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# The detection of solar proton produced $^{14}\text{CO}$

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Major solar eruptions (coronal mass ejections) are accompanied by massive ejections of protons. When these charged particles head for the Earth through the interplanetary magnetic field with high flux and energy, a solar proton event (SPE) is recorded. Strong SPEs, in which energetic protons penetrate the atmosphere in large numbers are rare, but do have chemical effects (Crutzen, 1975; Jackman et al., 1990, 2001). They also have nucleonic effects by which they can almost instantaneously increase the atmospheric production of radio-nuclides, including  $^{14}\text{C}$  (radiocarbon), but this has never been demonstrated. We show, using satellite observations and modeling, that the 2<sup>nd</sup> most intensive set of SPEs on record, that of August–December 1989, must have caused detectable increases in atmospheric  $^{14}\text{CO}$ . This is confirmed by a sequence of peaks in the Baring Head (NZ) time series of  $^{14}\text{CO}$  observations (Brenninkmeijer, 1993), providing a unique indication of production of  $^{14}\text{C}$  by solar protons, and demonstrating the use of SPE  $^{14}\text{CO}$  as an atmospheric tracer.

## 1. Introduction

The global annual, solar cycle averaged atmospheric production of  $\approx 7$  kg of  $^{14}\text{C}$  atoms via  $^{14}\text{N}(n,p)^{14}\text{C}$  by cosmic radiation (mainly from high energy protons) forms the basis for the widely applied radiocarbon dating (Libby, 1952).  $^{14}\text{C}$  enters the biosphere as  $^{14}\text{CO}_2$  through its assimilation by photosynthetic plants, and when for any living organism the exchange of  $^{14}\text{C}$  with the atmosphere ceases, the radiocarbon decay clock starts to tick. It is well known that  $^{14}\text{C}$  production has changed in the past, and that it varies with solar activity with a period of about 11 years. Yet, it has been shown that a 25% modulation (peak to peak) of the production of  $^{14}\text{C}$  due to the 11 year solar cycle, leaves only a very weak signal in atmospheric  $^{14}\text{CO}_2$  and in tree rings (Stuiver and Braziunas, 1993). The sudden increases associated with solar proton events can therefore certainly not be detected in  $^{14}\text{CO}_2$  or in tree rings, because the reservoir of

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$^{14}\text{CO}_2$  is large. However, before  $^{14}\text{C}$  enters the oceans and biosphere as  $^{14}\text{CO}_2$ , it resides in the atmosphere in the incompletely oxidized form of  $^{14}\text{CO}$ . Thus although a  $^{14}\text{C}$  atom instantaneously oxidizes to  $^{14}\text{CO}$  (Pandow et al., 1960; MacKay et al., 1963) the second chemical conversion,



takes on the average some months. This rapid but not immediate turnover leads to a small but measurable  $^{14}\text{CO}$  inventory; there are only 5 to 25  $^{14}\text{CO}$  molecules in  $1\text{ cm}^3$  of tropospheric air near the surface, compared to  $\approx 10^4$   $^{14}\text{CO}_2$  molecules. This then allows the actual detection of  $^{14}\text{C}$  production by SPEs, provided that observations of this ultra-low level tracer  $^{14}\text{CO}$  do exist.

Observations of  $^{14}\text{CO}$  have been made because it is the only natural atmospheric tracer available that can help to better quantify a fundamental property of the troposphere, namely its OH (hydroxyl) based oxidative activity or its self-cleansing capacity. With OH radicals constituting the main sink reaction of  $^{14}\text{CO}$ , and knowing its production rate, and measuring its levels in the atmosphere, OH abundance can be inferred according to equation (1) above (Volz et al., 1981; Brenninkmeijer et al., 1992; Mak and Southon, 1998; Jöckel et al., 1999, 2000, 2002; Jöckel and Brenninkmeijer, 2002). It is, however, in view of the rarity of SPEs, a coincidence that the second most powerful set of SPEs ever recorded, occurred just a few months after the first systematic observations of  $^{14}\text{CO}$  had started on the planet. In the following we first calculate using satellite data, the production of  $^{14}\text{C}$  during the fall 1989 SPEs. Next we calculate using a state of the art atmospheric chemistry transport model the expected temporal changes in  $^{14}\text{CO}$ . Finally we analyze the Baring Head (NZ)  $^{14}\text{CO}$  record for signs of these SPEs.

## 2. Model simulations of SPE produced atmospheric $^{14}\text{CO}$

About 50% of  $^{14}\text{C}$  production by cosmic rays takes place in the stratosphere, whereas for solar protons, which have lower energies, the  $^{14}\text{C}$  production maximum is shifted

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higher into the stratosphere (Fig. 1).

The energy spectrum of the 3 main 1989 SPEs (Fig. 2 gives the pulse shape) was measured by instruments of GOES-7 (Geostationary Operational Environmental Satellite). For solar protons the rigidity spectrum (momentum over charge) follows an exponential function with a characteristic rigidity  $P_0$  from 50 to 325 MV (Freier and Webber, 1963; Lingenfelter and Ramaty, 1970) and the total  $^{14}\text{C}$  production (Table 1) is determined by  $P_0$  and the flux  $I$  (Sauer et al., 1990; Shea and Smart, 1992; Feynman et al., 1993).

Because geomagnetic shielding might be weakened during geomagnetic storms occurring coeval to the SPEs,  $^{14}\text{C}$  production can be enhanced (Lingenfelter and Ramaty, 1970). Therefore Table 1 also includes estimates for an assumed 80% reduction in the cut-off rigidity (the minimum rigidity a proton needs to reach a particular point in the atmosphere). These estimates are considered as upper limits. Given both background GCR (Lingenfelter, 1963) and the SPE produced  $^{14}\text{C}$ , a 3-dimensional atmospheric transport and chemistry model (see Appendix) is used to calculate the resulting  $^{14}\text{CO}$  distribution (it is assumed that 95% of  $^{14}\text{C}$  is instantaneously converted to  $^{14}\text{CO}$  (Pandow et al., 1960; MacKay et al., 1963)). Essential is the spatial and temporal pattern of the production of  $^{14}\text{CO}$ , its transport in the atmosphere and its concurrent removal by OH. The zonally averaged enhancement ratio of  $^{14}\text{CO}$ , given as the ratio (SPE+GCR)/GCR, simulated for the lowest model layer (Fig. 3, left) shows that atmospheric  $^{14}\text{CO}$  temporarily increases up to 22% several months after the SPEs.

A distinct asymmetry between the two hemispheres is predicted by the model. The simulated downward transport of  $^{14}\text{CO}$  at this time of the year is almost as twice as effective in the northern hemisphere (NH) compared to the southern hemisphere (SH). Repeating the model simulation assuming an 80% reduced cut-off rigidity, the resulting signal shape of excess  $^{14}\text{CO}$  is similar, however with a maximum increase of 55%. The time lag between the largest SPE (No. 3 in Table 1) and maximum excess  $^{14}\text{CO}$  is dependent on the location as shown in Fig. 3 (right). The model does not resolve the individual events (the 2 main events are only 3 weeks apart), mainly because spatial

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resolution is limited.

### 3. Observations of SPE produced atmospheric $^{14}\text{C}$

With the model predicting significant increases in  $^{14}\text{C}$ , experimental verification is based on the Baring Head (41.4° S, 174.9° E, New Zealand)  $^{14}\text{C}$  record. This record (Fig. 4) commenced in June 1989 some months prior to the three strong SPEs, and comprises the only observational data available for this period.

The question whether  $^{14}\text{C}$  increases have been observed in the austral winter of 1989/90 will be answered by comparing this period to that one exactly one year later as reference (no SPEs).

A potential complication is that a fraction of the atmospheric  $^{14}\text{C}$  inventory is not directly due to cosmic radiation, but originates from  $^{14}\text{C}$  that has been recycled through the biosphere. All atmospheric CO produced from non-fossil organic compounds, by processes such as biomass burning or the oxidation of natural hydrocarbons in the atmosphere, contribute some  $^{14}\text{C}$ . This fraction can be estimated. Background CO at southern mid-latitudes is 90 to 95% of biogenic origin, as fossil fuel sources are sparse and little CO is imported from the NH (Manning et al., 1997). However, even this small fraction from the NH contains a substantial fraction of biogenic CO. Assuming that the background CO defined by the lower envelope (Fig. 4) (clear air, non-polluted conditions at Baring Head) is 100% biogenic, then 1.52  $^{14}\text{C}$  molecules/cm<sup>3</sup> STP, at the most, are estimated to be of biogenic origin in the SH (late) summer when the CO mixing ratio bottoms out at  $\approx 40$  nmol/mol. This increases to 2.47 molecules/cm<sup>3</sup> in spring (about 65 nmol/mol CO). Therefore, the fraction of biogenic  $^{14}\text{C}$  varies systematically between only 18% and 25%. Because the baseline CO mixing ratios at Baring Head (Fig. 4) for the consecutive (late) summer periods did not differ by more than 5%, biogenic CO cannot have affected observed  $^{14}\text{C}$  by more than about 1%. Accordingly, no correction for biogenic  $^{14}\text{C}$  in comparing the two consecutive years is necessary.

The  $^{14}\text{C}$  time series (Fig. 4) is smoothed with a low-pass convolution filter of the

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form  $f(t) = \exp(-t/T)^2$  using a time window  $T = 1$  week in the frequency domain (by Fourier transformation) after linear interpolation to obtain daily values. Because during 1989 and 1990 solar activity was at its maximum (solar cycle 22) it can be expected that the global average  $^{14}\text{C}$ O background production did not vary significantly between the two years. To totally rule out effects related to changing solar activity, other than the SPEs, the  $^{14}\text{C}$ O data further are normalized to equal conditions (Fig. 4) of solar activity using neutron monitor data of Mt. Wellington (42.92° S, 147.25° E, Tasmania) assuming that the relative changes in amplitude of neutron count rate and of  $^{14}\text{C}$  production during a solar cycle is the same. The resulting scaling of the global source strength is

$$c_q(t) = 1 + \left( \frac{n_{\text{smax}}}{n_{\text{smin}}} - 1 \right) \cdot \frac{(n(t) - n_{\text{smin}})}{(n_{\text{smax}} - n_{\text{smin}})}, \quad (2)$$

where  $n$  is the neutron count rate,  $t$  is the time, and the indices indicate solar minimum ( $n_{\text{smin}}$ ) and solar maximum ( $n_{\text{smax}}$ ). Further, an average response time  $\tau = 3.5$  months for atmospheric  $^{14}\text{C}$ O to changes in the global source strength is used for deducing the  $^{14}\text{C}$ O mixing ratio time dependence (Jöckel et al., 2000). The  $^{14}\text{C}$ O values are therefore scaled by

$$c_c(t) = \sum_{i=1}^{12} w_i c_q(t - i), \quad (3)$$

where  $i$  counts the months backwards,  $w_0 = 0.14$  is the estimated relative contribution of the current month's production rate,  $w_i \propto \exp(-i/\tau)$ , and  $\sum_{i=0}^{12} w_i = 1$ . After having corrected the  $^{14}\text{C}$ O record for modulation by solar activity, and having shown that the effect of changes in total CO, via biogenic  $^{14}\text{C}$ O, are negligible, the remaining differences between the 2 consecutive years have only 2 possible causes. One is that synoptic scale changes in atmospheric circulation affect  $^{14}\text{C}$ O levels. For instance, soon after the record had started in winter 1989, conditions of southerly air mass and southerly gales at the site were frequent in July and later again in October, leading to elevated  $^{14}\text{C}$ O levels, manifest in Fig. 5 as the first 2 clear peaks.

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However the three subsequent peaks (framed by the green box in Fig. 5) occur  $117 \pm 2$ ,  $89 \pm 2$  and  $114 \pm 2$  days after the three SPEs, within the period for which the model predicts enhanced  $^{14}\text{CO}$ . The observed relative peak height does correspond to the strength of the SPEs in terms of the  $^{14}\text{CO}$  production rate. Moreover, the time lag between each SPE and the respective measured maximum excess  $^{14}\text{CO}$  does correspond to the characteristic rigidity of the events (Table 1). With increasing characteristic rigidity the SPE induced  $^{14}\text{CO}$  production reaches deeper into the atmosphere and it takes less time to transport the excess signal to the surface. These are strong indications that these 3 peaks observed are indeed caused by the SPEs. Following these events, considerable oscillations in the ratio between the consecutive years occur related to strong changes in  $^{14}\text{CO}$  in February–June 1991, again related to meteorological conditions. Moreover these peaks occur during the steep increase in  $^{14}\text{CO}$  during fall. The rapid building up of a strong north south  $^{14}\text{CO}$  latitudinal gradient augments the impact of synoptic scale meteorological differences.

For testing the sensitivity of resolving the 3 peaks to the sequence and distribution of actual data points, the analyses were repeated with reduced data sets from which randomly 10% of the data points had been omitted. All results obtained were very similar to the one using the complete set.

#### 4. Cross tropopause transport of SPE produced $^{14}\text{CO}$

Because the timing of the transport of SPE  $^{14}\text{CO}$  from the stratosphere into the troposphere is critical, an independent estimate is derived from the mean downward motion of air in the lower stratosphere at high latitudes using meteorological data. This downward motion can be regarded as a part of the global stratospheric Brewer-Dobson circulation (Holton et al., 1995). We diagnosed the Transformed Eulerian Mean (TEM) vertical velocity  $\omega^*$  (Peixoto and Oort, 1992), using the 15 year (1979–1993) ECMWF reanalysis (ERA) data set (Gibson et al., 1997). The TEM circulation is the Eulerian mean circulation in which the part forced by eddy heat transport is removed. For each

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of the 180 months of the ERA period, the  $\omega^*$  in the lower stratosphere has been diagnosed separately, with a latitudinal resolution of 2.5 degrees and at the levels of 10, 30, 50, 70, 100, 150, 250 and 300 hPa. The results applied in this study are 15-year August to November averages over the regions poleward of  $60^\circ$ , along with the respective averages for the year 1989. From  $\omega^*$  the average subsidence time between the levels listed in Table 2 has been computed ( $\Delta t = \Delta p / \omega^*$ ).

The year-to-year variation of this subsidence time is computed (error propagation) as standard deviation  $\sigma_{\Delta t} = (\Delta p / \omega^{*2}) \sigma_{\omega^*}$ , where  $\Delta p$  is the pressure interval, and  $\sigma_{\omega^*}$  is the standard deviation of  $\omega^*$  from the 15-year average. Computations for the two hemispheres of the extra-tropical average of  $\omega^*$  (i.e. of the strength of the Brewer-Dobson circulation) using ERA data give generally somewhat larger (tens of %) values than found in similar studies. For the  $\omega^*$  as used in this study, averaged as described above, an educated guess of its accuracy is 50%. More accurate vertical velocity data are presently not available. Assuming a time lag of 14 days for the air that just crossed the tropopause (Fig. 1) to reach the surface (tropospheric vertical mixing time), the maximum observed excess  $^{14}\text{CO}$  at Baring Head induced by the three SPEs should then originate from a pressure level between 95 hPa and 120 hPa, not considering the uncertainty in the subsidence velocity. With a 50% overestimate of the subsidence velocity, the respective levels are 55 hPa and 70 hPa. The maximum of the SPE induced  $^{14}\text{CO}$  production rate is expected to be somewhat higher (see Fig. 1), at about 30 hPa, however depending further on the characteristic rigidity of the particular event and the cut-off rigidity. Furthermore, the production rate between 70 hPa and 100 hPa exhibits a significant contribution to the total SPE induced  $^{14}\text{CO}$  production.

## 5. Conclusions

Following the 3 major 1989 SPEs, increases in  $^{14}\text{CO}$  have been observed at Baring Head with a timing and intensity that corresponds to the flux and characteristic rigidity of these events. The magnitude of the increase in  $^{14}\text{CO}$  over the period of the 3 events

is comparable to that calculated by the 3D model, using a normal cut-off rigidity. The reliability of the transport time estimates is further supported by independent meteorological analyses, leaving little to no doubt that indeed SPE derived  $^{14}\text{C}$  has been detected.

## 5 Appendix: Model setup

Model simulations have been performed with the 3-dimensional global model of atmospheric transport and chemistry (MATCH) (Rasch et al., 1997) in the Max-Planck-Institute for Chemistry (MPICH) version 2.0 (Lawrence et al., 1999). The employed SPITFIRE advection scheme (Rasch and Lawrence, 1998) is driven by re-analysed wind data from NCEP (Kalnay et al., 1996) using the year 1993 (inter-annual variations in transport and chemistry are neglected). The oxidative removal of  $^{14}\text{C}$  by OH, its distribution in the troposphere (Lawrence, 1996) and the stratosphere (2-dimensional model calculations, Ch. Brühl, personal communication 2001) are described off-line (monthly averages). The small uptake of  $^{14}\text{C}$  by soils was also included (Conrad and Seiler, 1985; Sanhueza et al., 1998). The spatial distribution of GCR  $^{14}\text{C}$  production (Lingenfelter, 1963) has been transformed from geomagnetic to geographic coordinates using altitude adjusted geomagnetic coordinates (AACGM) (Bhavnani and Hein, 1994). The distribution pattern is assumed to be constant, since the effect of the varying  $^{14}\text{C}$  distribution with the solar cycle on atmospheric  $^{14}\text{C}$  is small (Jöckel et al., 1999). The varying global average source strength of GCR  $^{14}\text{C}$  is parameterized by linear interpolation between solar maximum and minimum using monthly mean sunspot numbers (Lingenfelter, 1963). A yield of  $^{14}\text{C}$  from  $^{14}\text{C}$  of 95% is assumed for both the GCR and the SPE component (MacKay et al., 1963; Pandow et al., 1960). The time dependence of the  $^{14}\text{C}$  production during the 3 SPEs is derived from solar proton measurements on board the GOES-7 satellite. The total amount of  $^{14}\text{C}$  production is distributed (daily averages) according to the flux of protons with energy greater than 30 MeV (Fig. 2). The model simulations were initialized with a global  $^{14}\text{C}$  distribution

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precalculated with the same model (2 year integration starting from zero mass mixing ratio) without SPEs. One model run including the 3 major SPEs (assuming normal cut-off rigidity) and one run without SPEs simulate the effect of the SPEs. Model output was archived as 5-day averages.

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**Table 1.** Characteristic rigidity  $P_0$ , flux  $I_{>30\text{MeV}}$  of protons with energy greater than 30 MeV, and estimated total  $^{14}\text{C}$ O production for normal cut-off rigidity ( $Q_{100\%}$ ) and an 80% reduced cut-off rigidity ( $Q_{20\%}$ ) of the 3 major SPEs in 1989. For comparison, the amount of  $^{14}\text{C}$ O produced during the SPEs is also listed as fraction (in %) of the annual global GCR produced background  $^{14}\text{C}$ O (normalized to  $1 \text{ molec cm}^{-2} \text{ s}^{-1}$  global average), both for normal cut-off rigidity ( $f_{100\%}$ ), and for the reduced cut-off rigidity ( $f_{20\%}$ ).

| Date             | $P_0$<br>MV | $I_{>30\text{MeV}}$<br>$\text{cm}^{-2}$ | $Q_{100\%}$<br>$\text{molec cm}^{-2}$ | $f_{100\%}$<br>$\%/ \text{molec cm}^{-2} \text{ s}^{-1}$ | $Q_{20\%}$<br>$\text{molec cm}^{-2}$ | $f_{20\%}$<br>$\%/ \text{molec cm}^{-2} \text{ s}^{-1}$ |
|------------------|-------------|---|---------------------------------------|--|--------------------------------------|---|
| 12 Aug.–07 Sept. | 60.6        | $1.53 \cdot 10^9$                       | $1.027 \cdot 10^6$                    | 3.3  | $2.433 \cdot 10^6$                   | 7.7   |
| 29 Sept.–13 Oct. | 102.0       | $1.42 \cdot 10^9$                       | $2.513 \cdot 10^6$                    | 8.0  | $6.210 \cdot 10^6$                   | 19.7  |
| 19 Oct.–09 Nov.  | 77.4        | $4.25 \cdot 10^9$                       | $4.378 \cdot 10^6$                    | 13.9   | $1.079 \cdot 10^7$                   | 34.2  |

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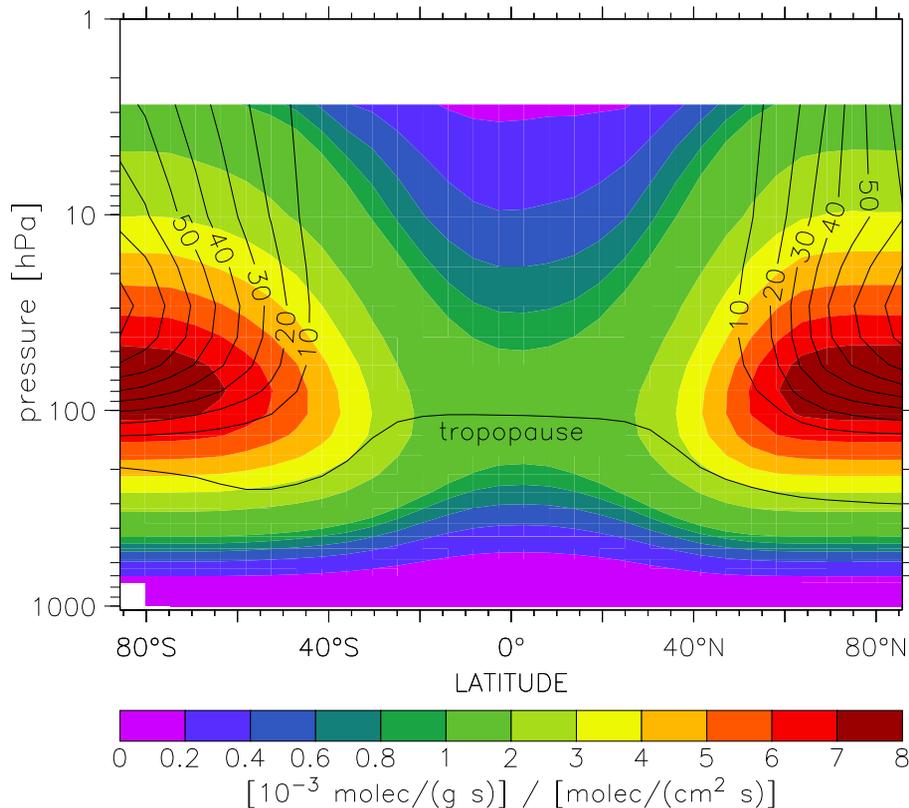
**Table 2.** ECMWF reanalysis based subsidence durations between the listed pressure levels in days for the northern hemisphere (NH) and the southern hemisphere (SH). The values are averaged over 60° to 90° latitude, and the months of August to November. For the averages of the years 1979 to 1993 the uncertainty is the standard deviation according to the year-to-year variation as described in the text. In addition, the respective subsidence times for the year 1989 are listed.

| Pressure [hPa] | subsidence duration [days] |      |           |      |
|----------------|----------------------------|------|-----------|------|
|                | SH                         |      | NH        |      |
|                | 1979–1993                  | 1989 | 1979–1993 | 1989 |
| 30 to 70       | 129 ± 39                   | 140  | 177 ± 53  | 222  |
| 70 to 100      | 64 ± 14                    | 60   | 75 ± 21   | 95   |
| 100 to 150     | 40 ± 6                     | 36   | 51 ± 9    | 58   |
| 150 to 200     | 19 ± 2                     | 17   | 23 ± 3    | 25   |
| 200 to 250     | 15 ± 2                     | 13   | 18 ± 2    | 19   |
| 250 to 300     | 18 ± 3                     | 17   | 16 ± 2    | 17   |

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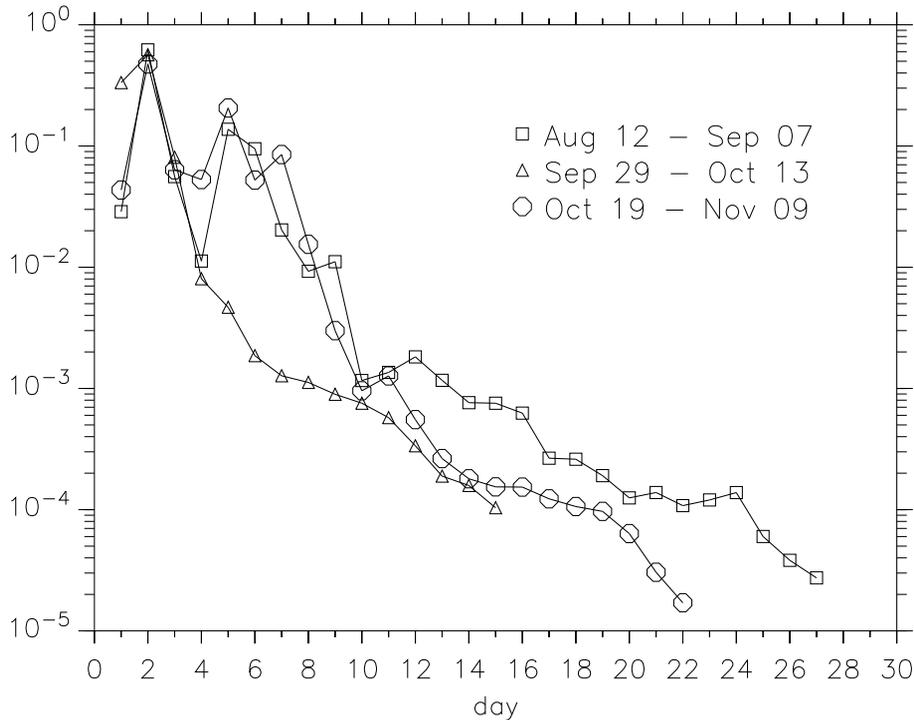
**Fig. 1.** Annual zonal mean galactic cosmic ray induced  $^{14}\text{C}$ O production rate (GCR, shaded) and annual zonal mean solar proton event induced  $^{14}\text{C}$ O production rate (SPE, contour lines). The unit is  $10^{-3} \text{ molec g}^{-1} \text{ s}^{-1}$  normalized to a global average production rate of  $1 \text{ molec cm}^{-2} \text{ s}^{-1}$ . The yield of  $^{14}\text{C}$ O from  $^{14}\text{C}$  oxidation is assumed to be 95%. The plotted zonal average tropopause level between August and November is calculated from the NCEP reanalysis data for 1993 according to the WMO definition.

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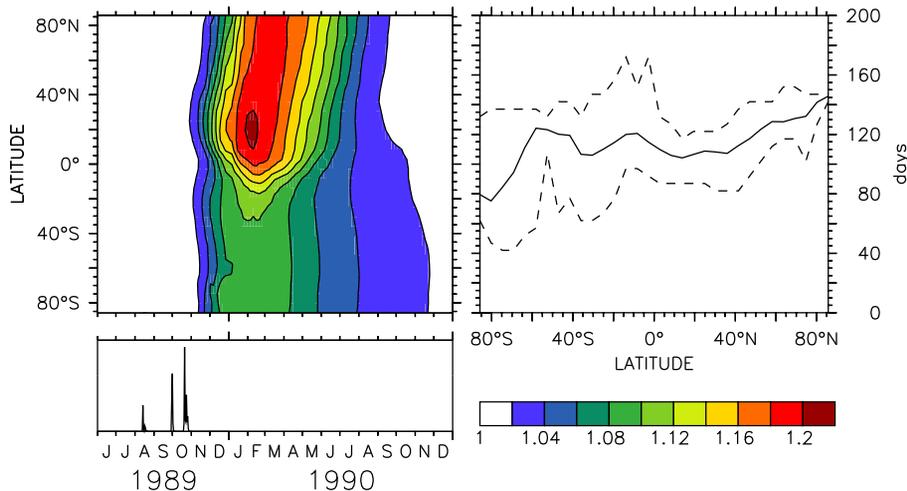
**Fig. 2.** Pulse shape of  $^{14}\text{C}$ O production for the 3 major SPEs in 1989, based on GOES-7 solar proton flux measurements. The vertical axis measures the fraction of the total  $^{14}\text{C}$ O production during the indicated time interval of the respective event.

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**Fig. 3.** Left: Zonally averaged enhancement of  $^{14}\text{C}$ O (zonal mean ratio of the  $^{14}\text{C}$ O mixing ratio calculated including SPEs to the  $^{14}\text{C}$ O GCR background) in the lowest model layer after the 3 SPEs, calculated with the 3-D model for normal cut-off rigidity. The peaks in the lower panel are proportional to the  $^{14}\text{C}$ O production of the respective SPEs. Right: Zonal average time lag between the largest SPE and the maximum excess  $^{14}\text{C}$ O (solid line). The dashed lines indicate the maximum and the minimum predicted time lag at a given latitude.

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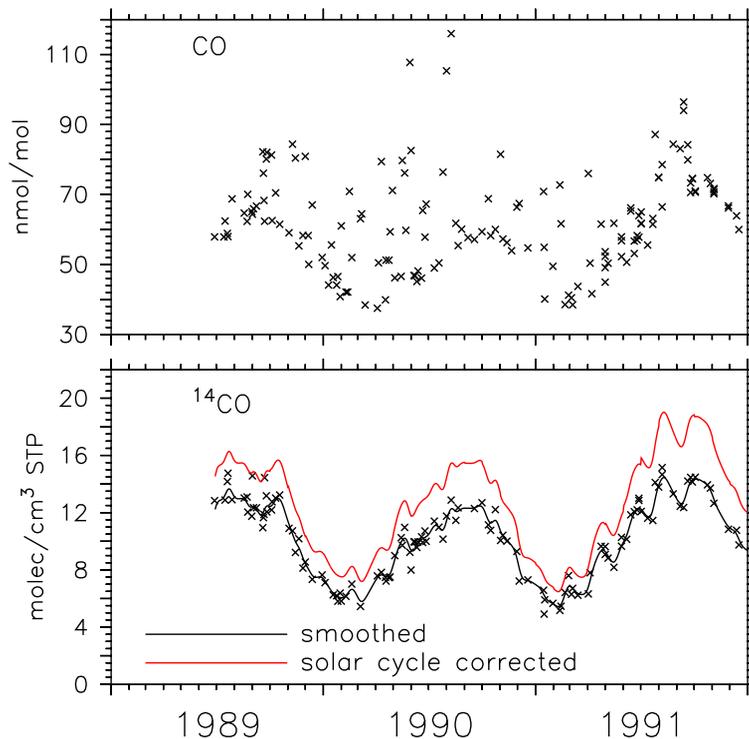
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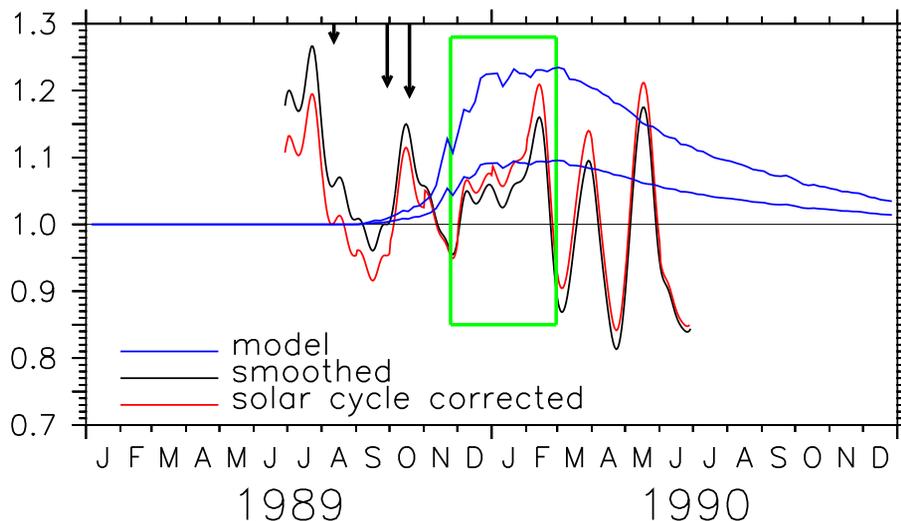
**Fig. 4.** CO (upper) and  $^{14}\text{CO}$  (lower) measurements at Baring Head ( $41.4^\circ\text{S}$ ,  $174.9^\circ\text{E}$ ), New Zealand for the years 1989 to 1991. The black solid line shows the result of the data smoothing, the red line indicates the smoothed data corrected with respect to solar variation of the GCR  $^{14}\text{CO}$  background production rate using neutron count rates of the Mt. Wellington ( $42.92^\circ\text{S}$ ,  $147.25^\circ\text{E}$ ) neutron monitor. The scatter in CO is the result of CO from local sources during non background conditions. This CO is free of  $^{14}\text{CO}$ . Non clean air conditions were included to not bias the  $^{14}\text{C}$  record towards specific meteorological conditions.

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**Fig. 5.** Observed (black), and solar cycle adjusted observed (red, cf. Fig. 4), enhancement of  $^{14}\text{C}$ O for June 1989 to June 1990 with June 1990 to June 1991 as reference year at Baring Head. The 3 SPEs in 1989 are indicated by the arrows at the top. The relative arrow-lengths correspond to the estimated total  $^{14}\text{C}$ O production of the respective SPE. The blue lines show the results of the model prediction for normal cut-off rigidity (lower line) and an 80% reduced cut-off rigidity (upper line).

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