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# Investigating the flyby scenario for the HD 141569 system

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## ABSTRACT

**Context.** HD 141569, a triple star system, has been intensively observed and studied for its massive debris disk. Until recently, it was rather regarded as a gravitationally bound triple system but recent measurements of the HD 141569A radial velocity seem to invalidate this hypothesis. The flyby scenario has therefore to be further investigated to test its compatibility with the observed disk structures.

**Aims.** We present a study of the flyby scenario for the HD141569 system, by considering 3 variants: a sole flyby, a flyby associated with one planet embedded in the disk and a flyby with two planets in the disk. We discuss the merits of each one to reproduce the scattered light observations of the disk.

**Methods.** We first use analytical calculations to reduce the parameter space of the 2 stellar companion's orbit and then perform N-body numerical simulations of the flyby encounter, using symplectic integration, taking into account the gravitational influence of the stars and the planets on massless test particles.

**Results.** The binary orbit is found to be almost fixed by the observational constraint on a edge-on plane with respect to the observers. If the binary has had an influence on the disk structure, it should have a passing time at the periapsis between 5000 and 8000 years ago and a distance at periapsis between 600 and 900 AU. It also appears that the best scenario for reproducing the disk morphology is a flyby with 1 planet embedded in the disk. For a  $2 M_J$  planet, its orbital eccentricity must be around 0.2 while for a  $8 M_J$  planet, it must be below 0.1. In the two cases, its apoapsis is about 130 AU.

**Conclusions.** Although th global disk shape is reasonably well reproduced, some observed features cannot be explain by the present model and the likelihood of the flyby event remains an issue for the scenario explored in this paper. Dynamically speaking, HD 141569 is still a puzzling system.

**Key words.** Celestial mechanics - (Stars:) planetary systems - Methods: N-body simulations - Methods: analytical - Stars: individual: HD 141569

## 1. Introduction

HD 141569, a triple star system, has been intensively observed and studied for its circumstellar dusty disk imaged in scattered light in the visible and at near-infrared wavelengths. This system is located at  $99 \pm 10$  pc and its age is estimated to  $5 \pm 3$  Myrs (Weinberger et al. 2000; Merín et al. 2004). The disk is associated to the central star (B9.5 Ve) while the two other companions (M2 and M4) form a binary.

The resolved images of the disk, showing a complicated morphology, have generated a lot of discussion in the literature about the systems dynamics (see Table 1 for a summary). First of all, this disk is in transition to a debris disk: although it is a dusty optically thin disk with a fractional disk luminosity of  $L_{disk}/L_* \simeq 8.4 \times 10^{-3}$  (Sylvester et al. 1996), the gas mass is not negligible. It could in fact represent most of the mass according to Jonkheid et al. (2006) ( $M_{gas} = 80 M_\oplus$ ,  $M_{dust} = 2.2 M_\oplus$  for grains with radius sizes between  $1 \mu\text{m}$  and  $1 \text{cm}$ ), although the spatial distribution of the gas and of the dust may differ. The influence of the gas on the dust dynamics has been therefore taken into account by Ardila et al. (2005), but other authors consider only classical N-Body simulations without gas (Augereau & Papaloizou 2004; Wyatt 2005). Quillen et al. (2005), on the other hand, only considered the gas in their simulations.

The dynamical status of the external binary is also questionable: in order to reproduce the external structures of the disk,

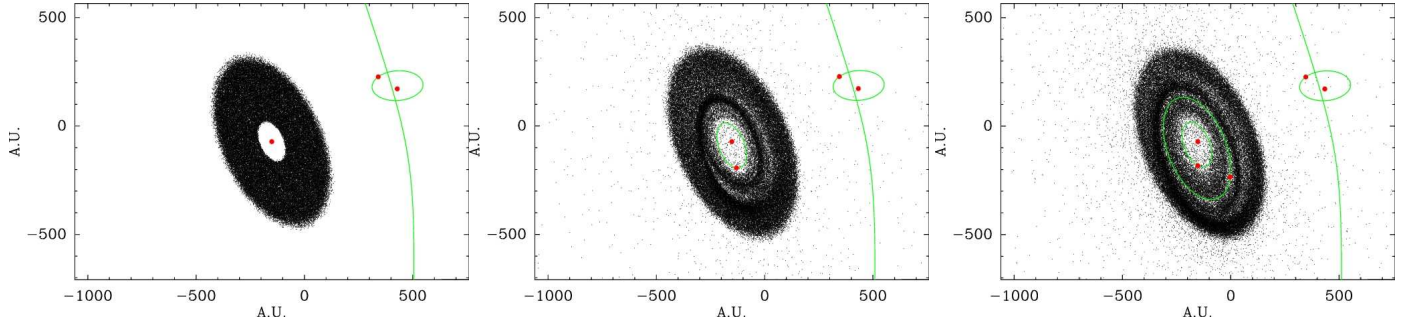
previous studies have either considered the case of gravitational bound companions (Augereau & Papaloizou 2004; Quillen et al. 2005) or the case of a flyby (Ardila et al. 2005). The internal disk structures might be explained by the gravitational perturbation of an unseen planet at large distance, around 250 AU (Wyatt 2005) but also by alternative mechanisms, such as dust migration in a gaseous disk (Takeuchi & Artymowicz 2001). The latter scenario is successful in producing annular structures, but cannot account for the observed non-axisymmetric of these features. The combined effect of a planet and external perturbers has only been considered in one study, without successfully reproducing all the structures (Ardila et al. 2005).

New millimeter observations of the gas disk (Dent et al. (2005) and Augereau et al., in prep) , which give better constraints on the radial velocity of the primary star, show that the differential velocity between the binary and the central star ( $5.8 \pm 0.3 \text{ km s}^{-1}$ ) is larger than the system escape velocity ( $2.6 \text{ km s}^{-1}$ ) so that the flyby scenario seems therefore to be the most plausible. Therefore, we propose in this paper to extensively study the flyby scenario, as it is the least studied up to now. In Sect. 2 we summarise the available observations of the disk and the stars astrometry constraints. In order to best reproduce the observations, we discuss 3 different scenarios (Fig. 1): a sole flyby (Sect. 3), a flyby with 1 planet embedded in the disk (Sect. 4) and a flyby with 2 planets embedded in the disk (Sect. 5). In Sect. 6, we compare our approach to previous studies and discuss in particular the likelihood of this scenario. We finally summarise and conclude in Sect. 7.

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**Table 1.** Summary of recent papers on the dynamical modelling of the HD 141569 system.

Authors	Modelled structures	External companions	Planet	Notes
Takeuchi & Artymowicz (2001)	gap & two rings	No	No	Dust migration due to gas friction
Augereau & Papaloizou (2004)	outer ring	bound	No	N-Body gravitational code
Quillen et al. (2005)	outer ring	bound	No	2D hydrodynamics code
Ardila et al. (2005)	spiral arms & inner depletion	flyby	Yes	N-body code, SPH code, collisions
Wyatt (2005)	gap & outer ring	No	Yes	N-body gravitation code

**Fig. 1.** Illustrations of the 3 scenarios studied in this paper: a sole flyby (left panel), a flyby + 1 planet (middle panel) and a flyby + 2 planets (right panel). The stars and planet locations, are represented by large red points, and the various orbits by a thin green line. See the electronic edition of the Journal for a color version of this figure.

## 2. Available Observations

### 2.1. Overview of the disk shape

The motivation for our study is to reproduce the dust disk structures that have been observed in scattered light. According to the visible, near-infrared and mid-infrared observations from Augereau et al. (1999), Weinberger et al. (1999), Mouillet et al. (2001), Marsh et al. (2002) and Clampin et al. (2003) the following facts appear:

- The main structure is the two rings shape of the disk. The two annuli peak at  $\sim 200$  AU and  $\sim 325$  AU from the star. The two bright rings show out of phase brightness asymmetries of up to factors of 2.5-3 for the outer ring in the visible. The outer ring moreover shows a tightly-wound spiral structure. The two annuli are separated by a darker ring or “gap”, which is radially wide compared to the two annuli.
- an extended diffuse emission associated with a faint spiral arm is present in the North-East of the disk and is detected up to more than 600 AU. According to the millimeter maps and using the hypothesis of forward scattering, this spiral would be a trailing structure with respect to the disk rotation. An other spiral arm is possibly observed, pointing toward the companions.
- the disk brightness sharply decreases between 200 and 150 AU from the central star until the background level is reached. This behaviour is suggestive of a strong depletion of dust inside 150 AU.

### 2.2. Radial velocities & astrometry

Weinberger et al. (2000) have measured the radial velocities of the HD 141569B and C stars, respectively  $-1.5 \pm 0.6$  km s $^{-1}$  and  $-2.4 \pm 0.7$  km s $^{-1}$  by cross-correlating their spectra to that of standard stars. This method was not suitable for HD 141569A due to its lack of lines for comparison with the radial velocity standard. Weinberger et al. (2000) also summarise the astrometry of the 3 stars from 1938 to 1999.

Augereau et al. (in prep) and Dent et al. (2005) have measured the double-peaked CO J=2-1 spectral profile of the disk around HD 141569A. It is thus possible to deduce the radial velocity of this star from the CO line and to obtain a far smaller uncertainty on this measure. To a first approximation, assuming a symmetric emission in the blue-shifted and red-shifted parts of the disk, the radial heliocentric velocity of HD 141569A is found to be  $-7.6 \pm 0.3$  km s $^{-1}$  (mean velocity in the line). Supposing that the star masses are 2.5, 0.5 and 0.25  $M_{\odot}$  for A, B and C respectively (Weinberger et al. 2000) and that the distance between HD141569A and the center of mass of the B-C binary is at least 800 AU (projected distance, Weinberger et al. (2000)), the maximum escape velocity is 2.6 km s $^{-1}$ . As the measured velocity is far above this limit, the two M star companions are not bound to HD 141569A. In the context of dynamical perturbation of the disk by the companions, the flyby scenario seems therefore to be the most plausible.

## 3. The flyby scenario

### 3.1. Kinematic constraints

We define a cartesian coordinate system  $(X, Y, Z)$  where  $X$  points toward the north,  $Y$  toward the east and  $Z$  toward the Earth and where the origin is HD141569A. In this referential frame, the position of the center of mass of the B-C binary will be described by the coordinates  $(X_b, Y_b, Z_b)$  and its velocity by  $V_{b_x}, V_{b_y}, V_{b_z}$ . We already know the present  $X_b, Y_b$  (right ascension and declination) and  $V_{b_z}$  (radial velocity). But the hyperbolic orbit of the binary around the primary is also defined by the 6 orbital elements  $(q, e, i, \omega, \Omega, u)$  where  $q$  is the pericenter distance,  $e$  the eccentricity,  $i$  the inclination with respect to the sky plane  $(OXY)$ ,  $\omega$  the periapsis argument,  $\Omega$  the longitude of the ascending node and  $u$  the eccentric anomaly at the observing time. The transformation from cartesian coordinates to orbital elements, or the inverse transformation, is defined by a system of 6 non-linear equations from classical keplerian formalism. It is thus possible to look for orbital element sets verifying the 3 observables  $(X_b, Y_b, V_{b_z})$ . But with 6 unknown and only 3 constraints

the equation system is underconstrained. One has therefore to consider 3 orbital elements as free parameters in order to solve the system for the 3 other orbital elements. This system is defined by the following equations, for an hyperbolic orbit:

$$X_b \cos \Omega - Y_b \sin \Omega = a \cos i \left( \sqrt{e^2 - 1} \cos \omega \sinh u + (e - \cosh u) \sin \omega \right) \quad (1)$$

$$X_b \sin \Omega + Y_b \cos \Omega = a \left( -\sqrt{e^2 - 1} \sin \omega \sinh u + (e - \cosh u) \cos \omega \right) \quad (2)$$

$$V_{b_z} \sqrt{a} (e \cosh u - 1) = \sqrt{GM} \sin i (\sinh u \sin \omega - \sqrt{e^2 - 1} \cosh u \cos \omega) \quad (3)$$

where  $G$  is the gravitational constant and  $M$  the sum of the 3 star masses. For this study, we choose to consider  $(i, \Omega, u)$  as free parameters and  $(q, e, \omega)$  as unknowns. Once the 3 parameters  $(i, \Omega, u)$  are fixed to an arbitrary set of values, the non-linear system generally admits none or one solution for the triplet  $(q, e, \omega)$ . For our problem, many triplets of the parameters  $(i, \Omega, u)$  give a solution: for each  $(i, u)$  pair, at least one value of  $\Omega$  gives a valid solution for  $(q, e, \omega)$ . The results are thus not significant because all the scenarios can be compatible with the observations: orbital plane perpendicular to the line of sight, or instead very inclined orbits, more or less eccentric hyperbolic orbits ...

However, an other constraint exists, not yet used in our method: the observed proper motion of the stars between 1938 and 1998 (Weinberger et al. 2000), and the associate velocity  $V_{sky} = 1.1 \pm 0.7 \text{ km s}^{-1}$ . The uncertainty on this observable is quite large and we consider it as an upper limit: all the solutions that predict a proper motion less than twice the observed proper motion are valid:

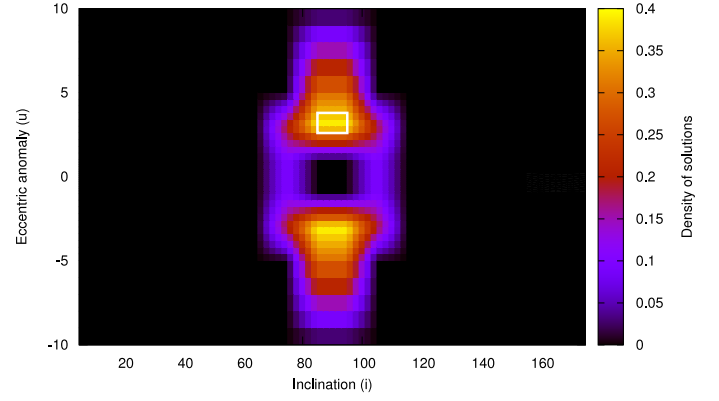
$$\sqrt{V_{b_x}^2 + V_{b_y}^2} < 2V_{sky} \quad (4)$$

This is enough to reduce significantly the number of solutions, as showed in Fig. 2. On this plot, the two axes are the eccentric anomaly  $u$  and the inclination  $i$ , while the colors correspond to the density of valid solutions (i.e. fulfilling the relation 1 to 4) for the longitude of the ascending node  $\Omega$  in the range  $[0, 2\pi]$ . The parameter space region where the valid solutions exist is indeed very much reduced by using the proper motion constraint as we find, with an inclination  $i$  between  $75^\circ$  and  $105^\circ$  and an eccentric anomaly  $u$  between  $-5$  and  $5$  radians. The fourth constraint has therefore eliminated all the solutions with lower inclination and larger proper motion.

The disk inclination is  $55^\circ \pm 1.2^\circ$  from pole-on (Mouillet et al. 2001), or rather  $180^\circ - 55^\circ$  in order to have the disk rotation (Goto et al. 2006) consistent with our convention (i.e.  $i < 90^\circ$  is an anticlockwise orbit on the sky). The binary orbits with an inclination below (resp. above)  $90^\circ$  are thus retrograde (resp. prograde) with respect of the disk rotation. The relative inclination  $i_{rel}$  between the orbital plane of the binary and that of the disk can be calculated by the following equation:

$$\cos i_{rel} = \cos i \cos i_{disk} + \sin i \sin i_{disk} \cos(\Omega - \Omega_{disk}) \quad (5)$$

with  $\Omega_{disk}$ , the position angle of the disk ( $357^\circ \pm 2^\circ$ ) and  $i, \Omega$  orbital elements of the binary orbit.



**Fig. 2.** Density of valid orbital configurations in the parameter space  $(i, u, \Omega)$ . The two axis are the eccentric anomaly  $u$  and the inclination  $i$ , while the colors correspond to the density of the valid solutions for the longitude of the ascending node  $\Omega$  in the range  $[0, 2\pi]$ . The small white box shows the final valid parameter space when the disk geometry constraints are also taken in account (Sect. 3.2). See the electronic edition of the Journal for a color version of this figure.

### 3.2. Disk geometry constraints

At this point, only the kinematic constraints have been used to calculate the valid orbital configurations for the binary. However, the encounter geometry is also constrained by the observed structures in the disk if one assumes that the current disk shape is a consequence of the gravitational perturbations of the disk by the binary. The comparison study of numerical simulations of the flyby scenario to the observations therefore reduces even more the range of valid orbital configurations as detailed below.

#### 3.2.1. Numerical model

We consider in our simulations a hierarchical system consisting of a central star associated with a debris disk and an external binary. The trajectory of this binary can be a bound orbit around the primary star or a flyby (hyperbolic orbit). We address the case of a disk consisting of large particles (planetesimals), which are insensitive to pressure forces (radiation, stellar wind or gas pressure). Importantly, we also do not take into account the gravitational interactions between planetesimals as they are negligible, nor mutual collisions. Dynamically speaking, the planetesimals are thus considered as test particles. In Sect. 4 and 5, one or several planets will be added to the simulations, orbiting the central star in the orbital plane of the disk. The central star is assumed to have a mass of  $2.5 M_\odot$  and the low mass companions  $0.5 M_\odot$  and  $0.25 M_\odot$  (Weinberger et al. 2000). The initial disk consists of 100 000 test particles with a surface density distribution proportional to  $r^{-1}$ .

To perform our simulations we use the symplectic package HJS (Beust 2003), a SWIFT variant (Wisdom & Holman 1991; Levison & Duncan 1994) for hierarchical systems. It allows us to integrate accurately and fastly the motion of the disk particles, the relative motion of the two companions and their orbit around the central star, although they are on different timescales. Numerous scenarios can be therefore tested and studied in a reasonable CPU time.



### 3.2.2. Results

Among the solutions consistent with the kinematic constraints (Sect 3.1), the periapsis of the orbit,  $d_0$ , (i.e. the closest approach) can range between 10 and several thousands of AU. If the flyby is however too close to the star, the disk can be totally dissipated or, at least, can show a spiral structure which strongly differs from the observations (Fig. 3 top left panel). We numerically explored a broad range of  $d_0$  values consistent with the kinematic constraints of Sect. 3.1. According to the simulations, we find that the distance at periapsis,  $d_0$  must be above about 600 AU. An other useful constraint is the time span between the periapsis passage and the present time,  $T = t - t_0$ . If  $T$  is too small compared to the keplerian period at the outer disk edge, the spiral structure does not have the time to develop (Fig. 3 top right panel). The simulations give an estimate of at least 5 000 years before present for the periapsis passage. All the well defined structures in the disk develop after the periapsis passage, never before, which eliminates all the configurations with  $u < 0$ , because the binary is still too far to generate any perturbations in the disk (Fig. 3 bottom left panel).

Few orbital configurations finally verify the system given by the equations 1 to 4 plus the additional constraints  $d_0 > 600$  AU and  $T = t - t_0 > 5\,000$  yrs. Valid solutions are only possible for  $i$  between  $85^\circ$  and  $95^\circ$  and  $u \approx 3$  and share the following properties:

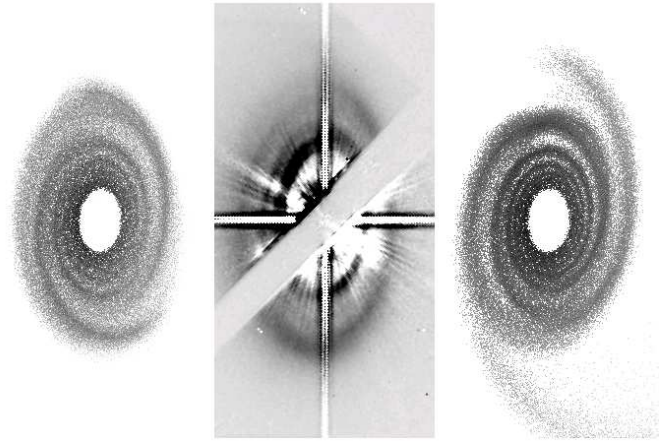
$$\begin{aligned} 5\,000 \text{ yrs} < t - t_0 < 8\,000 \text{ yrs} \\ 600 \text{ AU} < d_0 < 900 \text{ AU} \\ e \approx 10 \end{aligned}$$

Thanks to Eq. 5, it is also possible to determinate that:

$$35^\circ < i_{rel} < 145^\circ$$

Considering these solutions, the binary has at present time already reached its velocity at infinity. Within this range of solutions, the initial size of the disk can also be better constrain by our simulations in the range between 400–450 AU. If it is larger (Fig. 3 bottom right panel), the flyby produces large spiral arms which are too extended while if the disk size is smaller, the disk does not fit the observed size. We do not address here the question of the origin of the disk size, nor the shape of the external disk edge. But collisional evolution of the disk (Thébaud & Wu 2008) or gravitationnal truncation by the binary in a past bound state of the system (see Sect. 6.1) can both be used to explain them.

Overall, this scenario gives satisfactory results for the global shape of the disk and reduces significantly the range of possible orbital configurations for the companions. This is illustrated in Fig. 4 which shows a disk of correct size with a outer trailing spiral. The inner disk structure is nevertheless not properly reproduced and the next step is thus to reproduce more accurately the two ring-like structure at distances of about 200 and 325 AU from the star (Mouillet et al. 2001; Clampin et al. 2003). It appears however that none of the valid solutions for the flyby is able to well generate the two rings. In the case of retrograde orbit ( $85^\circ \leq i < 90^\circ$ ), the flyby generates a clear two arms spiral structure (Fig. 4, left panel). In the case of prograde orbit ( $90^\circ < i \leq 95^\circ$ ), as studied for instance by Pfalzner (2003), the two arms are more wound and more extended but it is still too different from the observations (Fig. 4, right panel). In the two cases, the flyby does not perturb the inner part of the disk, below 200 AU. It thus cannot explain the inner ring, nor the observed depletion around 150 AU. It means that we need to add some other hypothesis to the model.



**Fig. 4.** Best solutions for the flyby scenario on a retrograde orbit (left panel) and on a prograde (right panel) orbit, compared to the observations from Mouillet et al. (2001) (middle panel). A scattering anisotropic factor  $g$  of 0.2 is assumed to mimic real scattered light observations.

## 4. The flyby + planet scenario

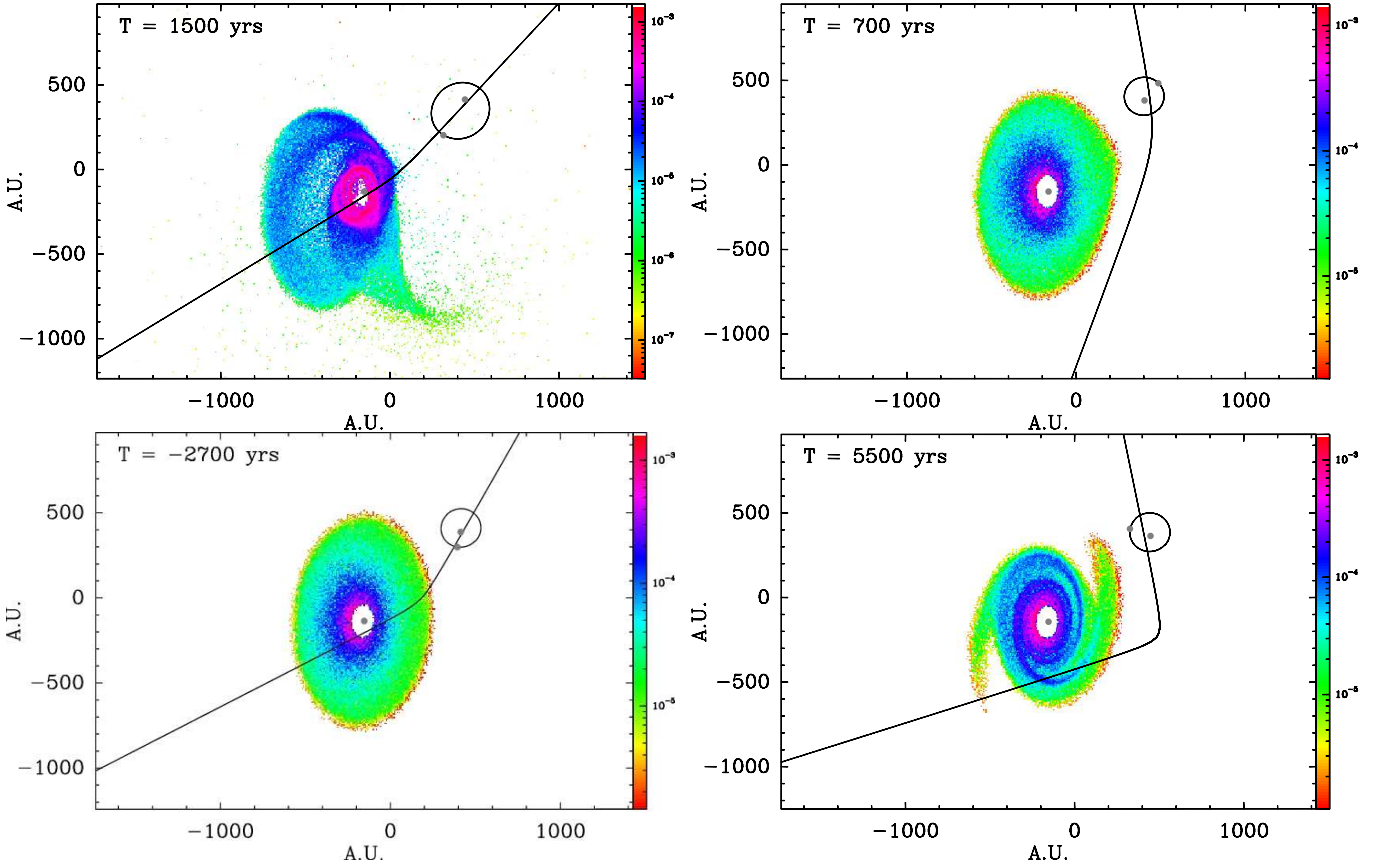
We can conclude from Sect. 3 that the sole presence of the two star companions is not enough to well explain the observed disk properties summarised in Sect. 2, especially the double-ring structure. The flyby can, at least, constrain the initial disk size between 400 and 450 AU but it is necessary to use other hypothesis to improve the model and to try to better fit the disk geometry. As the observations show a strong depletion in the dust surface density below 150 AU, a plausible hypothesis is the presence of an unseen planet which generates this gap.

In this section, we consider the case where we add a planet within 150 AU in our numerical models. The planet is supposed to be coplanar with the disk and its main parameters are its mass, its orbital semi-major axis and its orbital eccentricity. The simulations take into account the mutual gravitational interactions between the planet and the 3 stars, as well as the gravitational perturbation of the disk by the planet but not the feedback of the disk on the planet: the planetesimals are still considered as test particles. The system consisting of the central star, the disk and the planet is first evolved during 5 Myrs (the assumed age of the system), then the two star companions are introduced to simulate the flyby with one of the valid set of orbital parameters found in Sect. 3.

### 4.1. The inner ring

To constrain the planet parameters, we compare the simulations to the disk optical thickness derived from Marsh et al. (2002). The authors have combined their own observations in the mid-infrared and Weinberger et al. (1999) observations in the near-infrared (Fig. 5) to estimate the radial profiles of the vertical dust optical depth at  $1.1 \mu\text{m}$ . The observations in the mid-infrared cover the inner region of the disk, up to 130 AU while the observations in the near-infrared cover the region between 130 and 500 AU. The radial profile of the optical depth at  $1.1 \mu\text{m}$  is deduced, for the mid-infrared data, from the optical depth at  $12.5 \mu\text{m}$  and assuming the opacity scales inversely with wavelength.

The profiles show a ring which peaks around 200 AU, a depletion at 250 AU and an even more depleted region below 150 AU. We use in the following these data to constrain the mass and the orbital elements of the planet, by trying to reproduce some



**Fig. 3.** Examples of parameters for the flyby scenario inconsistent with the observations: the periastris is too close to the disk, at 170 AU, (top left panel) the crossing time at periastris is too recent, 700 years, (top right), the binary is still not yet at periastris (bottom left) and the disk is too large (bottom right). These figures plot the intensity, in logarithm scale, of scattered light by the planetesimal spatial distribution obtained with our simulations. The scattering anisotropic factor is 0.2. The time is relative to the crossing time at periastris,  $t_0$ .

key characteristics of the azimuthal averaged radial density profile of the dust:

- a peak between 190 and 210 AU
- a ratio  $\tau_{peak}/\tau_{250}$  of about 1.3 between the peak surface density at  $\sim 200$  AU and the depletion at  $\sim 250$  AU
- a ratio  $\tau_{peak}/\tau_{150}$  of about 2.7 between the peak at  $\sim 200$  AU and the inner depleted region at 150 AU. The uncertainty on this ratio is nevertheless rather large because of the large uncertainties in the optical depth between 100 and 150 AU.

The “1 planet + flyby” scenario is tested for 3 different planet masses: 0.2, 2 and  $8 M_J$ . For each planet mass, 3 apoapsis are considered (130, 150 and 180 AU) with 4 different values of the eccentricity: 0, 0.05, 0.1 and 0.2. This gives a total of 36 simulations and their results are summarised in Tables 2, 3 and 4. The peak position (Table 2) indicates the position of the largest over-density if several are present.  $\tau_{peak}/\tau_{150}$  is defined as the ratio between the dust density at the peak position and the maximum density between 100 and 150 AU.  $\tau_{peak}/\tau_{250}$  is defined as the ratio between the dust density at the peak position and the density around 250 AU.

According to these simulations, a  $0.2 M_J$  planet cannot reproduce the observational data: it does not clear out enough the inner part of the disk. For higher mass planets, the best solution is obtained for an apoapsis of 130 AU. For larger values of apoapsis, the peak position is shifted toward too large distance from the star, or the planet can even prevent the formation of an

over-density. For a  $2 M_J$  planet, the best fits are obtained for eccentricities around 0.2, while for a  $8 M_J$  planet, it must be below 0.1. The best solution between these two planets is for a  $2 M_J$  planet.

With an initial surface density distribution proportional to  $r^{-1}$ , Fig. 5 shows that there is a lack of particles in the outer part of the disk compared to the optical depth profile at  $\lambda = 1.1 \mu\text{m}$ . We therefore adjust the initial surface density to obtain the same maximum at the outer peak location, near 330 AU. We find that the best fit is obtained for an initial surface density distribution proportional to  $r^{-0.5}$  and that the outer peak position in our simulation is consistent with the observations. However, there is still a lack of particles between the two rings and outside the outer ring. This might reflect a limitation of our model as we only model the dynamics of planetesimals. One can expect that smaller dust particles, sensitive to radiation pressure and generated by collisions, will migrate outside the rings and flatten the profiles, hence better reproducing the optical depth profile.

#### 4.2. The global structure

The ratio between the two peaks in the optical depth profile depends of the observing wavelength, as the optical depth at the outer peak position is larger than the optical depth at the inner peak position at  $0.5 \mu\text{m}$  (Ardila et al. 2005), while it is the opposite of Fig. 5 ( $\lambda = 1.1 \mu\text{m}$ ). Our constraint on the initial surface density distribution should therefore be taken with care, but our

**Table 2.** Summary of the results for the peak position (in AU) in the planetesimal density profile, for all the simulations, as a function of the planet mass, eccentricity and apoapsis. “–” indicates the simulations where no peak is observable or where the peak is below 170 AU. † indicates the simulations where multiple peaks appear. In this case, the largest peak is used to calculate the peak position,  $\tau_{peak}/\tau_{150}$  and  $\tau_{peak}/\tau_{250}$ . The numbers boldface correspond to the best solutions.

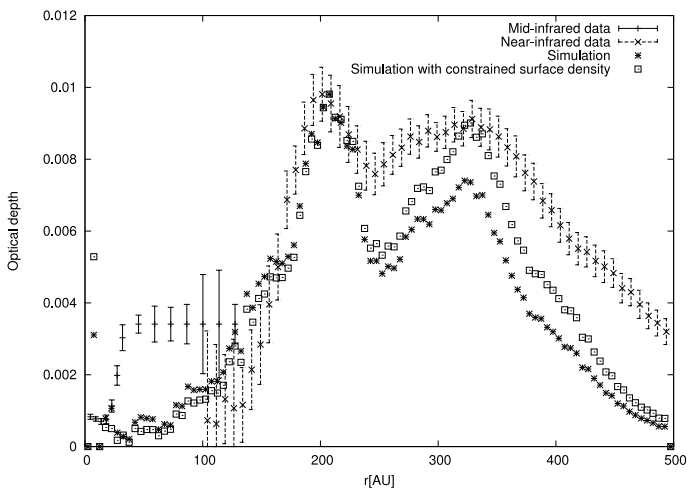
Planet mass	$0.2 M_J$				$2 M_J$				$8 M_J$			
	0	0.05	0.1	0.2	0	0.05	0.1	0.2	0	0.05	0.1	0.2
apoapsis of 180 AU	210	210	210	–	–	230	–	–	–	–	–	–
apoapsis of 150 AU	220†	–	–	230	–	210†	230	240	–	–	–	–
apoapsis of 130 AU	–	–	–	220	–	220†	–	<b>210</b>	<b>220</b>	<b>220</b>	230	220
apoapsis of 120 AU	–	–	–	–	–	–	200	<b>190</b>	<b>220</b>	<b>220</b>	230	230

**Table 3.** Same as Table 2, but for the ratio  $\tau_{peak}/\tau_{150}$  between the density at peak position and the maximum density between 100 and 150 AU.

Planet mass	$0.2 M_J$				$2 M_J$				$8 M_J$			
	0	0.05	0.1	0.2	0	0.05	0.1	0.2	0	0.05	0.1	0.2
apoapsis of 180 AU	0.5	0.5	0.5	–	–	–	0.4	–	–	–	–	–
apoapsis of 150 AU	0.5†	–	–	0.8	–	0.6†	1.6	3.0	–	–	–	–
apoapsis of 130 AU	–	–	–	1.1	–	2.4†	–	<b>2.3</b>	<b>3.0</b>	<b>5.0</b>	6.0	8.0
apoapsis of 120 AU	–	–	–	–	–	–	1.3	<b>2.1</b>	<b>2.7</b>	<b>4.7</b>	4.2	4.5

**Table 4.** Same as in Table 2, but for the ratio  $\tau_{peak}/\tau_{250}$  between the density at peak position and the density in the gap around 250 AU.

Planet mass	$0.2 M_J$				$2 M_J$				$8 M_J$			
	0	0.05	0.1	0.2	0	0.05	0.1	0.2	0	0.05	0.1	0.2
apoapsis of 180 AU	1.5	1.7	1.3	–	–	–	1.1	–	–	–	–	–
apoapsis of 150 AU	1.7†	–	–	1.8	–	1.4†	1.5	2.0	–	–	–	–
apoapsis of 130 AU	–	–	–	2.0	–	1.7†	–	<b>2.0</b>	<b>1.9</b>	<b>1.9</b>	1.7	1.3
apoapsis of 120 AU	–	–	–	–	–	–	2.2	<b>2.1</b>	<b>1.8</b>	<b>1.7</b>	1.8	1.4



**Fig. 5.** Estimated profiles of dust optical depth (normal to the disk plane) by combining mid-infrared data (Marsh et al. 2002) and near-infrared data (Weinberger et al. 1999), compared to the simulation output of the best scenario, a  $2 M_J$  planet with an eccentricity of 0.2.

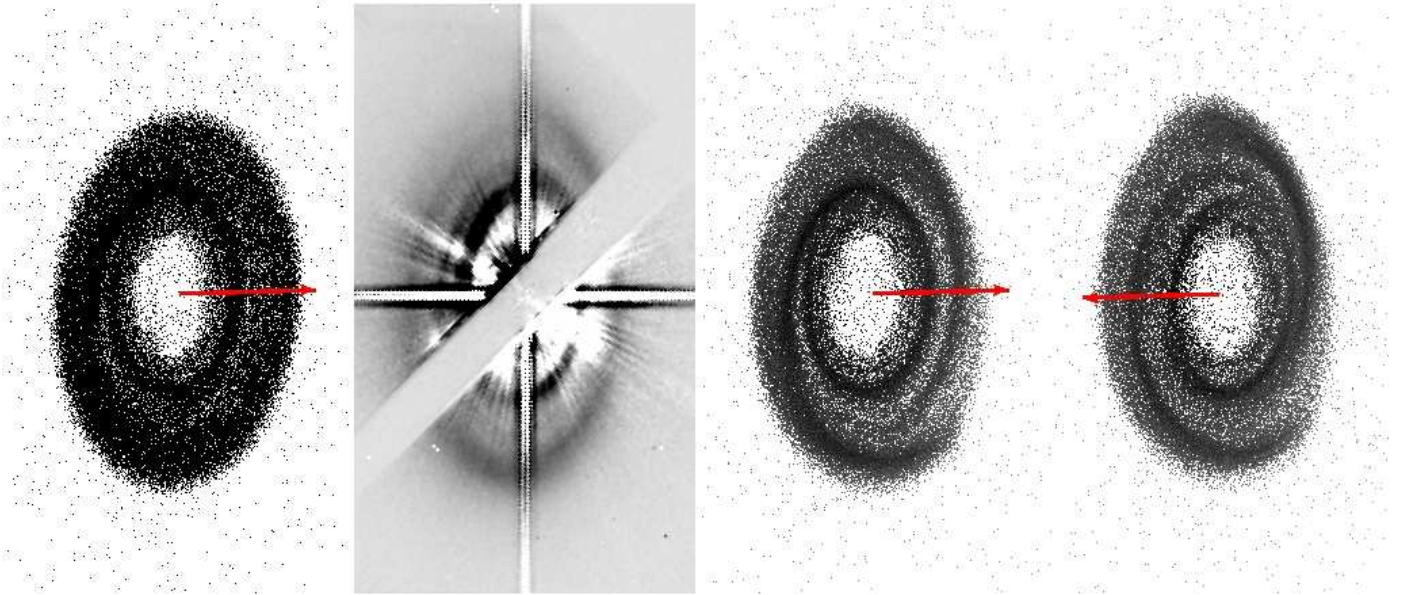
model is nevertheless very successful to predict the position of the outer peak in the optical depth profile. The planet parameters are strongly constrained by the inner ring, but our simulations show that the planet also perturb the outer parts of the disk as shown in Figs. 5 and 6.

In less than 5 Myrs, the planet has time enough to develop a strong, one-armed, anticlockwise wound spiral, due to the differential precession rate of the planetesimal orbits as explained by Augereau & Papaloizou (2004). The flyby does thus not occur with an initially smooth disk and this has strong consequences on the final shape of the structures produced by the flyby. For instance, in the case where the binary is on a retrograde orbit, we previously obtained a regular two armed spiral. But now, the planet not only produces the inner ring, but can also make disappear, or at least perturb, a significant fraction of one of the two spiral arms. As shown in Fig. 6, the disk morphology between 200 and 300 AU slightly depends of the direction of the planet periapsis with respect to the direction of the binary periapsis, but the differences remain too small to constrain the orientation of the planet orbit. The results are thus not perfect because the simulations do not still reproduce exactly the outer ring but they clearly show that we need a complete description of the system to properly account for it. It is not satisfactory to consider solely the external companions or the inner planet. Combining both, we clearly achieve a much better fit to the observations.

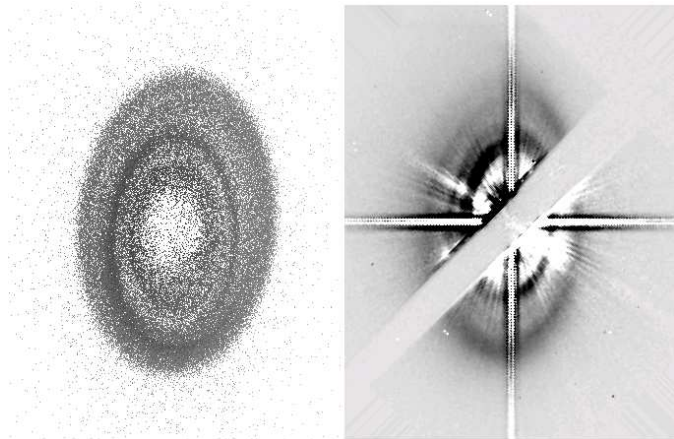
## 5. The flyby + 2 planets scenario

As the previous scenario is still not fully satisfactory because it still does not generate two rings clearly separated by a large gap, we have also investigated an alternative with two planets in the disk. Wyatt (2005) has already considered the case of a single planet at about 250 AU to explain the outer ring but without considering the perturbation due to the companion stars. We





**Fig. 6.** Same as in Fig. 4, but for the best solution for the flyby + planet scenario in the case of a retrograde orbit of the binary (third panel). The red arrow indicates the direction of the planet periapsis. The planet mass is equal to  $2 M_J$ , it apoapsis is 130 AU and its eccentricity 0.2. The first panel show the disk before the flyby but after 5 Myrs of gravitational perturbation by the planet. The fourth panel shows an alternative configuration where the initial direction of the planet periapsis is turned by  $180^\circ$ .



**Fig. 7.** Same as in Fig. 4, but for the flyby + 2 planets scenario. The inner planet mass is equal to  $2 M_J$ , it apoapsis is 130 AU and its eccentricity 0.2. The outer planet mass is  $0.2 M_J$ , it apoapsis is 250 AU and its eccentricity 0.05.

therefore first investigate the consequence of both the companions and the planet on the disk geometry. Using the solutions found in Sect. 3, it appears that adding the gravitational influence of the two companion stars does not improve the structure generated by the external planet but, on contrary, disrupts them because the planetesimals orbits are too destabilised by the perturbations of a planet and that of the two companions. For this scenario, the constraints used in Sect. 3.2 are no longer valid and the correct orbital elements for the binary orbit are those which prevent interactions between the disk and the two M stars: they do not still reach the periapsis location ( $u < 0$ ) or the periapsis itself is located at more than 1500 AU from the primary star.

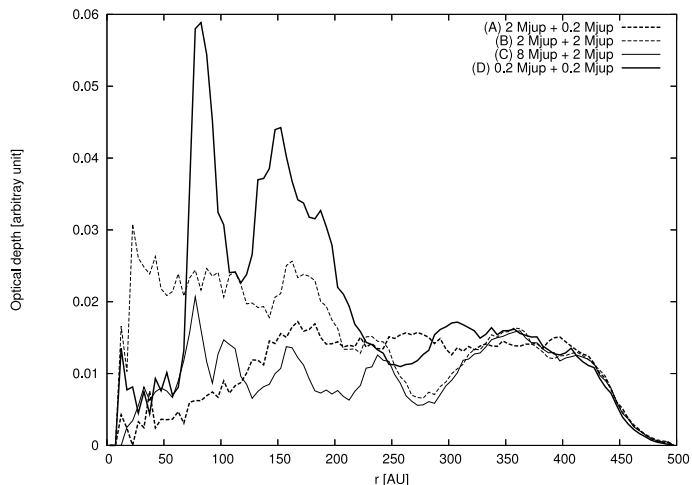
Such a scenario cannot, however, reproduce the inner depletion below 150 AU; another planet is thus necessary like in the previous section. According to Wyatt (2005), two scenarios can

be considered: a  $0.2 M_J$  planet evolving over 5 Myrs on an orbit with a semi-major axis of 250 AU and an eccentricity of 0.05, and a  $2 M_J$  planet perturbing the disk over 0.5 Myrs and on the same orbit. We add thus another planet in the internal part of the disk, following the best solutions found in the previous section, namely a  $2 M_J$  planet on an orbit with a semi-major axis of 110 AU and an eccentricity of 0.2 or a  $8 M_J$  planet on an orbit with a semi-major axis of 115 AU and an eccentricity of 0.05 (Fig. 7).

Like in Sect. 4, the estimated radial profiles of the dust optical depth obtained with these simulations are compared to the observed optical depth at  $1.1 \mu\text{m}$  (Fig. 5). Two major informations are therefore given by this method. First, and independently of the mass of the inner planet, a  $2 M_J$  at 250 AU generates a too deep gap in the disk compared to the observations (scenario B and C of Fig. 8). Contrary to scenario “flyby + 1 planet”, one cannot expect that the optical depth profile would be flattened by smaller dust particles sensitive to radiation pressure because the particles will be quickly removed from the gap, due to close encounter with this planet orbiting at 250 AU. Wyatt (2005) constrains the planet parameters considering the location and the winding of the spiral but there is a degeneracy between the planet mass and the time the planet is orbiting around the star. If one adds the constrain given by the observations on the depth of the gap around 250 AU, only the scenarios with smaller planet masses ( $0.2 M_J$ ) orbiting for a long time (5 Myrs) are satisfactory (scenario A and D of Fig. 8)). Second, the simulations show that it is difficult to obtain an inner ring massive enough between the two planets, around 200 AU. Due to the perturbations of these two objects, a lot of planetesimals are ejected from the disk. We have therefore tried other simulations with a smaller mass ( $0.2 M_J$ , scenario D) for the internal planet in order to less perturb the disk. However, the results are still not satisfactory because the peak position is not at the correct location (150 AU instead of 200 AU) and because, with a low mass planet, there are too much planetesimals below 100 AU.

A planet located at 250 AU from the primary star is an interesting alternative to the companions in order to explain the





**Fig. 8.** Optical depth obtained by simulations with 4 different scenarios for the mass of the 2 planets, to compare to Fig.5.

outer ring without the gravitational perturbation of the 2 stellar companions, but our work shows that combining it with an other planet around 120 AU to explain the inner ring and the dust depletion within 150 AU does not provide satisfactory results. Even if it is still not perfect, the scenario “flyby + 1 planet” is the one which gives the best results among the three studied scenarios in this paper.

## 6. Discussion

### 6.1. Probability of the encounter

The main drawback to the flyby scenario is that, according to our models, the binary passed to the periastris less than 8000 years ago. The probability of observing such an event is low, making the flyby scenario statistically unlikely. On the other hand, the similar physical properties between the central star and the low mass companions suggest a common origin. The 3 stars have been independently estimated to be around 5 Myrs (Weinberger et al. 2000; Merín et al. 2004) and the differential radial velocity between the HD 141569A and the binary is only a fraction of the total speed of the primary star. But the observed radial velocities clearly show that the binary is on an escape orbit and the open spiral structure observed in the disk is also in favour of a flyby scenario, as trailing structures are created by tidal effects (Augereau et al., in prep.).

One can however reconcile these positions with an hybrid scenario. HD 141569 could have been, in the past, a triple bound system until an other star of its vicinity approached and destabilised the system, putting the binary on its current orbit. In this case, we would be observing the last periastris passage of the binary and the consequences of a 4 star encounter. This scenario could explain the observed velocities, the common physical properties and the disk size, and it is more plausible that a sole flyby of the binary. The gravitationally bound scenario studied by Augereau & Papaloizou (2004) and Quillen et al. (2005) can possibly be used to constraint the orbit of the binary before the encounter: typically an orbital period of several hundreds of thousands years and a apoapsis of almost 10000 AU. It appears more probable that a fourth star destabilises such an orbit that the binary randomly cross the path of the primary star at less than 1000 AU. The study done by Scally & Clarke (2001) show that, in the Orion Nebula Cluster, almost all the stars undergo at least

one encounter between 1000 and 10000 AU. On the other hand, Aarnio et al. (2008) have found that HD 141569 formed in relative isolation, tens of parsecs away from the recent sites of star formation in the Ophiucus-Scorpius-Centaurus region. The likelihood of such a scenario has thus to be better estimated. Its dynamical modelling has also to be investigated by this is beyond the scope of this paper. Theoretical works are however already abundant in the literature (see, for example, the series discussing many aspects of the three-body scattering from Hut & Bahcall (1983) to Heggie et al. (1996)).

### 6.2. Comparison with previous works

Other theoretical studies of the HD 141569 system can be divided in 3 groups with respect to the reproduced disk structures: the outer ring, the gap between the two rings or the extended spiral arms.

Augereau & Papaloizou (2004) and Quillen et al. (2005) have both studied the outer ring of the disk, but using different tools: a pure N-body code for Augereau & Papaloizou (2004) and a 2D hydrodynamics code for Quillen et al. (2005) assuming in the latter case that the gas and the dust are fully mixed. Their main issue compared to our work is that they have considered bound orbits for the binary which is incompatible with the new measurement of radial velocities for the primary star. Interestingly, Augereau & Papaloizou (2004) well reproduce the outer ring but not the open spiral arms structure while it is the contrary for our work. One can therefore ask if the scenario with 4 stars can reconcile these two approaches. This has to be investigated in a future paper.

Wyatt (2005) and Takeuchi & Artymowicz (2001) have considered the opening of the gap between the two rings with very different mechanisms: a planet embedded in the disk (Wyatt 2005) or the sole interaction between the dust and the gas (Takeuchi & Artymowicz 2001). In the two cases, the authors do not consider the influence of the two companion stars which can be justified by our study of the kinematic constraints. It is indeed possible to obtain a binary orbit reproducing the observed radial velocities but which does not let the two M stars generate any perturbation to the disk because they pass too far from the disk or because they have not yet passed at the periastris. The planet used by Wyatt (2005) is efficient to open the gap and generate the outer ring but, as shown in Sect. 5, it is difficult to improve this scenario in order to also explain the inner ring and the inner disk depletion. Takeuchi & Artymowicz (2001), with an alternative scenario of dust migration due to the gas friction, generate a two ring structures but the hypothesis of axisymmetricity does not leave a lot of possibilities to extend this work and explain the tightly-wound spiral structure of the outer ring nor the spiral arm in the extended diffuse emission.

Ardila et al. (2005) consider several physical processes for the modelling of HD 141569: gravitation, interaction between the gas and the dust, dust generation by collisions. It is therefore, up to now, the most complete theoretical study about HD 141569. They considered the case of a flyby in order to obtain large spiral arms similar to those observed in the extended diffuse emission. They also added a  $5 M_J$  planet, on an orbit with a semi-major axis of 100 AU with an eccentricity of 0.6, to explain the depleted zone below 150 AU. Although the authors obtain good results for the the extended spiral arms and the inner hole, they did not succeed to reproduce the two rings nor the gap in between. One possible explanation is that, in our work, we obtain satisfactory results only if we let the disk evolves with the planet during several million years before the flyby. To go further in the

description of the HD 141569 system, it will be therefore interesting to use the model developed by these authors but with the initial conditions on the binary orbit and on the planet parameters established by our work.

## 7. Summary and conclusion

The binary orbit is finally found to be almost fixed by the observational constraint on a edge-on plane with respect to the observer. If the binary has had an influence on the disk structure, it should have a passing time at the periapsis between 5000 and 8000 years ago and a distance at periapsis between 600 and 900 AU. The scenarios with retrograde orbits better reproduce the observations but prograde orbits cannot be totally excluded. In order to reproduce the observed structures in the debris disk, it also appears that the best scenario is a flyby with 1 planet embedded in the disk. For a  $2 M_J$  planet, its orbital eccentricity must be around 0.2 while for a  $8 M_J$  planet, it must be below 0.1. In the two cases, its apoapsis is about 130 AU.

This scenario is able to both reproduce the observed radial velocities of the stars and the overall structure of the disk with the modelling of a flyby by the binary and a embedded planet in the disk. However, this scenario is not fully satisfactory because it does not reproduce all the structures, and in particular the entire gap between the two annulus. Future study can therefore extend this scenario, for example with the modelling of the gas and/or collisions, as we have not taken into account these effects on the dust dynamics.

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