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Particle formation at
a continental
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Particle formation at a continental background site: comparison of model results with observations

U. Uhrner¹, W. Birmili², F. Stratmann³, M. Wilck³, I. J. Ackermann⁴, and H. Berresheim⁵

¹Institute for Tropospheric Research, Permoserstr. 15, 04318 Leipzig, Germany

²Division of Environmental Health and Risk Management, Birmingham University, B15 2TT, UK

³Institute for Tropospheric Research, Permoserstr. 15, 04318 Leipzig, Germany

⁴Ford Research Centre Aachen, Süsterfeldstr. 200, 52072 Aachen, Germany

⁵German Weather Service – Deutscher Wetterdienst, Meteorological Observatory Hohenpeissenberg, Albin-Schwaiger-Weg 10, 83282 Hohenpeissenberg, Germany

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Correspondence to: U. Uhrner (uhrner@tropos.de)

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At Hohenpeissenberg (47°48' N, 11°07' E, 988 m asl), a rural site 200 - 300 m higher than the surrounding terrain, sulphuric acid concentrations, particle size distributions, and other trace gas concentrations were measured over a two and a half year period.

5 Measured particle number concentrations and inferred particle surface area concentrations were compared with box-model simulations based on a multimodal lognormal aerosol module that included a binary sulphuric acid water nucleation scheme. The calculated nucleation rates were corrected with a factor to match measured particle number concentrations. These corrections varied over a range of $10^{-3} - 10^{17}$. The correction factors were close to 1 for the measurements made in the winter, which represented stable thermal stratification and low wind conditions. In contrast, the correction factors were the largest for measurements made under strong convective conditions.

10 Our comparison of measured and simulated particle size distributions suggest a distant particle-formation process under convective conditions near the interface of the mixed layer and the entrainment zone, followed by downward transport and particle growth. For stable stratification and low winds, our comparisons suggest that particles formed close to the measurement site.

1. Introduction

20 Aerosols impact climate (e.g. Charlson and Heintzenberg, 1995) and human health (e.g. Dockery and Pope, 1994). New particle formation from gas-phase precursors, i.e. nucleation, is frequently observed in marine locations (e.g. Weber et al., 1999; O'Dowd et al., 2002) and in continental locations, such as forests (Mäkelä et al., 1997), remote (e.g. Weber et al., 1997) and polluted sites (e.g. Birmili and Wiedensohler, 2000; McMurry et al., 2000). However, the fundamental processes that cause nucleation and subsequent growth in the size-range of a few nanometers are still uncertain.

25 Sulphuric acid and water are believed to be the important constituents controlling

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atmospheric nucleation processes. Calculated nucleation rates using classical binary $\text{H}_2\text{SO}_4\text{-H}_2\text{O}$ theory substantially underestimate the observed nucleation rates for mid-latitude ambient conditions (e.g. [Weber et al., 1998](#); [O'Dowd et al., 1999](#)). Observed H_2SO_4 concentrations are about 1 order of magnitude too low to explain the ambient formation according to classical binary theory.

More complex nucleation mechanisms have been proposed, such as ion-mediated nucleation and growth ([Yu and Turco, 2000](#)) and the participation of a third molecule such as ammonia (NH_3) ([Coffman and Hegg, 1995](#); [Korhonen et al., 1999](#)), or the ubiquitous existence of thermodynamically stable 1–3 nm clusters probably formed by nucleation of H_2SO_4 , NH_3 , and H_2O ([Kulmala et al., 2000](#)). In comparison to binary systems, for ternary nucleation of H_2SO_4 , H_2O , and NH_3 , up to 2 orders of magnitude less gas-phase $\text{H}_2\text{SO}_{4(g)}$ is needed to achieve observed nucleation rates. However, the observed growth in the nanometer range can not be attributed to the subsequent condensation of H_2SO_4 , H_2O , and NH_3 . Therefore, to explain observed condensational growth rates, additional unknown species must participate in subsequent growth ([Kulmala et al., 2000](#)).

Other researchers noted that favourable atmospheric conditions, such as turbulence due to breaking Kelvin-Helmholtz waves ([Bigg, 1997](#)) or boundary layer mixing processes ([Easter and Peters, 1994](#); [Nilsson and Kulmala, 1998](#)), and atmospheric waves ([Nilsson et al., 2000](#)) can enhance nucleation rates by up to several orders of magnitude. [Weber et al. \(1999\)](#) suggested that nucleation mechanisms may vary with altitude. By comparing observed new-particle formation rates against modelled nucleation rates using classical binary nucleation, they concluded that at least for higher elevations in remote marine regions, new particles can be formed along cloud perimeters through $\text{H}_2\text{SO}_4\text{-H}_2\text{O}$ nucleation.

Even for the well characterized, binary $\text{H}_2\text{SO}_4\text{-H}_2\text{O}$ system, there are significant differences between theory and measurements. Potential weaknesses of current nucleation theories include the description of small liquid droplets in models by macroscopic thermodynamic properties, and condensational growth of cluster-sized droplets. [Mc-](#)

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Graw and Weber (1998) showed that models based on the liquid-drop approach over-predict the degree of H_2SO_4 hydration, which results in under-predicted nucleation rates. Moreover, thermodynamic properties used in nucleation models are not accurately known. Nucleation rates are extremely sensitive to small changes in $\text{H}_2\text{SO}_{4(g)}$ and water concentrations (e.g. Viisanen et al., 1997). In situ measurements of gas-phase $\text{H}_2\text{SO}_{4(g)}$ are difficult, require extensive know-how and instrumentation (Eisele and Tanner, 1993), and data are therefore rarely available. $\text{H}_2\text{SO}_{4(g)}$ concentrations have therefore been calculated in most studies (e.g. Andronache et al., 1997) by using a simple mass-balance approach considering particle formation via the reaction of hydroxyl radicals (OH) and sulphur dioxide (SO_2) and the loss onto preexisting aerosol particles. OH concentrations have been estimated or have been computed by using chemical models, such as used in the EURAD Model System (Haas, 1991; Ackermann et al., 1998), causing uncertainties in $\text{H}_2\text{SO}_{4(g)}$ concentrations and therefore uncertainties in calculated particle formation rates.

In this work, particle dynamics were simulated and compared to measured data that included recent new-particle formation data. Similar to many other studies, only ground-based measurements were available. For this work complete sets of ground-based measurements consisting of $\text{H}_2\text{SO}_{4(g)}$, RH, T, and size distributions were available as direct input for aerosol dynamics simulations for a H_2SO_4 - H_2O system. Instead of using computed $\text{H}_2\text{SO}_{4(g)}$ concentrations, measured time series of $\text{H}_2\text{SO}_{4(g)}$ were used as direct model input. Because we used a simple box-model approach, no spatial distributions of precursors or aerosol properties were needed to initialise the simulation. Uncertainties of $\text{H}_2\text{SO}_{4(g)}$ concentrations used in the model were reduced to measurement errors. Because an applicable and validated ternary nucleation scheme was not available, we used a binary nucleation model. The aim of this work was to investigate the differences between the measured aerosol properties and model predictions and to relate these differences to measured physical parameters such as the near-surface temperature gradient; wind speed and direction; trace gas concentrations; humidity and temperature profiles; and related processes such as vertical exchange and small-

scale turbulence and transport. Therefore, to avoid the strong effects of strong wind veering and inhomogeneities of the terrain on the aerosol dynamics processes, cases were selected to meet specific meteorological and aerosol dynamics criteria.

2. Field measurement data used for comparison

5 We give a brief description of the measured data that we used for comparison with the simulation results. More detailed information about the HAFEX (Hohenpeissenberg Aerosol Formation Experiment) measurements (1998–2000) can be found in [Birmili et al. \(2002\)](#) and in the case study by [Birmili et al. \(2000\)](#). For this comparison we used data from the Meteorological Observatory at Hohenpeissenberg (MOHp) located in
10 South Germany, which is run by the German Weather Service (DWD). Measurements were made on top of the Hohenpeissenberg (988 m), a single hill 200–300 m higher than the surrounding terrain. This remote site is located 60 km southwest of Munich and 30–40 km north of the Bavarian Alps.

Available long-term measurement data used in this work included dry, submicrometer particle size distributions ranging from 3–700 nm measured with a differential mobility particle sizer, gas-phase H_2SO_4 concentrations measured by using atmospheric-pressure chemical ionization mass spectrometry (AP/CIMS, see [Berresheim et al. \(2000\)](#) for details), NO_x , and routinely measured meteorological data. Ammonia measurements were not made. The time resolution was 15 min for the meteorological data and 5 min for the $\text{H}_2\text{SO}_{4(g)}$ data. Size distributions were measured every 15 min.
15 A CIMS measurement cycle consisted of $\text{H}_2\text{SO}_{4(g)}$ and OH concentration measurements. Therefore, there are gaps of several minutes in the data for $\text{H}_2\text{SO}_{4(g)}$, when OH concentrations were measured. Interpolated values for $\text{H}_2\text{SO}_{4(g)}$ were used for those periods. Hourly averaged NO_x data were used for interpretation of our results for
20 particle growth.

25 Local characteristics that must be considered in the interpretation of the data are inhomogeneities in heat flux, temperature, humidity, and local wind systems originat-

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ing from the position of the site. The site is affected by lee effects due to southerly airflows, in particular with foehn. Mountain winds can evolve particularly in summer due to intense irradiation onto the southern slopes of the Bavarian Alps, which creates a mesoscale circulation system that causes a compensating sinking air motion and northerly winds at the Hohenpeissenberg site (Fricke et al., 1997). Consequently, there is a local air mass change due to changes in the local flow pattern.

3. Modelling and comparison methodology

One of the goals of this study was to identify the atmospheric processes that control particle formation and growth. For the meteorological conditions corresponding to the measured data we used to compare with our simulation results, the following new-particle formation scenarios are possible:

- Particle nucleation occurs inside the boundary layer, where particle precursor concentrations are high. Nucleation is controlled by gas-phase chemical reactions (source of condensable material) and existing particles, which act as condensational sink (e.g. Pirjola et al., 1999).
- Particle nucleation occurs near inversion layers and new particles are mixed downward during the break-up of the inversion layer. Turbulent transport may control new-particle formation by (a) initiating nucleation by mixing air parcels with different chemical and thermodynamical properties (Nilsson and Kulmala, 1998; Jaenisch et al., 1998) and by (b) mixing newly formed particles down to the ground, where they can be detected.

Depending on the meteorological conditions, both scenarios are possible. Based on only point measurements of the particle size distribution and measurement of only a limited number of particle precursors, it is not possible to determine which of these scenarios dominated particle nucleation in the measured data sets. However, by combining measured size distribution and particle precursor data, known meteorological

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conditions at the time of such measurements, and by using a suitable aerosol dynamics model, it is possible to determine which of the particle nucleation scenarios was dominant. Therefore, based on the experimental data available from the HAFEX experiments, we used a box model to simulate particle nucleation and the development of the particle size distribution. The model is initialised by using measured particle size distributions at the start of each model run and the particle dynamics are driven by measured temperatures, relative humidities, and $\text{H}_2\text{SO}_{4(g)}$ concentrations. Simulating carefully selected particle nucleation events and relating the differences between measured and calculated particle size distributions to suitable meteorological parameters, the near-surface temperature gradient and the local Richardson number can be used to indicate where nucleation occurs.

Using a box model is sufficient because the data was measured at a single point and is therefore 0-dimensional. Furthermore, multidimensional simulations would involve unknown boundary conditions, such as the 3-D distribution of particle size distribution, gaseous precursor type, and emission rates, which would introduce an undesirable number of free parameters in the model.

3.1. Aerosol dynamics model

The simulations were made with a box model version of the lognormal MADMAcS model (Multicomponent Aerosol Dynamics Modal Approach System) (Wilck and Stratmann, 1997; Wilck, 1998). The model accounts for nucleation, condensation, and coagulation. Atmospheric mixing, sedimentation, and deposition were not considered. Nucleation rates were calculated by using the binary nucleation theory for the $\text{H}_2\text{SO}_4\text{-H}_2\text{O}$ system (Kulmala et al., 1998). The nucleation rate, J [$\#/(m^3s)$] can be expressed as

$$J = F_{\text{nuc}} \exp\left(-\frac{\Delta G^*}{k_B T}\right) \quad (1)$$

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where k_b is the Boltzmann constant, T is temperature, and ΔG^* is the energy required to form a cluster sufficiently large that it will not reevaporate (critical cluster). The kinetic pre-exponential factor, F_{nuc} , expresses the rate at which vapour is transported to critical clusters. A detailed description of the formulation of F_{nuc} and ΔG^* is given in Kulmala et al. (1998). The critical cluster composition according to Wilemski (1984) is solved numerically and the hydrate interaction (Jaeger-Voirol et al., 1987) is taken into account. The water content of the clusters and particles is computed from an equilibrium relationship with the relative humidity. Prognostic variables are “dry” particle size distribution moments, $M_{k,j}$, representing particle number concentration ($k = 0$), surface area ($k = 2$), and mass ($k = 3$) for each mode j . $M_{k,j}$ can be expressed as

$$M_{k,j}(\mathbf{x}, t) = \int_0^{\infty} \tilde{m}_{p,j}^k f_{lnj} d\tilde{m}_{p,j} \quad (2)$$

$$f_{lnj} = \frac{1}{\sqrt{2\pi} \tilde{m}_{p,j} \ln \sigma_{g,j}} \exp\left(-\frac{\ln^2(\tilde{m}_{p,j}/\tilde{m}_{gN,j})}{2 \ln^2 \sigma_{g,j}}\right) \quad (3)$$

where f_{lnj} represents the normalized lognormal frequency function of mode j , \tilde{m}_p is the “dry” particle mass, \tilde{m}_{gN} is the “dry” geometric mean particle mass, and σ_g is the geometric standard deviation. The time evolution of the particle size distribution can be described by the time evolution of the moments of the “dry” distribution as

$$\begin{aligned} \frac{\partial M_{k,j}}{\partial t} &= \delta_{j,nuc,1} (\tilde{m}_p^*)^k J \\ &+ k N_j \int_0^{\infty} \tilde{m}_p^{k-1} [C(m_{p,j}) - E(m_{p,j})] f_j^{ln}(\tilde{m}_p) d\tilde{m}_p \\ &- \sum_{i=1}^{n_m} N_j N_i \int_0^{\infty} \int_0^{\infty} \tilde{m}_p^k \beta(m_p, m'_p) f_j^{ln}(\tilde{m}_p) f_i^{ln}(\tilde{m}'_p) d\tilde{m}'_p d\tilde{m}_p \end{aligned} \quad (4)$$

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$$+ \sum_{i,v=1}^{n_m} C_{ivj} N_i N_v$$

$$\cdot \int_0^{\infty} \int_0^{\infty} (\tilde{m}_p + \tilde{m}'_p)^k \beta(m'_p, m_p) f_i^{\text{In}}(\tilde{m}_p) f_v^{\text{In}}(\tilde{m}'_p) d\tilde{m}'_p d\tilde{m}_p$$

where N_j represents the total number concentration in mode j , $C(m_p)$ and $E(m_p)$ represent the condensation and evaporation rates, respectively, and $\beta(m_p, m'_p)$ represents the coagulation coefficient for particles of mass m_p and m'_p . d_{gN} represents the geometric mean diameter of the number-weighted distribution. The tensor C_{ivj} expresses the coagulation convention for the coagulation gain term as suggested by Whitby and McMurry (1997). $C_{ivj}=1/2$ if $i=v=j$ (intramodal coagulation), $C_{ivj}=1$ if $i=j$ and $d_{gNi} > d_{gNv}$, $C_{ivj}=1$ if $v=j$ and $d_{gNi} \leq d_{gNv}$, and $C_{ivj}=0$ otherwise. The integrals must be evaluated numerically because they involve a nonlinear dependence of the “wet” particle mass m_p on the “dry” particle mass \tilde{m}_p . The numerical integrals were evaluated with an 8-point Gauss-Hermite quadrature technique.

Up to four modes were used. The model was initialised by using modal parameters for particle number concentration N , geometric mean diameter d_{gN} , and standard deviation σ for each mode; $\text{H}_2\text{SO}_{4(\text{g})}$, RH, and temperature. The measured time series of $\text{H}_2\text{SO}_{4(\text{g})}$, RH, and temperature were used as the thermodynamic forcing parameters.

3.2. Data selection

From the experimental data available from HAFEX for 1998–2000, data sets were selected that matched the following criteria:

- Data sets with strong veering of the wind prior to and during nucleation were rejected, because local wind systems such as upslope flows or mountain wind systems involve varying transport patterns and air mass changes. Therefore, data sets were chosen with either constant wind direction at wind speeds > 4 m/s or data sets with moderate changes ($< 90^\circ$) in wind direction for wind speeds

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< 4 m/s prior to and during particle nucleation.

- Because the site is in close proximity to the Alps, which reach an elevation of 2961 m to the south, data sets influenced by foehn conditions were rejected, such as data sets with strong winds from the southeast–southwest sector, low relative humidities, and relatively high temperatures.
- Smooth time evolution of d_{gN} for the nucleation mode (consisting of the smallest particles in the size distribution, ranging from 3–11 nm).
- A significant particle number concentration of particles smaller than 11 nm.

Therefore, only data sets classified as Type “I” or “II” (strong or medium events) in the article by [Birmili et al. \(2002\)](#) were chosen.

From 46 data sets classified as Type “I” or “II”, 12 data sets met these criteria, and these are summarized in Table 1. A further limiting criterion was the simultaneous availability of T, RH, and $H_2SO_{4(g)}$ measurements in the period of interest.

3.3. Comparison of measured and simulated particle size distributions

The simulated and measured particle size distributions were compared. To avoid artifacts due to the data-reduction procedure (e.g. fitting lognormal distributions to measured size distributions) and to reduce the influence of model assumptions such as the multi-modal lognormal size distributions used in MADMAcS, the integral moments of the measured and simulated particle size distributions were compared. The moments we compared were total number and surface area concentrations. Because the simulated particle size distributions start at the size of the critical cluster size of about 1 nm and because the measured particle size distributions start at the lower detection limit of the measurement instruments which was about 3 nm, the simulated particle size

distribution moments for the nucleation mode were integrated from 3 nm upwards:

$$M_{k,1} = \int_{3nm}^{\infty} \tilde{d}_p^k n_1(\tilde{d}_p) d\tilde{d}_p \quad (5)$$

where n_1 represents the lognormal size distribution of mode $j = 1$ and \tilde{d}_p represents the dry particle diameter.

This integration limit affects only the zeroth and first moment when nucleation occurs, and the impact on higher moments is negligible. The zeroth moment that represents the corrected particle number concentration is therefore represented by $N_{>3nm}$. To determine from the experimental data which of the two nucleation scenarios were active, the difference between the measured and modelled evolution of the size distribution moments was evaluated. As an indicator for these differences related to particle number concentration, a linear correction factor for the nucleation rate, cf , is defined as

$$J_c = cf \cdot J$$

cf is the only free parameter in the model and was determined iteratively by varying cf until the measured and simulated peak number concentrations matched within $\pm 25\%$. The result of this procedure is shown in Fig. 1 for data taken on 19 May 1998 and 26 January 2000. Figure 1 shows measured and simulated total particle number concentration as a function of time for two values of cf for each data set.

Figure 1 indicates that different values of cf were required to match particle number concentration for each of the data sets. For 19 May 1998 $cf = 10^5$, and for 26 January 2000 $cf = 0.2$.

The data sets we used for comparison exclude strong effects of local air masses affected by veering winds. However, days with clear-sky radiation exhibited a smooth, continuous evolution in temperature, humidity, particle number concentration, particle diameter, and other parameters, indicating that they were affected by vertical exchange and turbulent transport processes. These processes, which could not be accounted for in our box-model, may increase nucleation rates (see Easter and Peters, 1994; Nilsson

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and Kulmala 1998), requiring a high value of cf to achieve acceptable comparison between measured and simulated particle concentrations.

This suggests the use of regression analysis to relate the correction factor to meteorological properties such as the near-surface temperature gradient and the bulk Richardson number, Ri_B . These parameters are related to the meteorological conditions controlling new-particle formation and may indicate which nucleation scenario was active.

The model was initialised 2–4 h before the observed increase in particle number concentration, at times where fluctuations in particle surface area concentration were less pronounced and therefore represented the background aerosol before nucleation occurred¹. For data sets that displayed strong fluctuations in particle surface area concentration the model was initialised at times where the particle surface area concentration was at approximately the mean value (e.g. see Fig. 2p for 19 May 1998).

The input for the aerosol dynamics model were determined from the measured particle size distributions by using a least-squares fit for 2–4 modes, and yielded lognormal fit parameters N , d_{gn} and σ . For details on this fitting procedure, see Birmili et al. (2001).

4. Results

To illustrate the different types of aerosol dynamics observed during HAFEX, four data sets with different characteristic temperature, humidity, $H_2SO_{4(g)}$, particle number concentration, and particle surface area concentration were chosen. Then differences between measured and simulated results were related by using regression analysis to additional meteorological information, such as near-surface temperature gradients and

¹20 April 1998 featured a preceding ancillary maxima ($9000 \#/(cm^3)$). Therefore, the model was initialised after the first particle number increase at a value corresponding to the minimum in particle number concentration 1–1 1/2 hours before the second increase

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wind speed.

4.1. Data sets selected for detailed comparison:

Of the four data sets we selected for case studies, two were in winter and two were in spring. The measured time series of T and RH for these data sets are shown in Figs. 2a–d, for $\text{H}_2\text{SO}_4(\text{g})$ in Fig. 2e–h. In the lower panels of Fig. 2 the measured number concentration (Figs. 2i–l) and particle surface area concentration (Fig. 2m–p) are shown for 26 December 1998, 26 January 2000, 20 April 1998, and 19 May 1998. *26 December 1998 and 26 January 2000:* Both winter cases are characterized by low temperatures, stable stratification, relatively high relative humidities, and southwesterly winds prior to nucleation. On 26 December 1998 there were occasional clouds, sustained winds from the southwest, and a near-surface temperature gradient of about 0.3 Km^{-1} prior to increasing particle number concentration and about 0.7 Km^{-1} in the afternoon. 26 January 2000 had low, but variable local winds during and after particle nucleation, a near-surface temperature gradient of about $0.6 - 1.2 \text{ Km}^{-1}$ prior to the main particle number increase, and neutral shortly after the first maximum occurred in $N_{>3\text{nm}}$. In particular, for winter cases, the measured $N_{>3\text{nm}}$ resembles the evolution of $\text{H}_2\text{SO}_4(\text{g})$ with a time lag of about 1–2 h. *20 April 1998 and 19 May 1998:* For 20 April 1998 typical diurnal cycles for T and RH for a day with strong solar radiation were measured. Winds were low and predominantly from southerly directions. On this day two different maxima in $N_{>3\text{nm}}$ were observed, a relatively weak maximum at 08:00 and the main peak between 12:00–13:00. On 19 May 1998 there were scattered clouds and highly variable H_2SO_4 concentrations. However, T and RH displayed values typical of a day with strong solar radiation. The winds were low and variable. Both spring days were characterized by convective conditions. On both days there was a strong, negative near-surface temperature gradient up to -4 Km^{-1} at the measurement site, and the profiles for temperature and specific humidity indicated well-mixed boundary layers reaching up to 1700 m (see Fig. 4).

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On 19 May and 26 January the increases in $\text{H}_2\text{SO}_{4(g)}$ and $N_{>3\text{nm}}$ were rapid, in contrast to 20 April and 26 December for which $\text{H}_2\text{SO}_{4(g)}$ and $N_{>3\text{nm}}$ showed smooth changes and one major increase in particle number concentration occurred (Figs. 2e–l).

Figure 2 indicates that an increase of $N_{>3\text{nm}}$ is accompanied by a drop in relative humidity, especially preceding an increase in $\text{H}_2\text{SO}_{4(g)}$ concentration and preceding a minimum in particle surface area concentration. The time series for surface temperature are dominated by the diurnal cycle and all four data sets exhibit only minor changes prior to the observed nucleation burst.

The two lower panels of Fig. 2 show simulated results (dashed lines) as well as measured results (solid lines) for particle number concentration (Figs. 2i–l) and particle surface area concentration (Figs. 2m–p). The increase in $N_{>3\text{nm}}$ is reproduced well for 26 December 1998, which is characterized by $N_{>3\text{nm}}$ increasing continuously over a two-hour period (Fig. 2i). The onset of nucleation was well predicted.

For 26 January 2000, the simulation reproduces the distinct increases of the measured $N_{>3\text{nm}}$ very well (Fig. 2j). For 20 April and 19 May 1998, the simulated increase of $N_{>3\text{nm}}$ is earlier and steeper than the measured increase of $N_{>3\text{nm}}$ (Figs. 2k and l). Although the simulations seem to reproduce the measured data for both spring data sets, the different structures of $N_{>3\text{nm}}$ indicate that for 19 May (Fig. 2l) the simulated and observed maxima are out of phase. The simulated burst in $N_{>3\text{nm}}$ coincides with the steep increase in $\text{H}_2\text{SO}_{4(g)}$ (Fig. 2h), whereas the observed burst in $N_{>3\text{nm}}$ lags half an hour behind the simulated burst. For 20 April, the simulated increase in $N_{>3\text{nm}}$ (Fig. 2k) coincides with the small hump in RH at 10:00 (Fig. 2c) and the earlier increase in $\text{H}_2\text{SO}_{4(g)}$ (Fig. 2g), whereas the observed increase evolves over a longer time period coinciding with falling humidity.

To achieve agreement between measured and modelled particle number concentration, the nucleation rate was decreased for 26 January by a factor of 5, and increased by 10^4 for 26 December, 10^{13} for 20 April, and 10^5 for 19 May.

After 1 to 2 h of simulated increasing particle number concentration, $N_{>3\text{nm}}$ de-

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creases due to coagulation. On average, for all 12 cases, the simulated decrease in particle number concentration diverges about 20% 3 h after the maxima and about 40% 6 h after the maxima. We attribute these differences to neglected aerosol sinks associated with mixing, such as deposition.

5 The simulation results for particle surface area concentration are shown in Figs. 2m–p. For 26 December (Fig. 2m), 20 April (Fig. 2o), and 19 May (Fig. 2p), the model accurately reproduces the overall tendency, but the simulated particle surface area concentration is too low for 19 May. 26 January 2000 (Fig. 2n), shows large fluctuations and the measured and simulated results do not agree, probably because of the
10 effect of advection and mixing of polluted air, which is not represented in the model. Unfortunately, trace-gas measurements are missing for the time period of 13:00–18:00. However, strongly increasing NO concentrations around noon and enhanced levels of SO₂ and NO₂ after 18:00 indicate the impact of polluted air.

15 The evolution of simulated particle surface area concentration is smooth in all four cases. The measurements for particle surface area concentration show distinct fluctuations. The fluctuations in particle surface area can be caused by advection and mixing of either polluted or clean air.

4.2. Micrometeorological influence

20 The simulations for 20 April 1998 and for 19 May 1998 are characterized by simulated N_{>3nm} profiles that are steeper than the measured profiles and where the simulated onset of increasing N_{>3nm} occurs before the measured increase of N_{>3nm}. Dew-point temperature measurements taken at 1 min intervals were available for both days. As shown in Fig. 3 these data provide further indications of micrometeorological processes prior to the nucleation bursts. Both dew-point temperature time series show
25 strong fluctuations. Strong vertical exchange processes can be inferred for both cases because for 19 May 1998 the wind speed was lower than 1 m/s from 10:00 to 12:00 and for 20 April the wind speed was about 2 m/s from 10:00 to 13:00. In both cases the main increase in observed N_{>3nm} was accompanied by a significant reduction

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in dew-point temperature. This might indicate that drier air from aloft was entrained into the boundary layer and mixed downwards by so-called top-down diffusion. Specific humidity taken from radiosonde profiles can be used as an indicator for vertical exchange if there is a significant difference between the surface layer and the entrainment layer. Unfortunately, the closest radiosonde station was located 70 km northeast of the site at Munich-Oberschleissheim. However, for some days there were well-developed boundary layers several hundred meters higher than the MOHp site, and the existence of weak zonal flows allows us to rule out strong orographic effects, so that we can use the distant data 70 km away to interpret the measurements at the MOHp site. Figure 4 shows the temperature and specific humidity profiles taken at Munich-Oberschleissheim. The elevation of the MOHp site is indicated by the horizontal line at 988 m. The temperature lapse rate was nearly adiabatic for both cases, indicating that well-mixed boundary layers developed, reaching up to ≈ 1600 m for 19 May and up to ≈ 1700 m for 20 April. Assuming horizontally homogenous conditions and a shallow superadiabatic surface layer at the MOHp site, the temperature difference between the measurement site and the top of the boundary layer would have been between 7 and 8°C.

The dew-point temperature time series for May 19, 1998 indicates that marked variations occurred at frequencies from 30 to 60 minutes. These variations might indicate the existence of turbulent eddies and plumes. $\text{H}_2\text{SO}_{4(g)}$ fluctuations seemed to correspond to fluctuations of dew-point temperatures. In contrast, the measured N lagged about 1 to 1.5 hours behind the measured $\text{H}_2\text{SO}_{4(g)}$. For April 20, 1998, until 10:30 $\text{H}_2\text{SO}_{4(g)}$ fluctuations seemed to correspond to fluctuations of dew-point temperature. The time lag between $\text{N}_{>3\text{nm}}$ and $\text{H}_2\text{SO}_{4(g)}$ was about 2 hours in the morning, but thereafter it is difficult to associate N to $\text{H}_2\text{SO}_{4(g)}$. However, the variations in dew-point temperatures might indicate the existence of turbulent eddies and plumes that occurred at shorter cycles than the time-lag between N and $\text{H}_2\text{SO}_{4(g)}$.

Both measurements suggest that $\text{H}_2\text{SO}_{4(g)}$ and other condensible species and precursors, such as SO_2 , were first mixed upwards by bottom-up diffusion from the lower

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layers, and then particles formed higher up in the vicinity of the entrainment layer and at the top of the mixed layer. The reasons for this are probably that lower temperatures and entrainment of air with lower preexisting particle surface area concentration enhanced the formation of new particles. Then particles were mixed downward in the reverse top-down diffusion. Because the time lag between N and $H_2SO_{4(g)}$ was longer than the inferred cycles for turbulent eddies and plumes, it is possible that the detection of freshly-formed particles occurred in one of the mixing cycles after particles had already grown to a large size. However, neither turbulent particle flux measurements nor turbulent heat fluxes are available to validate this hypothesis.

These two spring data sets provide indications of the impact of vertical exchange processes on the differing model responses and of the complexity of interacting processes.

4.3. Comparisons of measured and simulated results

In the previous sections we compared aerosol dynamics simulations to four data sets that represent distinct atmospheric conditions. In this section we use regression analysis to compare simulation and measurement results for all 12 data sets, and discuss the origin of the nucleation-rate correction, cf , for all simulations. To assess the degree to which vertical exchange processes can explain the differences between the measured and simulated particle size distributions, the best available indicator of vertical exchange processes is the local temperature gradient. Temperature measurements were made at 0.05 and 2 m. Due to the complex terrain in the vicinity of the measurements, calculating the temperature gradient from radio-sonde data taken 70 km away may not represent local conditions. In particular, for winter days characterized by stable stratification, temperature profiles from Munich-Oberschleissheim do not represent the local conditions at MOHp. Because it is difficult to determine from measurements the time interval for the onset of particle formation, the values for the near-surface temperature gradient, wind speed and direction, and particle surface area concentration were taken at the onset of the simulated increase of $N_{>3nm}$ and are shown in Table 1.

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For all 12 cases the gradient between the surface and ground temperatures was related to the logarithm of cf as an indicator of the difference between measured and simulated results (Fig. 5). The logarithm of cf varied from -3 to 17.2 .

Figure 5 shows a distinct relation between cf and the near-surface temperature gradient. For the 12 cases we considered, a strong correlation was obtained. With increasing atmospheric instability (increasing negative temperature gradient) and related increased forcing of buoyancy-driven mixing processes, cf increased. For neutral conditions and stable stratification (positive temperature gradient) cf decreased with increasing atmospheric stability. In the stable regime the effect of shear-driven turbulence on atmospheric mixing decreases with increasing atmospheric stability. The interaction between wind shear and thermal stratification is described below.

For five out of six winter data sets, our box model reproduced the onset of increasing $N_{>3\text{nm}}$ within a few minutes. In contrast, on average the spring data sets show about a 1-h premature onset of increasing $N_{>3\text{nm}}$. This bias confirms that for convective conditions important time and length scales are omitted in the model, which strongly affect nucleation.

By analysing the atmospheric conditions that correspond to each of the data sets, the comparison between the measured and simulated evolution of $N_{>3\text{nm}}$ is the best for data sets corresponding to stable atmospheric conditions, and degrades for the data sets corresponding to progressively more unstable atmospheric conditions. For data sets corresponding to unstable atmospheric conditions, the simulated profile of $N_{>3\text{nm}}$ is typically too steep and the onset of the sharp increase of $N_{>3\text{nm}}$ occurs too soon.

Wind shear is another important atmospheric mixing process that affects particle nucleation and could explain part of the difference between our measured and simulated results. In particular, for stable, stratified atmospheric conditions, vertical mixing is driven by wind shear. If wind speed is related to cf , this would indicate that wind shear strongly affects particle nucleation. However, no correlation between wind speed and cf was found, which could mean that buoyancy-driven turbulence is the dominant process.

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Ri relates the effect of buoyancy-driven turbulence to shear-driven turbulence. To calculate Ri_B , only two temperature measurements at different heights and one wind-speed measurement (with the boundary condition $U = 0$ for $z = 0$) are required. Although the wind-speed measurements should be made at the same height as the temperature measurements, at MOHp temperatures were only measured at 0.05 and 2 m, and the wind speed (U) was measured at 40 m. We therefore calculated a bulk value for an auxiliary \tilde{Ri}_B as

$$\tilde{Ri}_B = \frac{g \Delta\theta / \Delta z_1}{\theta \Delta U / \Delta z_2}. \quad (6)$$

In Eq. (6) we set $\Delta z_1 = 2$ m and $\Delta z_2 = 40$ m. The resulting values for \tilde{Ri}_B are shown in Fig. 6.

The scatterplot shown in Fig. 6 indicates an inverse relation between cf and \tilde{Ri}_B for wind speeds higher than 3 m/s. For increasing \tilde{Ri}_B (i.e. increasing atmospheric stability) cf decreased (i.e. lower cf). These results may indicate that low wind speeds and stable atmospheric conditions produce large, positive \tilde{Ri}_B . In this regime wind shear is small and turbulence is suppressed by the stable atmospheric conditions, yielding laminar flow conditions. With increasing wind speed, and therefore increasing wind shear, laminar flow shifts to turbulent flow when Ri decreases to about 0.25 (Stull, 1988). Because we calculated the temperature gradient just for a height interval of 2 m, the critical value for \tilde{Ri}_B should be greater than 0.25. Two of the data sets that we considered may fall within this stable regime, and although four data sets had positive temperature gradients, which should yield thermally stable atmospheric conditions ($1.2 > \tilde{Ri}_B > 0.15$ indicated by stars), according to the criteria defined by Eq. (6), these four data sets are classified as unstable and may indicate some influence of wind shear in this regime. With decreasing \tilde{Ri}_B , for ($\tilde{Ri}_B < 0$) buoyancy-driven turbulence dominates the turbulent exchange processes, so that there is only 1 data set that was noticeably affected by wind shear. There are also two data sets that are characterized by calm winds and strong negative $\partial T / \partial z$. For one of these data sets $cf = 10^5$ and $Ri = -853$, which

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lay outside the range of Fig. 6, and therefore are not shown. The results shown in the scatterplot in Fig. 6 indicate the complexity of interacting atmospheric processes and the different flow regimes under which the measurements were made.

Similar to the scatterplots shown in Fig. 5 and 6, Fig. 7 shows the relationship between cf and particle surface area concentration. Figure 7 indicates a weak relation between cf and particle surface area concentration. For entrainment of aerosol with low preexisting particle surface area concentration, subsequent mixing and dilution could yield only a weak signal in the time series measured at ground level.

Regression analysis confirmed that the evolution of particle number concentration was significantly affected by micrometeorological processes, such as strong convection on some days. The main problem with this analysis is that the atmospheric conditions affecting particle formation might have been very different from the atmospheric conditions at the point of measurements. We therefore look in the data sets for indicators of mixing from layers above the measurement points down into the layers where measurements were made. The near-surface temperature gradient and \bar{Ri}_B were good indicators of the impact of atmospheric mixing on particle nucleation. Near the top of the boundary layer the temperature, specific humidity, and particle size distributions might be substantially different from those closer to the ground, and could favour enhanced particle formation rates. Even for stable atmospheric conditions, strong atmospheric inversions below the site could also increase particle formation rates. Particles may form at the discontinuity under the inversion, below the MOHp, and could be transported upwards by wind shear and by orographic lift.

4.4. Particle growth

The average of the difference between the measured and simulated particle surface area concentration was calculated for all 12 cases by integrating the difference between the measured and simulated particle surface area concentration over the six-hour measurement period, beginning with the simulated increase of particle number concentration. Similar to Fig. 5, Fig. 8 shows 6-h mean differences for particle surface

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area concentrations. The differences shown in Fig. 8 are normalized by the measured particle surface area concentration. We define this normalized difference between the measured and simulated particle surface area concentration as ΔS .

For particle nucleation, cf was used to determine the difference between particle nucleation and the evolution of particle number concentration, and could be correlated to the vertical exchange rate. However, for the behaviour of the particle surface area concentration, there is no identifiable relation between ΔS and the vertical exchange processes. This suggests that particle nucleation and growth are controlled by different processes. The processes most likely controlling ΔS are:

- Advection of polluted air
- Condensation and evaporation of unidentified species. Likely species are ammonia, nitric acid (HNO_3), oxygenated biogenic hydrocarbons, and other unidentified, condensable, organic species

Because these processes may occur simultaneously, it is difficult to determine their relative contributions to ΔS .

Hourly NO_2 and NO measurements were available for ten out of the twelve selected cases for the periods of interest. We used NO_x measurements as an indicator for the degree of pollution of the air.

The data sets for measured particle surface area concentration for 20 April (Fig. 2o), 19 May (Fig. 2p), and 26 December 1998 (Fig. 2m) show low values of ΔS , and were characterized by relatively low NO_x concentrations, ranging from 0.5 to 2 ppbv. For 26 January 2000 NO_x measurements are missing for the period of the most interest (see Fig. 9). However the remaining measurements for the late evening indicate high NO_x levels.

Two other data sets with large ΔS are also shown in Fig. 9. For both cases enhanced NO_x concentrations and enhanced particle surface area concentrations were measured. Moreover, the time series for NO_x indicates a time evolution similar to the

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time series for measured particle surface area concentration. Using NO_x as an indicator of combustion sources, this suggests that the site was affected by emissions, most likely from nearby roads and heating sources below the Hohenpeissenberg measurement site.

5. Conclusions

Selected data sets from a long-term set of atmospheric field measurements of particle size distributions, sulphuric acid concentrations, and meteorological parameters were used for comparisons with box model simulations of atmospheric particle nucleation and growth. The primary objective of this work was to study differences between the measured and simulated particle size distributions, and to identify atmospheric processes responsible for these differences. The differences for particle number concentration were evaluated by using a constant correction factor for the nucleation rate expression, which was the only free parameter used in the model. The benefit of the chosen approach was that we are able to compare a closed set of measurements against model results, and assumptions related to boundary conditions, background concentrations, particle composition, etc. were avoided.

Our results indicate that the measured and simulated particle number concentration compared well for data sets where the atmosphere was stably stratified and when the wind speed was relatively low. For wind speeds greater than 6 m/s and for stable stratification, the comparison degrades. Larger differences in particle number concentration were found for days with convective conditions, as indicated by negative near-surface temperature gradients and radio sonde profiles. For particle number concentration profiles where the simulated particle number concentration either rose faster than the measured increase, or where the onset of a sharp rise in particle number concentration occurred before the measured onset, the cause could be related to buoyancy-driven turbulent exchange processes. This indicates that under convective conditions the initial particle nucleation process occurs higher up in the atmosphere, where more

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favourable conditions occur followed by downward mixing and growth to detectable sizes. Therefore, a significant part of these differences and their variability is attributed to distant formation of particles and micrometeorological processes that cause them to be transported to the ground-based measurement site. Our results suggest that buoyancy-driven turbulence and wind shear are the micrometeorological processes accounting for such transport.

In contrast to particle formation, micrometeorological processes are not sufficient to explain the differences in simulated particle surface area concentration. This indicates that particle surface area concentration is affected by condensation and evaporation processes of additional species, which are strongly affected by transport of emissions, such as NO_x . A major element missing from the data sets to which we compared our simulation results is the lack of turbulent flux measurements or profile measurements, which could provide better insight into vertical exchange processes.

Our comparison of measured and simulated particle size distribution parameters indicate the complexity of various interacting processes, such as micrometeorology, particle nucleation, growth, coagulation, transport, and deposition. To further study the effect of transport processes on particle nucleation and growth, measurements up to the entrainment layer are needed.

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Table 1. Measured parameters for the 12 data sets used in this work. The values were taken at the begin of increasing simulated $N_{>3\text{nm}}$.

date	$\partial T/\partial z$ Km^{-1}	$T_{2\text{m}}$ T in $^{\circ}\text{C}$	surface area $\text{nm}^2\text{cm}^{-3}$	wind direction	$ U $ ms^{-1}
010498	-0.7	10.9	1.2e8	W	8.3
200498	-2.8	6.1	2.3e7	SW	2.1
150598	-3.6	13.6	5.2e7	ONO	7.1
160598	-4.4	11.6	3.1e7	NO	6.2
190598	-2.5	10.7	3.6e7	WSW	0.4
231298	1.7	-3.0	1.3e7	SSO	2.7
261298	0.3	4.1	6.0e6	WSW	5.3
030199	0.5	2.2	7.0e6	WSW	10.0
140499	-2.5	6.0	6.3e6	WSW	6.0
130100	0.6	-3.9	5.0e6	WSW	5.5
260100	1.0	-7.9	1.4e7	SW	2.0
270100	0.8	-3.3	7.1e6	WSW	7.3

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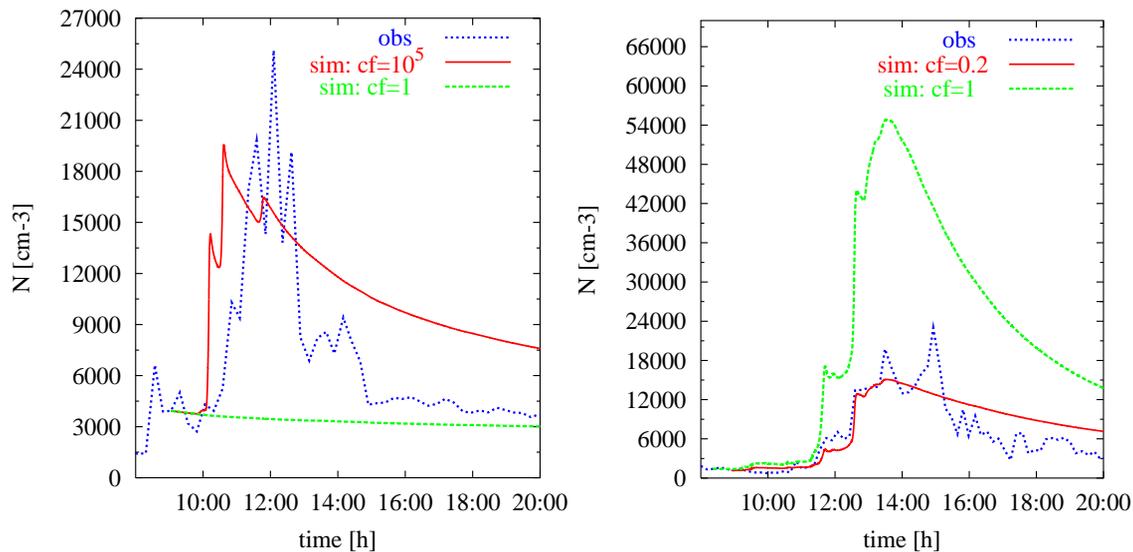


Fig. 1. Effect of cf on the simulated particle formation rates for 19 May 1998 (left, $cf = 10^5$) and 26 January 2000 (right $cf = 0.2$).

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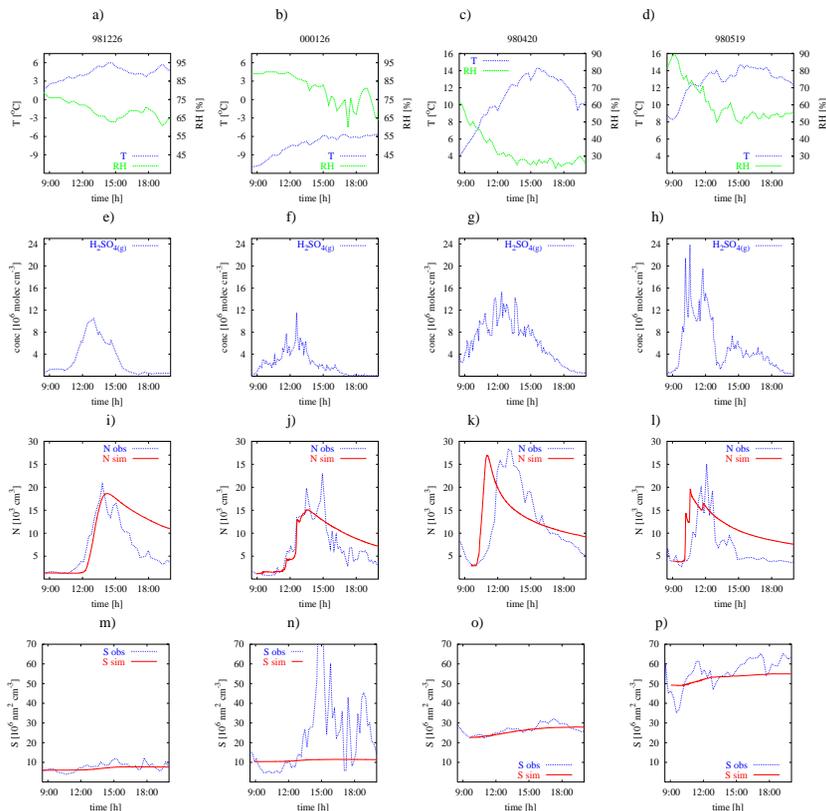


Fig. 2. Data sets (from left to right) for 26 December 1998, 26 January 2000, 20 April 1998, and 19 May 1998 (each panel). The panels show measured surface temperature and relative humidity (Figs. a–d), measured $\text{H}_2\text{SO}_4(\text{g})$ (e–h), measured and simulated time evolution of $\text{N}_{>3\text{nm}}$ (i–l), and measured and simulated time evolution of particle surface area concentration (m–p).

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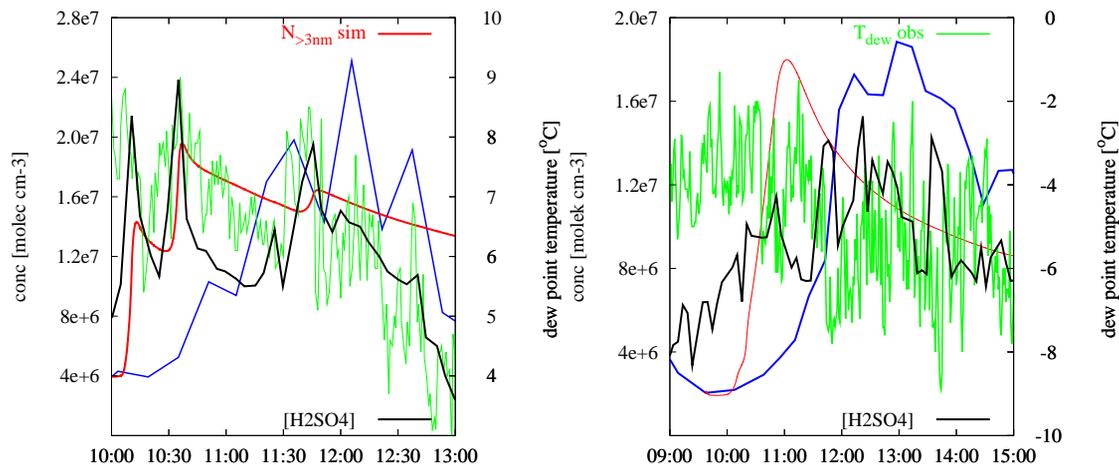


Fig. 3. Dew-point temperature taken at 1-min intervals and $H_2SO_{4(g)}$ concentrations measured at MOHp, 19 May 1998 (left) and 20 April 1998 (right). N_{obs} is indicated in green and $N_{>3nm}$ is indicated in red.

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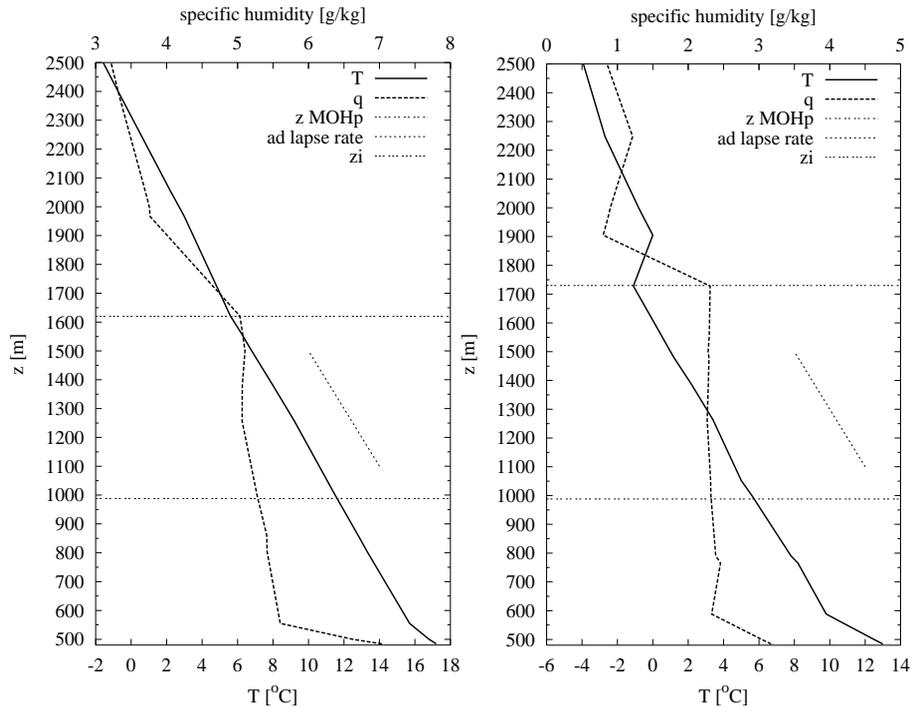


Fig. 4. Vertical profiles for specific humidity and temperature, Munich-Oberschleissheim, 19 May 1998 (left) and 20 April 1998 (right).

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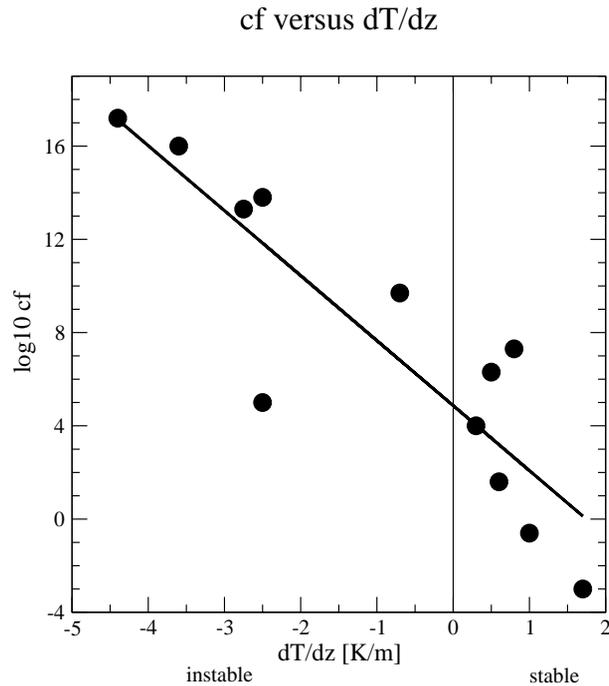


Fig. 5. \log_{10} correction factor (cf) versus $\partial T/\partial z$.

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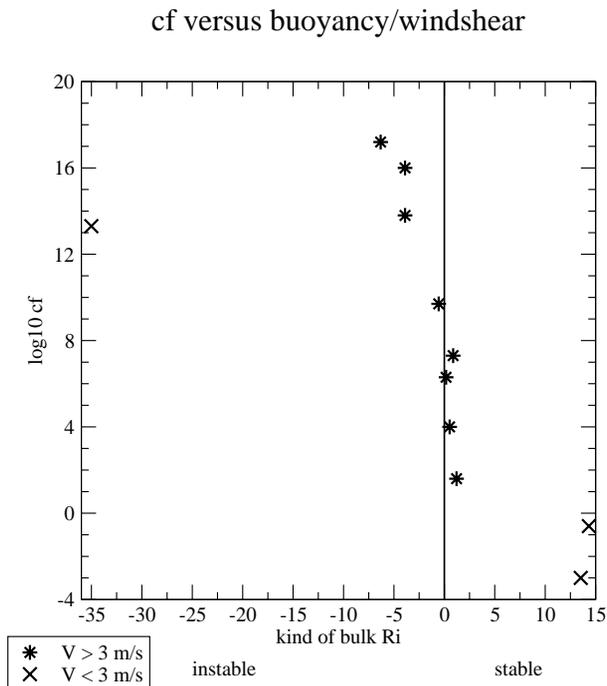


Fig. 6. Nucleation rate correction, cf , versus auxiliary bulk Richardson number, \tilde{Ri}_B . $\partial T/\partial z$ and $\partial U/\partial z$ are taken at different altitudes.

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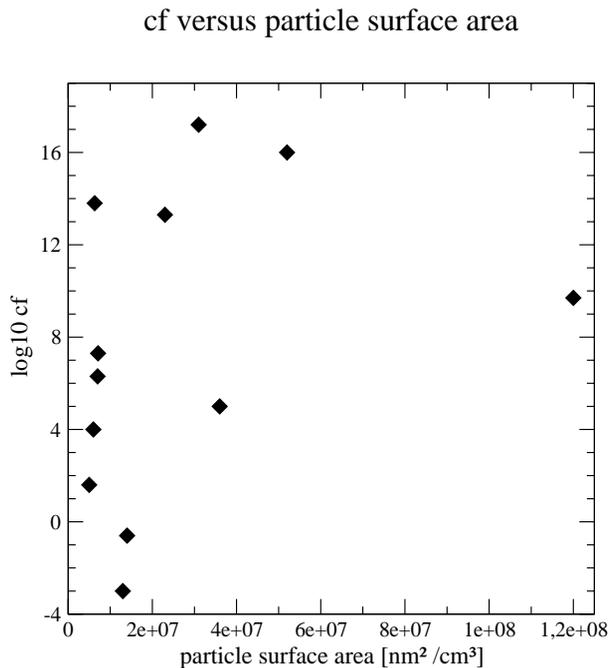


Fig. 7. Logarithm of nucleation rate correction, $\log_{10}(cf)$, versus particle surface area concentration.

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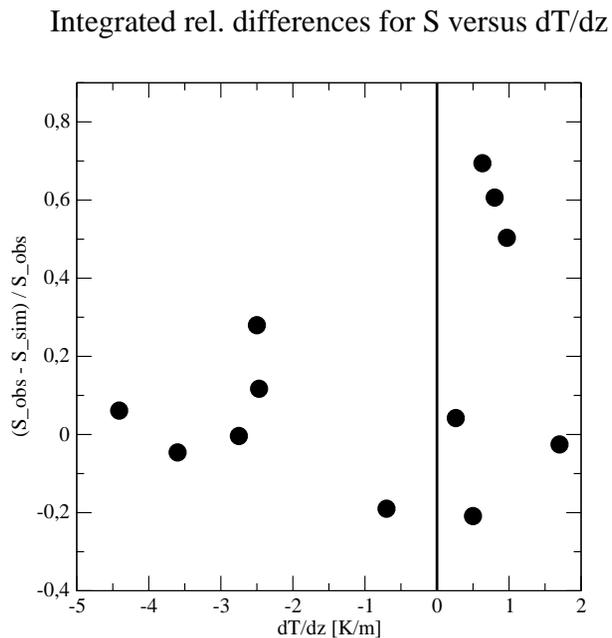


Fig. 8. Six-hour mean difference of the measured and simulated particle surface area concentration normalized by the measured particle surface area concentration, versus $\partial T/\partial z$.

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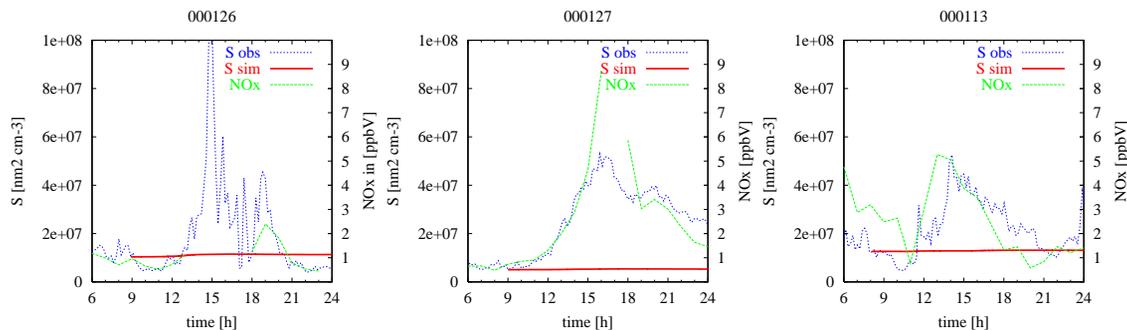


Fig. 9. Particle surface area and NO_x concentration for the three data sets with the highest difference between measured and simulated surface area concentration, ΔS .

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