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Climatological features of electron density in the polar ionosphere from long-term observations of EISCAT/ESR radar

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Abstract. In this paper, climatological features of the polar F2-region electron density (N_e) are investigated by means of statistical analysis using long-term observations from the European Incoherent Scatter UHF radar (called EISCAT in the following) and the EISCAT Svalbard radar (ESR) during periods of quiet to moderate geomagnetic activity. Field-aligned measurements by the EISCAT and ESR radars operating in CP-1 and CP-2 modes are used in this study, covering the years 1988–1999 for EISCAT and 1999–2003 for ESR. The data are sorted by season (equinox, summer and winter) and solar cycle phase (maximum, minimum, rising and falling). Some novel and interesting results are presented as follows: (1) The well-known winter anomaly is evident during the solar maximum at EISCAT, but it dies out at the latitude of the ESR; (2) The daytime peaks of N_e at EISCAT for all seasons during solar maximum lag about 1–2 h behind those at ESR, with altitudes about 10–30 km lower. (3) In addition to the daytime peak, it is revealed that there is another peak just before magnetic midnight at ESR around solar maximum, especially in winter and at equinox. The daytime ionization peak around magnetic noon observed by ESR can be attributed to soft particle precipitation in the cusp region, whereas the pre-midnight N_e maximum seems likely to be closely related to substorm events which frequently break out during that time sector, in particular for the winter case. (4) Semiannual variations are found at EISCAT during solar minimum and the falling phase of the solar cycle; at the rising phase, however, the EISCAT observations show no obvious seasonal variations.

Keywords. Ionosphere (Particle precipitation; Plasma convection; Polar ionosphere)

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1 Introduction

The Earth's ionosphere is formed by the ionization of the atmosphere, owing to the solar EUV and X-ray radiations and, mainly for the high-latitudes, by energetic particles impinging on the atmosphere from the magnetosphere. Once formed, the ionized particles of ionospheric plasma are affected by numerous processes, including chemical reactions and various transport due to neutral wind, electric and magnetic fields, etc. The ionosphere is coupled closely with both the magnetosphere and the thermosphere in which the ionosphere is immersed. These lead the ionosphere to vary significantly with solar cycle, season, local time, and geomagnetic activity. The background characteristics of ionospheric electron density are of basic importance in understanding and forecasting space weather.

Electron density variations of the mid- and low-latitude ionospheric F2-layer have been extensively investigated, and a number of anomalies, such as the winter anomaly and the semiannual anomaly, have been reported (Yonezawa and Arima, 1959; Yonezawa, 1971). A number of possible causes have been proposed for these anomalies, as reviewed by Rishbeth (1998). The best-accepted mechanism is the chemical explanation suggested by Rishbeth and Setty (1961). Millward et al. (1996) modeled the semiannual variation, confirming the role of the chemical composition changes. Recently, using ionosonde data, Rishbeth et al. (2000b) studied the semiannual and annual variations of the ionospheric F2-peak at middle and low latitudes. They found that $H_m F_2$ has a well-defined relationship with solar activity. Zou et al. (2000) numerically investigated the annual, seasonal and semiannual variation of the F2-layer electron density under geomagnetically quiet conditions, and claimed that the noon $N_m F_2$ is closely related to the ambient atomic/molecular concentration ratio. Rishbeth et al. (2000a) highlighted the role of global-scale thermospheric motions, which produce the composition changes, in the

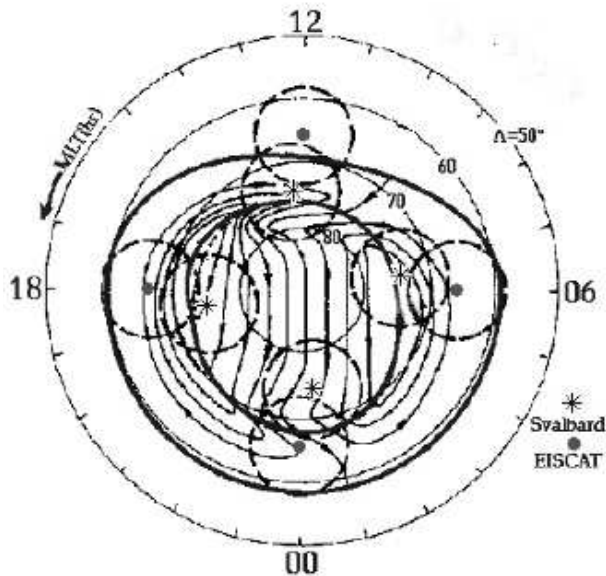


Fig. 1. The position of EISCAT/ESR radar in relation to with the auroral ellipse and polar cap. (The solid points represent the location of EISCAT; asterisks for the ESR, and the two thick black ellipse lines imply the boundaries of the auroral oval during lower geomagnetic activity).

formation of seasonal and semiannual effects in the mid-latitude F2-layer.

As the aforementioned anomalies are frequently observed at mid- and low-latitudes, they were thought to die away at latitudes higher than 60° . Baron and Wand (1983a), however, reported these anomalies in Alaska (65.1° N) on the basis of the Chatanika incoherent radar observations. A couple of years later, the winter anomaly at Tromsø (69.6° N, 19.2° E) was reported by Farmer et al. (1990). Liu (2001) investigated the climatological features of the high-latitude ionosphere using several years of EISCAT/ESR observations, but her research was only carried out in terms of two characteristic parameters, namely, the ionospheric F2-layer electron density peak, $N_m F_2$, and the corresponding altitude, $H_m F_2$. In the Southern Hemisphere, Zhu et al. (1996, 1998) reported the existence of an F-region winter anomaly over Zhongshan Station (77.1° S), based on ionosonde data around solar maximum (1990–1991). Low energy particle precipitation was proposed to be a possible cause of the ionospheric F-region N_e peak around magnetic noon. By analyzing observations from the same station, Shen et al. (2003) claimed that there is no winter anomaly around solar minimum, and highlighted the role of magnetospheric particle precipitation in the formation of averaged ionospheric N_e based on long-term observations at high latitudes, resulting in particular from the polar cusp/cap regions, are not as well reported in the literature.

In both hemispheres, there are funnel-shaped areas with near zero magnetic field magnitude but dense plasma, which are called the polar cusps or clefts. In these regions, solar plasma has direct access to the upper atmosphere. They provide an entry for particles from the solar wind and magnetosphere into near-Earth geospace without crossing geomagnetic field lines. The unique configuration of the geomagnetic field at polar cusps makes them among the most popular regions for research into the geospace environment. Generally, the cusp regions lie around $\pm 78^\circ$ MLAT, with about 5° extension in latitude. The ESR is located close to Longyearbyen on the island of Spitsbergen. The position of the EISCAT and the ESR in relation to the auroral oval and polar cap is shown in Fig. 1. It can be seen that EISCAT radar lies in the auroral oval for most of the time. The ESR radar, however, lies in the vicinity of the cusp region during the daytime and within the polar cap during the nighttime. The specific locations and powerful capabilities of these two radars provide an excellent opportunity for comparative studies of polar ionospheric characteristics, especially in the auroral oval and polar cusp/cap regions.

The EISCAT radar (69.6° N, 19.2° E) began taking data in 1981, while the ESR (78° N, 16° E) began operations in 1996. Hence a relatively long period of ionospheric measurements have now been accumulated. In this paper, field-aligned measurements from the EISCAT UHF radar during the 22nd solar cycle and from the ESR radar during the period of 1999–2003 are combined in a statistical study of the climatological characteristics of electron density in the polar ionospheric F2-layer.

2 Data analysis

Field-aligned measurements obtained from EISCAT Common Program One and Two (CP-1 and CP-2) during 1988–1999 and measurements from the ESR during 1999–2003 are used in this work. In order to eliminate the effects of extensive disturbances, only data with the maximum $K_p \leq 4$ are included. The annual averaged sunspot numbers during this period are shown in Fig. 2. According to the variations of averaged sunspot number, the period of selected ESR observations is around solar maximum. As for the EISCAT UHF data, the period of 1988–1999 can be divided into 4 phases: the rising phase (1988, 1998 and 1999), solar maximum (1989, 1990 and 1991), the falling phase (1992, 1993 and 1994) and solar minimum (1995, 1996 and 1997). Each phase is further divided into 3 seasons, namely the equinox, summer and winter. We do not distinguish spring and autumn seasons here because there are not enough observations to separate them statistically, especially for the falling phase of the solar cycle at EISCAT. The numbers of EISCAT/ESR radar observation events for each group are listed in Table 1. The entries correspond to the number of time sequences of

Table 1. Number of EISCAT/ESR radar observation events for each group.

Seasons	EISCAT				ESR Solar Max
	Rising Phase, Solar Max.,	Falling Phase, Solar Min.			
Equinox (Feb 5–May 4) (Aug. 4–Nov. 5)	25	14	6	12	150
Summer (May 5–Aug 4)	12	8	10	11	30
Winter (Nov 5–Feb 4)	8	12	11	15	105

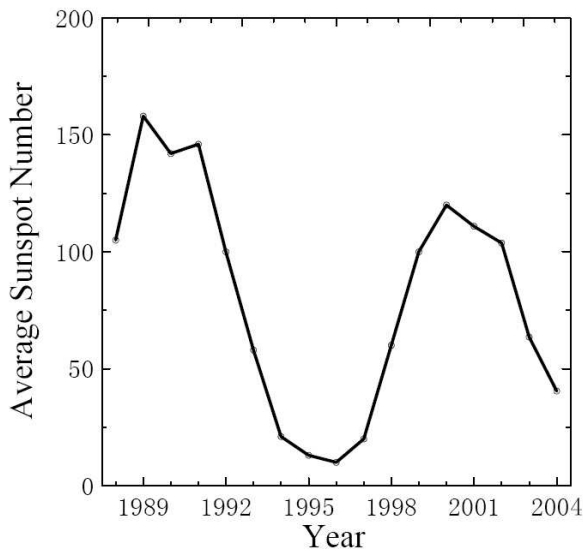


Fig. 2. Annual averaged sunspot number during 1988–2004.

successive measurements with lengths between 4 hours and one day, which have been used in the study.

Global-scale thermospheric motions have an important effect on the long-term variation of the F2-layer electron density (Rishbeth et al., 2000a), and Rishbeth (1998) has pointed out that thermospheric circulation may differ during quiet and disturbed times. Hence, it seems to be an oversimplification to assume that the daytime peak of electron density, $N_m F_2$, has no relation to geomagnetic activity, though Baron and Wand (1983b) claimed that there was no obvious relation between the nighttime $N_m F_2$ and the K_p index. Therefore, only data obtained under quiet to moderate geophysical conditions are chosen in our study to eliminate variations caused by extensive disturbances. This is different from the philosophy employed by Farmer et al. (1990), where measurements made under all conditions were used.

The signal-to-noise ratio (SNR) indicates the reliability of the data and is therefore important in the analysis. Unreliable observations are eliminated according to the quality of data fitting. Since the input data come from several different kinds of radar experiments, they usually have different time and range resolutions. For the convenience of analysis, a fixed 2-D grid has been specified, with resolutions of 10 km in range and 10 min in UT time. A linear interpolation method is employed to project all the data sets onto the aforementioned fixed grid. The resulting data sets are binned by height (H , every 10 km), hour of day (UT , every 10 min), and season of year (S). A sequence of N_e measurements representing the same conditions can be obtained by collecting all observations falling into the same bin. Within each bin, the final “mean” electron density is defined as the median value of the corresponding electron density sequence:

$$\langle N_e(H, UT, S) \rangle = \text{Median}(\{ N_{e_i}(H, UT, S) \}), \quad (1)$$

where, $\{ N_{e_i}(H, UT, S) \}$ corresponds to the N_e sequence obtained in the (H, UT, S) bin. The standard deviation σ is calculated by

$$\sigma = \sqrt{\frac{\sum(\langle N_e(H, UT, S) \rangle - N_{e_i}(H, UT, S))^2}{N(H, UT, S) - 1}}, \quad (2)$$

where $N(H, UT, S)$ is the number of available data in the (H, UT, S) bin, and plotted as an error bar on every 10th temporal grid points in the following figures. Since the UT coverage of input data is uneven, $N(H, UT, S)$ varies in different bins. In most bins, $N(H, UT, S)$ is no less than 70 percent of the number of observation events. The worst cases appear around dawn and dusk at equinox of solar minimum, with $N(H, UT, S)$ being slightly higher than 30 percent of the total number of observation events.

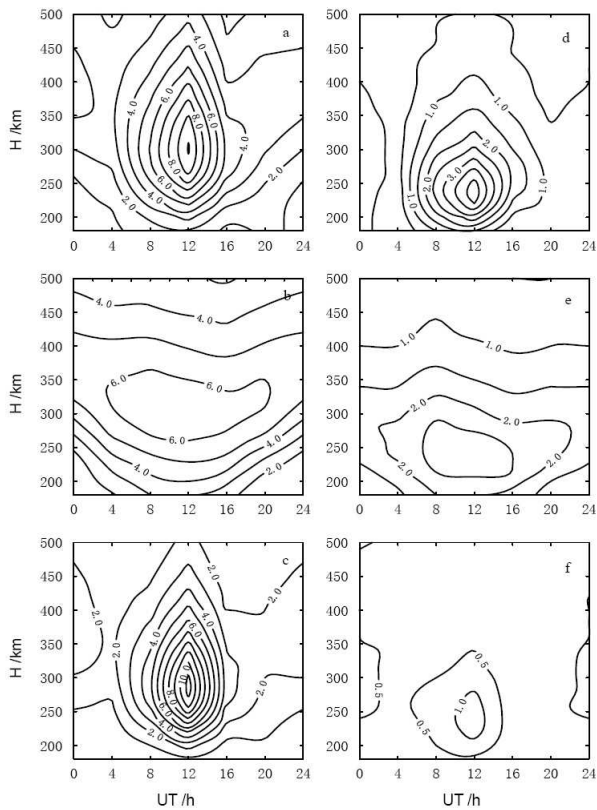


Fig. 3. 2-D distributions of mean N_e versus UT and height H around solar maximum (**a**, **b** and **c**) and solar minimum (**d**, **e** and **f**). (top panel: equinox; middle panel: summer; bottom panel: winter) (units: 10^{11} m^{-3}).

3 Climatological characteristics of N_e

3.1 Statistical results from the EISCAT (1988–1999)

3.1.1 Results around solar maximum and minimum

2-D distributions of mean N_e vs. UT and H for the 3 seasons around solar maximum (3a, 3b and 3c) and solar minimum (3d, 3e and 3f) are shown in Fig. 3. Diurnal variations of N_e at the altitudes of the F2-peak are plotted in Fig. 4 for solar maximum (4a, 4b and 4c, at 290 km) and solar minimum (4d, 4e and 4f, at 230 km), respectively.

Around the solar maximum, N_e in winter reaches a prominent peak of $11.5 \times 10^{11} \text{ m}^{-3}$ at 11:40 UT (12:57 LT), at about 290 km altitude. Note that there is a slight decrease during 04:00–06:00 UT and also a slight increase during 18:00–22:00 UT. The maximum N_e around equinox is $10.3 \times 10^{11} \text{ m}^{-3}$ at 11:50 UT (13:07 LT). Compared to that in winter, the density peak at equinox is wider (in time), covering the range of 04:00–16:00 UT. N_e near summer reaches a maximum of $6.7 \times 10^{11} \text{ m}^{-3}$ at 09:50 UT (11:07 LT). The peak becomes wider between 04:00–20:00 UT, resulting in a

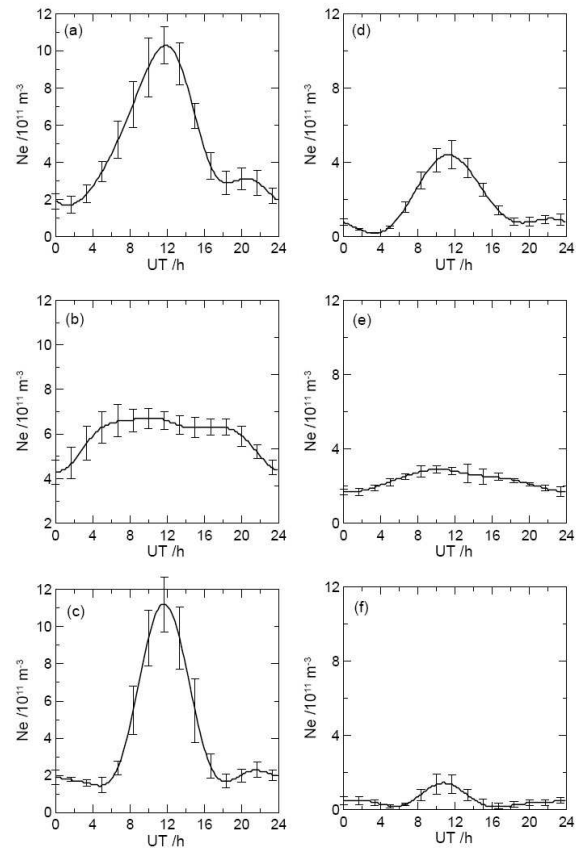


Fig. 4. Diurnal variation of mean N_e at 290 km around solar maximum (**a**, **b** and **c**) and at 230 km around solar minimum (**d**, **e** and **f**). (top panel: equinox; middle panel: summer; bottom panel: winter) (units: 10^{11} m^{-3}).

smoother daily variation. It shows no sign of a minimum in the morning and evening sectors. The maximum electron density in winter is nearly twice as high as that in summer, exhibiting the manifestation of the well-known winter anomaly around solar maximum.

Around solar minimum, N_e in winter is quite different from that around solar maximum. The maximum value is only $1.5 \times 10^{11} \text{ m}^{-3}$ at 10:50 UT (12:07 LT). The diurnal variation is smooth throughout all UT times. N_e at summer and equinox varies in a similar way to that near solar maximum, but with relatively lower values, $2.9 \times 10^{11} \text{ m}^{-3}$ for summer and $4.5 \times 10^{11} \text{ m}^{-3}$ for equinox. The higher density at equinox than in summer and winter demonstrates the existence of the semiannual anomaly at EISCAT latitudes around solar minimum.

Compared to solar maximum, the altitudes of the density peaks near solar minimum are about 60–70 km lower.

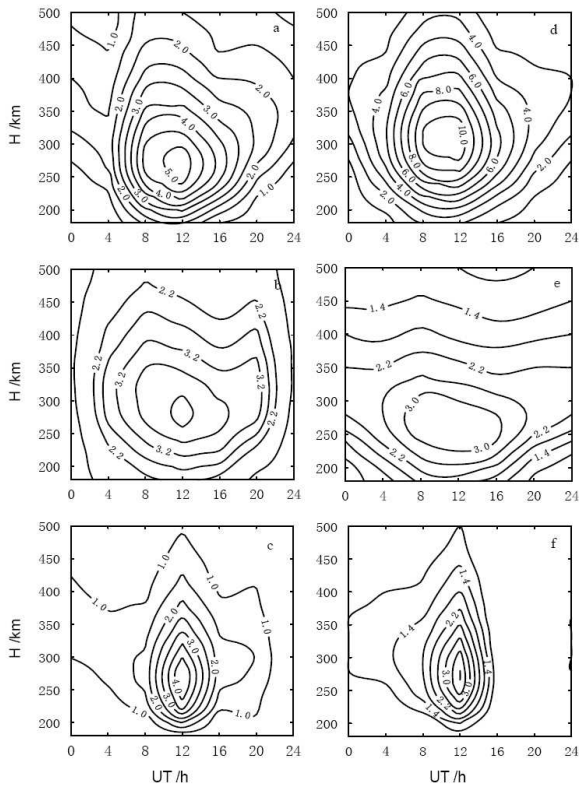


Fig. 5. The same as Figure 3, but for the rising (a, b and c) and the falling (d, e and f) phase of solar cycle. (top panel: equinox; middle panel: summer; bottom panel: winter) (units: 10^{11} m^{-3}).

3.1.2 Results at rising and falling phase

Seasonal variations of N_e exhibit quite different features in the rising and falling phases of the solar cycle. Mean 2-D distributions of N_e vs. UT and H for the rising phase (5a, 5b and 5c) and falling phases (5d, 5e and 5f) are shown in Fig. 5. Corresponding diurnal variations at 270 km (at about the altitude of the F2-peak) are plotted in Fig. 6 for the rising (6a, 6b and 6c) and falling phase (6d, 6e and 6f), respectively.

During the rising phase of the 22nd solar cycle, the maximum N_e is $4.5 \times 10^{11} \text{ m}^{-3}$ for winter, $4.8 \times 10^{11} \text{ m}^{-3}$ for summer and $5.6 \times 10^{11} \text{ m}^{-3}$ for equinox, exhibiting no obvious seasonal variations.

During the falling phase, however, N_e exhibits quite different seasonal behaviors. At equinox, N_e reaches a peak value of $11.0 \times 10^{11} \text{ m}^{-3}$ at 10:50 UT (12:07 LT). Compared to the dramatic peak near equinox, the maximum N_e near winter is much smaller, with a value of $4.0 \times 10^{11} \text{ m}^{-3}$ at 11:10 UT (12:27 LT). The peak is confined in the range of 08:00–16:00 UT. N_e near summer reaches a maximum of $3.4 \times 10^{11} \text{ m}^{-3}$ at 11:00 UT (12:17 LT), with a smooth daily variation. The N_e at equinox is more than 2 times higher than those in summer and/or winter, demonstrating obvious characteristics of the semiannual anomaly.

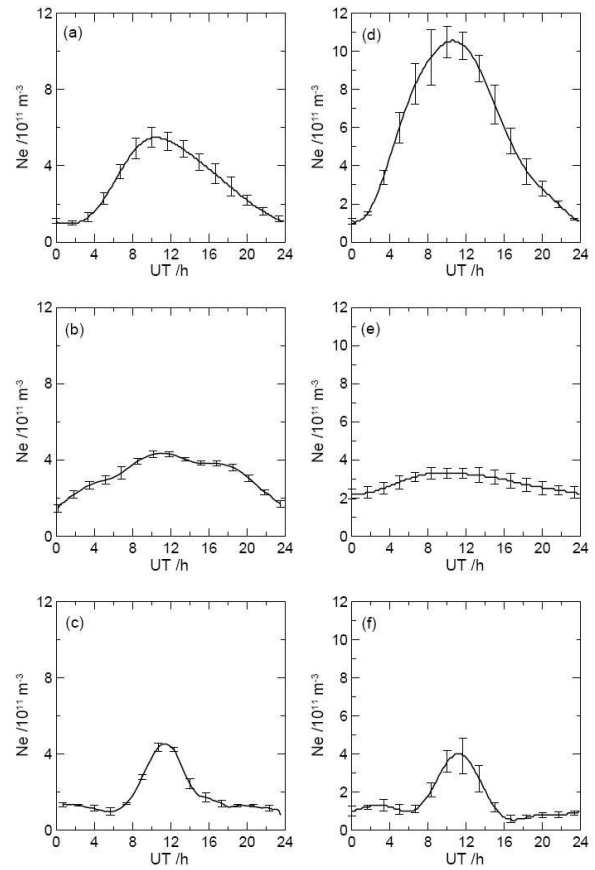


Fig. 6. The same as Fig. 4, but at 270 km for the rising (a, b and c) and the falling phase (d, e and f) of solar cycle. (top panel: equinox; middle panel: summer; bottom panel: winter) (units: 10^{11} m^{-3}).

3.2 The ESR radar observations around solar maximum

2-D mean N_e distributions derived from the ESR data around solar maximum are shown in the left-hand column of Fig. 7a, b and c. Corresponding diurnal variations of N_e at 310 km are shown in the right column (7d, 7e and 7f).

The variations of N_e are quite different to those at Tromsø around solar maximum (Figs. 3a, b, c and 4a, b, c). In winter, N_e reaches a maximum value of $5.2 \times 10^{11} \text{ m}^{-3}$ at 19:45 UT (22:35 MLT), with a minor peak at 09:20 UT (12:05 MLT). At equinox, N_e reaches a prominent maximum of $8.7 \times 10^{11} \text{ m}^{-3}$ at 09:40 UT (12:25 MLT), followed by a weaker but long-lasting peak around 18:00 UT (20:40 MLT). Note that the increase of N_e before 09:00 UT is very fast. The summertime N_e is about $4.0 \times 10^{11} \text{ m}^{-3}$ before it starts to increase to $5.3 \times 10^{11} \text{ m}^{-3}$ at 10:10 UT (12:30 MLT). The peak density is relatively unchanged up to 20:00 UT before N_e drops gradually back to $4.0 \times 10^{11} \text{ m}^{-3}$ at 22:00 UT.

Except for the pre-midnight peak, N_e in winter is less than that in summer, demonstrating no sign of the so-called winter anomaly at the latitude of the ESR around solar maximum.

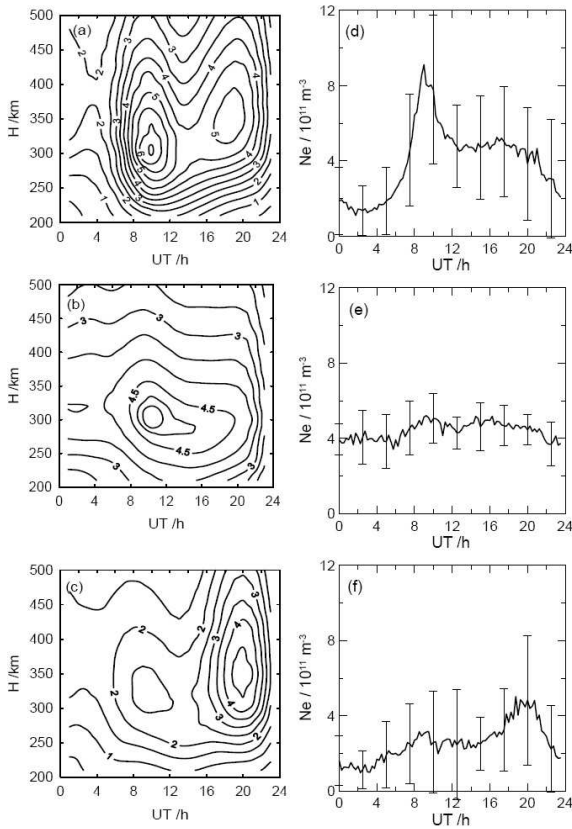


Fig. 7. Mean N_e around solar maximum at the ESR site: 2-D distribution vs. UT and H (a, b and c) and diurnal variations at 310 km (d, e and f). (top panel: equinox; middle panel: summer; bottom panel: winter) (units: 10^{11} m^{-3}).

The altitude of the daytime N_e peak at ESR is around 320 km for all seasons, about 10–30 km higher than that at the EISCAT site. Those for the minor peak at equinox and the maximum peak in winter before magnetic midnight are at about 350 km, slightly higher than that during the daytime.

Note that the standard deviations of the ESR density estimates are appreciably larger than those at Tromsø, implying larger day-to-day variations of the electron density. This may perhaps be related to the characteristic dynamic processes at the ESR latitude, which will be discussed in the following section.

4 Discussion

The aforementioned N_e distributions around solar maximum over the ESR site exhibit quite different features from those observed at auroral oval latitudes (EISCAT radar) and at mid-latitudes (Millstone Hill radar) (Lei et al., 2005), such as the two-peaked distribution of the electron density in winter and at equinox. The daytime density peaks over the ESR appear 1–2 h later and 10–30 km higher than that at EISCAT site, in-

dicating that N_e distributions observed by ESR may not be produced by solar irradiations alone. Plasma dynamic and electrodynamic processes which are characteristic at that latitude may contribute to the special features of electron density in ionospheric the F2-layer.

4.1 Soft particle precipitation in cusp region

Particle precipitation has been extensively investigated in the past several decades, with emphasis on particles' spectra and related parameters (e.g. integral energy flux, number flux and characteristic energy). Based on long-term satellite observations, various empirical models of particle precipitating have been established. The widely-accepted models are from Spiro et al. (1982) and Hardy et al. (1985, 1987, 1989). On the basis of FAST observations and EISCAT radar measurements, a new empirical model of auroral precipitating electrons, PEM-2004, has been established by our colleagues at Wuhan University (private communication). One of the agreements among these models is that soft electrons are dominant in the cusp region around magnetic midday. In the cusp region, auroral precipitating electrons usually have high number flux but low energy flux, resulting in low average energy. Around magnetic noon, the ESR radar moves near to the cusp region (Fig. 1). During this time period, an ionospheric ionization enhancement at altitudes around 320 km has been observed by ESR in all seasons (Fig. 7a, b and c).

Predicted results over the ESR site (75.3° ILAT) from PEM-2004 under conditions of $100 < AE < 300$ nT are shown in Fig. 8. It is shown that auroral electrons over the ESR site have a maximum number flux with the value of $6.3 \times 10^8 \text{ el cm}^{-2} \text{ s}^{-1}$ during the magnetic post-noon period (Fig. 8a). The minimum average energy, about 140 eV, appears at about 12:00 MLT (Fig. 8b). At magnetic noon (Fig. 8c), the average energy reaches the minimum value of 127 eV at 76.3° ILAT, which is rather similar to the characteristic energy of soft particles from the magnetosheath.

Millward et al. (1999) numerically simulated the ionospheric effects caused by soft particle precipitation in the cusp region. Ionization rates caused by precipitating electrons with different energies ranging from 100 to 1000 eV are shown in Fig. 9. The simulations imply that electrons with various energies precipitate at different altitudes. Electrons with higher energy can penetrate deeper into the atmosphere, while softer electrons stop at higher altitudes. Auroral electrons with energies of 100 and 200 eV induce the highest ionization rates at altitudes around 360 and 300 km, respectively. According to the calculations of Millward et al. (1999), the characteristics of soft particle precipitation in cusp region are reasonably consistent with the ionization enhancements around 320 km observed by the ESR.

Note that intense summertime electron density enhancements around magnetic noon have also been reported in recent case studies (Smith et al., 2000; Sims et al., 2005). In a case with southward IMF conditions, Smith et al. (2000)

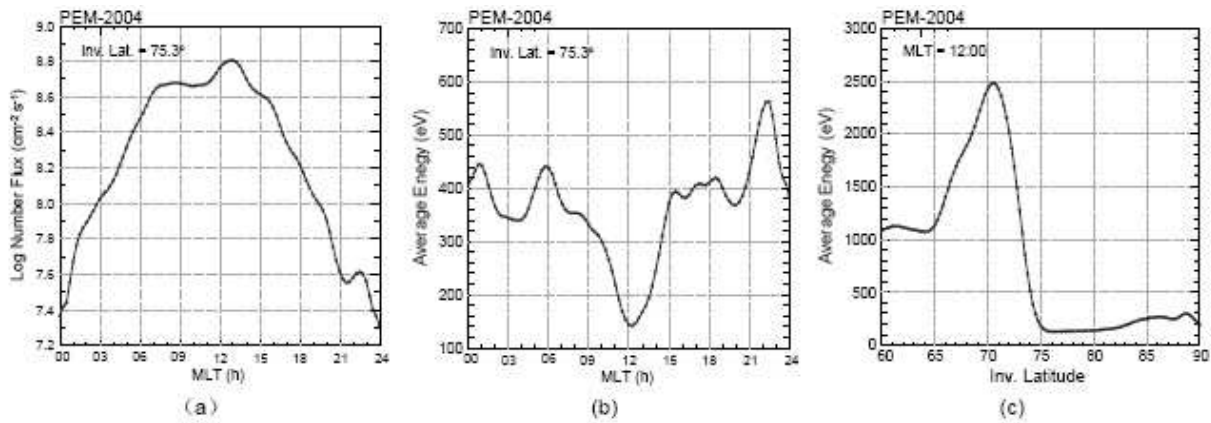


Fig. 8. Parameters forecasted by PEM-2004 model at the ESR site under conditions of $100 < AE < 300$ nT. (a) diurnal variation of integral number flux, (b) diurnal variation of average energy, (c) variations of average energy with respect to invariable magnetic latitude at 12:00 MLT.

concluded that the origin of electron density enhancements observed by ESR was soft-particle precipitation in the cusp region. In another event also under southward IMF conditions, however, Sims et al. (2005) argued that plasma convection may play a more important role than soft-particle precipitation in the formation of the summertime Tongue-of-Ionization (TOI).

4.2 Substorms in the winter-time evening sector

The occurrence of a density peak before magnetic midnight during winter is one of the outstanding features of the N_e variations over the ESR site. N_e at equinox varies in a similar way, but with its maximum peak in the daytime. Note that the ESR radar was within polar cap during that time period.

In Fig. 8a, it is shown that a minor peak in the number flux of auroral electrons appears around 22:00 MLT, when the corresponding energy exceeds 500 eV. According to Fig. 9, an electron precipitation source with such characteristics does not induce a strong ionization enhancement at altitudes around 350 km. At such times, the contribution of solar UV/EUV radiation can also be reasonably neglected. Therefore, the observed enhancement may be attributed to plasma transport processes, which are related to the electric field and neutral wind, including plasma convection and field-aligned transportation, chemical loss process, as well as mass exchanging process between topside ionosphere and magnetosphere.

FUV images from the IMAGE satellite during substorms have also been involved to investigate possible causes for the wintertime ionization enhancement before midnight MLT. Using the list of substorms produced by Frey et al. (2004), we simply compared the mean N_e with that observed by the ESR during the relevant time period of each day. A strong correlation was found between the F2-layer ionization enhancement and substorm events, implying a possible relationship

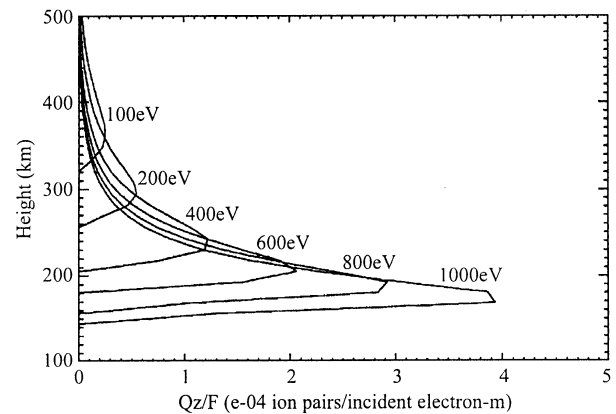


Fig. 9. Ionization production rate caused by precipitating electrons with energies ranging from 100 to 1000 eV. (From Millward et al., 1999).

between them. During the substorm expansion phase, the auroral oval extends poleward, and the ionization enhancements caused by particle precipitation may travel towards and be observed by the ESR. This hypothesis seems to be reasonably consistent with the much larger N_e fluctuations seen at the ESR (Fig. 7), especially around 20UT in winter, which may indicate the intermittent appearance of such substorm contributions.

Dayside ionization enhancements, transported across the pole in the anti-sunward convection may, however, provide an alternative interpretation of the midnight N_e enhancement. Soft-particle precipitation in sub-polar regions has been cited as a possible source for nighttime polar cap ionization, with the less energetic incoming particles giving rise to enhanced density at higher altitudes where the lifetime of plasma is sufficiently long for it to be transported over the polar cap (Weber et al., 1984). Using ground-based optical and digital

ionosonde measurements and DE-2 observations, Weber and coauthors concluded that the ionospheric structures in the nighttime polar cap F-layer were produced near or equatorward of the dayside auroral zone, and convected across the polar cap in the anti-sunward direction. Using incoherent scatter radar observations, Foster (1993) reported evidences which supported the hypothesis that photoionization convected poleward through the cleft as a source of polar-cap ionization. Unfortunately, this hypothesis finds it difficult to explain why the midnight N_e concentrations are much higher than the dayside densities observed by ESR field-aligned beam (see Fig. 7c and f) unless effective and persistent ionization enhancement mechanisms were present during the cross-polar convection of the plasma.

5 Summary

A statistical analysis of the seasonal characteristics of electron density in the polar ionospheric F2-layer has been carried out on the basis of long-term EISCAT and ESR observations. Around solar maximum, EISCAT observations show an obvious winter anomaly, but the ESR measurements exhibit no sign of it. Electron densities measured during the daytime at the EISCAT site are generally much higher than the corresponding ESR measurements, especially in winter, and the daytime N_e peak at EISCAT lags behind that at ESR by about 1–2 h.

In addition, characteristics of semiannual variation are found at EISCAT around solar minimum, as well as during the falling phase of the solar cycle. During the rising phase, however, the radar observations show no obvious seasonal variations.

For the first time, a prominent wintertime N_e peak before magnetic midnight has been observed by the ESR radar. Near equinox, there is a weaker but stable ionization enhancement in almost the same time period. The strong correlation suggested by a preliminary comparison of the EISCAT data with the list of substorms implies that this pre-midnight ionization enhancement may be closely related to substorm events which frequently break out during that MLT sector.

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