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# Growth and structural properties of CuAg and CoPt bimetallic nanoparticles

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Core/shell CuAg and alloyed CoPt have been synthesized using two vapor phase deposition <sup>10</sup> techniques. For CuAg prepared by Thermal Evaporation (TE), the size and the morphology of the Cu cores are the key parameters to promote the formation of the core/shell arrangement. For CoPt synthesized by Pulsed Laser Deposition (PLD), the growth kinetics of nanoparticles, depending on the deposition rate, the substrate nature and the temperature, controls the nanoparticles morphology. The competition between the growth and the ordering kinetics govern the <sup>15</sup> nanoparticles structure. By reducing the growth kinetics, as-grown L1<sub>0</sub> ordered nanoparticles are obtained according to the bulk phase diagram.

#### **A. Introduction**

Multi-components nanostructures have recently attracted much attention, from both experimental and theoretical point <sup>20</sup> of view, since it is possible to enhance or to obtain new physical and chemical properties that cannot be obtained in the single component nanoparticules. When metals are combined in core/shell configurations or alloyed structures, one can observed dramatic changes in chemical, <sup>1-4</sup> magnetic<sup>5,6</sup> <sup>25</sup> and optical properties<sup>7-9</sup> compared to those of the individual

- <sup>25</sup> and optical properties compared to those of the individual component. This comes in addition to the size which directly influences the physico-chemical properties when the characteristic length of a physical property becomes similar to or larger than the particle size.
- <sup>30</sup> For both fundamental and technological purposes, new developments are strongly dependent on the ability to have accurate control on the growth of nanoparticles. It is essential to control chemical composition, crystallography, mean size, size distribution and morphology to tailor the physical and
- <sup>35</sup> chemical properties of the nanoparticles. Up to now, bimetallic core/shell or alloyed nanostructures have been synthesized by several chemical and physical methods: colloids chemistry,<sup>10,11</sup> ion implantation,<sup>12</sup> as well as vapour phase deposition techniques like sputtering<sup>13,14</sup> and electron-

<sup>40</sup> beam evaporation.<sup>15,16</sup> Among the vapour phase deposition techniques, multiple targets sequential pulsed laser deposition (PLD) has appeared to be one of the most flexible and promising growth technique because of its capability to control the synthesis by varying parameters such as the laser
<sup>45</sup> energy, the pulse repetition rate and fluence (i.e. energy density), the target to substrate distance, *etc.* This technique has been widely and successfully used in the preparation of monometallic and semiconductor nanoparticules as well as thin oxides films with controlled composition due to its ability.
<sup>50</sup> to very well reproduce the stoichiometry of the ablated target. 17

In this paper, we investigate the growth and the structural properties of core/shell CuAg and alloyed CoPt nanoparticles prepared by thermal evaporation (TE) and pulsed laser <sup>555</sup> deposition techniques respectively. For the CuAg system, we report on the size effects governing the core/shell arrangement, through the study of morphological transitions of the Cu cores depending on their sizes. Concerning the CoPt nanoparticles, we present the influence of the substrate and the laser frequency of the PLD experiment on the growth and the structure of the nanoparticles. We will show how it is possible, by controlling the growth kinetics of nanoparticles, to control the structure and the morphology of the nanoparticles.

#### **B.** Experimental procedure

CuAg and CoPt nanoparticles have been prepared on amorphous substrates using the same ultra-high vacuum (UHV) chamber either by classical thermal evaporation and 70 pulsed laser deposition. During the evaporation, the residual pressure in the chamber is close to 10<sup>-8</sup> mbar. The growth of the nanoparticles has been performed on two different substrates: a 10 nm thick amorphous carbon film (a-carbon) of

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a commercial 3mm Cu grid or a 2 nm layer of amorphous  $_{75}$  alumina (a-Al<sub>2</sub>O<sub>3</sub>), deposited using PLD on the top of the amorphous carbon. The deposition rate of each element is controlled by an *in situ* quartz crystal monitor, which indicates the nominal thickness of deposited materials on the quartz surface, in a continuous thin film approximation. As

- <sup>80</sup> the metallic species do not wet the two amorphous substrates used here, nanoparticles are formed instead of a continuous thin films and the nominal thickness of a sample obviously do not corresponds to the particle thickness. During the deposition, the substrate is held at a fixed temperature T<sub>s</sub>, <sup>85</sup> between 25°C and 750°C. The heating system is a tungsten filament heated by Joule effect, positioned just behind the TEM grid. The heat transfer occurs through radiation from the filament to the substrate and conduction from substrate to particles. At the end of synthesis, samples were covered by a
- $_{90}$  3 nm thick a-Al<sub>2</sub>O<sub>3</sub> layer deposited by PLD, in order to protect the nanoparticles from air oxidation. Nevertheless, the high vacuum conditions of the deposition chamber offer the possibility to anneal *in situ* the synthesized particles before the deposition of the protective amorphous alumina layer.

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#### 1. Pulsed Laser Deposition

A typical target-substrate holder configuration is used to deposit separately the a-Al<sub>2</sub>O<sub>3</sub>, Co and Pt by PLD, using a 100 KrF excimer laser at 248 nm. The pulse duration is 25 ns and the repetition rate in the range of 1 to 15 Hz. The laser irradiation produces the plasma of the different species and vaporised atoms condense on the substrate placed at 5 cm from the target. The growth of the CoPt nanoparticles has 105 been performed on amorphous carbon and amorphous alumina. CoPt nanoparticles are obtained by alternated irradiation of two pure metal targets allowing a composition controlled synthesis.<sup>18,19</sup> To insure a good homogeneity of the particles composition, the nominal thickness of deposited 110 materials at each successive step (corresponding to one series of laser pulses on Co and Pt targets) is set to 0.1 nm. The number of pulses on Co and Pt targets at each step was fixed in order to obtain a composition closed to Co<sub>50</sub>Pt<sub>50</sub>. The laser energy can be chosen in the range of 150 to 250 mJ depending 115 on the ablation threshold of the selected target. The deposition rate depends on the target nature, the energy and the frequency of the laser. We used energies of 190 mJ and 250 mJ for the cobalt and the platinum targets respectively. These selected energies avoid the formation of micrometric droplets 120 and minimizes the deposition rate which is, as we will show, a key parameter to control the particles morphology. With this

configuration, deposition rates for Co and Pt were respectively 0.30 and 0.04 nm per minute.

#### 125 2. Thermal evaporation

The sources were Knudsen type cells containing high purity Cu and Ag pellets heated at 1040°C and 970°C respectively. Copper was first evaporated to form pure Cu nanoparticles at the surface of the a-carbon substrate. In a second step, Ag was <sup>130</sup> evaporated. During the deposition of Cu and Ag, the substrate is heated at a temperature of 400°C. This allows the diffusion of the two species on the surface of the substrate. The nanoparticles are annealed during half an hour at the same temperature. Typically, deposition rates for Cu and Ag were <sup>135</sup> around 0.25 nm per minute.

#### 3. Characterisation tools

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Structural characterisations of the nanoparticles were performed using several TEM techniques: conventional Bright <sup>140</sup> Field imaging (BF), High Resolution imaging (HRTEM), Energy Filtered imaging (EFTEM) and electron tomography. The microscopes used for these studies were a JEOL JEM-2100F and a ZEISS LIBRA 200 FE field-emission TEM operating at 200 keV.

#### C. CuAg Core/shell nanoparticles

Success in obtaining bimetallic core/shell nanoparticles by <sup>150</sup> physical route mainly depends on : 1/ the miscibility of the two metals, 2/ the surface energies of those metals and 3/ the interfacial energies between a substrate and the two metals. The CuAg system was chosen as a model because of silver and copper being immiscible elements (0.1 % solid solution at <sup>155</sup> 298 K). A lower surface energy for Ag (1210 mJ/m<sup>2</sup> compared to 2130 mJ/m<sup>2</sup> for Cu)<sup>20</sup> favours the silver surface segregation. CuAg core/shell nanoparticles are then expected to be formed on the amorphous carbon substrate, with Cu core and Ag shell. To improve the efficiency of the core/shell <sup>160</sup> formation, a sequential procedure has been used during the thermal evaporation, with deposition of Cu firstly, and Ag in a second step.

#### 1.6@3 system

- <sup>165</sup> Nominal thicknesses of 6 nm and 3 nm for Cu and Ag respectively were deposited and we obtained nanoparticles sizes around 25 nm (Fig. 1a). In order to study the growth of core/shell nanoparticles on the substrate, the samples (denoted 6@3 in the following) were fully characterised by
  <sup>170</sup> transmission electron microscopy. The EFTEM technique was particularly efficient to identify the Ag shells and the Cu cores. For each element, a chemical map was obtained on which the black and white contrast corresponds to the localisation of the selected element (Fig. 1b and 1c).
  <sup>175</sup> Gathering together the Cu and Ag chemical maps results in a color coded image as the one presented on figure 1d for the 6@3 system. Such chemical information cannot be obtained directly from the classical BF image (Fig. 1a) since it shows no contrast between the two elements.
- 80 Many different configurations regarding the silver localisation are noticed. Silver can form continuous shells

around the Cu cores, as expected from the thermodynamical considerations discussed before, but as well silver can grow either independently from the Cu cores, forming pure silver

- <sup>185</sup> nanoparticles, or onto the Cu cores, but only on one side of the Cu nanoparticles. The EFTEM images can hardly bring proof of a covering of the Cu cores by Ag because of the projection along the z axis perpendicular to the substrate to form the images. The complete capping of the Cu core by the
- <sup>190</sup> Ag shell is evidenced by a well-known optical effect called 'Moirés patterns' that can be seen on HRTEM or even conventional TEM images (Fig. 2a). Large dark stripes superimposed to the lattice fringes of the crystal arise when two crystals are overlapped. These patterns appear when <sup>195</sup> interferences occur between two slightly different spatial frequencies (in orientation and/or spacing) in the reciprocal space. Here, the Moiré patterns originates from the epitaxial relationship between the core and the shell.

The crystallographic orientation is the same for the core <sup>200</sup> and the shell, and this was noticed for all the observed nanoparticles. This can be seen on the power spectrum on Fig. 2b, on which Ag and Cu crystals are oriented along the same [110] zone axis of the Face Centred Cubic (FCC) structure, with the same radial orientation of the 200 reflections of the

<sup>205</sup> two crystals. The mirror corresponding to the twin plane is reported on the power spectrum. This is a striking example of the shell following exactly the underlying defect of the core.

On figure 3, showing a CuAg nanoparticle for which silver did not form a perfect shell around the core, the same <sup>210</sup> epitaxial relationship is reported on the [100] zone axis given in the inset of figure 3b. Because of the difference in lattice parameters between the two materials, interfacial edge dislocations would be expected. Since the interface between the two crystals lies on a (220) plane, we carried out a Bragg <sup>215</sup> filtering of the image using the 220 type reflections of Cu and Ag. This allows to evidence the localisation of the excess (220) Cu planes composing the misfit dislocations. On Braggfiltered images (Fig. 3b), the interfacial dislocations spacing is 1.1 nm, that is, around 8 (220)<sub>Ag</sub> interplanes distances. When <sup>220</sup> estimating the distance between interfacial dislocations for an Ag film grown on a Cu film along the [220] direction from the

simple formula<sup>21</sup>:

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$$n = d_{(200)Ag} \ / \ (d_{(200)Cu} - d_{(200)Ag})$$

one would find 8d<sub>(200)Ag</sub>. The rule of thumb used here is still valid at the nanoscale because we found the same dislocations periodicity on the image. This is a surprising situation since dislocations are usually drived to free surfaces in order to <sup>230</sup> decrease the elastic field around their core. Here, the epitaxial stress is important enough (a<sub>Ag</sub> is 11% larger than a<sub>Cu</sub>) to maintain 6 perfect dislocations inside approximately a 5x5x5 nm<sup>3</sup> volume. The presence of interfacial dislocations reveals also that Ag and Cu have recovered their own lattice <sup>235</sup> parameter. This has been verified by checking the lattice parameters on both sides of the nanoparticles. This is coherent with another simple model used for heteroepitaxial films, giving the critical thickness over which the interface is no more coherent and misfit dislocations appear.<sup>22</sup> This critical <sup>240</sup> thickness  $t_c$  is roughly given by :

$$t_c = b / \epsilon$$

where b is the Burger's vector norm (1/2 [110] for the FCC <sup>245</sup> structures) and  $\varepsilon$  the misfit between the two lattice parameters defined by  $\varepsilon = (a_{Ag} - a_{Cu})/a_{Cu}$ . Applying this model leads to a critical thickness close to 2 nm, to be compared with the thickness of the Ag layer on figure 3a which is around 2.5 nm.

#### 250 2. Cu morphological transitions

From figure 1d, it is worth noting that the core/shell nanoparticle sizes are the smallest (<12 nm) from the size distribution of the nanoparticles. Over this limit, most of the CuAg nanoparticles do not present the core/shell 255 configuration. The size of the pre-existing Cu cores being unambiguously an important parameter in the formation of the core/shell arrangement, we turned to the structural characterisation of the Cu cores alone, especially under the identified size limit. Controlling the size of the Cu 260 nanoparticles can be done by two different means. The first way would be to lower the substrate temperature during the formation of the Cu core, decreasing the Cu diffusion on the surface. Cu atoms would be then fixed more efficiently on the nucleating sites, resulting in an increased number of particles 265 per unit area, and subsequently smaller particle sizes. The other mean is to decrease the quantity of matter during deposition, keeping constant all the other parameters, especially the substrate temperature. The density of particles per square nanometer remains the same, and the mean 270 nanoparticle size decreases because of the lack of atoms available for the growth. As for gold, Cu exhibits some morphological transitions between three different types: 1/ icosahedral shape constituted by 20 slightly distorted FCC tetrahedra gathered to form a 20 face polyhedron, 2/ 275 decahedral shape formed by 5 identical FCC sectors separated by imperfect twins, giving a 5-fold symmetry polyhedron, and 3/ rounded or facetted monocrystalline FCC structure. These transitions have been predicted using tight-binding semiempirical potentials for metal-metal interactions<sup>23</sup> and 280 observed in a lot of published works, for clusters either free or supported on a substrate. We conducted HRTEM experiments on a sample with 2 nm of Cu deposited at 400°C to check whether the synthesis method used here reproduces the same morphological transitions. The work by Koga et al.<sup>24</sup> was very 285 helpful to identify the different morphologies by interpreting the lattice fringes contrast on the HRTEM images. Although a good statistical sampling is difficult to realize, we obtain the following results, gathered on figure 4.

Due to the relatively broad size distribution of the Cu <sup>290</sup> nanoparticles, it is clear that the three types of morphologies and structures are present on the substrate, with a strong overlapping between icosahedral and decahedral domains. The transition from decahedra to monocrystalline FCC structures is sharper, with FCC Cu nanoparticles sizes starting from 8 <sup>295</sup> nm. The size limit over which core/shell arrangements are no more formed corresponds well to the decahedra-to-FCC structural transition.

#### 3. Dependence with Cu core sizes

<sup>300</sup> We carried out two other experiments to confirm that a Cu core size under 12 nm is the key parameter to obtain core/shell nanoparticles. The parameters were adjusted to obtain on one hand a Cu core sizes distribution centred under 12 nm, and on the other hand a distribution largely over 12 <sup>305</sup> nm.

The first CuAg system has been synthesized with nominal thickness of 2 nm and 1 nm for respectively Cu and Ag (denoted 2@1 in the following), keeping a substrate temperature of 400°C and the ratio between the Cu and Ag

- <sup>310</sup> nominal thicknesses. An EFTEM image corresponding to this system is presented on figure 5a. A majority of the nanoparticles presents a core/shell configuration, with core sizes below 12 nm and shell thickness distribution centred on 2 nm (values ranging between 1 nm and 3.5 nm). On the same
- <sup>315</sup> area of the microscopy grid corresponding to the 2@1 system, size counting on the EFTEM and TEM images as well as statistical measurements using a micro-Raman probe<sup>13</sup> have been carried out. The statistical approach confirms the local measurements.

The second experiment aimed at the formation on the substrate of Cu cores with sizes over 12 nm. This was achieved by depositing the same Cu thickness (2 nm), but with a substrate temperature of 500°C (Fig. 5b). For Ag, no difference was noticed in the final microstructure between a

- <sup>325</sup> deposition at a substrate temperature of 500°C (as for Cu) and a deposition after cooling down to 400°C. Almost no particles present a core/shell structure, which confirms the assumption of a Cu core size limit for the core/shell formation. The explanation must be linked to the way silver grows on the Cu
- <sup>330</sup> nanoparticles. From crystal growth theory, it is well known that the growth rate is different depending on the Miller indices of the face exposed to incoming species. This remark must be related to the fact that for an icosahedron, all the faces present the same crystallographic orientation, that is
- 335 (111) faces. For a decahedron, the two bases are also (111) faces, and depending on the growth along the 5-fold axis, the decahedron show (100) lateral faces. The situation is completely different for the monocrystalline FCC structure: the nanoparticles exhibit either no faceting, or faces typical
- <sup>340</sup> from a cuboctahedron which show (111), (200) and (220) faces. For the two morphologies corresponding to sizes less than 10 nm, the growth of the Ag shell will be far more isotropic, and hence the thickness will be more homogeneous all around the Cu cores. This hypothesis is sustained by

<sup>345</sup> HRTEM observations on the 2@1 system deposited with a substrate temperature of 400°C. Figure 6 shows an example of a Cu decahedron with a complete Ag shell all over the particle. The Moiré patterns evidence the presence of Ag in epitaxial relationship with the 5 subjacent Cu sectors. The <sup>350</sup> [110] zone axis for the 5 different FCC sectors composing the decahedron are visible on the power spectrum presented on

decahedron are visible on the power spectrum presented on figure 6b.

Moreover, Moiré patterns stop before the edges of the particles, which means that only Ag is present on the edges, s hence forming a continuous shell all around the particle. The shell thickness on this image (Fig. 6a) can then be estimated to about 2 nm. This is the limit predicted by the critical thickness model (Section C.1) over which the Ag shell is relaxed to its bulk lattice parameter.

#### 360

#### D. CoPt bimetallic alloyed nanoparticles

This part focuses on the influence of the nature of the substrate (a-Al<sub>2</sub>O<sub>3</sub> or a-carbon) and the synthesis parameters on the growth and the structure of CoPt nanoparticles. Bulk <sup>365</sup> equiatomic CoPt exhibits a phase transition at 825°C, between a tetragonal ordered phase (L1<sub>0</sub>) at low temperature and a disordered FCC structure at high temperature.<sup>25</sup> The L1<sub>0</sub> ordered phase, due to the alternative atomic stacking of pure Co and pure Pt layers, has large magnetocrystalline <sup>370</sup> anisotropy. By means of this property, important research efforts are devoted to study L1<sub>0</sub> ordered CoPt nanoparticles which are expected to be the future information storage media for extremely high density recording (EHDR).<sup>26,27</sup> The diffraction pattern of this structure presents superstructure <sup>375</sup> reflections characterizing the ordered structure.

#### 1. Influence of the substrate on the nanoparticles growth

CoPt nanoparticles thin film with 2 nm nominal thickness have been prepared at a substrate temperature of 600°C with a <sup>380</sup> laser frequency of 5 Hz on a-carbon and a-Al<sub>2</sub>O<sub>3</sub> substrates. Figures 7a and 7b show TEM bright field images of the CoPt nanoparticles prepared on the a-carbon and the a-Al<sub>2</sub>O<sub>3</sub> substrates respectively. The mean size, the polydispersity (standard deviation divided by the mean size of the <sup>385</sup> nanoparticles), the particles density and the coverage ratio for each sample are given in table 1.

The coverage ratio, defined as the percentage of the covered surface of the substrate by the nanoparticles, has been determined from TEM bright field images using Digital <sup>390</sup> Micrograph software. Note that the contrast of particles smaller than 2 nm is too weak to consider them in the statistical characterisation.

The growth mechanisms involved in a PLD experiment are quite different from those observed in classical thermal <sup>395</sup> evaporation techniques because the flux of atoms is discontinuous. These mechanisms can be described as follow: firstly, nucleation and size increase of the *nuclei* under high flux of atoms during the deposition time (i.e. several µs) and secondly, coalescence and Ostwald ripening between two <sup>400</sup> pulses (i.e. several 0.1 s). The coalescence takes place by atomic rearrangement of two close enough nanoparticles, leading to the formation of atoms inter-diffusion bridges.<sup>28</sup> The Ostwald ripening implies the growth of large particles at the expense of the ones smaller than the critical nuclei.<sup>29</sup> In

- <sup>405</sup> the case of PLD, due to the rapid decrease of supersaturation after a laser pulse, these two mechanisms occur alone during a long time and influence the nanoparticles morphology. As a consequence, in the late stage of the synthesis, new nucleation sites become available on the substrate between the biggest
- <sup>410</sup> particles, making possible the growth of smaller CoPt nanoparticles. This explains the high polydispersity observed on the a-carbon substrate that shows a high number of very small particles in the range size from 2 nm to 4 nm (inset Fig. 7a). This phenomenon is less important in the case of the a-
- $_{415}$  Al<sub>2</sub>O<sub>3</sub> substrate (inset Fig. 7b). It suggests that the diffusion of Co and Pt on a-carbon substrate is higher than the ones on a-Al<sub>2</sub>O<sub>3</sub>.

Since the nominal thickness is the same for both samples, by comparing the mean size and the coverage ratio of CoPt 420 nanoparticles on a-carbon and on a-Al2O3, we can deduce that the nanoparticles are more flattened on a-Al<sub>2</sub>O<sub>3</sub> substrate than on a-carbon substrate. This difference is closely related to the kinetic barriers involved during the nanoparticles formation. Kinetic energy of the atoms, given by the substrate 425 temperature governs the ability to overcome the various diffusion barriers (metal/metal and metal/substrate upstepping and downstepping activation energies).<sup>30,31</sup> At a given temperature, these diffusion barriers control the evolution of film morphology during the growth. From our experimental 430 results, the two-dimensional growth on alumina may be then the signature either of strong interactions of CoPt with alumina (higher activation barrier for the 3-dimensional growth compared to carbon) or a lower effective substrate temperature in the case of alumina. To summarize, we show 435 here that the kind of interactions between the substrate and the species deposited drive both the size distribution and the morphologies of the nanoparticles: the 3-dimensional growth of the nanoparticles is promoted on a-carbon whereas a-Al<sub>2</sub>O<sub>3</sub> substrate minimizes the size dispersion. The following 440 question remains: how can we know if the diffusion and the substrate/metal upstepping energy of the Co and Pt atoms are retained by a stronger a-Al<sub>2</sub>O<sub>3</sub>/metals interaction, or slowed

down by the lower temperature of the alumina surface? It is thus important to emphasis the thermal properties of <sup>445</sup> the substrate, which are generally not enough taken into account in nanoparticles synthesis on either monocrystalline

or amorphous substrates. Contrary to carbon, alumina is a good thermal insulator. As a rough indication, the heat capacity, at room temperature, of the crystalline alumina <sup>450</sup> (79.45 J/mol.K) is almost ten times higher than the one of graphite (8.58 J/mol.K).<sup>32</sup> This property of alumina is well

known since in the electronic and optic industries, amorphous alumina thin films are promising materials for thin thermal insulator layers.<sup>33</sup> However, it is difficult to differentiate the 455 respective role of the strong interactions of the deposited

species with the substrate and low thermal conductivity of the

 $a-Al_2O_3$  layer, because they both have the same influence on the morphology of CoPt nanoparticles.

#### 460 2. Thermal insulator properties of a-Al<sub>2</sub>O<sub>3</sub> substrate

To demonstrate the strong influence of the thermal conductivity of the substrate, we have synthesized nanoparticles on a-Al<sub>2</sub>O<sub>3</sub> with the same experimental conditions than the ones used for the synthesis of 465 nanoparticles shown in figure 7b. However, we have increased the thickness of the a-Al<sub>2</sub>O<sub>3</sub> substrate to 10 nm. As it can be seen on figure 8, the particles are very close to percolation threshold with irregular shapes elongated in the plane of the substrate. These flat morphologies compared to figure 7b indicate the tendency of the metals to wet the surface because 470 the thermal energy brought to the system is not high enough to permit 3D growth mechanism. The coverage ratio of 80% is close to the one obtained on particles formed on a 2 nm thick a-Al<sub>2</sub>O<sub>3</sub> layer heated at 400°C. The particles morphologies studied as a function of the a-Al<sub>2</sub>O<sub>3</sub> layer thickness clearly show the thermal influence of this substrate.

## **3.** Thermal influence of the substrate on the structure of the nanoparticles

<sub>480</sub> The structure of previous samples ( $T_s = 600^{\circ}C$ ) was FCC disordered which is a non equilibrium phase compared to the bulk material phase diagram. This FCC structure appears using the PLD synthesis route because the growth kinetics of the particles is faster than the ordering kinetics. To obtain asgrown L10 ordered nanoparticles, the substrate has to be heated over 650°C.<sup>18</sup> At  $T_s = 600$ °C, the ordering mechanisms require a post-synthesis annealing. CoPt nanoparticles, prepared in the same conditions as the ones studied in part D1, were annealed during 2 hours at 600°C and were then covered 490 by a thin a-Al<sub>2</sub>O<sub>3</sub> layer after the annealing. Figure 9a and 9b show the morphology of the nanoparticles after the annealing on the a-carbon and a-Al<sub>2</sub>O<sub>3</sub> substrates respectively. Since substrates are amorphous, the nanoparticles are randomly oriented on the surface and the diffraction patterns display a 495 typical ring pattern characteristic of a powder-like diffraction. Because of the very weak intensity of the superstructure reflections (see inset of figure 9c), we used imaging plate detectors which have a very high sensitivity and a large dynamic range. Moreover, to highlight the recorded intensity 500 we performed a rotational averaging intensity profile using Process Diffraction software. This intensity profile is drawn as a function of the norm of the diffraction vector  $g=2\pi/d$ where d is the interplanar distance. Figure 9c and 9d show respectively averaging plots corresponding to the diffraction 505 patterns of each sample: a-carbon / CoPt and a-Al<sub>2</sub>O<sub>3</sub> / CoPt. All rings are indexed according to the FCC disordered or the L1<sub>0</sub> ordered structures of the CoPt system.

The diffraction pattern of the particles annealed during 2h after their synthesis on a-carbon presents the 001, 110 and 211 <sup>510</sup> superstructure reflections characteristic of the L1<sub>0</sub> ordered lattice. On the contrary, the diffraction pattern of particles

grown and annealed on a-Al<sub>2</sub>O<sub>3</sub> presents only the fundamental reflections of the FCC disordered structure. This result has been confirmed by HRTEM. Energy Dispersive X-ray (EDX) <sup>515</sup> analyses show that particles are close to the equiatomic composition on both samples. The annealed particles mean size is smaller on a-carbon (7.8 nm) than the one on a-Al<sub>2</sub>O<sub>3</sub>

(8.8 nm). We can then rule out the possibilities that the different structural states observed on both samples are due to <sup>520</sup> deviation of the composition of the CoPt particles or to size effects. During the annealing, the contact surface between the substrate and the particles is limited. Taking into account the

particles size, the proportion of atom linked to the substrate is too low to consider a strong influence of the chemical <sup>525</sup> interaction on the ordering of the particles. It remains one assumption to explain this difference: the thermal insulator properties of amorphous alumina.

During the annealing, the increase of the mean particles size on a-carbon (20%) is twice higher than the one of the 530 particles on a-alumina (10%). It shows the low diffusion of Co and Pt atoms on a-Al<sub>2</sub>O<sub>3</sub> substrate. These alumina properties have been already highlighted by Takahashi et al.34 who compared the morphology of FePt nanoparticles after an annealing of 1h at 600°C in MgF<sub>2</sub> and a-Al<sub>2</sub>O<sub>3</sub> matrix. They 535 showed that the covering alumina layer reduces the coalescence and Oswald ripening phenomena avoiding the size increase of particles. Then they explained by size effects the disordered structure of the particles in alumina compared with the ordered ones in MgF<sub>2</sub> but they did not consider the 540 eventually different effective particles temperature. The goal of our structural studies is not to deny their interpretation but to emphasize the strong influence of the thermal properties of the substrate.

## **545 4. Influence of the laser frequency on the growth and the structure of the nanoparticles**

The influence of the laser frequency in the PLD experiment can be seen by comparing the morphology and structure obtained with a laser frequency of 5 Hz (Fig. 7a) with the 550 ones obtained with a lower frequency. We have synthesized a sample of 2 nm nominal thickness prepared on a-carbon substrate heated at 600°C with a laser frequency of 1 Hz. Figure 10a shows a TEM bright field image of this nanoparticles thin film. By decreasing the laser frequency 555 from 5 to 1 Hz, we have increased the polydispersity from 45 to 66% and the mean size from 6.3 to 9.2 nm. With a 1 Hz laser frequency, the time between two pulses is five times longer than the time at 5 Hz. As a consequence, the histogram in the inset of figure 10a shows clearly a bimodal size 560 distribution, the first population has a mean size of 3 nm whereas the second one has a mean size of 13 nm. It is the signature of the Ostwald ripening and coalescence mechanisms which occurs a longer time to influence the particles morphology. The synthesis of samples lasts 30 565 minutes and 150 minutes, using a laser frequency of 5 Hz and 1 Hz respectively. To dissociate the laser frequency effect from the synthesis time effect, we compared nanoparticles synthesized with a laser frequency of 5 Hz and annealed at 600°C during 120 minutes (Fig. 9a), to the ones obtained with <sup>570</sup> a laser frequency of 1 Hz (Fig. 10a). The higher mean size and polydispersity of the latter clearly show the influence of the laser frequency. The Oswald ripening and coalescence phenomena during the deposition at 1 Hz (150 minutes) is more effective than the ones occuring during the synthesis at 5 <sup>575</sup> Hz (30 minutes) and the post synthesis annealing (120 minutes).

The shape of the nanoparticles, fabricated with a 1 Hz laser frequency, has been determined by means of a 3D electron tomography experiment. The electron tomography technique allows the 3D morphology to be reconstructed from 2D observations of the nanoparticles. One or several series of projections of the nanoparticles have to be recorded in bright field imaging conditions. Such a series of projections is obtained by tilting the specimen in a range from -70° to 70° with respect to the electron beam and taking an image for each 585 tilt angle. Once the series of projections is recorded, the reconstruction of the object is computed by using a backprojection algorithm.<sup>35</sup> A bright field image of the selected area is presented in figure 10c and its 3D representation 590 (tomogram) is presented in figure 10d. The thickness of the particles has been measured by slicing the tomogram in the two perpendicular x and y directions (Fig. 10d). The thickness of the nanoparticles are measured directly on the sliced images, corresponding to the (x,z) or the (y,z) planes, by the 595 size of the nanoparticles along the z direction. For 80% of the nanoparticles, the measured thickness is equal to the diameter in the plane of the substrate. It can be explained by the 3D growth mechanisms involved between two pulses (metal/metal and metal/substrate upstepping), discussed in the part D1. The 600 shape of the last 20% of the nanoparticles, which is larger than thick, results from coalescence mechanisms.

A low laser pulses frequency during the growth also influences the structure of the nanoparticles. With a laser frequency of 1 Hz and  $T_s=600$  °C, the L1<sub>0</sub> structure is observed from electron diffraction of the particles and figure 10b shows a HRTEM image of a single L1<sub>0</sub> ordered CoPt nanoparticle oriented along the [110] zone axis. In such images, the order is clearly evidenced by the succession of pure Co and Pt planes along the [001] direction as indicated by the power spectrum of the image given in the inset of figure 10b.

To conclude, when decreasing the laser frequency, we  $^{615}$  reduce the nanoparticles growth kinetics, which is then slower than the ordering one. In these conditions, the diffusion of the metallic atoms at 600°C allows to obtain as-grown nanoparticles with the equilibrium structure (L1<sub>0</sub>) according to the bulk phase diagram.

#### Conclusion

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Growth and structural properties of immiscible CuAg and alloyed CoPt nanoparticles have been studied. In CuAg, prepared by TE, Ag layer is completely relaxed when its

- 625 thickness is equal to 2 nm. The mechanism involved in the stress relaxation has been demonstrated to be the same as the one occurring in thin film epitaxy. We have shown that there is a critical size for the Cu core, up to which the formation of core/shell arrangement is promoted. This configuration is
- <sup>630</sup> probably favoured by the underlying morphologies of the Cu cores which present icosahedral and decahedral configuration under a size smaller than 8 nm. Such shapes are mainly constituted of (111) faces implying that the growth rate of the Ag layer is the same on all these faces.
- <sup>635</sup> In CoPt, prepared by PLD, the morphology and the structure of the nanoparticles has been investigated as a function of the nature of the substrate and the laser frequency. The diffusion of Co and Pt atoms is favoured on a-carbon compared to amorphous alumina, or by using a lower
- <sup>640</sup> frequency during the deposition process. These phenomena result from the competition between the kinetic energy of the atoms, given by the substrate temperature, and the kinetic barriers, depending on the metal/substrate interaction, which drive the formation of 3D morphologies. The thermal
- <sup>645</sup> properties of the substrate can have a strong influence on the structure of the nanoparticles. Finally, we have shown that it is possible to synthesize as-grown ordered  $L1_0$  CoPt nanoparticles by reducing the growth kinetics (*i.e.* decreasing the laser frequency).

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#### Notes and references

1 S. Sao-Joao, S. Giorgio, J. M. Penisson, C. Chapon, S. Bourgeois, C. Henry, J. Phys. Chem. B 2005, 109, 342.

- 2 B. Skarman, T. Nakayama, D. Grandjean, R. E. Benfield, E.
   <sup>655</sup> Olsson, K. Niihara, L. R. Wallenberg, *Chem. Mater.* 2002, 14, 3686.
  - 3 N. Toshima, Pure and Applied Chemistry 2000, 72, 317.
  - 4 N. Toshima, M. Harada, Y. Yamazaki, K. Asakura, *J. Phys. Chem.* 1992, **96**, 9927.
- <sup>660</sup> 5 L. Favre, S. Stanescu, V. Dupuis, E. Bernstein, T. Epicier,
  P. Melinon, A. Perez, *Applied Surface Science* 2004, 226, 265.
  - 6 H. Zeng, J. Li, Z. L. Wang, J. P. Liu, S. Sun, *Nano Lett.* 2004, **4**, 187.
- 665 7 S. Basu, D. Chakravorty, Journal of Non-Crystalline Solids 2006, 352, 380.
- 8 M. Gaudry, E. Cottancin, M. Pellarin, J. Lerme, L. Arnaud, J. R. Huntzinger, J. L. Vialle, M. Broyer, J. L. Rousset, M. Treilleux, P. Melinon, *Physical Review B (Condensed Matter* 670 and Materials Physics) 2003, **67**, 155409.

9 J. Zhu, Y. Wang, L. Huang, Y. Lu, *Physics Letters A* 2004, **323**, 455.

10D. Garcia-Gutierrez, C. Gutierrez-Wing, M. Miki-Yoshida, M. Jose-Yacaman, *Applied Physics A: Materials Science &* 675 Processing 2004, **79**, 481. 11S. Sun, C. B. Murray, D. Weller, L. Folks, A. Moser, *Science* 2000, **287**, 1989.

12S. P. Withrow, C. W. White, J. D. Budai, L. A. Boatner, K. D. Sorge, J. R. Thompson, R. Kalyanaraman, *Journal of* 

- Magnetism and Magnetic Materials 2003, 260, 319.
  13M. Cazayous, C. Langlois, T. Oikawa, C. Ricolleau, A.
- Sacuto, *Physical Review B* 2006, **73**, 113402.
- 14T. Miyazaki, O. Kitakami, S. Okamoto, Y. Shimada, Z. Akase, Y. Murakami, D. Shindo, Y. K. Takahashi, K. Hono, <sup>685</sup> *Physical Review B* 2005, **72**, 144419.
- 15K. Sato, B. Bian, T. Hanada, Y. Hirotsu, *Scripta Materialia* 2001, **44**, 1389.
- 16L. Castaldi, K. Giannakopoulos, A. Travlos, D. Niarchos, *Applied Physics Letters* 2004, **85**, 2854.
- <sup>690</sup> 17D. B. Chrisey, G. K. Hubler, *Pulsed Laser Deposition of Thin Films*, John Wiley & Sons, London 1994.
- 18D. Alloyeau, C. Langlois, C. Ricolleau, Y. Le Bouar, A. Loiseau, *submitted to nanotechnology* 2007.
- 19T. W. Trelenberg, L. N. Dinh, B. C. Stuart, M. Balooch, 695 Applied Surf. Sci. 2004, **229**, 268.
- 20L. Vitos, A. V. Ruban, H. L. Skriver, J. Kollar, *Surf. Sci.* 1998, **411**, 186.
- 21J. P. Hirth, J. Lothe, *Theory of dislocations*, Krieger Publishing Company, Berlin 1982.
- 700 22A. E. Romanov, W. Pompe, S. Mathis, G. E. Beltz, J. S. Speck, *Journal of Applied Physics* 1999, 85, 182.
  - 23C. Mottet, J. Goniakowski, F. Baletto, R. Ferrando, G. Treglia, *Phase Transitions: A Multinational Journal* 2004, **77**, 101.
- <sup>705</sup> 24K. Koga, K. Sugawara, *Surface Science* 2003, **529**.
   25Y. Le Bouar, A. Loiseau, A. Finel, *Physical Review B* 2003, **68**, 224203.
  - 26D. J. Sellmyer, M. Yu, -. R.D. Kirby 1999, 1021., Nanostructured Materials, **12**, 1021.
- <sup>710</sup> 27M. Yu, Y. Liu, D. J. Sellmyer, J. Applied Physics 2000, 87, 6959.
  - 28G. Palasantzas, T. Vystavel, S. A. Koch, L. T. D. Hosson, *J. Applied Physics* 2006, **99**, 024307.
- 29I. M. Lifshitz, V. V. Slyozov, *Soviet Physic JETP* 1959, **35**, 715 331.
  - 30C. T. Campbell, Surface Science Reports 1997, 27, 1.
  - 31C. R. Henry, Surface Science Reports 1998, 31, 231.
  - 32D. R. Lide, *CRC Handbook of Chemistry and Physics*, CRC press, London 2002-2003.
- 720 33B. Behkam, Y. Yang, M. Asheghi, International Journal of Heat and Mass Transfer 2005, 48, 2023.
  - 34Y. H. Takahashi, T. Ohkubo, M. Ohnuma, K. Hono, *Journal of Applied Physics* 2003, **93**, 7166.
- 35B. F. Mc Ewen, M. Marko, J. of Histochem. and Cytochem. 725 2001, **49**, 553.

Table captions for one column table:

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735

2 nm of CoPt on Substrate	Mean size (nm)	Poly- dispersity	Particles density (10 <sup>2</sup> part.µm <sup>-2</sup> )	Coverage ratio
a-Carbon	6.3	45%	110	42%
a-Al <sub>2</sub> O <sub>3</sub>	7.9	30%	90	50%

Table. 1. Mean size, polydispersity, particles density and coverage ratio of the samples prepared at 600°C, with a laser frequency of 5 Hz on a-carbon and  $a-Al_2O_3$  substrates.

### Figures captions for one column figure

Figure 2 – (a) High resolution image of a CuAg nanoparticle of the 6@3 system, oriented along the [110] zone axis. Moiré
<sup>740</sup> patterns show that the silver shell caps completely the Cu core. (b) Power spectrum of the previous image that emphasize the epitaxial relationship between the core and the shell. The power spectrum shows a symmetry mirror (highlighted by a white line on the figure) coming from the
<sup>745</sup> twin boundary.

Figure 3 – (a) High resolution image of a CuAg nanoparticle with Ag located on the upper side and Cu on the lower side of the particle. The (220) interface plane is represented by a <sup>750</sup> dotted line. The misfit dislocations are localized by white circles along the interface plane. (b) Bragg-filtered image using the 220 reflection that allows to evidence the excess Cu planes of the perfect misfit dislocations. The inset represents the power spectrum of micrograph (a) on which the two 220 <sup>755</sup> reflections used for the filtering are circled.

Figure 5 – (a) Ag chemical map obtained by EFTEM and corresponding to the 2@1 system deposited with a substrate temperature of 400°C. (b) Colour-coded chemical map with 760 Ag (blue) and Cu (yellow) obtained by EFTEM and corresponding to the 2@1 system deposited with a substrate temperature of 500°C.

Figure 6 – (a) High resolution micrograph of a CuAg <sup>765</sup> core/shell decahedron with Moiré patterns on each of the five FCC sectors. (b) Power spectrum of micrograph (a).

Figure 7. CoPt FCC disordered nanoparticles synthesized by 770 PLD at 600°C with in the insets the size distribution of the particles. (a) **Substrate**: 10 nm thick amorphous carbon film on a commercial TEM grid. (b) **Substrate**: 2 nm thick amorphous alumina film deposited by PLD on the carbon film of a commercial TEM grid.

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Figure 8. CoPt nanoparticles, with a 2 nm nominal thickness, synthesized by PLD on a 10 nm thick amorphous alumina layer deposited by PLD on the carbon film of a commercial TEM grid heated at 600°C.

## Figures captions for double columns figure

Figure 1 – (a) BF conventional TEM image of an area on the <sup>785</sup> 6@3 sample, (b) Cu chemical map of the same area obtained by EFTEM, (c) Ag chemical map and (d) Colour-coded chemical map with Ag (blue) and Cu (yellow). Note the presence of pure Ag nanoparticles that have nucleated separately on the substrate.

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Figure 4 – Size distributions of the three different morphologies identified on the 2@1 sample. These distributions were obtained from conventional and high resolution TEM images. For each size range, a typical 795 HRTEM image of a nanoparticle is shown.

Figure 9. CoPt nanoparticles deposited by PLD with a laser frequency of 5 Hz (30 minutes) on a substrate heated at  $600^{\circ}$ C and then annealed in the deposition chamber at  $600^{\circ}$ C during

<sup>800</sup> 2 hours. In the inset: size distribution of the particles. (a) **Substrate**: 10 nm thick amorphous carbon film of a commercial TEM grid. (b) **Substrate**: 2nm thick amorphous alumina film of deposited by PLD on the carbon film of a commercial TEM grid.

<sup>805</sup> Electron diffraction patterns of CoPt nanoparticles grown (30 minutes) and annealed (120 minutes) at 600 °C, and their corresponding rotational averaging profiles. (c) Ordered  $L1_0$  nanoparticles on amorphous carbon film. (d) Disordered FCC nanoparticles on amorphous alumina film.

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Figure 10. Ordered L1<sub>0</sub> CoPt nanoparticles synthesized by PLD with a laser frequency of 1 Hz (2h30) on amorphous carbon film at a temperature of 600°C. (a) BF image of the nanoparticles. (b) HRTEM image of a CoPt L1<sub>0</sub> ordered
<sup>815</sup> nanoparticle oriented along the [110] direction, with in the inset the power spectrum on the image where the superstructure reflections are encircled. (c) TEM images of the selected area for the tomography experiment, (d) 3D reconstruction of the selected area obtain by electron
<sup>820</sup> tomography. The correspondence between 2D images of the particles and the 3D volumes is emphasized by the numbers on these two images.