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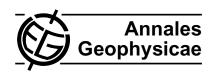
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Remote sensing of a near-Earth neutral line during the 5 October 2000 substorm

D. Nagata¹, S. Machida¹, S. Ohtani², S. B. Mende³, Y. Saito⁴, and T. Mukai⁴

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Abstract. In this paper we examined the continuous motions of a near-Earth neutral line during the recovery phase of the 5 October 2000 substorm. Estimation was based on the PSBL ion beam model proposed by Onsager (1991) and the Geotail observations. Estimated distances from the Earth ranged from 20 to $60 R_E$ and retreated tailward at velocities of 250 and 300 km/s. This event initiated with the arrival of solar wind discontinuity. Simultaneous observations of electromagnetic field and electrons indicate the existence of earthward propagating waves associated with field-aligned currents. Based on these observations, we suggest that the source of the PSBL ion beams was the retreating near-Earth neutral line formed by the compression of the magnetosphere. Two scenarios of near-Earth neutral line motion in the tail dynamics are also proposed. One is the formation of plural neutral lines to create a long plasmoid. The other is the oscillation of one neutral line between the near-Earth region and the mid-tail stagnant plasmoid.

Keywords. Magnetospheric configuration and dynamics; Plasma sheet; Storms and substorms

1 Introduction

Magnetic reconnection in the tail plasma sheet is one of the basic components of substorm dynamics and it changes the structure of the tail plasma sheet through the formation of neutral lines (Hones, 1977).

Neutral lines are observed in the near-Earth tail $(-20 > Xgsm > -30 R_E)$ at the onset of substorm (Nagai et al., 1998; Machida et al., 1999; Miyashita et al., 1999). After its formation, the neutral line retreats downtail during the recovery phase of the substorm (Russell and McPherron, 1973; Hones, 1977; Forbes et al., 1981). The speed of the neutral

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line retreat is important to understand the cycle of the substorm (Angelopoulos et al., 1996).

The relative position of the satellite to the neutral line is inferred from in-situ observations of the plasma flow direction and the north-south magnetic field direction (Nishida et al., 1996). The consecutive observation of tailward plasma flow with the southward magnetic field and then earthward plasma flows with the northward magnetic field indicates the retreat of the neutral line (Ueno et al., 1999). However, reliance on single point measurements requires a statistical approach, such as a timing analysis to discuss the speed of the neutral line retreat.

Another approach to estimate the motion of the neutral line is to use counter-streaming field-aligned beams in the PSBL (Forbes et al., 1981; Onsager et al., 1991; Elphic et al., 1995). Elphic et al. (1995) analyzed the PSBL measurements by ISEE-2 and estimated the locations of neutral lines based on the model proposed by Onsager. Estimated distances from the Earth were larger than $60\,R_E$, which correspond to the distant neutral lines.

In this paper we discuss the motion of a near-Earth neutral line by applying the Onsager model to the PSBL ion beam observations made by Geotail. We visually surveyed Geotail observations from 1995 to 2003, and found the longest-duration counter-streaming ion beam event during the 5 October 2000 substorm. Other events were short-lived for tracing the motion or indicated distant neutral lines. In addition to Geotail observations, we use ACE, GOES-8, Interball-Tail, and IMAGE observations to analyze this event from the global point of view.

This paper is structured as follows. In Sect. 2 we present an overview of the event and details of the PSBL observations. In Sect. 3 we estimate the location of a PSBL ion beam source based on ion velocity distributions. In Sect. 4 we discuss the implications of the observations and analysis, and our conclusions are summarized in Sect. 5.

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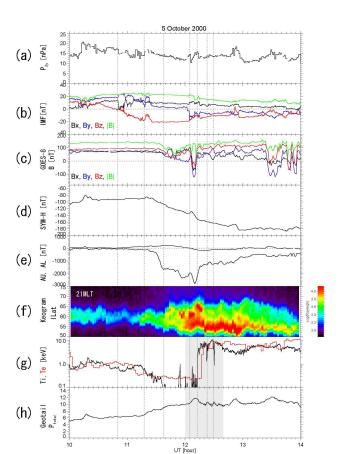


Fig. 1. Summary plot of the 5 October 2000 substorm: (a) ram pressure of solar wind, (b) interplanetary magnetic field, (c) magnetic field at the GOES-8, (d) SYM-H index, (e) AU and AL indices, (f) keogram from the IMAGE/FUV-WIC at 21:00 MLT, (g) ion temperature at the Geotail (black) and electron temperature at the Interball-Tail magnified by 5 (red), (h) total pressure at the Geotail.

2 Observations

2.1 Instrumentation

In this paper we utilize the following satellite data and geomagnetic indices:

- Solar wind and interplanetary magnetic field (IMF) data from the SWEPAM (McComas et al., 1998) and the MAG (Smith et al., 1998) instruments on board the ACE satellite.
- 2. Magnetic field data at geosynchronous altitude measured by the GOES-8 satellite (Singer et al., 1996).
- AL, AU (Davis and Sugiura, 1966), and SYM-H (Iyemori, 1990) indices from the World Data Center for Geomagnetism at Kyoto.
- 4. Two-min resolution auroral images from the FUV (Mende et al., 2000) on board the IMAGE satellite.

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 - Electron and magnetic field data from the ELECTRON (Sauvaud et al., 1997) and the FM-3I on board the Interball-Tail.
 - 6. Low energy particle, magnetic field, and electric field data from the LEP-EA (Mukai et al., 1994), the MGF (Kokubun et al., 1994), and the EFD (Tsuruda et al., 1994) on board the Geotail satellite, respectively. LEP-EA measures 3-D ion and electron velocity distributions from 21 eV to 44 keV with a time resolution of 12 s. All ions are assumed to be protons.

2.2 Overview

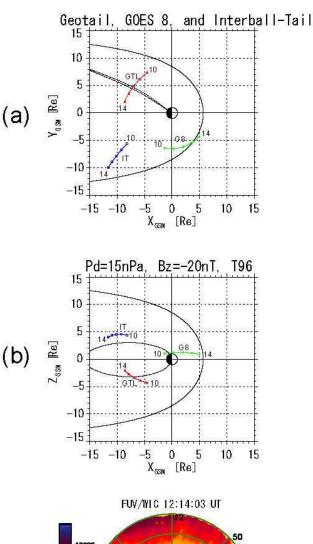
To acquire the perspective of the event, we first take an overview of the ground and satellite observations. Figure 1 plots, from the top, the ram pressure of solar wind, the interplanetary magnetic field (IMF) in GSM measured by the ACE satellite (ACE data are shifted by a 46-min travel time to the Earth), the magnetic field in GSM at geosynchronous altitude measured by the GOES-8 satellite, the SYM-H index, the AU and AL indices, the keogram at 21:00 MLT from the far-ultraviolet Wideband Imaging Camera (FUV-WIC) on board the IMAGE satellite, the ion temperature measured by the Geotail and electron mean energy measured by the Interball-Tail (magnified by 5), the total pressure measured by the Geotail satellite in the magnetotail. The ACE satellite was located around (225, -27, -9) R_E in GSE in the solar wind at 12:00 UT. The locations of GOES-8, Geotail, and Interball-Tail are shown in Fig. 2.

The growth phase of this substorm started around 10:50 UT when the IMF turned from northward to southward. Subsequently, the Interball-Tail escaped from the plasma sheet to the lobe. The onset of the expansion phase was around 11:18 UT after the northward turning of the IMF arrived. At 11:36 UT the auroral oval expanded poleward and Geotail escaped from the plasma sheet to the lobe.

At 12:05 UT the discontinuity of the solar wind arrived. At the same time the GOES-8 experienced the sudden change in the magnetic field, which indicates that the magnetopause crossed the GOES-8. A concurrently observed positive bay of the SYM-H index and a total pressure increase at Geotail (in the tail lobe) suggest the compression of the magnetosphere.

At 12:12 UT the compression signatures ceased and then Geotail and Interball-Tail were engulfed by the plasma sheet. At the same time a poleward boundary intensification (PBI) was observed near the footprint of Geotail (Fig. 2c). Afterward, the AL index started to recover while the main UV oval lasted until 14:00 UT.

In summary, PBIs and plasma sheet expansions were observed after the compression of the magnetosphere at 12:10 UT during the 5 October 2000 substorm.



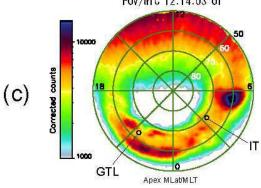


Fig. 2. Locations of the GOES-8, Geotail, and Interball-Tail during 10:00–14:00 UT on 5 October 2000, projected on the GSM (a) XY and (b) XZ plane. Magnetic field line by Tsyganenko T96 model (Tsyganenko, 1995) from the position of Geotail at 12:14 UT and the magnetopause model (Shue et al., 1997) are also shown for reference. (c) The far-ultraviolet image from the IMAGE satellite. Footprints of Geotail and Interball-Tail are mapped with the T96 model.

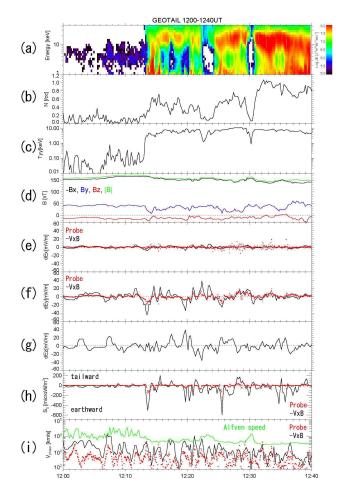


Fig. 3. (a) Energy-time diagram of ion, (b) ion density, (c) ion temperature, (d) magnetic field in GSM, (e-g) three components of electric field perturbation, (h) field-aligned Poynting flux, (i) phase velocity and Alfvén velocity.

2.3 PSBL crossing

In this section we closely examine the plasma sheet expansion from 12:00 UT to 12:40 UT with plasma and electromagnetic field measurements made by Geotail (Fig. 3). Earthward propagating waves and counter-streaming field-aligned ion beams were simultaneously observed.

From the lobe to the plasma sheet, the plasma temperature suddenly increased at 12:13 UT at Geotail. Near the lobe-plasma sheet interface, Geotail observed negative B_y perturbations. These perturbations indicate a pair of field-aligned currents in the outermost PSBL (Ueno et al., 2002).

The existence of a field-aligned current pair is supported by electron measurements. Though electron velocity distributions were almost isotropic (Fig. 4c), excesses of tailward electrons at energies $1.1-4.4\,\text{keV}$ and earthward electrons below $200\,\text{eV}$ are evident (Fig. 4d). These electrons are consistent with the field-aligned currents indicated by the B_y perturbations (Hoshino et al., 2001; Fujimoto et al., 2001; Nagai et al., 2003).

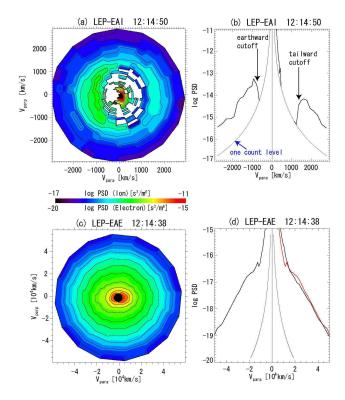
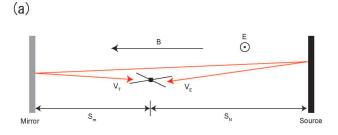


Fig. 4. (a) Ion velocity distribution in the B-C plane at 12:14 UT, (b) 1-D cut of ion velocity distribution at the B-axis, (c) 2-D cut of electron velocity distribution at 12:14 UT, (d) 1-D cut of the electron velocity distribution at the B-axis at 12:14 UT. Red line shows the mirror image of the opposite side.

Magnetic field perturbations, together with electric field perturbations, produced earthward field-aligned Poynting flux (Fig. 3h). The electric field was estimated in two ways: (1) the electric field measured by the probe of the EFD instrument (3-s resolution) with an assumption that Ez=0, (2) the electric field computed from $-V \times B$ (12-s resolution). Perturbations of fields are calculated by subtracting a 10-min running average from the original electric and magnetic field data. The phase velocity computed by the E-to-B ratio did not exceed the Alfvén velocity calculated by the ion density and magnetic field.

During this PSBL crossing, ion velocity distributions consisted of counter-streaming field-aligned beams with lower-energy cutoffs and sometimes cold ions (Fig. 4a), which were probably outflows from the auroral ionosphere. Tailward beams had higher energy and lower flux than earthward beams (Fig. 4b). Estimation of the beam source location in the next section is based on the lower-energy cutoffs of these counter-streaming beams.



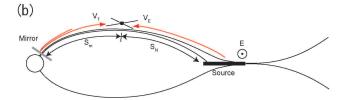


Fig. 5. Schematic pictures of the Onsager model: **(a)** an original configuration, **(b)** a more realistic configuration.

3 Analysis

In this section we estimate the distance from the Geotail to the ion beam source by the lower-energy cutoffs of the counter-streaming ion beams.

3.1 Onsager model

To estimate the distance from the source of ion beams, we utilize the PSBL model proposed by Onsager et al. (1991) which assumes:

- 1. a plasma sheet as a source of PSBL beams,
- 2. a neutral line located tailward of the spacecraft,
- 3. a near-Earth mirror point located earthward of the spacecraft,
- a quasi-steady convection toward the central plasma sheet,
- 5. the conservation of an energy and magnetic moment of particles.

As a consequence of convection toward the center of the plasma sheet and the finite tailward extent of the plasma sheet, which is caused by the neutral line formation, particles with lower cutoff energy are those which originate from the region just earthward of the neutral line.

In the PSBL, these particles come to the spacecraft location in two ways (Fig. 5a). The earthward traveling particles come directly from the plasma sheet to the spacecraft, while the tailward traveling particles are reflected once by the near-Earth magnetic mirror to reach the spacecraft. Flight time T_1

of the particles coming directly from the source to the spacecraft is given by

$$T_1 = S_N / V_E, \tag{1}$$

and for particles reflected by the magnetic mirror, is given by

$$T_2 = (S_N + 2S_m)/V_T,$$
 (2)

where S_N is the distance from the spacecraft to the source and S_m is the distance from the spacecraft to the near-Earth mirror point. V_T and V_E are the tailward and earthward cutoff velocities, respectively. Here we used the constant travel velocity approximation to calculate the flight time.

If T_1 and T_2 are equal to T_{tof} , which is the time for the field line being convected from the source to the spacecraft location, S_N is obtained as

$$S_N = 2S_m/(V_T/V_E - 1).$$
 (3)

 V_E and V_T are known by measurements and S_m is defined as the field line length calculated by Tsyganenko T96 and IGRF magnetic field models. Then we evaluated the downtail distance to the source $|X_N|$ from the Earth as

$$|X_N| = S_N + |X_{SC}|,\tag{4}$$

where $|X_{SC}|$ is the downtail distance of the spacecraft. The estimated value $|X_N|$ is the value at the time when the particles left the source. We define event time, T_{event} , by

$$T_{\text{event}} = T_{\text{obs}} - T_{\text{tof}},\tag{5}$$

where $T_{\rm obs}$ is the time when particles were observed at the spacecraft. Event time means the time when particles were launched from the source. The motion of the source is obtained by $|X_N|$ versus $T_{\rm event}$ plot.

3.2 Uncertainties and limitations

The finite energy resolution of the actual instrument results in an uncertainty of the source location. The lower bound S_N^+ and upper bound S_N^+ are expressed as

$$S_N^+ = 2S_m/(V_T^-/V_F^+ - 1), (6)$$

$$S_N^- = 2S_m/(V_T^+/V_F^- - 1), (7)$$

where $^+$ and $^-$ of V_T and V_E correspond to the upper and lower borders of each energy passband. This is the prime factor for the uncertainties in S_N .

To stabilize our estimate, we took a running average of the cutoff velocities and calculated the downtail distances of the neutral line based on them. If cutoff velocities change coherently, this procedure works properly (Elphic et al., 1995).

Because event time has also an uncertainty which comes from the flight time estimate, we fitted the obtained $|X_N|$ versus event time plot by the two-dimensional, least-squares fitting of Awaya (1983).

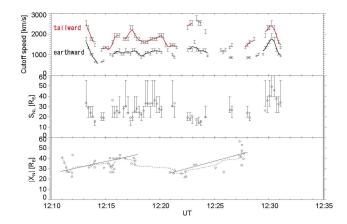


Fig. 6. (a) Time histories of lower-energy cutoffs, (b) estimated distance from the Geotail to the beam source, (c) estimated beam source location against the event time.

There are two limitations for this analysis based on the quasi-steady assumption. First, the temporal resolution of this analysis is regulated by the particle flight time from the source to the spacecraft. Second, this analysis is not suitable for rapid motion of the source due to the finite speed of the particle. Hence, this analysis is applicable for slow and long time scale motion of the particle source.

3.3 Application

Figure 6a shows the cutoff velocities of earthward and tailward beams from 12:13 UT to 12:31 UT. Lower-energy cutoffs were determined as the energy passband of non-zero flux before flux falls below the one-count level (Fig. 4b). Bars indicate the velocity width corresponding to the energy width of each energy step of the LEP instrument. Solid lines are the 36-s (3 samples) window running average of the cutoff velocities.

We could not identify lower-energy cutoffs in the following intervals:

- 1. Before 12:13 UT, Geotail was in the lobe.
- 2. From 12:22 UT to 12:24 UT, tailward low energy cutoffs were sometimes below the one-count level of the instrument, and only earthward cutoffs were obtained.
- 3. From 12:24 UT to 12:28 UT, cold ions sometimes merged with beam components.
- 4. After 12:31 UT, Geotail was embedded in the plasma sheet and flux did not fall before the one-count level.

Figure 6b shows the downtail distance S_N . Diamonds represent the distances from the smoothed cutoff velocities. S_m is set to be 8.5 R_E from the length of the field line calculated by the T96 model. Bars indicate the uncertainties of S_N due to the finite energy resolution of the instrument.

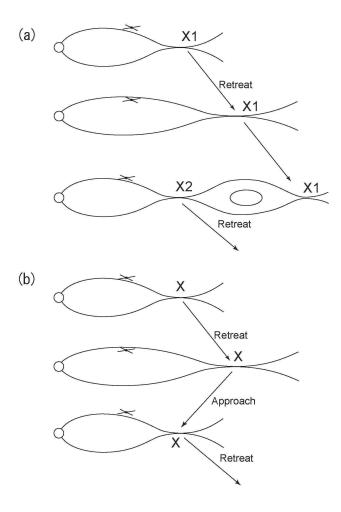


Fig. 7. Schematic picture of two possible neutral line motion scenarios. (a) Zigzag motion between the near-Earth region and midtail region, (b) formation of plural X-lines by successive reconnections.

Figure 6c shows $|X_N|$ against event time (diamonds). The downtail distance of Geotail, $|X_{SC}|$ was about 7.1 R_E during this interval. The dashed line shows the smoothed value of $|X_N|$. We divided points into two groups, from 12:10 UT to 12:18 UT and from 12:21 UT to 12:28 UT. Solid lines were linear fittings for the two intervals.

At 12:11 UT the estimated distance suddenly decreased from $40 R_E$ to $25 R_E$ at the speed of $2200 \,\mathrm{km/s}$. This short and rapid motion is beyond the limitation of the analysis because it is faster than the lower-energy cutoffs of the earthward beam and its duration is comparable to the particle flight time (1–2 min). This apparent motion is probably due to the temporal effect associated with the formation of the neutral line.

From 12:11 UT to 12:17 UT the estimated source retreated at the average speed of $250\pm15\,\mathrm{km/s}$ to about 45 R_E . At 12:17 UT the distance seemed to start to decrease. There is a gap of estimation between 12:18 UT and 12:21 UT, due to the escape of Geotail from the plasma sheet. At 12:21 UT the

distance retreated again at the speed of 300 ± 25 km/s from $25~R_E$ to $45~R_E$.

4 Discussion

Based on the above observations and analysis, we suggest that the source of PSBL ion beams was the near-Earth neutral line formed by the compression of the magnetosphere.

4.1 Existence of magnetic reconnection downtail

Since we took a remote sensing approach, we should discuss whether the source of PSBL ion beams was the reconnection neutral line. We list up the following five supporting evidences:

- 1. The compression of the magnetosphere could work as a trigger of magnetic reconnection (Horiuchi et al., 2001).
- 2. Initial positions of the beam source were around $X=-25 R_E$, where the near-Earth neutral line is considered to be formed at the onset of substorm (Nagai et al., 1998).
- 3. Reconnection creates field-aligned beams in the PSBL by a Speiser-type acceleration of enhanced electric fields around the diffusion region (Hoshino et al., 1998). These field-aligned beams fill the empty lobe flux tubes and cause the plasma sheet expansions (Forbes et al., 1981; Ohtani and Mukai, 2006).
- 4. Simultaneously observed field-aligned currents and earthward propagating waves are consistent with the impulsive reconnection located downtail of the Geotail (Nakamura et al., 2004). These waves are also considered to be the source of the PBIs observed at the Geotail footprint (Wygant et al., 2000; Mende et al., 2003; Dombeck et al., 2005).
- PBIs are related to the plasma flows burst in the plasma sheet (Lyons et al., 2002). These flows are generally considered as the major manifestation of magnetic reconnection.

4.2 Retreat speed of neutral line

As for the retreat speed, we can compare our result to the retreat speed of a plasmoid. The average retreat speed of a plasmoid is $350 \, \text{km/s}$ at $30 \, R_E$ (Ieda et al., 1998). This value is slightly larger than our results. If the neutral line retreat follows the ejection of a plasmoid, this magnitude relation of the retreat speed is reasonable.

The speed of the neutral line motion for an asymmetric boundary condition is derived as 10% of the local Alfvén speed from the numerical simulation (Fujimoto, 2005). If we assume a proton density of 0.1/cc and a magnetic field of $40 \, \mathrm{nT}$ at $X = -30 \, R_E$ (This event was under high solar wind

ram pressure and typical parameters could be inadequate), then the Alfvén speed is 2800 km/s and thus predicts the retreat speed of 280 km/s. This is comparable to our estimated values.

4.3 Interpretations of neutral line motion

From Fig. 6c there are two interpretations of the neutral line motion. One is that two near-Earth neutral lines were created and retreated (Fig. 7a). The other is that only one near-Earth neutral line went back and forth between the near-Earth region and mid-tail region (Fig. 7b). If the former is the case, the recurrent formation of the neutral lines forms a long plasmoid between the neutral lines. This could be the mechanism to create $20 R_E$ long plasmoids seen in the distant tail (Slavin et al., 2003).

The latter case is not denied observationally. Ohtani et al. (2004) reported the case study of the earthward moving neutral line. They discussed that the neutral line they analyzed did not reach the lobe and was confined in the plasma sheet

If the neutral line returns at $X=-60~R_E$, a maxima of the occurrence rate is expected there, due to the deceleration before the earthward motion. In fact, distribution of the neutral lines statistically obtained by Ueno et al. (1999) has a maximum occurrence rate at Xgsm= $-60~R_E$. A possible cause of the neutral line deceleration is the formation of the stagnant plasmoid under the southward IMF condition (Nishida et al., 1996). This event also occurred under the strong southward IMF condition.

Two cases give totally different results on the plasmoid release and hence mass and energy loss in the tail plasma sheet. The effect of IMF B_z on the plasmoid evolution is worth further investigation. In both cases, we would like to point out that the 10-min duration of retreat motion is comparable with the typical duration of BBFs and PBIs. This might suggest that the motion of neutral line is a part of the global ULF oscillation of the magnetotail (Lyons et al., 2002).

4.4 Future application of the Onsager model

In this paper we investigated the motion of near-Earth neutral line during the recovery phase of a substorm. Another interesting application of the Onsager model is traveling compression regions (TCRs) during the expansion phase of a substorm. These TCRs are likely related to the reconnection in the near-Earth tail region (Owen et al., 2005). There is a possibility to reveal the motion of the neutral line at the start of the substorm expansion by applying the Onsager model to the PSBL crossing associated with TCRs. In this case, high time resolution measurements of the electrons are necessary because the typical duration of TCRs in the near-Earth tail region is as short as 35 s (Slavin et al., 2005). We will hopefully investigate this issue in a future study.

5 Conclusions

From the lower-energy cutoffs of PSBL ion beams, we estimated the motion of the near-Earth neutral lines during the 5 October 2000 substorm. The retreat speeds of the neutral line were 250 and 300 km/s. The estimated motion of the neutral line suggests either the creation of plural X-lines or the oscillation of one neutral line between near-Earth and the midtail. Two cases are totally different for a plasmoid release but not distinguishable by our remote sensing analysis, thereby making coordinated multi-satellite studies necessary.

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