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Efficient modeling approach for optimization of a system based on passive diamagnetic levitation as a platform for bio-medical applications

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Abstract. Constrained optimization may use time-consuming analysis on standard methods such as the finite-element method (FEM). In order to decrease the optimization time and to enable flexible gradient based optimization, this paper reports, a semi-analytical modelling approach and optimization of two variants of a diamagnetic levitation system as contactless manipulation platform for bio-medical and bio-chemical applications.

Key words: Optimal design, Bio-medical systems, Microsystems, Diamagnetic levitation.

I. LEVITATION IN BIOMEDICAL APPLICATIONS

The magnetic biochips are in increasing development and evolution supported by several applications in the biochemical and biomedical fields [1-4]. This family of microsystems includes partial or complete integrations of elementary or complex tasks processed by traditional laboratories of our scale. The automation and parallelization of these tasks [2] in the small world avoids human handling which is frequently not precise and likely of errors, and reduces times and the costs necessary to achieve operations such as diagnostic, analyzes and particles sorting [2-4]. This evolution of micro-actuators based on magnetic actuation has lead to an increasing need for specific modeling methodologies especially for optimization requirements. Thus, in order to facilitate the use of gradient based optimization process [5], this work deals with a realistic semi-numerical modeling approach of a system based diamagnetic levitation declined in two variants. Indeed, these systems have more and more applications, especially in biomedical field [6-7]. Diamagnetic contactless handling of bioparticles brings amazing perspectives to biochips and avoids serious problems as adhesion, adsorption and contamination.

II. MAGNETOPHORESIS FORMULATION AND MODELING

Currently, existing modeling approaches treating diamagnetic levitation problems are numerical methods such as finite elements method or magnetic moments method [8]. These modeling approaches are not adequate for full design studies. Therefore, in order to permit an adaptable optimization with gradient based algorithms, semi-analytical modeling was chosen and then established to compute outputs and formally exact sensitivities according to the inputs parameters of design [9]. The proposed device comprises four fix magnets and a levitating diamagnetic object. In order to approach both possible diamagnetic materials: anisotropic and isotropic and both possible fluids for biological bodies: air and liquid, the models of a levitating anisotropic graphite plate in air, and a levitating isotropic latex sphere in fluid were established. The magnetic field generated by these four parallelepiped magnets was expressed using equivalent surface method [10]. In this formulation the induced induction of the diamagnetic levitating object was neglected. This induced field did not affect significantly the external field due to the low magnetic susceptibilities of the diamagnetic materials. Hereafter is reported the magnetophoretic force [11] formulations corresponding to both proposed cases.

A. Modeling of an anisotropic diamagnetic graphite plate in air

A diamagnetic anisotropic plate in air with (2a, 2b, 2c) dimensions, having magnetic susceptibilities (χ_x , χ_y , χ_z) respectively on x, y and z-axis and within a magnetic field induction expressed by its components B_x , B_y and B_z , experiences a magnetic force obtained from integral over whole diamagnetic plate volume (2). According to Ostrogradsky theorem (3) [12], magnetic force components (magFx, magFy and magFz) can be simplified leading to surface integrals (4-6). This simplification allows to establish a flexible semi-numerical model based on surface integrals to compute the magnetic force components and thus to enables a faster optimization.

$$G(x, y, z) = \chi_x B_x(x, y, z)^2 + \chi_y B_y(x, y, z)^2 + \chi_z B_z(x, y, z)^2 \quad (1)$$

$$\text{magF}(x, y, z) = \frac{1}{2\mu_0} \iiint_V \nabla G(x, y, z).dv \quad (2)$$

$$\iiint_V \nabla f .dv = \oint_S f .\vec{n}.ds \quad (3)$$

$$\text{magFx}(x, y, z) = \frac{1}{2\mu_0} \iiint_V \frac{\partial G(x, y, z)}{\partial x}.dv = \frac{1}{2\mu_0} \int_{-c}^c \int_{-b}^b (G(x+a, y+Y, z+Z) - G(x-a, y+Y, z+Z))dY.dZ \quad (4)$$

$$\text{magFy}(x, y, z) = \frac{1}{2\mu_0} \iiint_V \frac{\partial G(x, y, z)}{\partial y}.dv = \frac{1}{2\mu_0} \int_{-c}^c \int_{-a}^a (G(x+X, y+b, z+Z) - G(x+X, y-b, z+Z))dX.dZ \quad (5)$$

$$\text{magFz}(x, y, z) = \frac{1}{2\mu_0} \iiint_V \frac{\partial G(x, y, z)}{\partial z}.dv = \frac{1}{2\mu_0} \int_{-b}^b \int_{-a}^a (G(x+X, y+Y, z+c) - G(x+X, y+Y, z-c))dX.dY \quad (6)$$

where μ_0 is the vacuum magnetic permeability.

B. Modeling of an isotropic diamagnetic latex sphere in a liquid

A diamagnetic isotropic sphere having a radius R within a susceptible liquid experiences a magnetic force according to magnetic susceptibility contrast $\Delta\chi$ between the liquid and the sphere material. This force is obtained with an integral over the whole sphere volume (9). Then magnetic force components (magFx, magFy and magFz) can be simplified leading to surface integrals (10-12). This simplification allows to establish a flexible semi-numerical model to compute magnetic force components based on surface integrals and thus a faster optimization.

$$G(x, y, z) = B(x, y, z)^2 = Bx(x, y, z)^2 + By(x, y, z)^2 + Bz(x, y, z)^2 \quad (7)$$

$$g1(y, z) = \sqrt{R^2 - y^2 - z^2}, \quad g2(x, z) = \sqrt{R^2 - x^2 - z^2}, \quad g3(x, y) = \sqrt{R^2 - x^2 - y^2} \quad (8)$$

$$\text{magF}(x, y, z) = \frac{\nabla \chi}{2\mu_0} \iiint_V \nabla G(x, y, z).dv \quad (9)$$

$$\text{magFx}(x, y, z) = \frac{\nabla \chi}{2\mu_0} \int_{-R}^R \int_{-\sqrt{R^2 - Y^2}}^{\sqrt{R^2 - Y^2}} (G(x + g1(Y, Z), y + Y, z + Z) - G(x - g1(Y, Z), y + Y, z + Z))dZ.dY \quad (10)$$

$$\text{magFy}(x, y, z) = \frac{\nabla \chi}{2\mu_0} \int_{-R}^R \int_{-\sqrt{R^2 - Z^2}}^{\sqrt{R^2 - Z^2}} (G(x + X, y + g2(X, Z), z + Z) - G(x + X, y - g2(X, Z), z + Z))dX.dZ \quad (11)$$

$$\text{magFz}(x, y, z) = \frac{\nabla \chi}{2\mu_0} \int_{-R}^R \int_{-\sqrt{R^2 - X^2}}^{\sqrt{R^2 - X^2}} (G(x + X, y + Y, z + g3(X, Y)) - G(x + X, y + Y, z - g3(X, Y)))dY.dX \quad (12)$$

C. Automated modeling tool

A modeling tool [13] based on this diamagnetic computation method was developed. Designer has to describe the geometrical and physical system (size, magnetic parameters, etc) and the tool generates automatically these formulas. This tool (Fig. 1) computes forces with magnets, current conductors with several shapes (parallelepiped, sphere, cylinder, arc of cylinder, etc).

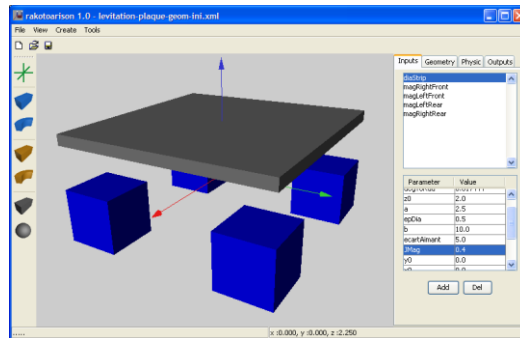


Fig. 1 – Interface screenshot of the tool used to generate designs computation formulas.

III. MODELING VALIDATION

The proposed device is composed of a diamagnetic anisotropic plate having a susceptibility $\chi_z = -4.510^{-4}$ and $\chi_x = \chi_y = -8.10^{-5}$, a density $\rho_g = 2200 \text{ kg/m}^3$. It levitates above four magnets configured as shown on Fig. 2-(a). The magnetic vertical force experienced by the plate was computed semi-analytically, and then it was compared to the magnetic moment method simulation for which the induced magnetization of diamagnetic object is not neglected. The results obtained from these two modeling methods are reported on Fig. 2-(b). These results show

a performing modeling and that approximation (no induced induction into the diamagnetic material) is realistic.

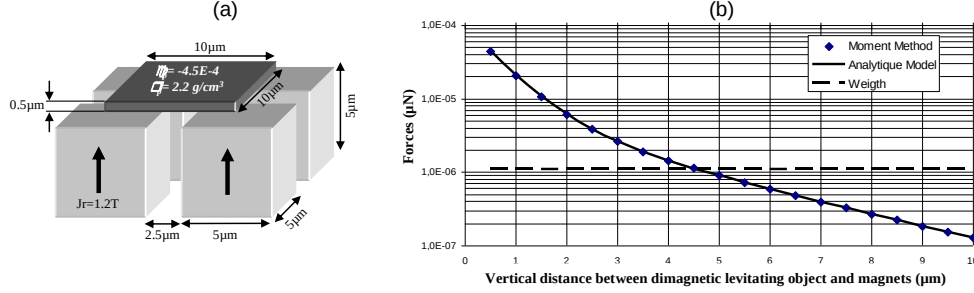


Fig. 2 – (a) Device configuration, (b) Simulation results of the magnetic force F_z obtained from the semi-analytical model and the magnetic moment method.

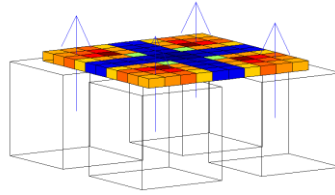


Fig. 3 – The design meshed with an tool based magnetic moments computation.

FEM computation was also compared with these modeling techniques but has lead to interpretation issues. Indeed, diamagnetic forces, which are about some picoNewtons, are drowned into meshing noise. Thus to compute forces with a FEM tool, meshing noise has to be removed by this technique.

- Compute the force on the diamagnetic plate with $\chi=0$. The force value result only on the noise.
- Compute the force on the diamagnetic plate with real χ value, without changing the meshes.
- Subtract the result value with preceding value to obtain the real force (Table. 1).

TABLE I: FORCE COMPUTATION WITH FEM METHOD

Vertical distance (μm)	Force computation with $\chi=0$ (1)	Computation with $\chi=10^{-5} \times [-8, -8, 4.5]$ (2)	Resulting force=(2)-(1)
3	13.5725018E-9	13.5698702E-9	2.6316E-12

IV. OPTIMIZATION

The optimizations performed tries to converge towards magnets structures for given levitating objects and levitation heights. Its objective is to maximise the spatial stability of the object. The stability of the levitating object is quantified by the total force gradient acting on corresponding axis. These optimizations respond to precise constrained specifications. A semi-analytical model characterizing the device was established. It enables computation of the forces exerted on the diamagnetic levitating object and theirs corresponding gradients as outputs. These outputs are required to express the stability of the levitating object according to structure geometry and physical parameters. In order to use an efficient optimization algorithm like SQP (Sequential Quadratic Programming), an automatic tool was used to generate formally exact sensitivities of the established semi-analytical model. This tool is called a generator and is part of the CADES design framework pattern shown Fig. 4, which describes how model and optimizer can interact in an iterative and reusable manner.

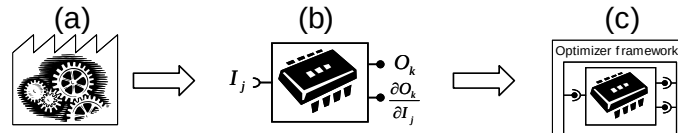


Fig. 4 – (a) Generators create, (b) components, (c) Services use.

The symbolic derivation of the model produces formal exact sensitivity (without noise like finite differences) but also quicker model. Indeed, force sensitivity depending on a diamagnetic size parameter, leads to a simple integral formula as describe in the following equations. Thus, as numerical integration take most of the model computation time, this may reduce consequently the global sizing time.

$$f(a, b) = \int_0^b \int_0^a g(x, y) dx dy \quad (13)$$

$$\frac{\partial f(a, b)}{\partial b} = \int_0^a [g(x, b) - g(x, 0)] dx \quad (14)$$

Using this design framework, the graphite plate and the polystyrene latex sphere devices were sized as shown in the following.

A. Diamagnetic anisotropic graphite plate in air

The device comprises a fixed part and a mobile one as shown on Fig. 5. The fix part of the device comprises four electroplated CoPt magnets [14] magnetized throughout the plane with a residual magnetization $J_z=0.3$ T. These magnets are identical and have as parameters *magsize*, *magheight* and *gap* which correspond respectively to the magnets squared base size, magnets height and distance between two adjacent magnets. The mobile part is a plate made of pyrolytic graphite which is an anisotropic material. Its magnetic susceptibilities on x,y and z axis are $[\chi_x, \chi_y, \chi_z] = 10^{-5} \times [-8, -8, -45]$ and its density $\rho_g=2200$ kg/m³. This plate has as parameters: *diasize* (the squared base size of the plate) and *diaheight* (height of the plate). Both parameters were fixed respectively at 10 μm and at 0.5 μm . The objective of optimization was to converge towards a solution giving an optimal stability satisfying a fixed levitation height of 2 μm . This levitation height is computed between top of magnets and bottom of diamagnetic plate at position where the total force on z axis become null. Neglecting the Archimedes force due to low density of the medium. Only two forces act vertically on the diamagnetic plate: the weight of the plate (*diaWeight*) and the magnetic force (*magFz*) (6). Therefore, the balance of vertical forces *FtotZ* can be expressed as:

$$F_{totZ} = \text{magFz} + \text{diaWeight} \quad (15)$$

The parameters *magsize*, *magheight* and *gap* were constrained as given on Table. 2. The optimization results corresponding to these constrained specifications are illustrated on Table. 3, where is given inputs parameters and the computed outputs values corresponding to two geometries. The diamagnetic plate at the first geometry shown on Fig. 5-(a) has a levitation height of 1.75 μm and is unstable on x axis and stable on vertical one. The second geometry shown on Fig. 5-(c) corresponds to an optimal stability obtained after optimization. The forces curves corresponding to both devices are shown on Fig. 6.

TABLE II: OPTIMIZATION SPECIFICATIONS

Constrained parameters (μm)	<i>magsize</i> =[1, 20]	<i>magheight</i> =[1, 10]	<i>gap</i> =[0.1, 5]	<i>Dist</i> =2
Objective	Optimal stability			

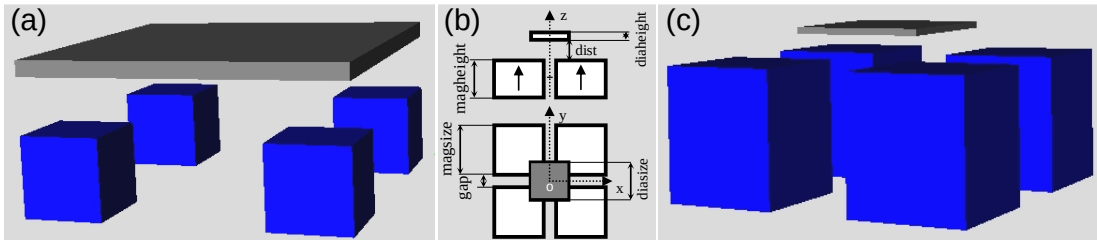


Fig. 5 – (a) Geometry of device with unstable plate on x axis, (b) Geometry parameters of the diamagnetic plate and the magnets, (c) Optimal design for axial stability of the diamagnetic anisotropic plate.

TABLE III: UNSTABLE DESIGN PARAMETERS AND OPTIMIZATION RESULTS

Name	Designation	Unstable design	Optimal design	unity
<i>diaWeight</i>	Weight of diamagnetic plate	-1.0788	-1.0788	pN
<i>magFz</i>	Vertical magnetic force	-1.0799	1.0791	pN
<i>FtotZ</i>	Balance of vertical forces	1.1611E-9 (~0 N)	3.0364E-10 (~0 N)	μN
<i>gradFx_x</i>	Gradient of <i>Fx</i> along x axis	0.1808(unstable)	-1.3668 (stable)	$\mu\text{N/m}$
<i>gradFz_z</i>	Gradient of <i>Fz</i> along z axis	-0.1916 (stable)	-1.672 (stable)	$\mu\text{N/m}$
<i>gap</i>	Distance between two magnets	5	6.4430	μm
<i>magheight</i>	Magnets height	2.5	10	μm
<i>magsize</i>	Magnets squared base size	2.5	8.5960	μm

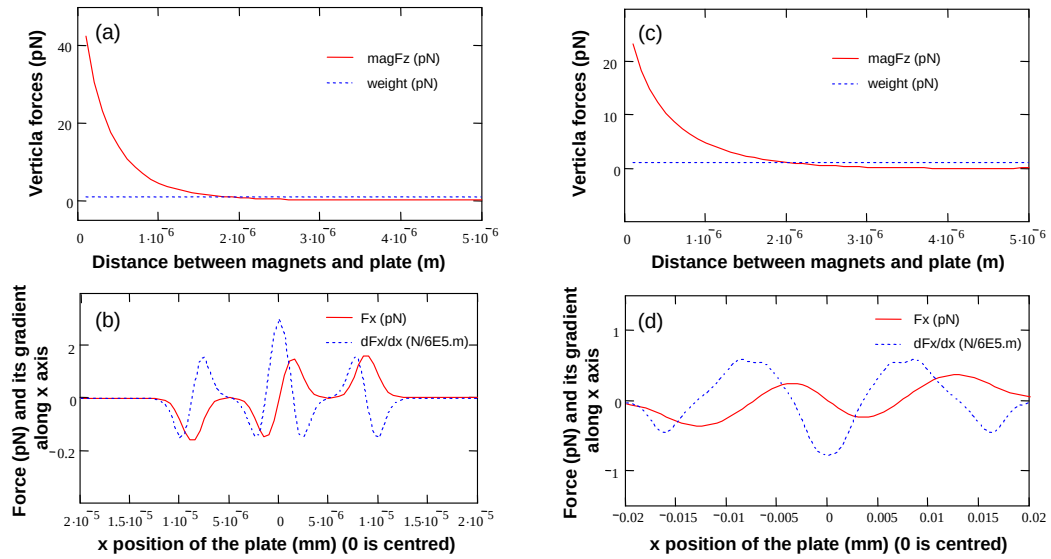


Fig. 6 – Forces experienced by the diamagnetic plate, (a) vertical forces according to the distance between the magnets and the plate at the case of the unstable design, (b) horizontal force and its gradient according to the x position of the plate at the case of the unstable design, (c, d) Case of the stable design.

The Pareto optimization results shown on Fig. 7 validate the obtained results of the classic optimization. This one was done in two steps. Firstly isovalues of the gradient of the force on x axis where plotted as function of the parameters gap and magsize. Secondly the curve of the parameter gap was plotted as a function of magsize with the total force on the vertical axis and its gradient constrained to be respectively null and negative to satisfy the stability criterion on vertical axis, where the force on x-axis and its gradient were free.

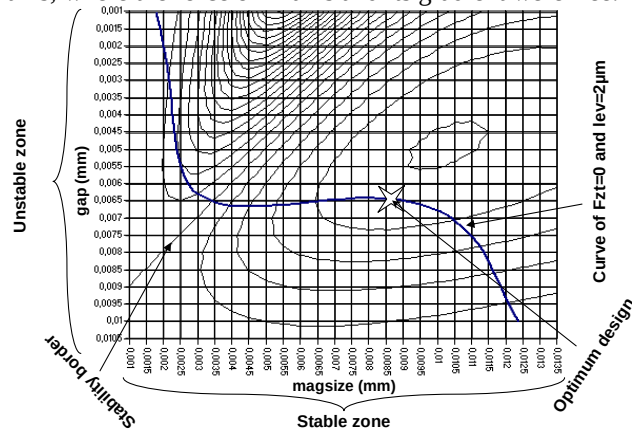


Fig. 7 – 2D cartography of stability on x axis and the optimal design localisation expressed according to the squared base size of the magnet and the gap between the magnets.

B. Diamagnetic latex bead in paramagnetic medium

The fix part of device is still the same used above for the diamagnetic graphite plate in air. However the mobile part is now a diamagnetic sphere made of polystyrene latex which is an isotropic material. Its susceptibility is $\chi_p = -8.2 \cdot 10^{-6}$ and its density $\rho_p = 1050 \text{ kg/m}^3$. The medium is a paramagnetic configured with about 0.47 M of manganese chloride tetrahydrate ($\text{MnCl}_2 \cdot 4\text{H}_2\text{O}$) in pure water, having a magnetic susceptibility $\chi_m = 7.10^{-5}$ and a density $\rho_m = 1000 \text{ kg/m}^3$. This sphere has as parameter diasize (the radius of the sphere) fixed at $1.5 \mu\text{m}$. The objective of optimization is to converge towards a solution giving an optimal stability satisfying a levitation height of $3 \mu\text{m}$. The parameters magsize, magheight and gap were constrained as given on Table. 4. The optimization results corresponding to these constrained specifications are shown on Table. 5. It illustrates the inputs parameters and the computed outputs values corresponding to two geometries. The diamagnetic sphere is unstable at the first geometry shown on Fig. 8-(a) on x axis and stable on vertical one and the levitation height is $3 \mu\text{m}$. The second one shown on Fig. 8-(c) corresponds to a solution giving an optimal horizontal stability which was obtained after optimization process. The levitation height is computed between top of magnets and bottom of diamagnetic sphere at the position where the total force on z axis become null. The levitation height is fixed at

dist=3 μm . This position is determined from the balance of forces experienced by the sphere on vertical axis. Due to spherical geometry and the density of the medium, three forces acts vertically: diaweight, diaArch (16) and FtotZ (17) respectively the weight of the sphere, the Archimedes force and the balance of forces. The forces curves corresponding to unstable and stable devices are shown on Fig. 9.

$$\text{diaArch} = \frac{4}{3}\pi\left(\frac{\text{diasize}^3}{8}\right)\rho_m g \quad (16)$$

$$\text{FtotZ} = \text{magFz} + \text{diaWeight} + \text{diaArch} \quad (17)$$

TABLE IV: OPTIMIZATION SPECIFICATIONS

Constrained parameters (μm)	magsize=[4, 12]	magheight=[4, 10]	gap=[1, 7]	Dist=3
Objective	Optimal stability on x axis			

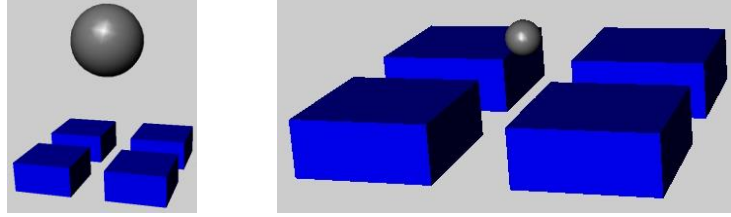


Fig. 8 – (a) Geometry of device with unstable plate on x axis, (b) Optimal design for axial stability of the diamagnetic plate.

TABLE V: UNSTABLE DESIGN PARAMETERS AND OPTIMIZATION RESULTS

Name	Designation	Unstable design	Optimal design	unity
gap	Distance between two magnets	1.4	6.0075	μm
Magheight	Magnets height	1.26552	4.7046	μm
Magsize	Magnets squared base size	2.40428	11.0468	μm
diaWeight	Weight of sphere	0.14561	0.14561	pN
diaArch	Archimedes force	0.13868	0.13868	pN
magFz	Vertical magnetic force	0.007	0.00697	pN
FtotZ	Balance of vertical forces	0.000066 (~0 N)	0.000037 (~0 N)	pN
gradFx_x	Gradient of Fx along x axis	0.31019 (unstable)	-2.8182 (stable)	nN/m
gradFz_z	Gradient of Fz along z axis	-4.0304 (stable)	-16.684 (stable)	nN/m

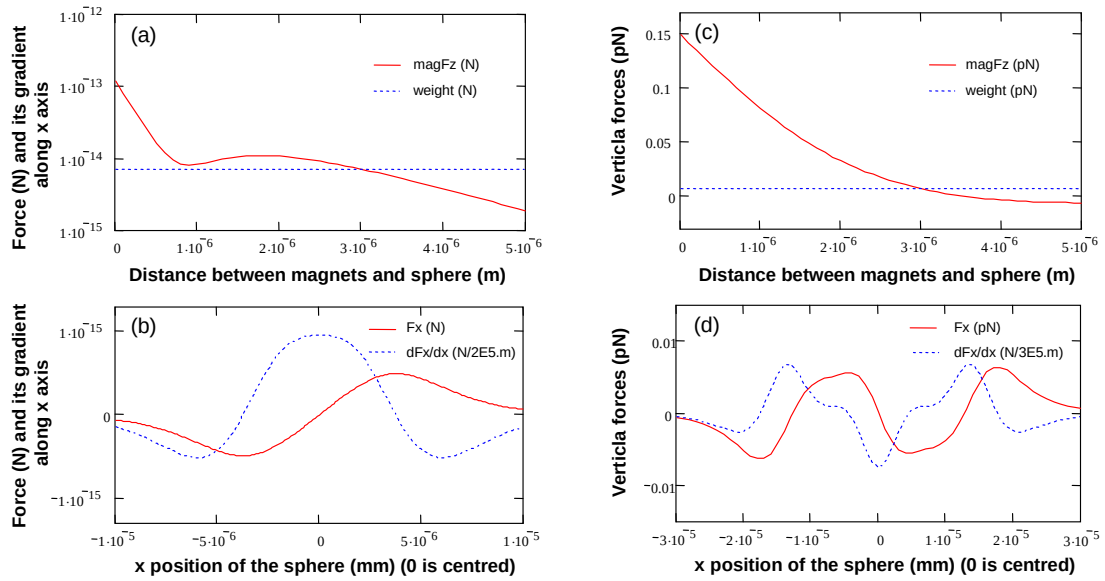


Fig. 9 – (a,b) Forces experienced by the diamagnetic plate, (a) vertical forces according to the distance between the magnets and the plate at the case of the unstable design, (b) horizontal force and its gradient according to the x position of the plate at the case of the unstable design, (c, d) Case of the stable design.

V. CONCLUSION

We report an efficient and accurate semi-analytical modeling and optimization approach of devices based on passive diamagnetic levitation. The studied devices are composed of a levitating object and fixed part of four permanent magnets dedicated for biochemical and biomedical applications. In order to establish semi-analytical modeling of the proposed device surface approach was used to compute analytically the magnetic field induction produced by the magnets. Forces components formulas experienced by the levitating objects were also simplified to surface integral expressions. These simplifications enabled faster optimizations which were thereafter performed by a SQP algorithm. The results demonstrate a solution to perform a fast and flexible optimization of devices based on diamagnetic levitation which can be generalized to other application fields. This avoids the usual approaches which are not adequate for the optimization task in this field of applications.

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