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### **► To cite this version:**

S. Gabersek, R. Sorgente, S. Natale, A. Ribotti, A. Olita, et al.. The Sicily Channel Regional Model forecasting system: initial boundary conditions sensitivity and case study evaluation. Ocean Science, 2007, 3 (1), pp.31-41. <hal-00298308>

**HAL Id: hal-00298308**

**<https://hal.science/hal-00298308v1>**

Submitted on 18 Jun 2008

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# The Sicily Channel Regional Model forecasting system: initial boundary conditions sensitivity and case study evaluation

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Received: 30 March 2006 – Published in Ocean Sci. Discuss.: 17 May 2006

Revised: 1 August 2006 – Accepted: 17 August 2006 – Published: 30 January 2007

**Abstract.** The Sicily Channel Regional Model forecasting system was tested using an optimization package for the initial and lateral boundary conditions. Spurious high frequency oscillations during the spin-up time were successfully reduced both in duration and magnitude by optimizing the time tendency of the free surface elevation using the Variational Initialization and Forcing Platform method developed in the framework of the Mediterranean Forecasting System Towards the Environmental Prediction project. The effect of optimization was most profound for the free surface elevation, where all oscillations with periods shorter than 4 h were suppressed.

The overall forecast skill was assessed on a 5 day case study starting on 6 April 2005, characterized by a fast passage of a deepening atmospheric low-pressure field with strong winds and marked wind direction change. We compared the predicted variables with in-situ and remotely sensed data. The forecasts of temperature, including the sea surface temperature, and salinity were quite successful, while the forecasted currents, especially within the surface layer, were not in good agreement with the measurements.

## 1 Introduction

The numerical simulation of the sea circulation over extended coastal shelf areas, such as the central Mediterranean basin, is extremely important in order to assess the short and long term environmental variability. However, in a short range forecast (order of days) of ocean properties, the spin-up time becomes a problem of great importance especially at the coastal time scales where spurious oscillations and propagating waves are present in results for several days after the initialization. (Auclair et al., 2001).

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Regional models instead of OGCM are used in order to capture the mesoscale scale phenomena. The use of Limited Area Model (LAM) as SCRM at regional scale embedded into a coarse resolution ocean model (OGCM) represents an efficient way to downscale the model solutions from the basin-scale (12.5 Km) to the regional scale (3 Km) through a one-way, off-line nesting at the open lateral boundaries. This method was found to be computationally efficient and sufficiently robust to transmit information across the lateral boundaries without excessive distortion (Sorgente et al., 2003).

The spin-up is the time needed by an ocean model to reach a state of physical equilibrium under the applied forcing. The results can not be trusted until this equilibrium is reached due to spurious noise in the numerical solution. For an ocean basin-scale circulation model the spin-up can take hundreds of years to achieve an equilibrium, while regional models, that do not include deep ocean basins, may be much quicker. In order to start a model run for operational forecast or research purposes, some variables need to be specified at initial time. These variables include potential temperature, salinity, free surface displacement and velocity fields.

One way to initialize models is by using climatological values of temperature and salinity from databases and assuming the velocity field equal to zero at the start time (Fox-Rabinovitz et al., 2002). The model will adjust the velocity field in balance with the density and the forcing field. As the forcing is applied, the velocity field will respond to the initial condition with a transient flow that may be unrealistic. For this reason, the results from the beginning of ocean circulation model runs are usually not used. This kind of approach, commonly used to initialize global or basin scale models, is named “cold start”.

Another common way is to initialize an ocean model by using fields from a previous run of the same model. If the saved fields from a recent forecast of the same model are used to initialize a new simulation, we call this approach a

“warm start”. The main advantage is that the ocean state has already been adjusted to the initial condition and is available for use in prediction.

One can also use dynamically balanced forecast fields from a model with coarser spatial resolution to initialize a model with higher resolution, but the phenomena of the smaller scale are not captured in the initial conditions. This is again a “cold start” and this approach has been used in the Sicily Channel Regional Model (hereafter SCRM) forecasting system, that will be discussed later.

In the “cold start” used in the SCRM we applied an innovative tool based on the Variational Initialization and Forcing Platform (hereafter VIFOP) method (Auclair et al., 2000) in order to initialize and force the regional forecast system from the coarse grid solution. The aim is to reduce the amplitude of the spurious external gravity wave generated during the spin-up time and to assess the sensitivity of the SCRM forecasting system to optimized initial and lateral boundary conditions. Furthermore, we want to perform a quantitative evaluation of forecasted fields in a case study.

The SCRM has been implemented in the framework of the Mediterranean Forecasting System Towards the Environmental Prediction project (hereafter MFSTEP). The SCRM has been realized for the Central Mediterranean sub-basin including the Sicily and Sardinia Channels (hereafter SiC and SaC, respectively) and the wide Tunisian/Libyan shelf. Although there exist modeling studies of the area (Molcard et al., 2002; Pierini and Rubino, 2001), the SCRM is the first operationally used model.

This study area is characterized by a complex bathymetry with wide continental shelves, deep and shallow channels and wide abyssal plains. Its central position in the Mediterranean basin plays a crucial role in a passage of the superficial and intermediate water masses in transit between the Eastern and the Western Mediterranean sub-basins (hereafter EM and WM, respectively). It also prevents the water masses from the deep and bottom layers of the two basins from mixing. The hydrology in the study area can be described as a multi-layer system composed of a series of adjacent water masses both in the horizontal and in the vertical direction. The upper layer is occupied by a surface layer (0–100 m) of relatively fresh water of Atlantic origin, frequently called Modified Atlantic Water (hereafter MAW). It is described as a broad homogeneous layer that undergoes progressive modifications becoming warmer and saltier as it spreads eastward toward the EM basin (Manzella, 1994; Moretti et al., 1993). The inflow of MAW from the SaC drives the so-called Algerian Current (hereafter AC) and its pathway is influenced by the topography, the surface forcing and by the density gradient between the Atlantic and the Mediterranean waters. The lateral variations in density with the surrounding water masses render the AC unstable, with the generation of instabilities represented by rather energetic mesoscale eddies. The surface circulation is characterized also by meanders, semi-permanent cyclonic and anticyclonic gyres and eddies,

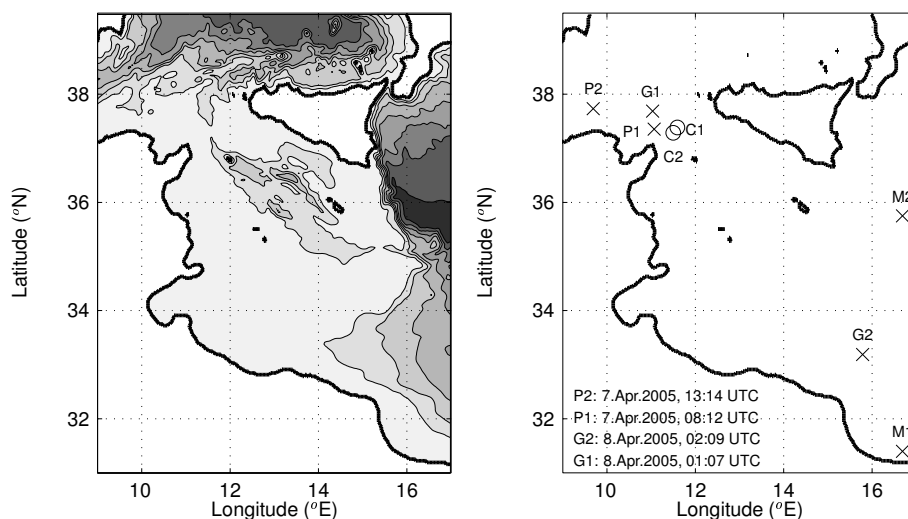
some of them associated with the instabilities of the AC and the meandering of the Atlantic Ionian Stream (hereafter AIS) (Lermusiaux, 1999; Lermusiaux and Robinson, 2001; Robinson et al., 1999). The AIS is a summer feature and is the main transporter of the MAW from the SaC to the open Ionian Sea, together with the Atlantic Tunisian Current flowing along the Tunisian-Libyan shelf break (Beranger et al., 2004; Sorgente et al., 2003). At the bottom part of the MAW there is a thin layer traced by means of a relative temperature minimum located between depths of about 100 and 200 m. It has been associated to the Winter Intermediate Water, previously formed in the WM by surface cooling of MAW during the winter (Fuda et al., 2000; Ribotti et al., 2004; Sammari et al., 1999). Between about 200 m and the bottom there are the Levantine Intermediate Water (hereafter LIW) and, below, the transitional Eastern Mediterranean Deep Water forming the Eastern Mediterranean Overflow Water (hereafter EOW). The LIW is characterised by a relatively high salinity and temperature. From the EM the EOW flows westward at an intermediate depth of approximately 300 m towards the shallowest depths of the SiC through the Maltese Sill (Fig. 1), where it passes two sills at 360 and 430 m deep before spreading into the Tyrrhenian Sea. Here its stream bifurcates under topographic constraints and the Coriolis effect, and partially mixes with the upper and lower local waters (Sparnocchia et al., 1999; Astraldi et al., 2002).

The paper is structured as follows: in Sect. 2, the model setup is described; sensitivity tests using VIFOP are in Sect. 3; a case study is in Sect. 4; conclusions follow in Sect. 5.

## 2 Model setup

The numerical model used to simulate the ocean circulation at a regional scale over the Central Mediterranean region is the Sicily Channel Regional Model (SCRM) developed within the MFSTEP project framework. It is based on the Princeton Ocean Model (POM) (Blumberg and Mellor, 1987), a three-dimensional, hydrostatic, free surface ocean model using the Boussinesq approximation. The horizontal mesh is orthogonal and the vertical coordinate is terrain following ( $\sigma$ ). The domain extends from 9° to 17° E and from 31° to 39.5° N, with horizontal resolution  $\Delta x$ ,  $\Delta y$  of 1/32° (~3.5 km) (Fig. 1). There are  $n_x=257 \times n_y=273$  mesh points with 24  $\sigma$  levels. The numerical model uses time splitting with the external (barotropic) and internal (baroclinic) mode time steps of 4 and 120 s, respectively. The bathymetry is based on a regular grid with spacing of 1 min, bilinearly interpolated onto a regular SCRM grid (3 km).

SCRM is driven through initial and lateral boundary conditions by a coarse resolution model using a one-way nesting of daily mean fields of temperature, salinity and velocity from a weekly forecast (MFS1671, Madec et al., 1998). Forcing at the surface is done using an air-sea coupling



**Fig. 1.** The SCRM domain with bathymetry in grey contours from light (shallow) to dark (deep), contour interval 500 m. The thick black line represents the coast line. Locations of in-situ measurements with XBT profilers are indicated with P1 and P2, Med-Argo gliders with G1 and G2, and with chained current meters with C1 and C2. Two additional points (M1,M2) indicate the location of model output points used for frequency analysis.

algorithm based on bulk parameterizations for the computations of the momentum, heat and fresh water fluxes at the air-sea interface. The atmospheric fields needed for the surface boundary conditions were obtained using Skiron (Kallos et al., 1997).

The SCRM operational forecast system performs a 5 day ocean forecast in slave mode, i.e. it is re-initialized for each operational weekly forecast from the coarse resolution model. The 5 day forecasts start at 00:00 UTC on each Wednesday. In this paper we discuss only a five day period from 6 April 2005 (00:00 UTC) to 11 April 2005 (00:00 UTC), during which there was an atmospheric low-pressure system passage over the SCRM domain. The rapid deepening was accompanied with strong winds at the surface.

Additional four experiments were performed using different dates (one 5 day forecast in each month from September to December 2004) with similar conclusions (results not shown in the paper). The April case was chosen because we initially planned to test the sensitivity to different atmospheric forcings as well.

### 3 VIFOP experiments

This section provides a description of sensitivity tests performed with VIFOP. VIFOP is a variational balanced initialization technique able to analyze the observations or the outputs of regional scale or basin scale circulation models used as initial field or open boundary forcing of high resolution ocean models. Such method drastically reduces the amplitude of numerical transient processes following the initial-

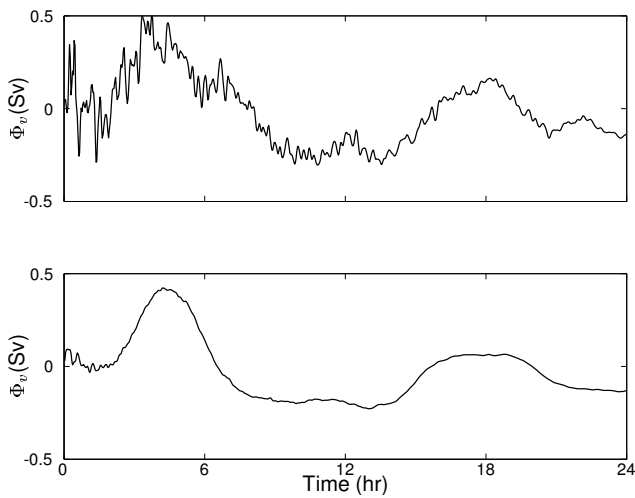
ization by reducing the misfit of the initial field. The method is based on the minimization of a cost function involving data constraints as well as a dynamical penalty involving the tangent linear model. The optimization focuses on the external mode, more specifically on the time tendency of the free surface elevation. The tendency can be decomposed into a physical and numerical part and the goal is to minimize the spurious numerical tendency. For a detailed description of the package see Auclair et al. (2000) and for a more general text on data assimilation see Daley (1991) or Kalnay (2003).

The VIFOP sensitivity tests included seven experiments, in which various combinations of initial and lateral boundary conditions were tested. Besides the prescribed fields of temperature, salinity and velocity, there is also an additional requirement for the volume transport (VT) across boundaries of the fine resolution model to match that of the coarse resolution model (Zavatarelli and Pinardi, 2003). The VIFOP sensitivity tests were performed for a 5 day forecast period. The atmospheric fields used to provide the surface boundary conditions were the same as in the operational forecasting (Skiron). The coarse scale fields are from the Ocean General Circulation Model (OGCM) OPA, suited for the MFSTEP project (MFS1671 – 1/16° horizontal resolution, 71 vertical levels). The sensitivity tests are described in Table 1, specifying the means of obtaining both the initial (IBC) and lateral (LBC) boundary condition, as well as the volume transport adjustment (VT).

During the full optimization, the initialization of the external mode is controlled by values of various parameters (e.g. weak or strong constraint during the optimization of the surface elevation tendency, optimization of the global divergence, optimization of the depth averaged velocities,

**Table 1.** List of sensitivity experiments. OPA and VIFOP stand for non-optimized and optimized IBC, LBC and VT, respectively. Subscript  $i$  stands for interpolation only (no optimization).

name	IBC	LBC	VT	comment
exp <sub>0</sub>	OPA	OPA	OPA	The control experiment, no optimization. Coarse grid fields are interpolated onto the finer grid using a bilinear interpolation to obtain IBC, LBC and VT.
exp <sub>1</sub>	VIFOP <sub>i</sub>	OPA	OPA	The IBC interpolation is done by VIFOP, where all points within the radius of influence ( $R=10$ km) are considered, their influence is determined by the Gaussian weight ( $G=6$ km).
exp <sub>2</sub>	VIFOP <sub>i</sub>	VIFOP <sub>i</sub>	OPA	In exp <sub>2</sub> , VIFOP performs interpolation also on LBC.
exp <sub>3</sub>	VIFOP <sub>i</sub>	VIFOP <sub>i</sub>	VIFOP <sub>i</sub>	VIFOP interpolation of IBC and LBC. VT is obtained from interpolated fields.
exp <sub>4</sub>	VIFOP	OPA	OPA	Full optimization of IBC. The rest is the same as in exp <sub>1</sub> .
exp <sub>5</sub>	VIFOP	VIFOP	OPA	Full optimization of IBC and LBC.
exp <sub>6</sub>	VIFOP	VIFOP	VIFOP	Full optimization (LBC, IBC, VT).



**Fig. 2.** The total volume transport  $\Phi_v$  (Sv) through the SCRM boundaries for exp<sub>0</sub> (top) and exp<sub>6</sub> (bottom).

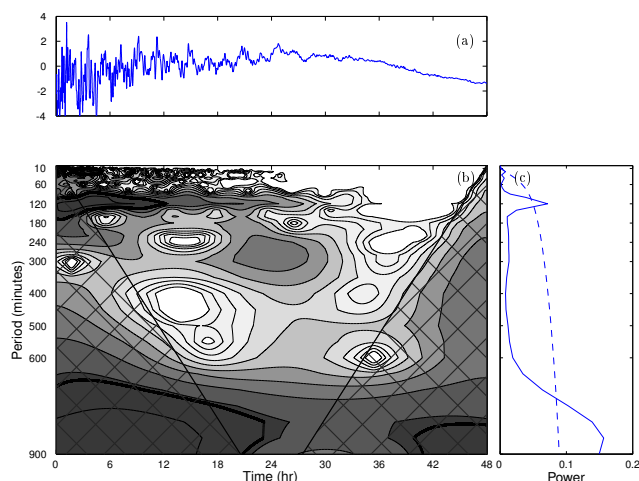
etc.). Various combinations were tested (not shown) and the best results were obtained by optimizing the global divergence and using the strong constraint on the free surface tendency. Unless otherwise noted this setup is used throughout the paper.

Surface or volume averaged quantities, like barotropic or baroclinic kinetic energy, respectively, do not reveal any differences among the experiments (not shown). Volume transport through lateral boundaries of the domain is very sensitive to the LBC. The total volume transport through the SCRM boundaries during the first day of simulation is shown in Fig. 2. The high frequency oscillations of fluxes are present in all experiments with interpolation only, regardless of method used (exp<sub>0</sub>,...,exp<sub>3</sub>). In fact, the volume transport was similar in all experiments and for clarity only exp<sub>0</sub> is shown (Fig. 2 top panel). When VIFOP is used for the full optimization, the high frequency fluctuations in the total vol-

ume transport are of much smaller magnitude to begin with and become almost nonexistent after 12 h of simulations. It is important to note the use of VIFOP for IBC only (exp<sub>4</sub>, not shown) to some extent decreases the spurious oscillations, but the additional optimization of LBC leads to even smoother volume transport (Fig. 2 bottom panel).

A further insight into impact of VIFOP was provided by a set of additional runs in which the model was integrated for two days only, instead of the regular 5 day forecast, but an array of variables was saved at each time step. The chosen variables were the free surface elevation (output every barotropic time step of 4 s), and  $u$  component of velocity (output every baroclinic time step of 120 s). There were two sets of ten adjacent points at a constant latitude, spaced at  $\Delta x$ , close to the eastern boundary at i)  $\phi=31.4^\circ$  N to observe the shallow water response, and ii)  $\phi=35.75^\circ$  N to observe a response over the deepest section of the domain (see Fig. 1 for location of points M1 and M2). The reason to monitor the chosen variables at ten adjacent points was to observe a transition from the boundary to the interior. The values at the outermost points ( $i=nx$ ) were determined by the lateral boundary conditions from the OGCM, and therefore changed linearly between the two subsequent daily values imposed through the LBC. The remaining nine points exhibited similar features during the simulation, so only values at the innermost point ( $i=nx-10$ ) are discussed.

Since the high frequency oscillation associated with the model initialization diminish in time, a Fourier analysis of the whole time series is not appropriate. An alternative would be to perform a windowed Fourier analysis, but because we were interested in time–frequency localization of the signal, we used a wavelet analysis. It is now widely used in geophysics (see reviews in Torrence and Compo, 1998, Kumar and Foufoula-Georgiou, 1997, Massel, 2001). During the analysis, the mother wavelet is scaled and translated along the original signal to perform a convolution. By scaling the wavelet function, signals at various scales from the data

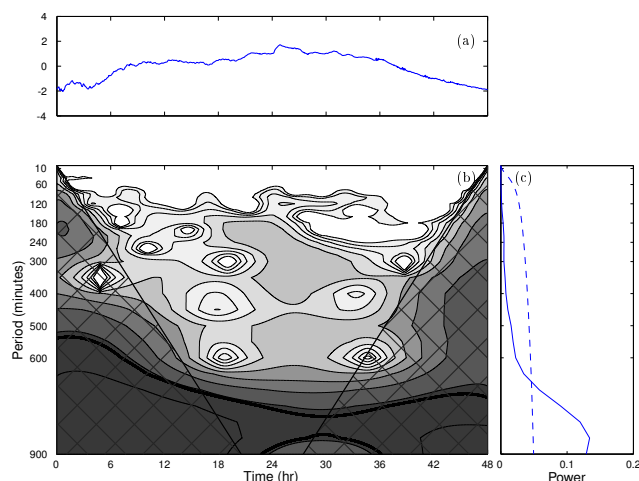


**Fig. 3.** (a) Time series of a detrended and normalized free surface elevation at  $i=248$ ,  $j=153$  ( $\phi=35.75^\circ$  N,  $\lambda=16.7^\circ$  E). (b) Wavelet power spectrum of the same time series. The contour interval values for normalized variance are from light to dark in range [1, 2, ..., 10]. The thick black line indicated the 95% confidence level for a red noise background, the hashed region represents the cone of influence. (c) The global power spectrum, with the 95% confidence level (dashed line). See text for a more detailed explanation.

series can be picked up, analogous to Fourier periods. By shifting the wavelet function along the data series and performing the transformation, the location of the signals in the original data are determined.

We chose the Morlet wavelet, a plane wave modulated by an exponential decay, with a nondimensional frequency  $\omega_0=6$ . An additional beneficial characteristic of this particular wavelet is its approximate one-to-one ratio between the Fourier period and wavelet scale (Torrence and Compo, 1998). The set of scales to be resolved depended on the analyzed variable, ranging from  $\mathcal{O}(1$  min) for the free surface displacement, to  $\mathcal{O}(1$  day) for the  $u$  component of velocity.

The data values from the model were first detrended and normalized by standard deviation, followed by the wavelet analysis. Each analysis presented in the paper (Figs. 3 and 4) consists of three panels. The top panel is detrended and normalized model output data time series. The bottom panel is the local wavelet power spectrum in shaded contours with darker shades representing increasing values of normalized variance. A cone of influence indicates possible errors related to the edge effects because the data series were finite, but the Fourier transform assumes the data to be cyclic (the transform is used for convolution purposes, see e.g. Torrence and Compo, 1998 for details). To the right is the global wavelet power spectrum, obtained by averaging in time. Lag-1 and lag-2 autocorrelation coefficients were calculated from the original time series in order to estimate the background red noise (increasing power with decreasing frequency/wavenumber) and determine the 95% confidence



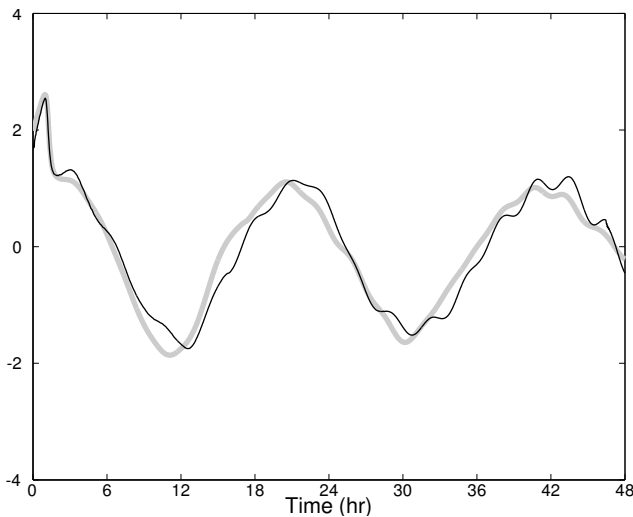
**Fig. 4.** Same as Fig. 3 except for a simulation using VIFOP optimized initial and boundary conditions.

level (Torrence and Compo, 1998). On the contour plot, the confidence level contours are shown with a thick black line while a thin dashed line is used on the global power spectrum.

The wavelet analysis of the free surface elevation at the northern point (M2,  $\phi=35.75^\circ$  N) with no optimization shows high frequency oscillations, with at least three distinct periods, all statistically significant (Fig. 3bc). The shortest periods are at around 20 and 40 min and they intermittently appear during the first 15 h of the simulation. The signal strength is barely above the red noise threshold. A stronger signal with a period of 120 min is present throughout the first 30 h of the simulations, visible as a thin, horizontal, dark region on the contour plot and a distinct peak on the global spectrum. The cone of influence for longer periods suggests the influence of edge effects on the analysis. Indeed, the overall peak at approximately 850 min is an artifact of “cyclic” data, which vanishes if the length of the record is shortened to e.g. 24 h only. One way to avoid the spurious low frequency contamination would be to remove the low-order harmonics (48, 24 h) after detrending and normalizing. Using optimized IBC and LBC, the wavelet analysis reveals an effective removal of all the significant high frequency signal (Fig. 4bc). Similar differences after applying optimization were observed also at the southern point (M1,  $\phi=31.4^\circ$  N).

Even though the timestep of the baroclinic mode was small enough to resolve the surface waves  $\mathcal{O}(10$  s), the corresponding wavelengths  $\mathcal{O}(100$  m) (Sverdrup et al., 1942) could not be resolved with the spatial resolution of the model, which was two orders of magnitude too big.

The most noticeable feature in the analysis of the  $u$  component of the velocity at M2 ( $\phi=35.75^\circ$  N) is the inertial oscillation, with a period of around 22 h (Fig. 5). The wavelet analysis is not shown, since all the signal power is contained



**Fig. 5.** Time series of a detrended and normalized  $u$  component of velocity optimized with VIFOP (thick gray line) and non-optimized (thin black line) at M2 (Fig. 1).

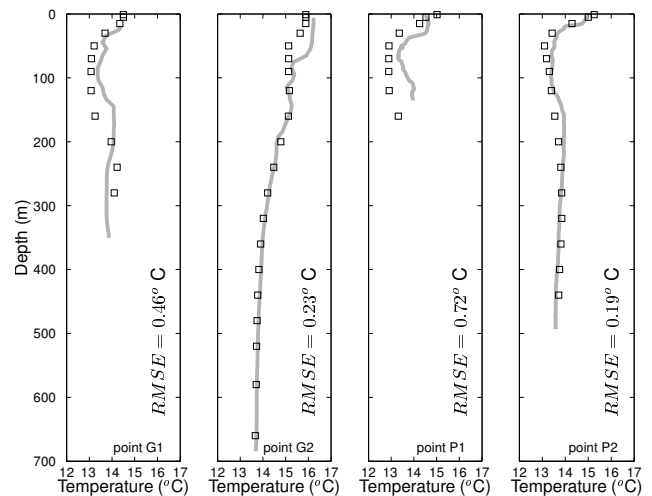
in the low-frequency oscillations in both optimized and non-optimized cases. The optimization of the initial and lateral boundary conditions using VIFOP does not affect the power spectrum. The only change is a reduction of small amplitude oscillations with an approximate period of 9 h towards the end of the 48 h period (wavelet power spectrum not shown). The inertial oscillation signal was also the most noticeable at the second point further south (M1,  $\phi=31.4^\circ$  N), although a smaller depth of the point (50 m compared to 100 m) and its proximity to shallower sea bottom ( $\sim 130$  m) influenced the dynamics.

The internal low-pass digital filter provided by VIFOP effectively removes oscillations faster than  $\sim 5$  h. If they are attributed to noise the removal is welcome. If, on the other hand, the signal is true, one should be aware of the side effects of using VIFOP. If the studied phenomena is on the time scale of 5 h or less, VIFOP at the current configuration should be used with caution.

#### 4 Model evaluation

The results of numerical model were compared to measurements, including in-situ profiles of temperature, salinity and currents, and remotely sensed sea surface temperature (SST). All the measured data was obtained directly from French Research Institute for Exploitation of the Sea, Ifremer ([http://www.ifremer.fr/mfstep/data\\_management/wp14\\_data\\_management](http://www.ifremer.fr/mfstep/data_management/wp14_data_management)), except for the currents, which were obtained from two chains moored in the strait and deep passages of the Sicily Strait.

The available in-situ measurements of temperature and salinity for the selected time period were quality controlled



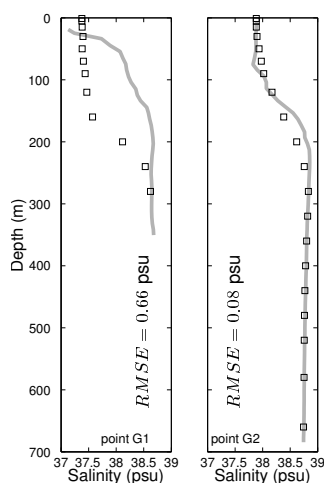
**Fig. 6.** Measured in-situ temperature profiles (thick grey line) obtained from Med-Argo gliders (G1, G2) and VOS-XBT (P1, P2), with interpolated model output (black squares). The root mean square error ( $RMSE$ ) is added for each plot. For spatial coordinates see Fig. 1.

and data points with incorrect quality flags were discarded. The original vertical coordinate of the measurements is pressure, which was transformed into depth using the hydrostatic relation. If the salinity was available (points G1 and G2), a polynomial relation was used to calculate the water density as a function of temperature and salinity (Mellor, 1991), otherwise a constant reference density was assumed ( $\rho_0=1025 \text{ kg m}^{-3}$ ), for the hydrostatic calculation.

The temperature profile was generally well captured by the model at all points (Fig. 6). The exceptions were at sharp transitions where the model tends to smooth out the profile (e.g. point G1 between 100 and 200 m). A discrepancy of  $\sim 1^\circ\text{C}$  was found also at G2 where the model underpredicted the depth of the mixed surface layer. The  $RMSE$  values are the lowest at points G2 and P2, because of the good agreement between the model and measurements at depths below 200 m.

The model derived salinity profile at G1 differs substantially from the measurements (Fig. 7). While the model predicts a  $\sim 100$  m thick top layer with constant salinity, the measured profile shows a large gradient in the surface layer, with almost constant salinity below 50 m. The depth of the eastward flowing less saline water in the top layer was overestimated by the model, but the origin of this error can be traced to the lateral boundary conditions provided by OGCM. If the model profile is retrieved at another location in proximity to G1, the salinity in the top layer starts resembling the measured values. The predicted salinity profile at G2 matches the measurements quite well.

The second type of comparison was performed using remotely sensed SST values. From the satellite data, 3-hourly

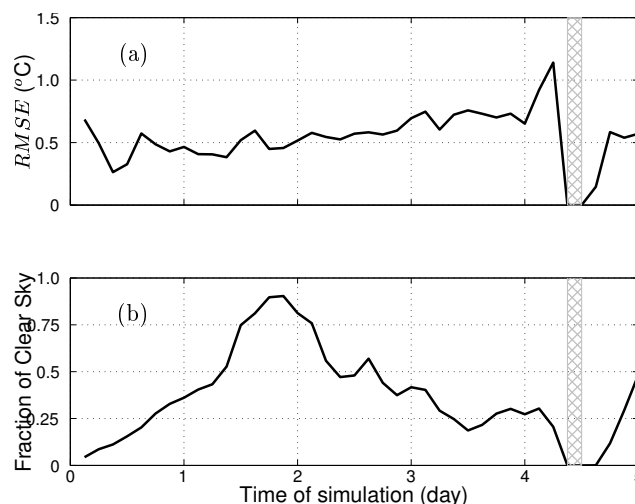


**Fig. 7.** Measured in-situ salinity profiles (thick grey line) obtained from Med-Argo gliders (G1, G2) with interpolated model output (black squares). For spatial coordinates see Fig. 1.

values ( $0.1^\circ$  spatial resolution) of SST were obtained from Ifremer ftp site ([ftp.ifremer.fr/pub/ifremer/cersat/SAFOS1/Products/ATLSST/2005](http://ftp.ifremer.fr/pub/ifremer/cersat/SAFOS1/Products/ATLSST/2005)). In order to retain high temporal resolution, the SST data used was not corrected for cloudiness, which is a regular procedure for 12 and 24-hourly products. Thus, for days with a total cloud overcast, there were no data points available. Besides temperature, each data point contains also a quality control flag and only those with excellent (5), good (4) or acceptable (3) values of confidence level were used. The third piece of information is the mean time at which the information was obtained (Brisson et al., 2001a). To proceed with the comparison, the model data was interpolated spatially onto the satellite grid and temporally to the time output determined by the time flag. An additional modification of the model postprocessing included extrapolation of the sea temperature profile to the depth of 1 cm (skin depth) by using the vertical temperature gradient between the top two model levels, which served as the SST for the comparison with the satellite measurements.

The temporal evolution of the root mean square error of SST shows a steady increase until it suddenly drops to zero, due to the total overcast with no visible pixels (Fig. 8a). The average value of  $RMSE$  for the whole period was  $0.54^\circ\text{C}$ . The percentage of visible sky is shown on (Fig. 8b), with a steady decline after day 2 associated with the passage of an atmospheric low-pressure system.

Another view of the impact of the number of pixels on the error is shown in Fig. 9. The  $RMSE$  has much smaller variability if there is approximately at least one third of the total number of pixels available. The overall value of  $RMSE$  does not change if only points above the one-third threshold are taken into account. It is interesting to note the reported error of the satellite-derived SST is around  $0.5^\circ\text{C}$  with the cloudiness threshold of 0.6 (Brisson et al., 2001b), comparable to our findings.

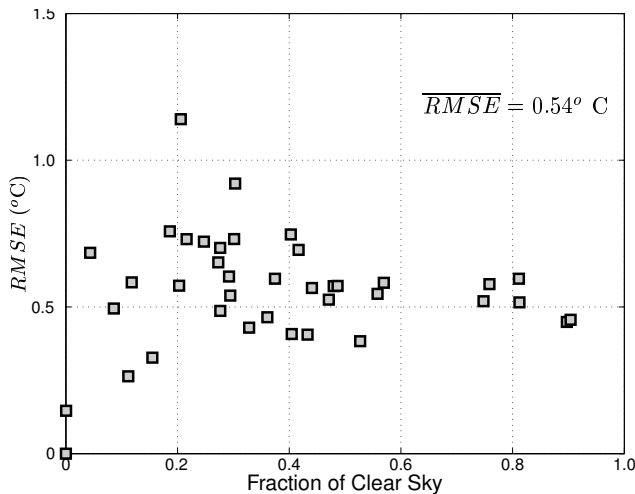


**Fig. 8.** (a) Root Mean Square Error of SCRIM predicted SST evaluated with satellite-derived values. (b) Ratio of visible pixels (clear sky) to the total number of pixels over the water over the SCRIM domain. Grey hatched area denotes the total cloud cover.

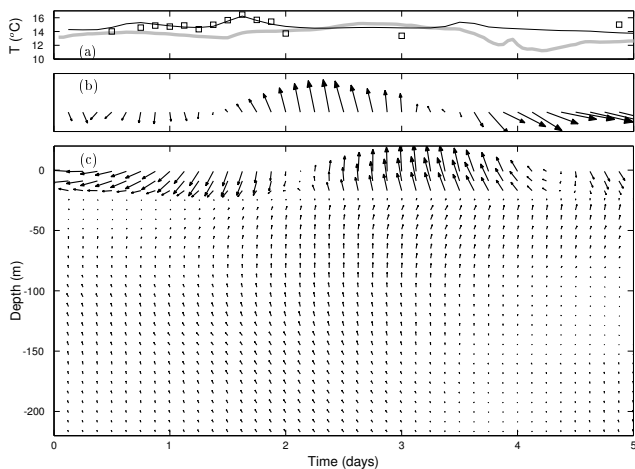
The model predicted SST has a small negative bias with the mean error of  $-0.08^\circ\text{C}$  (not shown). To assess the effects of the optimization of IBC and LBC, the same comparison procedure for SST was repeated with the model output using VIFOP. The differences regarding the non-optimized IBC and LBC were negligible for the comparison with the measurements available for the case study.

Data from two chained sea-current meters was also used for model validation (see location of points C1 and C2 on Fig. 1). Both instruments are located in the Sicilian Channel and provide hourly observations of sea-current profiles with a vertical resolution of 8 (C1) and 4 m (C2). The strongest signal in the data were lunar semi-diurnal tidal oscillations ( $M_2$ ) (Gill, 1982), which were removed by using a Doodson X0 numerical filter (Schureman, 1958). The resulting data set had a strong high frequency noise with a period of 2 h, which was subsequently filtered out. The filtered measurements were then interpolated spatially and temporally onto the model levels and times, respectively.

The main reason for the chosen period of the case study was a passage of deepening atmospheric low-pressure field over the domain. The forecasted atmospheric fields (interpolated to points C1 and C2) used for the surface boundary conditions indicate an increase in air temperature during the second and third day of the simulation, followed by a more rapid drop during the fourth day (thick gray line in Fig. 10a). The rising air temperature is accompanied by strong southerly winds, peaking at approximately  $18\text{ ms}^{-1}$  during the third day which were replaced by northwesterly winds (Fig. 10b) and temperature decrease. The evolution of the air temperature and wind resembles that of a classical warm/cold frontal passage. Due to insignificant differences



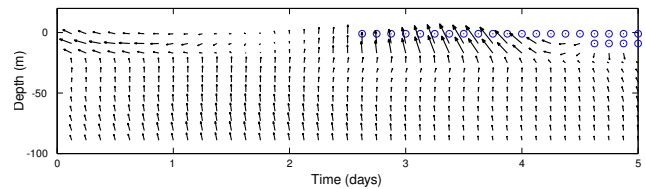
**Fig. 9.** Root Mean Square Error of the forecasted SST compared to satellite derived values as a function of cloud coverage.



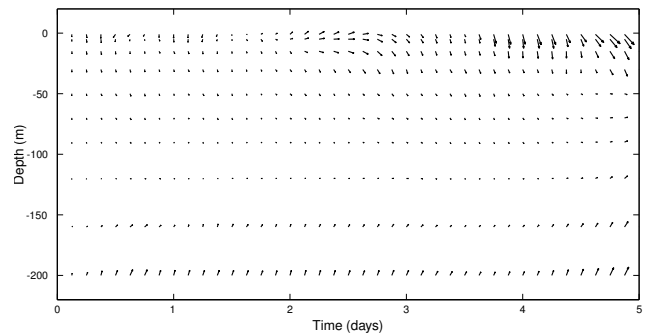
**Fig. 10.** Time series of (a) observed SST interpolated to point C2 (squares), forecasted by SCR model (black line), and air-temperature from the atmospheric model (thick grey line); (b) wind vector from the atmospheric model, interpolated to point C2. Upward pointing arrow represents southerly winds; maximum wind speed was approximately  $18 \text{ ms}^{-1}$ ; (c) sea-current profile at C2, plotted every other level and every 3 h. Upward pointing arrow represents a northward current; maximum amplitude was approximately  $1.2 \text{ ms}^{-1}$ .

in forecasted atmospheric parameters (wind and air temperature), and measured and forecasted SST between the two points (C1 and C2), data for only one point is discussed. The missing values of SST from the third to the fifth day are due to the predominantly overcast skies over the two points.

The oceanic response to the atmospheric forcing at C2 is visible comparing Figs. 10b and c. The observed sea currents at the topmost layer (depth of approximately 20 m) resemble the atmospheric forcing, albeit phase shifted by  $\sim 18 \text{ h}$ . Below 50 m, the response is much weaker and spread over



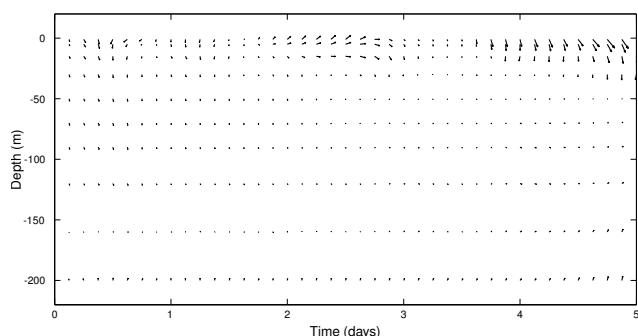
**Fig. 11.** Same as Fig. 10c, except for the point C1. Maximum velocity was approximately  $0.9 \text{ ms}^{-1}$ . Circles denote missing values.



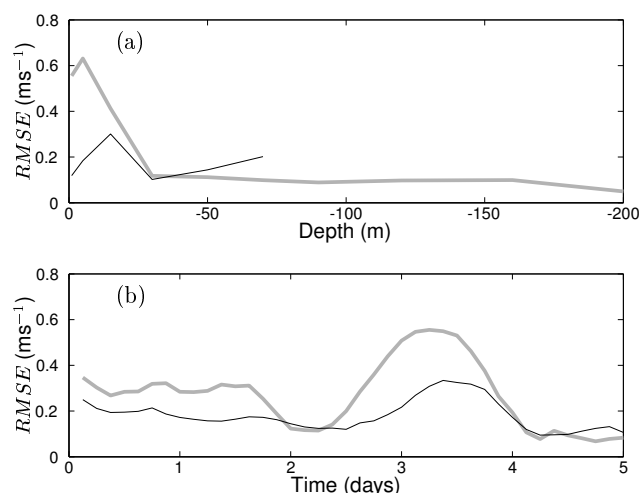
**Fig. 12.** The forecasted currents at C1. Maximum velocity was approximately  $0.58 \text{ ms}^{-1}$ . The scaling of the current vectors is kept the same as in Figs. 10 and 11 for comparison.

a longer time scale. Even though there were no direct measurements of salinity and temperature at C2, the general characteristics were deduced from points P1 and G1 due to their spatial proximity (Fig. 1), but note that the data refers only to one sample per location and not the whole time series. In the top layer the water temperature rapidly decreases (Fig. 6), while the salinity increases (Fig. 7). Below 50 m, the water is colder and saltier, and it flows generally towards NW. The decoupled flow structure is consistent with the general structure in the Sicily Channel, with the Atlantic Water (AW) at the surface and Levantine Intermediate Water (LIW) below (Manzella et al., 1988).

The differences in observed currents between C1 and C2 can be attributed to the local bathymetry. The Strait of Sicily, where both instruments are located, has a non-uniform bottom topography. From the coast of Tunisia (Cap Bon) to the southwestern tip of Sicily, there are two sills over which the currents are channeled (Manzella et al., 1990). Their position roughly corresponds to C1 and C2 according to the bathymetry used for the SCR model. A possible branching of LIW could be evidenced by measurements from both points. Note that measured currents at C1 between depths of 50 and 100 m are consistently stronger and more uniform in direction (Fig. 11), most likely due the channeling. The delay to the atmospheric forcing in the upper layer is similar to C2, but the direction-wise response is different, with currents never obtaining strong southerly component, as in C2 (Figs. 10c and 11).



**Fig. 13.** Same as Fig. 12, except for the point C2. Maximum velocity was approximately  $0.47 \text{ ms}^{-1}$ .



**Fig. 14.** Current velocity: (a) Time averaged *RMSE* as a function of depth, and (b) depth averaged *RMSE* as a function of time for C1 (thin black line) and C2 (thick gray line).

The response of the surface layer in the model is not so straightforward to interpret. During the first day of simulation when the winds are northerly, the currents are accordingly southward in both points (Figs. 12 and 13). However, when the strong forcing with southerly winds starts on the second day of simulation, lasting for approximately two days, the surface layer responds almost instantly, but with currents perpendicular to the wind. With the onset of northerly winds during the fourth day, the surface layer currents in both points respond properly (Figs. 12 and 13). Note that the forecasted currents at the depth of 200 m at C1 are again stronger than at C2. Optimizing the boundary conditions did not significantly affect the simulated currents at either of the two points.

Another possible reason for the discrepancy are largely underestimated surface geostrophic currents. The simulated surface currents are perpendicular in direction and in phase with the atmospheric forcing, while the observed currents are consistent with the along-channel pressure gradient. It

appears that the model was not able to properly reproduce sub-surface horizontal pressure gradients.

For the statistical evaluation of the currents, two skill scores were calculated. The first one is a time averaged *RMSE* at each level, the second one is a depth averaged temporal evolution of *RMSE*. We constrained our analysis only on the magnitude of the currents, since the forecasted direction is highly affected by the model resolved bathymetry.

High values of *RMSE* are found within the surface layer for both points (Fig. 14a). Note that the relatively lower values at C1 are caused by missing values in measurements (Fig. 11). The time series of depth averaged *RMSE* reveals maxima at both points during the period of strong surface currents forced by winds to which the model did not respond accordingly (Fig. 14b).

## 5 Conclusions

Unoptimized operational forecasts of SCRM contained spurious high frequency oscillations, lasting up to two days into the forecast. The oscillations mostly affected the barotropic variables, i.e. the free surface displacement and both components of barotropic velocity. The derived variables, such as the total volume transport, consisting of both baroclinic and barotropic modes, were consequently also affected by the noise. The amplitude was high enough to render the forecast useless. Baroclinic variables (components of velocity, temperature and salinity) did not contain the same type of noise.

In order to reduce the magnitude of the oscillations, an optimization tool (VIFOP) was developed within the scope of MFSTEP. Because of the nature of the noise, we attempted to optimize only the time tendency of the free surface displacement. During the optimization, the spurious numerical time tendency was minimized without affecting the physical time tendency, which is the real signal. We tested various combinations of optimized boundary conditions (i.e. initial, lateral) and concluded that the best results were achieved by optimizing all of them. Spurious initial oscillations were successfully suppressed both in duration and magnitude by using VIFOP. The time needed for the full optimization of SCRM is on the order of 8 min per day of simulation, using a Linux PC with 2 GB of memory, of which approximately 1 GB was needed for VIFOP. Thus for the whole 5 day forecast which takes around 3 hours, the overhead of 40 min was not negligible, but still acceptable compared to the total duration of the simulation.

We analyzed the effect of optimization by performing a wavelet analysis on one barotropic variable (free surface displacement) and one baroclinic variable (*u* component of velocity). The free surface oscillations with periods shorter than 4 h are effectively suppressed. The total volume transport which contained high amplitude oscillations during the first two days of forecast is vastly improved and useful for the

full 5 day duration of the simulation. The barotropic velocity field was dominated by inertial oscillations and the optimization process did not affect the main signal.

The forecast evaluation was performed by comparing the predicted variables with observations from in-situ and remotely sensed measurements. Both temperature and salinity were predicted quite well, the only discrepancy occurred in the surface layer at one station and the source of error can be traced back to the prescribed IBC/LBC. The SST prediction also had a small error, but in order to resolve high temporal resolution we used the data not corrected for cloudiness. We concluded that the error of the SST prediction stays rather constant when the fraction of visible sky is above 0.35. The optimization procedure had no impact on prediction of SST, temperature, salinity and currents.

The verification of predicted currents resulted in lower confidence in forecasted values. This was especially the case within the surface layer, where the observed currents clearly responded to the atmospheric forcing (passing cold front with a rapid shift in wind direction), while the model did not manage to create a similar response.

Ability of the SCRM to develop its own dynamics is inhibited by OGCM forcing. A more detailed comparison of various operational settings, in which forecasts will be preceded by hindcasts in order to develop high resolution features (both temporal and spatial) is left for the future.

**Acknowledgements.** This research has been realized in the framework of EC MFSTEP project (EVK3-2001-00174) with an additional support from the Marie Curie – ODASS fellowship (HPMD-CT-2001-00075 ODASS). The leading author is also grateful for the hospitality at the International Marine Center, and to his employer, University of Ljubljana, Department of Mathematics and Physics, for allowing the sabbatical leave. We are grateful to three anonymous referees for their suggestions which helped to improve the paper.

Wavelet software was provided by C. Torrence and G. Compo, and is available at <http://paos.colorado.edu/research/wavelets>.

Edited by: N. Pinardi

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