{i \nu}(\sqrt{4 \theta})$ is purely imaginary,
the development of $\Delta(\theta)$ as $\theta \rightarrow \infty$ is that of the real part of $\sqrt{4 \theta} K_{i \nu-1}(\sqrt{4 \theta})$. Then, we obtain :

$$
\begin{align*}
\sqrt{4 \theta} K_{i \nu-1}(\sqrt{4 \theta}) & =\sqrt{4 \theta} \frac{i \pi}{2} e^{\frac{i \pi}{2}(i \nu-1)} H_{i \nu-1}^{(1)}\left(e^{\frac{i \pi}{2}} \sqrt{4 \theta}\right) \\
& =\sqrt{4 \theta} \frac{\pi}{2} e^{-\frac{\nu \pi}{2}} H_{i \nu-1}^{(1)}\left(e^{\frac{i \pi}{2}} \sqrt{4 \theta}\right) \tag{3.9}
\end{align*}
$$

from (3.3), by replacing $i \nu$ by $i \nu-1$. This time, we use the second order development (3.4) with here : $z=e^{\frac{i \pi}{2}} \sqrt{4 \theta}$ and :

$$
\begin{align*}
i \nu-1=z(1-\epsilon) & =i \sqrt{4 \theta}(1-\epsilon), \text { i.e. : } \quad \epsilon z=\epsilon e^{\frac{i \pi}{2}} \sqrt{4 \theta} \\
\text { or } \quad \epsilon z & =\epsilon e^{\frac{i \pi}{2}} \sqrt{4 \theta}=1+\frac{1}{4 \sqrt{\theta}}+o\left(\frac{1}{\sqrt{\theta}}\right) \quad(\nu=\sqrt{4 \theta-1}) \tag{3.10}
\end{align*}
$$

We obtain :

$$
\begin{align*}
& \sqrt{4 \theta} K_{i \nu-1}(\sqrt{4 \theta})=\sqrt{4 \theta} \frac{\pi}{2} e^{-\frac{\nu \pi}{2}}\left(-\frac{2}{3} \pi\right) \times \cdots \\
& \quad \ldots\left\{C_{1} e^{\frac{2 i \pi}{3}} \frac{1}{\left(e^{\frac{i \pi}{2}} \sqrt{4 \theta}\right)^{\frac{1}{3}}}+C_{2} e^{\frac{4 i \pi}{3}}\left(1+\frac{i}{4 \sqrt{\theta}}\right) \cdot \frac{1}{\left(e^{\frac{i \pi}{2}} \sqrt{4 \theta}\right)^{\frac{2}{3}}}+o\left(\frac{1}{\theta^{\frac{1}{3}}}\right)\right\} \tag{3.11}
\end{align*}
$$

with $C_{1}=\left(\sin \frac{\pi}{3}\right) \Gamma\left(\frac{1}{3}\right) 6^{\frac{1}{3}}$ and $C_{2}=\left(\sin \frac{2 \pi}{3}\right) \Gamma\left(\frac{2}{3}\right) 6^{\frac{2}{3}} \quad\left(C_{1}, C_{2}>0\right)$

$$
\begin{equation*}
=-\frac{1}{3} \sqrt{4 \theta} e^{-\frac{\nu \pi}{2}}\left\{C_{1} i \frac{1}{(\sqrt{4 \theta})^{\frac{1}{3}}}-C_{2}\left(1+\frac{i}{4 \sqrt{\theta}}\right) \frac{1}{\left(\sqrt{4 \theta^{\frac{2}{3}}}\right.}+o\left(\frac{1}{\theta^{\frac{1}{3}}}\right)\right\} \tag{3.12}
\end{equation*}
$$

Considering now the real part of (3.12) to obtain :

$$
\begin{align*}
\Delta(\theta) & =\mathcal{R} e\left(\sqrt{4 \theta} K_{i \nu-1}(\sqrt{4 \theta})+i \sqrt{4 \theta-1} K_{i \nu}(\sqrt{4 \theta})\right) \\
& =\frac{1}{3} \sqrt{4 \theta} e^{-\frac{\nu \pi}{2}} C_{2}\left(\frac{1}{(\sqrt{4 \theta})^{\frac{2}{3}}}+o\left(\frac{1}{\theta^{\frac{1}{3}}}\right)\right) \\
& \sim \frac{2^{\frac{1}{3}} C_{2}}{3} e^{-\frac{\nu \pi}{2}} \theta^{\frac{1}{6}} \tag{3.13}
\end{align*}
$$

- We then gather $(3.2),(3.6)$ and (3.13), we obtain :

$$
\begin{align*}
& E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right) \underset{\theta \rightarrow \infty}{\sim} \frac{2 \frac{2^{\frac{2}{3}}}{3}}{\sim} C_{1} e^{-\frac{\nu \pi}{2}} \frac{1}{\theta^{\frac{1}{6}}} \\
& \frac{2^{\frac{2}{3}}}{3} C_{1} e^{-\frac{\nu \pi}{2}} \frac{1}{\theta^{\frac{1}{6}}}+\frac{2^{\frac{1}{3}}}{3} C_{2} e^{-\frac{\nu \pi}{2}} \theta^{\frac{1}{6}}  \tag{3.14}\\
& \underset{\theta \rightarrow \infty}{\sim} \frac{2^{\frac{4}{3}} C_{1}}{C_{2}} \frac{1}{\theta^{\frac{1}{3}}}=\frac{2}{3^{\frac{1}{3}}} \frac{\Gamma\left(\frac{1}{3}\right)}{\Gamma\left(\frac{2}{3}\right)} \frac{1}{\theta^{\frac{1}{3}}}
\end{align*}
$$

from the explicit formulae for $C_{1}$ and $C_{2}$. This is Theorem 3.1.
Remark 3.2 Using Theorem 1.3, we obtain, for $0<x \leq 1$ :

$$
E_{x}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right)=1+x \frac{\frac{1}{2} K_{i \sqrt{4 \theta-1}}(\sqrt{4 \theta})+\sqrt{\theta} K_{i \sqrt{4 \theta-1}}^{\prime}(\sqrt{4 \theta})}{\frac{1}{2} K_{i \sqrt{4 \theta-1}}(\sqrt{4 \theta})-\sqrt{\theta} K_{i \sqrt{4 \theta-1}}^{\prime}(\sqrt{4 \theta})}
$$

It now follows easily from (3.6) and (3.13) that :

$$
\begin{equation*}
E_{x}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right) \underset{\theta \rightarrow \infty}{\longrightarrow} 1-x \tag{3.15}
\end{equation*}
$$

Now, on the other hand :

$$
\begin{equation*}
E_{x}\left(e^{-\frac{\theta}{2} \Pi_{1}}\right) \underset{\theta \rightarrow \infty}{\longrightarrow} P_{x}\left(\Pi_{1}=0\right) \tag{3.16}
\end{equation*}
$$

We recall (and recover here) that $P_{x}\left(\Pi_{1}=0\right)=1-x$. This is point $i i$ ) of Theorem 1.2.

## 4 Extending the preceding results to the variables $\Pi_{k}^{(\rho, x)}$ with

$$
\begin{equation*}
\Pi_{k}^{(\rho, x)}:=\int_{0}^{\infty}\left(e^{\rho\left(\log x+B_{t}-\frac{t}{2}\right)}-k\right)^{+} d t \tag{4.1}
\end{equation*}
$$

$(\rho, x, k>0)$. In the preceding Sections, we have studied the case $\rho=1$.
The analogue of Theorem 1.3 may be stated as :
Theorem 4.1 The Laplace transform of $\Pi_{1}^{(\rho)}$ under $P_{x}$ is given by :

$$
E_{x}\left(e^{-\frac{\theta}{2} \Pi_{1}^{(\rho)}}\right) \begin{cases}\frac{\sqrt{x} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho} x^{\frac{\rho}{2}}\right)}{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho}\right)-\sqrt{\theta} K_{\gamma}^{\prime}\left(\frac{\sqrt{4 \theta}}{\rho}\right)} & \text { if } x \geq 1  \tag{4.2}\\ 1+x \frac{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho}\right)+\sqrt{\theta} K_{\gamma}^{\prime}\left(\frac{\sqrt{4 \theta}}{\rho}\right)}{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho}\right)-\sqrt{\theta} K_{\gamma}^{\prime}\left(\frac{\sqrt{4 \theta}}{\rho}\right)} & \text { if } 0<x \leq 1\end{cases}
$$

with $\gamma=\frac{\sqrt{1-4 \theta}}{\rho}$ if $\theta \leq \frac{1}{4}$ and $\gamma=i \frac{\sqrt{4 \theta-1}}{\rho}$ if $\theta>\frac{1}{4}$.
The proof of this Theorem 4.1 is quite similar to that of Theorem 1.3. It hinges upon the following Lemma 4.2.

## Lemma 4.2

i) Let, for $\theta \geq 0$ and $\rho>0$, the function $\varphi_{\theta}^{(\rho)}: \mathbb{R}_{+} \rightarrow \mathbb{R}$ defined by :

$$
\begin{equation*}
\varphi_{\theta}^{(\rho)}(y)=\sqrt{y} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho} y^{\frac{\rho}{2}}\right) \quad(y \geq 0) \tag{4.3}
\end{equation*}
$$

Then, it satisfies :

$$
\begin{equation*}
\varphi_{\theta}^{(\rho)^{\prime \prime}}(y)+\theta\left(-y^{\rho-2}+\frac{1}{y^{2}}\right) \varphi_{\theta}^{(\rho)}(y)=0 \tag{4.4}
\end{equation*}
$$

(see [Pet], p. 306, with $\rho-2=m$, formula (8), after taking care of replacing a by $-a$ ). ii) Let $\widetilde{\varphi}_{\theta}^{(\rho)}: \mathbb{R}_{+} \rightarrow \mathbb{R}$ defined by :

$$
\widetilde{\varphi}_{\theta}^{(\rho)}(y)= \begin{cases}\varphi_{\theta}^{(\rho)}(y) & \text { if } y \geq 1  \tag{4.5}\\ \varphi_{\theta}^{(\rho)}(1)-\varphi_{\theta}^{(\rho)^{\prime}}(1)+y \varphi_{\theta}^{(\rho)^{\prime}}(1) & \text { if } y \leq 1\end{cases}
$$

Then $\widetilde{\varphi}_{\theta}^{(\rho)}$ is positive, decreasing, convex and satisfies :

$$
\begin{equation*}
\widetilde{\varphi}_{\theta}^{(\rho)}(y)+\theta\left(-y^{\rho-2}+\frac{1}{y^{2}}\right) 1_{y \geq 1} \widetilde{\varphi}_{\theta}^{(\rho)}(y)=0 \tag{4.6}
\end{equation*}
$$

Remark 4.1 As a check, we note that formula (4.2) allows to recover the identity (1.8) :

$$
\int_{0}^{\infty} e^{\rho B_{t}-\frac{\rho t}{2}} d t \stackrel{(\mathrm{law})}{=} \frac{2}{\rho^{2} \gamma_{\frac{1}{\rho}}}
$$

Indeed, on one hand :

$$
E\left(e^{-\frac{\theta}{2} \frac{2}{\rho^{2} \gamma_{1}}}\right)=\frac{1}{\Gamma\left(\frac{1}{\rho}\right)} \int_{0}^{\infty} e^{-\frac{\theta}{\rho^{2} z}-z} z^{\frac{1}{\rho}-1} d z=\frac{2}{\Gamma\left(\frac{1}{\rho}\right)} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4 \theta}}{\rho}\right)\left(\frac{\theta}{\rho^{2}}\right)^{\frac{1}{2 \rho}} \text { (see [Leb], p. }
$$

On the other hand :

$$
\begin{aligned}
E_{1}\left(e^{-\frac{\theta}{2} \Pi_{0}^{(\rho)}}\right) & =\lim _{\epsilon \downharpoonright 0} E_{1}\left(\exp -\frac{\theta}{2} \int_{0}^{\infty}\left(e^{\rho B_{t}-\frac{\rho t}{2}}-\varepsilon\right)^{+} d t\right) \\
& =\lim _{\epsilon \downarrow 0} E_{1}\left(\exp -\frac{\theta \varepsilon}{2} \int_{0}^{\infty}\left(e^{\rho\left(\frac{\log \frac{1}{\varepsilon}}{\rho}+B_{t}\right)-\frac{\rho t}{2}}-1\right)^{+} d t\right) \\
& =\lim _{\epsilon \downarrow 0} \frac{\left(\frac{1}{\varepsilon}\right)^{\frac{1}{2 \rho}} K_{\sqrt{1-4 \theta \varepsilon}}^{\rho}}{\left.\frac{\sqrt{4 \theta \varepsilon}}{\rho} K_{\frac{\sqrt{1-4 \theta \varepsilon}}{\rho}}\left(\frac{\sqrt{4 \theta \varepsilon}}{\rho}\right)-\sqrt{\theta \varepsilon}\right)^{\frac{1}{\rho} \frac{\rho}{2}}} K_{\frac{\sqrt{1-4 \theta \varepsilon}}{\prime}}^{\prime}\left(\frac{\sqrt{4 \theta \varepsilon}}{\rho}\right)
\end{aligned}
$$

from (4.2), replacing $\theta$ by $\theta \varepsilon$ and $x$ by $\left(\frac{1}{\varepsilon}\right)^{\frac{1}{\rho}}$

$$
\begin{aligned}
& \underset{\varepsilon \rightarrow 0}{\sim} \frac{\left(\frac{1}{\varepsilon}\right)^{\frac{1}{2 \rho}} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4 \theta}}{\rho}\right)}{\frac{1}{2} K_{\frac{\sqrt{1-4 \theta \varepsilon}}{\rho}}\left(\frac{\sqrt{4 \theta \varepsilon}}{\rho}\right)-\frac{\rho}{2}\left(\frac{\sqrt{4 \theta \varepsilon}}{\rho}\right) K_{\frac{\sqrt{1-4 \theta \varepsilon}}{\rho}}^{\prime}\left(\frac{\sqrt{4 \theta \varepsilon}}{\rho}\right)} \\
& \underset{\varepsilon \rightarrow 0}{\sim} \frac{\left(\frac{1}{\varepsilon} \frac{\frac{1}{2 \rho}}{\frac{\rho}{2 \rho}} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4 \theta \theta}}{\rho}\right)\right.}{\rho} \quad\left(\text { since } z K^{\prime} \nu=\nu K_{\nu}(z)-z K_{\nu+1}(z)\right) \\
& \underset{\varepsilon \rightarrow 0}{\sim} \frac{\left(\frac{\sqrt{1 \theta \varepsilon}}{\rho}\right)^{\frac{1}{2 \rho}} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4 \theta \varepsilon}}{\rho}\right)}{\frac{\rho}{2}\left(\Gamma\left(\frac{1}{\rho}+1\right)\left(\frac{\sqrt{4 \theta \varepsilon}}{\rho}\right)^{-\frac{1}{\rho}}\right.} \quad\left(\text { since } K_{\nu}(z) \sim \frac{1}{2} \Gamma(\nu)\left(\frac{z}{2}\right)^{-1}\right) \\
& =\frac{\left(\frac{1}{\varepsilon}\right)^{\frac{1}{2 \rho}} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4 \theta}}{\rho}\right)}{\Gamma\left(\frac{1}{\rho}\right)} \frac{\varepsilon^{\frac{1}{2 \rho}} \theta^{\frac{1}{2 \rho}}}{\rho^{\frac{1}{\rho}}} \\
& \underset{\varepsilon \downarrow 0}{\longrightarrow}\left(\frac{\theta}{\rho^{2}}\right)^{\frac{1}{2 \rho}} \frac{2}{\Gamma\left(\frac{1}{\rho}\right)} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4 \theta}}{\rho}\right)=E\left(e^{-\frac{\theta}{2} \Pi_{0}^{(\rho)}}\right)
\end{aligned}
$$

## Remark 4.2

Taking up again the arguments of the proof of Theorem 3.1, it is not difficult to see that :

$$
E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}^{(\rho)}}\right) \underset{\theta \rightarrow \infty}{\sim} \frac{C(\rho)}{\theta^{\frac{1}{3}}}
$$

where $C(\rho)$ is a strictly positive constant, depending on $\rho$. We then deduce that, for $\alpha<0$ :

$$
E_{1}\left(\left(\Pi_{1}^{(\rho)}\right)^{\alpha}\right)<\infty \text { if and only if }|\alpha|<\frac{1}{3}
$$

On the other hand, it is not difficult to see that :

- If $x>1$, for all $\alpha<0, E_{1}\left(\left(\Pi_{1}^{(\rho)}\right)^{\alpha}\right)<\infty$.
- If $x<1$, for all $\alpha<0, E_{1}\left(\left(\Pi_{1}^{(\rho)}\right)^{\alpha}\right)=+\infty$.

Concerning the positive moments of $\Pi_{1}^{(\rho)}$ :

- If $0<\alpha<\frac{1}{\rho}$, then $E_{x}\left(\left(\Pi_{1}^{(\rho)}\right)^{\alpha}\right)<\infty$

Indeed (4.7), for $x=1$, follows from :

$$
\begin{aligned}
E_{1}\left(\left(\Pi_{1}^{(\rho)}\right)^{\alpha}\right) & \leq E_{1}\left(\left(\Pi_{0}^{(\rho)}\right)^{\alpha}\right)=\frac{1}{\Gamma\left(\frac{1}{\rho}\right)} \int_{0}^{\infty}\left(\frac{2}{\rho^{2} z}\right)^{\alpha} e^{-z} z^{\frac{1}{\rho}-1} d z \\
< & \infty \text { if } \alpha<\frac{1}{\rho} \\
(\text { since } \quad & \left.\Pi_{0} \stackrel{(\text { law })}{=} \frac{2}{\rho^{2} \gamma_{\frac{1}{\rho}}}, \quad \text { from (1.10) }\right)
\end{aligned}
$$

The fact that $E_{x}\left(\left(\Pi_{1}^{(\rho)}\right)^{\alpha}\right)<\infty$ for every $x>0$, and every $\alpha<\frac{1}{\rho}$ may be obtained by using arguments close to those used in the proof of Theorem 1.1.
We believe that, for all $\rho>0$ :

$$
\begin{equation*}
E_{x}\left(\left(\Pi_{1}^{(\rho)}\right)^{\frac{1}{\rho}}\right)=+\infty \tag{4.8}
\end{equation*}
$$

We now show (4.8), when $\rho>1$
It then suffices, by using the arguments of the proof of Theorem 1.1, to see that (4.8) is true when $x=1$.

- We first show that : $1-E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}^{(\rho)}}\right) \underset{\theta \rightarrow 0}{\sim} C \theta^{\frac{1}{\rho}}$

Indeed, from Theorem 4.1, we have :

$$
E_{1}\left(\exp -\frac{\theta}{2} \Pi_{1}^{(\rho)}\right)=1+\frac{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho}\right)+\sqrt{\theta} K_{\gamma}^{\prime}\left(\frac{\sqrt{4 \theta}}{\rho}\right)}{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho}\right)-\sqrt{\theta} K_{\gamma}^{\prime}\left(\frac{\sqrt{4 \theta}}{\rho}\right)}
$$

with $\gamma=\frac{\sqrt{1-4 \theta}}{\rho}\left(\right.$ and $\left.\theta \leq \frac{1}{4}\right)$. Thus :

$$
\begin{aligned}
1- & E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}^{(\rho)}}\right)=-\frac{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho}\right)+\frac{\rho}{2}\left\{\frac{\sqrt{4 \theta}}{\rho} K_{\gamma}^{\prime}\left(\frac{\sqrt{4 \theta}}{\rho}\right)\right\}}{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4 \theta}}{\rho}\right)-\frac{\rho}{2}\left\{\frac{\sqrt{4 \theta}}{\rho} K_{\gamma}^{\prime}\left(\frac{\sqrt{4 \theta}}{\rho}\right)\right\}} \\
& =-\frac{\frac{1}{2} K_{\frac{\sqrt{1-4 \theta}}{\rho}}\left(\frac{\sqrt{4 \theta}}{\rho}\right)\left(1-\frac{\sqrt{1-4 \theta}}{\rho}\right)-\frac{\sqrt{\theta}}{\rho} K_{\frac{\sqrt{1-4 \theta}}{\rho}-1}\left(\frac{\sqrt{4 \theta}}{\rho}\right)}{\frac{1}{2} K_{\frac{\sqrt{1-4 \theta}}{\rho}}\left(\frac{\sqrt{4 \theta}}{\rho}\right)\left(1+\frac{\sqrt{1-4 \theta}}{\rho}\right)+\frac{\sqrt{\theta}}{\rho} K_{\frac{\sqrt{1-4 \theta}}{\rho}}^{\rho}\left(\frac{\sqrt{4 \theta}}{\rho}\right)}
\end{aligned}
$$

(after using $\left.z K_{\mu}^{\prime}(z)=-\mu K_{\mu}(z)-z K_{\mu-1}(z)\right)$.
Since $\rho>1$, we replace $K_{\frac{\sqrt{1-4 \theta}}{\rho}-1}$ by $K_{1-\frac{\sqrt{1-4 \theta}}{\rho}}\left(\right.$ since $\left.K_{\mu}=K_{-\mu}\right)$ and we deduce, from :

$$
K_{1-\frac{\sqrt{1-4 \theta}}{\rho}} \underset{\theta \rightarrow 0}{\longrightarrow} K_{1-\frac{1}{\rho}}, \quad \text { with } \quad 1-\frac{1}{\rho}>0:
$$

that:
$1-E_{1}\left(e^{-\frac{\theta}{2} \Pi_{1}^{(\rho)}}\right) \underset{\theta \rightarrow 0}{\sim} \frac{a \theta^{1-\frac{1}{2 \rho}}+b \theta^{\frac{1}{2 \rho}}}{a^{\prime} \theta^{-\frac{1}{2 \rho}}+b^{\prime} \theta^{\frac{1}{2 \rho}}}\left(\right.$ since , for $\mu>0 K_{\mu}(z) \underset{z \rightarrow 0}{\sim} C_{\mu} z^{-\mu} ;$ see [Leb], p. 136)

$$
\underset{\theta \rightarrow 0}{\sim} \frac{b}{a^{\prime}} \theta^{\frac{1}{\rho}} \quad\left(\text { since } \rho \geq 1 \text { implies } 1-\frac{1}{2 \rho} \geq \frac{1}{2 \rho}\right)
$$

- From (4.9) we deduce :

$$
\begin{aligned}
& \int_{0}^{\infty} e^{-\frac{\theta t}{2}} P\left(\Pi_{1}^{(\rho)} \geq t\right) d t=\frac{2}{\theta}\left(1-E\left(e^{-\frac{\theta}{2} \Pi_{1}^{(\rho)}}\right)\right) \\
& \underset{\theta \rightarrow \infty}{\sim} C^{\prime} \theta^{\frac{1}{\rho}-1}
\end{aligned}
$$

Hence, from the Tauberian Theorem :

$$
\begin{equation*}
P\left(\Pi_{1}^{(\rho)} \geq t\right) \underset{t \rightarrow \infty}{\sim} C^{\prime \prime} \frac{1}{t^{\frac{1}{\rho}}} \tag{4.10}
\end{equation*}
$$

and

$$
\begin{aligned}
E\left(\left(\Pi_{1}^{(\rho)}\right)^{\frac{1}{\rho}}\right) & =\int_{0}^{\infty} P\left(\left(\Pi_{1}^{(\rho)}\right)^{\frac{1}{\rho}} \geq t\right) d t \\
& =\int_{0}^{\infty} P\left(\Pi_{1}^{(\rho)} \geq t^{\rho}\right) d t=+\infty, \quad \text { from (4.10). }
\end{aligned}
$$

We note that, for $\rho<1$, the preceding argument cannot be applied since :

$$
\int_{1}^{\infty} e^{-\frac{\theta t}{2}} \frac{d t}{t^{\frac{1}{\rho}}} \underset{\theta \rightarrow 0}{\longrightarrow} \int_{1}^{\infty} \frac{d t}{t^{\frac{1}{\rho}}}<\infty .
$$

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