



HAL
open science

SHARP DECAY ESTIMATES AND VANISHING VISCOSITY FOR DIFFUSIVE HAMILTON-JACOBI EQUATIONS

Said Benachour, Matania Ben-Artzi, Philippe Laurençot

► **To cite this version:**

Said Benachour, Matania Ben-Artzi, Philippe Laurençot. SHARP DECAY ESTIMATES AND VANISHING VISCOSITY FOR DIFFUSIVE HAMILTON-JACOBI EQUATIONS. 2008. hal-00337773

HAL Id: hal-00337773

<https://hal.science/hal-00337773>

Preprint submitted on 10 Nov 2008

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

SHARP DECAY ESTIMATES AND VANISHING VISCOSITY FOR DIFFUSIVE HAMILTON-JACOBI EQUATIONS

SAÏD BENACHOUR, MATANIA BEN-ARTZI, AND PHILIPPE LAURENÇOT

ABSTRACT. Sharp temporal decay estimates are established for the gradient and time derivative of solutions to the Hamilton-Jacobi equation $\partial_t v_\varepsilon + H(|\nabla_x v_\varepsilon|) = \varepsilon \Delta v_\varepsilon$ in $\mathbb{R}^N \times (0, \infty)$, the parameter ε being either positive or zero. Special care is given to the dependence of the estimates on ε . As a by-product, we obtain convergence of the sequence (v_ε) as $\varepsilon \rightarrow 0$ to a viscosity solution, the initial condition being only continuous and either bounded or non-negative. The main requirement on H is that it grows superlinearly or sublinearly at infinity, including in particular $H(r) = r^p$ for $r \in [0, \infty)$ and $p \in (0, \infty)$, $p \neq 1$.

1. Introduction

The purpose of this paper is to derive temporal decay estimates for the gradient and the time derivative of viscosity solutions to the Hamilton-Jacobi equation

$$(1.1) \quad \partial_t v + H(|\nabla_x v|) = 0, \quad (x, t) \in \mathbb{R}^N \times (0, \infty),$$

$$(1.2) \quad v(x, 0) = \varphi(x), \quad x \in \mathbb{R}^N.$$

and its diffusive counterpart

$$(1.3) \quad \partial_t v_\varepsilon - \varepsilon \Delta v_\varepsilon + H(|\nabla_x v_\varepsilon|) = 0, \quad (x, t) \in \mathbb{R}^N \times (0, \infty), \quad \varepsilon > 0,$$

$$(1.4) \quad v_\varepsilon(x, 0) = \varphi(x), \quad x \in \mathbb{R}^N,$$

under suitable assumptions on the Hamiltonian function H and for initial data φ which are continuous but not necessarily uniformly continuous (and in some cases not even bounded). The main feature of our analysis is that we carefully trace the dependence on the “viscosity” parameter ε in the estimates of the space-time gradients of v_ε . We obtain estimates which do not deteriorate as $\varepsilon \rightarrow 0$ and thus reflect the regularizing effect of the nonlinear term $H(|\nabla_x v_\varepsilon|)$. As a by-product of our analysis, we may perform the limit $\varepsilon \rightarrow 0$ (the so-called vanishing viscosity limit) and show the convergence of the solutions v_ε to the nonlinear parabolic equation (1.3)-(1.4) without requiring much on the initial condition (besides continuity and either boundedness or only non-negativity). The limiting solutions we obtain are “viscosity solutions” in the sense of Crandall & Lions [11], and we refer to [2, 5, 12, 18] for extensive discussions of these solutions and to [13, Chapter 10] for the connection between viscosity solutions and the “vanishing viscosity” approach.

Date: November 10, 2008.

1991 Mathematics Subject Classification. Primary 35F25; Secondary 35K15, 35K55, 49L25.

Key words and phrases. Hamilton-Jacobi equation, Bernstein estimates, general initial data.

M. Ben-Artzi thanks the Institut Elie Cartan, Université de Nancy I, for the hospitality during February 2000, when this work was started.

The main tool in this work consists of uniform (with respect to ε) estimates of the (space-time) gradient of v_ε . These estimates enable us to treat the more general initial data φ as mentioned above. Roughly speaking, the main requirement placed on our Hamiltonian function $H = H(r)$, $0 \leq r < \infty$, is that it grows either “superlinearly” or “sublinearly” as $r \rightarrow \infty$. More precisely, the basic set of assumptions (1.5)-(1.6) on H is the following.

$$(1.5) \quad \left\{ \begin{array}{l} \bullet H \text{ is continuous nonnegative on } [0, \infty) \text{ and } H(0) = 0. \\ \text{In addition, } H \text{ is locally Lipschitz continuous on } (0, \infty). \\ \bullet \text{ There exists a family of nonnegative smooth functions} \\ \{\Phi_\eta\}_{\eta>0} \text{ defined in } [0, \infty) \text{ such that} \\ (i) \quad \Phi_\eta(0) = 0 \text{ for all } \eta > 0. \\ (ii) \quad \Phi_\eta(r^2) \xrightarrow{\eta \rightarrow 0^+} H(r), \text{ uniformly in compact intervals of } [0, \infty). \end{array} \right.$$

Definition 1.1. Consider the family of functions $\{\Theta_\eta\}_{\eta>0}$ defined by

$$\Theta_\eta(r) = 2r\Phi'_\eta(r) - \Phi_\eta(r), \quad (r, \eta) \in [0, \infty) \times (0, \infty).$$

Let $p \in (0, 1) \cup (1, \infty)$. We say that H satisfies the p -condition if there exist $\gamma > 0$, $a > 0$, $b > 0$, such that, for $r > 0$ and sufficiently small $\eta > 0$,

$$\left\{ \begin{array}{l} (i) \quad \Theta_\eta(r) \geq ar^{\frac{p}{2}} - b\eta^\gamma, \quad \text{if } p \in (1, \infty), \\ (ii) \quad \Theta_\eta(r) \leq -ar^{\frac{p}{2}} + b\eta^\gamma, \quad \text{if } p \in (0, 1). \end{array} \right.$$

Our third basic assumption is

$$(1.6) \quad \text{For some } p \in (0, 1) \cup (1, \infty), \quad H \text{ satisfies the } p\text{-condition.}$$

As we show in Appendix A to this paper, the prototypical example

$$(1.7) \quad H(r) = r^p, \quad p \in (0, 1) \cup (1, \infty),$$

satisfies the above assumptions with $\Phi_\eta(r) = (r + \eta^2)^{\frac{p}{2}} - \eta^p$. In fact, the same argument shows that one can take

$$(1.8) \quad H(r) = \sum_{k=1}^m \mu_k r^{p_k}, \quad \mu_k > 0,$$

where either $\{p_1, \dots, p_m\} \in (0, 1)^m$, or $\{p_1, \dots, p_m\} \in (1, \infty)^m$.

We can easily extend further this special case, as shown by the following example.

Proposition 1.2. Let $p > 1$ and let G be a smooth function supported in $[r_0, \infty)$ for some $r_0 > 0$. Assume that for some $q \geq p$ and $\lambda > 0$ we have

$$\frac{d}{dr} \left(\frac{G(r)}{r} \right) \geq \lambda r^{q-2}, \quad r > r_0.$$

Then the function $H : r \mapsto r^p + G(r)$ satisfies all the assumptions (1.5)-(1.6). In particular, we can take $G(r) = (r - r_0)_+^q$.

Proof. It suffices, by the above remarks, to consider only the part of G . By taking $\Phi_\eta(r^2) \equiv G(r)$ we get

$$\Theta_\eta(r) = 2r\Phi'_\eta(r) - \Phi_\eta(r) = \sqrt{r}G'(\sqrt{r}) - G(\sqrt{r}) \geq \lambda(\sqrt{r})^q \geq \lambda r_0^{\frac{q-p}{2}} r^{\frac{p}{2}}.$$

□

From now on, we assume that all special Hamiltonians H satisfy a p -condition for some $p \in (0, 1) \cup (1, \infty)$. Our initial function φ is assumed only to be bounded from below and can be taken in $C(\mathbb{R}^N)$, the space of real-valued continuous functions on \mathbb{R}^N if $1 < p < \infty$ whereas, if $0 < p < 1$, it is taken in $C_b(\mathbb{R}^N)$, the space of bounded continuous functions.

Under these conditions we obtain the existence and uniqueness of a solution to (1.1)-(1.2) in $\mathbb{R}^N \times [0, \infty)$ in Theorems 3.1 and 3.3, provided a suitable comparison principle is available in the following sense.

Definition 1.3. (a) We say that Equation (1.1) satisfies the **(discontinuous) comparison principle** if the following condition holds: Let $v_1 \in USC(\mathbb{R}^N \times (0, \infty))$ (resp. $v_2 \in LSC(\mathbb{R}^N \times (0, \infty))$) be a viscosity subsolution (resp. supersolution) of (1.1). Assume that $v_1(x, 0) \leq v_2(x, 0)$ for $x \in \mathbb{R}^N$ and that $\inf_{\mathbb{R}^N \times (0, \infty)} v_2 > -\infty$. Then $v_1 \leq v_2$ in $\mathbb{R}^N \times (0, \infty)$.

(b) We say that Equation (1.1) satisfies the **comparison principle in $C_b(\mathbb{R}^N)$** if the following condition holds: Let $v_1 \in C_b(\mathbb{R}^N \times (0, \infty))$ (resp. $v_2 \in C_b(\mathbb{R}^N \times (0, \infty))$) be a viscosity subsolution (resp. supersolution) of (1.1). Assume that $v_1(x, 0) \leq v_2(x, 0)$ for $x \in \mathbb{R}^N$. Then $v_1 \leq v_2$ in $\mathbb{R}^N \times (0, \infty)$.

Here, $USC(\mathbb{R}^N \times (0, \infty))$ and $LSC(\mathbb{R}^N \times (0, \infty))$ denote the space of upper and lower semicontinuous functions in $\mathbb{R}^N \times (0, \infty)$, respectively. We refer to [16] for conditions that imply the (discontinuous) comparison principle. For instance, if H is convex, Equation (1.1) satisfies the (discontinuous) comparison principle. This applies in particular to $H(\xi) = |\xi|^p$ for $p > 1$. Concerning the case $H(\xi) = |\xi|^p$ for $p \in (0, 1)$, Equation (1.1) satisfies the comparison principle in $C_b(\mathbb{R}^N)$ as recalled in Appendix C [6].

While the comparison principle seems to be the most effective tool in the study of uniqueness (for equations of the type considered here) we mention the proof in [20] concerning the uniqueness of the solution obtained by the Lax-Hopf formula.

2. Notation

Throughout the paper, we shall make use of the following standard functional notation.

The space $C^{2,1}(\mathbb{R}^N \times (0, \infty))$ is the space of all functions $u = u(x, t)$ which are twice continuously differentiable with respect to $x \in \mathbb{R}^N$ and once with respect to $t \in (0, \infty)$.

The space $C_b^2(\mathbb{R}^N)$ is the space of all twice continuously differentiable functions f such that all their derivatives up to second order are bounded (i.e., in $C_b(\mathbb{R}^N)$).

The space $W^{1,\infty}(\mathbb{R}^N)$ is the space of functions having uniformly bounded (distributional) first order derivatives (i.e., using Rademacher's theorem, uniformly Lipschitz continuous functions).

The norm in $L^q(\mathbb{R}^N)$ is denoted by $\|\cdot\|_q$, $q \in [1, \infty]$.

3. Results

The existence and uniqueness results for solutions to (1.3)-(1.4) are recalled in Proposition 4.2 below. When the initial function is bounded, these solutions converge to a viscosity solution to (1.1)-(1.2), as expressed in the following theorem.

Theorem 3.1. *Let $p \in (0, \infty)$, $p \neq 1$, and let $\varphi \in C_b(\mathbb{R}^N)$. Assume that H satisfies the hypotheses (1.5) and (1.6) and that Equation (1.1) satisfies the comparison principle in $C_b(\mathbb{R}^N)$ (cf. Definition 1.3 (b)). Assume also that H satisfies the growth condition*

$$(3.1) \quad H(r) \leq \tilde{H}(r) := g_H (r^{\kappa_\infty} + r^{\kappa_0}), \quad 0 < r < \infty, \quad 0 < \kappa_\infty \leq \kappa_0,$$

where $g_H > 0$ is a constant. Then the solutions v_ε to (1.3)-(1.4) converge as $\varepsilon \rightarrow 0$ towards the unique (viscosity) solution v to (1.1)-(1.2), uniformly in every compact subset of $\mathbb{R}^N \times (0, \infty)$. The function v is differentiable a.e. in $\mathbb{R}^N \times (0, \infty)$ and satisfies (1.1) at any point of differentiability. Furthermore, we have, for a.e. $(x, t) \in \mathbb{R}^N \times (0, \infty)$,

$$(3.2) \quad |\nabla_x v(x, t)| \leq \lambda_p \|\varphi\|_\infty^{\frac{1}{p}} (at)^{-\frac{1}{p}},$$

$$(3.3) \quad 0 \geq \partial_t v(x, t) \geq -L t^{-\mu},$$

where $\lambda_p = 1$ if $p > 1$ and $\lambda_p = (2/p)^{\frac{1}{p}}$ if $p < 1$,

$$\mu = \begin{cases} \frac{\kappa_0}{p} & \text{if } 0 < t \leq 1, \\ \frac{\kappa_\infty}{p} & \text{if } 1 < t < \infty, \end{cases}$$

and

$$L = g_H \left\{ (2\lambda_p^p \|\varphi\|_\infty a^{-1})^{\frac{\kappa_0}{p}} + (2\lambda_p^p \|\varphi\|_\infty a^{-1})^{\frac{\kappa_\infty}{p}} \right\}.$$

As our aim in this paper is to derive estimates for the solutions v_ε to (1.3)-(1.4) (almost) independent of ε , the estimate (3.3) is obtained by passing to the limit as $\varepsilon \rightarrow 0$ in an analogous estimate for v_ε (see Proposition 4.3 below). However, an alternative and simpler proof (with a slightly better constant than L) relies on (3.1), (3.2), and the fact that v solves (1.1)-(1.2) almost everywhere. Indeed, we infer from these properties that

$$\begin{aligned} \partial_t v &= -H(|\nabla_x v|) \geq -g_H (|\nabla_x v|^{\kappa_\infty} + |\nabla_x v|^{\kappa_0}) \\ &\geq -g_H \left\{ (\lambda_p^p \|\varphi\|_\infty a^{-1})^{\frac{\kappa_0}{p}} + (\lambda_p^p \|\varphi\|_\infty a^{-1})^{\frac{\kappa_\infty}{p}} \right\} t^{-\mu}, \end{aligned}$$

the parameter μ being defined in Theorem 3.1. A further comment in that direction is that the vanishing viscosity approach used here (and already used in [19]) is not the only route towards gradient or time derivative estimates, see, e.g., [3, 4, 17, 18].

Remark 3.2. *If $H(r) = r^p$ for $p > 1$, we have $\kappa_\infty = \kappa_0 = p$ and (3.3) indicates that $\partial_t v \geq -C/t$ for $t \geq 1$ for some positive constant C depending on N , p and $\|\varphi\|_\infty$. It gives a temporal decay rate for large times of the same order as that obtained in [9] where the inequality $\partial_t v \geq -v/((p-1)t)$ (in the sense of distributions) is established by using the homogeneity of the Hamiltonian H .*

We now turn to the case where the initial function φ is continuous but not necessarily bounded. Thus, in contrast to the previous theorem, where the positivity of φ was not essential (as φ could be replaced by $\varphi + c$), the positivity assumption (or rather the requirement that φ be bounded from below) in the following theorem is essential. Also, we need to impose an additional growth assumption on H , namely that H fulfills the p -condition (1.6) with $p > 1$.

Theorem 3.3. *Let $0 \leq \varphi \in C(\mathbb{R}^N)$ and assume that H satisfies the hypotheses (1.5) and (1.6) with $p > 1$, together with (3.1). Assume also that Equation (1.1) satisfies the (discontinuous) comparison principle (cf. Definition 1.3 (a)). Then the solutions v_ε to (1.3)-(1.4) converge as $\varepsilon \rightarrow 0$ towards the unique (viscosity) solution v to (1.1)-(1.2), uniformly in every compact subset of $\mathbb{R}^N \times (0, \infty)$. The function v belongs to $W_{loc}^{1,\infty}(\mathbb{R}^N \times (0, \infty))$ and satisfies (1.1) as well as (3.2), (3.3) for a.e. $(x, t) \in \mathbb{R}^N \times (0, \infty)$.*

Remark 3.4. *When dealing with the Hamilton-Jacobi equation (1.1) it seems unavoidable (as is the case in the references cited in this paper) to have a rather long list of assumptions on the Hamiltonian H . Furthermore, some results depend only on partial lists of the assumptions, in addition to the interplay between the degree of generality assumed on the initial data (and solutions) and the corresponding assumptions. We therefore emphasize that Theorem 3.1 is applicable to the case of sums of powers as in (1.8), while Theorem 3.3 is applicable to the case (1.8) when $\{p_1, \dots, p_m\} \in (1, \infty)^m$.*

Remark 3.5. *Using a condition similar to our assumption (1.6) Lions [19, Section IV] obtains viscosity solutions for bounded, lower semicontinuous initial data.*

4. The viscous Hamilton-Jacobi equation

We first draw useful consequences of (1.6) and (3.1) on H .

Lemma 4.1. *Assume that H fulfills (1.5) and (1.6). Then*

$$(4.1) \quad H(r) \geq \frac{a}{|p-1|} r^p, \quad r \geq 0.$$

If H also satisfies (3.1) then $\kappa_\infty \leq p \leq \kappa_0$.

Proof. Assume first that $p > 1$ in (1.6). Then, if $r > 0$, $\delta \in (0, r)$, $s \in (\delta, r)$ and $\eta > 0$, we infer from (1.6) that

$$\frac{d}{ds} \left(\frac{\Phi_\eta(s^2)}{s} \right) = \frac{\Theta_\eta(s^2)}{s^2} \geq a s^{p-2} - b\eta^\gamma s^{-2}.$$

Integrating over (δ, r) with respect to s gives

$$\frac{\Phi_\eta(r^2)}{r} \geq \frac{\Phi_\eta(\delta^2)}{\delta} + \frac{a}{p-1} (r^{p-1} - \delta^{p-1}) + b\eta^\gamma \left(\frac{1}{r} - \frac{1}{\delta} \right),$$

whence, after letting $\eta \rightarrow 0$ and using the nonnegativity of H ,

$$\frac{H(r)}{r} \geq \frac{H(\delta)}{\delta} + \frac{a}{p-1} (r^{p-1} - \delta^{p-1}) \geq \frac{a}{p-1} (r^{p-1} - \delta^{p-1}).$$

Passing to the limit as $\delta \rightarrow 0$ gives (4.1). The proof is similar for $p \in (0, 1)$ except than one integrates over (r, A) for $A > r$ and then let $A \rightarrow \infty$.

Next, if H also satisfies (3.1), the claimed constraints on κ_0 and κ_∞ readily follow from (4.1) by looking at the behavior for small r and large r . \square

We next recall the basic existence and uniqueness theorem for regular initial data [1]. In fact, the result there refers to the special case (1.7). However the same method of proof can be used in order to obtain the following proposition [21].

Proposition 4.2. *Let $0 \leq \varphi \in C_b^2(\mathbb{R}^N)$ and H satisfy (1.5). Then the Cauchy problem (1.3)-(1.4) has a unique global solution v_ε such that*

- (1) $v_\varepsilon \in C^{2,1}(\mathbb{R}^N \times [0, \infty))$,
- (2) $0 \leq v_\varepsilon(x, t) \leq \|\varphi\|_\infty$, $(x, t) \in \mathbb{R}^N \times (0, \infty)$,
- (3) $\|\nabla_x v_\varepsilon(\cdot, t)\|_\infty \leq \|\nabla_x \varphi\|_\infty$, $t > 0$.

Observe that the assumption $0 \leq \varphi$ entails no loss of generality as φ can be replaced by $\varphi + \|\varphi\|_\infty$ without changing the equation. Proposition 4.2 is actually valid assuming only the first property in (1.5).

A remarkable fact (which is crucial in our study) concerning the solution v_ε is that its gradient can be estimated *independently* of ε while only a mild dependence on ε shows up for its time derivative. Such estimates are obtained by the ‘‘Bernstein method’’ [8], namely, using the comparison principle for a certain function of $\nabla_x v_\varepsilon$ or $\partial_t v_\varepsilon$. The estimates needed in this paper are gathered in the following proposition.

Proposition 4.3. *Let $p \neq 1$ and $0 \leq \varphi \in C_b^2(\mathbb{R}^N)$. Assume that H satisfies the hypotheses (1.5) and (1.6). Then the solution v_ε to (1.3)-(1.4) satisfies, with p as in (1.6),*

$$(4.2) \quad \|\nabla_x v_\varepsilon(\cdot, t)\|_\infty \leq \lambda_p \|\varphi\|_\infty^{\frac{1}{p}} (at)^{-\frac{1}{p}}, \quad t > 0,$$

where $\lambda_p = 1$ if $p > 1$ and $\lambda_p = (2/p)^{\frac{1}{p}}$ if $p < 1$.

In addition, if $p > 1$,

$$(4.3) \quad \left\| \nabla_x \left(v_\varepsilon^{\frac{p-1}{p}} \right) (\cdot, t) \right\|_\infty \leq \mu_p t^{-\frac{1}{p}}, \quad t > 0,$$

where $\mu_p = (p-1)a^{-\frac{1}{p}}/p$.

For the time derivatives, the following estimates hold. First, for $0 < \varepsilon < \varepsilon_0$, where ε_0 depends only on p and N ,

$$(4.4) \quad \partial_t v_\varepsilon(x, t) \leq 2^{\frac{p+1}{p}} N \lambda_p \|\varphi\|_\infty^{\frac{1}{p}} a^{-\frac{1}{p}} \varepsilon^{\frac{1}{2}} t^{-\frac{p+2}{2p}}$$

for $(x, t) \in \mathbb{R}^N \times (0, \infty)$.

Next, assume in addition that (3.1) is satisfied. Then, for all $0 < \varepsilon < \varepsilon_0$,

$$(4.5) \quad \partial_t v_\varepsilon(x, t) \geq -L t^{-\mu} - 2^{\frac{p+1}{p}} N \lambda_p \|\varphi\|_\infty^{\frac{1}{p}} a^{-\frac{1}{p}} \varepsilon^{\frac{1}{2}} t^{-\frac{p+2}{2p}}$$

for $(x, t) \in \mathbb{R}^N \times (0, \infty)$, the constants μ and L being defined in Theorem 3.1.

An estimate similar to (4.2) was obtained by Lions [19, Section I] but with a dependence upon ε which vanishes in the limit $\varepsilon \rightarrow 0$. The estimate (4.3) (for the special case (1.7)) was first derived in [7], and we follow it rather closely. Let us emphasize here that it is not only independent on ε but also on the initial data, a property we shall use in the proof of Theorem 3.3. The estimate (4.2) for the case $p < 1$ (again for the special form (1.7)) was first derived in [14]. However our proof seems to be simpler. The estimates derived here for the time derivatives generalize estimates obtained in [15] for the special case (1.7). In this latter case we have $\kappa_0 = \kappa_\infty = p$, hence $\mu = 1$.

Remark 4.4. *Remark that the estimates for the time derivative of v_ε are much more complicated (and depend explicitly on the behavior of H at $r = 0$ and $r = \infty$ as reflected in the additional assumption (3.1)). Such estimates are needed in order to ensure the convergence of the solution v_ε to (1.3)-(1.4) (as $\varepsilon \rightarrow 0$) in $W_{loc}^{1,\infty}(\mathbb{R}^N \times (0, \infty))$ to a viscosity solution of (1.1)-(1.2). In fact, in Appendix B (see also Remark 4.7) we show that the equicontinuity in t of the family $\{v_\varepsilon\}_{\varepsilon>0}$ can be obtained without the additional requirement (3.1). It follows, in view of the stability result for viscosity solutions (see, e.g., [5, Théorème 2.3] or [10, Theorem 1.4]) that the limit function is a viscosity solution to (1.1)-(1.2). However, without the uniform boundedness in $W_{loc}^{1,\infty}(\mathbb{R}^N \times (0, \infty))$ it is not possible to show that the limiting solution is differentiable a.e. and hence satisfies (1.1) a.e. (see [13, Section 10.1]).*

Proof of Proposition 4.3. Observe that we can assume without loss of generality that $v_\varepsilon > c$, where $c > 0$ is arbitrary, by adding c to φ , provided the estimates do not depend on c .

We take the regularized function Φ_η as in (1.5) and consider the solution v_ε^η to the modified equation

$$(4.6) \quad \partial_t v_\varepsilon^\eta - \varepsilon \Delta v_\varepsilon^\eta + \Phi_\eta(|\nabla v_\varepsilon^\eta|^2) = 0, \quad (x, t) \in \mathbb{R}^N \times (0, \infty),$$

$$(4.7) \quad v_\varepsilon^\eta(x, 0) = \varphi(x), \quad x \in \mathbb{R}^N$$

for $\varepsilon > 0$ and $\eta > 0$. Inspecting the proof in [1] we see that the solution v_ε^η exists globally, belongs to $C_b^2(\mathbb{R}^N)$ for all $t \geq 0$ and satisfies $v_\varepsilon^\eta > c$. In addition, it is smooth for $t > 0$ and we can differentiate it as many times as needed. Also, as $\eta \rightarrow 0$,

$$(v_\varepsilon^\eta, \nabla v_\varepsilon^\eta) \longrightarrow (v_\varepsilon, \nabla v_\varepsilon) \quad \text{uniformly in compact subsets of } \mathbb{R}^N \times (0, \infty).$$

It therefore suffices to prove the estimates in Proposition 4.3 for v_ε^η , provided these estimates are independent of the positive constants η and c . In what follows we simplify the notation by referring to v_ε^η as V .

We now consider a strictly monotone smooth function f and define the functions u and w by

$$(4.8) \quad u = f^{-1}(V) \quad \text{and} \quad w = |\nabla u|^2.$$

Clearly u and w both belong to $C^\infty(\mathbb{R}^N \times (0, \infty))$ and their first derivatives in x and t are uniformly bounded and continuous in $\mathbb{R}^N \times [0, \infty)$. From (4.6) we obtain that u solves

$$f'(u) \left\{ \partial_t u - \varepsilon \Delta u - \varepsilon \frac{f''(u)}{f'(u)} |\nabla u|^2 + \frac{1}{f'(u)} \Phi_\eta(f'(u)^2 |\nabla u|^2) \right\} = 0.$$

so that

$$(4.9) \quad \begin{aligned} & \partial_t w = 2 \nabla u \cdot \nabla \partial_t u \\ &= 2 \nabla u \cdot \left\{ \varepsilon \left[\Delta(\nabla u) + \nabla \left(\frac{f''(u)}{f'(u)} w \right) \right] - \nabla \left[\frac{1}{f'(u)} \Phi_\eta(f'(u)^2 w) \right] \right\} \\ &= 2\varepsilon \nabla u \cdot \left[\Delta(\nabla u) + \frac{f''(u)}{f'(u)} \nabla w \right] + 2\varepsilon \left(\frac{f''}{f'} \right)'(u) w^2 \\ &+ 2 \frac{f''(u)}{f'(u)^2} \Phi_\eta(f'(u)^2 w) w - 2 \Phi_\eta'(f'(u)^2 w) [f'(u) \nabla u \cdot \nabla w + 2f''(u) w^2]. \end{aligned}$$

Define the operator

$$\mathcal{L}z = z_t - \varepsilon \Delta z + 2\Phi'_\eta(f'(u)^2 |\nabla u|^2) f'(u) \nabla u \cdot \nabla z - 2\varepsilon \frac{f''(u)}{f'(u)} \nabla u \cdot \nabla z.$$

Noting that

$$(4.10) \quad \Delta w = 2 \sum_{j=1}^N \sum_{k=1}^N |\partial_{x_j} \partial_{x_k} u|^2 + 2 \nabla u \cdot \nabla (\Delta u) \geq 2 \nabla u \cdot \nabla (\Delta u),$$

we deduce from (4.9) that

$$(4.11) \quad \begin{aligned} \mathcal{L}w &\leq 2\varepsilon \left(\frac{f''}{f'} \right)'(u) w^2 + 2 \frac{f''}{(f')^2}(u) \Phi_\eta(f'(u)^2 w) w \\ &\quad - 4\Phi'_\eta(f'(u)^2 w) f''(u) w^2 \\ &= 2\varepsilon \left(\frac{f''}{f'} \right)'(u) w^2 - 2 \frac{f''}{(f')^2}(u) \Theta_\eta(f'(u)^2 w) w, \end{aligned}$$

where Θ_η is defined in Definition 1.1.

We now specify the function f and begin with the case $p > 1$. We choose

$$(4.12) \quad f(r) = r^{\frac{p}{p-1}}, \quad r \geq 0,$$

so that

$$-2 \left(\frac{f''}{f'} \right)'(r) = \frac{2}{p-1} r^{-2} \geq 0 \quad \text{and} \quad \frac{f''}{(f')^2}(r) = \frac{1}{p} r^{-\frac{p}{p-1}}.$$

Inserting these estimates in (4.11) we get

$$\mathcal{L}w + \frac{2}{p} u^{-\frac{p}{p-1}} \Theta_\eta(f'(u)^2 w) w \leq 0.$$

Owing to Definition 1.1 (i) we further obtain

$$\mathcal{L}w + \frac{2}{p} u^{-\frac{p}{p-1}} \left[a f'(u)^p w^{\frac{p}{2}} - b \eta^\gamma \right] w \leq 0,$$

and finally, inserting $f'(r) = p r^{\frac{1}{p-1}} / (p-1)$, we obtain

$$(4.13) \quad \mathcal{L}w + \frac{2}{p-1} \left(\frac{p}{p-1} \right)^{p-1} a w^{1+\frac{p}{2}} \leq \frac{2b}{p} \eta^\gamma u^{-\frac{p}{p-1}} w.$$

Note that $u^{-\frac{p}{p-1}} = V^{-1} \leq c^{-1}$, so that (4.13) yields, if we take $2b\eta^{\frac{\gamma}{2}} < pc$,

$$(4.14) \quad \mathcal{L}w + \frac{2a}{p-1} \left(\frac{p}{p-1} \right)^{p-1} w^{1+\frac{p}{2}} \leq \eta^{\frac{\gamma}{2}} w.$$

Now consider the function $h_\eta(t) = K_\eta t^{-\frac{2}{p}}$, where $K_\eta > 0$ is a constant to be determined. We require h_η to be a supersolution (in some time interval) to (4.14), namely,

$$\mathcal{L}h_\eta + \frac{2a}{p-1} \left(\frac{p}{p-1} \right)^{p-1} h_\eta^{1+\frac{p}{2}} \geq \eta^{\frac{\gamma}{2}} h_\eta.$$

This condition is satisfied in $0 < t < \eta^{-\frac{\gamma}{4}}$, if

$$-\frac{2}{p} K_\eta + \frac{2a}{p-1} \left(\frac{p}{p-1} \right)^{p-1} K_\eta^{1+\frac{p}{2}} = \eta^{\frac{\gamma}{4}} K_\eta,$$

or

$$(4.15) \quad K_\eta = \left(\frac{1}{a} + \frac{p}{2a} \eta^{\frac{\gamma}{4}} \right)^{\frac{2}{p}} \left(\frac{p-1}{p} \right)^2.$$

The comparison principle now implies that $w(x, t) \leq h_\eta(t)$ for $x \in \mathbb{R}^N$ and $0 < t < \eta^{-\frac{\gamma}{4}}$. Consequently,

$$(4.16) \quad \left\| \nabla \left((v_\varepsilon^\eta)^{\frac{p-1}{p}} \right) (\cdot, t) \right\|_\infty \leq K_\eta^{\frac{1}{2}} t^{-\frac{1}{p}}, \quad 0 < t < \eta^{-\frac{\gamma}{4}}, \quad 0 < \frac{2b}{p} \eta^{\frac{\gamma}{2}} < c.$$

In addition, combining (4.16) with the bound $v_\varepsilon^\eta \leq \|\varphi\|_\infty + c$ gives

$$(4.17) \quad \|\nabla v_\varepsilon^\eta(\cdot, t)\|_\infty \leq \frac{pK_\eta^{\frac{1}{2}}}{p-1} (\|\varphi\|_\infty + c)^{\frac{1}{p}} t^{-\frac{1}{p}}, \quad 0 < t < \eta^{-\frac{\gamma}{4}}, \quad 0 < \frac{2b}{p} \eta^{\frac{\gamma}{2}} < c.$$

These estimates are independent of $\varepsilon > 0$, and by letting $\eta \rightarrow 0$, and then $c \rightarrow 0$, we obtain (4.3) and (4.2) in the case $p > 1$.

We now turn to the case $0 < p < 1$. Our starting point is again the inequality (4.11) with the same function Φ_η but with a different choice of the function f . More precisely, instead of (4.12), we take

$$(4.18) \quad f(r) = 2\eta^{\frac{\gamma}{2}} + \|\varphi\|_\infty - \frac{1}{2}r^2.$$

Then

$$-2 \left(\frac{f''}{f'} \right)' (r) = \frac{2}{r^2} \geq 0 \quad \text{and} \quad \frac{f''}{(f')^2} (r) = -\frac{1}{r^2}.$$

Inserting these estimates in (4.11) we get

$$\mathcal{L}w - \frac{2}{u^2} \Theta_\eta(f'(u)^2 w) w \leq 0,$$

so that in conjunction with Definition 1.1 (ii) we obtain

$$\mathcal{L}w + 2au^{p-2}w^{1+\frac{p}{2}} \leq \frac{2b\eta^\gamma}{u^2} w.$$

Taking $0 < c < \eta^{\frac{\gamma}{2}}$ the maximum principle (for $V = v_\varepsilon^\eta$) implies that

$$(4.19) \quad \eta^{\frac{\gamma}{2}} \leq \frac{1}{2}u^2 \leq 2\eta^{\frac{\gamma}{2}} + \|\varphi\|_\infty.$$

This estimate provides an upper bound for the right-hand side of the above inequality and leads us to

$$(4.20) \quad \mathcal{L}w + 2au^{p-2}w^{1+\frac{p}{2}} \leq b\eta^{\frac{\gamma}{2}} w.$$

Also by (4.19), since $p-2 < 0$,

$$u^{p-2} \geq \left[2(2\eta^{\frac{\gamma}{2}} + \|\varphi\|_\infty) \right]^{\frac{p-2}{2}},$$

so that (4.20) yields

$$(4.21) \quad \mathcal{L}w + 2^{\frac{p}{2}} a (2\eta^{\frac{\gamma}{2}} + \|\varphi\|_\infty)^{\frac{p-2}{2}} w^{1+\frac{p}{2}} \leq b\eta^{\frac{\gamma}{2}} w.$$

As above, we now try a supersolution to (4.21) of the form $h_\eta(t) = K_\eta t^{-\frac{2}{p}}$ (in a certain time interval). We therefore need

$$-\frac{2}{p} + 2^{\frac{p}{2}} a (2\eta^{\frac{\gamma}{2}} + \|\varphi\|_\infty)^{\frac{p-2}{2}} K_\eta^{\frac{p}{2}} \geq b\eta^{\frac{\gamma}{2}} t,$$

hence for $0 < t < \eta^{-\frac{7}{4}}$ we can take

$$K_\eta = \left(\frac{2 + bp\eta^{\frac{7}{4}}}{2^{\frac{p}{2}} ap} \right)^{\frac{2}{p}} \left(2\eta^{\frac{7}{2}} + \|\varphi\|_\infty \right)^{\frac{2-p}{p}}.$$

The comparison principle then entails that $w(x, t) \leq K_\eta t^{-\frac{2}{p}}$ for $x \in \mathbb{R}^N$ and $0 < t < \eta^{-\frac{7}{4}}$. Using (4.18) and (4.19) we conclude that

$$\|\nabla v_\varepsilon^\eta(\cdot, t)\|_\infty \leq \|u(\cdot, t)\|_\infty \|\nabla u(\cdot, t)\|_\infty \leq \left(\frac{2 + bp\eta^{\frac{7}{4}}}{pa} \right)^{\frac{1}{p}} (2\eta^{\frac{p\gamma}{2}} + \|\varphi\|_\infty)^{\frac{1}{p}} t^{-\frac{1}{p}}$$

for $0 < t < \eta^{-\frac{7}{4}}$ and $0 < c < \eta^{\frac{7}{2}}$. This estimate is independent of $\varepsilon > 0$. Letting $c \rightarrow 0$ and then $\eta \rightarrow 0$ we obtain (4.2) for $0 < p < 1$ with $\lambda_p = (2/p)^{\frac{1}{p}}$.

We next turn to the proof of (4.4) and (4.5). We still work with the modified equation (4.6) and simplify as before the notation by setting $v_\varepsilon^\eta = V$. We follow the idea of proof in [15, Lemma 10].

Let $M > 0$ and $\vartheta > 0$ be positive constants (to be specified later) and define $\Gamma = M + N (\|\nabla\varphi\|_\infty^2 - |\nabla V|^2) / (4M)$ and $w = (\delta \partial_t V - \vartheta) / \Gamma$ for $\delta \in \{-1, 1\}$. From (4.6) we get readily

$$\partial_t w = -\frac{w}{\Gamma} \partial_t \Gamma + \frac{\varepsilon \Delta(\Gamma w)}{\Gamma} - \frac{2}{\Gamma} \Phi'_\eta(|\nabla V|^2) \nabla V \cdot \nabla(\Gamma w)$$

which we can rewrite as

$$(4.22) \quad \partial_t w = \frac{N}{4M} \frac{w}{\Gamma} A + B \cdot \nabla w + \varepsilon \Delta w,$$

where

$$B = 2\varepsilon \frac{\nabla \Gamma}{\Gamma} - 2\Phi'_\eta(|\nabla V|^2) \nabla V$$

is a bounded continuous function and

$$A = \partial_t (|\nabla V|^2) - \varepsilon \Delta (|\nabla V|^2) + 2\Phi'_\eta(|\nabla V|^2) \nabla V \cdot \nabla (|\nabla V|^2).$$

Recalling that (cf. (4.10))

$$\Delta (|\nabla V|^2) = 2\nabla V \cdot \nabla(\Delta V) + 2 \sum_{j=1}^N \sum_{k=1}^N (\partial_{x_j} \partial_{x_k} V)^2$$

and

$$\partial_t (|\nabla V|^2) - 2\varepsilon \nabla V \cdot \nabla(\Delta V) + 2\Phi'_\eta(|\nabla V|^2) \nabla V \cdot \nabla (|\nabla V|^2) = 0$$

by (4.9) (with $f(r) = r$ so that $u = V$), we obtain

$$A = 2\varepsilon \nabla V \cdot \nabla(\Delta V) - \varepsilon \Delta (|\nabla V|^2) = -2\varepsilon \sum_{j=1}^N \sum_{k=1}^N (\partial_{x_j} \partial_{x_k} V)^2.$$

Then (4.22) reads

$$\partial_t w - \varepsilon \Delta w - B \cdot \nabla w + \frac{N\varepsilon}{2M} \frac{w}{\Gamma} \left(\sum_{j=1}^N \sum_{k=1}^N (\partial_{x_j} \partial_{x_k} V)^2 \right) = 0$$

or

$$\partial_t w - \varepsilon \Delta w - B \cdot \nabla w + \frac{\varepsilon}{2M} A_1 \frac{w}{\Gamma} + \frac{\varepsilon}{2M} \frac{w}{\Gamma} |\Delta V|^2 = 0$$

with

$$A_1 = N \sum_{j=1}^N \sum_{k=1}^N (\partial_{x_j} \partial_{x_k} V)^2 - |\Delta V|^2 \geq 0$$

by the Cauchy-Schwarz inequality. Noting that

$$|\Delta V|^2 = \frac{1}{\varepsilon^2} (w\Gamma + \vartheta + \delta \Phi_\eta (|\nabla V|^2))^2$$

by (4.6) and introducing the differential operator

$$\begin{aligned} \mathcal{M}z &= \partial_t z - \varepsilon \Delta z - B \cdot \nabla z + \frac{\Gamma}{2M\varepsilon} z^3 + \frac{1}{M\varepsilon} (\vartheta + \delta \Phi_\eta (|\nabla V|^2)) z^2 \\ &+ \frac{\varepsilon}{2M\Gamma} \left\{ A_1 + \frac{(\vartheta + \delta \Phi_\eta (|\nabla V|^2))^2}{\varepsilon^2} \right\} z, \end{aligned}$$

we realize that w solves

$$(4.23) \quad \mathcal{M}w = 0 \quad \text{in} \quad \mathbb{R}^N \times (0, \infty).$$

We first take $\delta = -1$ and

$$\vartheta = \sup_{r \in [0, \|\nabla\varphi\|_\infty]} \{\Phi_\eta(r^2)\}$$

in the definition of w . As $\Gamma \geq M$, we infer from the nonnegativity of A_1 , ϑ , Φ_η and Proposition 4.2 (3) that $W(t) = (\varepsilon/t)^{\frac{1}{2}}$ satisfies

$$\mathcal{M}W \geq -\frac{\sqrt{\varepsilon}}{2} t^{-\frac{3}{2}} + \frac{\Gamma}{2M\varepsilon} W^3(t) \geq 0.$$

Therefore W is a supersolution to (4.23) and the comparison principle entails that $w \leq W$ in $\mathbb{R}^N \times (0, \infty)$. Consequently,

$$-\partial_t V(x, t) - \sup_{r \in [0, \|\nabla\varphi\|_\infty]} \{\Phi_\eta(r^2)\} \leq \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}} \Gamma(x, t) \leq \left(M + \frac{N}{4M} \|\nabla\varphi\|_\infty^2\right) \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}}.$$

Choosing $M = \|\nabla\varphi\|_\infty$, we end up with

$$(4.24) \quad \partial_t V(x, t) \geq - \sup_{r \in [0, \|\nabla\varphi\|_\infty]} \{\Phi_\eta(r^2)\} - \frac{N+4}{4} \|\nabla\varphi\|_\infty \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}}.$$

We next take $\delta = 1$ and $\vartheta = 0$ in the definition of w . As above, it follows from the nonnegativity of A_1 and Φ_η and the bound $\Gamma \geq M$ that the function W satisfies $\mathcal{M}W \geq 0$ in $\mathbb{R}^N \times (0, \infty)$, whence $w \leq W$ by the comparison principle. Therefore,

$$\partial_t V(x, t) \leq \left(M + \frac{N}{4M} \|\nabla\varphi\|_\infty^2\right) \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}},$$

and the choice $M = \|\nabla\varphi\|_\infty$ gives

$$(4.25) \quad \partial_t V(x, t) \leq \frac{N+4}{4} \|\nabla\varphi\|_\infty \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}}.$$

We then pass to the limit as $\eta \rightarrow 0$ and infer from (1.6) and the convergence of $(v_\varepsilon^\eta)_\eta$ towards v_ε that

$$- \sup_{r \in [0, \|\nabla\varphi\|_\infty]} \{H(r)\} - \frac{N+4}{4} \|\nabla\varphi\|_\infty \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}} \leq \partial_t v_\varepsilon(x, t) \leq \frac{N+4}{4} \|\nabla\varphi\|_\infty \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}}$$

for $(x, t) \in \mathbb{R}^N \times (0, \infty)$. We finally use (3.1) to conclude that

$$\begin{aligned} \partial_t v_\varepsilon(x, t) &\geq -g_H (\|\nabla\varphi\|_\infty^{\kappa_\infty} + \|\nabla\varphi\|_\infty^{\kappa_0}) - \frac{N+4}{4} \|\nabla\varphi\|_\infty \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}} \\ \partial_t v_\varepsilon(x, t) &\leq \frac{N+4}{4} \|\nabla\varphi\|_\infty \left(\frac{\varepsilon}{t}\right)^{\frac{1}{2}} \end{aligned}$$

for $(x, t) \in \mathbb{R}^N \times (0, \infty)$. But, since (4.6) is an autonomous equation, we also have

$$\begin{aligned} \partial_t v_\varepsilon(x, t) &\geq -g_H \left(\left\| \nabla v_\varepsilon \left(\frac{t}{2} \right) \right\|_\infty^{\kappa_\infty} + \left\| \nabla v_\varepsilon \left(\frac{t}{2} \right) \right\|_\infty^{\kappa_0} \right) \\ &\quad - \frac{N+4}{4} \left\| \nabla v_\varepsilon \left(\frac{t}{2} \right) \right\|_\infty \left(\frac{2\varepsilon}{t} \right)^{\frac{1}{2}}, \\ \partial_t v_\varepsilon(x, t) &\leq \frac{N+4}{4} \left\| \nabla v_\varepsilon \left(\frac{t}{2} \right) \right\|_\infty \left(\frac{2\varepsilon}{t} \right)^{\frac{1}{2}} \end{aligned}$$

for $(x, t) \in \mathbb{R}^N \times (0, \infty)$. Inserting (4.2) in the above estimates and using that $N+4 \leq 4\sqrt{2}N$ complete the proof of (4.4) and (4.5). \square

As already mentioned, in the particular case where H is given by (1.7), we have $\kappa_0 = \kappa_\infty = p$, and thus $\mu = 1$. We can then derive a better estimate for the time derivative, using a scaling argument as follows.

Corollary 4.5. *Let H be of the special form (1.7). Then for every $\rho > 0$ there exists a constant $C > 0$, depending only on $p, N, \|\varphi\|_\infty, \rho$ such that, for all $0 < \varepsilon < \varepsilon_0$,*

$$(4.26) \quad |\partial_t v_\varepsilon(x, t)| \leq Ct^{-1}, \quad \begin{cases} (x, t) \in \mathbb{R}^N \times (\rho, \infty), & p \in (0, 1) \cup (1, 2], \\ (x, t) \in \mathbb{R}^N \times (0, \rho), & p \in [2, \infty). \end{cases}$$

Proof. Note that v_ε satisfies, in view of (4.4)-(4.5), the estimate

$$\|\partial_t v_\varepsilon(\cdot, \rho)\|_\infty \leq c_0, \quad 0 < \varepsilon < \varepsilon_0,$$

where c_0 is independent of ε . Define the function

$$V(y, \tau) = v_\varepsilon(r^\beta y, r^\alpha \tau), \quad (y, \tau) \in \mathbb{R}^N \times (0, \infty), \quad r > 0.$$

It satisfies the equation (using the special form of H)

$$\partial_\tau V(y, \tau) + r^{\alpha-p\beta} |\nabla V(y, \tau)|^p = \varepsilon r^{\alpha-2\beta} \Delta V(y, \tau).$$

Assume first that $0 < p \leq 2$ and take $\alpha = 1$, $\beta = p^{-1}$ and $r > 1$. Then $\varepsilon r^{\alpha-2\beta} < \varepsilon < \varepsilon_0$ hence

$$\|\partial_\tau V(\cdot, \rho)\|_\infty \leq c_0,$$

and turning back to v_ε with $t = r\rho$ we obtain (4.26) for $0 < p \leq 2$. In the case $p > 2$ we repeat the same argument, but with $r < 1$. \square

In view of the fact that only $\|\varphi\|_\infty$ appears in the estimates, we can follow the methodology of [14] and extend the result of Proposition 4.2 as follows.

Corollary 4.6. *Let $0 \leq \varphi \in C_b(\mathbb{R}^N)$, and let H satisfy the hypotheses (1.5) and (1.6). Then (1.3)-(1.4) has a unique global solution v_ε such that*

- (i) $v_\varepsilon \in C^{2,1}(\mathbb{R}^N \times (0, \infty)) \cap C(\mathbb{R}^N \times [0, \infty))$,
- (ii) $0 \leq v_\varepsilon(x, t) \leq \|\varphi\|_\infty, \quad (x, t) \in \mathbb{R}^N \times (0, \infty)$,

(iii) v_ε satisfies in $\mathbb{R}^N \times (0, \infty)$ all the estimates of Proposition 4.3, the estimate (4.5) being only true if H fulfills the additional assumption (3.1).

In addition, if $\varphi \in W^{1,\infty}(\mathbb{R}^N)$, then

$$\|\nabla_x v_\varepsilon(\cdot, t)\|_\infty \leq \|\nabla_x \varphi\|_\infty, \quad t > 0.$$

Remark 4.7. In contrast to the rather involved proof of (4.4)-(4.5), it is quite easy to show that v_ε belongs to $C([0, \infty), L^\infty(\mathbb{R}^N))$ (and is in fact Hölder continuous with respect to time, see Proposition B.1 in Appendix B below). Such an estimate is sufficient for proving the uniform convergence (in compact subsets) of a subsequence $\{v_{\varepsilon_j}\}_{j \geq 1}$ (as $\varepsilon_j \rightarrow 0$), when the estimate (4.2) is known, using the Arzela-Ascoli theorem. It follows, in view of the stability result for viscosity solutions [5, Théorème 2.3] that the limit function is a viscosity solution to (1.1)-(1.2).

5. Proof of Theorem 3.1

Consider $0 \leq \varphi \in C_b(\mathbb{R}^N)$ and let v_ε be the solution to (1.3)-(1.4) given in Corollary 4.6. In view of Proposition 4.2 (2), (4.2) and (4.4)-(4.5), the family $\{v_\varepsilon\}_{\varepsilon \in (0, \varepsilon_0)}$ is uniformly bounded in $\mathbb{R}^N \times (0, \infty)$ and also bounded in $W_{loc}^{1,\infty}(\mathbb{R}^N \times (0, \infty))$. It follows that there exist a subsequence $\{v_{\varepsilon_j}\}$, $\varepsilon_j \rightarrow 0$ and a function $v \in C_b(\mathbb{R}^N \times (0, \infty)) \cap W_{loc}^{1,\infty}(\mathbb{R}^N \times (0, \infty))$ such that

$$(5.1) \quad v_{\varepsilon_j} \xrightarrow{\varepsilon_j \rightarrow 0} v, \quad \text{uniformly in every compact subset of } \mathbb{R}^N \times (0, \infty).$$

The differentiability (a.e. in $\mathbb{R}^N \times (0, \infty)$) and the inequalities (3.2), (3.3) now follow from Rademacher's theorem [13, Chapter 5] and Proposition 4.3.

The limit function v satisfies (1.1) a.e. in $\mathbb{R}^N \times (0, \infty)$. Indeed, the convergence (5.1) implies, as in [13, Chapter 10], that v is a "viscosity solution" to (1.1) and therefore it satisfies (1.1) at any point where it is differentiable.

Next, we need to show that v attains the assigned initial condition (1.2). In view of (1.1) and the nonnegativity of H we have $\partial_t v \leq 0$ a.e. in $\mathbb{R}^N \times (0, \infty)$, whence $v(x, t_1) \leq v(x, t_2)$ for $x \in \mathbb{R}^N$ and $t_2 > t_1 > 0$ owing to the continuity of v in $\mathbb{R}^N \times (0, \infty)$. Recalling that $0 \leq v_\varepsilon \leq \|\varphi\|_\infty$ by Proposition 4.2 (2), the function

$$(5.2) \quad v_0(x) = \sup_{t > 0} v(x, t) = \lim_{t \rightarrow 0} v(x, t) \in [0, \|\varphi\|_\infty]$$

is thus well-defined and satisfies

$$0 \leq v_0(x) \leq \|\varphi\|_\infty \quad \text{for } x \in \mathbb{R}^N.$$

We now identify v_0 . Assume first that $\varphi \in C_0^\infty(\mathbb{R}^N)$ (it actually suffices to assume $\varphi \in W^{1,\infty}(\mathbb{R}^N)$) and consider $t > 0$. Then, multiplying (1.3) by any $\psi \in C_0^\infty(\mathbb{R}^N)$ and integrating over $\mathbb{R}^N \times (0, t)$, we get

$$\begin{aligned} \left| \int_{\mathbb{R}^N} v_\varepsilon(x, t) \psi(x) dx - \int_{\mathbb{R}^N} \varphi(x) \psi(x) dx \right| &\leq \varepsilon t \|\varphi\|_\infty \|\Delta \psi(x)\|_1 \\ &\quad + t \|\psi\|_1 \max_{0 \leq s \leq \|\nabla \varphi\|_\infty} H(s), \end{aligned}$$

where we have used the estimates in Proposition 4.2. Letting $\varepsilon = \varepsilon_j$ and $j \rightarrow \infty$ we obtain

$$\left| \int_{\mathbb{R}^N} v(x, t) \psi(x) dx - \int_{\mathbb{R}^N} \varphi(x) \psi(x) dx \right| \leq t \|\psi\|_1 \max_{0 \leq s \leq \|\nabla \varphi\|_\infty} H(s)$$

which yields, by taking $t \downarrow 0$

$$(5.3) \quad v_0(x) = \varphi(x), \text{ if } \varphi \in C_0^\infty(\mathbb{R}^N).$$

Coming back to the general case $\varphi \in C_b(\mathbb{R}^N)$ we consider $0 \leq \psi \in C_0^\infty(\mathbb{R}^N)$ and $t > 0$. Multiplying (1.3) by ψ , integrating over $\mathbb{R}^N \times (0, t)$ and using the positivity of H we get

$$\int_{\mathbb{R}^N} v_\varepsilon(x, t) \psi(x) dx \leq \int_{\mathbb{R}^N} \varphi(x) \psi(x) dx + \varepsilon \int_0^t \int_{\mathbb{R}^N} v_\varepsilon(x, s) \Delta \psi(x) dx ds,$$

which yields, by taking the sequence $\varepsilon = \varepsilon_j$ and letting $j \rightarrow \infty$,

$$\int_{\mathbb{R}^N} v(x, t) \psi(x) dx \leq \int_{\mathbb{R}^N} \varphi(x) \psi(x) dx.$$

It follows that, by taking the limit $t \downarrow 0$

$$(5.4) \quad v_0(x) \leq \varphi(x), \quad \text{for a.e. } x \in \mathbb{R}^N.$$

To prove the opposite inequality, we first observe that, if $\varphi(x_0) = 0$ for some $x_0 \in \mathbb{R}^N$, then $v_0(x_0) = 0$ by (5.4). Next, let $x_0 \in \mathbb{R}^N$ be such that $\varphi(x_0) > 0$. For $\eta > 0$ sufficiently small let

$$B^\eta = B_{\delta(\eta)}(x_0) = \{x \in \mathbb{R}^N, \quad |x - x_0| < \delta(\eta)\}$$

be a ball such that

$$\varphi(x) \geq (1 - \eta)\varphi(x_0), \quad x \in B^\eta.$$

Consider now $0 \leq \psi_\eta \in C_0^\infty(B^\eta)$ such that $\psi_\eta(x) \leq (1 - \eta)\varphi(x_0)$ with equality at $x = x_0$. Let Ψ_η denote the solution to (1.1) (constructed as above), with initial condition ψ_η . By the comparison principle for viscosity solutions we have $\Psi_\eta(x, t) \leq v(x, t)$ for $(x, t) \in \mathbb{R}^N \times (0, \infty)$. However, in view of (5.3) it follows that

$$\Psi_\eta(x, t) \xrightarrow[t \rightarrow 0^+]{} \psi_\eta(x), \quad x \in \mathbb{R}^N$$

so that by the previous inequality $\psi_\eta(x) \leq v_0(x)$ for a.e. $x \in \mathbb{R}^N$. If now x_0 is a Lebesgue point of v_0 , the last inequality implies that $(1 - \eta)\varphi(x_0) = \psi_\eta(x_0) \leq v_0(x_0)$, and by sending η to 0 we get for such a point $\varphi(x_0) \leq v_0(x_0)$. Thus, finally

$$\varphi(x) \leq v_0(x) \quad \text{for a.e. } x \in \mathbb{R}^N.$$

Combining this inequality with (5.4) we get $\varphi(x) = v_0(x)$ for $x \in \mathbb{R}^N$. Finally, as $\varphi \in C(\mathbb{R}^N)$, the time monotonicity of v and the Dini theorem warrant that

$$v(x, t) \xrightarrow[t \rightarrow 0^+]{} \varphi(x) \quad \text{uniformly in compact subsets of } \mathbb{R}^N.$$

The uniqueness of the solution follows from the fact that Equation (1.1) satisfies the comparison principle in $C_b(\mathbb{R}^N)$.

6. Proof of Theorem 3.3

We begin by noting that since φ is only assumed to be continuous (but not necessarily bounded), we cannot invoke Corollary 4.6. The existence of a solution v_ε to (1.3)-(1.4) is therefore not guaranteed and must be addressed as a first step towards the study of a ‘‘vanishing viscosity solution’’.

For any integer $n \geq 1$ we set

$$\varphi_n = \min\{\varphi, n\} \in C_b(\mathbb{R}^N),$$

and let $v_{\varepsilon,n}$ be the solution to (1.3) subject to the initial condition $v_{\varepsilon,n}(x,0) = \varphi_n(x)$. In view of Corollary 4.6

$$\begin{aligned} v_{\varepsilon,n} &\in C^{2,1}(\mathbb{R}^N \times (0, \infty)) \cap C(\mathbb{R}^N \times [0, \infty)), \\ 0 &\leq v_{\varepsilon,n}(x, t) \leq \|\varphi_n\|_\infty, \quad (x, t) \in \mathbb{R}^N \times (0, \infty), \end{aligned}$$

and the estimate (4.3) is satisfied. Moreover, by Theorem 3.1 we have for any fixed n ,

$$v_{\varepsilon,n} \xrightarrow{\varepsilon \rightarrow 0^+} v_{0,n}, \quad \text{uniformly in every compact subset of } \mathbb{R}^N \times (0, \infty),$$

where the limit function $v_{0,n}$ is differentiable a.e. in $\mathbb{R}^N \times (0, \infty)$ with $v_{0,n}(\cdot, 0) = \varphi_n$ and satisfies (1.1) at any point of differentiability.

Next we show that the family $\{v_{\varepsilon,n}\}_{n \geq 1}$ is uniformly bounded in every compact subset of $\mathbb{R}^N \times (0, \infty)$. To this end we follow [22] and state the following lemma.

Lemma 6.1. *Let $z \in C^{2,1}(\mathbb{R}^N \times (0, \infty))$ be any classical solution to (1.3) for some $\varepsilon \in (0, 1)$. Assume that H satisfies the assumptions (1.5) and (1.6) with $p > 1$. Then, for any $t > s > 0$ and all $y \in \mathbb{R}^N$ and $R > 0$,*

$$(6.1) \quad \int_{B_R(y)} z(x, t) dx \leq \int_{B_{2R}(y)} z(x, s) dx + C(t-s)R^N(1 + R^{-\frac{p}{p-1}}),$$

where $B_r(y) = \{x \in \mathbb{R}^N, |x - y| < r\}$, and $C > 0$ depends only on N, p and a (but, in particular, not on ε).

The proof of the lemma is given at the end of this section.

We now continue with the proof of Theorem 3.3. In what follows we use $C > 0$ to denote various constants depending only on p, N and a unless explicit dependence on other parameters is indicated.

Let $t > 0$. In view of (4.3) we have, for any $x, y \in \mathbb{R}^N$,

$$v_{\varepsilon,n}(y, t)^{\frac{p-1}{p}} \leq v_{\varepsilon,n}(x, t)^{\frac{p-1}{p}} + \mu_p t^{-\frac{1}{p}} |x - y|,$$

hence, since $p > 1$,

$$v_{\varepsilon,n}(y, t) \leq C \left[v_{\varepsilon,n}(x, t) + t^{-\frac{1}{p-1}} |x - y|^{\frac{p}{p-1}} \right].$$

Integrating this inequality over $B_R(y)$ with respect to x we get

$$\int_{B_R(y)} v_{\varepsilon,n}(y, t) dx \leq C \left[\int_{B_R(y)} v_{\varepsilon,n}(x, t) dx + t^{-\frac{1}{p-1}} \int_{B_R(0)} |x|^{\frac{p}{p-1}} dx \right].$$

We now invoke the estimate (6.1) for $z = v_{\varepsilon,n}$ and $s = 0$ (which is possible since $v_{\varepsilon,n}$ is continuous at $s = 0$) in the right-hand side of the last inequality.

$$\begin{aligned} v_{\varepsilon,n}(y, t) &\leq CR^{-N} \left[\int_{B_{2R}(y)} \varphi_n(x) dx + tR^N \left(1 + R^{-\frac{p}{p-1}} \right) + t^{-\frac{1}{p-1}} R^{N+\frac{p}{p-1}} \right] \\ &\leq C \left[R^{-N} \int_{B_{2R}(y)} \varphi(x) dx + t \left(1 + R^{-\frac{p}{p-1}} \right) + t^{-\frac{1}{p-1}} R^{\frac{p}{p-1}} \right]. \end{aligned}$$

It follows that in any cylinder $Q = B_R(0) \times (t_1, t_2)$, $t_2 > t_1 > 0$,

$$\|v_{\varepsilon,n}\|_{L^\infty(Q)} \leq C(R, t_1, t_2, \varphi), \quad n \geq 1,$$

and, in view of (4.3),

$$\|\nabla_x v_{\varepsilon,n}\|_{L^\infty(Q)} \leq \frac{p\mu_p}{p-1} \|v_{\varepsilon,n}\|_{L^\infty(Q)}^{\frac{1}{p}} t_1^{-\frac{1}{p}} \leq C(R, t_1, t_2, \varphi), \quad n \geq 1.$$

In view of Theorem 3.1 we have, by passing to the limit $\varepsilon \rightarrow 0$

$$(6.2) \quad \partial_t v_{0,n} + H(|\nabla_x v_{0,n}|) = 0, \quad \text{for a.e. } (x, t) \in \mathbb{R}^N \times (0, \infty)$$

The last two estimates and (6.2) yield

$$(6.3) \quad \|v_{0,n}\|_{W^{1,\infty}(Q)} \leq C(R, t_1, t_2, \varphi), \quad n \geq 1.$$

Using a diagonal process, we obtain a subsequence $(n_j)_{j \geq 1}$ such that

$$v_{0,n_j} \xrightarrow{j \rightarrow \infty} v, \quad \text{uniformly in every compact subset of } \mathbb{R}^N \times (0, \infty),$$

where the limit function $v \in W_{loc}^{1,\infty}(\mathbb{R}^N \times (0, \infty))$ satisfies (6.3), hence is differentiable a.e. in $\mathbb{R}^N \times (0, \infty)$.

We now use the stability result for viscosity solutions [5, Théorème 2.3] in order to obtain the fact that v is indeed a viscosity solution to (1.1), satisfying (1.1) a.e. in $\mathbb{R}^N \times (0, \infty)$.

Finally, the proof that

$$v(x, t) \xrightarrow{t \rightarrow 0^+} \varphi(x), \quad \text{uniformly in every compact subset of } \mathbb{R}^N$$

follows essentially the same reasoning as the corresponding proof in the case of Theorem 3.1. As in the case of Theorem 3.1, the uniqueness assertion follows from the fact that Equation (1.1) satisfies the (discontinuous) comparison principle. This concludes the proof of the theorem.

Proof of Lemma 6.1. Let $\xi \in C_0^\infty(\mathbb{R}^N)$ such that $0 \leq \xi \leq 1$ and k an integer such that $k > p/(p-1)$. Multiplying (1.3) by ξ^k and integrating over \mathbb{R}^N we get

$$\begin{aligned} \frac{d}{dt} \int_{\mathbb{R}^N} \xi(x)^k z(x, t) dx &+ \int_{\mathbb{R}^N} \xi(x)^k H(|\nabla z(x, t)|) dx \\ &= -\varepsilon \int_{\mathbb{R}^N} \nabla(\xi(x)^k) \cdot \nabla z(x, t) dx. \end{aligned}$$

We now use Young's inequality to estimate the right-hand side,

$$\begin{aligned} \left| \int_{\mathbb{R}^N} \nabla(\xi(x)^k) \cdot \nabla z(x, t) dx \right| &\leq \frac{a}{p(p-1)} \int_{\mathbb{R}^N} (\xi(x))^k |\nabla z(x, t)|^p dx \\ &+ \frac{p-1}{p} \left(\frac{a}{p-1} \right)^{-\frac{1}{p-1}} \int_{\mathbb{R}^N} \xi(x)^{-\frac{k}{p-1}} |\nabla(\xi(x)^k)|^{\frac{p}{p-1}} dx, \end{aligned}$$

so that by (4.1) we get

$$\frac{d}{dt} \int_{\mathbb{R}^N} \xi^k(x) z(x, t) dx \leq \frac{p-1}{p} \left(\frac{a}{p-1} \right)^{-\frac{1}{p-1}} \int_{\mathbb{R}^N} \xi(x)^{-\frac{k}{p-1}} |\nabla(\xi(x)^k)|^{\frac{p}{p-1}} dx.$$

By the choice of k and ξ the first integral in the right-hand side is

$$\begin{aligned} \int_{\mathbb{R}^N} \xi(x)^{-\frac{k}{p-1}} |\nabla(\xi(x)^k)|^{\frac{p}{p-1}} dx &= k^{\frac{p}{p-1}} \int_{\mathbb{R}^N} \xi(x)^{k-\frac{p}{p-1}} |\nabla \xi(x)|^{\frac{p}{p-1}} dx \\ &\leq k^{\frac{p}{p-1}} \int_{\mathbb{R}^N} |\nabla \xi(x)|^{\frac{p}{p-1}} dx, \end{aligned}$$

so that

$$\frac{d}{dt} \int_{\mathbb{R}^N} \xi^k(x) z(x, t) dx \leq \frac{p-1}{p} \left(\frac{a}{p-1} \right)^{-\frac{1}{p-1}} k^{\frac{p}{p-1}} \int_{\mathbb{R}^N} \left[\xi^k(x) + |\nabla \xi(x)|^{\frac{p}{p-1}} \right] dx.$$

Let $\psi \in C_0^\infty(\mathbb{R}^N)$ be such that $0 \leq \psi \leq 1$ and $\psi(x) = 1$ (resp. $\psi(x) = 0$) if $|x| \leq 1$ (resp. $|x| \geq 2$). Taking $\xi(x) = \psi((x-y)/R)$ in the last estimate and integrating from s to t we obtain (6.1). \square

APPENDIX A. THE CASE $H(r) = r^p$

In this Appendix we establish (1.6) for the special case $H(r) = r^p$, where $p \in (0, \infty)$, $p \neq 1$. Here we set

$$(A.1) \quad \Phi_\eta(r) = (r + \eta^2)^{\frac{p}{2}} - \eta^p, \quad r \in [0, \infty)$$

so that by Definition 1.1

$$\Theta_\eta(r) = (p-1)(r + \eta^2)^{\frac{p}{2}} - p\eta^2(r + \eta^2)^{\frac{p}{2}-1} + \eta^p,$$

Assume first that $1 < p \leq 2$. Using

$$(A.2) \quad \Theta_\eta(r) \geq (p-1)(r + \eta^2)^{\frac{p}{2}} - (p-1)\eta^p \geq (p-1)[r^{\frac{p}{2}} - \eta^p],$$

Definition 1.1 (i) is established with $a = b = p-1$ and $\gamma = p$.

Consider next the case $p > 2$. Instead of (A.2) we now use the Young inequality to obtain

$$(A.3) \quad \begin{aligned} \Theta_\eta(r) &= (p-1)(r + \eta^2)^{\frac{p}{2}} - p\eta^2(r + \eta^2)^{\frac{p}{2}-1} + \eta^p \\ &\geq (p-1)(r + \eta^2)^{\frac{p}{2}} - \eta(r + \eta^2)^{\frac{p}{2}} - 2(p-2)^{\frac{p-2}{2}} \eta^{\frac{p+2}{2}}. \end{aligned}$$

Thus, for $p > 2$, Definition 1.1 (i) is established with $a = (p-1)/2$, $b = 2(p-2)^{\frac{p-2}{2}}$ and $\gamma = (p+2)/2$.

The case of a sum of powers (1.8) follows immediately from the above argument by taking $p = \min\{p_1, \dots, p_m\}$. Finally we turn to the case $0 < p < 1$. We now have

$$(A.4) \quad \begin{aligned} \Theta_\eta(r) &= pr(r + \eta^2)^{\frac{p}{2}-1} - (r + \eta^2)^{\frac{p}{2}} + \eta^p \\ &= (r + \eta^2)^{\frac{p}{2}} \left[p-1 - \frac{p\eta^2}{r + \eta^2} \right] + \eta^p \\ &\leq (p-1)r^{\frac{p}{2}} + \eta^p, \end{aligned}$$

which completes the proof of Definition 1.1 (i) for $0 < p < 1$ with $a = 1-p$, $b = 1$ and $\gamma = p$. As in the case $p > 1$, this treatment generalizes readily to the case of a sum of powers (all less than 1) as in (1.8), with $p = \max\{p_1, \dots, p_m\}$.

APPENDIX B. TIME EQUICONTINUITY

Proposition B.1. *For $t > 0$ and $h \in (0, 1)$, we have the following estimate.*

$$(B.1) \quad \|v_\varepsilon(\cdot, t+h) - v_\varepsilon(\cdot, t)\|_\infty \leq C_1 h^{\frac{1}{2}} \left\{ \sqrt{\varepsilon} \|\varphi\|_\infty^{\frac{1}{p}} t^{-\frac{1}{p}} + Q \left(\lambda_p \|\varphi\|_\infty^{\frac{1}{p}} (at)^{-\frac{1}{p}} \right) \right\},$$

where $Q(r) = \max_{0 \leq s \leq r} H(s)$, and C_1 depends only on p , a and N .

Proof. To establish this estimate, we consider, for $r > 0$, $y \in \mathbb{R}^N$ the ball $B_r(y) = \{x \in \mathbb{R}^N, |x - y| < r\}$. Integrating (1.3) over $B_r(y) \times [t, t+h]$ we get (where $|B_r|$ is the Euclidean volume of the ball)

$$\begin{aligned} (v_\varepsilon(y, t+h) - v_\varepsilon(y, t))|B_r| &= \left[\int_{B_r(y)} ((v_\varepsilon(y, \tau) - v_\varepsilon(x, \tau)) dx \right]_{\tau=t}^{\tau=t+h} \\ &+ \varepsilon \int_t^{t+h} \int_{\partial B_r(y)} \nabla v_\varepsilon \cdot \frac{x}{r} dS d\tau \\ &- \int_t^{t+h} \int_{B_r(y)} H(|\nabla v_\varepsilon(x, \tau)|) dx d\tau. \end{aligned}$$

Invoking (4.2) we obtain readily

$$\begin{aligned} |v_\varepsilon(y, t+h) - v_\varepsilon(y, t)| &\leq C \|\varphi\|_\infty^{\frac{1}{p}} \left\{ \frac{\varepsilon}{r} \left[(t+h)^{1-\frac{1}{p}} - t^{1-\frac{1}{p}} \right] + r \left[(t+h)^{-\frac{1}{p}} + t^{-\frac{1}{p}} \right] \right\} \\ &+ \int_t^{t+h} Q \left(\lambda_p \|\varphi\|_\infty^{\frac{1}{p}} (a\tau)^{-\frac{1}{p}} \right) d\tau, \end{aligned}$$

where $C > 0$ depends only on p , a and N . Taking $r = \sqrt{\varepsilon h}$ we obtain (B.1). \square

APPENDIX C. A COMPARISON PRINCIPLE FOR SUBADDITIVE AND NON-DECREASING H

Lemma C.1. *Assume that $H \in C([0, \infty))$ is a nonnegative and non-decreasing function such that $H(0) = 0$ and H is uniformly Lipschitz continuous in (δ, ∞) for each $\delta > 0$. Assume further that H is subadditive, that is,*

$$H(r+s) \leq H(r) + H(s), \quad (r, s) \in [0, \infty).$$

Then Equation (1.1) satisfies the comparison principle in $C_b(\mathbb{R}^N)$ as stated in Definition 1.3 (b).

The proof of Lemma C.1 which we give below was kindly indicated to us by G. Barles [6].

Proof. Let $v_1 \in C_b(\mathbb{R}^N \times (0, \infty))$ (resp. $v_2 \in C_b(\mathbb{R}^N \times (0, \infty))$) be a viscosity subsolution (resp. supersolution) of (1.1) and assume that $v_1(x, 0) \leq v_2(x, 0)$ for $x \in \mathbb{R}^N$. We first infer from the monotonicity and subadditivity of H that

$$H(|\xi_1|) + H(|\xi_2 - \xi_1|) \geq H(|\xi_1| + |\xi_2 - \xi_1|) \geq H(|\xi_2|), \quad (\xi_1, \xi_2) \in \mathbb{R}^N \times \mathbb{R}^N.$$

Setting $w = v_1 - v_2$, it readily follows from the properties of v_1, v_2 and the previous inequality that w is a subsolution to $\partial_t z - H(|\nabla z|) = 0$ in $\mathbb{R}^N \times (0, \infty)$ with $w(\cdot, 0) \leq 0$.

Now, on the one hand, if $\delta \in (0, 1)$, we have $-H\left(\sqrt{|\xi|^2 + \delta^2}\right) \leq -H(|\xi|)$ for $\xi \in \mathbb{R}^N$ by the monotonicity of H . Consequently, w is also a subsolution to

$$(C.1) \quad \partial_t z - H\left(\sqrt{|\nabla z|^2 + \delta^2}\right) = 0 \quad \text{in} \quad \mathbb{R}^N \times (0, \infty)$$

with $w(\cdot, 0) \leq 0$. On the other hand, $W_\delta : t \mapsto H(\delta)t$ clearly solves (C.1) with $W_\delta(0) = 0$ and the Hamiltonian $H_\delta : \xi \mapsto H\left(\sqrt{|\xi|^2 + \delta^2}\right)$ is uniformly Lipschitz continuous in \mathbb{R}^N . We are then in a position to apply [11, Theorem V.3] and

conclude that $w(x, t) \leq W_\delta(t)$ for $(x, t) \in \mathbb{R}^N \times (0, \infty)$. Since H is continuous and vanishes at zero, the claimed result then follows by letting $\delta \rightarrow 0$. \square

REFERENCES

- [1] L. Amour and M. Ben-Artzi, *Global existence and decay for viscous Hamilton-Jacobi equations*, *Nonlinear Anal.* **31** (1998), 621–628.
- [2] M. Bardi and I. Capuzzo-Dolcetta, “*Optimal Control and Viscosity Solutions of Hamilton-Jacobi-Bellman Equations*”, *Systems Control Found. Appl.*, Birkhäuser, Boston, 1997.
- [3] G. Barles, *Regularity results for first order Hamilton-Jacobi equations*, *Differential Integral Equations* **3** (1990), 103–125.
- [4] G. Barles, *Uniqueness and regularity results for first order Hamilton-Jacobi equations*, *Indiana Univ. Math. J.* **39** (1990), 443–466.
- [5] G. Barles, “*Solutions de Viscosité des Équations de Hamilton-Jacobi*”, *Mathématiques & Applications* **17**, Springer-Verlag, Paris, 1994.
- [6] G. Barles, *personal communication*, 2004.
- [7] S. Benachour and Ph. Laurençot, *Global solutions to viscous Hamilton-Jacobi equations with irregular initial data*, *Comm. Partial Differential Equations* **24** (1999), 1999–2021.
- [8] Ph. Bénilan, “*Evolution equations and accretive operators*”, lecture notes taken by S. Lenhardt, Univ. of Kentucky, 1981.
- [9] Ph. Bénilan and M.G. Crandall, *Regularizing effects of homogeneous evolution equations*, in “*Contributions to Analysis and Geometry (Baltimore, MD., 1980)*”, pp. 23–39, John Hopkins Univ. Press, Baltimore, Md., 1981.
- [10] M.G. Crandall, L.C. Evans and P.-L. Lions, *Some properties of viscosity solutions of Hamilton-Jacobi equations*, *Trans. Amer. Math. Soc.* **282** (1984), 487–502.
- [11] M.G. Crandall and P.-L. Lions, *Viscosity solutions to Hamilton-Jacobi equations*, *Trans. Amer. Math. Soc.* **277** (1983), 1–42.
- [12] M.G. Crandall, *Viscosity solutions: a primer*, in “*Viscosity Solutions and Applications (Montecatini Terme, 1995)*”, I. Capuzzo Dolcetta and P.-L. Lions (Eds.), pp. 1–43, *Lecture Notes in Math.* **1660**, Springer, Berlin, 1997.
- [13] L.C. Evans, “*Partial Differential Equations*”, *Grad. Stud. Math.* **19**, Amer. Math. Soc., Providence, RI, 1998.
- [14] B.H. Gilding, M. Guedda and R. Kersner, *The Cauchy problem for $u_t = \Delta u + |\nabla u|^q$* , *J. Math. Anal. Appl.* **284** (2003), 733–755.
- [15] B.H. Gilding, *The Cauchy problem for $u_t = \Delta u + |\nabla u|^q$, large-time behaviour*, *J. Math. Pures Appl.* **84** (2005), 753–785.
- [16] H. Ishii, *Comparison results for Hamilton-Jacobi equations without growth condition on solutions from above*, *Appl. Anal.* **67** (1997), 357–372.
- [17] M.A. Katsoulakis, *A representation formula and regularizing properties for viscosity solutions of second-order fully nonlinear degenerate parabolic equations*, *Nonlinear Anal.* **24** (1995), 147–158.
- [18] P.-L. Lions, “*Generalized Solutions of Hamilton-Jacobi Equations*”, *Res. Notes in Math.*, **69**, Pitman, Boston, 1982.
- [19] P.-L. Lions, *Regularizing effects for first-order Hamilton-Jacobi equations*, *Appl. Anal.* **20** (1985), 283–307.
- [20] T. Strömberg, *The Hopf-Lax formula gives the unique viscosity solution*, *Differential Integral Equations* **15** (2002), 47–52.
- [21] A. Tafia, *Nonlinear viscous Hamilton-Jacobi equations with gradient dependent function*, M. Sc. Thesis, Hebrew University of Jerusalem (2004, unpublished).
- [22] J.L. Vázquez and M. Walias, *Existence and uniqueness of solutions of diffusion-absorption equations with general data*, *Differential Integral Equations* **7** (1994), 15–36.

SAÏD BENACHOUR: INSTITUT ELIE CARTAN, UNIVERSITÉ HENRI POINCARÉ, F-54506 VANDŒUVRE-LÈS-NANCY CEDEX, FRANCE.

E-mail address: Said.Benachour@iecn.u-nancy.fr

MATANIA BEN-ARTZI: INSTITUTE OF MATHEMATICS, HEBREW UNIVERSITY OF JERUSALEM,
JERUSALEM 91904, ISRAEL.

E-mail address: `mbartzi@math.huji.ac.il`

PHILIPPE LAURENÇOT: INSTITUT DE MATHÉMATIQUES DE TOULOUSE, CNRS UMR 5219, UNI-
VERSITÉ DE TOULOUSE, 118 ROUTE DE NARBONNE, F-31062 TOULOUSE CEDEX 9, FRANCE.

E-mail address: `Philippe.Laurencot@math.univ-toulouse.fr`