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# UNIQUENESS FOR THE VORTEX-WAVE SYSTEM WHEN THE VORTICITY IS CONSTANT NEAR THE POINT VORTEX

CHRISTOPHE LACAVER & EVELYNE MIOT

ABSTRACT. We prove uniqueness for the vortex-wave system with a single point vortex introduced by Marchioro and Pulvirenti [7] in the case where the vorticity is initially constant near the point vortex. Our method relies on the Eulerian approach for this problem and in particular on the formulation in terms of the velocity.

## 1. INTRODUCTION

In this paper, we study a system occurring in two dimensional fluid dynamics. The motion of an ideal incompressible fluid in  $\mathbb{R}^2$  with divergence-free velocity field  $v = (v_1, v_2) : \mathbb{R}^+ \times \mathbb{R}^2 \rightarrow \mathbb{R}^2$  and vorticity  $\omega = \text{curl } v = \partial_2 v_1 - \partial_1 v_2 : \mathbb{R}^+ \times \mathbb{R}^2 \rightarrow \mathbb{R}$  is given by the Euler equations

$$\begin{cases} \partial_t \omega + v \cdot \nabla \omega = 0, \\ \omega = \text{curl } v, \text{ div } v = 0, \end{cases} \quad (1.1)$$

where  $\text{div } v = \partial_1 v_1 + \partial_2 v_2$ . For this system, Yudovich's Theorem states global existence and uniqueness in  $L^\infty(\mathbb{R}^+, L^1 \cap L^\infty(\mathbb{R}^2))$  for an initial vorticity  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$ . Equation (1.1) is a transport equation with field  $v$ , therefore one may solve it with the method of characteristics. When  $v$  is smooth, it gives rise to a flow defined by

$$\begin{cases} \frac{d}{dt} \phi_t(x) = v(t, \phi_t(x)) \\ \phi_0(x) = x \in \mathbb{R}^2. \end{cases} \quad (1.2)$$

In view of (1.1), we then have

$$\frac{d}{dt} \omega(t, \phi_t(x)) \equiv 0, \quad (1.3)$$

which means that  $\omega$  is constant along the characteristics. In the general case of a vorticity  $\omega \in L^\infty(\mathbb{R}^+, L^1 \cap L^\infty(\mathbb{R}^2))$ , these computations may be rigorously justified, so that the Eulerian formulation (1.1) and the Lagrangian one (1.2), (1.3) turn out to be equivalent.

Since equation (1.1) governs the evolution of the vorticity  $\omega$ , it is natural to express the velocity  $v$  in terms of  $\omega$ . This can be done by taking the orthogonal gradient in both terms in the relation  $\omega = \text{curl } v$

and using that  $v$  is divergence free. This yields  $\nabla^\perp \omega = -\Delta v$ , so that under the additional constraint that  $v$  vanishes at infinity, we have

$$v = K * \omega. \quad (1.4)$$

Here  $*$  denotes the convolution product and  $K : \mathbb{R}^2 \setminus \{0\} \rightarrow \mathbb{R}^2$  stands for the Biot-Savart Kernel defined by

$$K(x) = \frac{1}{2\pi} \frac{x^\perp}{|x|^2}, \quad x \neq 0, \quad (1.5)$$

where  $(x_1, x_2)^\perp = (-x_2, x_1)$ . When the vorticity tends to be concentrated at points, one may modify equation (1.1) according to formulas (1.4) and (1.5) into a system of ordinary differential equations, called point vortex system, which governs the motion of these points. A rigorous justification for this system has been carried out in [9]. It is proved there that if the initial vorticity  $\omega_0$  is close to the weighted sum of Dirac masses  $\sum d_i \delta_{z_i}$  in a certain sense, then  $\omega(t)$  remains close to  $\sum d_i \delta_{z_i(t)}$  for all time, where the vortices  $z_i(t)$  evolve according to the point vortex system.

In the early 90s, Marchioro and Pulvirenti [7, 8] investigated the mixed problem in which the vorticity is composed of an  $L^\infty$  part and a sum of Dirac masses. They obtained the so-called vortex-wave system, which couples the usual point vortex system and the classical Lagrangian formulation for the two-dimensional fluid dynamics. In the case of a single point vortex (which will be the case studied here), these authors obtained the global existence of solutions of the vortex-wave system in Lagrangian formulation.

**Definition 1.1** (Lagrangian solutions). *Let  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$  and  $z_0 \in \mathbb{R}^2$ . We say that the triple  $(\omega, z, \phi)$  is a global Lagrangian solution to the vortex-wave system with initial condition  $(\omega_0, z_0)$  if  $\omega \in L^\infty(\mathbb{R}^+, L^1 \cap L^\infty(\mathbb{R}^2))$ ,  $v = K * \omega \in C(\mathbb{R}^+ \times \mathbb{R}^2)$  and*

$$z : \mathbb{R}^+ \rightarrow \mathbb{R}^2, \quad \phi : \mathbb{R}^+ \times \mathbb{R}^2 \setminus \{z_0\} \rightarrow \mathbb{R}^2$$

are such that  $z \in C^1(\mathbb{R}^+, \mathbb{R}^2)$ ,  $\phi(\cdot, x) \in C^1(\mathbb{R}^+, \mathbb{R}^2)$  for all  $x \neq z_0$  and satisfy

$$\begin{cases} v(\cdot, t) = (K * \omega)(\cdot, t), \\ \dot{z}(t) = v(t, z(t)), \\ z(0) = z_0, \\ \dot{\phi}_t(x) = v(t, \phi_t(x)) + K(\phi_t(x) - z(t)), \\ \phi_0(x) = x, \quad x \neq z_0, \\ \omega(\phi_t(x), t) = \omega_0(x), \end{cases} \quad (\text{LF})$$

where  $\phi_t = \phi(t, \cdot)$ . In addition, for all  $t$ ,  $\phi_t$  is a diffeomorphism from  $\mathbb{R}^2 \setminus \{z_0\}$  into  $\mathbb{R}^2 \setminus \{z(t)\}$  that preserves Lebesgue's measure.

This system involves two kinds of trajectories. The point vortex  $z(t)$  moves under the velocity field  $v$  produced by the regular part  $\omega$  of the vorticity. This regular part and the vortex point give rise to a smooth flow  $\phi$  along which  $\omega$  is constant. The main difference with the classical Euler dynamics is the presence of the field  $K(x - z(t))$ , which is singular at the point vortex but smooth elsewhere. Marchioro and Pulvirenti [7] proved global existence for (LF). The proof mainly relies on estimates involving the distance between  $\phi_t(x)$  and  $z(t)$  and uses almost-Lipschitz regularity for  $v = K * \omega$  and the explicit form of  $K$ . It is shown in particular that a characteristic starting far apart from the point vortex cannot collide with  $z(t)$  in finite time. Consequently, the singular term  $K(\phi_t(x) - z(t))$  in (LF) remains well-defined for all time.

The notion of Lagrangian solutions is rather strong. One can define a weaker notion of solutions: solutions in the sense of distributions of the PDE (without involving the flow  $\phi$ ). We call these Eulerian solutions and we define them here below.

**Definition 1.2** (Eulerian solutions). *Let  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$  and  $z_0 \in \mathbb{R}^2$ . We say that  $(\omega, z)$  is a global Eulerian solution of the vortex-wave equation with initial condition  $(\omega_0, z_0)$  if*

$$\omega \in L^\infty(\mathbb{R}^+, L^1 \cap L^\infty(\mathbb{R}^2)), \quad z \in C(\mathbb{R}^+, \mathbb{R}^2)$$

and if we have in the sense of distributions

$$\begin{cases} \partial_t \omega + \operatorname{div}((v + H)\omega) = 0, \\ \omega(0) = \omega_0, \\ \dot{z}(t) = v(t, z(t)), \quad z(0) = z_0, \end{cases} \quad (\text{EF})$$

where  $v$  and  $H$  are given by

$$v(t, \cdot) = K *_x \omega(t), \quad H(t, \cdot) = K(\cdot - z(t)).$$

In other words, we have for any test function  $\varphi \in \mathcal{D}(\mathbb{R}^+ \times \mathbb{R}^2)$ <sup>1</sup>

$$-\int_{\mathbb{R}^2} \omega_0(x) \varphi(0, x) dx = \int_{\mathbb{R}^+} \int_{\mathbb{R}^2} \omega(\partial_t \varphi + (v + H) \cdot \nabla \varphi) ds dx,$$

and for all  $t \in \mathbb{R}^+$

$$z(t) = z_0 + \int_0^t v(s, z(s)) ds.$$

This kind of Eulerian solutions appeared for example in [4]. In that paper, a solution of the Euler equation with a fixed point-vortex is obtained as the limit of the Euler equations in the exterior of an obstacle

<sup>1</sup>By virtue of Lemma 2.2, the field defined by  $v = K * \omega$  belongs to  $L^\infty(\mathbb{R}^+ \times \mathbb{R}^2)$ . On the other hand,  $H$  belongs to  $L^1_{\text{loc}}(\mathbb{R}^+ \times \mathbb{R}^2)$ , so that this definition makes sense.

<sup>2</sup>We will see in Lemma 2.2 and Proposition 3.8 that  $v(t)$  is defined for all time and is continuous in the space variable.

that shrinks to a point. The regularity of the limit solution obtained in [4] is not better than the one given in Definition 1.2.

In this paper, we are concerned with the problems of uniqueness of Eulerian and Lagrangian solutions and with the related question of equivalence of Definitions 1.1 and 1.2.

We will first prove the following Theorem, clarifying that a Lagrangian solution is an Eulerian solution.

**Theorem 1.3.** *Let  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$  and  $z_0 \in \mathbb{R}^2$ . Let  $(\omega, z, \phi)$  be a global Lagrangian solution of the vortex-wave system with initial condition  $(\omega_0, z_0)$ . Then  $(\omega, z)$  is a global Eulerian solution.*

We turn next to our main purpose and investigate uniqueness for Lagrangian or Eulerian solutions.

Uniqueness for Lagrangian solutions can be easily achieved when the support of  $\omega_0$  does not meet  $z_0$ ; in that case, the support of  $\omega(t)$  never meets  $z(t)$  and the field  $x \mapsto K(\phi_t(x) - z(t))$  is Lipschitz on  $\text{supp } \omega_0$ .

Another situation that has been studied is the case where the vorticity is initially constant near the point vortex. Marchioro and Pulvirenti [7] suggested with some indications that uniqueness for Lagrangian solutions should hold in that situation. Starovoitov [10] proved uniqueness for Lagrangian solutions under the supplementary assumption that  $\omega_0$  is Lipschitz. In this paper, we treat the general case where the initial vorticity is constant near the point vortex  $z_0$  and belongs to  $L^1 \cap L^\infty(\mathbb{R}^2)$ . More precisely, we prove the following

**Theorem 1.4.** *Let  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$  and  $z_0 \in \mathbb{R}^2$  such that there exists  $R_0 > 0$  and  $\alpha \in \mathbb{R}$  such that*

$$\omega_0 \equiv \alpha \text{ on } B(z_0, R_0).$$

*Suppose in addition that  $\omega_0$  has compact support. Then there exists a unique Eulerian solution of the vortex-wave system with this initial data.*

In order to prove Theorem 1.4, we first show that if  $(\omega, z)$  is an Eulerian solution, then  $\omega$  is a renormalized solution in the sense of DiPerna-Lions [2] of its transport equation. This in turn implies that if the vorticity is initially constant near the point vortex, then this remains true for all time. We then take advantage of the weak formulation (EF) to derive a partial differential equation satisfied by the velocity  $v = K * \omega$ . In order to compare two solutions, one not only has to compare the two regular parts, but also possibly the diverging trajectories of the two vortices. Given two Eulerian solutions  $(\omega_1, z_1)$  and  $(\omega_2, z_2)$ , we therefore introduce the quantity

$$r(t) = |\tilde{z}(t)|^2 + \|\tilde{v}(t)\|_{L^2(\mathbb{R}^2)}$$

where  $\tilde{z} = z_1 - z_2$ ,  $\tilde{\omega} = \omega_1 - \omega_2$  and  $\tilde{v} = v_1 - v_2 = K * \tilde{\omega}$ . Since  $\tilde{\omega}$  vanishes in a neighborhood of the point vortex, the velocity  $\tilde{v}$  has to be

harmonic in this neighborhood. This provides in particular a control of its  $L^\infty$  norm (as well as the  $L^\infty$  norm for the gradient) by its  $L^2$  norm, which ultimately yields a Gronwall-type estimate for  $r(t)$  and allows to prove that it vanishes.

Clearly, since Lagrangian solutions are always Eulerian solutions and given the global existence of Lagrangian solutions proved by Marchioro and Pulvirenti [7], it follows from Theorems 1.3 and 1.4 that an Eulerian solution is also a Lagrangian solution. Therefore Definitions 1.1 and 1.2 are equivalent if the initial vorticity belongs to  $L^1 \cap L^\infty(\mathbb{R}^2)$ , is compactly supported and constant near  $z_0$ .

## 2. LAGRANGIAN IMPLIES EULERIAN

We first briefly recall some remarkable properties of the convolution by the Biot-Savart Kernel  $K$ . The proofs are standard and may be found in [5, 8]. We begin with the Calderón-Zygmund inequality.

**Lemma 2.1.** *Let  $f \in L^1 \cap L^\infty(\mathbb{R}^2)$  and  $g = K * f$  so that  $\text{curl } g = f$  and  $\text{div } g = 0$ . Then for all  $2 \leq p < +\infty$ , we have*

$$\|\nabla g\|_{L^p(\mathbb{R}^2)} \leq Cp \|f\|_{L^p(\mathbb{R}^2)},$$

where  $C$  is some universal constant.

The following Lemma will be very useful in our further analysis.

**Lemma 2.2.** *Let  $f \in L^1 \cap L^\infty(\mathbb{R}^2)$  and  $g = K * f$ . Then  $g$  satisfies*

$$\|g\|_{L^\infty} \leq C(\|f\|_{L^\infty} + \|f\|_{L^1}). \quad (2.1)$$

Moreover,

$$|g(x) - g(y)| \leq C(\|f\|_{L^\infty}, \|f\|_{L^1}) \varphi(|x - y|), \quad \forall x, y \in \mathbb{R}^2, \quad (2.2)$$

where  $\varphi$  is the continuous, concave and non-decreasing function defined by

$$\varphi(z) = \begin{cases} z(1 - \ln(z)) & \text{if } 0 \leq z < 1 \\ 1 & \text{if } z \geq 1. \end{cases}$$

From now on, we will denote by  $\mathcal{AL}$  the set of almost-Lipschitz functions from  $\mathbb{R}^2$  into  $\mathbb{R}^2$ , that is those for which

$$|g(x) - g(y)| \leq C\varphi(|x - y|), \quad x, y \in \mathbb{R}^2$$

with some constant  $C$ , and by  $L^\infty(\mathcal{AL})$  the set of functions  $v = v(t, x) : \mathbb{R} \times \mathbb{R}^2 \mapsto \mathbb{R}^2$  satisfying

$$|v(t, x) - v(t, y)| \leq C\varphi(|x - y|), \quad x, y \in \mathbb{R}^2, \quad t \in \mathbb{R}$$

for some constant  $C$  independent of  $t$ . Note that the uniform bound (2.1) holds true provided  $f$  belongs to  $L^p \cap L^q$  for some  $p < 2$  and  $q > 2$ . However, the almost-Lipschitz estimate requires the assumption  $f \in L^\infty$ .

In our situation, we will always deal with  $f = \omega(t)$  and  $v = g = K *_x \omega(t)$ , where  $\omega \in L^\infty(\mathbb{R}^+, L^1 \cap L^\infty(\mathbb{R}^2))$ . For this reason, the estimates above will actually hold uniformly with respect to time.

Finally, we define  $\chi_0 : \mathbb{R}^2 \rightarrow \mathbb{R}$  to be a smooth, *radial* cut-off map such that

$$\chi_0 \equiv 0 \text{ on } B(0, \frac{1}{2}), \quad \chi_0 \equiv 1 \text{ on } B(0, 1)^c, \quad 0 \leq \chi_0 \leq 1. \quad (2.3)$$

For a small and positive  $\delta$ , we set

$$\chi_\delta(z) = \chi_0\left(\frac{z}{\delta}\right),$$

so that as  $\delta$  goes to 0, we have

$$\chi_\delta \rightarrow 1 \text{ a.e.}, \quad \|\nabla \chi_\delta\|_{L^1(\mathbb{R}^2)} \rightarrow 0. \quad (2.4)$$

In the sequel, we denote by  $u$  the velocity field

$$u \equiv v + H.$$

It is composed of an almost-Lipschitz part  $v$  and of a part  $H$  which is singular at the point vortex  $z(t)$  and smooth outside. Clearly, multiplying any test function by  $\chi_\delta(x - z(t))$  provides a test function having compact support away from the singularity. This observation will allow us in the subsequent proofs to avoid the singularity and to first perform computations with smooth vector fields. In a second step, we will pass to the limit  $\delta \rightarrow 0$ . This will be readily achieved since we have, thanks to the explicit form of  $H$  and the fact that  $\chi_\delta$  is radial

$$H(t, x) \cdot \nabla \chi_\delta(x - z(t)) \equiv 0. \quad (2.5)$$

**Proof of Theorem 1.3.** Given  $(\omega, z, \phi)$  a solution of (LF), it actually suffices to show that

$$\partial_t \omega + \operatorname{div}((v + H)\omega) = 0 \quad (2.6)$$

in the sense of distributions on  $\mathbb{R}^+ \times \mathbb{R}^2$ .

We first give a formal proof of (2.6). Let us take a  $C^1$  function  $\psi(t, x)$  and define

$$f(t) = \int_{\mathbb{R}^2} \omega(t, y) \psi(t, y) dy.$$

We set  $y = \phi_t(x)$ . Since  $\phi_t$  preserves Lebesgue's measure for all time and since  $\omega$  is constant along the trajectories, we get

$$f(t) = \int_{\mathbb{R}^2} \omega_0(x) \psi(t, \phi_t(x)) dx.$$

Differentiating with respect to time and using the ODE solved by  $\phi_t(x)$ , this leads to

$$f'(t) = \int_{\mathbb{R}^2} \omega_0(x) (\partial_t \psi + u \cdot \nabla \psi)(t, \phi_t(x)) dx.$$

Using the change of variables  $y = \phi_t(x)$  once more yields

$$f'(t) = \int_{\mathbb{R}^2} \omega(t, y) (\partial_t \psi + u \cdot \nabla \psi)(t, y) dy,$$

which is (2.6) in the sense of distributions. In order to justify the previous computation, we need to be able to differentiate inside the integral, and we proceed as follows.

Let  $\psi = \psi(t, x)$  be any test function. For  $0 < \delta < 1$ , we set

$$\psi_\delta(t, x) \equiv \chi_\delta(x - z(t)) \psi(t, x),$$

where  $\chi_\delta$  is the map defined in (2.3). Since we have  $\psi_\delta(t) \equiv 0$  on the ball  $B(z(t), \frac{\delta}{2})$ , we may apply the previous computation to  $\psi_\delta$ , which yields for all  $t$

$$\begin{aligned} \int_{\mathbb{R}^2} \omega(t, x) \psi_\delta(t, x) dx - \int_{\mathbb{R}^2} \omega_0(x) \psi_\delta(0, x) dx \\ = \int_0^t \int_{\mathbb{R}^2} \omega (\partial_t \psi_\delta + u \cdot \nabla \psi_\delta) dx ds. \end{aligned} \quad (2.7)$$

We first observe that thanks to the pointwise convergence of  $\psi_\delta(t, \cdot)$  to  $\psi(t, \cdot)$  as  $\delta \rightarrow 0$ , we have

$$\begin{aligned} \int_{\mathbb{R}^2} \omega(t, x) \psi_\delta(t, x) dx - \int_{\mathbb{R}^2} \omega_0(x) \psi_\delta(0, x) dx \\ \rightarrow \int_{\mathbb{R}^2} \omega(t, x) \psi(t, x) dx - \int_{\mathbb{R}^2} \omega_0(x) \psi(0, x) dx \end{aligned} \quad (2.8)$$

by Lebesgue's dominated convergence Theorem. Then, we compute

$$\begin{aligned} \partial_t \psi_\delta + u \cdot \nabla \psi_\delta &= \chi_\delta(x - z) (\partial_t \psi + u \cdot \nabla \psi) \\ &\quad + \psi(-\dot{z} + v + H) \cdot \nabla \chi_\delta(x - z). \end{aligned}$$

Using (2.5) and that  $v$  is uniformly bounded, we obtain

$$\begin{aligned} \left| \int_{\mathbb{R}^2} \omega [\partial_t \psi_\delta + u \cdot \nabla \psi_\delta - \chi_\delta (\partial_t \psi + u \cdot \nabla \psi)] dx \right| \\ \leq C \|\psi\|_{L^\infty} \|v\|_{L^\infty} \|\omega\|_{L^\infty} \int_{\mathbb{R}^2} |\nabla \chi_\delta|(x) dx. \end{aligned} \quad (2.9)$$

We now let  $\delta$  tend to zero. We observe that

$$\begin{aligned} \int_0^t \int_{\mathbb{R}^2} \omega \psi_\delta (\partial_t \psi + u \cdot \nabla \psi) dx ds \\ \rightarrow \int_0^t \int_{\mathbb{R}^2} \omega (\partial_t \psi + u \cdot \nabla \psi) dx ds, \end{aligned}$$

so that the conclusion finally follows from (2.4), (2.7), (2.8) and (2.9).



### 3. UNIQUENESS OF EULERIAN SOLUTIONS

This section is devoted to the proof of Theorem 1.4. A first step in this direction consists in proving that if the vorticity is initially constant near the point vortex, this remains true for all time. This is proved in [7] for any Lagrangian solution by estimating the distance between the flow and the point vortex. In the present situation where Eulerian solutions are considered, this is achieved by proving that the vorticity of an Eulerian solution is a renormalized solution in the sense of DiPerna and Lions [2] of its transport equation.

We recall from Lemmas 2.1 and 2.2 that if  $(\omega, z)$  is an Eulerian solution of (EF), then the velocity field defined by  $v = K * \omega$  satisfies

$$v \in L^\infty(\mathbb{R}^+ \times \mathbb{R}^2) \cap L^\infty(\mathbb{R}^+, W_{\text{loc}}^{1,1}(\mathbb{R}^2)) \cap L^\infty(\mathcal{AL}).$$

**3.1. Renormalized solutions.** In what follows, we consider equation (EF) as a linear transport equation with given velocity field  $u = v + H$  and trajectory  $z$ . Our purpose here is to show that if  $\omega$  solves this linear equation, then so does  $\beta(\omega)$  for a suitable smooth function  $\beta$ . When there is no point vortex, this directly follows from the theory developed in [2] (see also [1] for more details). The results stated in [2] hold for velocity fields having enough Sobolev regularity; a typical relevant space is  $L_{\text{loc}}^1(\mathbb{R}^+, W_{\text{loc}}^{1,1}(\mathbb{R}^2))$ . These results can actually be extended to our present situation, thanks to the regularity of  $H$  away from the point vortex and to its special form.

We define

$$\Sigma = \{(t, z(t)), t \in \mathbb{R}^+\}$$

and denote by  $\mathcal{G}$  its complement in  $\mathbb{R}^+ \times \mathbb{R}^2$ .

The starting point in [2] and [1] is to look at mollifiers

$$\omega_\varepsilon = \rho_\varepsilon *_x \omega, \quad \omega_{\varepsilon,\eta} = \omega_\varepsilon *_t \theta_\eta,$$

where  $\rho_\varepsilon$  and  $\theta_\eta$  are standard regularizing kernels on  $\mathbb{R}^2$  and  $\mathbb{R}^+$  respectively. We also set, for  $f \in L_{\text{loc}}^1(\mathbb{R}^+ \times \mathbb{R}^2)$ ,  $f_\eta = f *_x \rho_\eta *_t \theta_\eta$ . The following Lemma is a direct consequence of the Sobolev regularity of  $v$  and the regularity of  $H$  in  $\mathcal{G}$ .

**Lemma 3.1** (Commutators). *Let  $(\omega, z)$  be an Eulerian solution of (EF). Then we have*

$$\partial_t \omega_{\varepsilon,\eta} + u_\eta \cdot \nabla \omega_{\varepsilon,\eta} = r_{\varepsilon,\eta} \tag{3.1}$$

in the sense of distributions, where the remainder  $r_{\varepsilon,\eta}$  is defined by

$$r_{\varepsilon,\eta} = u_\eta \cdot \nabla \omega_{\varepsilon,\eta} - (u \cdot \nabla \omega)_{\varepsilon,\eta}$$

and satisfies

$$\lim_{\varepsilon \rightarrow 0} \left( \lim_{\eta \rightarrow 0} r_{\varepsilon,\eta} \right) = 0 \quad \text{in } L_{\text{loc}}^1(\mathcal{G}).$$

*Proof.* Let  $K$  be a compact subset of  $\mathcal{G}$ . Then there exists  $a > 0$  such that

$$|x - z(\tau)| \geq a, \quad \forall (\tau, x) \in K.$$

We set  $\chi(t, x) = \chi_0\left(\frac{x-z(t)}{a/2}\right)$ , where  $\chi_0$  is defined in (2.3). Clearly, we have for  $\eta$  and  $\varepsilon$  sufficiently small with respect to  $a$  and for  $(x, \tau) \in K$

$$r_{\varepsilon, \eta}(\tau, x) = (u\chi)_\eta \cdot \nabla \omega_{\varepsilon, \eta} - ((u\chi) \cdot \nabla \omega)_{\varepsilon, \eta}.$$

Firstly, the velocity  $v$ , and hence  $v\chi$  belongs to  $L_{\text{loc}}^\infty(\mathbb{R}^+, W_{\text{loc}}^{1,1}(\mathbb{R}^2))$ . Secondly, thanks to the equality

$$|H(\tau, x) - H(\tau, y)| = \frac{|x - y|}{|x - z(\tau)||y - z(\tau)|},$$

we infer that  $H\chi \in L_{\text{loc}}^\infty(\mathbb{R}^+, W_{\text{loc}}^{1,1}(\mathbb{R}^2))$ . Invoking Lemma II.1 in [2] or Lemma 1 in [1], we obtain

$$\lim_{\varepsilon \rightarrow 0} \left( \lim_{\eta \rightarrow 0} r_{\varepsilon, \eta} \right) = 0 \quad \text{in } L_{\text{loc}}^1(K).$$

The Lemma is proved.  $\square$

The second step is to use the explicit form of  $H$ .

**Lemma 3.2.** *Let  $(\omega, z)$  be a solution of (EF). Let  $\beta : \mathbb{R} \rightarrow \mathbb{R}$  be a smooth function such that*

$$|\beta'(t)| \leq C(1 + |t|^p), \quad \forall t \in \mathbb{R},$$

for some  $p \geq 0$ . Then for all test function  $\psi \in \mathcal{D}(\mathbb{R}^+ \times \mathbb{R}^2)$ , we have

$$\frac{d}{dt} \int_{\mathbb{R}^2} \psi \beta(\omega) dx = \int_{\mathbb{R}^2} \beta(\omega) (\partial_t \psi + u \cdot \nabla \psi) dx \text{ in } L_{\text{loc}}^1(\mathbb{R}^+).$$

*Proof.* We consider the mollifier  $\omega_{\varepsilon, \eta}$  defined above. Since  $\omega_{\varepsilon, \eta}$  is smooth, equality (3.1) actually holds almost everywhere in  $\mathbb{R}^+ \times \mathbb{R}^2$ . Multiplying (3.1) by  $\beta'(\omega_{\varepsilon, \eta})$  yields

$$\partial_t \beta(\omega_{\varepsilon, \eta}) + u_\eta \cdot \nabla \beta(\omega_{\varepsilon, \eta}) = \beta'(\omega_{\varepsilon, \eta}) r_{\varepsilon, \eta} \quad \text{a. e. in } \mathbb{R}^2. \quad (3.2)$$

We proceed now as in the proof of Theorem 1.3. Let  $\psi \in \mathcal{D}(\mathbb{R}^+ \times \mathbb{R}^2)$  be any test function. Let  $\delta > 0$  denote a small parameter and let  $\chi_\delta$  be the radial, smooth map on  $\mathbb{R}^2$  defined in (2.3).

We next set

$$\psi_\delta(t, x) = \chi_\delta(x - z(t)) \psi(t, x),$$

and

$$\psi_{\delta, n}(t, x) = \chi_\delta(x - z^n(t)) \psi(t, x),$$

where  $z^n(t)$  is a smooth approximation of  $z(t)$ . The functions  $\psi_\delta$  and  $\psi_{\delta, n}$  are supported in  $\mathcal{G}$ . Note that since  $z$  is Lipschitz with Lipschitz constant given by  $\|v\|_{L^\infty(\mathbb{R}^+ \times \mathbb{R}^2)}$ , we may choose  $z^n(t)$  so that

$$\sup_{t \in \mathbb{R}^+} |z(t) - z^n(t)| \leq \frac{1}{n} \|v\|_{L^\infty}, \quad \sup_{t \in \mathbb{R}^+} |\dot{z}^n(t)| \leq \|v\|_{L^\infty}. \quad (3.3)$$

Multiplying (3.2) by  $\psi_{\delta,n}$  and integrating in space gives for all  $t$

$$\begin{aligned} \frac{d}{dt} \int_{\mathbb{R}^2} \psi_{\delta,n}(t) \beta(\omega_{\varepsilon,\eta}(t)) dx &= \int_{\mathbb{R}^2} \beta'(\omega_{\varepsilon,\eta}) r_{\varepsilon,\eta} \psi_{\delta,n} dx \\ &+ \int_{\mathbb{R}^2} \beta(\omega_{\varepsilon,\eta}) (\partial_t \psi_{\delta,n} + u_\eta \cdot \nabla \psi_{\delta,n}) dx d\tau. \end{aligned}$$

Given  $\delta > 0$ , we find  $n$  sufficiently large so that  $\frac{1}{n} \|v\|_{L^\infty}$  is small with respect to  $\delta$ . We infer from the definition of  $\chi_\delta$  and (3.3) that  $\psi_{\delta,n}$  is compactly supported in  $\mathcal{G}$ . We may thus invoke Lemma 3.1, the assumption on  $\beta$  and uniform  $L^\infty$  bounds for  $\omega_{\varepsilon,\eta}$  to deduce that for fixed  $\delta$  and  $n$

$$\lim_{\varepsilon \rightarrow 0} \left( \lim_{\eta \rightarrow 0} \int_{\mathbb{R}^2} \beta'(\omega_{\varepsilon,\eta}) r_{\varepsilon,\eta} \psi_{\delta,n} dx \right) = 0 \quad \text{in } L^1_{\text{loc}}(\mathbb{R}^+). \quad (3.4)$$

Besides,

$$\lim_{\varepsilon \rightarrow 0} \left( \lim_{\eta \rightarrow 0} \|\omega_{\varepsilon,\eta} - \omega\|_{L^1_{\text{loc}}(\mathbb{R}^+, L^1(\mathbb{R}^2))} \right) = 0,$$

so that using the uniform bounds on  $\partial_t \psi_{\delta,n} + u_\eta \cdot \nabla \psi_{\delta,n}$  with respect to  $\eta, \varepsilon$ , we are led to

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \left( \lim_{\eta \rightarrow 0} \int_{\mathbb{R}^2} \beta(\omega_{\varepsilon,\eta}) (\partial_t \psi_{\delta,n} + u_\eta \cdot \nabla \psi_{\delta,n}) dx d\tau \right) \\ = \int_{\mathbb{R}^2} \beta(\omega) (\partial_t \psi_{\delta,n} + u \cdot \nabla \psi_{\delta,n}) dx d\tau \quad \text{in } L^1_{\text{loc}}(\mathbb{R}^+). \end{aligned} \quad (3.5)$$

Finally, since

$$\lim_{\eta, \varepsilon \rightarrow 0} \frac{d}{dt} \int_{\mathbb{R}^2} \beta(\omega_{\varepsilon,\eta}) \psi_{\delta,n} dx = \frac{d}{dt} \int_{\mathbb{R}^2} \beta(\omega) \psi_{\delta,n} dx$$

in the sense of distributions on  $\mathbb{R}^+$ , we infer from (3.4) and (3.5)

$$\frac{d}{dt} \int_{\mathbb{R}^2} \beta(\omega) \psi_{\delta,n} dx = \int_{\mathbb{R}^2} \beta(\omega) (\partial_t \psi_{\delta,n} + u \cdot \nabla \psi_{\delta,n}) dx \quad \text{in } \mathcal{D}'(\mathbb{R}^+). \quad (3.6)$$

On the other hand, we compute

$$\begin{aligned} \partial_t \psi_{\delta,n} + u \cdot \nabla \psi_{\delta,n} &= \chi_\delta(x - z^n) (\partial_t \psi + u \cdot \nabla \psi) \\ &+ \psi (v - z^n + H) \cdot \nabla \chi_\delta(x - z^n). \end{aligned}$$

We first let  $n$  go to  $+\infty$ . Since  $\chi_\delta$  is radially symmetric, we have

$$H \cdot \nabla \chi_\delta(x - z^n) \rightarrow H \cdot \nabla \chi_\delta(x - z) \equiv 0 \quad \text{in } L^1_{\text{loc}}(\mathbb{R}^+ \times \mathbb{R}^2).$$

Thanks to the pointwise convergence of  $\psi_{\delta,n}$  as  $n$  goes to  $+\infty$  to  $\psi_\delta$  and to the uniform  $L^\infty$  bounds for the velocity and the vorticity, we

deduce

$$\begin{aligned} \left| \int_{\mathbb{R}^2} \beta(\omega)(\partial_t \psi_\delta + u \cdot \nabla \psi_\delta) dx - \int_{\mathbb{R}^2} \beta(\omega) \chi_\delta(x-z)(\partial_t \psi + u \cdot \nabla \psi) dx \right| \\ \leq C \int_{\mathbb{R}^2} |\nabla \chi_\delta| dx. \end{aligned}$$

Letting  $\delta$  go to zero and using (2.4) and (3.6) yields

$$\frac{d}{dt} \int_{\mathbb{R}^2} \beta(\omega) \psi(t, x) dx = \int_{\mathbb{R}^2} \beta(\omega) (\partial_t \psi + u \cdot \nabla \psi) dx$$

in the sense of distributions on  $\mathbb{R}^+$ . Since the right-hand side in the previous equality belongs to  $L^1_{\text{loc}}(\mathbb{R}^+)$ , the equality holds in  $L^1_{\text{loc}}(\mathbb{R}^+)$  and the Lemma is proved.  $\square$

*Remark 3.3.* (1) Lemma 3.2 actually still holds when  $\psi$  is smooth, bounded and has bounded first derivatives in time and space. In this case, we have to consider smooth functions  $\beta$  which in addition satisfy  $\beta(0) = 0$ , so that  $\beta(\omega)$  is integrable. This may be proved by approximating  $\psi$  by smooth and compactly supported functions  $\psi_n$  for which Lemma 3.2 applies, and by letting then  $n$  go to  $+\infty$ .

(2) We let  $1 \leq p < +\infty$ . Approximating  $\beta(t) = |t|^p$  by smooth functions and choosing  $\psi \equiv 1$  in Lemma 3.2, we deduce that for an Eulerian solution  $\omega$  to (EF), the maps  $t \mapsto \|\omega(t)\|_{L^p(\mathbb{R}^2)}$  are continuous and constant. In particular, we have

$$\|\omega(t)\|_{L^1(\mathbb{R}^2)} + \|\omega(t)\|_{L^\infty(\mathbb{R}^2)} \equiv \|\omega_0\|_{L^1(\mathbb{R}^2)} + \|\omega_0\|_{L^\infty(\mathbb{R}^2)},$$

and we denote by  $\|\omega_0\|$  this last quantity.

**3.2. Conservation of the vorticity near the point vortex.** Specifying our choice for  $\beta$  in Lemma 3.2, we are led to the following

**Proposition 3.4.** *Let  $(\omega, z)$  be an Eulerian solution of (EF) such that*

$$\omega_0 \equiv \alpha \quad \text{on } B(z_0, R_0)$$

*for some positive  $R_0$ . Then there exists a continuous and positive function  $t \mapsto R(t)$  depending only on  $t, R_0$  and  $\|\omega_0\|$  such that  $R(0) = R_0$  and*

$$\forall t \in \mathbb{R}^+, \quad \omega(t) \equiv \alpha \quad \text{on } B(z(t), R(t)).$$

*Proof.* We set  $\beta(t) = (t - \alpha)^2$  and use Lemma 3.2 with this choice. Let  $\Phi \in \mathcal{D}(\mathbb{R}^+ \times \mathbb{R}^2)$ . We claim that for all  $T$

$$\begin{aligned} \int_{\mathbb{R}^2} \Phi(T, x) (\omega - \alpha)^2(T, x) dx - \int_{\mathbb{R}^2} \Phi(0, x) (\omega - \alpha)^2(0, x) dx \\ = \int_0^T \int_{\mathbb{R}^2} (\omega - \alpha)^2 (\partial_t \Phi + u \cdot \nabla \Phi) dx dt. \end{aligned}$$

This is actually an improvement of Lemma 3.2, in which the equality holds in  $L^1_{\text{loc}}(\mathbb{R}^+)$ . Indeed, we have  $\partial_t \omega = -\text{div}(u\omega)$  (in the sense of

distributions) with  $\omega \in L^\infty$  and  $u \in L^\infty(\mathbb{R}^+, L_{\text{loc}}^q(\mathbb{R}^2))$  for all  $q < 2$ , which implies that  $\partial_t \omega$  is bounded in  $L_{\text{loc}}^1(\mathbb{R}^+, W_{\text{loc}}^{-1,q}(\mathbb{R}^2))$ . Hence,  $\omega$  belongs to  $C(\mathbb{R}^+, W_{\text{loc}}^{-1,q}(\mathbb{R}^2)) \subset C(\mathbb{R}^+, L_{\text{loc}-w}^2)$ . Since on the other hand  $t \mapsto \|\omega(t)\|_{L^2}$  is continuous by Remark 3.3, we have  $\omega \in C(\mathbb{R}^+, L^2(\mathbb{R}^2))$ , and the previous integral equality holds for all  $T$ .

Now, we choose a test function  $\Phi$  centered at  $z(t)$ . More precisely, we let  $\varphi_0$  be a non-increasing function on  $\mathbb{R}$ , which is equal to 1 for  $s \leq 1/2$  and vanishes for  $s \geq 1$  and we set  $\Phi(t, x) = \varphi_0(|x - z(t)|/R(t))$ , with  $R(t)$  a smooth, positive and decreasing function to be determined later on, such that  $R(0) = R_0$ . We should regularize  $z$  as in the proof of Lemma 3.2 to ensure enough regularity for  $\Phi$ , but we omit the details here for the sake of clarity. For this choice of  $\Phi$ , we have  $(\omega_0(x) - \alpha)^2 \Phi(0, x) \equiv 0$ .

We compute then

$$\nabla \Phi = \frac{x - z}{|x - z|} \frac{\varphi_0'}{R(t)}$$

and

$$\partial_t \Phi = -\frac{R'(t)}{R^2(t)} |x - z| \varphi_0' + \frac{\dot{z} \cdot (z - x)}{|x - z|} \frac{\varphi_0'}{R(t)}.$$

Since  $u \cdot \nabla \Phi = (v + H) \cdot \nabla \Phi = v \cdot \nabla \Phi$ , we obtain

$$\begin{aligned} & \int_{\mathbb{R}^2} \Phi(T, x) (\omega - \alpha)^2(T, x) dx \\ &= \int_0^T \int_{\mathbb{R}^2} (\omega - \alpha)^2 \frac{\varphi_0'(|x-z|/R)}{R} \left( (v(x) - \dot{z}) \cdot \frac{(x-z)}{|x-z|} - \frac{R'}{R} |x-z| \right) dx dt. \end{aligned} \quad (3.7)$$

Without loss of generality, we may assume that  $R_0 \leq 1$ , so that  $R \leq 1$ . As  $\varphi_0'(\frac{|x-z|}{R})$  is negative for  $R/2 \leq |x-z| \leq R$  and vanishes elsewhere and  $R' < 0$ , we can estimate the right-hand side term of (3.7) by:

$$\begin{aligned} & \int_0^T \int_{\mathbb{R}^2} (\omega - \alpha)^2 \frac{\varphi_0'(|x-z|/R)}{R} \left( (v(x) - \dot{z}) \cdot \frac{(x-z)}{|x-z|} - \frac{R'}{R} |x-z| \right) dx dt \\ & \leq \int_0^T \int_{\mathbb{R}^2} (\omega - \alpha)^2 \frac{|\varphi_0'|(\frac{|x-z|}{R})}{R} \left( |v(x) - v(z)| + \frac{R'}{2} \right) dx dt. \end{aligned}$$

Using that  $v \in L^\infty(\mathcal{AL})$  and recalling that  $\varphi$  is non-decreasing (see Lemma 2.2), we deduce from (3.7)

$$\begin{aligned} & \int_{\mathbb{R}^2} \Phi(T, x) (\omega - \alpha)^2(T, x) dx \\ & \leq \int_0^T \int_{\mathbb{R}^2} (\omega - \alpha)^2 \frac{|\varphi_0'|(\frac{|x-z|}{R})}{R} \left( C\varphi(|x-z|) + \frac{R'}{2} \right) dx dt \\ & \leq \int_0^T \int_{\mathbb{R}^2} (\omega - \alpha)^2 \frac{|\varphi_0'|(\frac{|x-z|}{R})}{R} \left( CR(1 - \ln R) + \frac{R'}{2} \right) dx dt, \end{aligned}$$

where  $C$  only depends on  $\|\omega_0\|$ . Taking  $R(t) = \exp(1 - (1 - \ln R_0)e^{2Ct})$ , we arrive at

$$\int_{\mathbb{R}^2} \Phi(T, x) (\omega - \alpha)^2(T, x) dx \leq 0,$$

which ends the proof.  $\square$

Proposition 3.4 provides the following

**Corollary 3.5.** *Let  $(\omega_1, z_1)$  and  $(\omega_2, z_2)$  be two Eulerian solutions to (EF) starting from  $(\omega_0, z_0)$ . Assume in addition that*

$$\omega_0 \equiv \alpha \quad \text{on } B(z_0, R_0).$$

*Let  $T^* > 0$  be fixed. Then there exists a time  $T_C \leq T^*$  depending only on  $T^*$ ,  $\|\omega_0\|$  and  $R_0$  such that*

$$\omega_1(t) \equiv \omega_2(t) = \alpha \quad \text{on } B\left(z(t), \frac{R(t)}{2}\right), \quad \forall t \in [0, T_C],$$

*where  $z(t)$  is the middle point of  $[z_1(t), z_2(t)]$ . Moreover, we have*

$$z_1(t), z_2(t) \in B\left(z(t), \frac{R(t)}{8}\right).$$

*Proof.* Let us define  $R := \min_{t \in [0, T^*]} R(t) > 0$ , where  $R(t)$  is given in Proposition 3.4. Since  $|z_1(0) - z_2(0)| = 0$  and  $\|v_1\|_{L^\infty}, \|v_2\|_{L^\infty} \leq C\|\omega_0\|$ , we have

$$|z_1(t) - z_2(t)| \leq 2C\|\omega_0\|t \leq \frac{R}{4}, \quad \forall t \in [0, T_C],$$

where  $T_C = \min(R(8C\|\omega_0\|)^{-1}, T^*)$ . Hence, we get

$$|z_1(t) - z_2(t)| \leq \frac{R(t)}{4}, \quad \forall t \in [0, T_C],$$

and this yields

$$B\left(z(t), \frac{R(t)}{2}\right) \subset B(z_1(t), R(t)) \cap B(z_2(t), R(t)).$$

The conclusion follows from Proposition 3.4.  $\square$

*Remark 3.6.* We assume that  $\omega_0$  has compact support. Considering  $\beta(t) = t^2$  in Lemma 3.2 and adapting the proof of Proposition 3.4, we obtain that  $\omega(t)$  remains compactly supported and its support grows at most linearly. Indeed, if we choose  $\Phi(t, x) = 1 - \varphi_0(|x - z(t)|/R(t))$ , with  $R(t)$  a smooth, positive and increasing function such that  $R(0) = R_1$ , where  $\text{supp } \omega_0 \subset B(z_0, R_1)$ , then (3.7) becomes

$$\begin{aligned} & \int_{\mathbb{R}^2} \Phi(T, x) \omega^2(T, x) dx \\ &= \int_0^T \int_{\mathbb{R}^2} \omega^2 \frac{-\varphi'_0\left(\frac{|x-z|}{R}\right)}{R} \left( (v(x) - z) \cdot \frac{(x-z)}{|x-z|} - \frac{R'}{R}|x-z| \right) dx dt \\ &\leq \int_0^T \int_{\mathbb{R}^2} \omega^2 \frac{|\varphi'_0\left(\frac{|x-z|}{R}\right)|}{R} \left( 2C - \frac{R'}{2} \right) dx dt, \end{aligned}$$

where  $C$  depends only on  $\|\omega_0\|$ . The right-hand side is identically zero for  $R(t) = R_1 + 4Ct$ , and we conclude that  $\text{supp}(\omega(t)) \in B(0, R(t))$ .

**3.3. Weak formulation for the velocity.** We now turn to the equation satisfied by the velocity  $v$  for an Eulerian solution  $(\omega, z)$  of (EF). This equation is established in [4] in the situation where the point vortex is fixed at the origin. It can be easily extended to our case, and we obtain the following

**Proposition 3.7.** *Let  $(\omega, z)$  be a global solution to (EF) with initial condition  $(\omega_0, z_0)$ . Then we have in the sense of distributions on  $\mathbb{R}^+ \times \mathbb{R}^2$* <sup>3</sup>

$$\begin{cases} \partial_t v + v \cdot \nabla v + \text{div}(v \otimes H + H \otimes v) - v(z(t))^\perp \delta_{z(t)} = -\nabla p \\ \text{div } v = 0 \\ \dot{z}(t) = v(t, z(t)) \\ v(x, 0) = K * \omega_0 \text{ and } z(0) = z_0, \end{cases}$$

where  $\delta_{z(t)}$  is the Dirac mass centered at  $z(t)$  and  $H(t, x) \equiv K(x - z(t))$ .

In the sequel, we will denote by  $W_\sigma^{1,4}(\mathbb{R}^2)$  the set of functions belonging to  $W^{1,4}(\mathbb{R}^2)$  and which are divergence-free in the sense of distributions, and by  $W_\sigma^{-1,4/3}(\mathbb{R}^2)$  its dual space.

Given two solutions  $(\omega_1, z_1)$  and  $(\omega_2, z_2)$  of (EF), we define  $\tilde{v} = K * (\omega_1 - \omega_2) = v_1 - v_2$ . As a consequence of Proposition 3.7, we obtain the following properties for  $\tilde{v}$ .

**Proposition 3.8.** *Let  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$  be compactly supported,  $z_0 \in \mathbb{R}^2$  and  $(\omega_1, z_1), (\omega_2, z_2)$  be two Eulerian solutions of (EF) with initial condition  $(\omega_0, z_0)$ . Let  $\tilde{v} = v_1 - v_2$ . Then we have*

$$\tilde{v} \in L_{\text{loc}}^2(\mathbb{R}^+, W_\sigma^{1,4}(\mathbb{R}^2)), \quad \partial_t \tilde{v} \in L_{\text{loc}}^2(\mathbb{R}^+, W_\sigma^{-1,4/3}(\mathbb{R}^2)).$$

In addition, we have  $\tilde{v} \in C(\mathbb{R}^+, L^2(\mathbb{R}^2))$  and for all  $T \in \mathbb{R}^+$ ,

$$\|\tilde{v}(T)\|_{L^2(\mathbb{R}^2)}^2 = 2 \int_0^T \langle \partial_t \tilde{v}, \tilde{v} \rangle_{W_\sigma^{-1,4/3}, W_\sigma^{1,4}} ds, \quad \forall T \in \mathbb{R}^+.$$

*Proof.* We define  $\tilde{\omega} = \omega_1 - \omega_2$ , so that  $\tilde{v} = K * \tilde{\omega}$  and we have for all  $t$

$$\int_{\mathbb{R}^2} \omega_1(t, x) dx \equiv \int_{\mathbb{R}^2} \omega_0(x) dx \equiv \int_{\mathbb{R}^2} \omega_2(t, x) dx.$$

To see this, we may for instance choose  $\beta(t) \equiv t$  and  $\psi \equiv 1$  in Lemma 3.2. Therefore,  $\int \tilde{\omega}(t) \equiv 0$ . On the other hand,  $\omega_1$  and  $\omega_2$  are compactly supported in view of Remark 3.6, so we first infer that  $\tilde{v}(t) \in L^2(\mathbb{R}^2)$  for

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<sup>3</sup>In view of the equation satisfied by  $v$ , we have  $v \in C(\mathbb{R}^+, W_\sigma^{-1,4/3}(\mathbb{R}^2))$  (see the proof of Proposition 3.8), and therefore  $v(0)$  is well-defined.

all  $t$  (see [5] for more details). Using that  $\|\omega_i\|_{L^1(\mathbb{R}^2) \cap L^\infty(\mathbb{R}^2)} \in L^\infty(\mathbb{R}^+)$ , we even obtain

$$\tilde{v} \in L_{\text{loc}}^\infty(\mathbb{R}^+, L^2(\mathbb{R}^2)). \quad (3.8)$$

We now turn to the first assertion in Proposition 3.8. We apply Proposition 3.7 to  $(v_i, z_i)$  for  $i = 1$  and  $i = 2$ . First, we infer from Lemmas 2.1 and 2.2 that  $v_i = K * \omega_i$  belongs to  $L^\infty(\mathbb{R}^+ \times \mathbb{R}^2)$  and its gradient  $\nabla v_i$  to  $L^\infty(\mathbb{R}^+, L^4(\mathbb{R}^2))$ . On the other hand, since the vorticity  $\omega_i$  is compactly supported, we have for large  $|x|$

$$|v_i(t, x)| \leq \frac{C}{|x|} \int_{\mathbb{R}^2} |\omega_i(t, y)| dy,$$

hence  $v_i$  belongs to  $L_{\text{loc}}^\infty(\mathbb{R}^+, L^p(\mathbb{R}^2))$  for all  $p > 2$ . It follows in particular that

$$v_i \in L_{\text{loc}}^\infty(\mathbb{R}^+, W^{1,4}(\mathbb{R}^2))$$

and also that  $v_i \otimes v_i$  belongs to  $L_{\text{loc}}^\infty(L^{4/3})$ . Since  $v_i$  is divergence-free, we have  $v_i \cdot \nabla v_i = \text{div}(v_i \otimes v_i)$ , and so  $v_i \cdot \nabla v_i \in L_{\text{loc}}^2(\mathbb{R}^+, W^{-1, \frac{4}{3}}(\mathbb{R}^2))$ . Moreover,  $v_i(t) \otimes H_i(t)$  belongs to  $L_{\text{loc}}^{4/3}$ , whereas at infinity,  $H_i$  and  $v_i$  are bounded by  $C/|x|$  which belongs to  $L^{8/3}$ . This yields

$$\text{div}(v_i \otimes H_i), \text{div}(H_i \otimes v_i) \in L_{\text{loc}}^2(\mathbb{R}^+, W^{-1, \frac{4}{3}}(\mathbb{R}^2)).$$

Besides, we deduce from the embedding of  $W^{1,4}(\mathbb{R}^2)$  in  $C_0^0(\mathbb{R}^2)$  that  $\delta_{z_i}$  belongs to  $L_{\text{loc}}^2(W^{-1, \frac{4}{3}})$ . Therefore,  $v_i \delta_{z_i} \in L_{\text{loc}}^2(\mathbb{R}^+, W^{-1, \frac{4}{3}}(\mathbb{R}^2))$ .

According to Proposition 3.7, we finally obtain

$$\langle \partial_t v_i, \Phi \rangle = \langle \partial_t v_i - \nabla p_i, \Phi \rangle \leq C \|\Phi\|_{L^2(W_\sigma^{1,4})}$$

for all divergence-free smooth vector field  $\Phi$ . This implies that

$$\partial_t v_i \in L_{\text{loc}}^2(\mathbb{R}^+, W_\sigma^{-1, 4/3}(\mathbb{R}^2)), \quad i = 1, 2,$$

and the same holds for  $\partial_t \tilde{v}$ . Now, since  $\tilde{v}$  belongs to  $L_{\text{loc}}^2(\mathbb{R}^+, W_\sigma^{1,4})$ , we deduce from (3.8) and Lemma 1.2 in Chapter III of [11] that  $\tilde{v}$  is almost everywhere equal to a function continuous from  $\mathbb{R}^+$  into  $L^2$  and we have in the sense of distributions on  $\mathbb{R}^+$ :

$$\frac{d}{dt} \|\tilde{v}\|_{L^2(\mathbb{R}^2)}^2 = 2 \langle \partial_t \tilde{v}, \tilde{v} \rangle_{W_\sigma^{-1, 4/3}, W_\sigma^{1,4}}.$$

We finally conclude by using the fact that  $\tilde{v}(0) = 0$ .  $\square$

**3.4. Proof of Theorem 1.4.** In this paragraph, we provide the proof of Theorem 1.4 by making use of the equation for the velocity. To that aim, we let  $(\omega_i, z_i)$ ,  $i = 1, 2$  be two Eulerian solutions of (EF), and we follow the same notations as in the previous paragraph. From Proposition 3.8, we may introduce

$$r(t) \equiv \|v_1(t, \cdot) - v_2(t, \cdot)\|_{L^2(\mathbb{R}^2)}^2 + |z_1(t) - z_2(t)|^2.$$



Let us fix a positive time  $T^*$ . We will show that  $r$  is identically zero on  $[0, T^*]$  by mean of a Gronwall type argument. Since  $T^*$  is arbitrary, this will provide uniqueness on the whole of  $\mathbb{R}^+$ . Let

$$R = \min_{t \in [0, T^*]} R(t), \quad R_M = \max_{t \in [0, T^*]} R(t)$$

where  $R(t)$  is the function defined in Proposition 3.4 and let  $T_C$  be the time introduced in Corollary 3.5. The quantities  $R$  and  $T_C$  only depend on  $T^*$ ,  $\|\omega_0\|$  and  $R_0$ . Since  $r(0) = 0$ , there exists  $T_0$  such that

$$r(t) \leq 1, \quad \forall t \leq T_0,$$

and we take  $T_0 \leq T_C$  maximal for this property to hold true.

First of all, we take advantage of the fact that  $\omega_i$  is constant around the point vortex to state harmonic regularity estimates on  $\tilde{v}(t)$  in a neighborhood of  $z_1(t), z_2(t)$ . We recall that  $z(t)$  is the middle point of  $[z_1(t), z_2(t)]$ .

**Lemma 3.9.** *For all  $t \leq T_C$ ,  $\tilde{v}(\cdot, t)$  is harmonic on  $B(z(t), R(t)/2)$ , with  $R(t) > 0$  and  $T_C$  given in Corollary 3.4. In particular, we have the following estimates:*

$$\begin{aligned} (1) \quad & \|\tilde{v}(t, \cdot)\|_{L^\infty(B(z(t), R(t)/4))} \leq C \|\tilde{v}(t, \cdot)\|_{L^2}, \\ (2) \quad & \|\nabla \tilde{v}(t, \cdot)\|_{L^\infty(B(z(t), R(t)/4))} \leq C \|\tilde{v}(t, \cdot)\|_{L^2}, \end{aligned}$$

where  $C$  only depends on  $R(t)$ .

*Proof.* Indeed, Corollary 3.5 states that on  $B(z(t), R(t)/2)$ ,  $\text{curl } v_i = 1$ , then  $\text{curl } \tilde{v} = \text{div } \tilde{v} = 0$  which means that  $\tilde{v}$  is harmonic on this ball:  $\Delta \tilde{v} = 0$ . Estimates (1) and (2) follow then from the mean-value Theorem (see [3]).  $\square$

The Gronwall estimate for  $r(t)$  reads as follows.

**Proposition 3.10.** *For all  $T \in [0, T_0]$ , for all  $p \geq 2$ , we have*

$$r(T) \leq C \int_0^T \left[ \varphi(r(t)) + p r(t)^{1-\frac{1}{p}} \right] dt,$$

where  $C$  depends only on  $T^*$  and  $\|\omega_0\|$ , and with the function  $\varphi$  defined in Lemma 2.2.

*Proof.* We proceed in several steps. Throughout the proof,  $C$  will stand for a constant depending only on  $R$  and  $R_M$ , therefore on  $\|\omega_0\|$  and  $T^*$ .

**Step 1.** We have for the velocities

$$\|\tilde{v}(T, \cdot)\|_{L^2}^2 \leq C \int_0^T \left[ r(t) + \sqrt{r(t)} \varphi(\sqrt{r(t)}) + p r(t)^{1-1/p} \right] dt, \quad (3.9)$$

$$\forall p \geq 2, \quad \forall T \leq T_C.$$

Indeed, subtracting the two equations given by Proposition 3.7 for  $(v_i, z_i)$ , we find

$$\begin{aligned} \partial_t \tilde{v} + \tilde{v} \cdot \nabla v_1 + v_2 \cdot \nabla \tilde{v} + \operatorname{div}(\tilde{v} \otimes H_1 + v_2 \otimes \tilde{H} + H_1 \otimes \tilde{v} + \tilde{H} \otimes v_2) \\ - (v_1(z_1)^\perp \cdot \delta(z_1) - v_2(z_2)^\perp \cdot \delta(z_2)) = -\nabla \tilde{p}. \end{aligned} \quad (3.10)$$

We then consider smooth and divergence-free functions  $\Phi_n \in C_c^\infty(\mathbb{R}^+ \times \mathbb{R}^2)$  converging to  $\tilde{v}$  in  $L_{\text{loc}}^2(\mathbb{R}^+, W^{1,4}(\mathbb{R}^2))$  as test functions in (3.10), and let  $n$  goes to  $+\infty$ . First, we have for all  $T \in \mathbb{R}^+$

$$\int_0^T \langle \partial_t \tilde{v}, \Phi_n \rangle_{W_\sigma^{-1,4/3}, W_\sigma^{1,4}} ds \rightarrow \int_0^T \langle \partial_t \tilde{v}, \tilde{v} \rangle_{W_\sigma^{-1,4/3}, W_\sigma^{1,4}} ds,$$

and we deduce the limit in the other terms from the several bounds for  $v_i$  stated in the proof of Proposition 3.8. This yields

$$\frac{1}{2} \|\tilde{v}(T, \cdot)\|_{L^2}^2 = I + J + K, \quad (3.11)$$

where

$$\begin{aligned} I &= - \int_0^T \int_{\mathbb{R}^2} \tilde{v} \cdot (\tilde{v} \cdot \nabla v_1 + v_2 \cdot \nabla \tilde{v}) dx dt, \\ J &= \int_0^T \int_{\mathbb{R}^2} (\tilde{v} \otimes H_1 + v_2 \otimes \tilde{H} + H_1 \otimes \tilde{v} + \tilde{H} \otimes v_2) : \nabla \tilde{v} dx dt, \\ K &= \int_0^T (v_1(z_1)^\perp \cdot \tilde{v}(z_1) - v_2(z_2)^\perp \cdot \tilde{v}(z_2)) dt. \end{aligned}$$

The next step is to estimate all the terms in the right-hand side. We now consider times  $T \leq T_0$  in (3.11). In order to simplify the notation, we set  $B = B(z(t), R(t)/4)$ .

For the first term  $I$  in (3.11), we begin by noticing that

$$\int_{\mathbb{R}^2} (v_2 \cdot \nabla \tilde{v}) \cdot \tilde{v} dx = \frac{1}{2} \int_{\mathbb{R}^2} v_2 \cdot \nabla |\tilde{v}|^2 dx = -\frac{1}{2} \int_{\mathbb{R}^2} |\tilde{v}|^2 \operatorname{div} v_2 dx = 0.$$

Moreover, Hölder's inequality gives

$$\left| \int_{\mathbb{R}^2} (\tilde{v} \cdot \nabla v_1) \cdot \tilde{v} dx \right| \leq \|\tilde{v}\|_{L^2} \|\tilde{v}\|_{L^q} \|\nabla v_1\|_{L^p},$$

with  $\frac{1}{p} + \frac{1}{q} = \frac{1}{2}$ . On the one hand, Lemma 2.1 states that  $\|\nabla v_1\|_{L^p} \leq Cp \|\omega_1\|_{L^p}$  for  $p \geq 2$ . On the other hand, we write by interpolation  $\|\tilde{v}\|_{L^q} \leq \|\tilde{v}\|_{L^2}^a \|\tilde{v}\|_{L^\infty}^{1-a}$  with  $\frac{1}{q} = \frac{a}{2} + \frac{1-a}{\infty}$ . We have that  $a = 1 - \frac{2}{p}$ , so we are led to

$$|I| \leq Cp \int_0^T \|\tilde{v}\|_{L^2}^{2-2/p} dt. \quad (3.12)$$

We now estimate  $J$ . We have

$$\begin{aligned} \int_{\mathbb{R}^2} (\tilde{v} \otimes H_1) : \nabla \tilde{v} \, dx &= \int_{\mathbb{R}^2} \sum_{i,j} \tilde{v}_i H_{1,j} \partial_j \tilde{v}_i \, dx = \frac{1}{2} \sum_i \int_{\mathbb{R}^2} \sum_j H_{1,j} \partial_j \tilde{v}_i^2 \, dx \\ &= -\frac{1}{2} \sum_i \int_{\mathbb{R}^2} \tilde{v}_i^2 \operatorname{div} H_1 \, dx = 0, \end{aligned}$$

since  $H_1$  is divergence-free, and

$$\begin{aligned} \left| \int_0^T \int_{\mathbb{R}^2} (H_1 \otimes \tilde{v}) : \nabla \tilde{v} \, dx \, dt \right| &\leq \left| \int_0^T \int_B (H_1 \otimes \tilde{v}) : \nabla \tilde{v} \, dx \, dt \right| \\ &\quad + \left| \int_0^T \int_{B^c} (H_1 \otimes \tilde{v}) : \nabla \tilde{v} \, dx \, dt \right|. \end{aligned} \quad (3.13)$$

We perform an integration by part for the second term in the right-hand side of (3.13). Arguing that  $\operatorname{div} \tilde{v} = 0$ , we obtain

$$\begin{aligned} \left| \int_0^T \int_{\mathbb{R}^2} (H_1 \otimes \tilde{v}) : \nabla \tilde{v} \, dx \, dt \right| &\leq \left| \int_0^T \int_B (H_1 \otimes \tilde{v}) : \nabla \tilde{v} \, dx \, dt \right| \\ &\quad + \left| - \int_0^T \left( \int_{B^c} (\tilde{v} \cdot \nabla H_1) \cdot \tilde{v} \, dx + \int_{\partial B} (H_1 \cdot \tilde{v})(\tilde{v} \cdot \nu) \, ds \right) dt \right| \\ &\leq \int_0^T \|H_1\|_{L^1(B)} \|\tilde{v}\|_{L^\infty(B)} \|\nabla \tilde{v}\|_{L^\infty(B)} \, dt \\ &\quad + \int_0^T \|\nabla H_1\|_{L^\infty(B^c)} \|\tilde{v}\|_{L^2}^2 \, dt \\ &\quad + \int_0^T \|H_1\|_{L^\infty(\partial B)} \|\tilde{v}\|_{L^\infty(\partial B)}^2 |\partial B| \, dt. \end{aligned}$$

According to Lemma 3.9, this gives

$$\left| \int_0^T \int_{\mathbb{R}^2} (H_1 \otimes \tilde{v}) : \nabla \tilde{v} \, dx \, dt \right| \leq C \int_0^T \|\tilde{v}\|_{L^2}^2 \, dt.$$

In the same way, we obtain by integration by part

$$\begin{aligned} \left| \int_0^T \int_{\mathbb{R}^2} (v_2 \otimes \tilde{H}) : \nabla \tilde{v} \, dx \, dt \right| &\leq \left| \int_0^T \int_B (v_2 \otimes \tilde{H}) : \nabla \tilde{v} \, dx \, dt \right| \\ &\quad + \left| - \int_0^T \left( \int_{B^c} (\tilde{H} \cdot \nabla v_2) \cdot \tilde{v} \, dx + \int_{\partial B} (v_2 \cdot \tilde{v})(\tilde{H} \cdot \nu) \, ds \right) dt \right|. \end{aligned}$$

Therefore,

$$\begin{aligned} \left| \int_0^T \int_{\mathbb{R}^2} (v_2 \otimes \tilde{H}) : \nabla \tilde{v} \, dx \, dt \right| &\leq \int_0^T \|\tilde{H}\|_{L^1(B)} \|v_2\|_{L^\infty} \|\nabla \tilde{v}\|_{L^\infty(B)} \, dt \\ &\quad + \int_0^T \|\tilde{H}\|_{L^\infty(B^c)} \|\tilde{v}\|_{L^2} \|\nabla v_2\|_{L^2} \, dt \\ &\quad + \int_0^T \|\tilde{H}\|_{L^\infty(\partial B)} \|\tilde{v}\|_{L^\infty(\partial B)} \|v_2\|_{L^\infty} |\partial B| \, dt. \end{aligned}$$

Using again Calderón-Zygmund inequality for  $v_2$  and Lemma 3.9, we get

$$\left| \int_0^T \int_{\mathbb{R}^2} (v_2 \otimes \tilde{H}) : \nabla \tilde{v} \, dx \, dt \right| \leq C \int_0^T \left( \|\tilde{H}\|_{L^1(B)} + \|\tilde{H}\|_{L^\infty(B^c)} \right) \|\tilde{v}\|_{L^2} \, dt.$$

A very similar computation yields

$$\begin{aligned} \left| \int_0^T \int_{\mathbb{R}^2} (\tilde{H} \otimes v_2) : \nabla \tilde{v} \, dx \, dt \right| \\ \leq C \int_0^T \left( \|\tilde{H}\|_{L^1(B)} + \|\nabla \tilde{H}\|_{L^2(B^c)} + \|\tilde{H}\|_{L^\infty(\partial B)} \right) \|\tilde{v}\|_{L^2} \, dt. \end{aligned}$$

We need here some estimates for  $\tilde{H}$ . We recall that  $\tilde{H}$  is defined by  $\tilde{H} = H_1 - H_2$ , so that

$$\|\tilde{H}\|_{L^\infty(B^c)} = \sup_{z \in B^c} \frac{|z_1 - z_2|}{2\pi|z - z_1||z - z_2|} \leq (8/R)^2 \frac{|z_1 - z_2|}{2\pi}.$$

On the other hand, it follows from potential theory estimates (see e.g. [5]) that

$$\int_B |\tilde{H}(x)| \, dx = C \int_B |K(x - z_1) - K(x - z_2)| \, dx \leq C\varphi(|z_1 - z_2|).$$

Concerning the  $L^2$  norm, we observe that for  $x \in B^c$ ,

$$|\nabla \tilde{H}(x)| \leq |z_1 - z_2| \sup_{[x-z_1, x-z_2]} |D^2 K| \leq \frac{C}{|x-z|^3} |z_1 - z_2|,$$

which implies that  $\|\nabla \tilde{H}\|_{L^2(B^c)} \leq C|z_1 - z_2|$ . Therefore, we arrive at

$$|J| \leq 2C \int_0^T \left( \|\tilde{v}\|_{L^2} + |z_1 - z_2| + \varphi(|z_1 - z_2|) \right) \|\tilde{v}\|_{L^2} \, dt.$$

Since  $\varphi$  is increasing, this implies

$$|J| \leq C \int_0^T \left[ r(t) + \sqrt{r(t)}\varphi(\sqrt{r(t)}) \right] \, dt. \quad (3.14)$$

Finally, we decompose the third term  $K$  in (3.11) as follows:

$$\begin{aligned} v_1(z_1)^\perp \cdot \tilde{v}(z_1) - v_2(z_2)^\perp \cdot \tilde{v}(z_2) &= (v_1(z_1)^\perp - v_1(z_2)^\perp) \cdot \tilde{v}(z_1) \\ &\quad + (v_1(z_2)^\perp - v_2(z_2)^\perp) \cdot \tilde{v}(z_1) \\ &\quad + v_2(z_2)^\perp \cdot (\tilde{v}(z_1) - \tilde{v}(z_2)). \end{aligned}$$

Applying Lemma 2.2 to  $v_1$ , we obtain

$$\begin{aligned} |v_1(z_1)^\perp \cdot \tilde{v}(z_1) - v_2(z_2)^\perp \cdot \tilde{v}(z_2)| &\leq C|\tilde{v}(z_1)|\varphi(|z_1 - z_2|) + |\tilde{v}(z_2)|\|\tilde{v}(z_1)\| \\ &\quad + \|v_2\|_{L^\infty} \|\nabla \tilde{v}\|_{L^\infty([z_1, z_2])} |z_1 - z_2|, \end{aligned}$$

so that Lemma 3.9 finally yields

$$|K| \leq C \int_0^T \left[ r(t) + \sqrt{r(t)}\varphi(\sqrt{r(t)}) \right] \, dt. \quad (3.15)$$

Estimates (3.12), (3.14) and (3.15) complete the proof of (3.9).

**Step 2.** We have for the points vortex

$$|z_1(T) - z_2(T)|^2 \leq C \int_0^T \left[ r(t) + \sqrt{r(t)}\varphi(\sqrt{r(t)}) \right] dt, \quad \forall T \leq T_C.$$

Indeed, since  $z_1$  and  $z_2$  are Lipschitz, their derivatives exist for almost every time  $t$ , and we have at these points

$$\begin{aligned} \frac{d}{dt}|z_1 - z_2|^2 &= 2\langle z_1 - z_2, v_1(z_1) - v_2(z_2) \rangle \\ &= 2\langle z_1 - z_2, v_1(z_1) - v_1(z_2) \rangle + 2\langle z_1 - z_2, \tilde{v}(z_2) \rangle. \end{aligned}$$

This yields in view of Lemmas 3.9 and 2.2

$$\frac{d}{dt}|z_1 - z_2|^2 \leq C|z_1 - z_2|\varphi(|z_1 - z_2|) + |z_1 - z_2|\|\tilde{v}\|_{L^2},$$

and we conclude by integrating the previous inequality.

Finally, we observe that for  $z \leq 1$ , we have

$$z\varphi(z) \leq \varphi(z^2), \quad z \leq \varphi(z).$$

Proposition 3.10 then directly follows from Steps 1 and 2 and the definition of  $T_0$ .  $\square$

Theorem 1.4 is now an easy consequence of the following

**Lemma 3.11.** *We have for all  $p \geq 1$*

$$\varphi(t) \leq p t^{1-\frac{1}{p}}, \quad \forall t \geq 0,$$

where  $\varphi$  is defined in Lemma 2.2.

Let us assume Lemma 3.11 for a moment and finish the proof of Theorem 1.4. We deduce from Proposition 3.10 that for all  $T \leq T_0$ ,

$$r(T) \leq C \int_0^T p r(t)^{1-\frac{1}{p}} dt.$$

This implies

$$r(T) \leq (CT)^p, \quad \forall p \geq 2.$$

Denoting  $T_1 = 1/C$  and letting  $p$  tend to infinity, we conclude that  $r(T) \equiv 0$  for all  $T < \min(T_0, T_1)$ . Iterating the proof above, we deduce that  $r$  vanishes on  $[0, T_0]$ , therefore  $T_0 = T_C$ . Finally, since  $T_C$  depends only on  $\|\omega_0\|$  and  $T^*$ , we can pass from the interval  $[0, T_C]$  to  $[0, T^*]$  by considering a finite number of translations in time of size  $T_C$ . This concludes the proof of Theorem 1.4.

**Proof of Lemma 3.11.** The result is obvious for  $t \geq 1$ . Let  $f_p(t) = t^{\frac{1}{p}}(1 - \ln t)$ . It suffices to show that  $f_p(t) \leq p$ . Computing

$f'_p(t) = t^{\frac{1}{p}-1}(\frac{1}{p}(1 - \ln t) - 1)$ , we observe that  $f'_p(t) \leq 0$  if and only if  $t \leq e^{1-p}$ . Then  $f_p$  is maximal when  $t = e^{1-p}$  and we infer that

$$f_p(t) \leq f_p(e^{1-p}) = pe^{\frac{1}{p}-1} \leq p$$

for  $p \geq 1$ . This completes the proof.

#### 4. FINAL REMARKS AND COMMENTS

**The case of several point vortices.** In this paper, we have only considered the vortex-wave system with one single point vortex. In the case of a finite number  $N$  of vortices  $z_i$  with real intensities  $d_i, i = 1, \dots, N$ , the vortex-wave system (LF) modifies as follows:

$$\left\{ \begin{array}{l} v(\cdot, t) = (K * \omega)(\cdot, t), \\ \dot{z}_i(t) = v(t, z_i(t)) + \sum_{\substack{j=1 \\ j \neq i}}^N d_j K(z_i(t) - z_j(t)), \\ z_i(0) = z_{i,0}, \\ \dot{\phi}_t(x) = v(t, \phi_t(x)) + \sum_{j=1}^N d_j K(\phi_t(x) - z_j(t)), \\ \phi_0(x) = x, \quad x \neq z_{j,0}, \\ \omega(\phi_t(x), t) = \omega_0(x), \quad t \in \mathbb{R}. \end{array} \right. \quad (\text{LFN})$$

In this situation, every vortex trajectory  $z_i(t)$  is submitted to the fields generated by the other vortices and to the regular field  $v$ , and the regular part moves under the action of the field created by itself and by the  $N$  vortices. The velocity fields appearing in (LFN) are well-defined as long as the flow and the vortices remain separated. If the intensities  $d_i$  all have the same sign, it has been established in [7] that no collision among the vortices and the flow can occur in finite time, and global existence for (LFN) has been proved for an arbitrary initial vorticity  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$  and  $N$  distinct vortices  $z_{i,0}$ .

In particular, given any time  $T > 0$ , there exists a positive  $a$  such that up to time  $T$ , we have  $|z_i(t) - z_j(t)| \geq a$ . So, the field created by a vortex near the other vortices is Lipschitz and bounded. Localizing then the test functions used throughout the proofs in Sections 2 and 3 near each vortex, we may extend Theorems 1.3 and 1.4 to the case of many vortices and obtain uniqueness for the corresponding Eulerian formulation to (LFN) when the vorticity is initially constant near each point vortex. This gives more precisely

**Theorem 4.1.** *Let  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$  and  $z_{1,0}, \dots, z_{N,0}$  be  $N$  distinct points in  $\mathbb{R}^2$  with positive intensities  $d_i$ . Assume that there exist a positive  $R_0$ , which is smaller than the minimal distance between the*

initial vortices, and  $\alpha_i \in \mathbb{R}$  such that for all  $i$ ,

$$\omega_0 \equiv \alpha_i \text{ on } B(z_{i,0}, R_0).$$

Suppose in addition that  $\omega_0$  has compact support. Then there exists a unique Eulerian solution of the vortex-wave system with this initial data.

It follows in particular from Theorem 4.1 that equivalence between Eulerian and Lagrangian formulations also holds in the case of several point vortices.

**Uniqueness when the vortex point is fixed.** In this subsection, we address the problem of uniqueness to a slightly different equation from (EF). The main difference is that the point vortex is fixed (for instance at the origin) instead of moving under the action of the velocity. It reads

$$\begin{cases} \partial_t \omega + u \cdot \nabla \omega = 0, \\ u(t, x) = v(t, x) + \gamma K(x), \quad v = K * \omega, \end{cases} \quad (4.1)$$

where  $\gamma \in \mathbb{R}$ . This system is obtained by Iftimie, Lopes Filho and Nussenzveig Lopes in [4] as an asymptotical equation for the classical Euler equations on exterior domains. More precisely, they consider a family of obstacles  $\Omega_\varepsilon \equiv \varepsilon \Omega$  contracting to a point as  $\varepsilon \rightarrow 0$ , where  $\Omega$  is a smooth, bounded, open, connected, simply connected subset of the plane. Throughout [4], the authors assume that the initial vorticity  $\omega_0$  is independent of  $\varepsilon$ , smooth, compactly supported outside the obstacles  $\Omega_\varepsilon$  and that the circulation  $\gamma$  of the initial velocity on the boundary is independent of  $\varepsilon$ . The authors prove that as  $\varepsilon$  goes to 0, the flow converges to a global solution of equation (4.1). Of course, this system reduces to the classical Euler equations when  $\gamma = 0$ , for which uniqueness is known in the class  $L^\infty(\mathbb{R}^+, L^1 \cap L^\infty(\mathbb{R}^2))$  [12].

Equations (4.1) have also been considered in the Lagrangian formulation by Marchioro [6] in the case where the support of  $\omega_0$  does not intersect the origin and in a smooth setting. In this paper, it is proved that for a  $C^2$  vorticity, a trajectory starting away from the origin never reaches it. This provides in particular uniqueness in the Lagrangian formulation in this case.

According to Section 3.1 of this work, it is actually possible to adapt the key idea used in [6] to equations (4.1) without relying on the trajectories. In particular, we first prove that if the initial vorticity  $\omega_0$  vanishes in a neighborhood of the origin, then this holds for all time.

**Proposition 4.2.** *Let  $\omega$  be a global Eulerian solution of (4.1) such that*

$$\text{supp } \omega_0 \subset B(0, R_0^{-1}) \setminus B(0, R_0)$$

for some  $0 < R_0 < 1$ . Then there exist positive constants  $C_1$  and  $C_2$  depending only on  $R_0$  and  $\|\omega_0\|$  such that

$$\omega(t) \equiv 0 \quad \text{on} \quad B(0, C_1 e^{-C_2 t}), \quad \forall t \geq 0.$$

*Proof.* As already mentioned, Section 3.1 of this work applies equally well to equations (4.1) by replacing the moving point vortex by the origin. In particular, we infer that  $\omega^2$  is also a weak solution of the linear transport equation corresponding to (4.1). According to Remark 3.6, we have

$$\text{supp } \omega(t) \subset B(0, K(1+t)), \quad \forall t \geq 0, \quad (4.2)$$

where  $K$  only depends on the initial conditions  $R_0$  and  $\|\omega_0\|$ . We aim to apply Lemma 3.2 with the choice  $\beta(t) = t^2$  and we set

$$\Phi(t, x) = \chi_0 \left( \frac{-\ln|x| - \int \ln|x-y|\omega(t, y) dy}{2\pi R(t)} \right),$$

where  $\chi_0$  is a smooth, increasing function  $:\mathbb{R} \mapsto \mathbb{R}$  which is identically zero for  $|x| \leq 1/2$  and identically one for  $|x| \geq 1$ , and  $R(t)$  is an increasing continuous function to be determined later on. We set

$$g(t, x) = \frac{1}{2\pi} \int_{\mathbb{R}^2} \ln|x-y|\omega(t, y) dy,$$

it follows from (4.2) that  $|g(t, x)| \leq C_0(1 + \ln(1+t))$  for some constant  $C_0$ . It is proved by Marchioro in [6] that if  $\omega$  is a smooth solution of (4.1),

$$\partial_t g(t, x) = - \int_{\mathbb{R}^2} K^\perp(y-x) \cdot (v(y) + K(y)) \omega(t, y) dy.$$

Marchioro's paper also states that

$$\|\partial_t g\|_{L^\infty} \leq C_1, \quad (4.3)$$

where  $C_1$  only depends on  $\|\omega_0\|$  and  $R_0$ . This can be extended to weak Eulerian solutions of (4.1) by replacing  $\ln$  by  $\ln_\varepsilon$  in the definition of  $g$ , where  $\ln_\varepsilon|z|$  coincides with  $\ln|z|$  on  $B(0, \varepsilon)^c$  and is identically equal to  $\ln \varepsilon$  in  $B(0, \varepsilon)$ . Letting  $\varepsilon$  then go to zero, we deduce that for all  $x$ ,  $g(x, \cdot)$  is Lipschitz and has a time derivative for almost every time; moreover the bound (4.3) holds at those times. We omit the details here and may consider that  $\Phi$  is  $C^1$ .

On the other hand, we have

$$\nabla_x g(t, x) = - \int_{\mathbb{R}^2} K^\perp(x-y)\omega(t, y) dy = -v^\perp(t, x),$$

therefore

$$(v + K) \cdot \nabla \Phi = (v + K) \cdot (v^\perp + K^\perp) \frac{\chi'_0}{R} \equiv 0.$$



Besides,

$$\partial_t \Phi(t, x) = \left( \frac{R'(t)}{R^2(t)} \left( \frac{\ln|x|}{2\pi} + g(t, x) \right) - \frac{1}{R} \partial_t g(t, x) \right) \chi'_0 \left( \frac{-\frac{\ln|x|}{2\pi} - g(t, x)}{R(t)} \right).$$

In view of Lemma 3.2, this yields

$$\begin{aligned} \int_{\mathbb{R}^2} \Phi(T, x) \omega^2(T, x) dx - \int_{\mathbb{R}^2} \Phi(0, x) \omega_0^2(x) dx \\ = \int_0^T \int_{\mathbb{R}^2} \omega^2 \frac{\chi'_0 \left( \frac{-\ln|x| - 2\pi g}{2\pi R} \right)}{R} \left( \frac{R'}{R} \left( \frac{\ln|x|}{2\pi} + g \right) - \partial_t g \right) dx dt. \end{aligned}$$

Since  $\chi'_0$  is non negative and has support in  $B(0, 1) \setminus B(0, \frac{1}{2})$ , we obtain

$$\int_{\mathbb{R}^2} \Phi(T, x) \omega^2(T, x) dx - \int_{\mathbb{R}^2} \Phi(0, x) \omega_0^2(x) dx \leq \int_0^T \int_{\mathbb{R}^2} \omega^2 \frac{\chi'_0}{R} \left( -\frac{R'}{2} - \partial_t g \right) dx dt.$$

Using (4.3) leads to

$$\int_{\mathbb{R}^2} \Phi(T, x) \omega^2(T, x) dx - \int_{\mathbb{R}^2} \Phi(0, x) \omega_0^2(x) dx \leq \int_0^T \int_{\mathbb{R}^2} \omega^2 \frac{\chi'_0}{R} \left( -\frac{R'}{2} + C_1 \right) dx dt.$$

We now choose

$$R(t) = C_2 + 2C_1 t,$$

with  $C_2$  to be determined later on, so that

$$\int_{\mathbb{R}^2} \Phi(T, x) \omega^2(T, x) dx \leq \int_{\mathbb{R}^2} \Phi(0, x) \omega_0^2(x) dx.$$

Since  $|g(0, x)| \leq C_0$  for all  $x$  and since  $\omega_0$  vanishes on  $B(0, R_0)$ , we have for all  $x \in \text{supp } \omega_0$

$$-\frac{\ln|x|}{2\pi} - g(0, x) \leq -\frac{\ln R_0}{2\pi} + C_0 = C_4.$$

We finally choose  $C_2$  so that

$$\frac{C_4}{C_2} \leq \frac{1}{2}.$$

For this choice, we have

$$\Phi(0, x) \omega_0^2(x) = \chi_0 \left( \frac{-\frac{\ln|x|}{2\pi} - g(0, x)}{C_2} \right) \omega_0^2(x) \equiv 0.$$

We deduce that for all  $T$ ,  $\Phi(T, x) \omega^2(T, x) \equiv 0$ . We consider  $x$  such that

$$-\ln|x| \geq 2\pi R(T) + C_0(1 + \ln(1 + T)),$$

then

$$\frac{-\ln|x| - g}{2\pi R} \geq 1,$$

therefore  $\omega(T, x) = 0$ . So finally

$$\omega(T) \equiv 0 \quad \text{on} \quad B(0, e^{-C_0(1+\ln(1+T))-2\pi R(T)}),$$

and the conclusion follows.  $\square$

Using Proposition 4.2, it is then straightforward to adapt the proof of Theorem 1.4 with a fixed vortex point instead of a moving vortex point and finally conclude that uniqueness holds for (4.1).

**Theorem 4.3.** *Let  $\omega_0 \in L^1 \cap L^\infty(\mathbb{R}^2)$  be compactly supported such that*

$$\text{supp}(\omega_0) \cap \{0\} = \emptyset.$$

*Then there exists a unique Eulerian solution of equation (4.1) with this initial data.*

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#### REFERENCES

- [1] Desjardins B., *A few remarks on ordinary differential equations*, Commun. in Part. Diff. Eq., 21:11, 16667-1703, 1996.
- [2] DiPerna R. J. and Lions P. L., *Ordinary differential equations, transport theory and Sobolev spaces*, Invent. Math. 98, 511-547, 1989.
- [3] Gilbarg D. and Trudinger N. S., *Elliptic Partial Differential Equations of Second Order*, Springer, Berlin 1998
- [4] Iftimie D., Lopes Filho M. C. and Nussenzveig Lopes H. J., *Two Dimensional Incompressible Ideal Flow Around a Small Obstacle*, Comm. Partial Diff. Eqns. 28 (2003), no. 1&2, 349-379.
- [5] Majda A. J. and Bertozzi A.L., *Vorticity and Incompressible flow*, Cambridge Texts in Applied Mathematics, 2002.
- [6] Marchioro C., *On the Euler equations with a singular external velocity field*, Rend. Sem. Mat. Univ. Padova, Vol. 84, 61-69, 1990.
- [7] Marchioro C. and Pulvirenti M., *On the vortex-wave system*, Mechanics, analysis, and geometry: 200 years after Lagrange, M. Francaviglia (ed), Elsevier Science, Amsterdam, 1991.
- [8] Marchioro C. and Pulvirenti M., *Mathematical Theory of Incompressible Non-viscous Fluids*, Springer-Verlag, 1991.
- [9] Marchioro C. and Pulvirenti M., *Vortices and Localization in Euler Flows*, Commun. Math. Phys. 154,49-61, 1993.
- [10] Starovoitov V.N., *Uniqueness of a solution to the problem of evolution of a point vortex*, Siberian Mathematical Journal, Vol. 35, no. 3, 1994.
- [11] Temam R., *Navier-Stokes Equations, Theory and Numerical Analysis*, North-Holland, Amsterdam, 1979.
- [12] Yudovich V.I., *Non-stationary flows of an ideal incompressible fluid*, Zh Vych Mat, 3:1032-1066, 1963.

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