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Non-Gaussian velocity shears in the environment of low mass dense cores[★]

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Abstract. We report on a novel kind of small scale structure in molecular clouds found in IRAM-30m and CSO maps of ¹²CO and ¹³CO lines around low mass starless dense cores. These structures come to light as the locus of the extrema of velocity shears in the maps, computed as the increments at small scale (~0.02 pc) of the line velocity centroids. These extrema populate the non-Gaussian wings of the shear probability distribution function (shear-PDF) built for each map. They form elongated structures of variable thickness, ranging from less than 0.02 pc for those unresolved, up to 0.08 pc. They are essentially pure velocity structures. We propose that these small scale structures of velocity shear extrema trace the locations of enhanced dissipation in interstellar turbulence. In this picture, we find that a significant fraction of the turbulent energy present in the field would be dissipating in structures filling less than a few % of the cloud volume.

Key words. ISM: evolution – ISM: kinematics and dynamics – ISM: molecules – ISM: structure – turbulence

1. Introduction

Star formation proceeds at vastly different rates, in space and time, within a given galaxy and from one galaxy to another. These rates are known to be governed by the local conditions prevailing in dense and cold gas, but also depend on large scale environments, up to extragalactic scales. The only two processes, together with rotation, able to mitigate the effects of gravity are magnetic fields and supersonic turbulence because they involve energies of the same order of magnitude as the gravitational energy in the densest phases of the interstellar medium (ISM). More precisely, turbulent energy, because of its steep power spectrum, can stabilize the largest masses, first prone to gravitational instability, a property not shared by thermal energy which is scale-free (Panis & Pérault 1998). On the other hand, it has also been proposed that supersonic turbulence might trigger star formation in shocks (Klessen et al. 2000). The possible role of turbulence in the star formation process is the motivation behind the plethora of recent studies on interstellar turbulence.

Important insights to the field have been provided by the determination of the multi-scale properties of interstellar clouds and their comparison to the scaling laws of turbulence. A broad variety of statistical tools have been used: wavelet analysis (Gill & Henriksen 1990; Langer et al. 1993), Δ -variance

(Bensch et al. 2001), auto-correlation function (Kleiner & Dickman 1984, 1985; Dickman & Kleiner 1985; Pérault et al. 1986), structure functions (Miesch & Bally 1994; Miesch & Scalo 1995; Miesch et al. 1999; Padoan et al. 2003a), analysis in terms of fractal structures (Bazell & Désert 1988; Falgarone et al. 1991; Stutzki et al. 1998). The principal component analysis has been used to diagnose the large-scale flows of atomic gas into which turbulence in molecular clouds is embedded (Brunt 2003). Observations of dense cores and their environment have revealed a break of the scaling properties of molecular clouds at the scale of the dense cores (Falgarone et al. 1998; Goodman et al. 1998). Direct numerical simulations of compressible turbulence have also been used to compare the real observables of the interstellar clouds to those simulated. The first attempt by Falgarone et al. (1994), using the high resolution simulations of mildly compressible turbulence of Porter et al. (1994) has been followed by more sophisticated comparisons such as those based on the spectral correlation function method (Rosolowsky et al. 1999). Numerical simulations of turbulence including magnetic fields, e.g. Ostriker et al. (2001), and self-gravity (Klessen et al. 2000; Heitsch et al. 2001) provided further support to the fact that interstellar turbulence bears many of the statistical properties of supersonic magneto-hydrodynamical (MHD) turbulence, as simulated.

One major surprise brought to the field by direct numerical simulations of MHD turbulence has been that magnetic fields do not delay the dissipation of supersonic turbulence (Mac Low et al. 1998). This unexpected result calls for further numerical

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[★] Appendices A, B and Figs. 5, 13–15, 17 and 18 are only available in electronic form at <http://www.edpsciences.org>

and observational approaches. The present paper is observational: it is an attempt to disclose kinematic signatures of turbulent dissipation in molecular clouds.

To discuss what these kinematic signatures may be, we first need to briefly recall the main drivers of dissipation in interstellar turbulence, assuming infinite conductivity and thus neglecting Ohmic dissipation. The primary source of dissipation of supersonic turbulence is shocks, but shock interaction generates vorticity, as do non-planar shocks (Porter et al. 1994). Since the divergence of the velocity field eludes direct detection in space, shocks cannot be detected by this kinematic signature, and remanent vorticity is a plausible signature of fossil shocks. Viscous dissipation is also present and is due primarily to elastic collisions. Dissipation due to neutral–neutral collisions follows the shear of the velocity field, and therefore the vorticity (Landau & Lifshitz 1987), and that due to ion–neutral collisions increases with the drift velocity of the ions relative to the neutrals (Kulsrud & Pearce 1969). This drift causes a force on the ions which, over timescales longer than the ion–neutral collision time (Zweibel 1988), is balanced by the Lorentz force $(\nabla \times \mathbf{B}) \times \mathbf{B}$. The dissipation driven by ion–neutral collisions therefore involves $\mathbf{J} = \nabla \times \mathbf{B}$, the current density.

Now, both vorticity, in hydrodynamical turbulence, and current density, in plasma turbulence, are known to be intermittent in space and time. Laboratory experiments of incompressible turbulence show the formation of transient long and thin coherent vortices at the edge of which the velocity shear is so large that a significant fraction of the viscous dissipation occurs there (Douady et al. 1991). In plasma turbulence, intermittency also exists and is at the origin of non-Gaussian probability distribution functions (PDFs) of current density. It is seen in Tokamak plasma turbulence (Wang et al. 1999). Observations in the solar wind reveal an intermittent dissipation as well. The intermittency of dissipation of plasma turbulence has been modelled by Politano & Pouquet (1995), who propose that the dissipative structures are sheetlike structures of intense current density, in agreement with recent 3-dimensional numerical simulations of MHD turbulence (Politano et al. 1995; Biskamp & Müller 2000). A recent attempt at estimating the sizescales of dissipation in incompressible MHD turbulence (Cho et al. 2002) shows that magnetic structures develop at scales much smaller than the viscous damping scale.

Turbulence in molecular clouds is neither laboratory turbulence nor plasma turbulence, but the above elements suggest that dissipation follows the vorticity and the current density, and that these quantities are intermittent, i.e. exhibit large fluctuations at small scales. The origin of currents in molecular clouds is not known, but is likely associated to differential rotation within clouds at the origin of toroidal or helical fields (Joulain et al. 1998). Such helical fields have been inferred from various observations of molecular clouds (Hanawa et al. 1993; Joulain et al. 1998; Carlqvist et al. 1998; Falgarone et al. 2001; Matthews et al. 2002, 2001). These fields have been invoked to explain the polarization patterns of the dust continuum emission of filaments of matter (Harjunpää et al. 1999; Fiege & Pudritz 2000c) and their gravitational stability (Fiege & Pudritz 2000a,b). Enhanced vorticity in molecular clouds should therefore be a good tracer of several major dissipative

processes of turbulence in weakly ionized molecular clouds. Its kinematic signature, if available in spectral observations of molecular clouds, is essential to search for because it may trace fossil shocks, regions of large viscous dissipation in the neutrals, or local gas differential rotation at the origin of currents.

In this paper, we analyze the velocity structure of several fields in molecular clouds in order to locate and characterize the regions of enhanced vorticity, likely to be regions of enhanced dissipation. We focus on the environment of low mass dense cores almost thermally supported, where dissipation is anticipated to occur, or has occurred in a recent past. Section 2 gives the method used to trace vorticity. Section 3 is a description of the target fields. The statistical analysis of the velocity fields, and the manifestations of their non-Gaussian features are given in Sect. 4. The spatial distribution of these non-Gaussian events is compared to that of the dense gas in Sect. 5. The possible biases of the method are discussed in Sect. 6. The implications of this study in terms of turbulence dissipation are discussed in Sect. 7.

2. Method for measuring vorticity

Measuring vorticity in interstellar turbulence is not straightforward because (1) the only velocity component accessible to measurement is its projection on the line of sight (v_x), and (2) measured spatial variations are limited to those in the plane of the sky (y, z). We thus have access only to incomplete components of the vorticity ($\omega_\zeta = \partial v_x / \partial \zeta$ where $\zeta = y$ or z). By analyzing the outputs of numerical simulations of mildly supersonic turbulence by Porter et al. (1994) and Lis et al. (1996) have shown that vorticity extrema may be localized in a map of molecular lines with high enough spectral resolution. Instead of computing the vorticity (an essentially inaccessible quantity in astronomy), they study statistical properties of line centroid velocity increments. More precisely, the centroid velocity (C) of a line is its first order moment

$$C = \int T(v_x) v_x dv_x / \int T(v_x) dv_x. \quad (1)$$

It is only in the optically thin case, that a line can be interpreted as the probability distribution function of the radial velocity along the observed line of sight. In this case, the centroid velocity is just the column density weighted average of the velocity along the line of sight. Centroid velocity increments associated to a lag l , are the differences of centroid velocities for any pair of lines of sight separated by the distance l in the plane of sky. Thus, if \mathbf{r} is the vector position and \mathbf{l} is a vector of modulus l lying in the plane of sky, centroid velocity increments are defined as

$$\delta C(\mathbf{r}, \mathbf{l}) \equiv C(\mathbf{r} + \mathbf{l}) - C(\mathbf{r}). \quad (2)$$

Lis et al. (1996) show that the PDFs of centroid velocity increments have non-Gaussian wings at small lags, those wings gradually disappearing when the lag increases. More interestingly, they show that the spatial distribution of the positions populating these non-Gaussian wings is the same as that of the extrema of $\langle (\nabla \times \mathbf{v})_y \rangle^2 + \langle (\nabla \times \mathbf{v})_z \rangle^2$, where the brackets hold for

Table 1. Characteristics of the maps: telescope, line observed, half-power beamwidth, sampling step (or pixel size) in arcsec and in AU at the distance of the sources (150 pc), the effective resolution in pixels and in AU.

Telescope	Line	HPBW	Pixel size	Resolution
IRAM (30 m)	CO ($J = 1-0$)	22''	7.5'' (1125 AU)	3 pixels (3375 AU)
CSO (10.4 m)	CO ($J = 2-1$)	28''	16'' (2400 AU)	2 pixels (4800 AU)

Table 2. Characteristics of the mapped fields: sizes in pc, number of spectra and average signal-to-noise ratios in peak temperature and integrated area (The associated uncertainties are the rms scatter of the values over the maps).

Source	Line	Field size (pc)	Nb. Spectra	SNR	SNR_a
Polaris	^{12}CO ($J = 1-0$)	0.26×0.35	3300	11 ± 3	51 ± 13
Polaris	^{13}CO ($J = 1-0$)	0.22×0.30	1650	12 ± 4	40 ± 12
L1512	^{12}CO ($J = 1-0$)	0.22×0.44	3200	15 ± 3	46 ± 8
L1512	^{13}CO ($J = 1-0$)	0.22×0.36	2560	23 ± 5	61 ± 13
L1512	^{12}CO ($J = 2-1$)	1.05×1.12	8300	10 ± 4	27 ± 12

the line-of-sight average. In the optically thin case, this result simply follows from the linearity of the operations of derivation, difference and summation along a line of sight. The largest centroid increments thus trace a subset of the regions of largest vorticity, integrated along the line of sight. It has to be appreciated here that the averaging along the line of sight occurs first, over vorticity components which have a sign. An extremum of the above quantity means the occurrence along the line of sight of one (or a few) events far above the average.

This method has been challenged by Klessen (2000) who computes the PDFs of line centroid increments (Δv -PDFs) in numerical simulations of compressible turbulence with and without self-gravity. He finds that, when computed in simulations of decaying turbulence, Δv -PDFs have shapes in disagreement with those observed. He shows that the inclusion of self-gravity in the simulations leads to better agreement with the Δv -PDFs observed in molecular clouds. This discussion is valuable but its bearing is somewhat limited by the small size of the simulations (64^3). In addition, the confrontation with the observations is restricted to star forming regions. It is therefore not surprising that the agreement with the numerical simulations be better when self-gravity is included in the simulations. Klessen concludes that “one should not rely on analyzing velocity PDFs alone to disentangle the different physical processes influencing interstellar turbulence”. We fully agree with him. Our scope here is far more limited. We search for regions of enhanced vorticity in turbulent interstellar clouds far from star forming regions, with the perspective of tracing bursts of dissipation of turbulence and characterizing them, on the basis of their statistical properties, spatial distribution and morphology.

The spectral correlation function (SCF) introduced by Rosolowsky et al. (1999) is a general tool that is able to detect velocity and/or column density variations through an analysis of line shape differences. The flexibility of the SCF formalism makes it easy to use for direct comparison between observations and simulations (Padoan et al. 2003b). For instance, Ballesteros-Paredes et al. (2002) use the local form of

the SCF to detect small scale velocity structures in HI observations of the north celestial pole loop. However, as recognized by those authors, this method cannot disentangle between the various origins of the small scale velocity variations (e.g. caustics versus shocks). The method we propose here is complementary as centroid velocity increments directly relate to the velocity shear (and thus vorticity) in the optically thin case.

3. Description of the observed fields

The fields analyzed in this paper fulfill several requirements: (1) the observed maps are large enough to provide significant statistics on the velocity field, (2) the fields are far from star forming regions, in order not to include internal energy sources (outflows, HII regions in expansion, ...) in the statistics, (3) they comprise low mass, almost thermally supported dense cores.

The two fields are nearby regions of low average column density, except in the small projection area of the dense cores. One, MCLD123.5+24.9 (hereafter referred to as Polaris), is located in the high latitude cirrus cloud of the Polaris Flare and the other, (hereafter L1512), is located at the eastern edge of the Taurus-Auriga complex. The average H_2 column densities (at the arcmin scale) deduced from ^{12}CO observations and/or star counts (Cambr esy et al. 2001) in these fields are respectively $N(\text{H}_2) \sim 10^{21}$ and $2.5 \times 10^{20} \text{ cm}^{-2}$ or about 1 and 0.25 mag of visual extinction, respectively. For comparison, the visual extinction in the Taurus molecular cloud reaches 33 mag in the most opaque regions like TMC1 or L1495 (Padoan et al. 2002), about one hundred times more opaque than the transparent areas that we analyze in this paper.

These fields have been mapped in the two lowest rotational transitions of ^{12}CO and ^{13}CO as well as in C^{18}O ($J = 1-0$) at the IRAM-30m telescope (Falgarone et al. 1998). A larger field around L1512 has been mapped with the 10.4 m antenna of the Caltech Submillimeter Observatory (CSO) in the ^{12}CO ($J = 2-1$) line and is described in Falgarone et al. (2001). The characteristics of the observations are summarized in Tables 1 and 2. Note that all the maps are Nyquist sampled

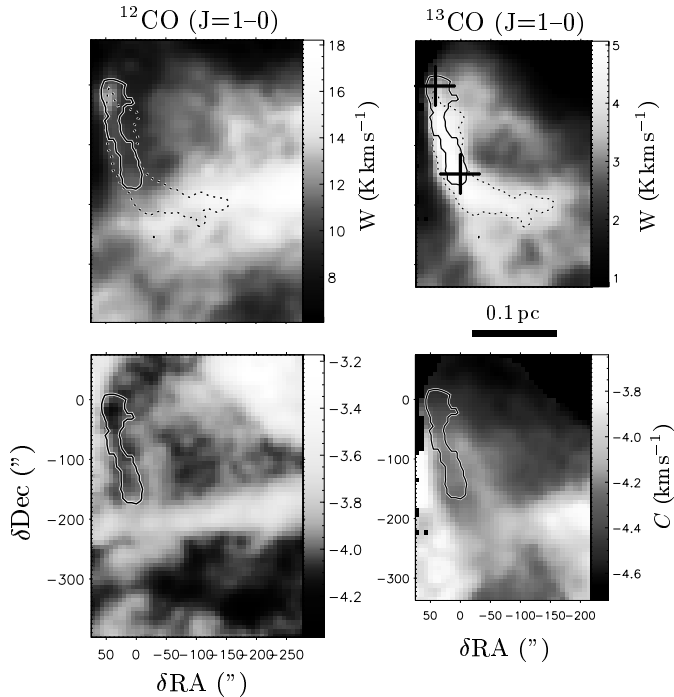


Fig. 1. Maps of the Polaris (MCLD123+24.9) field in the $^{12}\text{CO}(J=1-0)$ (left) and $^{13}\text{CO}(J=1-0)$ (right) lines: (top) line integrated areas, (bottom) line centroid velocities. The offsets are in arcsec respectively to the (0,0) position: $l_{\text{II}} = 123.68^\circ$, $b_{\text{II}} = 24.93^\circ$. The two upper-left corners of the ^{12}CO and ^{13}CO maps coincide. The dashed contour level at 4 K km s^{-1} in the $W_{13}\text{CO}$ map is meant to localize the region of largest column density and is reproduced in the $W_{12}\text{CO}$ map. The black contour is the 0.6 K km s^{-1} -level of the CS(3–2) map (see Fig. 4) and is meant to localize the dense core in the rest of the paper. The two crosses show the positions of the HC_3N peaks (Heithausen et al. 2002). The linear scale is shown.

or better. The average signal-to-noise ratios (and their rms dispersion within the maps) are given in the last two columns of Table 2, for the line peak temperatures and the line integrated areas. These numbers illustrate the high quality of the data. The maps of the line integrated areas (zero order moment), and line centroid velocities (first order moment) are displayed in Figs. 1 through 3. We have inserted the high angular resolution IRAM data at their relevant position in the large scale CSO field of L1512 to illustrate the agreement between the two data sets (Fig. 3).

In spite of their low average extinction, each field harbours a low mass dense core. The dense core in the Polaris field has been mapped in the CS(2–1), (3–2) and (5–4) lines (Heithausen 1999) and in a number of molecular species at a few positions only by Gerin et al. (1997). Recent observations of the dust thermal emission (Heithausen et al. 2002) show a clear peak of submillimeter dust continuum emission coinciding with the C^{18}O peak of Falgarone et al. (1998). Well-defined and barely resolved peaks of HC_3N emission are found to be clearly displaced relative to the dust continuum maximum. They are indicated by two crosses in the ^{13}CO map of Fig. 1. H_2 densities as large as a few 10^5 cm^{-3} in this dense core are derived from the multi-line analysis of Heithausen et al. (2002)

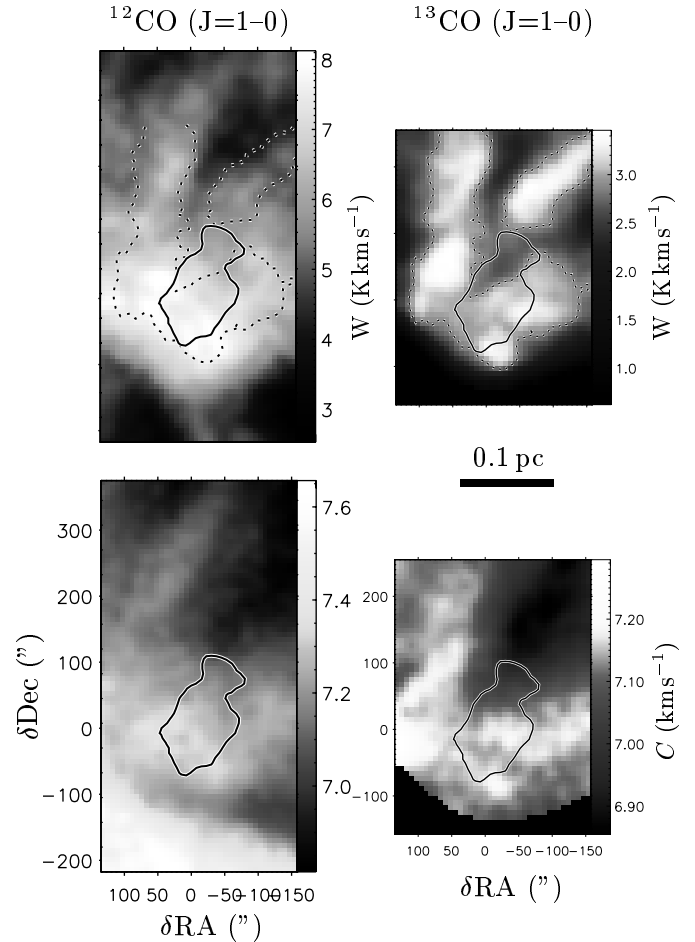


Fig. 2. Same as Fig. 1 for the L1512 field observed at IRAM-30m. Offsets are in arcsec relative to the (0,0) position: RA(1950) = $05^{\text{h}}00^{\text{m}}54.5^{\text{s}}$, Dec(1950) = $32^{\circ}39'00''$. The dashed contour level (3 K km s^{-1} in the $W_{13}\text{CO}$ map) is reproduced in the $W_{12}\text{CO}$ map and helps localize the regions of largest column density. The dense core here (black contour) is identified with the 0.15 K km s^{-1} contour of the $\text{HCO}^+(3-2)$ emission (see Fig. 4) and is reproduced in the rest of the figures. The linear scale is shown.

and Gerin et al. (1997). Our unpublished map of CS(3–2) integrated emission is displayed in Fig. 4 and delineates the region of largest density. The contour level at 0.6 K km s^{-1} of the CS(3–2) emission is shown in Fig. 1 to help localize the dense core. The $^{13}\text{CO}(J=1-0)$ contour level of 4 K km s^{-1} is also drawn in Fig. 1 to indicate the region of column density larger than $\sim 4 \times 10^{21} \text{ cm}^{-2}$.

In L1512, the dense core has been mapped in CS(2–1) by Fuller (1989) over several arcmin. Lee et al. (2001) have mapped it over a more restricted area in the CS(2–1) and N_2H^+ lines. A map of the $\text{HCO}^+(3-2)$ emission has been performed at the CSO in November 2002. The $\text{HCO}^+(3-2)$ line integrated emission is displayed in Fig. 4. As in other dense cores, the N_2H^+ emission is more concentrated than the CS emission and the N_2H^+ and $\text{HCO}^+(3-2)$ emissions have very similar boundaries. Again, in the following, we adopt the 0.15 K km s^{-1} contour of the $\text{HCO}^+(3-2)$ emission as the boundary of the L1512 dense core. It is drawn in Figs. 2 and 3 to help locate the dense core. The maps of Figs. 1 to 3

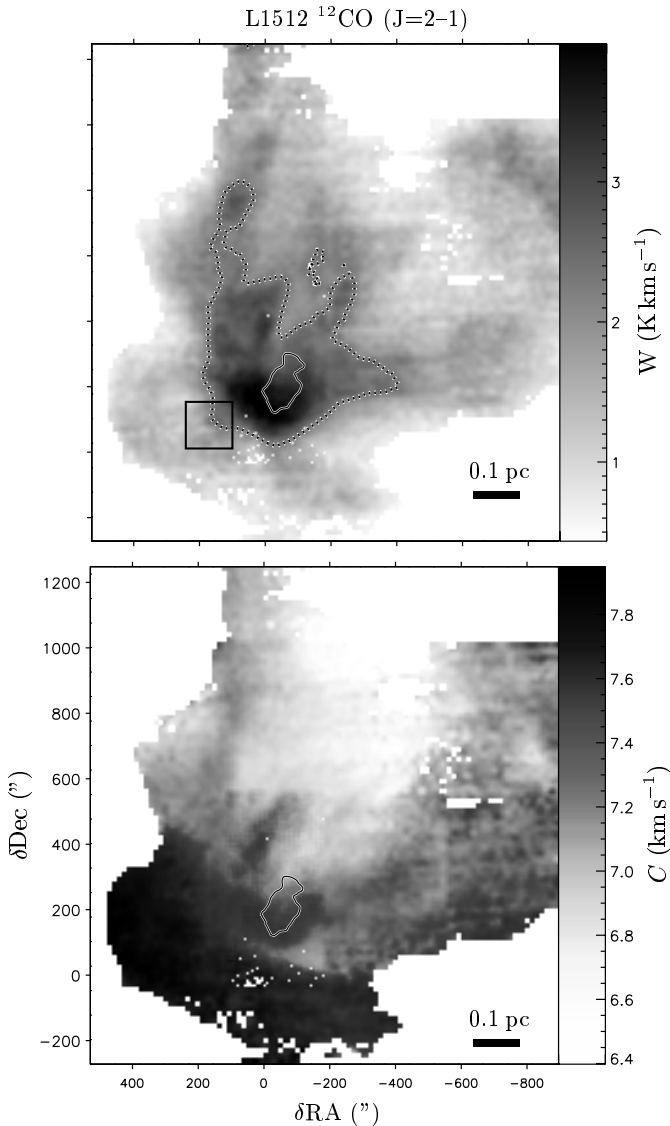


Fig. 3. Maps of the $^{12}\text{CO}(J=2-1)$ line integrated area W (top) and line centroid velocities C (bottom) of the L1512 field observed at the CSO. In both panels, a map of the same quantities computed with the high angular resolution data of the IRAM-30m is inserted in the CSO map at the relevant position. The dense core is identified as in Fig. 4 (black contour). The dashed contour corresponds to $W_{^{12}\text{CO}} = 2.5 \text{ K km s}^{-1}$ and will be used in Fig. 11. The box is the location of the spectra map shown in Fig. 5. The (0,0) position of the CSO map is $\text{RA}(1950) = 05\text{h}00\text{m}54.5\text{s}$, $\text{Dec}(1950) = 32^\circ 36'00''$. The linear scale is shown.

confirm that dense cores, as traced by molecular lines such as CS, N_2H^+ or HCO^+ at millimeter wavelengths, are embedded in larger structures of moderate column density, traced by the ^{13}CO lines, and further down in column density, by the ^{12}CO lines. Note that the ^{12}CO integrated emission in both fields is far from being isotropically distributed around the dense cores.

The analysis presented in the next sections has been conducted on the ^{12}CO and ^{13}CO data sets, when available. ^{12}CO is the only available tracer of molecular gas down to H_2 column densities as low as a few 10^{20} cm^{-2} . Reaching these limits

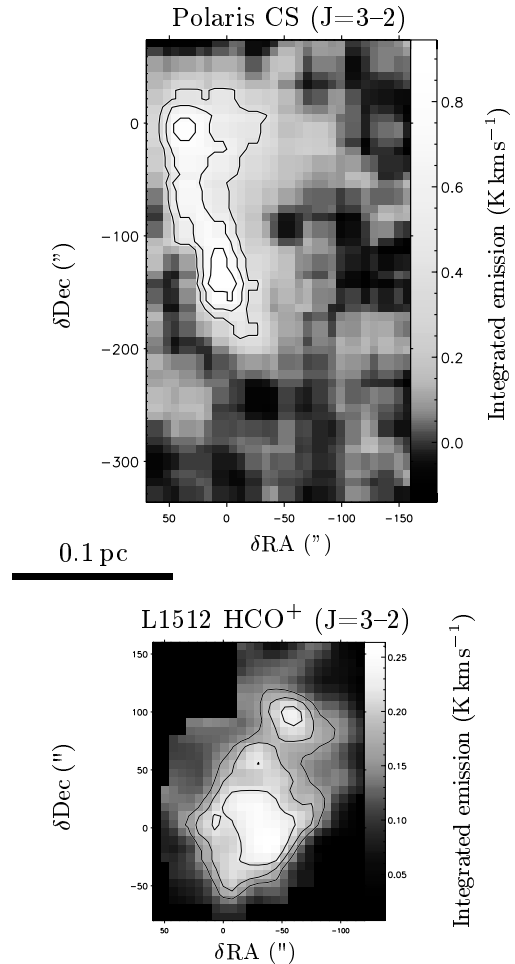


Fig. 4. Map of integrated emission of the dense cores: IRAM-30m map of Polaris in the CS(3-2) line (top) and CSO map of L1512 in the $\text{HCO}^+(3-2)$ line (bottom).

is mandatory in order to be able to make a link with statistical works performed on HI data (e.g. Miville-Deschênes et al. 2003; Brunt 2003), and to obtain large maps avoiding star forming regions. However, the large opacity of the ^{12}CO line prevents the sampling of all the gas on the line of sight, most severely in regions close to the dense cores. The ^{13}CO lines are thus used to analyze the gas in regions of intermediate column density between the transparent cloud edges traced by ^{12}CO only and the dense cores. The areas available for statistical analysis in this line are smaller than those available in the ^{12}CO line because of the different abundances of the two isotopomers. The two lines are therefore complementary.

4. Statistical analysis

4.1. Computation of line centroids and their increments

Centroid velocities have been computed according to the algorithm described in Pety (1999). The line window used to compute the line centroid is critical (see Appendix A.1) and we have adjusted it locally to maximize the signal-to-noise ratio of the integrated area defined as $SNR_a = \sum_i^m T_i / \sqrt{m\sigma}$

Table 3. Quantities used to normalize the δC -PDFs for each field and line and for the lag $l = 3$ pixels: the offset applied to the increments to get a centered PDF (Col. 3) and the standard deviation of the δC values (Col. 4). The ratio of these two quantities is given in Col. 5.

Field	Line	$\langle \delta C_3 \rangle$ km s ⁻¹	$\sigma_{\delta C_3}$ km s ⁻¹	$\frac{\langle \delta C_3 \rangle}{\sigma_{\delta C_3}}$
Polaris	¹² CO ($J = 1-0$)	8.3×10^{-3}	0.11	0.07
Polaris	¹³ CO ($J = 1-0$)	1.5×10^{-2}	0.10	0.15
L1512	¹² CO ($J = 1-0$)	1.2×10^{-2}	0.05	0.24
L1512	¹³ CO ($J = 1-0$)	1.8×10^{-3}	0.04	0.04
L1512	¹² CO ($J = 2-1$)	1.3×10^{-2}	0.09	0.015

where T_i are the temperatures of the m channels within the window and σ the rms noise level of the spectrum. As the optimal window is expected to vary softly from one position to another, we smoothed the variations of the window edges with a median, moving boxcar filter of size 5×5 pixels. The first moment is then computed on each spectrum using the optimal window found for its location. Figure 5 illustrates the optimal windows found by this algorithm for a subset of the CSO map. This method also ensures that there is no significant emission out of the final optimal window. The line centroid increments are then computed over a lag l expressed in pixels of the map. These increments are computed only between spectra that (i) have a SNR_a greater than 10.0 and (ii) have at least 4 neighbors with $SNR_a \geq 10.0$. We show in Appendix A.2 how this threshold $SNR_a = 10.0$ ensures that our statistical analysis of velocity increments is not contaminated by thermal noise. For a given lag l , the increments are computed for half the possible orientations of the lag vector because if all the orientations are kept for each lag, each increment in the map appears twice with opposite signs. In the case of $l = 3$ pixels, for instance, there are 16 different orientations, defined by the 16 neighbors of any position lying within the rings of radii 2.5 and 3.5 pixels. We kept only 9 different orientations for each position in the maps. Note that the smallest significant lag is $l = 2$ pixels for the CSO map and $l = 3$ pixels for the IRAM maps because the maps are oversampled.

4.2. Non-Gaussian wings in the centroid increment PDFs

We have built the PDFs of the computed centroid increments for four different lags in each field and for each line. Hereafter, the PDFs of centroid increments will be called δC -PDFs or increment PDFs. They may also be seen as shear-PDFs since, as said in Sect. 2, the measured projection of the velocity is orthogonal to the plane in which displacements are measured (i.e. the plane of the sky).

Figures 6 and 7 show the evolution of the δC -PDFs with increasing values of the lags for each map. The PDFs are normalized to zero mean and unit dispersion and a normalized Gaussian PDF with the same dispersion is shown, for comparison, as a dotted line. The quantities used for this normalisation

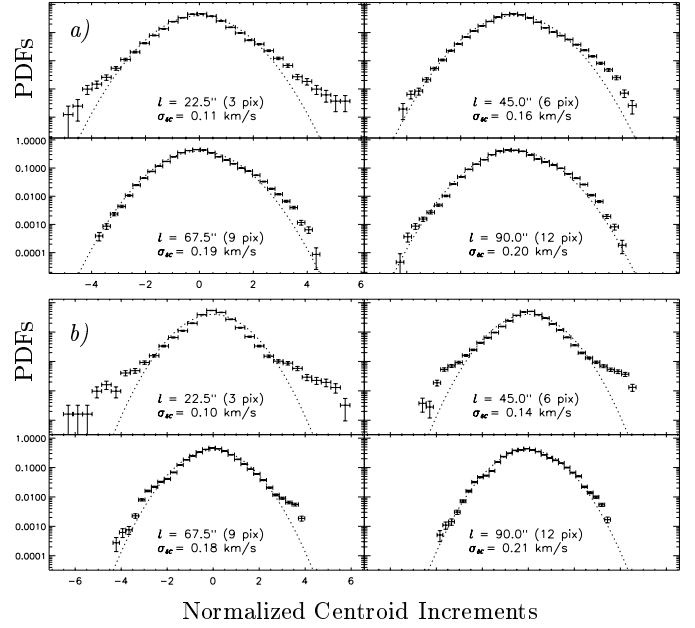


Fig. 6. PDFs of centroid velocity increments for increasing lag values l given in arcsec and pixels within each frame. The PDFs have been computed from **a)** the ¹²CO ($J = 1-0$) and **b)** ¹³CO ($J = 1-0$) spectra maps of Polaris. The centroid increment scales have been normalized so that $\langle \delta C \rangle = 0$ and $\sigma_{\delta C} = 1$. The actual velocity dispersion of each distribution $\sigma_{\delta C}$ is given in km s⁻¹ in each frame. The dotted curves represent a Gaussian of zero mean and unit standard deviation.

are given in Table 3 for each map and a lag of 3 pixels. $\langle \delta C_3 \rangle$ is the offset applied to the increments at lag $l = 3$ pixels to get a centered PDF and $\sigma_{\delta C_3}$ is the standard deviation of the increments, for the same lag. The ratios $\frac{\langle \delta C_3 \rangle}{\sigma_{\delta C_3}}$ are given in Col. 5: they are all significantly smaller than unity. We therefore feel confident that the normalisation we apply does not affect the estimate of the δC -PDFs computed at small lags. The error bars drawn simply reflect the number of elements in each bin of the PDF. Note that the statistical significance of the largest increments is best for the CSO field because the number of *independent* spectra available in this map is about nine times larger than in the IRAM maps. This illustrates the importance of the size of the maps for this kind of statistical study where departures from Gaussianity are searched for and occur only with probabilities close to 10^{-2} or below. In Appendix B, we compute the largest statistically significant lags for each map and two values of the numbers of bins in the PDFs, 30 and 8 bins. These numbers confirm that PDFs built with 30 bins, as in Figs. 6 and 7, are statistically significant up to lags of ~ 6 pixels for the small fields and 12 pixels for the large one. They also show that the broad features in the PDFs, such as non-Gaussian wings, spreading over 4 bins or more, are significant up to lags of 12 pixels for all the fields.

All the sets of increment-PDFs exhibit non-Gaussian wings with departure from Gaussianity being more pronounced as the lag becomes small. The effect is the most visible in the large scale L1512 field, because the large number of independent spectra in that field reduces effects due to the poor sampling at large lags and makes the PDFs there the most symmetrical.

Table 4. Influence of large scale velocity gradients. Difference between the observed extrema of line centroids in each map (Col. 3), the associated scale s (Col. 4) and the corresponding observed large scale velocity gradients (Col. 5), the contribution of this large scale gradient to a lag $l = 3$ pixels (Col. 6), the ratio of this contribution to the internal velocity dispersion at the same scale (Col. 7), the exponent of the scaling of the standard deviation of the centroid increments with the lag at which they are computed (Col. 8), the velocity gradients at scale s inferred from this scaling and ascribed to turbulence (Col. 9).

Field	Line	$ C_{\min} - C_{\max} $ km s ⁻¹	s "/pc	∇v_{obs} km s ⁻¹ pc ⁻¹	δv_3 km s ⁻¹	$\frac{\delta v_3}{\sigma_{\delta C_3}}$	ζ_2	∇v_{turb} km s ⁻¹ pc ⁻¹
Polaris	¹² CO(1–0)	1.1	500/0.37	3.0	0.04	0.35	0.45	3.8
Polaris	¹³ CO(1–0)	1.0	400/0.29	3.4	0.06	0.60	0.5	4.5
L1512	¹² CO(1–0)	1.2	600/0.44	2.7	0.05	1.0	0.62	2.7
L1512	¹³ CO(1–0)	0.4	300/0.22	1.8	0.03	0.77	0.7	3.5
L1512	¹² CO(2–1)	1.4	1200/0.9	1.5	0.03	0.33	0.5	1.6

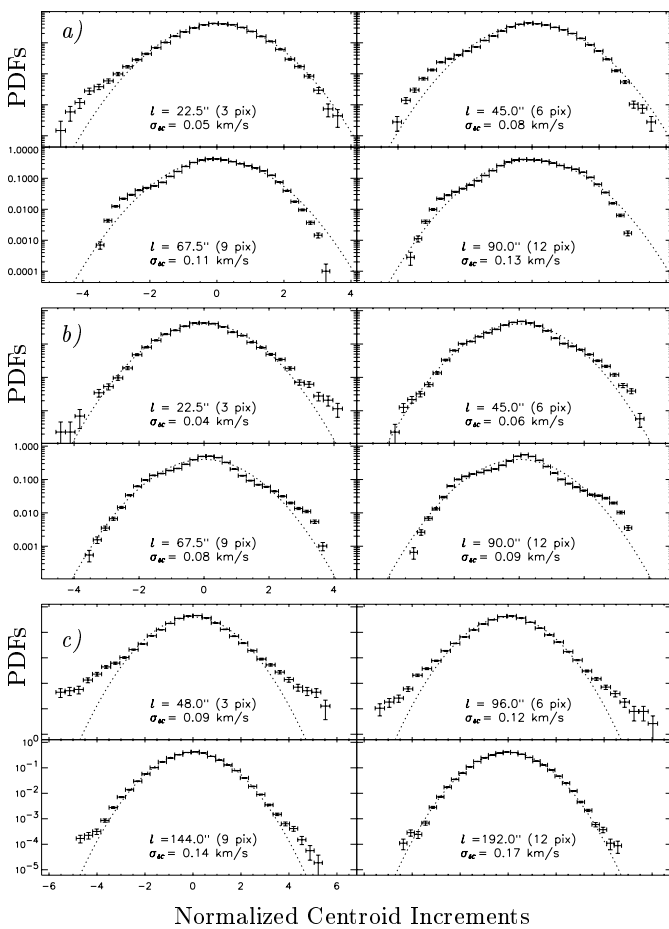


Fig. 7. Same as Fig. 6 for the L1512 field. The PDFs have been computed from the a) ¹²CO($J = 1-0$) and b) ¹³CO($J = 1-0$) IRAM-30m maps and c) ¹²CO ($J = 2-1$) CSO map.

4.3. The influence of large-scale velocity gradients

Large scale variations of the line centroid velocities are visible in Figs. 1 to 3. They most likely trace large scale velocity gradients. The observed values are derived from the two extreme values of the centroid velocity C_{\min} and C_{\max} and the size scale over which they are measured s (Table 4). These numbers provide an estimate of the observed velocity gradients $\nabla v_{\text{obs}} \approx |C_{\max} - C_{\min}|/s$ given in Col. 5 of Table 4. We

note here that the values derived from two different lines in the same field are not necessarily the same because the lines do not quite sample the same gas, as discussed above. Unlike in other studies (Grosdidier et al. 2001; Miesch et al. 1999), we did not remove the contribution of large scale velocity gradients in our centroid increment computations. In this section, we explain why.

In the analysis of turbulence in HII regions, the large scale velocity gradients are related to the expansion of the warm ionized gas, driven by the pressure gradient of the HII region relative to the surrounding medium. It is thus justified to remove large scale gradients, prior to any statistical analysis of the turbulence within HII regions, as done by Grosdidier et al. (2001), because the HII region expansion is an ordered large scale motion and is not part of the inertial range of the turbulent cascade.

Miesch et al. (1999), who analyze the turbulence in molecular clouds, have different goals from ours. Their goal is to extract an homogeneous and isotropic subset of interstellar turbulence and compare it to laboratory experiments or direct numerical simulations. Instead, we search for regions of enhanced shear *at small scales* to localize bursts of dissipation. We show below that we are not limited in our analysis by large scale anisotropy. We first argue, following the results of numerical simulations of Burkert & Bodenheimer (2000) that turbulence itself with its steep power spectrum (i.e. most of the power is in the large scales) generates velocity gradients which may be interpreted as rotation at any scale. We show in Fig. 8 and Table 4 that this is indeed the case. The scaling of the standard deviation of the line centroid increments $\sigma_{\delta C_l}$ with lag l is shown in Fig. 8. Up to lags of ~ 100 arcsec or about 12 pixels, $\sigma_{\delta C_l} \propto l^{\zeta_2}$. The values of ζ_2 are given in Table 4. This scaling provides that of the turbulent velocity shear $\nabla v_{\text{turb}} \propto l^{\zeta_2 - 1}$ with l . The turbulent velocity shear, ∇v_{turb} , inferred from this scaling law at the scale s of the largest centroid differences is given in Col. 9 of Table 4. Comparison of Cols. 5 and 9 shows that in all cases: $\nabla v_{\text{obs}} \lesssim \nabla v_{\text{turb}}$. The large scale observed velocity gradients are therefore smaller than or comparable to those expected from the turbulent shear at the same scale, suggesting that they are part of the turbulent dynamics.

Last, we show that the large-scale gradients do not affect the small scale statistics which are of importance to the present analysis. The contribution of the large-scale gradients to the

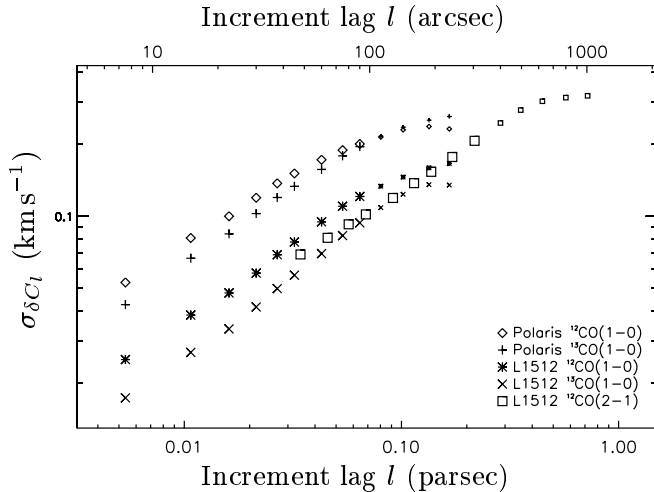


Fig. 8. Scaling of the standard deviations of centroid increments with the lag l used to compute the increments for the five fields and lines. Large symbols are used for the lags which give statistically meaningful results as explained in Appendix B. The exponents ζ_2 listed in Table 4 have been determined for the range of lags which is statistically meaningful.

centroid velocity increments computed as $\delta v_3 = \nabla v_{\text{obs}} \times l$ for $l = 3$ pixels are given in Table 4 for each map (Col. 6). These values are then compared to the standard deviation of the δC -PDFs distribution (Col. 7). In all cases, the ratio $\delta v_3/\sigma_{\delta C_3}$ is smaller than or equal to unity. The contribution of the large-scale gradients to the centroid increments at small scale is therefore much smaller than the increments populating the non-Gaussian wings which extend up to 4 or 5 $\sigma_{\delta C}$ (see Figs. 6 and 7). Furthermore, the large scale gradient has a well-defined direction and therefore preferentially affects the increments computed along the same direction i.e. one increment out of 9 in the case of $l = 3$ pixels (out of more for larger lags). The contribution of the large scale velocity gradient to increments computed over lags of different orientations is further reduced by the appropriate cosine. For these reasons, we are confident that keeping the large scale velocity gradients in the maps does not significantly distort the δC -PDFs of Figs. 6 and 7.

In summary, we did not remove large-scale velocity gradients prior to our statistical analysis for two reasons: (1) there is evidence for the large-scale gradients to be part of the turbulent dynamics we analyze and (2) their value is such that their contamination of the small scale statistics on which we focus here, is negligible.

5. Spatial distribution of the non-Gaussian centroid increments: A new kind of small scale structure

5.1. The positions populating the non-Gaussian wings of the increments PDFs are not randomly distributed: They form elongated structures

We have seen in the previous section that, in most cases, the departure of the increment PDFs from a Gaussian distribution is as large as the lag is small. We are here interested in tracing the

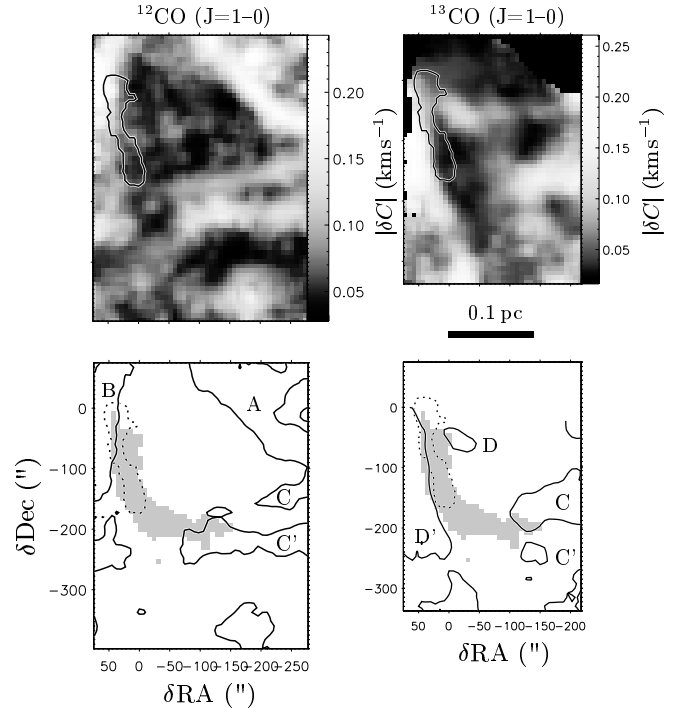


Fig. 9. Top panels: Maps of the averaged centroid increments $|\delta C|$ in Polaris computed with a lag of 3 pixels from the $^{12}\text{CO}(J=1-0)$ lines (left) and the $^{13}\text{CO}(J=1-0)$ lines (right). The solid contour localizes the dense core (see Fig. 1). The bottom panels are sketches meant to emphasize the relative orientation of the region of large column density traced by ^{13}CO (grey area), the dense core (dashed contour) and regions of large centroid increments where $|\delta C| > |\delta C_0|$ (black contours). The thresholds $|\delta C_0|$ are those listed in Table 5. The letters are the labels used in the text.

locations of the positions for which the small scale increment values build up the non-Gaussian wings. We display in Figs. 9 through 11 the spatial distribution of the increments computed for lags of three pixels. The lag value $l = 3$ pixels is the smallest significant lag in the IRAM maps because the maps are Nyquist sampled or better: adjacent pixels are therefore not independent (see Table 1).

The spatial distributions of the centroid increments are shown as azimuthal averages over the 16 possible directions because it allows a more compact presentation of the results. In the rest of the paper, we note $|\delta C|$ the azimuthal average of the modulus of all the oriented centroid increments computed at a given position, for lags of three pixels. The thresholds, $|\delta C_0|$, given in Table 5, have been chosen to isolate the regions where more than half of the oriented increments belong to the non-Gaussian wings of the δC -PDFs. These thresholds are close to the standard deviations $\sigma_{\delta C_3}$ of the increment-PDFs for $l = 3$ pixels. Note that the averaging produces increment values smaller by factors up to a few than the original values used to build the PDFs. This effect is illustrated in the cuts of Figs. 13–16 where the variations of the centroid and of the centroid increments before and after azimuthal averaging are shown as a function of position.

In Figs. 9–11, the uniform black areas (white in Fig. 3) correspond to positions where increments are not computed

Table 5. Characteristics of the maps of averaged centroid increments for each field and line: thresholds used to define the subsets of positions populating the non-Gaussian wings of δC -PDFs (Col. 3), peak increment (Col. 4) and background (Col. 5) values, square of the contrast of the peak to background value (Col. 6), average linewidth in each field (Col. 7), surface filling factor of regions where $|\delta C| > |\delta C_0|$ (Col. 8) and fraction of dissipation in these regions (Col. 9, see Sect. 7).

Field	Line	$ \delta C_0 $ km s ⁻¹	$ \delta C_{\max} $ km s ⁻¹	$ \delta C_{\text{bg}} $ km s ⁻¹	$(\delta C_{\max}/\delta C_{\text{bg}})^2$	$\langle \Delta v \rangle$ km s ⁻¹	f_s	f_ϵ
Polaris	¹² CO ($J = 1-0$)	0.1	0.4	0.06	39	0.8	0.29	0.65
Polaris	¹³ CO ($J = 1-0$)	0.1	0.4	0.05	82	0.4	0.20	0.62
L1512	¹² CO ($J = 1-0$)	0.05	0.1	0.03	15	0.3	0.29	0.61
L1512	¹³ CO ($J = 1-0$)	0.03	0.08	0.02	17	0.2	0.39	0.74
L1512	¹² CO ($J = 2-1$)	0.09	0.4	0.06	45	0.3	0.18	0.53

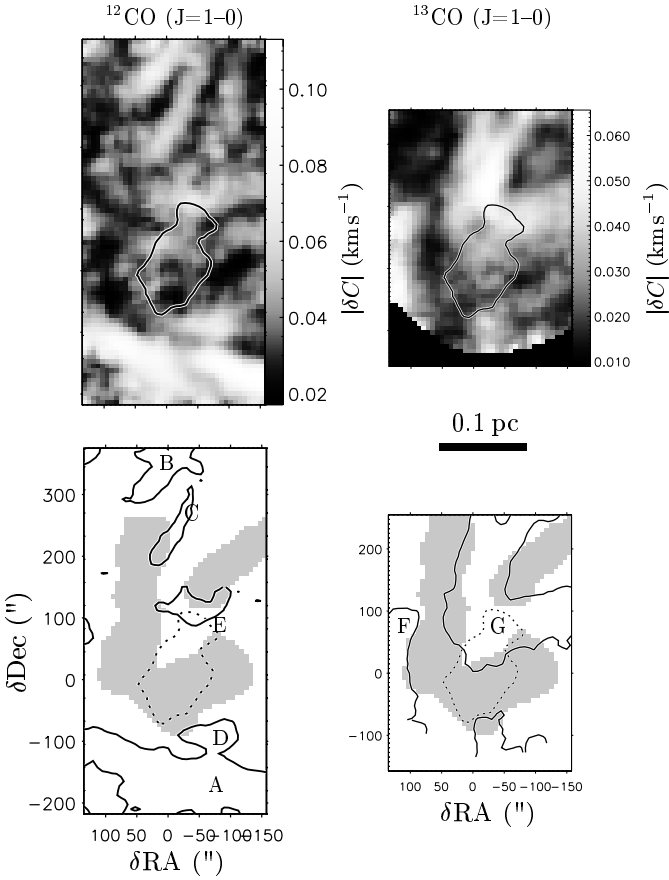


Fig. 10. Same as Fig. 9 for the L1512 field observed at IRAM.

because at least one of two spectra used either has a SNR_a value below 10.0 or is too isolated (i.e. with less than 4 contiguous neighbors with $SNR_a > 10.0$). The black contours drawn after the CS(3-2) or HCO⁺(3-2) maps help locate the dense cores.

In the sketches of Figs. 9–11, the black contours labelled by letters follow the thresholds $|\delta C_0|$. These contours are therefore intended to trace the structures delineated by the positions where the original, non averaged, increments are larger than a few $\sigma_{\delta C_3}$ and belong to the non-Gaussian wings of the δC -PDFs.

Two characteristics of these structures may be derived from Figs. 9–11. Firstly, the spatial distribution of the largest

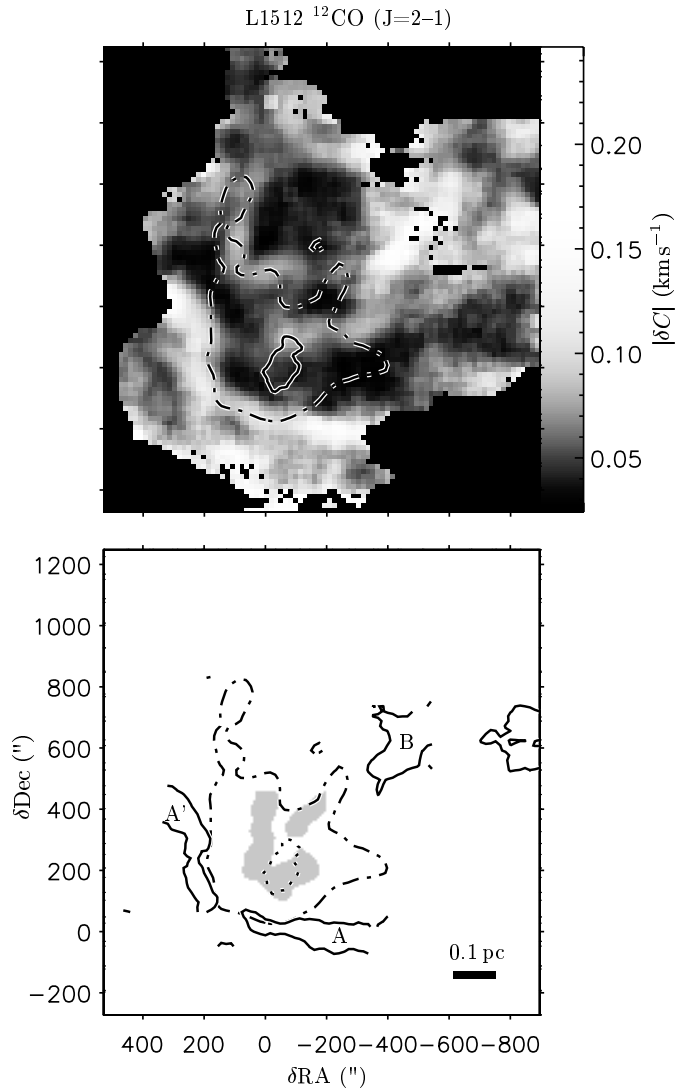


Fig. 11. Same as Fig. 9 for the large scale L1512 field. The dot-dashed contour is the same as in Fig. 3.

averaged increments is not random. They are concentrated in specific regions and appear to form elongated structures, almost straight in several cases, and covering only a small fraction of the maps. Secondly, the $|\delta C|$ values in these structures are well above the background fluctuations of the increments

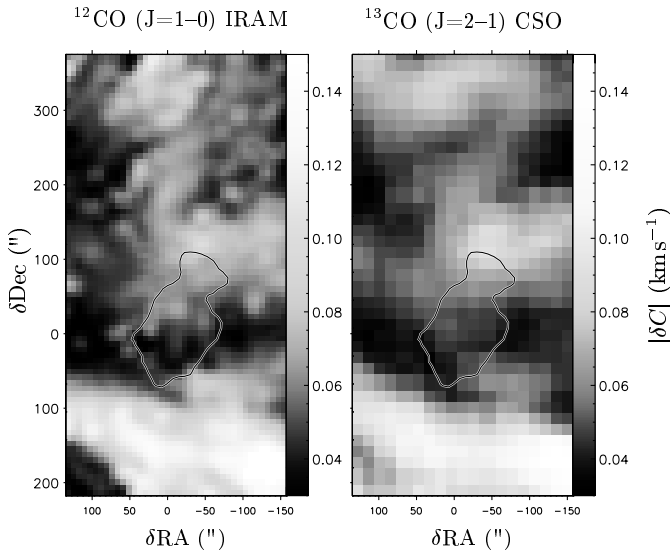


Fig. 12. Maps of average centroid increments of the same region of the L1512 field computed from the IRAM-30m $^{12}\text{CO}(J=1-0)$ map (left) and the CSO $^{13}\text{CO}(J=2-1)$ map (right). Note the remarkable quantitative agreement between the two maps.

in the maps. Table 5 gives the background values $|\delta C_{\text{bg}}|$ computed as the median of all the values smaller than the thresholds $|\delta C_0|$ and the largest values of increments $|\delta C_{\text{max}}|$ in the maps. It shows that the largest values are between 4 and 9 times above the background values, far above statistical fluctuations.

Last, it is remarkable that similar spatial distributions are found for the large increments computed in the same field with two independent data sets i.e. from two different telescopes and two different lines, $^{12}\text{CO}(J=1-0)$ and $(J=2-1)$. Figure 12 allows a comparison of the patterns found for the centroid increments in the L1512 field from the IRAM-30m and CSO maps. The morphological and quantitative agreement between the two maps fully supports the reality of these structures. In Sect. 6, we discuss why they do trace the shear of the velocity field.

The smallest lags of three pixels introduce an effective resolution, ascribing a size of three pixels to unresolved variations of the line centroid. But some of the structures traced by the largest increments are resolved by the observations. This last characteristic is also seen in the cuts of Figs. 13–16. In the RA cut across the L1512 field, the peak of increment visible at RA offset $\sim 200''$ is due to the steepening of the variation of the line centroid over 6 consecutive pixels, from offset $\sim 300''$ to $200''$ or more than three CSO beamwidths. The half-power width of this feature (above the background value), ~ 0.05 pc, is therefore resolved by the procedure of computing increments over lags of 3 pixels. We illustrate with the other cuts that several of these structures are actually resolved, with thicknesses of the order of 0.05 pc, up to 0.08 pc.

5.2. Comparison with the results of Miesch et al. (1999)

The non-random distribution of the largest centroid velocity increments seems at odd with the spotty distribution found

by Miesch et al. (1999) in their maps of centroid velocity differences of several star forming regions. However, we consider that our results agree with theirs for the following reason. Their method, as explained in Sect. 4.3, is different from ours because they first remove a smoothed value from the line centroid values in their maps. The centroid velocity fluctuations they obtain with this procedure are already of the same nature as a velocity increment, as shown by Lis et al. (1996). It is therefore meaningful to compare their maps of centroid fluctuations with our maps of centroid increments. Their maps of centroid velocity fluctuations (their Figs. 4 and 5) do indeed exhibit spatial structures. The largest values appear in several cases to form narrow and elongated structures similar to those we find. The spotty distribution is found only once they compute the increments of these former maps of velocity fluctuations. What they call centroid velocity differences cannot be compared to what we compute.

A more detailed comparison of the two sets of results would not be meaningful for several reasons. The signal-to-noise ratios of our data given in Table 2 are significantly better than those of Miesch et al. (1999). Their SNR values range between 4.5 and 8.9, except for HH83, which has $SNR = 19.7$ but is their smallest field. The nature of the fields is also different. While we analyze only quiescent fields in the present paper, Miesch et al. have analyzed active star forming regions where many shocks, driven by the young stars, interact and probably generate considerable small scale structure in the patterns of the vorticity. Last, most of the fields being at distances larger than the nearby fields studied here, the angular characteristic scales of the structures (if any) are expected to be smaller.

5.3. The regions of largest centroid increments are not coinciding with density nor column density peaks

The sketches of Figs. 9–11 allow a comparison of the structures of largest increments with the concentrations of matter, whether they are large column density or large density structures or both. The dotted contours delineate the dense cores and the grey areas those where the $^{13}\text{CO}(J=1-0)$ integrated emission is larger than 3 and 4 K km s^{-1} , for L1512 and Polaris respectively, which corresponds to $N(\text{H}_2) > \text{a few } 10^{21} \text{ cm}^{-2}$, about 10 times the average column density of the large scale environment. The main result visible in those figures is that none of the regions of largest increments coincides, even in projection, with either a dense core (i.e. a density peak) or a column density peak.

This result is better seen in the cuts of Figs. 13–16: the large variations at small scale of the line centroid, at the origin of the large increment values, do not coincide with column density extrema. Such a coincidence would be expected if the line centroid variations were due to the overlap on the line of sight of unrelated gas components. In that case, an extremum of column density would coincide with the region of large centroid increment, which is that over which the two components are intercepted, producing a modified (locally broader) line profile. This lack of correlation is also seen in the scatter plots of Fig. 17 where the largest increments are in no case associated

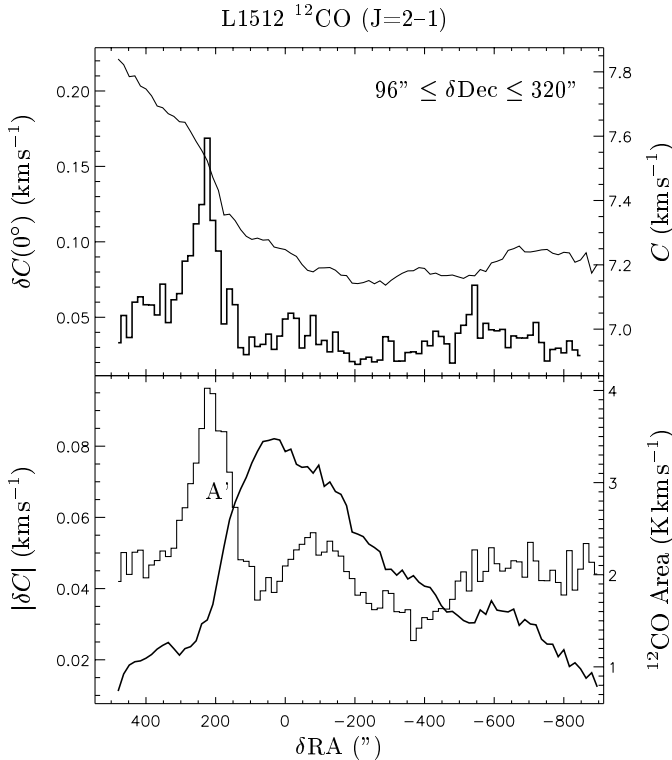


Fig. 16. Cuts in the ^{12}CO line of the large scale L1512 field in RA, averaged over the indicated Dec range. The histograms are the cuts of increment centroids: raw values in the top panel, averaged in azimuth as explained in Sect. 5.1 in the bottom panel. The other cuts are the centroid velocities (*top*) and integrated area (*bottom*). The letters correspond to those of Fig. 9.

with extrema of column density. Conversely, the largest increments tend to occur at the positions where the integrated areas are the smallest.

5.4. Link between the orientation of these structures and that of density structures

The previous section shows that the structures traced by the regions of largest centroid increments at small scale do not coincide (in projection) with the dense cores nor the column density extrema in the maps.

Nonetheless, this statement does not mean that these small scale structures are unrelated to the dense cores. The orientation of the most elongated structures tend to be related with that of the dense cores. In the Polaris field, the structure labelled C' in the ^{12}CO map, appears as a prolongation of the bright ^{13}CO structure, which seems to be an extension of the dense core itself (see sketches of Fig. 9). The structures B and D' are adjacent and parallel to the edge of the dense core, as is seen also in the cuts of Fig. 14. In the L1512 field (Fig. 11), similar patterns are visible. The two southern structures A and A' frame in projection the region of bright ^{12}CO lines (dot-dashed contour), and the structure B is elongated in the same direction as that of the filament of matter seen in ^{13}CO and ^{12}CO and the orientation of the dense core itself as mapped in HCO^+ . Since the velocity is continuous between the

structures of large increments and those of matter (see maps of line centroids in Figs. 1–3), it is most likely that they are connected in the 3-dimensional space.

5.5. Volume filling factor of the largest increment structures

For each field, in addition to the maximum value of the increments in the map, $|\delta C_{\text{max}}|$ and the background value $|\delta C_{\text{bg}}|$ defined in Sect. 5.1, Table 5 also gathers the fraction f_s of pixels in the map where the centroid increments are larger than the thresholds $|\delta C_0|$. f_s may be understood as their surface filling factor in the maps. From its values, one derives volume filling factors of the regions of largest increments of at most a few %. Assuming that the depth of the structures along the line of sight is comparable to their projected thickness, d , and that the extent of the cloud along the line of sight, L , is provided by the large scale maps (Falgarone et al. 1998), $f_v \sim f_s d/L$. For $d \sim 0.05$ pc and $L \sim 1$ pc, $f_v \sim$ a few % or less.

6. Discussion

The subset of spectra shown in Fig. 5 encompasses one of the regions where the centroid increments are the largest. Figure 5 shows that the optimal window found to compute the line centroids varies only by very small amounts from one spectrum to the next. This illustrates that the measured centroid increments are due to small variations in the shapes of the line profiles from one line of sight to the next. These variations may have several origins, unrelated to the structure of the velocity field: optical depth effects, gas temperature or density fluctuations. We show in the following section that our interpretation of these variations as tracers of small scale velocity shears is justified.

6.1. Optical depth effects

The optical depth of the lines used introduces a possible bias in the sampling of the velocity field performed by the spectra. Indeed, optically thick lines tend to sample preferentially foreground material.

However, the maps of centroid increments built from the ^{13}CO and ^{12}CO lines, although not exactly the same, are reminiscent of each other, suggesting that the two lines sample about the same gas (see Figs. 9–10). This is also suggested by the similarity of the maps obtained in L1512 with two different ^{12}CO transitions of different optical depths (Fig. 12). These results support the concept of low effective optical depth of the ^{12}CO lines discussed in Martin et al. (1984). Falgarone et al. (1998) also inferred a low effective optical depth in the ^{12}CO lines of the fields analyzed here (except in the projected area of the dense cores), on the basis of the uniformity of the excitation temperature. Note that a low effective optical depth of the ^{12}CO lines is to be expected in these fields (as opposed to brighter molecular clouds) because a large velocity dispersion of the gas, leading to a large escape probability of the photons in velocity space, is combined with low gas column densities. The same similarity is found between maps of centroid increments built from the ^{13}CO and C^{18}O spectra of the environment

of the L1689B dense core (Pety et al. in preparation), suggesting that optical depth effects are not seriously affecting the centroid increment determinations.

This result illustrates a major strength of the line centroid analysis. On the one hand, centroids are velocity moments and as such provide more weight to velocities far from the line centroids, i.e. the line wings. The wings are the parts of the line profiles the least affected by optical depth effects. On the other hand, our method which limits the range of velocities used to compute the centroid (see Sect. 4.1), limits the noise contribution, particularly in the line wings. We thus argue that there is a velocity range in spectrally well-sampled line profiles, neither too far in the line wings (to avoid noise artefacts) nor too close to the line centroid (to avoid optical depth artefacts), which carries enough information on the velocity field to be used on statistical grounds.

Last, self-absorption is also unlikely in those low brightness core environments because the excitation temperature of the ^{12}CO and ^{13}CO lines is remarkably uniform as was shown by the scatter plots of the ($J = 1-0$) to ($J = 2-1$) ^{12}CO and ^{13}CO lines (Falgarone et al. 1998). For self-absorption to occur the optical depth must exceed unity and the foreground material must have a lower excitation temperature.

6.2. Role of temperature fluctuations

It might be argued that the variations of line centroids traced by the ^{12}CO lines are due only to variations of the gas kinetic temperature. This might be the case in bright sources where the ^{12}CO lines are thermalized. But it has to be appreciated that in the fields under study here the ^{12}CO lines are weak (less than 5 K) while the gas is poorly shielded from the ambient UV field (H_2 column densities less than a few 10^{21} cm^{-2}) and thus likely to be much warmer than 5 K. It means that either the lines are subthermally excited, in which case their excitation is mostly radiative, or that they are thermally excited in small structures with a large beam dilution factor. In either case, the line profile variations are not determined by variations of the gas kinetic temperature but by those of the radiative excitation and/or beam filling factor, which both depend on the velocity field (Falgarone et al. 1998).

6.3. Role of density fluctuations unrelated to the velocity field

In this section, we discuss the influence of column density or density variations, independent of the velocity field, on the velocity increments that we measure. Under the optically thin hypothesis, $dN(u) \propto T(u) du$ where $dN(u)$ is the column density of gas at the projected velocity u within du . The centroid velocity defined in Eq. (1) is therefore also

$$C = \frac{\int u dN(u)}{\bar{N}} \quad \text{with} \quad \bar{N} = \int dN(u). \quad (3)$$

The centroid increment between lines of sight 1 and 2 can thus be written

$$\delta C_{1-2} = \left[\frac{\int u_1 dN_1(u_1)}{\bar{N}_1} \right] - \left[\frac{\int u_2 dN_2(u_2)}{\bar{N}_2} \right]. \quad (4)$$

From this expression, it is obvious that $\delta C_{1-2} = 0$ if the two lines of sight differ only by their total column density, the velocity structure being the same i.e. $N_2(u_2) \propto N_1(u_1)$ for all velocities. In this case the two integrated line profiles are simply homothetic.

Let us consider another simple case in which the total column density N along the two lines of sight is approximately the same, the only difference being that a fraction of the gas at projected velocity u_1 on the first line of sight is at u_2 on the other. This is possible, for instance, in the perspective of matter distributed in a large number of very small structures, possibly virialized in the potential well of the cloud. Then $\delta C_{1-2} = (u_1 - u_2) \delta N/N$ where δN is the column density of the gas which is at different projected velocities on each line of sight. We estimate below the value of $\delta N/N$ required to produce centroid increments of the order of $|\delta C_{\text{max}}|$ in each field. We assume that the velocity difference $u_1 - u_2$ is of the order of magnitude of the average linewidth in each field $\langle \Delta v \rangle$ given in Table 5, $u_1 - u_2 \sim \langle \Delta v \rangle$. Therefore an estimate of the required column density fluctuations is: $\delta N/N \sim |\delta C_{\text{max}}| / \langle \Delta v \rangle$. The latter ratio can be computed from the values in Table 5 and is equal to 1.25, 0.5 and 1.05 for the three fields of Fig. 18, respectively. Now, if we assume that the relative variations of the line integrated areas trace those of gas column density, we can check how the observed fluctuations of column density compare with the above values. This may be seen in the scatter plots of the averaged increments versus the relative variations of the line integrated area, $|\delta A|/A$ (Fig. 18). These scatter plots are shown only for the three maps where the statistics are the least affected by the dense core. The values of $|\delta C_{\text{max}}| / \langle \Delta v \rangle$ and of the threshold $|\delta C_0|$ are indicated for each field. The scatter plots show that, except for a few lines of sight of the Polaris ^{13}CO map, the observed relative variations of column density (or $|\delta A|/A$) at the positions of increments populating the non-Gaussian wings (i.e. increments above the thresholds $|\delta C_0|$) are smaller than the minimum value estimated above. In the large scale L1512 field and the Polaris field as seen in ^{12}CO , it is therefore unlikely that the observed centroid increments trace small scale column density fluctuations due for instance to randomly distributed clumps. This is also supported by the fact that the structures traced by the large centroid increments form spatially coherent patterns and are not randomly distributed as would be expected for randomly moving clumps. We thus conclude that the contribution of small scale random column density variations to the observed centroid increments is not dominant.

These small scale column density fluctuations could trace density fluctuations. They occur over lags of 3 pixels or $l \sim 5000 \text{ AU}$ (see Table 1). The corresponding density enhancements are therefore expected to occur over the same small scale l . Thus, $\delta n/n \sim (|\delta A|/A)(L/l)$ where L is the depth of the cloud. For $|\delta A|/A = 0.5$ and $L = 1 \text{ pc}$, $\delta n/n = 25$. These might be supersonic shocks with Mach numbers of about 5,

in which case the density and velocity fields are closely correlated. However, such large density enhancements would produce variations of the $\text{CO}(J = 2-1)/\text{CO}(J = 1-0)$ line ratio in the weak regions of the fields, variations which are not observed (Falgarone et al. 1998).

The tenuous variations of the line profiles at the origin of the large centroid increments at small scale are therefore most likely ascribed to large velocity shears at small scale.

7. Possible link with the dissipation of turbulence

Our statistical analysis of the environment of two low mass dense cores leads to two main results: (i) the PDFs of line centroid increments (or velocity shears) have non-Gaussian wings as more pronounced as the lag over which they are computed is small, and (ii) the spatial distribution of the largest shears is not random and reveals often resolved and elongated structures not coinciding with density nor column density peaks. These are predominantly velocity structures.

The first property is a characteristic shared by incompressible turbulence, compressible turbulence and compressible, magnetized turbulence and is a signature of the intermittency of the velocity or current density fields. Our results therefore suggest that the observed distributions of large shears may trace regions of intense vorticity reminiscent of those responsible for the intermittency of turbulence. Note that these regions of intense vorticity are embedded in the bulk of the flow and that the large centroid increments are only tracers of their presence in the flow. In other words, a small scale structure of large vorticity, embedded in the bulk of the flow, can be at the origin of a small variation of the line centroids, along the boundary of two large scale regions of more uniform centroid velocity, i.e. the large eddies.

We have argued that the large shears are unlikely to trace shocks but we cannot rule out the fact that the vorticity extrema are fossil structures of shocks, already disrupted.

If the largest shears that we have detected trace local enhancements of the vorticity and therefore dissipation, as discussed in the Introduction, an estimate of the local enhancement of the dissipation rate in those regions is provided by the ratio $(\delta C_{\text{max}}/\delta C_{\text{bg}})^2$ (Table 5) because the dissipation rate follows the square of the vorticity (or of the velocity shear) (Landau & Lifshitz 1987). The local enhancements are large, between 15 and 80 (Table 5). It is interesting to note that the smallest values of $(\delta C_{\text{max}}/\delta C_{\text{bg}})^2$ are obtained in the close vicinity of the L1512 dense core (IRAM field) and the largest in Polaris and in the large scale environment of L1512 (CSO field). This may reflect the fact that turbulent dissipation is a less violent process within a dense core which is already formed. We also note that even the weakest structures in the maps of $|\delta C|$ have a contrast larger than 2 above the background and therefore correspond to a dissipation rate locally 4 times larger than in the bulk of the field (see for instance the weak structures, labelled B and C, in the North of the IRAM L1512 ^{12}CO map).

Last, the quantity $f_\epsilon = \sum_{|\delta C| > |\delta C_0|} |\delta C|^2 / \sum_{\text{all}} |\delta C|^2$ (Table 5) may be understood as the fraction of the dissipation taking place in the regions of large increments (those

populating the non-Gaussian wings). The underlying assumption is that the ratio of the actual dissipation rate to $|\delta C|^2$ is the same in the regions where the increments populate the Gaussian cores of the PDFs (the bulk of the fields) as in those populating the non-Gaussian wings. Under this assumption, the values of Table 5 show that a significant fraction (more than half) of the total dissipation takes place in the ensemble of small scale structures filling a volume close to a few % of the cloud volume, as estimated above.

8. Conclusion

We have found a new kind of small scale structure in the environment of low mass starless dense cores.

They emerge in the maps of line centroid increments between adjacent spectra as regions between 4 and 9 times brighter than the background fluctuations. They are essentially velocity structures because they are not coinciding with any detected sufficient increase in the column density or density. They are small scale, often elongated, structures. A few are resolved by the observations with thicknesses up to 0.08 pc. Most of them are not resolved i.e. smaller than 0.02 pc, the effective resolution due to the smallest lag available to compute the increments.

The measured projected velocity being orthogonal to the plane of the sky, in which lags are measured, the centroid velocity increments trace one component of the shear of the velocity field. The shear-PDFs exhibit non-Gaussian wings which are the most prominent for the smallest lag. For this reason, we ascribe these structures to turbulence by analogy with the behavior of such PDFs in laboratory flows or numerical simulations of compressible turbulence. The new structures delineate the locus of positions where the increments are the largest and populate the non-Gaussian wings.

These structures therefore possibly trace regions of enhanced dissipation of turbulence in the observed fields. The nature of the dissipation process is not determined by the observations: it could be either viscous dissipation, ion-neutral collisions or enhanced current densities. These structures are unlikely to trace shocks, but they may trace the remanent vorticity generated by shocks already disrupted. In this picture, a significant fraction of the turbulent energy present in the field would be dissipating in those structures filling less than a few % of the cloud volume.

Last, the structures display some connexion with the tracers of dense gas, such as shared orientations and velocity and space pattern continuity. Before we can make a link between such structures and the formation of dense cores, we need to investigate control fields i.e. molecular clouds without dense cores (Hily-Blant et al. in preparation) and with moderate star formation activity (Pety et al., in preparation).

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