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The dominant role of structure for solute transport in soil: Experimental evidence and modelling of structure and transport in a field experiment

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Abstract

A classical transport experiment was performed in a field plot of 2.5 m² using the dye tracer brilliant blue. The measured tracer distribution demonstrates the dominant role of the heterogeneous soil structure for solute transport. As with many other published experiments, this evidences the need of considering the macroscopic structure of soil to predict flow and transport.

We combine three different approaches to represent the relevant structure of the specific situation of our experiment: i) direct measurement, ii) statistical description of heterogeneities and iii) a conceptual model of structure formation. The structure of soil layers was directly obtained from serial sections in the field. The sub-scale heterogeneity within the soil horizons was modelled through correlated random fields with estimated correlation lengths and anisotropy. Earthworm burrows played a dominant role at the transition between the upper soil horizon and the subsoil. A model based on percolation theory is introduced that mimics the geometry of earthworm burrow systems.

The hydraulic material properties of the different structural units were obtained by direct measurements where available and by a best estimate otherwise. From the hydraulic structure, the 3-dimensional velocity field of water was calculated by solving Richards' Equation and solute transport was simulated. The simulated tracer distribution compares reasonably well with the experimental data. We conclude that a rough representation of the structure and a rough representation of the hydraulic properties might be sufficient to predict flow and transport, but both elements are definitely required.

1. Introduction

Solute transport is important to plant uptake, chemical reactions and transformation processes, and is a key factor in controlling groundwater contamination.

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Despite its fundamental importance, solute transport in soil is still far from being quantitatively understood. This means we are typically not in the position to predict solute transport with respect to the first arrival time in a given depth or even with respect to the mean velocity.

The major hurdle is the heterogeneity of soils at various spatial scales with the consequence that material properties and model parameters are not well defined. In other words, for a given soil sample, averaged effective properties only exist if the scale of heterogeneities is much smaller than the size of the sample, which is the prerequisite for the existence of a representative elementary volume (REV). In the following, we want to use the term *“structure”* for that type of heterogeneity which is of similar size as the scale of interest (i.e. size of sample) and hence, can not be represented by effective properties.

A well known phenomenon, which can be considered as a consequence of heterogeneous material properties, is “preferential flow”, in which some fraction of a solute moves much faster than expected, through a small fraction of the medium. In this case, wrong expectations typically originate from the measurement of ‘effective’ material properties assuming homogeneity in the sense of a REV while the preferential flow paths are not captured by the measurement. This phenomenon of preferential flow was first attributed to continuous macropores (Bouma, 1981; Beven and Germann, 1982, e.g.) and has been evidenced for years in both laboratory and field experiments. The use of dye tracers by Flury et al. (1994) allowed a detailed study of the phenomenology of solute transport thereby revealing that preferential flow as a consequence of various soil structural features was the rule rather than the exception. It is not only produced by macropores but also, more generally, by the heterogeneous structure of hydraulic properties.

The common ad-hoc approach for modelling preferential flow is to separate different zones of different mobility with the idea that solutes are translocated only in a restricted part of the porous network (van Genuchten and Wierenga, 1976; Gerke and van Genuchten, 1993; Jarvis, 1994). This approach adds considerable flexibility to the

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classical convection-dispersion model, so that experimental findings can be described quite accurately. However these models are only effective descriptions of the natural complexity, and the additional parameters – the volume fraction of zones with different mobility and the rate parameter controlling the exchange between these zones – cannot be measured independently. Therefore, only a retrospective description of a specific experiment is feasible through the fitting of the model parameters to measured breakthrough curves or spatial distributions of concentration data. Moreover, the obtained parameters are expected to change with the initial and boundary conditions due to the changing geometry of the flow field (Clothier et al., 1995) and with the transport distance due to the heterogeneity of the material (Vanderborght et al., 2001). Hence, these models typically cannot be used to predict flow and transport for conditions other than those used for calibration.

In the context of flow and transport, relevant structures are ubiquitous at a wide range of spatial scales, which has been formalised for soils using the concept of fractal geometry (Bavey and Boast, 1998) or the approach of discrete or continuous hierarchy (Cushman, 1990). Vogel and Roth (2003) proposed the “scaleway” as a discrete hierarchical approach for modelling flow and transport, which is based on the explicit consideration of the macro structure while microscopic heterogeneities are averaged and replaced by effective descriptions. The major drawback of this approach is the fact that a full three-dimensional representation of the relevant structure is required. Given the heterogeneity of natural soil, there seems to be no loophole out of this dilemma, if we would like to come up with a reliable prediction.

In this paper we intend to present three different approaches to incorporating the relevant 3D-structure of a given soil for modelling flow and transport, and we discuss meanings of the term “relevant” in this context.

The first approach, which will be referred to as the “explicit” approach, is based on the direct measurement of the structure using appropriate instruments. This can be done using destructive methods as e.g. serial sectioning. Today, however, there are powerful instruments available for the non-invasive measurement of structure using

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the concept of tomography. At the small scale X-ray tomography is used to identify the structure of macropores and zones of contrasting bulk densities (Hopmans et al., 1994; Wildenschild et al., 2002). On larger scales geophysical tools might be applied, such as geoelectric (Puvance and Andricevic, 2000, e.g.) or georadar (Kowalsky et al., 2004, e.g.) techniques.

The second approach is a “statistical” approach. It is based on limited information on relevant structural properties such as the variance and the correlation scale of material properties. Given this information, a continuous representation of the structure can be represented by equivalent random fields. This approach is typically used in the field of stochastic continuum theory (Dagan, 1986; Gelhar, 1986).

The third approach is termed the “genetic” approach. It uses the available knowledge on the processes of structure formation to generate the structure itself. This approach is especially attractive for structural units which are highly relevant but sparse and/or highly anisotropic, so that the first approach is not adequate and the second is insufficient in terms of spatial resolution. Examples are fissures or earthworm burrows.

We apply these different strategies to predict the 3D spatial distribution of a dye tracer, which was measured in a field experiment at steady state flow conditions. The structure of the soil horizons is directly measured by photography (explicit approach), the sub-scale heterogeneity is modelled through estimated correlation lengths within the soil horizons (statistical approach), and the macroporous structure of earthworm burrows connecting the topsoil and the subsoil is modelled based on a few typical properties of the burrow system of anecic earthworms (genetic approach). Ideally, the hydraulic properties for each structural unit are measured directly. Where this was not possible we applied a best estimate. Then, based on the continuous 3-dimensional representation of the structure of the experimental domain and the hydraulic properties of the individual structural units, we simulated steady state water flow using Richards’ equation. Subsequently, the resulting velocity field was used to simulate solute transport assuming convection-dispersion type of transport at the small scale. Finally, the simulated dye-tracer distribution is directly compared to experimental data. The sim-

ulation is considered to be a prediction of solute transport, in the sense that no result of the transport experiment is used for model calibration. It is based on a rather rough description of the spatial heterogeneity and a rather rough approximation of the hydraulic properties. One aim of this study was to evaluate how far we can get in face of such rough approximations – a typical situation in field experiments – by combining structural information with material properties.

2. Experiment

The experiment was conducted in the Beauce region, in France on a 80 cm deep loamy-clay soil classified as a haplic Calcisol (FAO-UNESCO, 1975). It is organised in 3 horizons (Ould Mohammed and Bruand, 1994): from 0 to 35 cm, a formerly ploughed horizon, *A*, with a plough-pan at 35 cm depth; from 35 to 60 cm, a structured *B* horizon, with less organic matter; and from 60 to 80 cm, a cryoturbated calcareous *C* horizon including calcareous stones. At the date of the experiment, in April 2000, the soil was covered with rape. During the 5 years prior to our experiment the soil was treated with minimum tillage restricted to the upper 10 centimetres using a gruber. Before, the soil was periodically ploughed to a depth of 35 cm. Due to the change in tillage practice, the upper *A* horizon was partitioned into two parts: a loose, crumbly structured surface layer, *A*₁, down to 10 cm, and a lower part, *A*₂, from 10 to 35 cm where the soil was markedly more compact with a polyhedral structure and vertical earthworm burrows.

The infiltration experiment was performed on a plot of 1.6×1.6 m. Water and dye tracer were supplied using an automated sprinkler device with 5 nozzles spaced 33 cm apart (Kasteel et al., 2002). The distance between the sprinkler and the soil surface was adjusted to 36 cm to ensure a uniform irrigation of the soil surface. The experiment was performed inside a tent to avoid any disturbance by wind and natural rainfall. The kinetic energy of the small water droplets produced by the nozzles was small enough to avoid modifications of the soil surface. The applied steady state infiltration rate was adjusted to 13 mm h⁻¹ to avoid the formation of puddles at the soil surface. To stain the

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flow paths, we used the food dye brilliant blue FCF which is frequently used because of the supposedly low interaction of the anionic molecule with the soil matrix (Flury and Flühler, 1995). However, it was found that the retardation of brilliant blue might be significant depending on the ionic strength and the soil type (Ketelsen and Meyer-Windel, 1999; Germàn-Heins and Flury, 2000; Kasteel et al., 2002). Consequently we interpret the stained part of the soil as minimum travel distance. The experiment consisted of 3 consecutive periods while the flow rate was kept constant:

1. The soil profile was irrigated for 16 h by 210 mm of tracer-free water. After this period we assumed steady state flow in the upper part of the soil profile.
2. During the second period, 19 mm of the dye tracer (6 g/l brilliant blue) was applied as an extended pulse.
3. Finally, 55 mm of tracer-free water were applied to translocate the dye pulse.

After the infiltration experiment, a vertical soil pit (140 cm width, 90 cm depth) was excavated within the irrigated area. This vertical section was photographed together with a grey frame containing a metric colour scale (Fig. 1). Then, to produce serial sections, a slice of 4 cm was removed from the pit using a knife and another image was taken. This procedure was repeated 23 times to produce a set of 24 digital serial sections with a separation of 4 cm. The individual images cover an area of about 100 cm width and 70 cm height at a resolution of 960×745 pixel with a pixel size of 1 mm².

At a close distance from the experiment, undisturbed soil cores (16.2 cm diameter, 10 cm height) were taken from the A_1 and B horizon. These samples were used for classical multistep-outflow experiments (Hopmans et al., 2002) to determine the hydraulic properties of the different horizons. The sample taken from the compacted A_2 horizon had to be rejected because of a significant disturbance of the structure.

3. Image analysis

Extraction of structural units

First the serial images were corrected for distortion using the spatial coordinates of the edges which are clearly identified by the fixed frame. Then, the main horizons (A_1 , A_2 , B) could be identified based on different colours. After combining the serial sections we got the 3-dimensional structure of the soil horizons at a resolution of 1 mm×40 mm ×1 mm in the x, y, and z direction respectively (Fig. 1). Earthworm burrows were clearly detected as an additional structural unit especially in the compacted horizon A_2 . Evidently, the resolution of the serial sections was far too low to measure the 3-dimensional geometry of such burrows explicitly.

Distribution of dye tracer

The 3-dimensional distribution of the dye tracer was obtained at the same resolution as the soil horizons. The images were recorded in RGB format with brightness values in the range 0–255, and we used the red (r) and the green (g) channel to discriminate blue stained areas without distinguishing different concentrations. A voxel was considered to be stained if

$$\frac{r}{g} < 1.45 \quad \text{and} \quad r + g < [\bar{r} + 60], \quad (1)$$

where \bar{r} is the mean red value of the image. This rule was applied to each image individually. Because the r/g ratio is insensitive to brightness, and the threshold value is adjusted for mean image brightness, differences in illumination between images during the excavation and photography period are accounted for. No further image correction was considered necessary.

4. Modelling soil structure

Soil horizons – explicit description

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A description of the soil horizons was obtained from the serial sections described above, meaning the boundaries between the different horizons which were found to be not exactly horizontal, are directly measured by the serial sections. Together with the experimental determination of the hydraulic properties, a complete 3-dimensional description of the hydraulic conditions should be possible. Such an approach would have to rely on the assumption of macroscopic homogeneity of the different horizons, specifically, that the soil cores used to measure the hydraulic properties include a representative elementary volume (REV). In our experiment, this requirement was definitely violated by the macro-pores within the A_2 horizon formed by earthworms. Moreover, the A_1 and the B horizon exhibited some mesoscopic heterogeneities which were smaller but close to the size of the cores. In the A_1 horizon there were aggregates and clods resulting from tillage; the B horizon had regions of variable bulk density due to variable biological activity (roots, earthworms). Because the characteristic size of these structural features was not much smaller than the thickness of the soil horizons, the assumption of macroscopic homogeneity was critical also in the A_1 and B horizon. For a quantitative understanding of flow and transport, all these mesoscopic heterogeneities might be significant. On the other hand, no direct measurement of the mesoscopic structure is available, which could be used in a numerical model to predict flow and transport. In the following we use a classical statistical approach to represent mesoscopic heterogeneities in the A_1 and the B horizon, and a new genetic approach to represent the macroporous structure within the A_2 horizon.

Mesosopic heterogeneity – statistical approach

In the A_1 and the B horizons a mesoscale heterogeneity was introduced by varying the measured hydraulic properties $h^*(\theta)$ and $K^*(\theta)$ where h is the soil water potential, θ the volumetric water content, and K the hydraulic conductivity. We used a scaling factor c according to (Roth, 1995),

$$c = \frac{h^*(\theta)}{h(\theta)}, \tag{2}$$

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so that the actual water characteristic at a given location, $h(\theta)$, is related to the measured reference properties through c . For the hydraulic conductivity we use

$$c^2 = \frac{K(\theta)}{K^*(\theta)}. \quad (3)$$

This approach implies similarity of the region with different hydraulic properties according to Miller and Miller (1956). The scaling factor c reflects a characteristic length at the pore scale such as grain size or pore radius, which is typically assumed to be log-normally distributed. We considered $\log(c)$ to be normally distributed with expectation 0 and variance $\sigma_c^2=0.25$. We use the same scaling factor to describe both, $h(\theta)$ and $K(\theta)$, while $h(\theta)$ is linearly related to the characteristic length according to Eq. (2) and $K(\theta)$ is proportional to c^2 according to Eq. (3). The spatial correlation of c is described by a Gaussian random field with predefined covariance function and related correlation lengths $\lambda_{x,y,z}$ for the three spatial directions. The random fields were generated according to Robin et al. (1993).

In the A_1 horizon we chose the mesoscopic heterogeneity to be slightly anisotropic with a preferred horizontal orientation, i.e. $\lambda_{x,y}=45$ mm and $\lambda_z=30$ mm. This choice was based on field observation for the orientation of soil clods formed by tillage. In the B horizon, a preferred vertical orientation was assumed with a correlation length of $\lambda_{x,y}=60$ mm in the horizontal and $\lambda_z=120$ mm in the vertical direction, which was also based on field observations.

Macroporous structure – genetic approach

The phenomenology of flow and transport through the compacted soil layer A_2 is dominated by the macroporous structure due to earthworm activity and roots. Direct measurements of the detailed geometry of this structural unit was not possible since the resolution of the field images is much too low, because the distance between the serial sections is too large. However, the macroporous structure cannot be ignored for a quantitative and even qualitative understanding of flow and transport in our experiment. The huge difference between the characteristic size of a structural element λ , in this

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case the diameter of the macropores, and the functional correlation length of the structural unit Λ , here the continuous lengths of the macropores, is a general difficulty. This is true for the measurement as well as for the spatial discretisation of structural components for modelling purposes. In both cases the resolution, i.e. the pixel size, needs to be smaller than λ while the region of measurement should be larger than Λ . The technical difficulties to achieve this requirement increase with increasing anisotropy of the structure.

An alternative to the direct measurement of the geometry of earthworm burrow systems, is to model this geometry using the available knowledge on the way these burrow systems are created. Such an approach was recently proposed by Bastardie et al. (2002) who simulated the burrowing activity of different earthworm species based on the statistics of characteristic types of earthworm movements defined by (Capowiez, 2000). They used this “animal-based” model to investigate the effect of various structural aspects, such as volume, topology and orientation, on the hydraulic properties of earthworm channel systems.

Naturally, this approach cannot reproduce any details related to the multitude of different species having specific habits, which additionally might be modified by the type of substrate they are living in. Also the diameter of earthworm burrows varies considerably. For our experiment, however, the exact diameter is less important with respect to the overall pattern of solute transport because, in any case, the saturated hydraulic conductivity of the macropores is very much higher than that of the surrounding soil matrix. For our application we only used some general information on the basic geometry of earthworm channels available from the literature. Field measurements of borrow systems created by the large anecic species *Lumbricus terrestris*, *Aporrectodea spec.* and *Allolobophora spec.* indicated a vertical main channel with an increased tendency for branching towards the top soil which leads to a funnel like structure (Lighthart et al., 1993; Jegou et al., 1999). As direct information from our field experiment, we estimated the number density of continuous earthworm burrows in the horizontal plane to be 9 m^{-2} within the compacted A_2 horizon. This was done through visual inspection

during the excavation of the soil.

To model a single macropore we used the concepts of percolation theory (Stauffer, 1985). We started from an uncorrelated field of random numbers equally distributed in $(0, 255)$ on a cubic voxel grid. The horizontal dimensions (256 voxel) of this field corresponds to the area attributed to one macropore (1/9 of the total area) and the vertical dimension corresponds to the depth of the compacted horizon A_2 . Then, the random grey field was binarized just above the percolation threshold so that a percolation cluster with at least one continuous path in the vertical direction is produced and we deleted all voxels that did not belong to this continuous path to isolate the backbone of the percolation cluster. Then the most efficient path was determined through numerical simulation of transport through the backbone using a particle tracking algorithm: After reducing the backbone to its skeleton (i.e. a path of 1 voxel thickness) a number of particles were introduced at a randomly chosen top end of the backbone. Then, for each particle a random walk through the backbone was calculated to obtain the shortest path and “critical neck” of the backbone. The latter is defined as the location within the backbone where most of the particles have passed. Below this point only the shortest path was considered while above that point all paths inside the backbone were conserved. Hence, a funnel like macropore was generated reflecting some geometric properties of natural earthworm burrows. Finally the skeleton of the simulated macropore was dilated by a spherical structuring element to adjust the width of the channels to a constant value of 7.5 mm.

This approach for modelling the macroporous structure of the compacted soil layer also reflects the tendency of roots to pass through such layers along earthworm channels (Sveistrup et al., 1997). Once a root has “found” an earthworm channel it follows that channel to reach the more favourable layers below the compacted zone. Figure 2 shows different realizations of generated macropores.

Figure 3 shows the entire 3-D representation of the modelled domain including iso-surfaces of the scaling factors describing the mesoscopic heterogeneity in the A_1 and B horizon and the earthworm burrows in the A_2 horizon.

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5. Numerical simulation of flow and transport

Water flow

Based on the complete three-dimensional representation of the structure within the experimental domain (Fig. 3), we calculated the heterogeneous velocity field during the infiltration experiment using Richards' equation:

$$\frac{\partial}{\partial t}\theta - \nabla \cdot [K(\theta) [\nabla h - \rho_w g]] = 0 \quad (4)$$

where ρ_w is the density of water and g acceleration due to gravity. To solve this equation the hydraulic properties $h(\theta)$ and $K(\theta)$ are required for each location. This can be obtained, given the measured structure of the different horizons and assuming macroscopic homogeneity for each horizon with or without mesoscopic heterogeneity. However, it is not obvious how to handle the macroporous structure of the A_2 horizon. Clearly, Richards' equation does not apply to the flow of free water in macropores. Actually, the problem of coupling flow in macropores with flow in the soil matrix within a single model framework is still not solved. As an approximation, we considered the macropores to be a very coarse porous medium having a low air entry value and a high conductivity.

The hydraulic properties for each structural unit – soil horizons and macropores – are described using a modified van Genuchten model according to [Vogel et al. \(2001\)](#). This model introduces an explicit air entry pressure $h_e > 0$ to the classical model of [van Genuchten \(1980\)](#) and is given by

$$S_e = \begin{cases} \frac{1}{S_e^*} [1 + [\alpha h]^n]^{-m} & : h > h_e \\ 1 & : h \leq h_e \end{cases}, \quad (5)$$

with the effective water saturation

$$S_e = \frac{\theta - \theta_r}{\theta_s - \theta_r}. \quad (6)$$

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θ_s and θ_r are the saturated and the residual water content respectively, n and α are the empirical van Genuchten parameters, and $m=1-1/n$. $S_e^*=[1+[\alpha h_e]^{n-1}]^{-m}$ denotes the saturation at the air entry point using the classical van Genuchten model.

This formulation was introduced to avoid an unrealistically steep decline of the hydraulic conductivity close to water saturation when using the Mualem-model to derive the hydraulic conductivity function (Mualem, 1976). It can be shown, that such an unrealistic shape of the conductivity function has to be expected for $n < 2$ and that it hampers the numerical solution of Richards' equation considerably. For a detailed discussion we refer to Ippisch et al. (2005).

The relative hydraulic conductivity is then given by

$$K_r = \begin{cases} S_e^\tau \left[\frac{1-(1-(S_e S_e^*)^{1/m})^m}{1-(1-(S_e^*)^{1/m})^m} \right]^2 & : S_e < S_e^* \\ 1 & : S_e \geq S_e^* \end{cases}, \quad (7)$$

and the hydraulic conductivity function as

$$K(S_e) = K_s K_r(S_e), \quad (8)$$

where K_s is the saturated hydraulic conductivity and τ is an empirical parameter for tortuosity.

The hydraulic parameters were measured for the A_1 and the B horizons. In contrast, no measurements were available for the compacted A_2 horizon. The texture of this horizon was comparable to the horizons above and below. To account for the increased compactness of this layer we assumed a lower hydraulic conductivity and a larger air entry value. The hydraulic parameters for all structural units are given in Table 1.

The numerical simulations were conducted with the model $\mu\phi$ using a cell-centred finite-volume scheme with full-upwinding in space and an implicit Euler scheme in time. Linearisation of the nonlinear equations is done by an incomplete Newton-Method with line search. The linear equations are solved with an algebraic multi-grid solver. For the time solver the time step is adapted automatically. A constant flux was imposed at the

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upper boundary of the domain and no-flux boundary conditions were used at the side boundaries. As there was no influence of a ground water table at this site, gravity flow was assumed at the lower boundary. After reaching a quasi-steady-state the flux field was interpolated using RT0 Raviart Thomas elements (Raviart and Thomas, 1975; Brezzi and Fortin, 1991).

Once the steady state velocity field was calculated, we simulated the infiltration of the dye tracer assuming a local convection-dispersion model. As the spatial heterogeneity of the soil was already considered explicitly in the calculation of the flux field, the dispersion coefficient D_0 was calculated as $D_0 = [\theta^{7/3} / \phi^2] D_m$ using the second model of Millington and Quirk (1960) to account for the changes in water content, where D_m is the molecular diffusion coefficient. The potential nonlinear sorption of brilliant blue was represented by a Langmuir isotherm $C_s^\infty \frac{k\theta C_w}{1+k\theta C_w}$. As the coefficients k and C_s^∞ are unknown for this soil we used the values obtained by Kasteel et al. (2002) as an orientation, which was $0.006 \text{ m}^3/\text{g}$ and 1.2 g/l , respectively. The solute transport was discretized using a second-order Godonov method (LeVeque, 2002) with a minmod slope limiter for the convective part and a finite-volume scheme for the diffusive term.

The quasi-steady-state flux field was computed with a resolution of $128 \times 128 \times 128$ on a Pentium-4 Linux workstation, the solute transport was solved with a resolution of $256 \times 256 \times 256$ in parallel on the Linux cluster Helix (with 128 dual processor nodes). To speed up the calculation of the flux field we first calculated the quasi-steady state flux field with a resolution of $64 \times 64 \times 64$ starting with an initial condition of gravity flow and used the result as initial condition of the $128 \times 128 \times 128$ flux field.

The simulated time was the same as the duration of the experiment until the infiltration was stopped and the excavation started. We did not consider the subsequent period required to prepare all 24 serial sections. There was no marked difference in depth distribution of the tracer between the first and the last serial section, so that a direct comparison of the simulated tracer distribution and the experimental data should be possible.

To evidence the effect of the mesoscopic heterogeneity, the sorption of the solute

and the boundary conditions we simulated four different scenarios:

- $SC_{\text{reference}}$ a reference scenario including mesoscopic heterogeneity and sorption of the solute,
- $SC_{\text{no-sorp}}$ the same as $SC_{\text{reference}}$ but without sorption,
- $SC_{\text{homogeneous}}$ the same as $SC_{\text{reference}}$ but without mesoscopic heterogeneity,
- $SC_{\text{low-flow}}$ the same as $SC_{\text{reference}}$ but with a flow rate of 0.13 cm/h which is one order of magnitude lower than the rate in our experiment.

6. Results and Discussion

The measured 3D distribution of dye after the experiment in comparison to the simulation ($SC_{\text{reference}}$ scenario) is illustrated in Fig. 4. For the simulation an iso-surface with a concentration of 0.1 g/l is plotted. This threshold was chosen according to the analysis of [Forrer et al. \(2000\)](#) who found a clear visibility of brilliant blue above this value. This choice is expected to depend on the type of soil; however, the measured dye distribution patterns do not change significantly if we threshold at a slightly different value. The resolution of the experimental data is limited in one horizontal direction due to the separation of the serial sections of 4 cm.

The simulated distribution of brilliant blue is much smoother than the measurements. However, it qualitatively reflects one striking characteristics of the measured data, i.e. the reduced dye coverage within the compacted part of the soil profile which increases again in the lower B horizon. This pattern is an obvious consequence of the hydraulic structure of the soil (Table 1) especially of the fact that the saturated conductivity of the compacted A_2 horizon is lower than the infiltration rate. Consequently, most of the water is directed towards the macropores to bypass the compacted layer. For this configuration the absolute value of the saturated conductivity and the other hydraulic parameters of the A_2 horizons are of minor importance because the soil matrix is close

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to saturation. So the uncertainty of these parameters as a consequence of the estimation procedure is less critical for our experiment. On the other hand the ratio between K_s and the flow rate triggers the incidence of preferential flow along the worm holes as discussed further below.

5 For a more quantitative comparison of the different scenarios and the experiment with respect to the vertical distribution, the brilliant blue coverage was averaged over the horizontal cross section of the experimental domain. The resulting depth distributions are shown in Fig. 5 together with two-dimensional vertical sections showing the coverage of brilliant blue for the different scenarios.

10 The predicted depth distribution overestimates the amount of solute in the lower B horizon and conversely, the dye concentration in the compacted A_2 horizon is underestimated. Obviously, more dye moved from the worm holes into the surrounding matrix. This is clearly due to the over simplification of the compacted horizon which is partitioned into a dense matrix and open macropores. Although, this roughly reflects the most relevant hydraulic structure of the material, the partitioning is more fuzzy in reality.

15 The overall pattern of transport is probably not sensitive to the detailed morphology of the macropores which were modelled to mimic the shape of earthworm burrows. A number of straight channels would have produced a similar pattern of solute transport in our experiment. But the more detailed morphology of macropores may be relevant for other processes where the interaction with the solid phase is stronger and hence, the surface density of macropores is more important. We think that the genetic approach used to represent the macropores is especially attractive in case the structure itself is difficult to measure directly. The example of earthworm burrows presented here is thought as a demonstration of this concept.

25 The effect of the mesoscopic heterogeneity is actually not visible in the overall phenomenology of solute transport, since the difference between $SC_{\text{homogeneous}}$ and $SC_{\text{reference}}$ is negligible. In contrast, as expected, the depth distribution is significantly affected by the non-linear sorption of brilliant blue. For $SC_{\text{no-sorp}}$ the vertical translocation of brilliant blue is markedly increased while the overall pattern of solute distribution

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is not changed.

The scenario with the low flow rate shows a completely different behaviour. During $SC_{\text{low-flow}}$ the saturated conductivity of the compacted A_2 horizon is well above the simulated flow rate and consequently, the water potential remains well below water saturation. Under such conditions the conductivity of the macropores vanishes and the flow regime turns into a homogeneous matrix flow. The $SC_{\text{low-flow}}$ scenario has the same cumulative infiltration as the other scenarios; however, the mean vertical translocation of brilliant blue is much smaller compared to the scenarios with higher flow rate. This difference is due to the non-linear sorption characteristics of the dye. Although a corresponding experiment is missing, $SC_{\text{low-flow}}$ highlights the predictive potential of the approach to model solute transport based on the structure of hydraulic properties. Using this approach, predictions for arbitrary initial and boundary conditions are possible. This is an advantage compared to the common approach of fitting an effective transport model to measured tracer distributions. Actually, there might exist some sophisticated, multi-domain model that is in the position to reproduce the measured depth distribution accurately. However, the interpretation of the resulting model parameters is in principle restricted to the specific initial and boundary conditions of the experiment. The difference between $SC_{\text{reference}}$ and $SC_{\text{low-flow}}$ is just one striking example of how transport parameters might change with changing boundary conditions.

7. Conclusions

In this study we aimed at predicting the vertical translocation of a dye tracer in a heterogeneous field soil. The prediction was based on the representation of the complete three-dimensional structure of the hydraulic properties and the resulting velocity field which was calculated using Richards' equation. To represent the structure we used the direct measurement of soil horizons, a statistical approach to estimate the sub-scale fluctuations of hydraulic properties and a genetic approach to model the formation of macroporous originating from earthworm activity.

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The agreement between our experimental findings and the prediction of tracer distribution is far from being perfect. However, the overall pattern including the infiltration depth is reasonably matched. Given the fact, that the prediction is based on a only rough approximation of the structure and the related material properties we come to the following hypothesis which potentially could be corroborated by future experiments:

- Based on the structure of hydraulic properties, the approximate pattern of solute transport including the phenomenon of preferential flow can be predicted. The macroscopic dispersion is an immediate consequence of the macroscopic structure of the material. Given this structure, no “effective” dispersion coefficient is required.
- Given the relevant and continuous structure of the material, the hydraulic properties of the single structural units may either be measured directly (here: measurements of $h(\theta)$ and $K(\theta)$ for the different horizons), or they may be derived from proxy variables (here: estimated correlation length of bulk density within the horizons), or even derived from a plausible guess (here: $h(\theta)$ and $K(\theta)$ for earthworm burrows).
- To predict transport in heterogeneous soil, quantitative information on structure including the three-dimensional topology and on the related hydraulic properties is indispensable. If both components are available a rough approximation might be sufficient (see also [Vogel and Roth, 2003](#)). If one component is missing, the estimation of transport parameters is futile.
- To represent the structure of a model domain, different approaches such as direct measurements, statistical modelling, and modelling of structure formation can be combined. The latter is especially attractive for structural components that are hardly accessible by any measurement technique.

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Table 1. Hydraulic parameters for the different structural units. For the A_1 and the B horizons they were directly measured, for the A_2 and the macropores we used a plausible guess.

	α (cm^{-1})	n (–)	θ_s (cm^3/cm^3)	θ_r (cm^3/cm^3)	K_s (cm/h)	τ (–)	h_e (cm)
A_1 horizon	33.0	1.24	0.34	0.1	18.0	0.5	10.2
A_2 horizon	0.15	1.33	0.36	0.1	0.29	0.5	44.7
macropores	0.2	8.0	0.5	0.0	3600	0.5	0.58
B horizon	0.15	1.33	0.36	0.1	2.88	0.5	10.3

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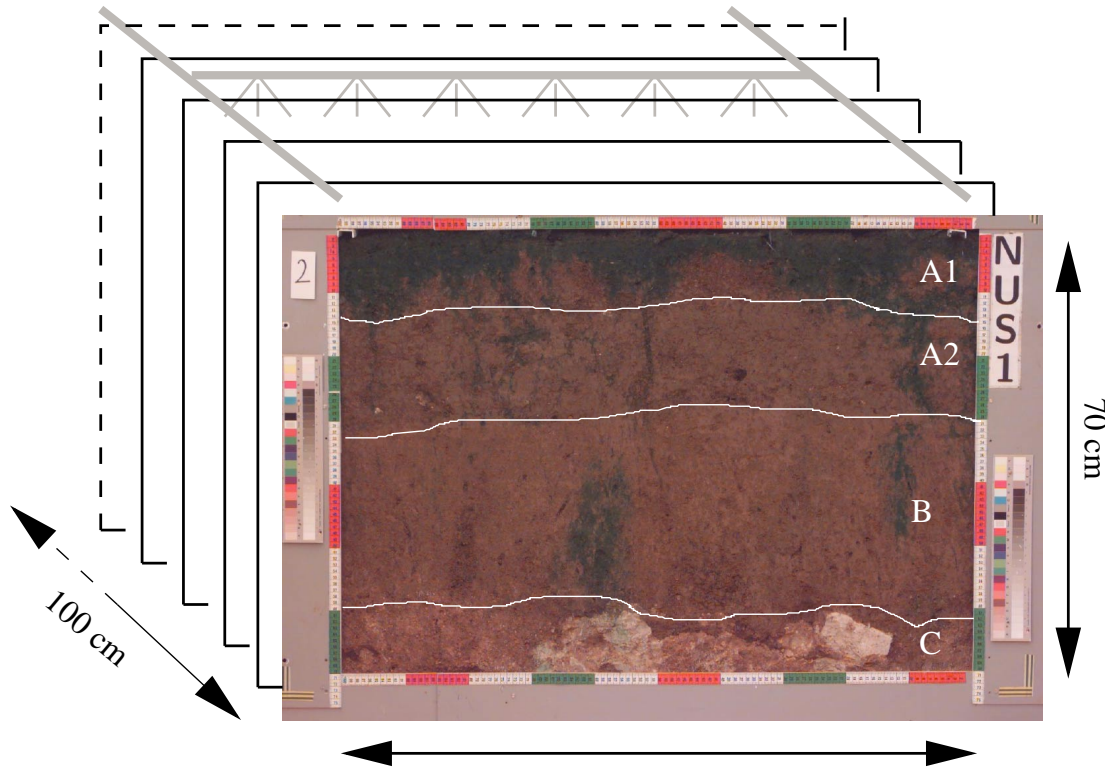


Fig. 1. Setup of the experiment: After the application of the dye tracer, the experimental domain ($1.0 \times 1.0 \times 0.7 \text{ m}^3$) was explored by 24 serial sections. Notice the irregular distribution of the dye.

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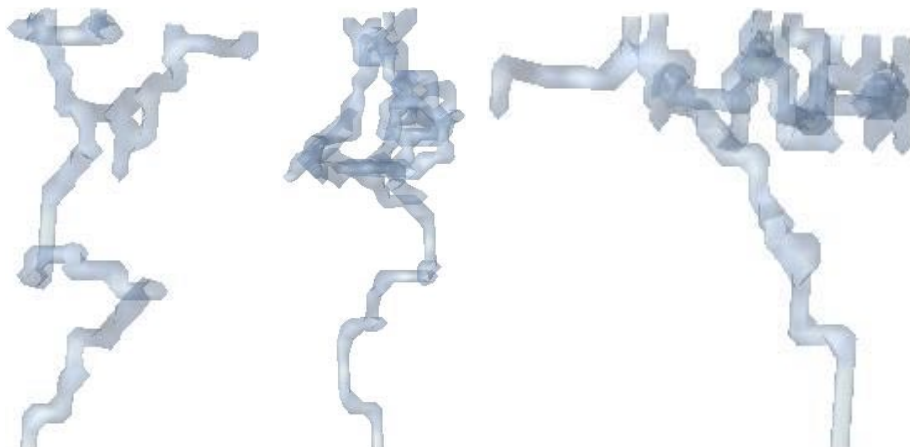


Fig. 2. Three realizations of earthworm burrows using a genetic model.

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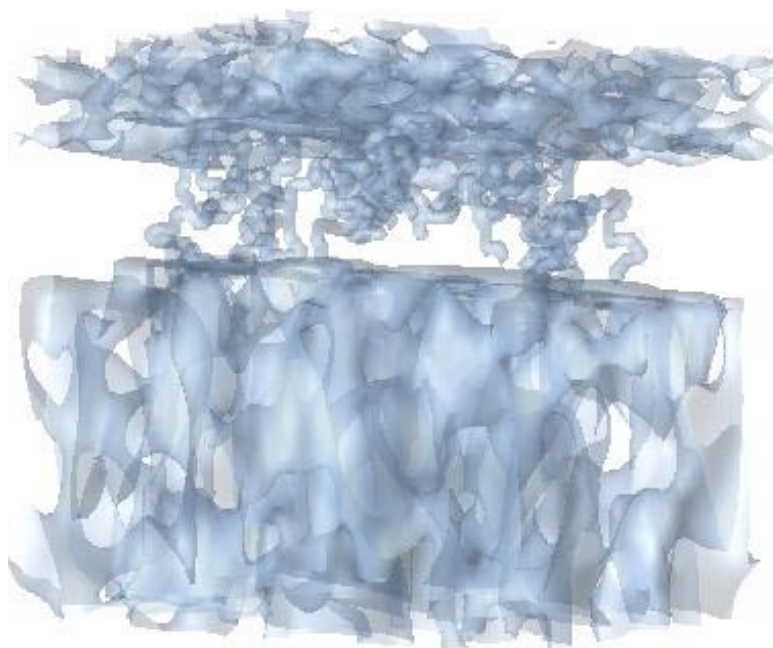


Fig. 3. Structure of the entire experimental domain (Fig. 1): Iso-surfaces of the mean log-scaling factors in the A_1 and B horizon, together with the modelled structure of macropores in the A_2 horizon.

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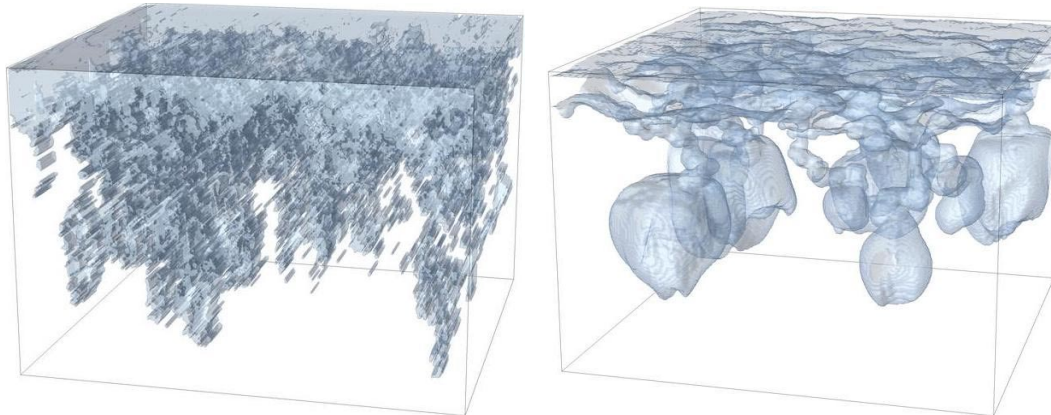


Fig. 4. Distribution of Brilliant Blue measured after the experiment (left) and simulated based on the heterogeneous structure of hydraulic properties (right) for the same initial and boundary conditions. The images cover the entire experimental domain shown in Fig. 1.

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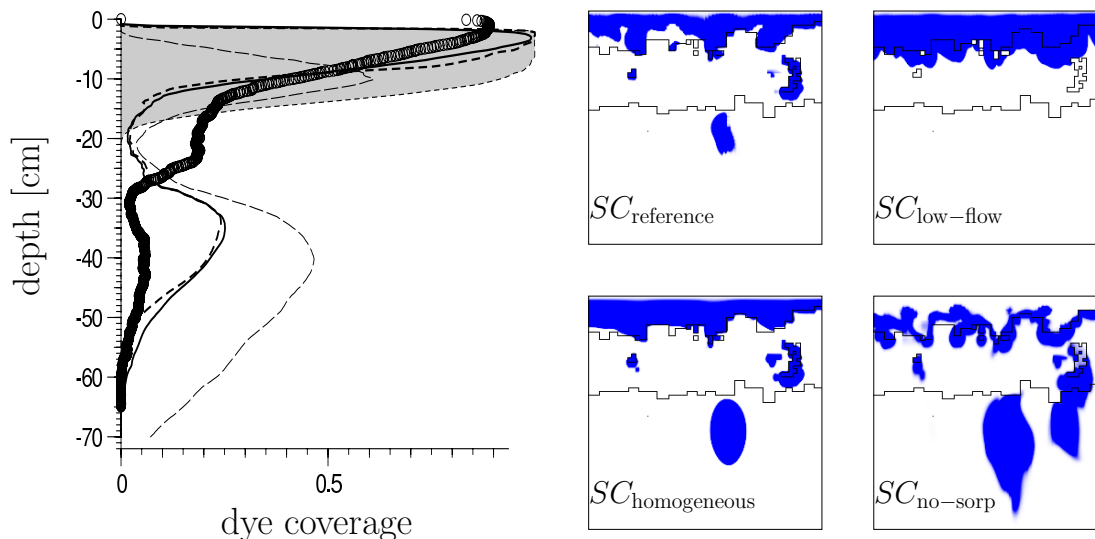


Fig. 5. Averaged vertical distribution of Brilliant Blue coverage measured after the experiment (symbols), simulated based on the heterogeneous structure of hydraulic properties including sorption ($SC_{reference}$ thick line), without mesoscopic heterogeneities ($SC_{homogeneous}$ thick dashed line), without sorption ($SC_{no-sorp}$ dashed line) and with the low flow rate but same amount of infiltrated water ($SC_{low-flow}$ grey). Notice that dye coverage is not necessarily related to total mass of dye. Right: vertical sections through the centre of the simulated domain are shown including the boundaries between the different materials (horizons and earthworm burrows). To compare with a measured section see Fig. 1.

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