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M. Ekström, P. Eriksson, B. Rydberg, D. P. Murtagh. First Odin sub-mm retrievals in the tropical upper troposphere: humidity and cloud ice signals. *Atmospheric Chemistry and Physics*, 2007, 7 (2), pp.459-469. hal-00296127

**HAL Id: hal-00296127**

**<https://hal.science/hal-00296127>**

Submitted on 18 Jun 2008

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# First Odin sub-mm retrievals in the tropical upper troposphere: humidity and cloud ice signals

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Received: 31 May 2006 – Published in Atmos. Chem. Phys. Discuss.: 13 September 2006

Revised: 7 December 2006 – Accepted: 22 January 2007 – Published: 25 January 2007

**Abstract.** Odin-SMR is a limb-sounder operating in the 500 GHz region with the capability of performing measurements down to altitudes of about 10 km with relatively low influence of ice clouds. Until now spectra from tropospheric tangent altitudes have been disregarded due to inadequate handling of scattering. A first method to extract upper tropospheric quantities has now been developed, yielding the humidity in two layers around 200 and 130 hPa and information on cloud ice content above 200 hPa. First results are compared with in situ MOZAIC measurements and presented to give a global view of the horizontal distribution. The seasonal structures are in agreement with other satellite measurements.

The main concern for these retrievals is the calibration performance. A careful analysis indicates a systematic calibration error of about 1 K, but also a random component that differs between the two bands. The random calibration uncertainty results in retrieval errors of 10–60% depending on humidity and band. Presently this prohibits use of single retrievals, but averages can be presented with good accuracy. The fixed calibration error can largely be removed, leaving the spectroscopic uncertainties to dominate the humidity retrieval accuracy, with a worst case estimate of 30%. However, the comparison of MOZAIC data and the measurements for the 200 hPa layer shows a systematic difference of <10%. This indicates that the actual systematic error is low and gives further confidence in the capability of Odin-SMR to measure humidity in the upper tropical troposphere.

ation budget. In the upper troposphere (UT) water vapour and clouds are known to be the dominating factor for the out-going longwave radiation (OLR), see e.g. Harries (1996). The role of water in the climate system is complicated by the fact that both positive and negative feedback processes as a result of an increased surface temperature have been suggested (e.g. Ramanathan and Collins, 1991; Lindzen et al., 2001; Minschwaner and Dessler, 2004). The net impact of these processes is to a large extent still uncertain. A review of the status of research on possible feedback processes connected to water vapour and cloud formation is provided by Harries (2000).

Long-term measurements of upper tropospheric humidity (UTH) with complementary instruments have been recommended to increase our understanding of its role for the climate (Kley et al., 2000). To achieve global coverage satellite observations are the only practical option. The UT is frequently associated with cloud cover. For example, Mergenthaler et al. (1999) reported occurrence frequencies above 60% for large parts of the tropical region in 1992/1993 based on CLAES observations. Also the SAGE II 6-year climatology (Wang et al., 1996) and the later OSIRIS 3-year climatology (Bourassa et al., 2005) give mean cloud occurrences of more than 25% for the tropical UT. The frequent cloud coverage restricts the application of visible and infrared limb sounding and observations in these wavelengths are mainly performed in down-looking geometries. This results in a low vertical resolution and measurements that are most sensitive to a layer in the lower part of the UT. For example, infrared measurements in the 6.3  $\mu\text{m}$  band are sensitive to the water vapour column over a few hundred hPa around 300 hPa (e.g. Schmetz and Turpeinen, 1988; Soden and Bretherton, 1996). The microwave region exhibits a lower influence of cloud scattering, and therefore limb sounding can be performed to give improved vertical resolution in the UT. The small group of microwave limb-sounders deployed in orbit consists of UARS-MLS, Odin-SMR and Aura-MLS. The

## 1 Introduction

Water is the strongest greenhouse agent in the Earth's atmosphere and as such is an important factor for the Earth's radi-

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best existing dataset of UTH was provided by UARS-MLS (Read et al., 2001; Kley et al., 2000), partly for the reasons discussed above, while the UTH retrieval for Aura-MLS (Froidevaux et al., 2006) and Odin-SMR is in development.

Odin-SMR operates in the range 480–580 GHz, thus at sub-millimetre wavelengths, and was originally intended for upper atmospheric research. However the measurements have sensitivity down to  $\sim 10$  km, that should allow for a possibility to retrieve UTH in the tropical region. This so far unused potential of Odin-SMR should be exploited, considering the importance of UTH and the advantages of microwave limb sounding in the field. There is also potential for synergy with the MLS instruments as Odin-SMR covers partly the time gap between the UARS and Aura missions. In addition Odin-SMR would provide an independent microwave UTH data set as the instrument designs and the observed frequency bands differ between the missions.

The retrieval of UTH requires that the influence of cloud scattering can be quantified, and this can now be achieved with recent development in the area of radiative transfer calculations. This paper describes a first method (Sect. 3) to retrieve UTH together with the ice cloud signal in the tropical upper troposphere. The results (Sect. 4) will serve as an indication whether Odin radiances (Sect. 2) can be utilised to retrieve water vapour and other constituents in the upper troposphere and lower stratosphere (UTLS) region. Comparisons with in-situ measurements are used as a quality check.

## 2 Data

### 2.1 Odin-SMR

Odin is a small Swedish research satellite built in cooperation with Canada, France and Finland (Nordh et al., 2003). It was launched into a sun-synchronous orbit in February 2001. The altitude of the orbit is 600 km and the period is 1.5 h. The satellite carries a Sub-Millimetre Radiometer (SMR) which is time-shared between astronomy and aeronomy observations. For the aeronomy observations there is also a complementary instrument that combines the optical and infrared region, the Optical Spectrograph and Infrared Imaging System (OSIRIS). The design lifetime of 2 years has now been surpassed by twice as much, and the instruments are continuously providing measurements. Level 1 data are available from the later half of 2001 up to present time.

The objective of the aeronomy mission is to provide global distributions of atmospheric species in the stratosphere and the mesosphere (Murtagh et al., 2002). This is achieved by limb-sounding measurements of thermal emission from the atmosphere. The satellite performs a scanning motion to cover the atmosphere from a few kilometres above the ground up to around 100 km. In the aeronomy mission the observations are grouped into different observation modes

associated with different frequency bands. Most of the observation time of the aeronomy mission is dedicated to the stratospheric mode and data are available for approximately every third day. Data from this observation mode is used in this study, both due to the amount of data and to the fact that it has been thoroughly inspected for the operational inversions of stratospheric species (Urban et al., 2005). Stratospheric mode observations use two frequency bands around 501 and 544 GHz. The spectra cover a total of 800 MHz with a resolution of 1 MHz, either in a split-mode (501) or as a continuous frequency band (544). In the UT the signal in the window channels of these frequency bands consists mainly of far-wing absorption of the 557 GHz H<sub>2</sub>O line. These measurements have sensitivity down to about 10 km in the 501 band which is farther away from the 557 GHz H<sub>2</sub>O line centre, and down to 14 km for the frequencies around 544 GHz.

### 2.2 MOZAIC

The Measurement of Ozone and Water Vapour by Airbus In-Service Aircraft (MOZAIC) program has been utilising commercial aircraft to perform measurements of water vapour and other species in the upper troposphere since 1993 (Marengo et al., 1998). Its third phase, MOZAIC III, started in 2000 and continued to 2004, overlaps the first years of measurements performed by Odin. MOZAIC consists of automatic and regular measurements on several long range passenger airliners. Of these flights roughly 20% pass into or through the tropical region. Most of the measurements (90%) correspond to cruise altitudes around 9–12 km. This corresponds to the lowest altitude layer measured by Odin in the 501 GHz band. The instrument for water vapour is based on the Vaisala Humicap-H capacitive relative humidity sensor. The sensor is calibrated every 500 flight hour in an environmental simulation chamber and is shown to have a mean total uncertainty of  $\pm 7\%$  RH in the upper troposphere (Helten et al., 1998). The horizontal resolution at cruise altitude is about 15–50 km.

During 2001–2004 a total number of 1793 flights carrying the MOZAIC instrument entered or passed the tropical region.

## 3 Method

The retrieval method described here is similar to the methods used to derive UTH from METEOSAT (Schmetz and Turpeinen, 1988) and TOVS (Soden and Bretherton, 1996) radiances, but with the important differences that in this case the surface have no influence, clouds interact less with the retrieval and that spectra are handled individually. Temperature fields from ECMWF are used to compute the brightness temperatures,  $T_B$ , for a set of relative humidities. The result is used as a mapping function to retrieve relative humidity with respect to ice (RH<sub>i</sub>) from the measured brightness

temperatures. In cases of significant extinction the clear-sky RHi is cloud corrected using spectral information from the observation.

### 3.1 Radiative transfer

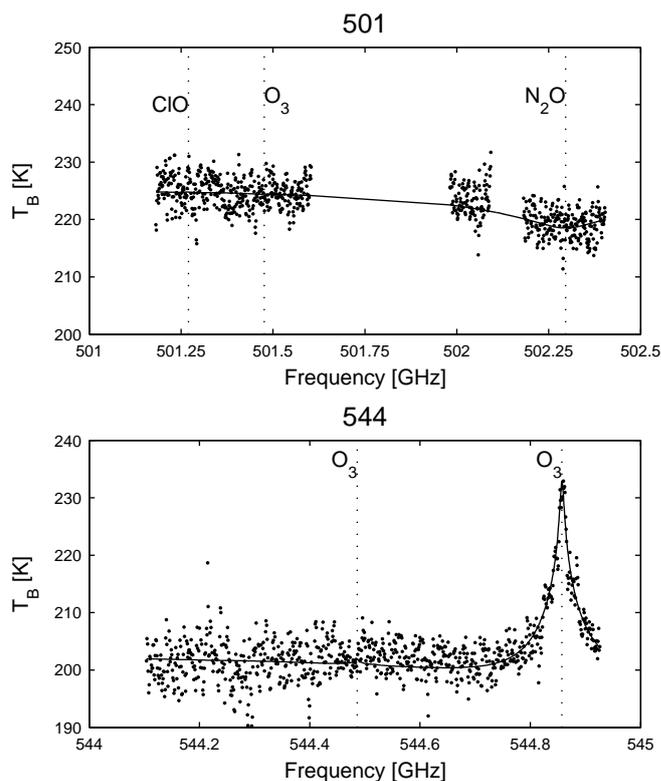
The radiative transfer calculations are performed using the Atmospheric Radiative Transfer System (ARTS) forward model, which is a general forward model for the millimetre and sub-millimetre spectral region (<http://www.sat.uni-bremen.de/arts/>). It handles atmospheric geometries from 1-D to 3-D and also any observation geometry and sensor position needed. Detailed sensor characteristics are also handled effectively with an innovative sensor matrix method (Eriksson et al., 2006). In this study the sideband filter and antenna characteristics are included in the forward simulations. Some details on the calculation of gaseous absorption are given in Sect. 4.1.2.

In ARTS the vector radiative transfer equation can be solved for all four Stokes components, at the same time as spherical geometry is considered. Together with two different modules to perform scattering calculations (Emde et al., 2004; Davies et al., 2005) it is now possible to rigorously study the impact of ice clouds on limb sounding spectra.

### 3.2 Tropospheric retrieval aspects

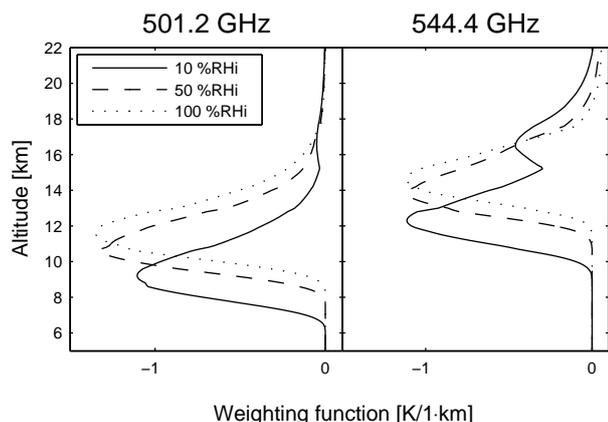
Odin spectra are saturated for the lowest tangent altitudes of the scan. The atmosphere has here very high optical thickness and with the 557 GHz H<sub>2</sub>O line outside the frequency range covered the brightness temperature spectra become spectrally flat (Fig. 1). The retrieval is restricted to absolute brightness temperatures, since no shape information is available. Consequently it is important to accurately know the temperature and the contribution of different species. At the frequencies considered in this retrieval water vapour completely dominates the absorption in most of the troposphere. The other strong absorbers in the UT are N<sub>2</sub>, O<sub>2</sub> and N<sub>2</sub>O. In the troposphere these species have very long lifetimes and their abundances are well known, hence they have no impact on the retrieval. More important, other species such as ClO and O<sub>3</sub>, which could impact with the retrieval become significant only above 16 km at these frequencies. The signal in the windows channels of the tropospheric spectra can therefore be considered to depend only on the amount of water vapour, temperature and clouds. Temperature is known with high accuracy from ECMWF analysis and with an initial assumption of clear-sky conditions a retrieval method relying on the absolute brightness temperature is feasible.

Figure 2 shows water vapour weighting functions for the two frequency bands. The weighting functions are defined as the sensitivity of the measurement to changes of H<sub>2</sub>O at different altitudes. A vertical resolution of 3–4 km can be estimated from the weighting functions. The sounding altitude of the measurements can be estimated as the weighted mean



**Fig. 1.** Example of spectra from the troposphere tangent altitudes in the stratospheric observation mode. The solid line represents simulated spectra for the retrieved humidity. The spectra shown correspond to an integration time of 0.875 s with a resolution of 1 MHz, giving a thermal noise of  $\sigma=3$  K. Vertical dotted lines represent the frequency of the stronger lines considered in the spectral band and the associated species.

altitude of these functions. Due to the shape of the weighting function the weighted mean altitude is slightly higher than the maximum of the weighting function. It can be seen that the sounding altitude is dependant on the UTH. For dry conditions the sensitivity will move toward lower altitudes and vice versa. For the 501 GHz band the sounding altitude goes from 10 km up to 13 km, and for 544 GHz the range is 14–15.5 km. Simulations have revealed that the sounding altitude throughout is found at more or less the same optical depth  $\tau$  along the line of sight, however at different  $\tau$  for the different bands. This is exploited to determine the sounding altitude of the individual measurements, where optical depth of  $\tau=0.45$  and  $\tau=0.7$  are used for 501.2 and 544.4 GHz, respectively. These values have been determined empirically by simulations of different humidity profiles, tangent altitudes and temperature profiles. The sensitivity to water vapour of this first retrieval scheme drops as the sounding altitude approaches the tropopause and as a consequence only the tropical region (30° S–30° N) is considered.



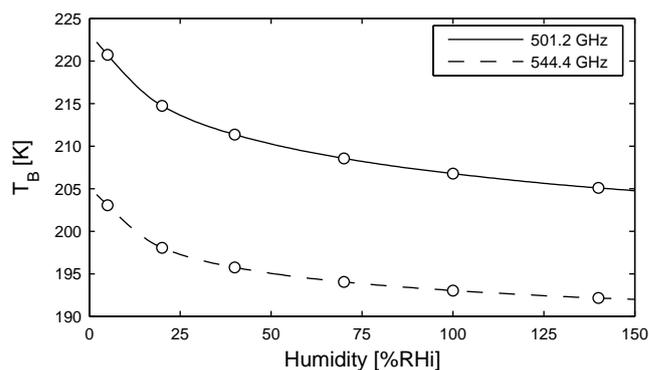
**Fig. 2.** Weighting functions for a doubling of  $\text{H}_2\text{O}$  normalized over the altitude resolution. Three cases of RH<sub>i</sub> in the troposphere are considered and calculations made for a tangent altitude of 7 km and a tropical atmosphere. The secondary maximum seen at 544.4 GHz for a dry atmosphere is caused by the transition to the  $\text{H}_2\text{O}$  climatology.

### 3.3 Clear-sky conditions

As a first step clear-sky conditions are assumed and a transfer function from  $T_{\text{B}}$  to RH<sub>i</sub> is calculated by forward simulations for different RH<sub>i</sub>. This can basically be seen as creating a look-up table between  $T_{\text{B}}$  and RH<sub>i</sub>. In the simulations the relative humidity levels are assumed constant throughout the whole troposphere with a 2 km transition layer around the tropopause that connects the tropospheric humidity with a  $\text{H}_2\text{O}$  climatology. The forward calculations are performed in a 2-D atmosphere, where temperature data are taken from ECMWF. Volume mixing ratios for other species than water vapour are taken from the Odin operational climatology.

The procedure starts by extracting the sensor position and geographic location of the targeted tangent point from the calibrated level 1 data. Temperature and mixing ratios are fetched for the time of the orbit. Odin measures in the orbital plane so the 2-D temperature field is extracted from the original 3-D grid along the satellite sub-track. In this way, the temperature gradients in the ECMWF field are preserved. With all forward model input set up, the brightness temperatures are calculated for a set of RH<sub>i</sub> using ARTS. The set of RH<sub>i</sub> is chosen to cover humidities from very dry (5 %RH<sub>i</sub>) to supersaturated (140 %RH<sub>i</sub>), and to be dense enough to allow for interpolation of brightness temperatures between the fixed values. These simulated  $T_{\text{B}}$  form together with their corresponding RH<sub>i</sub> level, the transfer function by which the measured  $T_{\text{B}}$  is mapped to RH<sub>i</sub> (Fig. 3).

From the measured spectra a band of 100 MHz is averaged, 501.2–501.3 GHz and 544.35–544.45 GHz, respectively. This reduces the noise and also any contamination of periodic patterns, which can be discerned in averaged tropospheric spectra with a period of  $\sim 100$  MHz. The averaged



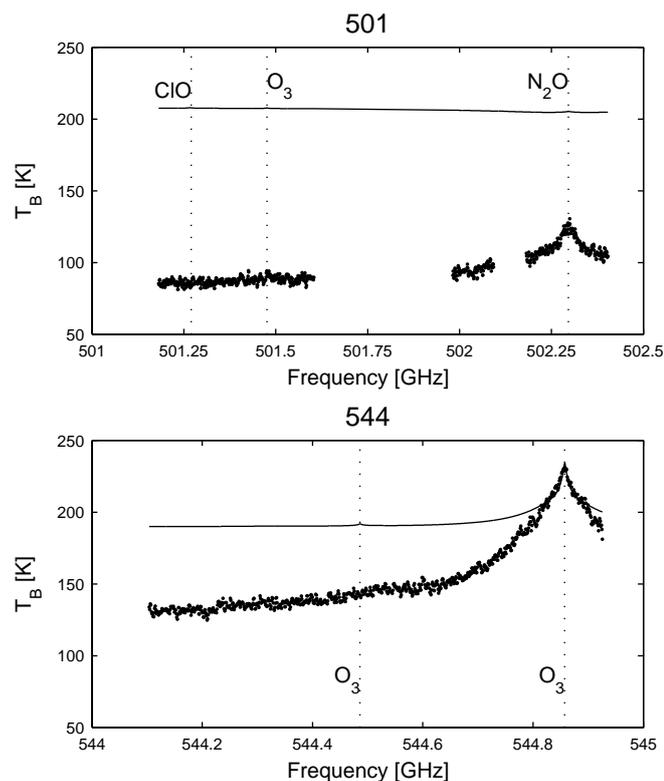
**Fig. 3.** Typical transfer functions calculated for tangent altitude of 8 km and standard tropical atmosphere. The lines show the inter- and extrapolation around the simulated RH<sub>i</sub> levels, shown as circles.

brightness temperature is mapped to a RH<sub>i</sub> using the transfer function. Cubic interpolation is used for  $T_{\text{B}}$  inside the calculated set of RH<sub>i</sub>, and for the values outside the range a linear extrapolation maintaining continuity up to the first derivative at the end points is applied.

The clear-sky retrieval procedure does not restrict how high the outgoing RH<sub>i</sub> can be. Very high values of supersaturation are for example found for spectra with ice cloud interference. These spectra are handled by the cloud correction, described below.

### 3.4 Cloud interference

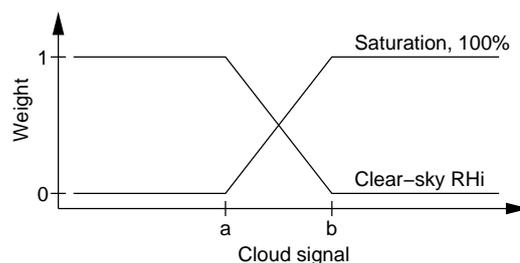
Cloud extinction consists of both absorption and scattering. Here, cases of significant cloud signal are strongly dominated by scattering (Eriksson et al., 2007). In the upper troposphere clouds act to both scatter away the up-welling radiation from the lower warmer atmosphere and to scatter into the line-of-sight the radiation from the region surrounding the ice cloud. For spectra measured at the lowest tangent altitudes the former dominates and the net effect throughout is a brightness temperature depression (Fig. 4). The emission lines in the spectra, where the atmosphere becomes optical thick at altitudes above the cloud region, are less affected by the cloud extinction. These lines are used as a reference for verifying the  $T_{\text{B}}$  depression derived from the window channels. The difference between the line and the window channels used for the clear-sky retrieval is here denoted  $\Delta T_{\text{B}}$ . In the frequency bands considered the strong emission lines are the 502.3 GHz  $\text{N}_2\text{O}$  line in the 501 GHz band, and the 544.86 GHz  $\text{O}_3$  line in the 544 GHz band. For cases of really strong scattering, i.e. thick ice clouds, the humidity is assumed to be saturated. By interpreting the  $\Delta T_{\text{B}}$  as a measure of the extinction of the cloud, it is used to weight the retrieved RH<sub>i</sub> with saturation (Fig. 5). The weights are based on statistics of the outcome of a large number of clear-sky retrievals and the corresponding  $\Delta T_{\text{B}}$ . The lowest  $\Delta T_{\text{B}}$  level for a cloud signal is determined by linear regression of the clear-sky retrieved RH<sub>i</sub> for strong



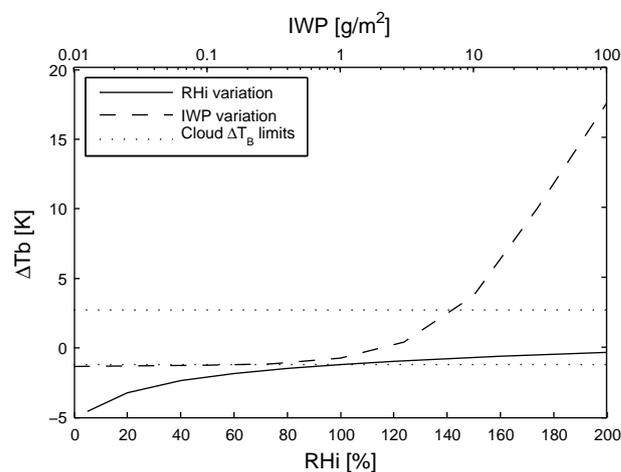
**Fig. 4.** Example of tropospheric spectra with strong ice cloud signal. Solid line represents simulated spectra for 100%RH<sub>i</sub>. Note the emission lines of N<sub>2</sub>O in 501 and O<sub>3</sub> in 544 that have become visible when the up-welling warmer radiation is redirected by cloud scattering.

ice cloud signals. The upper limit in  $\Delta T_B$  for the weighting is determined by the 95 percentile of the maximum  $\Delta T_B$  for retrieved clear-sky RH<sub>i</sub> below 180%. Thus allowing a supersaturation of 180%RH<sub>i</sub> for clear sky. The 180%RH<sub>i</sub> limit was chosen from the humidity distribution of MOZAIC measurements.

A comparison of  $\Delta T_B$  for variations in humidity and cloud ice amount (Fig. 6) show that the two quantities have different  $\Delta T_B$  signatures. For the cloud ice water path (IWP) it should be noted that Fig. 6 only serves as a schematic illustration since the cloud is crudely represented as a 1 km homogenous layer with a mono-dispersive spherical particle distribution. The figure shows that thin clouds (here IWP below  $\sim 1 \text{ g/m}^2$ ) are not detected. However, these clouds can be considered as transparent to the Odin-SMR observations, as they only have a weak impact on the  $T_B$  depression (Eriksson et al., 2007). On the other hand, thick clouds (here IWP above  $\sim 10 \text{ g/m}^2$ ) are easily detected by observing the  $\Delta T_B$ , but as the humidity information is lost 100%RH<sub>i</sub> is assumed (Fig. 5). The problematic region is where the effect of clouds and high RH<sub>i</sub> can not be distinguished in  $\Delta T_B$ . Here the measurement is treated as a combination of both clear-sky and cloud. This will introduce some uncertainty, but since the



**Fig. 5.** Principle of weighting of clear-sky retrieved humidity with saturation for cloud correction. For a cloud signal below  $a$  the retrieved RH<sub>i</sub> is equal to the clear-sky RH<sub>i</sub>, whereas for a sufficiently strong signal, greater than  $b$ , the retrieved RH<sub>i</sub> is set to 100%RH<sub>i</sub>. The  $\Delta T_B$  limits are determined from statistics, for 501 GHz  $a=-1.2 \text{ K}$  and  $b=2.7 \text{ K}$  and for 544 GHz  $a=14.6 \text{ K}$  and  $b=21.1 \text{ K}$ .



**Fig. 6.** Illustration of  $\Delta T_B$  in the 501 GHz spectra for different cases of humidity, compared to  $\Delta T_B$  due to different IWP. Scattering calculations were performed for 100%RH<sub>i</sub> with a 1 km thick cloud centered at 14 km in altitude. Cloud particles were assumed to be spherical with radius  $70 \mu\text{m}$ . The dotted lines show the limits of the cloud weighting, denoted by  $a$  and  $b$  in Fig. 5.

cloud compensation can act in both directions the resulting error is considered to be of random nature only. For supersaturated RH<sub>i</sub> that are classified as partly cloud influenced, the cloud weight is low, e.g. less than 25% for 200%RH<sub>i</sub>. Moderately supersaturated RH<sub>i</sub> will therefore be slightly underestimated.

Initial studies have shown that the extinction signatures in the spectra are closely related to the integrated cloud ice content above the sounding altitude (Rydberg, 2004). The cloud ice signal is here given as the  $T_B$  depression in the windows channel with respect to retrieved humidity. To accurately estimate the amount of ice from the cloud ice signal, assumptions on particle shape and size distribution are needed. This lies outside the scope of this paper, but is dealt with in an accompanying paper (Eriksson et al., 2007). As a rule of

thumb, the amount of ice needed to cause a 0.5 K  $T_B$  depression in the window channel is  $\sim 1 \text{ g/m}^2$  above 12.5 km. Here the cloud ice signal is presented as it is, without transformation to ice content.

### 3.5 Quality criteria

Measurements are discarded when either the spectrum is considered of bad quality or the retrieval conditions are poor. The determination of spectra quality is complicated by the fact that extinction signatures can be difficult to distinguish from errors in calibration or spectrometer read out. To test the quality of the spectra a series of tests are performed based on features of the spectra that are less affected by ice cloud scattering. The quality bounds are based on statistics of a large number of spectra, both clear-sky and with apparent cloud signature.

The retrieval method relies on the temperature gradient in the upper troposphere. For atmospheric conditions where the temperature gradient is low the uncertainty of the retrieved humidity becomes too large. To account for this, it is demanded that the  $T_B$  range of the transfer function is at least 15 K for 501 GHz and 11 K for 544 GHz. The few cases of negative retrieved RH<sub>i</sub> are rejected.

## 4 Results and discussion

### 4.1 Error estimation

Errors in the retrieval arise from uncertainties in the data and parameters used, both measured and modelled. The most important parts of the error budget are the uncertainties in calibration and therefore these are discussed separately. The error sources are divided into systematic and random. Due to the non-linearity of the transfer function used in the retrieval (Fig. 3), the sensitivity to uncertainties increases with relative humidity. The errors are here presented for the 100% RH<sub>i</sub> level to obtain conservative estimates.

#### 4.1.1 Calibration accuracy

The retrieval performance is very sensitive to the calibration of the measured radiances, partly as the retrieval method deals with absolute brightness temperatures. The accuracy of the current version of the calibration algorithm has been thoroughly investigated using in-flight observations (Ekström and Eriksson, 2006). The study was performed by comparing simulated and measured Odin-SMR spectra, exploiting the fact that measured values are completely determined by the atmospheric temperature if the atmosphere is isothermal and has a high optical thickness. Highly accurate simulations can be performed for such conditions even if there are considerable uncertainties regarding the amount of different gas species and instrument responses such as sideband ratio and pointing offsets. This is the case as the atmosphere

will act as a blackbody independently of gas variations, as long as a sufficiently high optical thickness is preserved, and that measured  $T_B$  depends on neither tangent altitude nor frequency. The only demand is that the atmospheric temperature is known accurately. A completely isothermal atmosphere is not required in practice, it is sufficient that the temperature gradient is close to zero in the altitude range corresponding to optical thicknesses in the order of one. These conditions are most closely met for Odin-SMR in the sub-arctic region for tangent altitudes below 8 km. ECMWF is used throughout as a priori temperature information for Odin-SMR, and the calibration was then assessed with respect to ECMWF temperatures, instead of in completely absolute terms. Detailed simulations confirmed that significant differences between observed and simulated  $T_B$  are directly related to calibration errors. For more details see Ekström and Eriksson (2006).

It was found that the calibration error for tropospheric spectra ( $T_B \approx 200 \text{ K}$ ) consists of a systematic and a random component. The systematic part was estimated to 1 K. The random calibration error was found to be  $\sim 2 \text{ K}$ , with probably a smaller value for the 501 band and a larger in the 544 band. These values were later found to be consistent with the distribution of UTH values obtained that lie outside the expected range (Sect. 4.2).

#### 4.1.2 Systematic errors

A systematic calibration error of 1 K was established (Sect. 4.1.1) and Odin-SMR values are corrected before the inversion with this value. This procedure should remove the main part of this error source. A rough estimate for the systematic error in the calibration study is 0.2 K and this value is used to represent the remaining retrieval error. It is worth repeating that the calibration accuracy was evaluated only with respect to ECMWF temperatures.

Other major systematic errors are associated with uncertainties in spectroscopic data used. The 557 GHz H<sub>2</sub>O line dominates the spectral region, and changes to its spectroscopic parameters have been investigated. Line intensities in the microwave region are generally assumed to be known with 2% accuracy, and this value is adopted here. The air broadened width is derived from measurements of the corresponding line in a vibrational band at  $1.39 \mu\text{m}$  (Moretti et al., 2001) and a direct error estimate is difficult. The value is 3.20 MHz/hPa. The temperature exponent for the air broadening is set to 0.64, taken from HITRAN. Kuhn (2003) gives a complete review of absorption models for the microwave region. Reported values for air broadening parameters discussed vary between 3.00–3.24 MHz/hPa and 0.60–0.69, respectively. The values of 3.00 MHz/hPa and 0.60 give the highest difference to the values assumed here, and are used to establish an error estimate. These values are both from the 1987 version of the Millimeter-wave Propagation Model, and gives a 8% lower line width at 200 K. This combined

**Table 1.** Estimated upper limit for systematic uncertainties and corresponding retrieval error. Calculated for 100 %RH<sub>i</sub> and a 8 km tangent altitude.

Error source	$\Delta$	$\Delta_{\text{RH}_i}$ [%RH <sub>i</sub> ]	
		501.2 GHz	544.4 GHz
557 GHz H <sub>2</sub> O $i_0$	2%	1.8	2.5
557 GHz H <sub>2</sub> O $a_\gamma$	8%	4.8	5.2
H <sub>2</sub> O continua	30%	3.9	0.6
N <sub>2</sub> absorption	30%	8.4	2.4
Calibration	0.2 K	3.2	6.4
Temperature	1 K	9.7	18

uncertainty is in Table 1 covered by a 8% error in the air broadened width,  $a_\gamma$ .

The calculation of nitrogen absorption follows Liebe et al. (1993). The maximum uncertainty for this term is set to 30%. This value is based on discussions in Kuhn (2003). For example, Pardo et al. (2001) found that the best fit to measured data was found with a 29% increase of the nitrogen absorption. The discrepancy found is not necessarily associated with nitrogen, it could be caused by gases and molecular interactions normally not considered in the microwave region.

The calculation of water vapour includes the addition of an empirically determined term, normally denoted as continuum absorption. This term is taken from Kuhn (2003). The uncertainty of this term at 500 GHz is not known and a value of 30% is applied here as well.

A systematic bias in the temperature fields used to calculate the transfer functions is included among the systematic errors. A bias of 1 K makes this the dominating systematic error.

The systematic retrieval errors are summarised in Table 1. Estimating the total systematic error by simple addition of the worst case errors gives an accuracy of 30%. In practice the total error is lower since the individual errors should be less than their worst case, and may also to some extent counteract each other. To estimate the total systematic error it is therefore important to compare or validate the retrieved data with other measurements. In Sect. 4.2 the data are compared to in situ MOZAIC measurements.

#### 4.1.3 Random errors

The random error sources considered are given in Table 2, together with the associated UTH errors. Uncertainties and errors are throughout assumed to be normally distributed and reported as one standard deviation.

The retrievals are based on 100 MHz wide spectral averages and the variation due to thermal noise is for Odin-SMR at this frequency resolution 0.5 K. However, the thermal noise is here overshadowed of the random part of the calibration error, estimated to be 2 K (Sect. 4.1.1). Increasing the

**Table 2.** Summary of random retrieval errors. Estimates are given at the 1 standard deviation level. Evaluated for 100 %RH<sub>i</sub>, cloud-free conditions and a 8 km tangent altitude. The random retrieval error is smaller for both lower RH<sub>i</sub> values and cases of thick ice clouds.

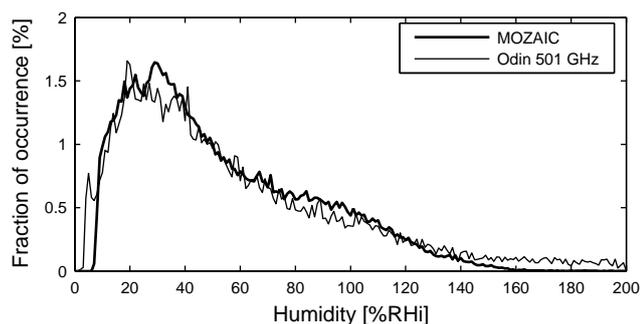
Error source	$\sigma$	$\sigma_{\text{RH}_i}$ [%RH <sub>i</sub> ]	
		501.2 GHz	544.4 GHz
Thermal noise, T <sub>B</sub>	0.5 K	8.0	16
Calibration, T <sub>B</sub>	2 K	30	50
Pointing	0.01°	6.8	3.8
Temperature	1 K	9.7	18
Constant %RH <sub>i</sub>	–	2	2
Scattering	–	20	20
Total		39	59

width of the spectral averaging would decrease the thermal noise, but not the calibration uncertainty as it is completely correlated between the spectrometer channels.

Another instrumental error source is pointing off-set. The standard processing has shown that present version of level 1 data has a pointing uncertainty of 0.01°. The potential to decrease this uncertainty still exists. All level 1 data do not include latest updates of the attitude data, and a reprocessing will be performed before next version of the standard retrieval products. A further improvement would be obtained by incorporating the pointing off-set determined in the operational processing, having a precision of  $\sim 0.0025^\circ$ . The nature of the observation gives a higher impact of a pointing off-set towards higher tangent altitudes, and this direction has been assumed when mapping the pointing uncertainty to an UTH error.

The precision in ECMWF temperatures was assumed to be 1 K ( $1\sigma$ ). It should be noted that an error in atmospheric temperature affects the retrieval error differently than an error in measured brightness temperature. This is a consequence of the fact that atmospheric temperature affects the relationship between relative and absolute humidity. The retrieval assumes a constant RH<sub>i</sub> through the troposphere, and the impact of this assumption has been determined empirically. This was done by simulating a large set of measurements with randomly selected linear varying humidity profiles in the troposphere. The result was a standard deviation of 2 %RH<sub>i</sub>.

Cloud scattering induces additional uncertainty in the UTH retrieval. A characterisation of this error would require a better knowledge on ice clouds in the upper troposphere than that provided by existing observation systems. However, assuming that the sounding region overlaps with the cloud, it is estimated that the cloud correction algorithm does not cause any substantial systematic errors. The correction can result in both too low and too high retrieved RH<sub>i</sub>, with similar probability and size in both directions. If the



**Fig. 7.** Histogram of individual measurements of UTH from Odin at 200 hPa and MOZAIC. Odin-SMR cases of strong scattering are here included following a random normal distribution of mean 100 %RH<sub>i</sub> and standard deviation 20 %RH<sub>i</sub> (see further the text). In this plot 79 659 MOZAIC measurements are used together with 17 208 Odin-SMR retrievals.

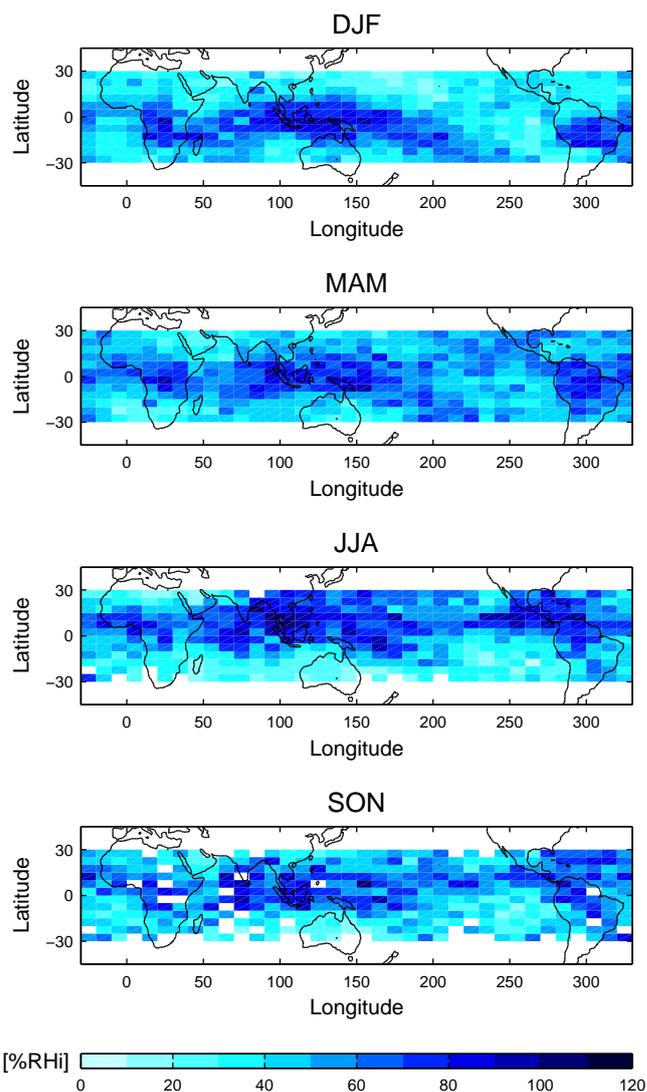
extinction is strong, it is expected that the water vapour is in equilibrium with the cloud ice particles causing the extinction and UTH is accordingly set to 100 %RH<sub>i</sub>. The retrieval can thus be said to only use a priori information. The humidity variation inside clouds at low temperatures (below  $-40^{\circ}\text{C}$ ) was reported by Ovarlez et al. (2002) to be 20 %RH<sub>i</sub>. A conservative estimate of the overall precision of the cloud correction is also 20 %RH<sub>i</sub>, and this value is applied generally.

An interesting consequence of the assumptions around the cloud correction is that the retrieval error is relatively small for cases with strong cloud extinction. The retrieval error is then equal to the a priori uncertainty for humidity inside ice clouds (20%).

Treating the random errors as independent variables, the total error at 100 %RH<sub>i</sub> can be estimated to 39 %RH<sub>i</sub> for 501.2 GHz and 59 %RH<sub>i</sub> for 544.4 GHz. The retrieval precision is hence too poor for an useful analysis of individual data points. However, these errors are expected to be of pure random nature and averaged values will have considerably smaller errors. For long term averages the retrieval error will accordingly be of primarily systematic nature.

## 4.2 Upper tropospheric humidity

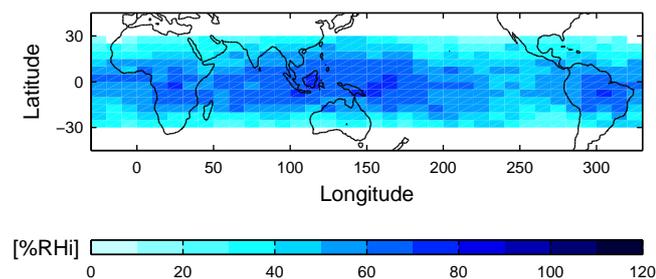
The sounding range of the 501 GHz band coincides with the altitude of the bulk of MOZAIC data, as noted in Sect. 2.2. The MOZAIC data then offer an interesting option for validation. It was found that the MOZAIC flights into the tropical region give too few coincident measurements to allow a comparison with data separated geographically. The two data sets are instead compared statistically by their distributions. All MOZAIC measurements at cruising altitude within the tropical region are included in the comparison, while for Odin-SMR only data from regions around the flight routes are considered. The comparison is shown in Fig. 7. Both distributions show the same overall pattern, but with higher occur-



**Fig. 8.** Seasonal averaged maps of the 200 hPa cloud corrected retrieved UTH for 2001–2004. The data are averaged over  $5^{\circ}$  in latitude and  $10^{\circ}$  in longitude. White areas are either outside considered latitude range or grid boxes with relatively few data points.

rence of very low ( $<10\%$  RH<sub>i</sub>) and very high ( $>150\%$  RH<sub>i</sub>) humidities for Odin-SMR, and consequently lower occurrence of intermediate RH<sub>i</sub>. However, the differences at both ends of the distribution can be explained as an effect of the established random calibration uncertainty. The averages of the two distributions differ with less than 5 %RH<sub>i</sub> and the differences are thus primarily caused by the relatively large random retrieval error for Odin-SMR. This indicates further that no larger systematic errors exist, which was discussed in Sect. 4.1.2.

As already stated, the low precision of the calibration makes the individual results less interesting and the purpose of this paper is not to investigate the data set in detail. Only some long term averages are here presented, with the aim



**Fig. 9.** Average 130 hPa cloud corrected retrieved UTH for 2001–2004. Otherwise as Fig. 8.

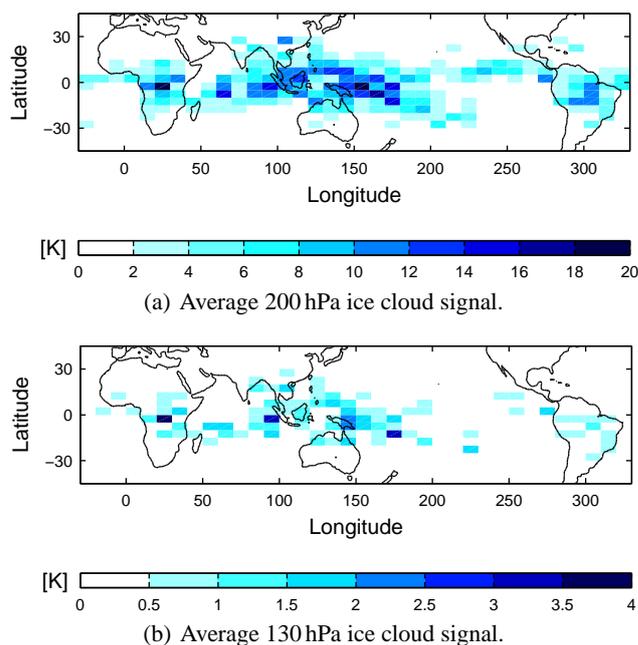
of checking if the data have the expected geographical and seasonal variations. The highest number of measurements is found for 501.2 GHz during December–February (DJF) and March–May (MAM). This as the aeronomy part of the mission is given priority during these seasons in order to study Arctic ozone depletion. This is compensated by more time given to astronomy during the other seasons, June–August (JJA) and September–November (SON). The amount of data retrieved is also limited by the quality of the measured spectra. With the current calibration scheme, the 501 GHz spectra have shown to be more stable than those at 544 GHz, and therefore more data are at hand for the lower altitude layer.

In Fig. 8 multi-year seasonal averages of the 200 hPa relative humidity are shown. The random error of the averaged humidity in the gridboxes is below 7 %RH. High relative humidities are found over Central Africa, the maritime continent and Central- and South America. These regions are the centra of the three tropical monsoon systems, regions with strong convection. The band of high humidity associated with the intertropical convergence zone is also visible over the Pacific Ocean. The seasonal patterns are consistent with what has been seen at 215 hPa by e.g. UARS-MLS (Chen et al., 1999). The regions of high humidity move with the season from the the South side of the Equator during DJF, southern summer, to the North side during JJA, northern summer.

For the 544 GHz band, only the average humidity from December 2001–August 2004 is shown here (Fig. 9). At 544 GHz the measured layer is close to the tropical tropopause. In this region high relative humidity and supersaturation is expected. The 130 hPa field measured by Odin shows also a higher mean humidity than at 200 hPa. The meridional structure is also much less pronounced, which is consistent with Sassi et al. (2001). The random error of the 544 GHz multi-year averages is less than 8 %RH.

#### 4.3 Cloud ice signals

The cloud ice signal is calculated as the brightness temperature depression from the  $T_B$  corresponding to retrieved humidity. Figure 10 shows the averaged  $T_B$  depression for the whole time period of measurements. The higher 130 hPa



**Fig. 10.** Cloud ice signals averaged for December 2001–August 2004. Gridboxes are  $5^\circ$  wide in latitude and  $10^\circ$  in longitude.

layer shows a much lower mean  $T_B$  depression than the lower 200 hPa layer. This is probably an effect of that both ice water content and particle size are decreasing with altitude. Relatively thick ice clouds are needed to give a noticeable scattering, and such clouds should mainly be found in the areas of strong convection. The cloud field at 130 hPa follows also the pattern of convection, while for humidity it was found to be more smooth (Fig. 9). Further work on ice cloud retrievals from Odin-SMR observations is described in detail by Eriksson et al. (2007), including comparisons of the results with both satellite and model data.

## 5 Conclusions

A first method to retrieve UTH from Odin radiances has been developed. Humidity is retrieved in two altitude layers around 200 and 130 hPa from the 501 and 544 GHz bands, respectively. When ice cloud interference is detected the retrieved humidities are corrected using cloud extinction signals from the spectra. Beside the UTH, the method gives also a measure of the strength of the cloud signal, that can be used to retrieve the amount of cloud ice for altitudes above 12 km. The mapping of cloud signals to ice contents is discussed separately.

The main UTH retrieval error source is the calibration uncertainty. The individual retrieval results have large random errors with the present version of the calibration algorithm. On the other hand, the systematic retrieval errors were found

to be low, which is the main concern when investigating seasonal means etc. A comparison with MOZAIC data confirmed the error characterisation performed, and a first investigation of long term averages gave the expected geographical and altitude patterns. The accuracy of such long term Odin-SMR averages should make these results a valuable tool for e.g. validation of climatological models, especially when it is considered that “all-weather” UTH is provided. The measurements are not restricted to clear-sky conditions.

A higher precision in the calibration algorithm would be very beneficial for the scientific usage of the data and further studies on how to improve the calibration scheme are therefore strongly recommended. In addition, the presented method is just a first step towards making full use of the potential of Odin-SMR in the UTLS region. The standard processing uses data effectively down to about 18 km, while here only spectra with tangent altitudes below 8 km were considered. This study has shown that Odin-SMR has a low systematic calibration error and that the impact of scattering can be handled. These are requirements for performing UTLS retrievals and more advanced retrieval schemes can now be considered. This includes the possibility of performing 2-D tomographic retrievals, utilising the overlapping of successive line of sights of the antenna around the lower turning point of the scanning motion. An initial study of such retrievals has shown a good measurement sensitivity to water vapour between 10 and 25 km.

*Acknowledgements.* The authors acknowledge M. Olberg at Onsala Space Observatory and U. Frisk at Swedish Space Corporation for their work on the Odin data and calibration. MOZAIC data were provided from the MOZAIC Data Base. This paper has been made possible through the support and previous work of the entire Odin-SMR team. The work has been funded by the Swedish National Graduate school of Space Technology and the Swedish National Space Board.

Edited by: U. Lohmann

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