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# SEMI-LINEAR SUB-ELLIPTIC EQUATIONS ON THE HEISENBERG GROUP WITH A SINGULAR POTENTIAL 

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#### Abstract

In this work, we study the Dirichlet problem for a class of semi-linear subelliptic equations on the Heisenberg group with a singular potential. The singularity is controlled by Hardy's inequality, and the nonlinearity is controlled by Sobolev's inequality. We prove the existence of a nontrivial solution for a homogenous Dirichlet problem.


Key words: Heisenberg group, Hardy's inequality, Sobolev's inequality, Singular potential

## A.M.S. Classification

## 1. Introduction

In this work, we study the partial differential equations on the Heisenberg group $\mathbb{H}^{d}$. Let us recall that the Heisenberg group is the space $\mathbb{R}^{2 d+1}$ of the (non commutative) law of product

$$
(x, y, s) \cdot\left(x^{\prime}, y^{\prime}, s^{\prime}\right)=\left(x+x^{\prime}, y+y^{\prime}, s+s^{\prime}+2\left(\left(y \mid x^{\prime}\right)-\left(y^{\prime} \mid x\right)\right)\right) .
$$

The left invariant vector fields are

$$
X_{j}=\partial_{x_{j}}+2 y_{j} \partial_{s}, \quad Y_{j}=\partial_{y_{j}}-2 x_{j} \partial_{s}, \quad j=1, \cdots, d \quad \text { and } \quad S=\partial_{s}=\frac{1}{4}\left[Y_{j}, X_{j}\right] .
$$

In all that follows, we shall denote by $Z_{j}=X_{j}$ and $Z_{j+d}=Y_{j}$ for $j \in\{1, \cdots, d\}$. We fix here some notations :

$$
z=(x, y) \in \mathbb{R}^{2 d}, w=(z, s) \in \mathbb{H}^{d}, \rho(z, s)=\left(|z|^{4}+|s|^{2}\right)^{1 / 4}
$$

where $\rho$ is the Heisenberg distance. Moreover, the Laplacian-Kohn operator on $\mathbb{H}^{d}$ and Heisenberg gradient is given by

$$
\Delta_{\mathbb{H}^{d}}=\sum_{j=1}^{n} X_{j}^{2}+Y_{j}^{2} ; \quad \nabla_{\mathbb{H}^{d}}=\left(Z_{1}, \cdots, Z_{2 d}\right) .
$$

Let $\Omega$ be an open and bounded domain of $\mathbb{H}^{d}$, we define thus the associated Sobolev space as following

$$
H^{1}\left(\Omega, \mathbb{H}^{d}\right)=\left\{f \in L^{2}(\Omega) ; \nabla_{\mathbb{H}^{d}} f \in L^{2}(\Omega)\right\}
$$

and $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ is the closure of $C_{0}^{\infty}(\Omega)$ in $H^{1}\left(\Omega, \mathbb{H}^{d}\right)$.

We consider the following semi-linear Dirichlet problem

$$
\left\{\begin{array}{l}
-\Delta_{\mathbb{H} d} u-\mu V u=\lambda u+|u|^{p-2} u \quad \text { in } \quad \Omega  \tag{1.1}\\
\left.u\right|_{\partial \Omega}=0
\end{array}\right.
$$

where $V$ is a positive potential function which admits the singularity on $\Omega, \lambda$ is a real constant and $2<p<2+\frac{2}{d}$; the index $2^{*}=2+\frac{2}{d}$ is the critical index of Sobolev's inequality on the Heisenberg group $[9,15,16,17,22]$

$$
\begin{equation*}
\|u\|_{L^{2^{*}}(\Omega)} \leq C_{\Omega}\|u\|_{H^{1}\left(\Omega, \mathbb{H}^{d}\right)} \tag{1.2}
\end{equation*}
$$

for all $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
The potential function $V$ is controlled by the following Hardy's inequality,

$$
\begin{equation*}
\int_{\Omega} V(w)|u(w)|^{2} d w \leq\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}^{2} \tag{1.3}
\end{equation*}
$$

for all $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
We will prove in the next section the following results

- If $0 \in \Omega$, Hardy's inequality (1.3) holds for

$$
V(z, s)=\frac{d^{4}}{(d+1)^{2}} \rho(z, s)^{-2}
$$

- If $0 \in \Omega$, Hardy's inequality (1.3) also holds for a softer potential

$$
V(z, s)=d^{2} \frac{|z|^{2}}{\rho(z, s)^{4}}
$$

see also $[12,18]$.

- In the Lemma 2.7, we prove that Hardy's inequality (1.3) holds for

$$
V(z, s)=\bar{\mu} \rho_{c}(z, s)^{-2}
$$

where $\rho_{c}$ defined in (2.8) is the distance to a sub-manifold $\Sigma_{c}$ of codimension $\geq 2$ and $\bar{\mu}$ is a constant.

Theorem 1.1. Assume that the potential function $V$ satisfies (1.3), then for any $\lambda>0$ and any $0 \leq \mu<1$, the Dirichlet problem (1.1) admits a nontrivial solution in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.

The Dirichlet problem (1.1) on the Heisenberg group is a natural generalization of the classical problem on $\mathbb{R}^{d}$, see $[5,7,8,10,11,14,20]$ and their references. The subellipticity of the operator $\Delta_{\mathbb{H}^{d}}$ implies that any weak solution of the Dirichlet problem (1.1) belongs to $C^{\infty}\left(\Omega \backslash \Sigma_{c}\right)$ (see [25]). The issue of regularity of a weak solution near $\partial \Omega \cup \Sigma_{c}$ is a very delicate problem.

This paper is organized as follows. In section 2, we recall Poincaré's inequality and prove Hardy's inequality on the Heisenberg group; Section 3 deals with the study of the eigenvalue problem; finally, we prove the existence of a weak solution in Sections 4 and 5 by using Rabinowitz's Theorem and the Palais-Smale Theorem.

## 2. HARDY'S INEQUALITIES ON $\mathbb{H}^{d}$

The following density theorem is very useful. Let us give here an other simple proof than [2] (see also [3]),
Theorem 2.1. We have that $\mathcal{C}_{0}^{\infty}(\Omega \backslash\{(0,0)\})$ is dense in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ for $d \geq 1$.

Proof : By definition of $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, it suffices to show that

$$
\mathcal{C}_{0}^{\infty}(\Omega) \subset \overline{\mathcal{C}_{0}^{\infty}\left(\Omega \backslash\{(0,0)\}, \mathbb{H}^{d}\right)}{ }^{\|} \cdot \|_{H^{1}}
$$

Let $\varphi$ be a cut-off function for which

$$
\varphi(\eta)=\left\{\begin{array}{lll}
0 & \text { if } & 0<\eta \leq 1  \tag{2.1}\\
1 & \text { if } & \eta \geq 2
\end{array}\right.
$$

For $u \in \mathcal{C}_{0}^{\infty}\left(\Omega, \mathbb{H}^{d}\right)$, let $\varepsilon>0$ small enough, and we set $u_{\varepsilon}(z, s)=\varphi\left(\frac{1}{\varepsilon} \rho(z, s)\right) u(z, s)$. So $u_{\varepsilon} \in \mathcal{C}_{0}^{\infty}\left(\Omega \backslash\{(0,0)\}, \mathbb{H}^{d}\right)$ and we have

$$
\left\|u_{\varepsilon}-u\right\|_{H^{1}(\Omega)}^{2}=\left\|\nabla_{\mathbb{H}^{d}}\left(u_{\varepsilon}-u\right)\right\|_{L^{2}(\Omega)}^{2}+\left\|u_{\varepsilon}-u\right\|_{L^{2}(\Omega)}^{2}
$$

Dominated convergence theorem implies that

$$
\left\|u_{\varepsilon}-u\right\|_{L^{2}(\Omega)}^{2} \rightarrow 0
$$

and

$$
\int_{\Omega}\left|\varphi\left(\frac{1}{\varepsilon} \rho(z, s)\right)-1\right|^{2}\left|\nabla_{\mathbb{H}^{d}} u(z, s)\right|^{2} d z d s \rightarrow 0, \text { when } \varepsilon \rightarrow 0
$$

On the other hand, we have

$$
\begin{aligned}
& \int_{\Omega}\left|\nabla_{\mathbb{H}^{d}}\left(\frac{1}{\varepsilon} \rho(z, s)\right)\right|^{2}\left|\varphi^{\prime}\left(\frac{1}{\varepsilon} \rho(z, s)\right)\right|^{2}|u(z, s)|^{2} d z d s \\
& =\frac{1}{\varepsilon^{2}} \int_{\Omega} \frac{|z|^{2}}{\rho(z, s)^{2}}\left|\varphi^{\prime}\left(\frac{1}{\varepsilon} \rho(z, s)\right)\right|^{2}|u(z, s)|^{2} d z d s \\
& \leq \frac{1}{\varepsilon^{2}}\|u\|_{L^{\infty}(\Omega)}^{2}\left\|\varphi^{\prime}\right\|_{L^{\infty}(\Omega)}^{2} \int_{\{(z, s) ; \varepsilon \leq \rho(z, s) \leq 2 \varepsilon\}} d z d s \\
& \leq C \frac{1}{\varepsilon^{2}} \varepsilon^{2 d+2} \rightarrow 0, \text { as } \varepsilon \rightarrow 0 .
\end{aligned}
$$

Remark that this proof also show that the density theorem is not true for classical Sobolev space in a 2 dimensional case.

Now, we state the following precise Poincaré inequality,
Theorem 2.2. Let $\Omega$ be a sub-domain of $\mathbb{H}^{d}$ bounded in some direction of $\left(z_{1}, \cdots, z_{2 d}\right)$, that is, there exist $R>0$ and $1 \leq j_{0} \leq 2 d$ such that $0<r=\left|z_{j_{0}}\right| \leq R$ for all $(z, s) \in \Omega$. Then for any $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$,

$$
\begin{equation*}
\int_{\Omega}|u|^{2} d z d s \leq 4 R^{2} \int_{\Omega}\left|\nabla_{\mathbb{H}^{d}} u\right|^{2} d z d s \tag{2.2}
\end{equation*}
$$

Remark : By using the inequality (2.2), we can use $\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}$ as a norm on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$. The Poincaré inequality (2.2) holds for any $\Omega \subset\left\{(z, s) \in \mathbb{H}^{d} ; \rho(z, s) \leq R\right\}$. We can also obtain the Poincaré inequality from Bony's maximum principle for general Hörmander's vector fields but with a non-precise constant (see [6, 21]). The proof given here is a modification of L. D'Ambrosio [13].
Proof : Using the density results of Theorem 2.1, take $u \in \mathcal{C}_{0}^{\infty}(\Omega \backslash\{(0,0)\})$ and let $T(z, s)=\left(T_{1}(z, s), \cdots, T_{2 d}(z, s)\right)$ be a $C^{1}$ vector function on $\Omega$. Denote by

$$
d i v_{\mathbb{H}} T=\sum_{j=1}^{2 d} Z_{j} T_{j},
$$

we have

$$
\begin{aligned}
\int_{\Omega}\left(d i v_{\mathbb{H}^{d}} T\right)|u|^{2} d z d s & =-2 \int_{\Omega}\left\langle T, \nabla_{\mathbb{H}^{d}} u\right\rangle u d z d s \\
& \leq 2 \int_{\Omega}\left|\left\langle T, \nabla_{\mathbb{H}^{d}} u\right\rangle u\right| d z d s \\
& \leq 2\left[\int_{\Omega}\left|\nabla_{\mathbb{H}^{d}} u\right|^{2} d z d s\right]^{\frac{1}{2}}\left[\int_{\Omega}|T|^{2}|u|^{2} d z d s\right]^{\frac{1}{2}} \\
& \leq \int_{\Omega}\left|\nabla_{\mathbb{H}^{d}} u\right|^{2} d z d s+\int_{\Omega}|T|^{2}|u|^{2} d z d s .
\end{aligned}
$$

Thus

$$
\int_{\Omega}\left(d i v_{\mathbb{H}^{d}} T-|T|^{2}\right)|u|^{2} d z d s \leq \int_{\Omega}\left|\nabla_{\mathbb{H}^{d}} u\right|^{2} d z d s .
$$

For $\varepsilon>0$, let us choose $T:=T_{\varepsilon}=-\frac{1}{2} \frac{\nabla_{\mathbb{H} d} r_{\varepsilon}}{r_{\varepsilon}}$ where $r_{\varepsilon}=\left(r^{2}+\varepsilon^{2}\right)^{\frac{1}{2}}$ and $r=\left|z_{j_{0}}\right|$. Then

$$
\begin{gathered}
d i v_{\mathbb{H}^{d}} T_{\varepsilon}=-\frac{1}{2} \frac{1}{r_{\varepsilon}^{2}}\left[r_{\varepsilon} \Delta_{\mathbb{H}^{d}} r_{\varepsilon}-\left|\nabla_{\mathbb{H}^{d}} r_{\varepsilon}\right|^{2}\right] \\
d i v_{\mathbb{H}^{d}} T_{\varepsilon}-\left|T_{\varepsilon}\right|^{2}=-\frac{1}{2} \frac{\Delta_{\mathbb{H}^{d}} r_{\varepsilon}}{r_{\varepsilon}}+\frac{1}{4} \frac{\left|\nabla_{\mathbb{H}^{d}} r_{\varepsilon}\right|^{2}}{r_{\varepsilon}^{2}}
\end{gathered}
$$

since

$$
\begin{gathered}
\nabla_{\mathbb{H} d} r_{\varepsilon}=\frac{\nabla_{\mathbb{H}^{d}} r^{2}}{2 r_{\varepsilon}}=\frac{r}{r_{\varepsilon}}, \quad\left|\nabla_{\mathbb{H}^{d}} r_{\varepsilon}\right|^{2}=\frac{r^{2}}{r_{\varepsilon}^{2}} \\
\Delta_{\mathbb{H}^{d} d} r_{\varepsilon}=\frac{\Delta_{\mathbb{H}^{d}} r^{2}}{2 r_{\varepsilon}}-\frac{\nabla_{\mathbb{H}^{d} d} r^{2}}{2 r_{\varepsilon}^{2}} \nabla_{\mathbb{H}^{d} d} r_{\varepsilon}=\frac{1}{r_{\varepsilon}}-\frac{r^{2}}{r_{\varepsilon}^{3}} .
\end{gathered}
$$

So,

$$
d i v_{\mathbb{H} d} T_{\varepsilon}-\left|T_{\varepsilon}\right|^{2}=-\frac{1}{2 r_{\varepsilon}^{2}}+\frac{3 r^{2}}{4 r_{\varepsilon}^{4}},
$$

and for any $(z, s) \in \Omega$,

$$
\lim _{\varepsilon \rightarrow 0} d i v_{\mathbb{H}^{d}} T_{\varepsilon}-\left|T_{\varepsilon}\right|^{2}=\frac{1}{4} \frac{1}{r^{2}} \geq \frac{1}{4} \frac{1}{R^{2}} .
$$

From the dominated convergence theorem, we have

$$
\lim _{\varepsilon \rightarrow 0} \int_{\Omega}\left[d i v_{\mathbb{H}^{d}} T_{\varepsilon}-\left|T_{\varepsilon}\right|^{2}\right]|u|^{2} \quad d z d s \geq \frac{1}{4} \frac{1}{R^{2}} \int_{\Omega}|u|^{2} d z d s
$$

thus

$$
\frac{1}{4} \frac{1}{R^{2}} \int_{\Omega}|u|^{2} \quad d z d s \leq \int_{\Omega}\left|\nabla_{\mathbb{H}^{d}} u\right|^{2} d z d s .
$$

Let us give a very easy proof of the classical Hardy inequality by using a radial vector field. Let $U$ be a bounded domain of $\mathbb{R}^{d}, d>2, H_{0}^{1}(U)$ is the usual Sobolev space.
Lemma 2.3. We have, for any $u \in H_{0}^{1}(U)$,

$$
\begin{equation*}
\left(\frac{d}{2}-1\right)^{2} \int_{U} \frac{u^{2}}{|x|^{2}} d x \leq\left\|\nabla_{x} u\right\|_{L^{2}(U)}^{2} \tag{2.3}
\end{equation*}
$$

Proof : As $C_{0}^{\infty}(U \backslash\{0\})$ is dense $H_{0}^{1}(U)$, we have restricted ourselves to a function $u$ in $C_{0}^{\infty}(U \backslash\{0\})$. The proof mainly consists of an integration by parts with respect to the radial vector field $R$,

$$
R=\sum_{j=1}^{d} x_{j} \partial_{x_{j}}
$$

We notice that $R\left(|x|^{-2}\right)=-2|x|^{-2}$ and $\operatorname{div} R=d$, so

$$
\begin{aligned}
\int_{U} \frac{u^{2}}{|x|^{2}} d x & =-\frac{1}{2} \int_{U} R\left(|x|^{-2}\right) d x \\
& =\frac{1}{2} \int_{U} d i v(R) \frac{u^{2}}{|x|^{2}} d x+\frac{1}{2} \int_{\Omega} \frac{1}{|x|^{2}} R\left(u^{2}\right) d x \\
\left(1-\frac{d}{2}\right) \int_{U} \frac{u^{2}}{|x|^{2}} d x & =\frac{1}{2} \int_{U} \frac{1}{|x|^{2}} R\left(u^{2}\right) d x=\int_{U} \sum_{j=1}^{d} \frac{u}{|x|} \frac{x_{j}}{|x|} \partial_{x_{j}} u d x .
\end{aligned}
$$

By Cauchy-Schwarz,

$$
\left(\frac{d}{2}-1\right) \int_{U} \frac{|u|^{2}}{|x|^{2}} d x \leq\left(\int_{U} \sum_{j=1}^{d}\left|\partial_{x_{j}} u\right|^{2} d x\right)^{\frac{1}{2}}\left(\int_{U} \frac{|u|^{2}}{|x|^{2}} d x\right)^{\frac{1}{2}} .
$$

On the Heisenberg group, if we introduce the radial vector field

$$
R=\sum_{j=1}^{d}\left(x_{j} \partial_{x_{j}}+y_{j} \partial_{y_{j}}\right)=\sum_{j=1}^{d}\left(x_{j} X_{j}+y_{j} Y_{j}\right),
$$

then we immediately obtain for $d>1$,

$$
(d-1)^{2} \int_{\Omega} \frac{u^{2}}{\rho(z, s)^{2}} d z d s \leq(d-1)^{2} \int_{\Omega} \frac{u^{2}}{|z|^{2}} d z d s \leq\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}^{2}
$$

for any $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
By using the idea inspired from the radial vector field, we now prove the following Hardy inequality.

Lemma 2.4. For $d \geq 1$, we have that

$$
\begin{equation*}
\left(\frac{d^{2}}{d+1}\right)^{2} \int_{\Omega} \frac{u^{2}}{\rho^{2}} d z d s \leq\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}^{2} \tag{2.4}
\end{equation*}
$$

for any $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
Proof : By using the density theorem, we prove the inequality (2.4) for the function $u \in C_{0}^{\infty}(\Omega \backslash\{(0,0)\})$. Then the proof mainly consists of an integration by parts with respect to the radial vector field $R_{\mathbb{H}^{d}}$ adapted to the structure of $\mathbb{H}^{d}$, namely

$$
R_{\mathbb{H}^{d}}=2 s \partial_{s}+\sum_{j=1}^{d}\left(x_{j} \partial_{x_{j}}+y_{j} \partial_{y_{j}}\right)=\frac{s}{2 d} \sum_{j=1}^{d}\left[Y_{j}, X_{j}\right]+\sum_{j=1}^{d}\left(x_{j} X_{j}+y_{j} Y_{j}\right) .
$$

We notice that $R_{\mathbb{H}^{d}}\left(\rho^{-2}\right)=-2 \rho^{-2}$ and $\operatorname{div} R_{\mathbb{H}^{d}}=2 d+2$. We have

$$
\begin{aligned}
\int_{\Omega} \frac{u^{2}}{\rho(z, s)^{2}} d z d s & =-\frac{1}{2} \int_{\Omega} R_{\mathbb{H}^{d}}\left(\rho(z, s)^{-2}\right) u^{2} d z d s \\
& =\frac{1}{2} \int_{\Omega} \rho^{-2} R_{\mathbb{H}^{d}}\left(u^{2}\right) d z d s+\frac{1}{2} \int_{\Omega} \rho^{-2} u^{2} \operatorname{div} R_{\mathbb{H}^{d}} d z d s .
\end{aligned}
$$

This gives

$$
\begin{aligned}
-d \int_{\Omega} \frac{u^{2}}{\rho^{2}} d z d s & =\int_{\Omega} \sum_{j=1}^{d} \frac{u}{\rho}\left(\frac{x_{j}}{\rho} X_{j}+\frac{y_{j}}{\rho} Y_{j}\right) u d z d s-\frac{1}{2 d} \int_{\Omega} Y_{j}\left(\frac{s}{\rho^{2}}\right) u\left(X_{j} u\right) d z d s \\
& +\frac{1}{2 d} \int_{\Omega} X_{j}\left(\frac{s}{\rho^{2}}\right) u\left(Y_{j} u\right) d z d s \\
& =\left(1+\frac{1}{d}\right) \int_{\Omega} \sum_{j=1}^{d}\left(\frac{x_{j} u}{\rho^{2}} X_{j} u+\frac{y_{j} u}{\rho^{2}} Y_{j} u\right) d z d s \\
& +\frac{1}{d} \int_{\Omega} \sum_{j=1}^{d} \frac{s}{\rho^{6}}\left[|z|^{2} y_{j}-s x_{j}\right] u X_{j} u d z d s \\
& -\frac{1}{d} \int_{\Omega} \sum_{j=1}^{d} \frac{s}{\rho^{6}}\left[|z|^{2} x_{j}+s y_{j}\right] u Y_{j} u d z d s,
\end{aligned}
$$

then

$$
\begin{aligned}
-d^{2} \int_{\Omega} \frac{u^{2}}{\rho^{2}} d z d s & =\int_{\Omega} \sum_{j=1}^{d}\left[(d+1)-\frac{s^{2}}{\rho^{4}}\right]\left(\frac{x_{j} u}{\rho^{2}} X_{j} u+\frac{y_{j} u}{\rho^{2}} Y_{j} u\right) d z d s \\
& +\int_{\Omega} \sum_{j=1}^{d} \frac{s|z|^{2}}{\rho^{4}}\left[\frac{y_{j} u}{\rho^{2}} X_{j} u-\frac{x_{j} u}{\rho^{2}} Y_{j} u\right] d z d s \\
& =\int_{\Omega} \sum_{j=1}^{d}\left[\left((d+1)-\frac{s^{2}}{\rho^{4}}\right) \frac{x_{j}}{\rho}+\frac{s|z|^{2}}{\rho^{4}} \frac{y_{j}}{\rho}\right] \frac{u}{\rho} X_{j} u d z d s \\
& +\int_{\Omega} \sum_{j=1}^{d}\left[\left((d+1)-\frac{s^{2}}{\rho^{4}}\right) \frac{y_{j}}{\rho}-\frac{s|z|^{2}}{\rho^{4}} \frac{x_{j}}{\rho}\right] \frac{u}{\rho} Y_{j} u d z d s .
\end{aligned}
$$

Setting

$$
A(z, s)=\sum_{j=1}^{d}\left\{\left[\left((d+1)-\frac{s^{2}}{\rho^{4}}\right) \frac{x_{j}}{\rho}+\frac{s|z|^{2}}{\rho^{4}} \frac{y_{j}}{\rho}\right]^{2}+\left[\left((d+1)-\frac{s^{2}}{\rho^{4}}\right) \frac{y_{j}}{\rho}-\frac{s|z|^{2}}{\rho^{4}} \frac{x_{j}}{\rho}\right]^{2}\right\},
$$

Cauchy-Schwarz inequality implies

$$
d^{2} \int_{\Omega} \frac{|u|^{2}}{\rho^{2}} d z d s \leq\left(\int_{\Omega} A(z, s) \frac{|u|^{2}}{\rho^{2}} d z d s\right)^{\frac{1}{2}}\left(\int_{\Omega} \sum_{j=1}^{d}\left(\left|X_{j} u\right|^{2}+\left|Y_{j} u\right|^{2}\right) d z d s\right)^{\frac{1}{2}}
$$

For $A(z, s)$, we have

$$
\begin{aligned}
A(z, s) & =\left((d+1)-\frac{s^{2}}{\rho^{4}}\right)^{2} \frac{|z|^{2}}{\rho^{2}}+\frac{s^{2}|z|^{4}}{\rho^{8}} \frac{|z|^{2}}{\rho^{2}} \\
& =\frac{|z|^{2}}{\rho^{2}}\left[(d+1)^{2}-(2 d+1) \frac{s^{2}}{\rho^{4}}\right] \\
& =\frac{|z|^{2}}{\rho^{6}}\left[(d+1)^{2}|z|^{4}+d^{2} s^{2}\right] \\
& =\frac{|z|^{2}}{\rho^{6}}\left[(2 d+1)|z|^{4}+d^{2}\left(|z|^{4}+s^{2}\right)\right] \\
& \leq(d+1)^{2} \frac{|z|^{2}}{\rho^{2}} \leq(d+1)^{2}
\end{aligned}
$$

So, we deduce the inequality (2.4).
The Hardy inequality on the Heisenberg group $\mathbb{H}^{d}$ is first proven in $[18,12]$ for a softer potential.
Lemma 2.5. We have, for any $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$,

$$
\begin{equation*}
d^{2} \int_{\Omega} \frac{|z|^{2}}{|z|^{4}+s^{2}}|u|^{2} \quad d z d s \leq\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}^{2} \tag{2.5}
\end{equation*}
$$

The singularity of potential in the Hardy inequalities (2.3), (2.4) and (2.5) is a isolate point of domain. We consider now the general case when the singularity is on a submanifold. We have first the following density result:

Lemma 2.6. Let $\Omega$ be a bounded domain of $\mathbb{R}^{2 d+1}$ and $\Sigma_{c}$ a sub-manifold of $\Omega$ such that $\operatorname{dim} \Sigma_{c} \leq 2 d-1$. Then $C_{0}^{\infty}\left(\Omega \backslash \Sigma_{c}\right)$ is dense in the space $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.

Proof : As $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ is a Hilbert space, it is enough to prove that the orthogonal of $C_{0}^{\infty}\left(\Omega \backslash \Sigma_{c}\right)$ in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ is $\{0\}$. Let $u$ be in this space. For any $v$ in $C_{0}^{\infty}\left(\Omega \backslash \Sigma_{c}\right)$, we have

$$
(u, v)_{L^{2}}+\left(\nabla_{\mathbb{H}^{d}} u, \nabla_{\mathbb{H}^{d}} v\right)_{L^{2}}=0
$$

By integration by part,

$$
\forall v \in C_{0}^{\infty}\left(\Omega \backslash \Sigma_{c}\right), \quad\left\langle u-\Delta_{\mathbb{H}^{d}} u, v\right\rangle=0
$$

this implies that, as a distribution,

$$
\operatorname{Supp}\left(u-\Delta_{\mathbb{H}^{d}} u\right) \subset \Sigma_{c}
$$

Since $u-\Delta_{\mathbb{H}^{d}} u$ belong to the classical Sobolev space $H^{-1}(\Omega)$ and except 0 , no distribution of $H^{-1}(\Omega)$ can be supported in a submanifold of dimension $\leq(2 d+1)-2$. Thus $u-\Delta_{\mathbb{H}^{d}} u=$ 0 on $\Omega$. Taking the $L^{2}$ scalar product with $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ implies that $u \equiv 0$. This completes the proof of Lemma 2.6.

We consider now the hyper-surface $\Sigma=\{(x, y, s) \in \Omega: g(x, y, s)=s+f(x, y)=0\}$ where $\Omega$ is a neighborhood of 0 in $\mathbb{H}^{d}$. Assume that

$$
\begin{equation*}
\Sigma_{c}=\left\{w \in \Omega: g(w)=0, \nabla_{\mathbb{H}^{d}} g(w)=0\right\} \tag{2.6}
\end{equation*}
$$

is a sub-manifold of dimension $(2 d+1)-r-1, r \geq 1$.

Lemma 2.7. Assume that $\Sigma_{c}$ is a sub-manifold of dimension $2 d-r$ and $r \geq 1$. Then, there exists $\bar{\mu}>0$ such that for any $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$,

$$
\begin{equation*}
\bar{\mu} \int_{\Omega} \frac{u^{2}}{\rho_{c}^{2}} d w \leq\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}^{2} \tag{2.7}
\end{equation*}
$$

with

$$
\begin{equation*}
\rho_{c}(w)=\left(g^{2}(w)+\left|\nabla_{\mathbb{H}^{d}} g(w)\right|^{4}\right)^{1 / 4} \tag{2.8}
\end{equation*}
$$

We refer to the proof of this lemma to [4] and also [2]. The constant $\bar{\mu}$ depends, of course on $\Sigma_{c}$, but in many interesting cases, it depends only on the dimension of $\Sigma_{c}$.

Here we present a proof for a model case in $\mathbb{H}^{1}$ to precise the constant $\bar{\mu}$. We take $g(x, y, s)=s+2 x y$, then

$$
\begin{align*}
\Sigma & =\left\{(x, y, s) \in \mathbb{H}^{1}: s+2 x y=0\right\} \\
\Sigma_{c} & =\{(x, 0,0), x \in \mathbb{R}\} \tag{2.9}
\end{align*}
$$

Lemma 2.8. Let $\Sigma_{c}$ as in (2.9), then, we have for any $u \in H^{1}\left(\mathbb{H}^{1}\right)$,

$$
\begin{equation*}
\frac{2^{2}}{5+2^{8}} \int_{\mathbb{H}^{1}} \frac{u^{2}}{\rho_{c}^{2}} d w \leq\left\|\nabla_{\mathbb{H}^{1}} u\right\|_{L^{2}\left(\mathbb{H}^{1}\right)}^{2} \tag{2.10}
\end{equation*}
$$

Proof : We rectify $\Sigma$ by setting $x^{\prime}=x, y^{\prime}=y, s^{\prime}=s+2 x y$, so the vector fields X and Y change to $X^{\prime}=\partial_{x^{\prime}}+4 y^{\prime} \partial_{s^{\prime}}, Y^{\prime}=\partial_{y^{\prime}}$ and

$$
\rho_{c}\left(x^{\prime}, y^{\prime}, s^{\prime}\right)=\left(\left(4 y^{\prime}\right)^{4}+s^{2}\right)^{1 / 4}
$$

Let $R$ be a radial vector field

$$
\begin{aligned}
R & =X^{\prime}\left(s^{\prime}\right) Y^{\prime}+2^{3} s^{\prime} \partial_{s^{\prime}} \\
& =4 y^{\prime} Y^{\prime}+2^{3} s^{\prime} \partial_{s^{\prime}} \\
& =4 y^{\prime} Y^{\prime}+2 s^{\prime}\left[Y^{\prime}, X^{\prime}\right]
\end{aligned}
$$

where $R\left(\rho_{c}^{-2}\right)=-8 \rho_{c}^{-2}$ and $\operatorname{div} R=12$. Using the density Lemma 2.6 , we have for $u \in C_{0}^{\infty}\left(\mathbb{H}^{1} \backslash\{(0,0)\}\right)$,

$$
\begin{aligned}
\int_{\mathbb{H}^{1}} \frac{u^{2}}{\rho_{c}^{2}} d z^{\prime} d s^{\prime} & =-\frac{1}{8} \int_{\mathbb{H}^{1}} u^{2} R\left(\rho_{c}^{-2}\right) d z^{\prime} d s^{\prime} \\
-\frac{1}{2} \int_{\mathbb{H}^{1}} \frac{u^{2}}{\rho_{c}^{2}} d z^{\prime} d s^{\prime} & =\int_{\mathbb{H}^{1}} \frac{y^{\prime}}{\rho_{c}^{2}} u Y^{\prime} u d z^{\prime} d s^{\prime}-\frac{1}{2} \int_{\mathbb{H}^{1}} Y^{\prime}\left(\frac{s^{\prime}}{\rho_{c}^{2}}\right) u X^{\prime} u d z^{\prime} d s^{\prime} \\
& +\frac{1}{2} \int_{\mathbb{H}^{1}} X^{\prime}\left(\frac{s^{\prime}}{\rho_{c}^{2}}\right) u Y^{\prime} u d z^{\prime} d s^{\prime} \\
& =\int_{\mathbb{H}^{1}} \frac{y^{\prime}}{\rho_{c}^{2}} u Y^{\prime} u d z^{\prime} d s^{\prime}+\int_{\mathbb{H}^{1}} 2^{8} \frac{y^{\prime 3} s^{\prime}}{\rho_{c}^{6}} u X^{\prime} u d z^{\prime} d s^{\prime} \\
& +\int_{\mathbb{H}^{1}}\left[\frac{2 y^{\prime}}{\rho_{c}^{2}}-\frac{2 y^{\prime} s^{\prime 2}}{\rho_{c}^{6}}\right] u Y^{\prime} u d z^{\prime} d s^{\prime} \\
& =\int_{\mathbb{H}^{1}}\left[\frac{3 y^{\prime}}{\rho_{c}}-\frac{2 y^{\prime} s^{\prime 2}}{\rho_{c}^{5}}\right] \frac{u}{\rho_{c}} Y^{\prime} u d z^{\prime} d s^{\prime}+\int_{\mathbb{H}^{1}} 2^{8} \frac{y^{\prime 3} s^{\prime}}{\rho_{c}^{5}} \frac{u}{\rho_{c}} X^{\prime} u d z^{\prime} d s^{\prime}
\end{aligned}
$$

We then obtain

$$
\frac{1}{2} \int_{\mathbb{H}^{1}} \frac{|u|^{2}}{\rho_{c}^{2}} d z^{\prime} d s^{\prime} \leq\left(\int_{\mathbb{H}^{1}} A\left(z^{\prime}, s^{\prime}\right) \frac{|u|^{2}}{\rho^{2}} d z^{\prime} d s^{\prime}\right)^{\frac{1}{2}}\left\|\nabla_{\mathbb{H}^{1}} u\right\|_{L^{2}\left(\mathbb{H}^{1}\right)},
$$

with

$$
\begin{aligned}
A\left(z^{\prime}, s^{\prime}\right) & =\left[\frac{3 y^{\prime}}{\rho_{c}}-2 \frac{y^{\prime} s^{\prime 2}}{\rho_{c}^{5}}\right]^{2}+2^{16} \frac{y^{\prime 6} s^{\prime 2}}{\rho_{c}^{10}} \\
& =3^{2} \frac{y^{\prime 2}}{\rho_{c}^{2}}-12 \frac{y^{\prime 2}\left(\rho_{c}^{4}-2^{8} y^{\prime 4}\right)}{\rho_{c}^{6}}+4 \frac{y^{\prime 2}\left(\rho_{c}^{4}-2^{8} y^{\prime 4}\right)^{2}}{\rho_{c}^{10}}+2^{16} \frac{y^{\prime 6}\left(\rho_{c}^{4}-2^{8} y^{\prime 4}\right)}{\rho_{c}^{10}} \\
& =\frac{y^{\prime 2}}{\rho_{c}^{2}}\left[1+2^{8}\left(4+2^{8}\right) \frac{y^{4}}{\rho_{c}^{4}}+2^{8}\left(2^{10}-2^{16}\right) \frac{y^{\prime 8}}{\rho_{c}^{8}}\right] \\
& \leq \frac{1}{2^{4}}\left[1+\frac{1}{2^{8}} 2^{8}\left(4+2^{8}\right)\right] \\
& \leq \frac{1}{2^{4}}\left(5+2^{8}\right) .
\end{aligned}
$$

## 3. Variational formulation and eigenvalue problem

Thanks to Hardy's inequality (1.3) and Poincaré's inequality (2.2),

$$
\begin{equation*}
\|u\|_{\mu}=\left(\int_{\Omega}\left[\left|\nabla_{\mathbb{H}^{d}} u(z, s)\right|^{2}-\mu V(z, s)|u(z, s)|^{2}\right] d z d s\right)^{\frac{1}{2}} \tag{3.1}
\end{equation*}
$$

is equivalent to the norm on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ for all $0 \leq \mu<1$, so that we will use $\|\cdot\|_{\mu}$ as the norm of $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.

We will use the variational method to study the Dirichlet problem (1.1). We define the following energy functional on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ :

$$
\begin{equation*}
I_{\mu, \lambda}(u)=\frac{1}{2} \int_{\Omega}\left[\left|\nabla_{\mathbb{H}^{d}} u\right|^{2}-\mu V|u|^{2}\right] d z d s-\frac{1}{p} \int_{\Omega}|u|^{p} d z d s-\frac{\lambda}{2} \int_{\Omega}|u|^{2} d z d s \tag{3.2}
\end{equation*}
$$

Similar to the classical case, $I_{\mu, \lambda}(\cdot)$ is well-defined on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ and belongs to $C^{1}\left(H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right) ; \mathbb{R}\right)$.
We say that $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ is a weak solution of the Dirichlet problem (1.1), if for any $v \in \mathcal{C}_{0}^{\infty}(\Omega)$, there holds

$$
\int_{\Omega}\left[\nabla_{\mathbb{H}^{d}} u \overline{\nabla_{\mathbb{H}^{d}} v}-\mu V u \bar{v}\right] d z d s-\int_{\Omega}|u|^{p-2} u \bar{v} d z d s-\lambda \int_{\Omega} u \bar{v} d z d s=0
$$

So a weak solution $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ of the Dirichlet problem (1.1) is a critical point of $I_{\mu, \lambda}$. The Euler-Lagrange equation of the variational problem (3.2) is exactly the semilinear equation in (1.1), and we have

$$
\left\langle I_{\mu, \lambda}^{\prime}(u), v\right\rangle=\int_{\Omega}\left[\nabla_{\mathbb{H}^{d}} u \overline{\nabla_{\mathbb{H}^{d}} v}-\mu V u \bar{v}-|u|^{p-2} u \bar{v}-\lambda u \bar{v}\right] d z d s=0
$$

for any $v \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
Since we consider the Dirichlet problem (1.1) for any $\lambda>0$, we cannot use the direct method to prove the existence of the critical point for $I_{\mu, \lambda}$. We need to use the Mountain Pass Theorem and the Linking Theorem of Rabinowitz (see [23, 24, 26]).

Thusthat we study firstly the spectral decomposition of $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ with respect to the operator $-\Delta_{\mathbb{H}^{d}}-\mu V$ where the singular potential $V$ satisfies Hardy's inequality (1.3). This eigenvalue problem has also its independent interest. We have the following proposition.

Proposition 3.1. Let $0 \leq \mu<1$. Then there exist $0<\lambda_{1}<\lambda_{2} \leq \lambda_{3} \leq \ldots \leq \lambda_{k} \leq \ldots \rightarrow$ $+\infty$, such that for each $k \geq 1$, the following Dirichlet problem

$$
\left\{\begin{array}{l}
-\Delta_{\mathbb{H}^{d}} \phi_{k}-\mu V \phi_{k}=\lambda_{k} \phi_{k}, \quad \text { in } \Omega  \tag{3.3}\\
\phi_{k} \mid \partial \Omega=0
\end{array}\right.
$$

admits a nontrivial solution in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$. Moreover, $\left\{\phi_{k}\right\}_{k \geq 1}$ constitutes an orthonormal basis of Hilbert space $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.

Remark that the first eigenvalue $\lambda_{1}$ is characterized by the following Poincaré inequality

$$
\begin{equation*}
\|u\|_{L^{2}(\Omega)}^{2} \leq \frac{1}{\lambda_{1}} \int_{\Omega}\left(\left|\nabla_{\mathbb{H}^{d} d} u\right|^{2}-\mu V|u|^{2}\right) d z d s \tag{3.4}
\end{equation*}
$$

for all $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
The first step of the proof is the following compact embedding result
Lemma 3.2. Let $\Omega \in \mathbb{H}^{d}$ be a bounded open domain. Then $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ is compactly embedded to $L^{2}(\Omega)$.

We can prove this result by the continuous embedding of $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ into usual the Sobolev space $H_{0}^{1 / 2}(\Omega)$, then the compact embedding of $H^{1 / 2}(\Omega)$ into $L^{2}(\Omega)$. But the first embedding requires some careful extension results. We refer to [19] for a complete and elegant proof of this compact embedding result.

## Proof of Proposition 3.1

Denote by $L_{\mu}=-\Delta_{\mathbb{H}^{d}}-\mu V$ the operator defined on the Hilbert space $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ with the norm $\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}$, then Hardy's inequality (1.3) implies

$$
\left(L_{\mu} u, u\right)_{L^{2}(\Omega)} \geq(1-\mu)\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}^{2}>0, \quad \forall u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)
$$

and

$$
\left(L_{\mu} u, v\right)=\left(u, L_{\mu} v\right), \forall u, v \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)
$$

Hence it is positive, definite and self-adjoint on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$. The Lax-Milgram Theorem implies that for any $g \in H^{-1}\left(\Omega ; \mathbb{H}^{d}\right)$, the following Dirichlet problem

$$
\left\{\begin{array}{clc}
L_{\mu} u & =g \quad \text { in } \Omega \\
u & =0 & \text { on } \partial \Omega
\end{array}\right.
$$

admits a unique solution $u$ belonging to $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, where $H^{-1}\left(\Omega ; \mathbb{H}^{d}\right)$ is the dual space of $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right), g \in H^{-1}\left(\Omega ; \mathbb{H}^{d}\right)$ if $g \in \mathscr{D}^{\prime}(\Omega)$ and there exists $C>0$ such that

$$
|\langle g, \varphi\rangle| \leq C\|\varphi\|_{H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)}
$$

for all $\varphi \in C_{0}^{\infty}(\Omega)$ with the norm

$$
\|g\|_{H^{-1}\left(\Omega, \mathbb{H}^{d}\right)}=\sup _{\varphi \in C_{0}^{\infty}(\Omega)} \frac{|\langle g, \varphi\rangle|}{\|\varphi\|_{H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)}}
$$

Then

$$
\nabla_{\mathbb{H}^{d}}: L^{2}(\Omega) \rightarrow H^{-1}\left(\Omega ; \mathbb{H}^{d}\right) \quad \text { and } \quad \Delta_{\mathbb{H}^{d}}: H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right) \quad \rightarrow \quad H^{-1}\left(\Omega ; \mathbb{H}^{d}\right)
$$

are continuous. The inverse operator $L_{\mu}^{-1}$ of $L_{\mu}$ is well defined and it is a continuous map from $H^{-1}\left(\Omega, \mathbb{H}^{d}\right)$ into $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.

The compact embedding $i: H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right) \rightarrow L^{2}(\Omega)$ and the continuous embedding $i^{*}$ : $L^{2}(\Omega) \rightarrow H^{-1}\left(\Omega, \mathbb{H}^{d}\right)$ imply that $K_{\mu}=L_{\mu}^{-1} \circ i^{*} \circ i: H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right) \rightarrow H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ is a compact
and self adjoint operator. So the spectrum of the compact operator $K_{\mu}$ is $\left\{\eta_{k}\right\}$ such that $\eta_{k}>0, k \geq 1$ and $\eta_{k} \rightarrow 0$. If $\left\{\phi_{k}\right\}$ are the associated normal eigenvectors, we have that

$$
K_{\mu} \phi_{k}=\eta_{k} \phi_{k}, \quad \forall k \geq 1,
$$

and $\left\{\phi_{k}\right\}$ form a complete basis of Hilbert space $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, which completes the proof of Proposition 3.1.

## 4. Existence of critical points

We prove now the following existence result of critical points for the variational functional $I_{\mu, \lambda}$ which gives the weak solution for the Dirichlet problem (1.1).
Theorem 4.1. Let $0 \leq \mu<1, \lambda>0$, then $I_{\mu, \lambda}$ admits at last one nontrivial critical point on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.

We recall now the well-known Palais-Smale condition.
Definition 4.2. Let $E$ be a Banach space, $I \in C^{1}(E, \mathbb{R})$ and $c \in \mathbb{R}$. We say that $I$ satisfies the $(P S)_{c}$ condition, if for any sequence $\left\{u_{n}\right\} \subset E$ with the properties :

$$
I\left(u_{n}\right) \rightarrow c \quad \text { and }\left\|I^{\prime}\left(u_{n}\right)\right\|_{E^{\prime}(\Omega)} \rightarrow 0
$$

there exists a subsequence which is convergent, where $I^{\prime}(\cdot)$ is the Frechet differentiation of $I$ and $E^{\prime}$ is the dual space of $E$. If this holds for any $c \in \mathbb{R}$, we say that I satisfies the $(P S)$ condition.

We will prove in the next section the following result
Theorem 4.3. Let $0 \leq \mu<1, \lambda>0$, then $I_{\mu, \lambda}$ satisfies the $(P S)$ condition on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
Let $0<\lambda_{1}<\lambda_{2} \leq \lambda_{3} \leq \ldots \leq \lambda_{k} \leq \ldots \rightarrow+\infty$ be the eigenvalues of $-\Delta_{\mathbb{H}^{d}}-\mu V$ in Proposition 3.1. We consider firstly the case $0<\lambda<\lambda_{1}$ and we use the following Mountain Pass Theorem to prove the existence of a critical point for $I_{\mu, \lambda}$ :
Theorem 4.4. (see [1, 23])
Let $E$ be a Banach space and $I \in C^{1}(E, \mathbb{R})$. We suppose that $I(0)=0$ and satisfies that
(i) there exist $R>0, a>0$ such that if $\|u\|_{E}=R$, then $I(u) \geq a$;
(ii) there exists $e \in E$ such that $\|e\|>R$ and $I(e)<a$. If I satisfies the $(P S)_{c}$ condition with

$$
c=\inf _{h \in \Gamma} \max _{t \in[0,1]} I(h(t)), \text { where } \Gamma=\{h \in C([0,1] ; E) ; h(0)=0 \text { and } h(1)=e\},
$$

then $c$ is a critical value of $I$ and $c \geq a$.
We check the above conditions for $I=I_{\mu, \lambda}$ on $E=H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$. We have $I_{\mu, \lambda}(0)=0$. For $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, Sobolev's inequality (1.2) and Hardy's inequality imply that

$$
\|u\|_{L^{p}(\Omega)} \leq C_{\Omega}\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)} \leq \frac{C_{\Omega}}{(1-\mu)^{1 / 2}}\|u\|_{\mu}
$$

Then, for $0<\lambda<\lambda_{1}$

$$
\begin{equation*}
I_{\mu, \lambda}(u) \geq \frac{1}{2}\left(1-\frac{\lambda}{\lambda_{1}}\right)\|u\|_{\mu}^{2}-\frac{C_{1}}{p}\|u\|_{\mu}^{p} \geq C_{1}\|u\|_{\mu}^{2}\left(\frac{1}{2 C_{1}}\left(1-\frac{\lambda}{\lambda_{1}}\right)-\frac{1}{p}\|u\|_{\mu}^{p-2}\right) \tag{4.1}
\end{equation*}
$$

where

$$
C_{1}=\left(\frac{C_{\Omega}}{(1-\mu)^{1 / 2}}\right)^{p}>1
$$

Let

$$
R_{0}=\left(\frac{p}{2 C_{1}}\left(1-\frac{\lambda}{\lambda_{1}}\right)\right)^{\frac{1}{p-2}}>0
$$

Then for any $0<R<R_{0}$,

$$
\begin{equation*}
\inf _{\|u\|_{\mu}=R} I_{\mu, \lambda}(u)=a(R)>0 \tag{4.2}
\end{equation*}
$$

So $I_{\mu, \lambda}$ satisfies condition (i) of Theorem 4.4.
For condition (ii) of Theorem 4.4, take $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ such that $\|u\|_{\mu}=R>0$, then for $\theta \geq 0$,

$$
\begin{align*}
I_{\mu, \lambda}(\theta u) & =\frac{\theta^{2}}{2} \int_{\Omega}\left[\left|\nabla_{\mathbb{H}^{d}} u\right|^{2}-\mu V(z, s)|u|^{2}\right] d z d s  \tag{4.3}\\
& -\frac{\theta^{p}}{p} \int_{\Omega}|u|^{p} d z d s-\frac{\lambda \theta^{2}}{2} \int_{\Omega}|u|^{2} d z d s \tag{4.4}
\end{align*}
$$

Since $p>2$, thus

$$
\lim _{\theta \rightarrow+\infty} I_{\mu, \lambda}(\theta u)=-\infty
$$

Then, there exists $\theta_{1}>0$ large enough such that for $e=\theta_{1} u$, we have $\|e\|_{\mu}>R$ and $I_{\mu}\left(\theta_{1} u\right)<0<a(R)$. Set now

$$
\Gamma=\left\{h \in C\left([0,1] ; H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)\right) ; h(0)=0 \text { and } h(1)=e\right\}
$$

then by continuity, we have

$$
c=\inf _{h \in \Gamma} \max _{t \in[0,1]} I_{\mu, \lambda}(h(t)) \geq a(R)>0
$$

and $c$ is a local minimum. Theorem 4.3 implies that the $(P S)_{c}$ condition is satisfied. So $c>0$ is a critical value by using Theorem 4.4 and the critical point is $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, which is nontrivial. We have proved Theorem 4.4 for $0<\lambda<\lambda_{1}$.

We need now the following Linking theorem from Rabinowitz [23].
Theorem 4.5. Let $E$ be a Banach space with $E=Y \oplus X$, where $\operatorname{dim} Y<\infty$. Suppose that $I \in C^{1}(E, \mathbb{R})$ and satisfies
(i)there exist $\rho, \alpha>0$ such that $\left.I\right|_{\partial B_{\rho} \cap X} \geq \alpha$;
(ii) there exist $e \in \partial B_{1} \cap X$ and $R>\rho$ such that if $A \equiv\left(\bar{B}_{R} \cap Y\right) \oplus\{r e, 0<r<R\}$, then $\left.I\right|_{\partial A} \leq 0$.

If I satisfies the $(P S)_{c}$ condition with

$$
c=\inf _{h \in \Gamma} \max _{u \in A} I(h(u)), \text { where } \Gamma=\left\{h \in C(\bar{A}, E) ;\left.h\right|_{\partial A}=i d\right\}
$$

then $c$ is a critical value of $I$ and $c \geq \alpha$.
Remark 4.6. Suppose $\left.I\right|_{Y} \leq 0$ and there are an $e \in \partial B_{1} \cap X$ and $\bar{R}>\rho$ such that $I(u) \leq 0$ for $u \in Y \oplus \operatorname{span}\{e\}$ and $\|u\| \geq \bar{R}$, then for any large $R$, we have $\left.I\right|_{\partial A} \leq 0$ where $A=\left(\bar{B}_{R} \cap Y\right) \oplus\{r e, 0<r<R\}$.

We assume now that there is $k$ such that $\lambda_{k} \leq \lambda<\lambda_{k+1}$, where $\lambda_{k}$ is the $k$-th eigenvalue of the operator $-\Delta_{\mathbb{H}^{d}}-\mu V$. Let $Y=\operatorname{span}\left\{\phi_{1}, \ldots, \phi_{k}\right\}$, where $\phi_{k}$ is the eigenfunction corresponding to $\lambda_{k}$. Then Proposition 3.1 implies that $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)=Y \bigoplus X$ where $X=$ span $\left\{\phi_{l} ; l>k\right\}$. Thus we have

$$
\begin{equation*}
\int_{\Omega}\left|\nabla_{\mathbb{H}^{d}} y\right|^{2}-\mu|y|^{2} \quad d z d s \leq \lambda_{k} \int_{\Omega}|y|^{2} d z d s, \quad \forall y \in Y \tag{4.5}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{\Omega}\left|\nabla_{\mathbb{H} d} u\right|^{2}-\mu V|u|^{2} d z d s \geq \lambda_{k+1} \int_{\Omega}|u|^{2} d z d s, \quad \forall u \in X . \tag{4.6}
\end{equation*}
$$

We will show that $I_{\mu, \lambda}$ satisfies the conditions (i), (ii) of Theorem 4.5 on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
Proposition 4.7. Assume that $0 \leq \mu<1$ and $\lambda_{k} \leq \lambda<\lambda_{k+1}$. There exist $\rho, \alpha>0$ such that $\left.I_{\mu, \lambda}\right|_{\partial B_{\rho} \cap X} \geq \alpha$ where $X=\operatorname{span}\left\{\phi_{l} ; l>k\right\}$.
Proof : For any $u \in X, \lambda_{k} \leq \lambda<\lambda_{k+1}$, we obtain from equation (4.6), Hardy's inequality and Poincare's inequality that

$$
\begin{aligned}
I_{\mu, \lambda}(u) & =\frac{1}{2} \int_{\Omega}\left[\left|\nabla_{\mathbb{H}^{d}} u\right|^{2}-\mu V|u|^{2}\right] d z d s \\
& -\frac{1}{p} \int_{\Omega}|u|^{p} d z d s-\frac{\lambda}{2} \int_{\Omega}|u|^{2} d z d s \\
& \geq \frac{1}{2} \frac{\lambda_{k+1}-\lambda}{\lambda_{k+1}}\|u\|_{\mu}^{2}-\frac{C_{1}}{p}\|u\|_{\mu}^{p} \\
& \geq C_{1}\|u\|_{\mu}^{2}\left(\frac{1}{2 C_{1}} \frac{\lambda_{k+1}-\lambda}{\lambda_{k+1}}-\frac{1}{p}\|u\|_{\mu}^{p-2}\right) .
\end{aligned}
$$

Let

$$
\rho_{0}=\left(\frac{p}{2 C_{1}}\left(1-\frac{\lambda}{\lambda_{k+1}}\right)\right)^{\frac{1}{p-2}}>0
$$

Then for any $0<\rho<\rho_{0}$,

$$
\begin{equation*}
\inf _{u \in X ;\|u\|_{\mu}=\rho} I_{\mu, \lambda}(u)=\alpha(\rho)>0 . \tag{4.7}
\end{equation*}
$$

Thus $I_{\mu, \lambda}$ satisfies the condition (i) of Theorem 4.5.
Proposition 4.8. Assume that $0 \leq \mu<1$ and $\lambda_{k} \leq \lambda<\lambda_{k+1}$. Then $I_{\mu, \lambda}$ verifies (ii) of Theorem 4.5 with $e=\phi_{k+1}$ and $Y=\operatorname{span}\left\{\phi_{1}, \cdots, \phi_{k}\right\}$.
Proof : We prove Proposition 4.8 using Remark 4.6. For any $y \in Y$, we have from (4.5) that

$$
\begin{aligned}
I_{\mu, \lambda}(y) & =\frac{1}{2} \int_{\Omega}\left[\left|\nabla_{\mathbb{H}^{d}} y\right|^{2}-\mu V|y|^{2}\right] d z d s \\
& -\frac{1}{p} \int_{\Omega}|y|^{p} d z d s-\frac{\lambda}{2} \int_{\Omega}|y|^{2} d z d s \\
& \leq \frac{1}{2} \frac{\lambda_{k}-\lambda}{\lambda_{k}}\|y\|_{\mu}^{2}-\frac{1}{p}\|y\|_{L^{p}(\Omega)}^{p} . \\
& \leq 0
\end{aligned}
$$

Let $e=\phi_{k+1}$ the $(k+1)$-th eigenfunction of $L_{\mu}$ and $y \in Y$, let us the following claim

$$
\begin{equation*}
I_{\mu, \lambda}\left(y+\theta \phi_{k+1}\right) \rightarrow-\infty \text { as } \theta \rightarrow+\infty \tag{4.8}
\end{equation*}
$$

which will prove Remark 4.6.
Since $\left\{\phi_{j}\right\}$ is an orthonormal basis of $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, we have for all $i, j \in \mathbb{N}, i \neq j$

$$
\begin{equation*}
\int_{\Omega}\left[\nabla_{\mathbb{H}^{d}} \phi_{i} \overline{\nabla_{\mathbb{H}^{d}} \phi_{j}}-\mu V \phi_{i} \overline{\phi_{j}}\right] d z d s=\lambda_{i} \int_{\Omega} \phi_{i} \overline{\phi_{j}} d z d s=0 \tag{4.9}
\end{equation*}
$$

Let $y=\sum_{i=1}^{k} \alpha_{i} \phi_{i} \in Y$, then for $\theta \geq 1$,

$$
\begin{aligned}
I_{\mu, \lambda}\left(y+\theta \phi_{k+1}\right) & =I_{\mu, \lambda}(y)+I_{\mu, \lambda}\left(\theta \phi_{k+1}\right)+\int_{\Omega}\left[\nabla_{\mathbb{H}^{d}} y \overline{\nabla_{\mathbb{H}^{d}}\left(\theta \phi_{k+1}\right)}-\mu V y \overline{\left(\theta \phi_{k+1}\right)}\right] d z d s \\
& -\lambda \int_{\Omega} y \overline{\left(\theta \phi_{k+1)}\right.} d z d s-\frac{1}{p} \int_{\Omega}\left|y+\theta \phi_{k+1}\right|^{p} d z d s \\
& +\frac{1}{p} \int_{\Omega}|y|^{p} d z d s+\frac{1}{p} \int_{\Omega}\left|\theta \phi_{k+1}\right|^{p} d z d s \\
& =I_{\mu, \lambda}(y)+I_{\mu, \lambda}\left(\theta \phi_{k+1}\right)-\frac{1}{p} \int_{\Omega}\left|y+\theta \phi_{k+1}\right|^{p} d z d s \\
& +\frac{1}{p} \int_{\Omega}|y|^{p} d z d s+\frac{1}{p} \int_{\Omega}\left|\theta \phi_{k+1}\right|^{p} d z d s .
\end{aligned}
$$

By using the following inequality

$$
\begin{equation*}
|a+b|^{p} \geq|a|^{p}+|b|^{p}-c_{p}\left(|a|^{p-1}|b|+|a||b|^{p-1}\right), \quad \forall a, b \in \mathbb{R}, p>1 \tag{4.10}
\end{equation*}
$$

and from the fact $I_{\mu, \lambda}(y) \leq 0$, we have

$$
\begin{aligned}
I_{\mu, \lambda}\left(y+\theta \phi_{k+1}\right) & \leq I_{\mu, \lambda}\left(\theta \phi_{k+1}\right)+c_{p} \int_{\Omega}\left(|y|^{p-1}\left|\theta \phi_{k+1}\right|+|y|\left|\theta \phi_{k+1}\right|^{p-1}\right) d z d s \\
& \leq \frac{\theta^{2}}{2}\left(\left\|\phi_{k+1}\right\|_{\mu}^{2}-\lambda\left\|\phi_{k+1}\right\|_{L^{2}(\Omega)}^{2}\right)-\frac{\theta^{p}}{p}\left\|\phi_{k+1}\right\|_{L^{p}(\Omega)}^{p} \\
& +c_{p} \int_{\Omega}\left(|y|^{p-1}\left|\theta \phi_{k+1}\right|+|y|\left|\theta \phi_{k+1}\right|^{p-1}\right) d z d s \\
& \leq \frac{\theta^{2}}{2}\left\{\left(\lambda_{k+1}-\lambda\right)\left\|\phi_{k+1}\right\|_{L^{2}(\Omega)}^{2}-\frac{2 \theta^{p-2}}{p}\left\|\phi_{k+1}\right\|_{L^{p}(\Omega)}^{p}\right. \\
& \left.+2 c_{p} \int_{\Omega}\left(\theta^{-1}|y|^{p-1}\left|\phi_{k+1}\right|+\theta^{p-3}|y|\left|\phi_{k+1}\right|^{p-1}\right) d z d s\right\} .
\end{aligned}
$$

Now $p>2$ and $\left\|\phi_{k+1}\right\|_{L^{p}(\Omega)}^{p}>0$ imply

$$
I_{\mu, \lambda}\left(y+\theta \phi_{k+1}\right) \rightarrow-\infty \text { as } \theta \rightarrow+\infty
$$

We have proved Proposition 4.8.

Now Proposition 4.7 and Proposition 4.8 imply that, if $0 \leq \mu<1$ and $\lambda_{k} \leq \lambda<\lambda_{k+1}$, $I_{\mu, \lambda}$ satisfies the assumptions of Theorem 4.5 and the Palais-Smale condition with

$$
\begin{aligned}
& \quad c=\inf _{h \in \Gamma} \max _{u \in A} I_{\mu, \lambda}(h(u)) \\
& \text { and } A \equiv\left(\bar{B}_{T} \cap Y\right) \oplus\{\theta e, 0<\theta<T\}>0
\end{aligned}
$$

Then $I_{\mu, \lambda}$ has a critical value $c$ and a nontrivial critical point $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, since $I_{\mu, \lambda}(u)=c \geq \alpha>0$. We have proved Theorem 4.1.

## 5. The Palais-Smale condition

We prove now Theorem 4.3, i.e. we verify the $(P S)$ condition for $I_{\mu, \lambda}$ on $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
Lemma 5.1. Let $0 \leq \mu<1, \lambda>0, c \in \mathbb{R}$, then any sequence $\left\{u_{n}\right\} \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ satisfying

$$
\begin{equation*}
I_{\mu, \lambda}\left(u_{n}\right) \rightarrow c \quad \text { and } \quad\left\|I_{\mu, \lambda}^{\prime}\left(u_{n}\right)\right\|_{H^{-1}\left(\Omega, \mathbb{H}^{d}\right)} \rightarrow 0 \tag{5.1}
\end{equation*}
$$

is bounded in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
Proof: We use

$$
\|u\|_{\mu}=\left(\int_{\Omega}\left[\left|\nabla_{\mathbb{H}^{d}} u\right|^{2}-\mu V(z, s)|u|^{2}\right] d z d s\right)^{1 / 2} \approx\left\|\nabla_{\mathbb{H}^{d}} u\right\|_{L^{2}(\Omega)}
$$

as norm of $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$. Let $\left\{u_{n}\right\} \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ be the sequence of Lemma 5.1, then

$$
\begin{equation*}
\frac{1}{2}\|u\|_{\mu}^{2}-\frac{1}{p}\left\|u_{n}\right\|_{L^{p}(\Omega)}^{p}-\frac{\lambda}{2}\left\|u_{n}\right\|_{L^{2}(\Omega)}^{2} \rightarrow c \tag{5.2}
\end{equation*}
$$

and

$$
\begin{align*}
\left\langle I_{\mu, \lambda}^{\prime}\left(u_{n}\right), \varphi\right\rangle & =\int_{\Omega}\left[\nabla_{\mathbb{H}^{d}} u_{n} \nabla_{\mathbb{H}^{d}} \bar{\varphi}-\left(\mu V(z, s) u_{n}+\left|u_{n}\right|^{p-2} u_{n}+\lambda u_{n}\right) \bar{\varphi}\right] d z d s  \tag{5.3}\\
& =\circ(1)\|\varphi\|_{H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)} .
\end{align*}
$$

We suppose that $\left\{u_{n}\right\}$ is not bounded in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, i. e.

$$
\left\|\nabla_{\mathbb{H}^{d}} u_{n}\right\|_{L^{2}(\Omega)} \rightarrow+\infty
$$

and let $w_{n}=\frac{u_{n}}{\left\|\nabla_{\mathbb{H}^{d}} u_{n}\right\|_{L^{2}(\Omega)}}$, then $\left\|\nabla_{\mathbb{H}^{d}} w_{n}\right\|_{L^{2}(\Omega)}=1$. So, there exist $w \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ and a subsequence still denoted by $\left\{w_{n}\right\}$ such that

$$
w_{n} \rightharpoonup w \text { in } H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right) .
$$

Now, for any $q=\nu 2+(1-\nu) 2^{*}$ with $0<\nu<1$, we have

$$
\left\|w_{n}-w\right\|_{L^{q}(\Omega)}^{q} \leq\left\|w_{n}-w\right\|_{L^{2}(\Omega)}^{\nu}\left\|w_{n}-w\right\|_{L^{2^{*}}(\Omega)}^{(1-\nu) 2^{*}} .
$$

Then, Sobolev's inequality (1.2) and the compact embedding result of Lemma 3.2 imply the following strong limit

$$
w_{n} \rightarrow w \text { in } L^{q}(\Omega) \text { for } 1 \leq q<2^{*},
$$

which implies also the weak limit, for $2<p<2^{*}$,

$$
\int_{\Omega}\left|w_{n}\right|^{p-2} w_{n} \bar{\varphi} d z d s \rightarrow \int_{\Omega}|w|^{p-2} w \bar{\varphi} d z d s
$$

for any $\varphi \in \mathcal{C}_{0}^{\infty}(\Omega)$. Divide (5.3) by $\left\|\nabla_{\mathbb{H}^{d}} u_{n}\right\|_{L^{2}(\Omega)}$, we get for any $\varphi \in \mathcal{C}_{0}^{\infty}(\Omega)$,

$$
\begin{align*}
& \int_{\Omega}\left[\nabla_{\mathbb{H}^{d}} w_{n} \nabla_{\mathbb{H}^{d}} \bar{\varphi}-\mu V(z, s) w_{n} \bar{\varphi}\right] d z d s  \tag{5.4}\\
& -\int_{\Omega}\left|u_{n}\right|^{p-2} w_{n} \bar{\varphi} d z d s-\lambda \int_{\Omega} w_{n} \bar{\varphi} d z d s \\
& =\frac{\circ(1)}{\left\|\nabla_{\mathbb{H}^{d}} u_{n}\right\|_{L^{2}(\Omega)}\|\varphi\|_{H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)}} .
\end{align*}
$$

Passing to the limit, we obtain for any $\varphi \in \mathcal{C}_{0}^{\infty}(\Omega)$,

$$
\begin{aligned}
& \int_{\Omega}\left[\nabla_{\mathbb{H}^{d}} w \nabla_{\mathbb{H}^{d}} \bar{\varphi}-\mu V(z, s) w \bar{\varphi}\right] d z d s-\lambda \int_{\Omega} w \bar{\varphi} d z d s \\
& =\lim _{n \rightarrow \infty} \int_{\Omega}\left|u_{n}\right|^{p-2} w_{n} \bar{\varphi} d z d s \\
& =\lim _{n \rightarrow \infty}\left\|\nabla_{\mathbb{H}^{d}} u_{n}\right\|_{L^{2}(\Omega)}^{p-2} \int_{\Omega}\left|w_{n}\right|^{p-2} w_{n} \bar{\varphi} d z d s,
\end{aligned}
$$

which implies that

$$
\int_{\Omega}|w|^{p-2} w \bar{\varphi} d z d s=0 \quad \forall \varphi \in \mathcal{C}_{0}^{\infty}(\Omega) .
$$

Then $w=0$ a.e. in $\Omega$ and so

$$
\lim _{n \rightarrow \infty} \int_{\Omega}\left|w_{n}\right|^{2} d z d s=\int_{\Omega}|w|^{2} d z d s=0
$$

Divide (5.2) by $\left\|\nabla_{\mathbb{H}^{d}} u_{n}\right\|_{L^{2}(\Omega)}^{2}$, we obtain

$$
\begin{align*}
& \frac{1}{2} \int_{\Omega}\left[\left|\nabla_{\mathbb{H}^{d}} w_{n}\right|^{2}-\mu V(z, s)\left|w_{n}\right|^{2}\right] d z d s  \tag{5.5}\\
& -\frac{1}{p} \int_{\Omega}\left|u_{n}\right|^{p-2}\left|w_{n}\right|^{2} d z d s-\frac{\lambda}{2} \int_{\Omega}\left|w_{n}\right|^{2} \quad d z d s=\circ(1),
\end{align*}
$$

where

$$
\circ(1)=\frac{c}{\left\|\nabla_{\mathbb{H}^{d}} u_{n}\right\|_{L^{2}(\Omega)}^{2}} .
$$

Passing to the limit in (5.5),

$$
\begin{equation*}
\frac{1}{2}-\frac{1}{2} \lim _{n \rightarrow+\infty} \int \mu V(z, s)\left|w_{n}\right|^{2} d z d s=\frac{1}{p} \lim _{n \rightarrow+\infty} \int_{\Omega}\left|u_{n}\right|^{p-2}\left|w_{n}\right|^{2} d z d s, \tag{5.6}
\end{equation*}
$$

and taking $\varphi=w_{n}$ in (5.4), we have

$$
\begin{equation*}
1-\lim _{n \rightarrow+\infty} \int \mu V(z, s)\left|w_{n}\right|^{2} \quad d z d s=\lim _{n \rightarrow+\infty} \int_{\Omega}\left|u_{n}\right|^{p-2}\left|w_{n}\right|^{2} d z d s \tag{5.7}
\end{equation*}
$$

We multiply (5.6) by $p$, we obtain from (5.7) and Hardy's inequality that

$$
\begin{aligned}
\frac{p}{2}-1 & =\left(\frac{p}{2}-1\right) \lim _{n \rightarrow+\infty} \int \mu V(z, s)\left|w_{n}\right|^{2} d z d s \\
& \leq\left(\frac{p}{2}-1\right) \lim _{n \rightarrow+\infty} \mu \int\left|\nabla_{\mathbb{H}^{d}} w_{n}\right|^{2} d z d s \\
& \leq\left(\frac{p}{2}-1\right) \mu .
\end{aligned}
$$

That is impossible since $\mu<1, p>2$ and this implies that $\left\{u_{n}\right\}$ is bounded in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.
Lemma 5.2. Under the assumption of Lemma 5.1, $\left\{u_{n}\right\}$ possesses a convergent subsequence in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$.

Proof : Since $\left\{u_{n}\right\}$ is a bounded sequence in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$, there exists $u \in H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$ such that for any $2<p<2^{*}$,

$$
\begin{aligned}
u_{n} & \rightharpoonup u \quad \text { in } H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right), \\
u_{n} & \rightarrow u \text { in } L^{p}(\Omega) \\
\left|u_{n}\right|^{p-2} u_{n} & \rightarrow|u|^{p-2} u \quad \text { in } L^{\frac{p}{p-1}}(\Omega)
\end{aligned}
$$

By passing to the limit in (5.3), we have

$$
\left\langle I_{\mu, \lambda}^{\prime}(u), \varphi\right\rangle=0
$$

for any $\varphi \in \mathcal{C}_{0}^{\infty}(\Omega)$ which is dense in $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$. Then

$$
I_{\mu, \lambda}^{\prime}\left(u_{n}\right) \rightarrow 0 \text { in } H^{-1}\left(\Omega, \mathbb{H}^{d}\right)
$$

implies that

$$
\lim _{n \rightarrow \infty}\left\langle I_{\mu, \lambda}^{\prime}\left(u_{n}\right)-I_{\mu, \lambda}^{\prime}(u), u_{n}-u\right\rangle=0
$$

We rewrite the norm of $H_{0}^{1}\left(\Omega, \mathbb{H}^{d}\right)$

$$
\begin{aligned}
\left\|u_{n}-u\right\|_{\mu}^{2} & =\left\langle I_{\mu, \lambda}^{\prime}\left(u_{n}\right)-I_{\mu, \lambda}^{\prime}(u), u_{n}-u\right\rangle \\
& +\int_{\Omega}\left[\left|u_{n}\right|^{p-2} u_{n}-|u|^{p-2} u\right] \overline{\left(u_{n}-u\right)} d z d s+\lambda \int_{\Omega}\left|u_{n}-u\right|^{2} d z d s
\end{aligned}
$$

From Hölder's inequality, we have

$$
\begin{aligned}
& \int_{\Omega}\left|\left[\left|u_{n}\right|^{p-2} u_{n}-|u|^{p-2} u\right] \overline{\left(u_{n}-u\right)}\right| d z d s \\
& \leq\left\|\left|u_{n}\right|^{p-2} u_{n}-|u|^{p-2} u\right\|_{L^{\frac{p}{p-1}}(\Omega)}\left\|u_{n}-u\right\|_{L^{p}(\Omega)}
\end{aligned}
$$

so we can deduce that $\left\|u_{n}-u\right\|_{\mu}^{2} \rightarrow 0, n \rightarrow \infty$.

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