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An interesting family of Black-Scholes perpetuities

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Abstract We obtain the Laplace transform and integrability properties of the integral over \mathbb{R}_+ of the call quantity associated with geometric Brownian motion with negative drift, thus adding a new element to the list of already studied Brownian perpetuities.

Key words Bessel-McDonald functions, Sturm-Liouville equation, perpetuities.

AMS Classification 60 J 25, 60 J 65, 60 G 44.

1 Motivation, introduction

1.1 Let $(B_t, t \ge 0)$ denote 1-dimensional Brownian motion starting from 0; associated to B, one considers the geometric Brownian motion :

$$\left(\mathcal{E}_t := \exp\left(B_t - \frac{t}{2}\right), \ t \ge 0\right)$$

which is a positive martingale converging a.s. to 0 as $t \to \infty$. In recent years, the following questions have been asked to the second author, in connection with European option pricing :

i) to express as simply as possible the quantity :

$$\int_0^\infty \lambda \, e^{-\lambda t} \, E\big((\mathcal{E}_t - k)^+ \big) dt \qquad (k \ge 0)$$

ii) to find the law of $\int_0^s (\mathcal{E}_t - k)^+ dt$, for fixed s

iii) to find the law of $\int_0^\infty \lambda e^{-\lambda t} (\mathcal{E}_t - k)^+ dt$

i) may be solved quite explicitly, and indeed this has led the authors, jointly with D. Madan, to write a series of papers relating European option prices and laws of last passage times of continuous positive martingales converging to 0 as $t \to \infty$ (see, e.g. [MRY1], [MRY2], [MRY3]).

Questions ii) and iii) are harder to solve explicitly, as, indeed, one may look for double Laplace transforms of either quantities, e.g. :

$$\int_0^\infty e^{-\mu t} E(e^{-\lambda \int_0^t (\mathcal{E}_s - k)^+ ds}) dt$$

In the present paper, we are studying thoroughly the law of $\int_0^\infty (\mathcal{E}_t - k)^+ dt$ which seems to be slightly less difficult than either question *ii*) or *iii*). However, the results we obtain are not particularly simple.

1.2 These motivations having been presented, we now concentrate exclusively on the law of $\int_0^\infty (\mathcal{E}_t - k)^+ dt$. It will be convenient to consider the two-parameter process :

$$\left(\mathcal{E}_t^{(x)} := x \exp\left(B_t - \frac{t}{2}\right) = \exp\left(\log x + B_t - \frac{t}{2}\right), \ t \ge 0\right)$$
(1.1)

 $(\mathcal{E}_t^{(x)}, t \ge 0, x > 0)$ is a Markov process taking values in \mathbb{R}_+ which, most often, we shall denote as : $((\mathcal{E}_t, t \ge 0; P_x, x > 0)$ since $\mathcal{E}_0^{(x)} = x$.

To simplify notation, we shall write P for P_1 and \mathcal{E}_t for $\mathcal{E}_t^{(1)}$. For any k > 0, let us define :

$$\Pi_k^{(x)} := \int_0^\infty (\mathcal{E}_t^{(x)} - k)^+ dt \tag{1.2}$$

Again, to simplify, we shall write Π_k for $\Pi_k^{(1)}$. Since $\mathcal{E}_t^{(x)} \xrightarrow[t \to \infty]{} 0$ a.s., the integral which defines (1.2) is a.s. convergent, as $(\mathcal{E}_t^{(x)} - k)^+ = 0$ for $t \ge \mathcal{G}_k := \sup\{t \ ; \ \mathcal{E}_t^{(x)} = k\}$, and $\mathcal{G}_k < \infty$ a.s.

1.3 We now explain about some reductions of the study of the laws of $\Pi_k^{(x)}$, to closely related problems.

i) From Itô's formula, we deduce :

$$\mathcal{E}_{t}^{(x)} = x + \int_{0}^{t} \mathcal{E}_{s}^{(x)} dB_{s} = \beta_{A_{t}^{(x)}}$$
(1.3)

where $(\beta_u, u \ge 0)$ denotes the Dubins-Schwarz Brownian motion associated with $(\mathcal{E}_t^{(x)}, t \ge 0)$ (and $\beta_0 = x$) and :

$$A_t^{(x)} = \langle \mathcal{E}^{(x)} \rangle_t = \int_0^t (\mathcal{E}_s^{(x)})^2 ds$$
(1.4)

Hence :

$$\Pi_k^{(x)} = \int_0^\infty (\beta_{A_s^{(x)}} - k)^+ ds = \int_0^{T_0(\beta)} \frac{(\beta_v - k)^+}{\beta_v^2} dv$$
(1.5)

(after the change of variable $A_s^{(x)} = v$ and where $T_0(\beta) = \inf\{u \ge 0 ; \beta_u = 0\}$)

$$= \int_{k}^{\infty} \frac{(y-k)}{y^2} L_{T_0(\beta)}^y dy \qquad \left(= \int_{k}^{\infty} dz \int_{z}^{\infty} \frac{dy}{y^2} L_{T_0(\beta)}^y \right)$$
(1.6)

from the density of occupation formula, and where $L_{T_0}^y$ denotes the local time at time T_0 and level y of Brownian motion ($\beta_u, u \ge 0$).

ii) When x = k, the first Ray-Knight Theorem allows us to write, from (1.6) :

$$\Pi_k^{(k)} = \int_k^\infty \frac{(y-k)}{y^2} \,\lambda_{y-k} \,dy \tag{1.7}$$

where $(\lambda_z, z \ge 0)$, conditionally on $\lambda_0 = l$, is a 0-dimensional squared Bessel process starting at l, and where λ_0 is an exponential variable with parameter $\frac{1}{2k}$, i.e. with expectation 2k.

1.4 From the elementary relations :

$$(x e^{B_t - \frac{t}{2}} - k)^+ = x \left(e^{B_t - \frac{t}{2}} - \frac{k}{x} \right)^+ = k \left(\frac{x}{k} \left(e^{B_t - \frac{t}{2}} \right) - 1 \right)^+$$
(1.8)

valid for every x, k > 0, we deduce that the law of Π_k under P_x is that of $x \prod_{\frac{k}{x}}$ under P, or that of $k \Pi_1$ under $P_{\frac{x}{k}}$. In other words, for every Borel positive function φ , we have :

$$E_x\left[\varphi(\Pi_k)\right] = E\left[\varphi(x\,\Pi_{\frac{k}{x}})\right] = E_{\frac{x}{k}}\left(\varphi(k\Pi_1)\right) \tag{1.9}$$

These relations allow us to reduce our study of the law of Π_k under P_x to that of Π_1 under $P_{\frac{x}{k}}$. We might as well limit ourselves to the study of Π_k under P.

1.5 It is proven in Dufresne [Duf] and Yor [Y] (see also P. Salminen and M. Yor [SY1]) that, for every $a \neq 0$ and $\nu > 0$:

$$\int_0^\infty \exp\left(aB_t - \nu t\right) dt \stackrel{\text{(law)}}{=} \frac{2}{a^2 \gamma_{\frac{2\nu}{a^2}}} \tag{1.10}$$

where γ_b denotes a gamma variable with parameter b, i.e. :

$$P(\gamma_b \in dt) = \frac{1}{\Gamma(b)} e^{-t} t^{b-1} dt \qquad (t \ge 0)$$
(1.11)

In particular, for a = 1 and $\nu = \frac{1}{2}$

$$\Pi_0 \stackrel{\text{(law)}}{=} \frac{2}{\gamma_1} = \frac{2}{\mathfrak{e}} \tag{1.12}$$

where \mathfrak{e} is a standard exponential variable.

1.6 Let $\nu \neq 0$ and :

$$\mathcal{E}_t^{(x,\nu)} := \exp\left(\nu\left(\log x + B_t\right) - \frac{\nu^2 t}{2}\right) = x^{\nu} \exp\left(\nu B_t - \frac{\nu^2 t}{2}\right)$$

and define, for k > 0 :

$$\Pi_k^{(x,\nu)} := \int_0^\infty (\mathcal{E}_s^{(x,\nu)} - k)^+ ds \tag{1.13}$$

Since, by scaling, $(\mathcal{E}_t^{(x,\nu)}, t \ge 0) \stackrel{(\text{law})}{=} (x^{\nu} \mathcal{E}_{\nu^2 t}, t \ge 0)$, we have :

$$\Pi_k^{(x,\nu)} \stackrel{\text{(law)}}{=} \int_0^\infty (x^\nu \mathcal{E}_{\nu^2 s} - k)^+ ds = \frac{x^\nu}{\nu^2} \int_0^\infty \left(\mathcal{E}_v - \frac{k}{x^\nu}\right)^+ dv$$

$$\stackrel{\text{(law)}}{=} \frac{x^\nu}{\nu^2} \Pi_{\frac{k}{x^\nu}} \tag{1.14}$$

Thus, the study of the law of $\Pi_k^{(x,\nu)}$ may be reduced very simply to that of $\Pi_{\frac{k}{x^{\nu}}}$. This is the reason why we have chosen, in this paper, to limit ourselves to $\nu = 1$.

1.7 Here are our results :

Theorem 1.1 Let $\alpha \ge 0$. Then, for every x > 0, $E_x((\Pi_1)^{\alpha}) < \infty$ if and only if $\alpha < 1$ **Theorem 1.2** Let $\alpha < 0$ and x > 0. Then :

- i) For every x > 1, $E_x((\Pi_1)^{\alpha}) < \infty$.
- ii) For every x < 1, $P_x(\Pi_1 = 0) = 1 x$; hence $E_x((\Pi_1)^{\alpha}) = +\infty$.

iii) For
$$x = 1$$
, $E_1((\Pi_1)^{\alpha}) < \infty$ if and only if $|\alpha| < \frac{1}{3}$.

Theorem 1.3 (Laplace transform of Π_1). For every $\theta \ge 0$:

$$E_x(e^{-\frac{\theta}{2}\Pi_1}) = \begin{cases} \frac{\sqrt{x} K_{\gamma}(\sqrt{4\theta}x)}{\frac{1}{2} K_{\gamma}(\sqrt{4\theta}) - \sqrt{\theta} K_{\gamma}'(\sqrt{4\theta})} & \text{if } x \ge 1\\ 1 + x \frac{\frac{1}{2} K_{\gamma}(\sqrt{4\theta}) + \sqrt{\theta} K_{\gamma}'(\sqrt{4\theta})}{\frac{1}{2} K_{\gamma}(\sqrt{4\theta}) - \sqrt{\theta} K_{\gamma}'(\sqrt{4\theta})} & \text{if } 0 < x \le 1 \end{cases}$$
(1.15)

where K_{γ} denotes the Bessel-McDonald function with index γ (see [Leb], p. 108) and where $\gamma = \sqrt{1-4\theta}$ if $4\theta \leq 1$ and $\gamma = i\sqrt{4\theta-1}$ if $4\theta \geq 1$. We shall also prove, as a consequence of Theorem 1.3 :

Theorem 1.4 Let $(\lambda_x, x \ge 0)$ denote a squared Bessel process with dimension 0 started at l and denote its law by $Q_l^{(0)}$. Then : • If $4\theta \ge 1$:

$$Q_l^{(0)} \left(\exp -\frac{\theta}{2} \int_1^\infty \frac{(x-1)}{x^2} \lambda_{x-1} dx \right) \left(= Q_l^{(0)} \left(\exp -\frac{\theta}{2} \int_0^\infty \frac{x}{(x+1)^2} \lambda_x dx \right) \right.$$
$$\left. = \exp \frac{l}{2} \left(\frac{1}{2} + \frac{\sqrt{\theta} K'_{i\nu}(\sqrt{4\theta})}{K_{i\nu}(\sqrt{4\theta})} \right)$$
(1.16)

with $\nu = \sqrt{4\theta - 1}$

• If $4\theta \leq 1$:

$$Q_l^{(0)}\left(\exp-\frac{\theta}{2}\int_1^\infty \frac{x-1}{x^2}\,\lambda_{x-1}\,dx\right) = \exp\frac{l}{2}\left(\frac{1}{2} + \frac{\sqrt{\theta}\,K_\nu'(\sqrt{4\theta})}{K_\nu(\sqrt{4\theta})}\right)$$

with $\nu = \sqrt{1 - 4\theta}$.

In Section 3 of this work, we shall study the asymptotic behavior, as $\theta \to \infty$, of $E(e^{-\frac{\theta}{2}\Pi_1})$ and we shall obtain :

$$E(e^{-\frac{\theta}{2}\Pi_1}) \mathop{\sim}\limits_{\theta \to \infty} \frac{C}{\theta^{\frac{1}{3}}} \qquad (C > 0)$$
(1.17)

Finally, in a short Section 4, we shall indicate how Theorems 1.1, 1.2 and 1.3 extend when we replace Π_k by $\Pi_k^{(\rho)}$, with :

$$\Pi_{k}^{(\rho)} := \int_{0}^{\infty} \left(e^{\rho(B_{t} - \frac{t}{2})} - k \right)^{+} dt \qquad (\rho, k > 0)$$

$$\stackrel{(\text{law})}{=} \frac{1}{\rho^{2}} \int_{0}^{\infty} \left(e^{(B_{u} - \frac{u}{2\rho})} - k \right)^{+} du \qquad (\text{by scaling})$$
(1.18)

which, from (1.18), corresponds to consider an extension of our previous perpetuities relative to $(B_t - \frac{t}{2}, t \ge 0)$ to Brownian motion with drift $-\frac{1}{2\rho}$, i.e. $(B_t - \frac{t}{2\rho}, t \ge 0)$.

1.8 For studies of other perpetuities related to Brownian motion with drift, we refer the reader to Salminen-Yor ([SY1], [SY2], [SY3]).

2 Proofs of Theorems 1.1, 1.2 and 1.3

2.1 A first proof of Theorem 1.1

2.1 *i*) We now prove that, for every x > 0, $E_x(\Pi_1) = \infty$ Indeed, from (1.6) :

$$E_{1}(\Pi_{1}) = \int_{1}^{\infty} \frac{(y-1)}{y^{2}} E_{1}(L_{T_{0}}^{y}) dy$$

$$= \int_{1}^{\infty} \frac{(y-1)}{y^{2}} E_{1}(L_{T_{0}}^{1}) dy$$

$$= 2 \int_{1}^{\infty} \frac{(y-1)}{y^{2}} dy = +\infty$$
(2.1)

(since, for $y \ge 1$, $E_1(L_{T_0}^y) = E_1(L_{T_0}^1) = 2$, as $L_{T_0}^1$ is an exponential variable with parameter $\frac{1}{2}$).

2.1 *ii*) We now prove that, for every x > 0, $E_x(\Pi_1) = \infty$

• For $x \ge 1$, this is clear since $\Pi_k^{(x)}$ is an increasing function of x (and a decreasing function of k). Hence :

$$E_x(\Pi_1) \ge E_1(\Pi_1) = +\infty \qquad \text{(from (2.1))}$$

• For x < 1, from the Markov property :

$$\Pi_{1}^{(x)} \stackrel{(\text{law})}{=} 1_{\left\{T_{\log(\frac{1}{x})}^{(-\frac{1}{2})} < \infty\right\}} \cdot \Pi_{1}$$
(2.2)

with $T_{\log(\frac{1}{x})}^{(-\frac{1}{2})} := \inf\left\{t \ge 0 \ ; \ B_t - \frac{t}{2} = \log\left(\frac{1}{x}\right)\right\}$, and in (2.2), Π_1 and $T_{\log(\frac{1}{x})}^{(-\frac{1}{2})}$ are assumed independent. Hence :

$$E(\Pi_1^{(x)}) = P(T_{\log(\frac{1}{x})}^{(-\frac{1}{2})} < \infty) \ E(\Pi_1) = +\infty$$

from (2.1) and since $P(T_{\log(\frac{1}{x})}^{(-\frac{1}{2})} < \infty) = x$ (see Lemma 2.1 below).

2.1 *iii)* We prove that, for $0 \le \alpha < 1$, $E_x((\Pi_1)^{\alpha}) < \infty$ We have :

$$E_x((\Pi_1)^{\alpha}) = x^{\alpha} E((\Pi_{\frac{1}{x}})^{\alpha}) \quad (\text{from } (1.9))$$

$$\leq x^{\alpha} E((\Pi_0)^{\alpha})$$

(since Π_k is a decreasing function of k)

$$\leq x^{\alpha} \int_{0}^{\infty} \frac{2}{y^{\alpha}} e^{-y} dy < \infty \qquad (\text{from}(1.12))$$

2.2 Second proof of Theorem 1.1

It hinges upon :

Lemma 2.1

Let $G_a^{(\nu)} := \sup\{u \ge 0; B_u + \nu u = a\}$ (= 0 if this set is empty) and $T_a^{(\nu)} := \inf\{u \ge 0; B_u + \nu u = a\}$ (= + ∞ if this set is empty).

i) If ν and a have the same sign :

•
$$E(e^{\mu G_a^{(\nu)}}) < \infty$$
 if and only if $\mu < \frac{\nu^2}{2}$ (2.3)

• for every real α , $E((G_a^{(\nu)})^{\alpha}) < \infty$ (2.4)

ii) If ν and a have opposite signs :

•
$$E(e^{\mu G_a^{(\nu)}}) < \infty$$
 if and only if $\mu < \frac{\nu^2}{2}$ (2.5)

• for every real $\alpha > 0$, $E((G_a^{(\nu)})^{\alpha}) < \infty$ (2.6)

•
$$P(T_a^{(\nu)} = \infty) = P(G_a^{(\nu)} = 0) = 1 - e^{2\nu a}$$
 (2.7)

Hence, for $\alpha < 0$, $E((G_a^{(\nu)})^{\alpha}) = +\infty$

The proof of this Lemma is obvious. It hinges on the well-known formulae : • If $\mu a > 0$:

• If
$$\nu, a > 0$$
:

$$P(G_a^{(\nu)} \in dt) = \frac{\nu}{\sqrt{2\pi t}} \exp\left(-\frac{1}{2t}(a-\nu t)^2\right) dt$$

• If
$$\nu, a > 0$$
 :

$$P(G_a^{(-\nu)} \in dt) = \frac{\nu}{\sqrt{2\pi t}} \exp\left(-\frac{1}{2t}(a+\nu t)^2\right) dt$$

• If ν and a have opposites signs :

$$P(G_a^{(\nu)} > 0) = P(T_a^{(\nu)} < \infty) = e^{2\nu a}$$

• $G_a^{(\nu)} \stackrel{(\text{law})}{=} G_{(-a)}^{(-\nu)} \stackrel{(\text{law})}{=} \frac{1}{T_{\nu}^{(a)}}$

We now give a second proof of Theorem 1.1.

• We first show that, for $0 \leq \alpha < 1$, $E_x((\Pi_1)^{\alpha}) < \infty$. From the relation :

$$\Pi_1^{(x)} = x \int_0^\infty \left(\mathcal{E}_t - \frac{1}{x} \right)^+ dt = x \int_0^{G_{-\log x}^{\left(-\frac{1}{2}\right)}} \left(\mathcal{E}_t - \frac{1}{x} \right)^+ dt$$
$$\leq x \left(\sup_{t \ge 0} \mathcal{E}_t \right) \cdot G_{-\log x}^{\left(-\frac{1}{2}\right)}$$

we deduce that, for $0 \le \alpha < 1$, $\alpha p < 1$ and $\frac{1}{p} + \frac{1}{q} = 1$ (p, q > 1):

$$E_{x}((\Pi_{1})^{\alpha}) \leq x^{\alpha} E\left[\left(\sup_{t\geq 0} \mathcal{E}_{t}\right)^{\alpha} \left(G_{-\log x}^{\left(-\frac{1}{2}\right)}\right)^{\alpha}\right]$$

$$\leq x^{\alpha} E\left(\left(\sup_{t\geq 0} \mathcal{E}_{t}\right)^{\alpha p}\right)^{\frac{1}{p}} E\left(\left(G_{(-\log x)}^{\left(-\frac{1}{2}\right)}\right)^{\alpha q}\right)^{\frac{1}{q}}$$

$$\leq C x^{\alpha} E\left(\left(\sup_{t\geq 0} \mathcal{E}_{t}\right)^{\alpha p}\right)^{\frac{1}{p}} \qquad \text{(from Lemma 2.1)}$$

As $(\mathcal{E}_t, t \ge 0)$ is a positive martingale, starting from 1, and converging a.s. to 0 as $t \to \infty$, we get (see [NY] or [RY], Chap. II, Ex. 3.12, p. 73) :

$$\sup_{s \ge 0} \mathcal{E}_s \stackrel{\text{(law)}}{=} \frac{1}{U} \tag{2.8}$$

with U uniform on [0, 1]. Thus :

$$E\left(\left(\sup_{t\geq 0} \mathcal{E}_t\right)^{\alpha p}\right) = \int_0^1 \frac{1}{u^{\alpha p}} \, du < \infty \tag{2.9}$$

since $\alpha p < 1$.

• We then show that $E_x(\Pi_1) = \infty$. First of all, it follows from [MRY1] that :

$$E_1((\mathcal{E}_t - k)^+) = P(G_{\log k}^{(\frac{1}{2})} \le t)$$
(2.10)

Hence :

$$E_1(\Pi_1) = \int_0^\infty P(G_0^{(\frac{1}{2})} \le t) dt = E\left(\int_{G_0^{(\frac{1}{2})}}^\infty dt\right) = +\infty$$

since $G_0^{(\frac{1}{2})} < +\infty$ a.s. Likewise :

$$E_x(\Pi_1) = x E_1(\Pi_{\frac{1}{x}}) \quad (\text{from (1.9)})$$
$$= x E\left(\int_{G_{\log \frac{1}{x}}}^{\infty} dt\right) = \infty$$

from (2.10) and since $G_{\log \frac{1}{x}}^{(\frac{1}{2})} < \infty$ a.s.

2.3 Proof of Theorem 1.2

i) We already prove point i) :

Let x > 1 and x' such that 1 < x' < x. It is then obvious that :

$$\Pi_1^{(x)} \ge (x'-1)T_{\log(\frac{x'}{x})}^{(-\frac{1}{2})} \tag{2.11}$$

(with $T_a^{(-\frac{1}{2})} := \inf\{t \ge 0 ; B_t - \frac{t}{2} = a\}$, since, if $t < T_{\log \frac{x'}{x}}^{(-\frac{1}{2})}$, then $e^{\log x + B_t - \frac{t}{2}} - 1 \ge x' - 1$). Hence, with $\gamma > 0$:

$$E_{x}\left(\frac{1}{(\Pi_{1})^{\gamma}}\right) = \int_{0}^{\infty} P_{x}\left(\frac{1}{(\Pi_{1})^{\gamma}} \ge t\right) dt = \int_{0}^{\infty} P_{x}(\Pi_{1} < v) \frac{\gamma}{v^{1+\gamma}} dv$$

$$\leq \int_{0}^{\infty} \frac{\gamma}{v^{1+\gamma}} dv P\left(T_{\log\frac{x}{x}} \le \frac{v}{x'-1}\right) dv \quad (\text{from (2.11)})$$

$$= \int_{0}^{\infty} \frac{\gamma}{v^{1+\gamma}} dv \left(P\left(G_{\frac{1}{2}}^{(\log(\frac{x}{x'}))} \ge \frac{x'-1}{v}\right) dv$$

$$(\text{since} \quad T_{a}^{(\nu)} \stackrel{(\text{law})}{=} \frac{1}{G_{\nu}^{(a)}} \quad \text{and} \quad G_{\nu}^{(a)} \stackrel{(\text{law})}{=} G_{-\nu}^{(-a)})$$

$$= \int_{0}^{\infty} \frac{\gamma}{(x'-1)^{\gamma}} u^{\gamma-1} du P\left(G_{\frac{1}{2}}^{(\log(\frac{x}{x'}))} \ge u\right) du$$

$$= \frac{1}{(x'-1)^{\gamma}} E\left(\left(G_{\frac{1}{2}}^{(\log\frac{x}{x'})}\right)^{\gamma}\right) < \infty$$

$$(2.12)$$

(from point i) of Lemma 2.1).

ii) We now prove point *ii)* : It is clear that, for x < 1 :

$$\{\Pi_1^{(x)} = 0\} = \{T_{\log \frac{1}{x}}^{(-\frac{1}{2})} = \infty\}$$

Thus, $P\{\Pi_1^{(x)} = 0\} = P\{T_{\log \frac{1}{x}}^{(-\frac{1}{2})} = \infty\} = P\{G_{\frac{1}{2}}^{(\log x)} = 0\} = 1 - x > 0$ from Lemma 2.1 and (1.26).

iii) We now prove point iii) :

For this purpose, we write, for $\gamma > 0$:

$$E_1\left(\frac{1}{(\Pi_1)^{\gamma}}\right) = \frac{1}{\Gamma(\gamma)} \int_0^\infty E(e^{-\theta \Pi_1}) \theta^{\gamma-1} d\theta$$
(2.13)

and we show, in the next Section 3, that :

$$E_1(e^{-\theta\Pi_1}) \underset{\theta \to \infty}{\sim} \frac{C}{\theta^{\frac{1}{3}}} \qquad (C > 0)$$

Thus, $E\left(\frac{1}{(\Pi_1)^{\gamma}}\right) < \infty$ if and only if $\int_1^{\infty} \theta^{\gamma - 1 - \frac{1}{3}} d\theta < \infty$, that is, if and only if $\gamma < \frac{1}{3}$.

2.4 Proof of Theorem 1.3

a) A useful Lemma

Lemma 2.2 Let $\theta \ge 0$ and K_{γ} the Bessel McDonald function with index γ , such that $\gamma = \sqrt{1-4\theta}$ if $\theta \le \frac{1}{4}$ and $\gamma = i\sqrt{4\theta-1}$ if $\theta \ge \frac{1}{4}$. **1)** Define the function $\varphi_{\theta} : \mathbb{R}_{+} \to \mathbb{R}$ by :

$$\varphi_{\theta}(y) = \sqrt{y} K_{\gamma}(\sqrt{4\theta y}) \qquad (y \ge 0)$$
(2.14)

Then:

i) φ_{θ} is a real valued function which satisfies :

$$\varphi_{\theta}''(y) + \left(-\frac{\theta}{y} + \frac{\theta}{y^2}\right)\varphi_{\theta}(y) = 0$$
(2.15)

ii) φ_{θ} , restricted to $[1, \infty]$ is positive, convex, bounded and decreasing.

2) We define the function $\widetilde{\varphi}_{\theta} : \mathbb{R}_+ \to \mathbb{R}_+$ by :

$$\widetilde{\varphi}_{\theta}(y) = \begin{cases} \varphi_{\theta}(y) & \text{if } y \ge 1\\ \left(\varphi_{\theta}(1) - \varphi_{\theta}'(1)\right) + y\varphi_{\theta}'(1) & \text{if } 0 \le y \le 1 \end{cases}$$

$$(2.16)$$

Then $\widetilde{\varphi}_{\theta}$ is a bounded, positive, convex, decreasing function which satisfies :

$$\widetilde{\varphi}_{\theta}^{\prime\prime}(y) + \left(-\frac{\theta}{y} + \frac{\theta}{y^2}\right) \mathbf{1}_{y \ge 1} \cdot \widetilde{\varphi}_{\theta}(y) = 0 \tag{2.17}$$

b) Proof of Lemma 2.2

i) Relation (2.15) (as well as relation (2.17)) follows from a direct computation, using the equation $K_{\gamma}''(z) + \frac{1}{z} K_{\gamma}'(z) - \left(1 + \frac{\gamma^2}{z^2}\right) K_{\gamma}(z) = 0$ (see [Leb], p. 110) and the fact that $\gamma^2 = 1 - 4\theta$ (see Petiau [Pet], p. 306, formula (8), which needs to be corrected by replacing a by -a, or Kamke [Kam], p. 440).

We distinguish two cases :

Case 1 : $4\theta \leq 1$, $\gamma = \sqrt{1 - 4\theta}$. The function K_{γ} , hence also φ_{θ} , is positive. Furthermore φ_{θ} is bounded on \mathbb{R}_+ since :

$$\varphi_{\theta}(y) \underset{y \to 0}{\sim} C y^{\frac{1-\gamma}{2}}, \ \varphi_{\theta}(y) \underset{y \to \infty}{\sim} C' y^{\frac{1}{4}} e^{-\sqrt{4\theta y}} \qquad \text{(see [Leb], p. 123 and 136)}$$

On the other hand, from (2.15), the function φ_{θ} is convex on the interval $[1, \infty]$. As it is convex, positive, and bounded, it is decreasing. Lemma 2.2 is thus proven in this case.

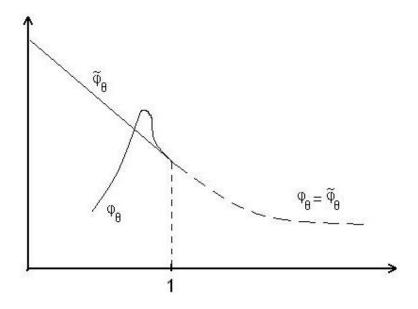


Figure 1: Graphs of φ_{θ} and $\widetilde{\varphi}_{\theta}$

Case 2: $4\theta \ge 1$, $\gamma = i\sqrt{4\theta - 1}$ and here $\varphi_{\theta}(y) = \sqrt{y} K_{i\nu}(\sqrt{4\theta y})$ with $\nu = \sqrt{4\theta - 1}$. From the integral representation formula (see [Leb], p. 119) :

$$K_{i\nu}(y) = \int_0^\infty e^{-y\cosh u} \cos(\nu u) du \qquad (y \ge 0)$$
(2.18)

we deduce that $K_{i\nu}(y)$ is real valued, hence so is $\varphi_{\theta}(y)$. On the other hand, φ_{θ} is bounded on $[0, \infty[$. Indeed, for $y \ge 0$, from (2.18) :

$$\left|\sqrt{4\theta}\,\varphi_{\theta}\left(\frac{y^2}{4\theta}\right)\right| = y \left|K_{i\nu}(y)\right| \le y \int_0^\infty e^{-y\frac{e^u}{2}} du \le y \int_{\frac{y}{2}}^\infty e^{-v} \frac{dv}{v}$$

Thus :

$$\left| \sqrt{4\theta} \,\varphi_{\theta} \left(\frac{y^2}{4\theta} \right) \right| \leq e^{-\frac{y}{4}} \quad \text{for } y \text{ large enough, hence} \left| \sqrt{4\theta} \,\varphi_{\theta} \left(\frac{y^2}{4\theta} \right) \right| \underset{y \to \infty}{\longrightarrow} 0 \quad \text{and} \\ \left| \sqrt{4\theta} \,\varphi_{\theta} \left(\frac{y^2}{4\theta} \right) \right| \leq y \left(C + |\log(y)| \right) \underset{y \to 0}{\longrightarrow} 0$$

On the other hand, it is clear, from (2.18), that $K_{i\nu}(y) > 0$ hence $\varphi_{\theta}(y) > 0$, for y large enough.

• We now show that $K_{i\nu}$ is decreasing on $[\nu, \infty[$. If not, there would exist a point $y_0 > \nu$ such that :

$$K_{i\nu}(y_0) > 0, \ K'_{i\nu}(y_0) = 0, \ K''_{i\nu}(y_0) \le 0.$$

However,

$$K_{i\nu}''(y_0) + \frac{1}{y_0} K_{i\nu}'(y_0) = \left(1 - \frac{\nu^2}{y_0^2}\right) K_{i\nu}(y_0)$$

Thus $K_{i\nu}''(y_0) > 0$, which is absurd. Since $K_{i\nu}$ is decreasing on $[\nu, \infty[$, and positive near ∞ , $K_{i\nu}$ is positive on $[\nu, \infty[$. Thus, φ_{θ} is positive on $[1, \infty[$ (since, if $y \ge 1$, $\sqrt{4\theta y} \ge \sqrt{4\theta} \ge \sqrt{4\theta - 1} = \nu$). From the relation (2.15), we then deduce that φ_{θ} is convex on $[1, \infty[$. Since it is bounded, convex and positive, it is decreasing. Lemma 2.2 is proven.

c) End of the proof of Theorem 1.3. Let $\left(M_t^{\theta} := \widetilde{\varphi}_{\theta}(B_t) \exp{-\frac{\theta}{2} \int_0^t \frac{\widetilde{\varphi}_{\theta}''}{\varphi_{\theta}}(B_s) ds}, t \ge 0\right)$. Then $(M_t^{\theta}, t \ge 0)$ is a local martingale. It is equal, from (2.17) to :

$$M_t^{\theta} = \widetilde{\varphi}_{\theta}(B_t) \exp{-\frac{\theta}{2} \int_0^t \frac{(B_s - 1)^+}{B_s^2} ds}$$
(2.19)

and, from Lemma 2.2, for every $x \ge 0$, $(M_{t \land T_0}^{\theta}, t \ge 0)$ is bounded. Thus, from Doob's optional stopping Theorem :

$$\widetilde{\varphi}_{\theta}(x) = E_x(M_0^{\theta}) = E_x \left[\widetilde{\varphi}_{\theta}(B_{T_0}) \exp{-\frac{\theta}{2} \int_0^{T_0(B)} \frac{(B_s - 1)^+}{B_s^2} ds} \right]$$
$$= \widetilde{\varphi}_{\theta}(0) E_x \left[\exp{-\frac{\theta}{2} \Pi_1} \right] \qquad (\text{from (1.5)})$$

Thus :

$$E_x\left(e^{-\frac{\theta}{2}\Pi_1}\right) = \frac{\widetilde{\varphi}_{\theta}(x)}{\widetilde{\varphi}_{\theta}(0)}$$

This is precisely Theorem 1.3, owing to formula (2.16) which yields $\tilde{\varphi}_{\theta}$ explicitly.

Remark 2.3 Theorem 1.3 allows to recover formula (1.9) : $\Pi_0 \stackrel{\text{(law)}}{=} \frac{2}{\mathfrak{e}} \cdot$ Indeed, on one hand :

$$E_1\left(e^{-\frac{\theta}{2}\Pi_0}\right) = E\left(e^{-\frac{\theta}{2}\frac{2}{\mathfrak{c}}}\right) = \int_0^\infty e^{-\frac{\theta}{z}-z} dz = 2\sqrt{\theta} K_1(\sqrt{4\theta}) \qquad \text{(see [Leb], p. 119)}$$

and, on the other hand :

$$E_{1}(e^{-\frac{\theta}{2}\Pi_{0}}) = \lim_{\epsilon \downarrow 0} E_{1}(e^{-\frac{\theta}{2}\Pi_{\epsilon}}) = \lim_{\epsilon \downarrow 0} E\left(\exp -\frac{\theta}{2} \int_{0}^{\infty} (e^{B_{t}-\frac{t}{2}} - \epsilon)^{+} dt\right)$$

$$= \lim_{\epsilon \downarrow 0} E\left(e^{-\frac{\theta\epsilon}{2} \int_{0}^{\infty} (\frac{1}{\epsilon}(e^{B_{t}-\frac{t}{2}} - 1)^{+} dt)\right) = \lim_{\epsilon \downarrow 0} E_{\frac{1}{\epsilon}}(e^{-\frac{\theta\epsilon}{2}\Pi_{1}})$$

$$= \lim_{\epsilon \downarrow 0} \frac{\frac{1}{\sqrt{\epsilon}} K_{\sqrt{1-4\theta\epsilon}}(\sqrt{4\theta})}{\frac{1}{2} K_{\sqrt{1-4\theta\epsilon}}(\sqrt{4\theta\epsilon}) - \frac{1}{2}\sqrt{4\theta\epsilon} K'_{\sqrt{1-4\theta\epsilon}}(\sqrt{4\theta\epsilon})}$$
(from (1.5), replacing θ by $\theta\epsilon$ and x by $\frac{1}{\epsilon}$)
$$= \lim_{\epsilon \downarrow 0} \frac{\frac{1}{\sqrt{\epsilon}} K_{1}(\sqrt{4\theta}) + O(1)}{\frac{1}{2} K_{\sqrt{1-4\theta\epsilon}}(\sqrt{4\theta\epsilon}) [1 - \sqrt{1-4\theta\epsilon}] + \sqrt{\theta\epsilon} K_{\sqrt{1-4\theta\epsilon+1}}(\sqrt{4\theta\epsilon})}$$
(since $zK'_{\nu}(z) = \nu K_{\nu}(z) - z K_{\nu+1}(z)$; see [Leb], p. 110)
$$= \lim_{\epsilon \downarrow 0} \frac{\frac{1}{\sqrt{\epsilon}} K_{1}(\sqrt{4\theta}) + O(1)}{O(1) + \sqrt{\theta\epsilon} (\frac{1}{2\theta\epsilon} + O(1))} = 2\sqrt{\theta} K_{1}(\sqrt{4\theta})$$

2.5 Proof of Theorem 1.4 From Theorem 1.3, we know that, for $\theta \geq \frac{1}{4}$:

$$E_{1}(e^{-\frac{\theta}{2}\Pi_{1}}) = \frac{2K_{i\nu}(\sqrt{4\theta})}{K_{i\nu}(\sqrt{4\theta}) - \sqrt{4\theta}K_{i\nu}'(\sqrt{4\theta})}$$
$$= \frac{1}{\frac{1}{\frac{1}{2} - \sqrt{\theta}\frac{K_{i\nu}'(\sqrt{4\theta})}{K_{i\nu}(\sqrt{4\theta})}}} = \frac{1}{2}\int_{0}^{\infty} e^{-\frac{l}{2} + \frac{l}{2}\left(\frac{1}{2} + \sqrt{\theta}\frac{K_{i\nu}'(\sqrt{4\theta})}{K_{i\nu}(\sqrt{4\theta})}\right)} dl$$
(2.20)

On the other hand, from (1.7):

$$E_{1}(e^{-\frac{\theta}{2}\Pi_{1}}) = E_{1}\left(\exp{-\frac{\theta}{2}}\int_{1}^{\infty}\frac{(y-1)^{+}}{y^{2}}\lambda_{y-1}\,dy\right)$$

$$= \frac{1}{2}\int_{0}^{\infty}e^{-\frac{l}{2}}dl\,E\left(\exp{-\frac{\theta}{2}}\int_{1}^{\infty}\frac{(y-1)^{+}}{y^{2}}\lambda_{y-1}dy\big|\lambda_{0}=l\right)$$

$$= \frac{1}{2}\int_{0}^{\infty}e^{-\frac{l}{2}}dl\,Q_{l}^{(0)}\left(\exp{-\frac{\theta}{2}}\int_{0}^{\infty}\frac{y}{(1+y)^{2}}\lambda_{y}\,dy\right)$$
(2.21)

where $Q_l^{(0)}$ denotes the expectation relative to a squared Bessel process of dimension 0, starting from l. The comparison of (2.21) and (2.20) implies Theorem 1.4 in the case

 $\theta \ge \frac{1}{4}$, since the Laplace transform is one-to-one. The proof of Theorem 1.4 in the case $\theta \le \frac{1}{4}$ is the same.

Remark 2.4 Let μ denote a positive σ -finite measure on \mathbb{R}_+ . Let Φ denote the unique decreasing, positive solution on \mathbb{R}_+ of the Sturm-Liouville equation $\Phi'' = \mu \Phi$ (and such that $\Phi(0) = 1$). It is well known (see [RY], Chap. IX, p. 444) that :

$$Q_l^{(0)}\left(\exp\left(-\frac{1}{2}\int_0^\infty \lambda_y \,\mu(dy)\right)\right) = \exp\left(\frac{l}{2}\,\Phi'(0_+)\right) \tag{2.22}$$

(Observe that, as $(\lambda_y, y \ge 0)$ has compact support a.s., $\int_0^\infty \lambda_y \mu(dy) < \infty$ if μ is σ -finite on \mathbb{R}_+).

Theorem 1.4 may be recovered easily by applying formula (2.22) with $\mu(dy) = \frac{y}{1+y^2} dy$. On the other hand, if $\mu(dy) = a(y)dy$, with

 $c_1 1_{[0,\gamma_1]}(y) \le a(y) \le c_2 1_{[0,\gamma_2]}(y) \qquad (0 < c_1 \le c_2, \ 0 < \gamma_1 \le \gamma_2),$

we deduce from [RY], Chap. XI, Corollary 1.8, that :

$$\exp -\frac{l}{2} c_2' \sqrt{\theta} \le E_l^{(0)} \left(e^{-\frac{\theta}{2} \int_0^\infty \lambda_y \, \mu(dy)} \right) \le \exp -\frac{l}{2} c_1' \sqrt{\theta}$$
(2.23)

for θ large enough, whereas, as we shall show in Section 3 :

$$Q_l^{(0)} \left(\exp -\frac{\theta}{2} \int_0^\infty \frac{y}{1+y^2} \lambda_y \, dy \right) \underset{\theta \to \infty}{\sim} C \, e^{-\frac{l}{2} \, \theta^{\frac{1}{3}}} \tag{2.24}$$

3 Asymptotic behavior of $E_1(e^{-\frac{\theta}{2}\Pi_1})$ as $\theta \to \infty$

We shall now end the proof of Theorem 1.2 by showing : **Theorem 3.1** There is the equivalence result:

$$E_1\left(e^{-\frac{\theta}{2}\Pi_1}\right) \underset{\theta \to \infty}{\sim} \frac{C}{\theta^{\frac{1}{3}}} \tag{3.1}$$

Proof of Theorem 3.1

We recall that, from Theorem 1.3 :

$$E_1\left(e^{-\frac{\theta}{2}\Pi_1}\right) = \frac{2K_{i\nu}(\sqrt{4\theta})}{K_{i\nu}(\sqrt{4\theta}) - \sqrt{4\theta}K'_{i\nu}(\sqrt{4\theta})}$$
(3.2)

with $\nu = \sqrt{4\theta - 1} \left(\theta \geq \frac{1}{4}\right)$. We shall successively find an equivalent of the numerator and the denominator of (3.2), the difficulty arising from the fact that, in $K_{i\nu}(\sqrt{4\theta})$ (and $K'_{i\nu}(\sqrt{4\theta})$ the argument $\sqrt{4\theta}$ and index $i\nu = i\sqrt{4\theta - 1}$ tend both to infinity as $\theta \to \infty$. To overcome this difficulty, we shall use some results about Bessel functions found in Watson ([Wat], p. 245-248), which we now recall.

i) Let $H_{i\nu}^{(1)}$ the first Hankel function (see [Leb], p. 120); it is related to $K_{i\nu}$ via the formula:

$$K_{i\nu}(z) = \frac{i\pi}{2} e^{-\frac{\nu\pi}{2}} H_{i\nu}^{(1)}(z e^{\frac{i\pi}{2}})$$
(3.3)

We define ϵ by the formula : $i\nu = iz(1-\epsilon)$ and assume that, as $z \to \infty$, ϵ remains bounded (with, of course, ν depending on z). Then, there is the second order asymptotic expansion :

$$H_{\nu}^{(1)}(z) = -\frac{2}{3\pi} \left\{ e^{\frac{2}{3}\pi i} \left(\sin \frac{\pi}{3} \right) \frac{\Gamma\left(\frac{1}{3}\right)}{\left(\frac{1}{6}z\right)^{\frac{1}{3}}} + e^{\frac{4}{3}\pi i}(\epsilon z) \left(\sin \frac{2\pi}{3} \right) \frac{\Gamma\left(\frac{2}{3}\right)}{\left(\frac{1}{6}z\right)^{\frac{2}{3}}} + o\left(\frac{1}{z^{\frac{2}{3}}}\right) \right\} (3.4)$$

ii) Let us study the numerator of (3.2):

$$N = 2K_{i\nu}(\sqrt{4\theta}) = 2\frac{i\pi}{2}e^{-\frac{\nu\pi}{2}}H_{i\nu}^{(1)}\left(e^{\frac{i\pi}{2}}\sqrt{4\theta}\right)$$

$$= (i\pi)e^{-\frac{\nu\pi}{2}}\left(-\frac{2}{3\pi}\right)\left\{C_{1}e^{\frac{2}{3}\pi i}\frac{1}{\left(e^{\frac{i\pi}{2}}\sqrt{4\theta}\right)^{\frac{1}{3}}} + o\left(\frac{1}{\theta^{\frac{1}{6}}}\right)\right\}$$

$$(3.5)$$

$$(\text{with } C_{1} = \left(\sin\frac{\pi}{3}\right)\Gamma\left(\frac{1}{3}\right)6^{\frac{1}{3}})$$

$$\underset{\theta \to \infty}{\sim} \frac{2^{\frac{2}{3}}}{3}C_{1}e^{-\frac{\nu\pi}{2}}\frac{1}{\theta^{\frac{1}{6}}}$$

$$(3.6)$$

Here, we have used the first order expansion (3.4) with : $e^{\frac{2}{3}i\pi} \cdot e^{-\frac{i\pi}{6}} = i\left(\frac{4\pi}{6} - \frac{\pi}{6} = \frac{3\pi}{6} = \frac{\pi}{2}\right)$ and the fact that here

$$\epsilon z = z - \nu$$

= $i\sqrt{4\theta} - i\sqrt{4\theta - 1} = i\sqrt{4\theta} \left(1 - \sqrt{1 - \frac{1}{4\theta}}\right) \xrightarrow[\theta \to \infty]{} 0$

iii) We now study the denominator of (3.2)

$$D = K_{i\nu}(\sqrt{4\theta}) - \sqrt{4\theta} K'_{i\nu}(\sqrt{4\theta})$$

= $K_{i\nu}(\sqrt{4\theta}) + \sqrt{4\theta} K_{i\nu-1}(\sqrt{4\theta}) + i\sqrt{4\theta-1} K_{i\nu}(\sqrt{4\theta})$ (3.7)

(after using $\frac{d}{dz}(z^{\mu}K_{\mu}(z) = -z^{\mu}K_{\mu-1}(z))$; [Leb], p. 110). Since we have already studied the asymptotic behavior of $K_{i\nu}(\sqrt{4\theta})$, it remains to study that of

$$\Delta(\theta) := \sqrt{4\theta} K_{i\nu-1}(\sqrt{4\theta}) + i\sqrt{4\theta-1} K_{i\nu}(\sqrt{4\theta})$$
(3.8)

 $\sqrt{4\theta} K_{i\nu-1}(\sqrt{4\theta}) + i\sqrt{4\theta-1} K_{i\nu}(\sqrt{4\theta})$ is real (from (2.18)) Now, since

and since $i\sqrt{4\theta-1} K_{i\nu}(\sqrt{4\theta})$ is purely imaginary, the development of $\Delta(\theta)$ as $\theta \to \infty$ is that of the real part of $\sqrt{4\theta} K_{i\nu-1}(\sqrt{4\theta})$. Then, we obtain :

$$\sqrt{4\theta} K_{i\nu-1}(\sqrt{4\theta}) = \sqrt{4\theta} \frac{i\pi}{2} e^{\frac{i\pi}{2}(i\nu-1)} H^{(1)}_{i\nu-1} \left(e^{\frac{i\pi}{2}} \sqrt{4\theta} \right)
= \sqrt{4\theta} \frac{\pi}{2} e^{-\frac{\nu\pi}{2}} H^{(1)}_{i\nu-1} \left(e^{\frac{i\pi}{2}} \sqrt{4\theta} \right)$$
(3.9)

from (3.3), by replacing $i\nu$ by $i\nu - 1$. This time, we use the second order development (3.4) with here : $z = e^{\frac{i\pi}{2}}\sqrt{4\theta}$ and :

$$i\nu - 1 = z(1 - \epsilon) = i\sqrt{4\theta}(1 - \epsilon), \text{ i.e. }: \quad \epsilon z = \epsilon e^{\frac{i\pi}{2}}\sqrt{4\theta}$$

or $\epsilon z = \epsilon e^{\frac{i\pi}{2}}\sqrt{4\theta} = 1 + \frac{1}{4\sqrt{\theta}} + o\left(\frac{1}{\sqrt{\theta}}\right) \quad \left(\nu = \sqrt{4\theta - 1}\right)$ (3.10)

We obtain :

$$\sqrt{4\theta} K_{i\nu-1}(\sqrt{4\theta}) = \sqrt{4\theta} \frac{\pi}{2} e^{-\frac{\nu\pi}{2}} \left(-\frac{2}{3}\pi\right) \times \cdots$$

$$\cdots \left\{ C_1 e^{\frac{2i\pi}{3}} \frac{1}{\left(e^{\frac{i\pi}{2}}\sqrt{4\theta}\right)^{\frac{1}{3}}} + C_2 e^{\frac{4i\pi}{3}} \left(1 + \frac{i}{4\sqrt{\theta}}\right) \cdot \frac{1}{\left(e^{\frac{i\pi}{2}}\sqrt{4\theta}\right)^{\frac{2}{3}}} + o\left(\frac{1}{\theta^{\frac{1}{3}}}\right) \right\} \quad (3.11)$$
with $C_1 = \left(\sin\frac{\pi}{3}\right) \Gamma\left(\frac{1}{3}\right) 6^{\frac{1}{3}}$ and $C_2 = \left(\sin\frac{2\pi}{3}\right) \Gamma\left(\frac{2}{3}\right) 6^{\frac{2}{3}} \quad (C_1, C_2 > 0)$

$$= -\frac{1}{3}\sqrt{4\theta} e^{-\frac{\nu\pi}{2}} \left\{ C_1 i \frac{1}{(\sqrt{4\theta})^{\frac{1}{3}}} - C_2 \left(1 + \frac{i}{4\sqrt{\theta}}\right) \frac{1}{(\sqrt{4\theta}^{\frac{2}{3}}} + o\left(\frac{1}{\theta^{\frac{1}{3}}}\right) \right\} \quad (3.12)$$

Considering now the real part of (3.12) to obtain :

$$\Delta(\theta) = \mathcal{R}e\left(\sqrt{4\theta} K_{i\nu-1}(\sqrt{4\theta}) + i\sqrt{4\theta-1} K_{i\nu}(\sqrt{4\theta})\right)$$

$$= \frac{1}{3}\sqrt{4\theta} e^{-\frac{\nu\pi}{2}} C_2\left(\frac{1}{(\sqrt{4\theta})^2} + o\left(\frac{1}{\theta^{\frac{1}{3}}}\right)\right)$$

$$\underset{\theta \to \infty}{\sim} \frac{2^{\frac{1}{3}}C_2}{3} e^{-\frac{\nu\pi}{2}} \theta^{\frac{1}{6}}$$
(3.13)

• We then gather (3.2), (3.6) and (3.13), we obtain :

$$E_{1}\left(e^{-\frac{\theta}{2}\Pi_{1}}\right) \sim \frac{2\frac{2^{\frac{3}{3}}}{3}C_{1}e^{-\frac{\nu\pi}{2}}\frac{1}{\theta^{\frac{1}{6}}}}{\frac{2^{\frac{2}{3}}}{3}C_{1}e^{-\frac{\nu\pi}{2}}\frac{1}{\theta^{\frac{1}{6}}} + \frac{2^{\frac{1}{3}}}{3}C_{2}e^{-\frac{\nu\pi}{2}}\theta^{\frac{1}{6}}}$$
$$\sim \frac{2^{\frac{4}{3}}C_{1}}{\theta^{-\frac{1}{3}}}\frac{1}{\theta^{\frac{1}{3}}} = \frac{2}{3^{\frac{1}{3}}}\frac{\Gamma(\frac{1}{3})}{\Gamma(\frac{2}{3})}\frac{1}{\theta^{\frac{1}{3}}}$$
(3.14)

from the explicit formulae for C_1 and C_2 . This is Theorem 3.1.

Remark 3.2 Using Theorem 1.3, we obtain, for $0 < x \le 1$:

$$E_x\left(e^{-\frac{\theta}{2}\Pi_1}\right) = 1 + x \frac{\frac{1}{2}K_{i\sqrt{4\theta-1}}(\sqrt{4\theta}) + \sqrt{\theta}K'_{i\sqrt{4\theta-1}}(\sqrt{4\theta})}{\frac{1}{2}K_{i\sqrt{4\theta-1}}(\sqrt{4\theta}) - \sqrt{\theta}K'_{i\sqrt{4\theta-1}}(\sqrt{4\theta})}$$

It now follows easily from (3.6) and (3.13) that :

$$E_x\left(e^{-\frac{\theta}{2}\Pi_1}\right) \underset{\theta \to \infty}{\longrightarrow} 1 - x \tag{3.15}$$

Now, on the other hand :

$$E_x\left(e^{-\frac{\theta}{2}\Pi_1}\right) \underset{\theta \to \infty}{\longrightarrow} P_x(\Pi_1 = 0) \tag{3.16}$$

We recall (and recover here) that $P_x(\Pi_1 = 0) = 1 - x$. This is point *ii*) of Theorem 1.2.

4 Extending the preceding results to the variables $\Pi_k^{(\rho,x)}$ with

$$\Pi_k^{(\rho,x)} := \int_0^\infty \left(e^{\rho(\log x + B_t - \frac{t}{2})} - k \right)^+ dt \tag{4.1}$$

 $(\rho, x, k > 0)$. In the preceding Sections, we have studied the case $\rho = 1$. The analogue of Theorem 1.3 may be stated as : **Theorem 4.1** The Laplace transform of $\Pi_1^{(\rho)}$ under P_x is given by :

$$E_{x}\left(e^{-\frac{\theta}{2}\Pi_{1}^{(\rho)}}\right) \begin{cases} \frac{\sqrt{x} K_{\gamma}\left(\frac{\sqrt{4\theta}}{\rho} x^{\frac{\theta}{2}}\right)}{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4\theta}}{\rho}\right) - \sqrt{\theta} K_{\gamma}'\left(\frac{\sqrt{4\theta}}{\rho}\right)} & \text{if } x \ge 1\\ 1 + x \frac{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4\theta}}{\rho}\right) + \sqrt{\theta} K_{\gamma}'\left(\frac{\sqrt{4\theta}}{\rho}\right)}{\frac{1}{2} K_{\gamma}\left(\frac{\sqrt{4\theta}}{\rho}\right) - \sqrt{\theta} K_{\gamma}'\left(\frac{\sqrt{4\theta}}{\rho}\right)} & \text{if } 0 < x \le 1 \end{cases}$$

$$(4.2)$$

with $\gamma = \frac{\sqrt{1-4\theta}}{\rho}$ if $\theta \leq \frac{1}{4}$ and $\gamma = i \frac{\sqrt{4\theta-1}}{\rho}$ if $\theta > \frac{1}{4}$. The proof of this Theorem 4.1 is quite similar to that of Theorem 1.3. It hinges upon the following Lemma 4.2.

Lemma 4.2

i) Let, for $\theta \ge 0$ and $\rho > 0$, the function $\varphi_{\theta}^{(\rho)} : \mathbb{R}_+ \to \mathbb{R}$ defined by :

$$\varphi_{\theta}^{(\rho)}(y) = \sqrt{y} K_{\gamma} \left(\frac{\sqrt{4\theta}}{\rho} y^{\frac{\rho}{2}}\right) \qquad (y \ge 0)$$
(4.3)

Then, it satisfies :

$$\varphi_{\theta}^{(\rho)''}(y) + \theta \left(-y^{\rho-2} + \frac{1}{y^2}\right)\varphi_{\theta}^{(\rho)}(y) = 0$$

$$(4.4)$$

(see [Pet], p. 306, with $\rho - 2 = m$, formula (8), after taking care of replacing a by -a). ii) Let $\tilde{\varphi}_{\theta}^{(\rho)} : \mathbb{R}_+ \to \mathbb{R}$ defined by :

$$\widetilde{\varphi}_{\theta}^{(\rho)}(y) = \begin{cases} \varphi_{\theta}^{(\rho)}(y) & \text{if } y \ge 1\\ \varphi_{\theta}^{(\rho)}(1) - \varphi_{\theta}^{(\rho)'}(1) + y \varphi_{\theta}^{(\rho)'}(1) & \text{if } y \le 1 \end{cases}$$

$$(4.5)$$

Then $\widetilde{\varphi}_{\theta}^{(\rho)}$ is positive, decreasing, convex and satisfies :

$$\widetilde{\varphi}_{\theta}^{(\rho)}(y) + \theta \left(-y^{\rho-2} + \frac{1}{y^2}\right) \mathbf{1}_{y\geq 1} \ \widetilde{\varphi}_{\theta}^{(\rho)}(y) = 0 \tag{4.6}$$

Remark 4.1 As a check, we note that formula (4.2) allows to recover the identity (1.8):

$$\int_0^\infty e^{\rho B_t - \frac{\rho t}{2}} dt \stackrel{\text{(law)}}{=} \frac{2}{\rho^2 \gamma_{\frac{1}{\rho}}}$$

Indeed, on one hand :

$$E\left(e^{-\frac{\theta}{2}\frac{2}{\rho^{2}\gamma_{\frac{1}{\rho}}}}\right) = \frac{1}{\Gamma\left(\frac{1}{\rho}\right)} \int_{0}^{\infty} e^{-\frac{\theta}{\rho^{2}z}-z} z^{\frac{1}{\rho}-1} dz = \frac{2}{\Gamma\left(\frac{1}{\rho}\right)} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4\theta}}{\rho}\right) \left(\frac{\theta}{\rho^{2}}\right)^{\frac{1}{2\rho}} (\text{see [Leb], p. 115})$$

On the other hand :

$$E_{1}\left(e^{-\frac{\theta}{2}\Pi_{0}^{(\rho)}}\right) = \lim_{\epsilon \downarrow 0} E_{1}\left(\exp -\frac{\theta}{2}\int_{0}^{\infty}\left(e^{\rho B_{t}-\frac{\rho t}{2}}-\varepsilon\right)^{+}dt\right)$$
$$= \lim_{\epsilon \downarrow 0} E_{1}\left(\exp -\frac{\theta \varepsilon}{2}\int_{0}^{\infty}\left(e^{\rho\left(\frac{\log \frac{1}{\varepsilon}}{\rho}+B_{t}\right)-\frac{\rho t}{2}}-1\right)^{+}dt\right)$$
$$= \lim_{\epsilon \downarrow 0}\frac{\left(\frac{1}{\varepsilon}\right)^{\frac{1}{2\rho}}K_{\frac{\sqrt{1-4\theta\varepsilon}}{\rho}}\left(\frac{\sqrt{4\theta\varepsilon}}{\rho}\right)\left(\frac{1}{\varepsilon}\right)^{\frac{1}{\rho}\frac{\theta}{2}}}{\frac{1}{2}K_{\frac{\sqrt{1-4\theta\varepsilon}}{\rho}}\left(\frac{\sqrt{4\theta\varepsilon}}{\rho}\right)-\sqrt{\theta\varepsilon}K'_{\frac{\sqrt{1-4\theta\varepsilon}}{\rho}}\left(\frac{\sqrt{4\theta\varepsilon}}{\rho}\right)}$$

from (4.2), replacing θ by $\theta \varepsilon$ and x by $\left(\frac{1}{\varepsilon}\right)^{\frac{1}{\rho}}$

$$\begin{array}{l} \underset{\varepsilon \to 0}{\sim} & \frac{\left(\frac{1}{\varepsilon}\right)^{\frac{1}{2\rho}} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4\theta}}{\rho}\right)}{\frac{1}{2} K_{\frac{\sqrt{1-4\theta\varepsilon}}{\rho}}\left(\frac{\sqrt{4\theta\varepsilon}}{\rho}\right) - \frac{\rho}{2}\left(\frac{\sqrt{4\theta\varepsilon}}{\rho}\right) K'_{\frac{\sqrt{1-4\theta\varepsilon}}{\rho}}\left(\frac{\sqrt{4\theta\varepsilon}}{\rho}\right)}{\frac{\rho}{2} \left[\frac{\sqrt{4\theta\varepsilon}}{\rho} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4\theta}}{\rho}\right)}{\frac{\rho}{2} \left[\frac{\sqrt{4\theta\varepsilon}}{\rho} K_{\frac{1}{\rho}+1}\left(\frac{\sqrt{4\theta\varepsilon}}{\rho}\right)\right]} & \text{(since } zK'\nu = \nu K_{\nu}(z) - z K_{\nu+1}(z)) \\ \underset{\varepsilon \to 0}{\sim} & \frac{\left(\frac{1}{\varepsilon}\right)^{\frac{1}{2\rho}} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4\theta}}{\rho}\right)}{\frac{\rho}{2} \left(\Gamma\left(\frac{1}{\rho}+1\right)\left(\frac{\sqrt{4\theta\varepsilon}}{\rho}\right)^{-\frac{1}{\rho}}} & \text{(since } K_{\nu}(z) \sim \frac{1}{2} \Gamma(\nu)\left(\frac{z}{2}\right)^{-1}\right) \\ = & \frac{\left(\frac{1}{\varepsilon}\right)^{\frac{1}{2\rho}} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4\theta}}{\rho}\right)}{\Gamma\left(\frac{1}{\rho}\right)} \frac{\varepsilon^{\frac{1}{2\rho}} \theta^{\frac{1}{2\rho}}}{\rho^{\frac{1}{\rho}}} \\ \underset{\varepsilon \downarrow 0}{\leftarrow} & \left(\frac{\theta}{\rho^{2}}\right)^{\frac{1}{2\rho}} \frac{2}{\Gamma\left(\frac{1}{\rho}\right)} K_{\frac{1}{\rho}}\left(\frac{\sqrt{4\theta}}{\rho}\right) = E\left(e^{-\frac{\theta}{2} \Pi_{0}^{(\rho)}}\right) \end{array}$$

Remark 4.2

Taking up again the arguments of the proof of Theorem 3.1, it is not difficult to see that :

$$E_1\left(e^{-\frac{\theta}{2}\Pi_1^{(\rho)}}\right) \underset{\theta \to \infty}{\sim} \frac{C(\rho)}{\theta^{\frac{1}{3}}}$$

where $C(\rho)$ is a strictly positive constant, depending on ρ . We then deduce that, for $\alpha < 0$:

$$E_1((\Pi_1^{(\rho)})^{\alpha}) < \infty$$
 if and only if $|\alpha| < \frac{1}{3}$

On the other hand, it is not difficult to see that :

- If x > 1, for all α < 0, E₁((Π₁^(ρ))^α) < ∞.
 If x < 1, for all α < 0, E₁((Π₁^(ρ))^α) = +∞.
 Concerning the positive moments of Π₁^(ρ):

• If
$$0 < \alpha < \frac{1}{\rho}$$
, then $E_x\left((\Pi_1^{(\rho)})^{\alpha}\right) < \infty$ (4.7)

Indeed (4.7), for x = 1, follows from :

$$E_1((\Pi_1^{(\rho)})^{\alpha}) \leq E_1((\Pi_0^{(\rho)})^{\alpha}) = \frac{1}{\Gamma(\frac{1}{\rho})} \int_0^\infty \left(\frac{2}{\rho^2 z}\right)^{\alpha} e^{-z} z^{\frac{1}{\rho}-1} dz$$
$$< \infty \quad \text{if} \quad \alpha < \frac{1}{\rho}$$
$$\left(\text{since} \qquad \Pi_0 \stackrel{(\text{law})}{=} \frac{2}{\rho^2 \gamma_{\frac{1}{\rho}}}, \quad \text{from (1.10)}\right)$$

The fact that $E_x((\Pi_1^{(\rho)})^{\alpha}) < \infty$ for every x > 0, and every $\alpha < \frac{1}{\rho}$ may be obtained by using arguments close to those used in the proof of Theorem 1.1. We believe that, for all $\rho > 0$:

$$E_x((\Pi_1^{(\rho)})^{\frac{1}{\rho}}) = +\infty.$$
 (4.8)

We now show (4.8), when $\rho > 1$

It then suffices, by using the arguments of the proof of Theorem 1.1, to see that (4.8) is true when x = 1.

• We first show that : $1 - E_1 \left(e^{-\frac{\theta}{2} \Pi_1^{(\rho)}} \right) \underset{\theta \to 0}{\sim} C \theta^{\frac{1}{\rho}}$ (4.9)

Indeed, from Theorem 4.1, we have :

$$E_1\left(\exp-\frac{\theta}{2}\Pi_1^{(\rho)}\right) = 1 + \frac{\frac{1}{2}K_\gamma\left(\frac{\sqrt{4\theta}}{\rho}\right) + \sqrt{\theta}K_\gamma'\left(\frac{\sqrt{4\theta}}{\rho}\right)}{\frac{1}{2}K_\gamma\left(\frac{\sqrt{4\theta}}{\rho}\right) - \sqrt{\theta}K_\gamma'\left(\frac{\sqrt{4\theta}}{\rho}\right)}$$

with $\gamma = \frac{\sqrt{1-4\theta}}{\rho}$ (and $\theta \le \frac{1}{4}$). Thus :

$$1 - E_1\left(e^{-\frac{\theta}{2}\Pi_1^{(\rho)}}\right) = -\frac{\frac{1}{2}K_\gamma\left(\frac{\sqrt{4\theta}}{\rho}\right) + \frac{\rho}{2}\left\{\frac{\sqrt{4\theta}}{\rho}K_\gamma'\left(\frac{\sqrt{4\theta}}{\rho}\right)\right\}}{\frac{1}{2}K_\gamma\left(\frac{\sqrt{4\theta}}{\rho}\right) - \frac{\rho}{2}\left\{\frac{\sqrt{4\theta}}{\rho}K_\gamma'\left(\frac{\sqrt{4\theta}}{\rho}\right)\right\}}$$
$$= -\frac{\frac{1}{2}K_{\frac{\sqrt{1-4\theta}}{\rho}}\left(\frac{\sqrt{4\theta}}{\rho}\right)\left(1 - \frac{\sqrt{1-4\theta}}{\rho}\right) - \frac{\sqrt{\theta}}{\rho}K_{\frac{\sqrt{1-4\theta}}{\rho}-1}\left(\frac{\sqrt{4\theta}}{\rho}\right)}{\frac{1}{2}K_{\frac{\sqrt{1-4\theta}}{\rho}}\left(\frac{\sqrt{4\theta}}{\rho}\right)\left(1 + \frac{\sqrt{1-4\theta}}{\rho}\right) + \frac{\sqrt{\theta}}{\rho}K_{\frac{\sqrt{1-4\theta}}{\rho}-1}\left(\frac{\sqrt{4\theta}}{\rho}\right)}$$

(after using $zK'_{\mu}(z) = -\mu K_{\mu}(z) - z K_{\mu-1}(z)$). Since $\rho > 1$, we replace $K_{\frac{\sqrt{1-4\theta}}{\rho}-1}$ by $K_{1-\frac{\sqrt{1-4\theta}}{\rho}}$ (since $K_{\mu} = K_{-\mu}$) and we deduce, from :

$$K_{1-\frac{\sqrt{1-4\theta}}{\rho}} \xrightarrow[\theta \to 0]{} K_{1-\frac{1}{\rho}}, \quad \text{with} \quad 1-\frac{1}{\rho} > 0 \; :$$

that :

$$1 - E_1 \left(e^{-\frac{\theta}{2} \Pi_1^{(\rho)}} \right) \sim \frac{a \, \theta^{1 - \frac{1}{2\rho}} + b \, \theta^{\frac{1}{2\rho}}}{a' \, \theta^{-\frac{1}{2\rho}} + b' \, \theta^{\frac{1}{2\rho}}} \left(\text{since, for } \mu > 0 \, K_\mu(z) \mathop{\sim}_{z \to 0} C_\mu \, z^{-\mu} \, ; \, \text{see [Leb], p. 136} \right)$$
$$\sim \frac{b}{\theta \to 0} \quad \frac{b}{a'} \, \theta^{\frac{1}{\rho}} \qquad \left(\text{since } \rho \ge 1 \text{ implies } 1 - \frac{1}{2\rho} \ge \frac{1}{2\rho} \right)$$

• From (4.9) we deduce :

$$\int_0^\infty e^{-\frac{\theta t}{2}} P(\Pi_1^{(\rho)} \ge t) dt = \frac{2}{\theta} \left(1 - E(e^{-\frac{\theta}{2} \Pi_1^{(\rho)}}) \right)$$
$$\sim_{\theta \to \infty} C' \theta^{\frac{1}{\rho} - 1}$$

Hence, from the Tauberian Theorem :

$$P(\Pi_1^{(\rho)} \ge t) \underset{t \to \infty}{\sim} C'' \frac{1}{t^{\frac{1}{\rho}}}$$
(4.10)

and

$$E((\Pi_{1}^{(\rho)})^{\frac{1}{\rho}}) = \int_{0}^{\infty} P((\Pi_{1}^{(\rho)})^{\frac{1}{\rho}} \ge t) dt$$
$$= \int_{0}^{\infty} P(\Pi_{1}^{(\rho)} \ge t^{\rho}) dt = +\infty, \quad \text{from (4.10)}$$

We note that, for $\rho < 1$, the preceding argument cannot be applied since :

$$\int_{1}^{\infty} e^{-\frac{\theta t}{2}} \frac{dt}{t^{\frac{1}{\rho}}} \xrightarrow[\theta \to 0]{} \int_{1}^{\infty} \frac{dt}{t^{\frac{1}{\rho}}} < \infty.$$

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