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Relaxed model for the hysteresis in micromagnetism

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► **To cite this version:**

Gilles Carbou, Messoud Efendiev, Pierre Fabrie. Relaxed model for the hysteresis in micromagnetism. Proceedings of the Royal Society of Edinburgh: Section A, Mathematics, 2008, à paraître. hal-00293625

HAL Id: hal-00293625

<https://hal.science/hal-00293625>

Submitted on 7 Jul 2008

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Abstract : in this paper we study a model of ferromagnetic material with hysteresis effects. The magnetic moment behaviour is described by the non-linear Landau-Lifschitz equation with an additional term modelling the hysteresis. This term takes the form of a maximal monotone operator acting on the time derivative of the magnetic moment. In our model, it is approximated via a relaxing heat equation. For this relaxed model we prove local existence of regular solutions.

1 Introduction

The hysteresis properties of the ferromagnetic materials are a very wide domain in physics (see E. Della Torre [4]). The Preisach model describing the magnetic hysteresis is obtained by a phenomenological approach (see [9]). It is explained from the mathematical point of view by A. Visintin in [12]. Other models for rate-independent evolution in ferromagnetic materials are given in [10]. The same kind of models are used for other applications in [7] and [8].

With a physical approach, W. F. Brown developed in [2] the micromagnetism theory. The model described by Landau and Lifschitz in [6] is the following. The magnetic moment m is a unitary vector field linking the magnetic field and the magnetic induction by the relation $B = H + m$. The variations of m are described by the Landau-Lifschitz equation:

$$\frac{\partial m}{\partial t} = -m \wedge \mathcal{H}_{eff} - m \wedge (m \wedge \mathcal{H}_{eff}), \quad (1.1)$$

where the effective field is given by $\mathcal{H}_{eff} = \Delta m + h_d(m) + H_a + \Psi(m)$, and the demagnetizing field $h_d(m)$ is solution of the magnetostatic equations

$$\operatorname{div} (h_d(m) + m) = 0 \text{ and } \operatorname{curl} h_d(m) = 0, \quad (1.2)$$

where H_a is an applied magnetic field and where $\Psi(m)$ is an anisotropic term.

Micromagnetic modeling and Preisach modeling are two complementary approaches but the links between these two models are not clear. Using a two time-scales asymptotic method, J. Starynkévitch [11] gives a first answer to bring to the fore the hysteresis in Landau-Lifschitz model. We study here a model due to M. Effendiev. The hysteresis effect in Landau-Lifschitz equation is reinforced by an additional term in the effective field. This term is described with the maximal monotone operator β defined as follows

$$\beta(\xi) = \begin{cases} \frac{\xi}{|\xi|} & \text{if } \xi \neq 0, \\ B(0,1) & \text{if } \xi = 0. \end{cases} \quad (1.3)$$

In this model the effective field is given by:

$$\mathcal{H}_{eff} = \Delta m + h_d(m) + H_a + \Psi(m) - \beta\left(\frac{\partial m}{\partial t}\right). \quad (1.4)$$

The existence of regular solutions for the system (1.1)-(1.4) is open. We propose here a relaxation model for this system:

$$\left\{ \begin{array}{l} \frac{\partial m}{\partial t} = m \wedge (\Delta m + h_d(m) + H_a + \Psi(m) - v) - m \wedge (m \wedge (\Delta m + h_d(m) + H_a + \Psi(m) - v)) \\ \frac{\partial v}{\partial t} = \Delta v + \frac{1}{\varepsilon}(\beta\left(\frac{\partial m}{\partial t}\right) - v) \\ \frac{\partial m}{\partial n} = \frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega, \\ m(t=0) = m_0 \text{ and } v(t=0) = v_0 \text{ on } \Omega. \end{array} \right. \quad (1.5)$$

We prove an existence result of strong solutions for this relaxed system for $\varepsilon > 0$ fixed.

We assume that the initial data satisfies the following conditions:

$$\left\{ \begin{array}{l} m_0 \in H^2(\Omega) \text{ and } v_0 \in H^1(\Omega), \\ |m_0| = 1 \text{ on } \Omega, \\ \frac{\partial m_0}{\partial n} = 0 \text{ on } \partial\Omega. \end{array} \right. \quad (1.6)$$

For regular solutions, the equation (1.5) with initial data satisfying (1.6) is equivalent to the following system (see [3]):

$$\left\{ \begin{array}{l} \frac{\partial m}{\partial t} - \Delta m = m|\nabla m|^2 + m \wedge \Delta m + m \wedge (h_d(m) + H_a + \Psi(m) - v) \\ \quad - m \wedge (m \wedge (h_d(m) + H_a + \Psi(m) - v)) \\ \frac{\partial v}{\partial t} - \Delta v = \frac{1}{\varepsilon}(\beta\left(\frac{\partial m}{\partial t}\right) - v) \\ \frac{\partial m}{\partial n} = \frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega, \\ m(t=0) = m_0 \text{ and } v(t=0) = v_0 \text{ on } \Omega. \end{array} \right. \quad (1.7)$$

Indeed if (m, v) is a regular solution of (1.5) then the punctual norm of m is preserved and so $|m| = 1$. Then we have $\Delta|m|^2 = 0 = m \cdot \Delta m + |\nabla m|^2$.

So $m \wedge (m \wedge \Delta m) = (m \cdot \Delta m)m - |m|^2 \Delta m = -\Delta m - m|\nabla m|^2$.

In addition if (m, v) is a regular solution of (1.7) then $|m|^2$ satisfies a parabolic equation which unique solution is $|m|^2 \equiv 1$. Then the previous computation remains valid and (m, v) satisfies (1.5).

Our main result is the following theorem:

Theorem 1.1 We fix $\varepsilon > 0$. Let (m_0, v_0) satisfying (1.6). Then there exists $T^* > 0$, there exists (m, v) solution of (1.7) such that for all $T < T^*$,

$$m \in \mathcal{C}^0(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega)), \quad v \in \mathcal{C}^0(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega)).$$

In the following section we recall technical lemmas about equivalent norms in the H^p spaces, about the demagnetizing field h_d and about the maximal monotone operator β . In the last section we prove Theorem 1.1.

2 Technical lemmas

2.1 Estimates tolls

The results of this subsection are proved in [3].

Lemma 2.1 Let Ω be a bounded regular open set. There exists a constant C such that for all $u \in H^2(\Omega)$ satisfying $\frac{\partial u}{\partial \nu} = 0$ on $\partial\Omega$, we have

$$\|u\|_{H^2(\Omega)} \leq C \left(\|u\|_{L^2(\Omega)}^2 + \|\Delta u\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}}, \quad (2.1)$$

$$\|\nabla u\|_{H^1(\Omega)} \leq C \left(\|\nabla u\|_{L^2(\Omega)}^2 + \|\Delta u\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}}, \quad (2.2)$$

and for $u \in H^3(\Omega)$ such that $\frac{\partial u}{\partial \nu} = 0$ on $\partial\Omega$,

$$\|\nabla u\|_{H^2(\Omega)} \leq C \left(\|\nabla u\|_{L^2(\Omega)}^2 + \|\Delta u\|_{L^2(\Omega)}^2 + \|\nabla \Delta u\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}}. \quad (2.3)$$

Using Lemma 2.1 and the classical interpolation inequality, we rewrite Sobolev and Gagliardo-Nirenberg inequalities on the following form:

Lemma 2.2 Let Ω be a regular bounded domain of \mathbb{R}^3 . There exists a constant C such that for all $u \in H^2(\Omega)$ such that $\frac{\partial u}{\partial n} = 0$ on $\partial\Omega$,

$$\|u\|_{L^\infty(\Omega)} \leq C \left(\|u\|_{L^2(\Omega)}^2 + \|\Delta u\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}}, \quad (2.4)$$

$$\|\nabla u\|_{L^6(\Omega)} \leq C \left(\|u\|_{L^2(\Omega)}^2 + \|\Delta u\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}}, \quad (2.5)$$

$$\|\nabla u\|_{L^4(\Omega)}^2 \leq C \|u\|_{L^\infty(\Omega)} \left(\|u\|_{L^2(\Omega)}^2 + \|\Delta u\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}}, \quad (2.6)$$

and for all $u \in H^3(\Omega)$ such that $\frac{\partial u}{\partial \nu} = 0$ on $\partial\Omega$,

$$\|D^2 u\|_{L^3(\Omega)} \leq C \left(\left(\|u\|_{L^2(\Omega)}^2 + \|\Delta u\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}} + \left(\|u\|_{L^2(\Omega)}^2 + \|\Delta u\|_{L^2(\Omega)}^2 \right)^{\frac{1}{4}} \|\nabla \Delta u\|_{L^2(\Omega)}^{\frac{1}{2}} \right). \quad (2.7)$$

2.2 Demagnetizing field

We consider the operator $m \mapsto h_d(m)$ defined by (1.2). It satisfies

$$\begin{cases} h_d(m) \in L^2(\mathbb{R}^3), \\ \operatorname{curl} h_d(m) = 0 & \text{in } \mathbb{R}^3, \\ \operatorname{div} (h_d(m) + \bar{m}) = 0 & \text{in } \mathbb{R}^3, \end{cases}$$

where \bar{m} is the extension of m by zero outside $\bar{\Omega}$.

We observe that $m \mapsto -h_d(m)$ is the orthogonal projection of \bar{m} on the vector fields of gradients in $L^2(\mathbb{R}^3)$. We prove in [3] the following estimates concerning the operator h_d :

Lemma 2.3 *Let $p \in]1, +\infty[$. Then, if m belongs to $W^{1,p}(\Omega)$ (resp. $W^{2,p}(\Omega)$), the restriction of $h_d(m)$ to Ω belongs to $W^{1,p}(\Omega)$ (resp. $W^{2,p}(\Omega)$) and there exists a constant C such that*

$$\|h_d(m)\|_{L^p(\Omega)} \leq c\|m\|_{L^p(\Omega)}, \quad 1 < p < +\infty. \quad (2.8)$$

$$\|h_d(m)\|_{W^{1,p}(\Omega)} \leq C\|m\|_{W^{1,p}(\Omega)}, \quad (2.9)$$

and

$$\|h_d(m)\|_{W^{2,p}(\Omega)} \leq C\|m\|_{W^{2,p}(\Omega)}. \quad (2.10)$$

2.3 Maximal monotone operators tools

We remark that β is a maximal monotone operator. We recall usefull results proved in [1]. The first proposition is about the approximation of β by a continuous operator:

Proposition 2.1 *For $\lambda > 0$ we define β_λ by*

$$\beta_\lambda(\xi) = \begin{cases} \frac{\xi}{|\xi|} & \text{for } |\xi| \geq \lambda \\ \frac{\xi}{\lambda} & \text{for } |\xi| \leq \lambda. \end{cases}$$

Then if ξ_λ tends to ξ uniformly on $[0, T] \times \Omega$ then extracting a subsequence, $\beta_\lambda(\xi_\lambda)$ tends to $\beta(\xi)$ in L^∞ weak.*

In order to take the limit in a maximal monotone operator, we have the following lemma:

Proposition 2.2 *If A is a maximal monotone operator, if $y_n \in A(x_n)$, if $x_n \rightharpoonup x$ and $y_n \rightharpoonup y$, if $\limsup \langle x_n | y_n \rangle \leq \langle x | y \rangle$, then $y \in A(x)$ and $\langle x_n | y_n \rangle \rightarrow \langle x | y \rangle$.*

3 Proof of Theorem 1.1

3.1 First Step : Galerkin Approximation

We denote by V_n the finite dimension space built on the n first eigen-functions of $-\Delta + Id$ with domain $D(A) = \left\{ u \in H^2(\Omega), \frac{\partial u}{\partial \nu} = 0 \text{ on } \partial\Omega \right\}$, and by \mathbf{P}_n the orthogonal projection from $L^2(\Omega)$ on V_n .

We first solve the Galerkin approximation for system (1.7). We fix n and we want to build (m_n, v_n) the solution of the following approximate problem:

$$\left\{ \begin{array}{l} m_n \in \mathcal{C}^1([0, T_n[; V_n), \quad v_n \in \mathcal{C}^1([0, T_n[; V_n) \\ \frac{\partial m_n}{\partial t} - \Delta m_n = \mathbf{P}_n \left(m_n |\nabla m_n|^2 + m_n \wedge \Delta m_n + m_n \wedge (h_d(m_n) + H_a + \Psi(m_n) - v_n) \right) \\ \quad - \mathbf{P}_n \left(m_n \wedge (m_n \wedge (h_d(m_n) + H_a + \Psi(m_n) - v_n)) \right) \\ \frac{\partial v_n}{\partial t} - \frac{1}{\varepsilon} \mathbf{P}_n \left(\beta \left(\frac{\partial m_n}{\partial t} \right) \right) - \Delta v_n = -\frac{1}{\varepsilon} v_n \\ m_n(t=0) = \mathbf{P}_n(u_0), \quad v_n(t=0) = \mathbf{P}_n(v_0) \end{array} \right. \quad (3.1)$$

In order to solve this problem and to take into account the specificity of the maximal monotone operator β , we consider the approximation β_λ of β , described in the previous section, and we solve the following equation:

$$\left\{ \begin{array}{l} m_n^\lambda \in \mathcal{C}^1([0, T_n^\lambda[; V_n), \quad v_n^\lambda \in \mathcal{C}^1([0, T_n^\lambda[; V_n) \\ \frac{\partial m_n^\lambda}{\partial t} - \Delta m_n^\lambda = \mathbf{P}_n \left(m_n^\lambda |\nabla m_n^\lambda|^2 + m_n^\lambda \wedge \Delta m_n^\lambda + m_n^\lambda \wedge (h_d(m_n^\lambda) + H_a + \Psi(m_n^\lambda) - v_n^\lambda) \right) \\ \quad - \mathbf{P}_n \left(m_n^\lambda \wedge (m_n^\lambda \wedge (h_d(m_n^\lambda) + H_a + \Psi(m_n^\lambda) - v_n^\lambda)) \right) \\ \frac{\partial v_n^\lambda}{\partial t} - \frac{1}{\varepsilon} \mathbf{P}_n \left(\beta_\lambda \left(\frac{\partial m_n^\lambda}{\partial t} \right) \right) - \Delta v_n^\lambda = -\frac{1}{\varepsilon} v_n^\lambda \\ m_n^\lambda(t=0) = \mathbf{P}_n(u_0), \quad v_n^\lambda(t=0) = \mathbf{P}_n(v_0) \end{array} \right. \quad (3.2)$$

This equation can be written on the following form:

$$\left\{ \begin{array}{l} m_n^\lambda \in \mathcal{C}^1([0, T_n^\lambda[; V_n), \quad v_n^\lambda \in \mathcal{C}^1([0, T_n^\lambda[; V_n) \\ \frac{\partial m_n^\lambda}{\partial t} = F_n(m_n^\lambda, v_n^\lambda) \\ \frac{\partial v_n^\lambda}{\partial t} - \frac{1}{\varepsilon} \mathbf{P}_n \left(\beta_\lambda \left(\frac{\partial m_n^\lambda}{\partial t} \right) \right) = G(v_n^\lambda) \\ m_n^\lambda(t=0) = \mathbf{P}_n(u_0), \quad v_n^\lambda(t=0) = \mathbf{P}_n(v_0) \end{array} \right. \quad (3.3)$$

where $F_n : V_n \times V_n \rightarrow V_n$ and $G : V_n \rightarrow V_n$ are smooth. Since we can replace the second equation by

$$\frac{\partial v_n^\lambda}{\partial t} = \frac{1}{\varepsilon} \mathbf{P}_n(\beta_\lambda(F_n(m_n^\lambda, v_n^\lambda))) + G(v_n^\lambda) \quad (3.4)$$

for a fixed λ we can apply the Cauchy-Lisfchitz theorem on the finite dimensional space $V_n \times V_n$: there exists a unique solution for equation (3.3) defined on the maximal interval $[0, T_n^\lambda[$.

Since $\|\beta_\lambda(\xi)\|_{L^\infty(\Omega)} \leq 1$, there exists K depending only on n such that for all $w \in V_n$ we have:

$$\|\mathbf{P}_n(\beta_\lambda(w))\|_{V_n} \leq K.$$

Since G is linear, we can obtain from (3.4) that there exists a constant C depending on n such that for all $t \leq T_n^\lambda$ we have:

$$\|v_n^\lambda\|_{V_n} \leq Ce^{Ct}.$$

Now, there exists a constant K depending on n such that for $(u, v) \in V_n \times V_n$ we have

$$\|F_n(u, v)\|_{V_n} \leq K'_n(\|u\|_{V_n}^4 + \|v\|_{V_n}^2).$$

By comparison lemma we then obtain that there exists a time $T^n > 0$ such that for all $\lambda > 0$, $T_n^\lambda \geq T^n$, and there exists a constant K_n such that for all λ ,

$$\|m_n^\lambda\|_{L^\infty(0, T^n)} + \|v_n^\lambda\|_{L^\infty(0, T^n)} \leq K_n. \quad (3.5)$$

Using (3.5) in (3.3), we obtain a bound for $\frac{\partial m_n^\lambda}{\partial t}$ and $\frac{\partial v_n^\lambda}{\partial t}$, and derivating the first equation of (3.3) with respect to t , we obtain a bound of $\frac{\partial^2 u_n^\lambda}{\partial t^2}$. Thus there exists a constant K such that for all λ ,

$$\|m_n^\lambda\|_{L^\infty(0, T^n)} + \left\| \frac{\partial m_n^\lambda}{\partial t} \right\|_{L^\infty(0, T^n)} + \left\| \frac{\partial^2 m_n^\lambda}{\partial t^2} \right\|_{L^\infty(0, T^n)} + \|v_n^\lambda\|_{L^\infty(0, T^n)} + \left\| \frac{\partial v_n^\lambda}{\partial t} \right\|_{L^\infty(0, T^n)} \leq K_n. \quad (3.6)$$

For a fixed n we take the limit when λ tends to zero. From (3.6) we obtain that there exists m_n and v_n such that

$$m_n^\lambda \longrightarrow m_n \text{ in } L^\infty(0, T_n)$$

$$\frac{\partial m_n^\lambda}{\partial t} \longrightarrow \frac{\partial m_n}{\partial t} \text{ in } L^\infty(0, T_n)$$

$$v_n^\lambda \longrightarrow v_n \text{ in } L^\infty(0, T_n)$$

In addition using Proposition 2.1, we have that $\beta^\lambda(\frac{\partial m_n^\lambda}{\partial t})$ tends to w_n and $w_n \in \beta(\frac{\partial m_n}{\partial t})$.

Furthermore we can take the limit when λ tends to zero in Equation (3.2) and we obtain that there exist $T_n > 0$, $m_n \in \mathcal{C}^1([0, T_n]; V_n)$ and $v_n \in \mathcal{C}^1([0, T_n]; V_n)$ satisfying (3.1).

3.2 Estimates for m_n and v_n

Taking the inner product in $L^2(\Omega)$ of the first equation in (3.1) with m_n we obtain that

$$\frac{1}{2} \frac{d}{dt} \left(\|m_n\|_{L^2(\Omega)}^2 \right) + \|\nabla m_n\|_{L^2(\Omega)}^2 \leq \|m_n\|_{L^\infty(\Omega)}^2 \|\nabla m_n\|_{L^2(\Omega)}^2. \quad (3.7)$$

Taking the inner product in $L^2(\Omega)$ of the second equation in (3.1) with v_n we obtain that

$$\frac{1}{2} \frac{d}{dt} \left(\|v_n\|_{L^2(\Omega)}^2 \right) + \|\nabla v_n\|_{L^2(\Omega)}^2 \leq K(1 + \|v_n\|_{L^2(\Omega)}^2), \quad (3.8)$$

since $\left\| \beta\left(\frac{\partial m_n}{\partial t}\right) \right\|_{L^2(\Omega)} \leq K$.

We take the inner product in $L^2(\Omega)$ of the second equation in (3.1) with Δv_n . Integrating by part the right hand side, and absorbing $\|\Delta v_n\|_{L^2(\Omega)}$ using that $\left\| \beta \left(\frac{\partial m_n}{\partial t} \right) \right\|_{L^2(\Omega)} \leq K$, we obtain that

$$\frac{1}{2} \frac{d}{dt} \left(\|\nabla v_n\|_{L^2(\Omega)}^2 \right) + \|\Delta v_n\|_{L^2(\Omega)}^2 \leq K(1 + \|\nabla v_n\|_{L^2(\Omega)}^2). \quad (3.9)$$

We take the inner product in $L^2(\Omega)$ of the second equation in (3.1) with $\Delta^2 m_n$. We obtain that:

$$\frac{1}{2} \frac{d}{dt} \left(\|\Delta m_n(t)\|_{L^2(\Omega)}^2 \right) + \|\nabla \Delta m_n(t)\|_{L^2(\Omega)}^2 = I_1 + I_2 + I_3 + I_4 + I_5$$

with

$$I_1 = \int_{\Omega} \nabla (|\nabla m_n|^2 m_n) \nabla \Delta m_n dx,$$

$$I_2 = \int_{\Omega} \nabla (m_n \wedge \Delta m_n) \nabla \Delta m_n dx,$$

$$I_3 = \int_{\Omega} \nabla (m_n \wedge h_d(m_n) - m_n \wedge (m_n \wedge h_d(m_n))) \nabla \Delta m_n dx,$$

$$I_4 = \int_{\Omega} \nabla (m_n \wedge (H_a + \Psi(m_n)) - m_n \wedge (m_n \wedge (H_a + \Psi(m_n)))) \nabla \Delta m_n dx,$$

$$I_5 = \int_{\Omega} \nabla (m_n \wedge v_n - m_n \wedge (m_n \wedge v_n)) \cdot \nabla \Delta m_n.$$

We bound separately each term.

- Estimate on I_1

$$\begin{aligned} |I_1| &\leq \int_{\Omega} |\nabla m_n|^3 |\nabla \Delta m_n| dx + \int_{\Omega} |D^2 m_n| |\nabla m_n| |m_n| |\nabla \Delta m_n| dx, \\ &\leq \|\nabla m_n\|_{L^6(\Omega)}^3 \|\nabla \Delta m_n\|_{L^2(\Omega)} + \|m_n\|_{L^\infty(\Omega)} \|D^2 m_n\|_{L^3(\Omega)} \|\nabla m_n\|_{L^6(\Omega)} \|\nabla \Delta m_n\|_{L^2(\Omega)} \end{aligned}$$

hence using the Sobolev embedding and Lemmas 2.1 and 2.2 we obtain that there exists a constant K independant of n such that

$$\begin{aligned} |I_1| &\leq K \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right)^{\frac{3}{2}} \|\nabla \Delta m_n\|_{L^2(\Omega)} \\ &\quad + K \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right)^{\frac{5}{4}} \|\nabla \Delta m_n\|_{L^2(\Omega)}^{\frac{3}{2}}. \end{aligned} \quad (3.10)$$

- Estimate on I_2

By Sobolev embeddings and interpolation, we obtain that

$$\begin{aligned} |I_2| &\leq \|\nabla m_n\|_{L^6(\Omega)} \|\Delta m_n\|_{L^3(\Omega)} \|\nabla \Delta m_n\|_{L^2(\Omega)} \\ &\leq K \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right) \|\nabla \Delta m_n\|_{L^2(\Omega)} + \\ &\quad K \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right)^{\frac{3}{4}} \|\nabla \Delta m_n\|_{L^2(\Omega)}^{\frac{3}{2}}. \end{aligned} \quad (3.11)$$

- Estimate on I_3

We have

$$|I_3| \leq (1 + \|m_n\|_{L^\infty(\Omega)}) \left(\|\nabla m_n\|_{L^2(\Omega)} \|h_d(m_n)\|_{L^2(\Omega)} + \|m_n\|_{L^2(\Omega)} \|\nabla h_d(m_n)\|_{L^2(\Omega)} \right) \|\nabla \Delta m_n\|_{L^2(\Omega)},$$

and using Lemmas 2.1, 2.2 and 2.3 we obtain that there exists a constant K such that

$$|I_3| \leq K \left(1 + \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right) \right) \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right) \|\nabla \Delta m_n\|_{L^2(\Omega)}. \quad (3.12)$$

- Estimate on I_4

From the linearity of Ψ we obtain that there exists a constant K such that

$$|I_4| \leq K \left(1 + \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right) \right) \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right) \|\nabla \Delta m_n\|_{L^2(\Omega)} \quad (3.13)$$

- Estimate on I_5

We have

$$|I_5| \leq (\|m_n\|_{L^\infty(\Omega)} + \|m_n\|_{L^\infty(\Omega)}^2) (\|v_n\|_{L^2(\Omega)} + \|\nabla v_n\|_{L^2(\Omega)}) \|\nabla \Delta m_n\|_{L^2(\Omega)}$$

thus there exists a constant K such that

$$|I_5| \leq K(1 + \|m_n\|_{L^\infty(\Omega)}^2) (\|v_n\|_{L^2(\Omega)} + \|\nabla v_n\|_{L^2(\Omega)}) \|\nabla \Delta m_n\|_{L^2(\Omega)}. \quad (3.14)$$

Using Gronwall lemma with the estimates (3.9) and (3.8) we obtain that for all T there exists a constant $C(T)$ such that for all n

$$\|v_n\|_{L^\infty(0,T;H^1(\Omega))} + \|v_n\|_{L^2(0,T;H^2(\Omega))} \leq C(T) \quad (3.15)$$

Thus plugging this estimate on (3.14), adding up estimates (3.7), (3.10), (3.11), (3.12), (3.13) and (3.14), for all T there exists a constant $C(T)$ such that:

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right) + \|\nabla \Delta m_n\|_{L^2(\Omega)}^2 &\leq C(T) \|\nabla \Delta m_n\|_{L^2(\Omega)} \\ &+ C(T) \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right)^2 \|\nabla \Delta m_n\|_{L^2(\Omega)} \\ &+ K \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right)^{\frac{3}{4}} \|\nabla \Delta m_n\|_{L^2(\Omega)}^{\frac{3}{2}} \end{aligned}$$

and after absorption of $\|\nabla \Delta m_n\|_{L^2(\Omega)}$ in the right hand side term we obtain that for all T there exists a constant $C(T)$ such that

$$\frac{d}{dt} \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right) + \|\nabla \Delta m_n\|_{L^2(\Omega)}^2 \leq C(T) \left(1 + \left(\|m_n\|_{L^2(\Omega)}^2 + \|\Delta m_n\|_{L^2(\Omega)}^2 \right)^3 \right) \quad (3.16)$$

We consider the solution of the following ordinary differential equation :

$$\begin{cases} \frac{d}{dt} \xi = C(T)(1 + \xi^3) \\ \xi(0) = \left(\|m_0\|_{L^2(\Omega)}^2 + \|\Delta m_0\|_{L^2(\Omega)}^2 \right) \end{cases}$$

Since for all n , $\left(\|\mathbf{P}_n(m_0)\|_{L^2(\Omega)}^2 + \|\Delta \mathbf{P}_n(m_0)\|_{L^2(\Omega)}^2\right) \leq \left(\|m_0\|_{L^2(\Omega)}^2 + \|\Delta m_0\|_{L^2(\Omega)}^2\right)$ we obtain that for all t and for all n , we have:

$$\left(\|m_n(t)\|_{L^2(\Omega)}^2 + \|\Delta m_n(t)\|_{L^2(\Omega)}^2\right) \leq \xi(t),$$

and if we denote by T^* the lifespan of ξ , for all $T < T^*$, for all n , we have:

$$\|m_n\|_{L^\infty(0,T;H^2(\Omega))} + \|m_n\|_{L^2(0,T;H^3(\Omega))} + \|v_n\|_{L^\infty(0,T;H^1(\Omega))} + \|v_n\|_{L^2(0,T;H^2(\Omega))} \leq C(T) \quad (3.17)$$

In addition using the equation (3.1) we obtain a bound for $\frac{\partial m_n}{\partial t}$ and $\frac{\partial v_n}{\partial t}$:

$$\left\|\frac{\partial m_n}{\partial t}\right\|_{L^2(0,T;H^1(\Omega))} + \left\|\frac{\partial v_n}{\partial t}\right\|_{L^2(0,T;L^2(\Omega))} \leq C(T). \quad (3.18)$$

3.3 Limit when n tends to $+\infty$

From (3.17) we obtain a uniform bound for m_n in $L^\infty(0,T;H^2(\Omega)) \cap L^2(0,T;H^3(\Omega))$ and using the first equation of (3.1) we obtain a uniform bound for $\frac{\partial m_n}{\partial t}$ in $L^2(0,T;H^1(\Omega))$. Thus we can extract a subsequence such that

$$\begin{cases} m_n \rightharpoonup m \text{ in } L^\infty(0,T;H^2(\Omega)) \text{ weak*} \\ m_n \rightharpoonup m \text{ in } L^2(0,T;H^3(\Omega)) \text{ weak} \\ \frac{\partial m_n}{\partial t} \rightharpoonup \frac{\partial m}{\partial t} \text{ in } L^2(0,T;H^1(\Omega)) \text{ weak} \end{cases}$$

In addition, concerning v_n we have by (3.17) a uniform bound in $L^\infty(0,T;H^1(\Omega)) \cap L^2(0,T;H^2(\Omega))$ and using the second equation of (3.1) we obtain a uniform bound for $\frac{\partial v_n}{\partial t}$ in $L^2(0,T;L^2(\Omega))$. Thus we can extract a subsequence such that

$$\begin{cases} v_n \rightharpoonup v \text{ in } L^\infty(0,T;H^1(\Omega)) \text{ weak*} \\ v_n \rightharpoonup v \text{ in } L^2(0,T;H^2(\Omega)) \text{ weak} \\ \frac{\partial v_n}{\partial t} \rightharpoonup \frac{\partial v}{\partial t} \text{ in } L^2(0,T;L^2(\Omega)) \text{ weak} \end{cases}$$

Since $\mathbf{P}_n(\beta(\frac{\partial m_n}{\partial t}))$ is uniformly bounded in $L^\infty(0,T;L^2(\Omega))$ we can assume that

$$\mathbf{P}_n(\beta(\frac{\partial m_n}{\partial t})) \rightharpoonup w \text{ in } L^\infty(0,T;L^2(\Omega)) \text{ weak*}$$

Taking the limit in (3.1) we obtain that m , v and w satisfy the following system on the time interval $[0, T^*]$:

$$\begin{cases} \frac{\partial m}{\partial t} - \Delta m = m|\nabla m|^2 + m \wedge \Delta m + m \wedge (h_d(m) + H_a + \Psi(m) - v) \\ \quad - m \wedge (m \wedge (h_d(m) + H_a + \Psi(m) - v)) \\ \frac{\partial v}{\partial t} - \frac{1}{\varepsilon} w - \Delta v = -\frac{1}{\varepsilon} v \\ m(t=0) = m_0, \quad v(t=0) = v_0 \end{cases} \quad (3.19)$$

It remains to prove that $w \in \beta(\frac{\partial m}{\partial t})$. We will prove that $\frac{\partial m_n}{\partial t}$ tends to $\frac{\partial m}{\partial t}$ strongly in $L^2(0, T \times \Omega)$. Then we will apply Proposition 2.2: since $\langle \frac{\partial m_n}{\partial t} | \beta(\frac{\partial m_n}{\partial t}) \rangle \longrightarrow \langle \frac{\partial m}{\partial t} | w \rangle$, then $w \in \beta(\frac{\partial m}{\partial t})$.

We know that $\frac{\partial m_n}{\partial t}$ is bounded in $L^\infty(0, T; L^2(\Omega)) \cap L^2(0, T; H^1(\Omega))$. In order to obtain compactness for $\frac{\partial m_n}{\partial t}$, we seek a bound on $\frac{\partial^2 m_n}{\partial t^2}$. We have :

$$\frac{\partial^2 m_n}{\partial t^2} = T_1 + \dots + T_7$$

where

$$T_1 = \Delta \frac{\partial m_n}{\partial t}$$

$$T_2 = \mathbf{P}_n \left(m_n \wedge \Delta \frac{\partial m_n}{\partial t} \right)$$

$$T_3 = \mathbf{P}_n \left(\frac{\partial m_n}{\partial t} |\nabla m_n|^2 + \frac{\partial m_n}{\partial t} \wedge \Delta m_n \right)$$

$$T_4 = \mathbf{P}_n \left(2m_n \nabla m_n \nabla \frac{\partial m_n}{\partial t} \right)$$

$$T_5 = \mathbf{P}_n \left(\frac{\partial m_n}{\partial t} \wedge (H(m_n) - v_n) - \frac{\partial m_n}{\partial t} \wedge (m_n \wedge (H(m_n) - v_n)) - m_n \wedge \left(\frac{\partial m_n}{\partial t} \wedge (H(m_n) - v_n) \right) \right)$$

where $H(m_n) = h_d(m_n) + H_a + \Psi(m_n)$

$$T_6 = \mathbf{P}_n \left(m_n \wedge H\left(\frac{\partial m_n}{\partial t}\right) - m_n \wedge (m_n \wedge H\left(\frac{\partial m_n}{\partial t}\right)) \right)$$

$$T_7 = \mathbf{P}_n \left(m_n \wedge \frac{\partial v_n}{\partial t} - m_n \wedge (m_n \wedge \frac{\partial v_n}{\partial t}) \right)$$

From (3.17) and (3.18) we estimate each term on the following way:

- $\|T_1\|_{L^2(0, T; H^{-1}(\Omega))} \leq K$
- We estimate the H^{-1} norm of T_2 by duality arguments: for $\varphi \in \mathcal{C}^1([0, T[; H_0^1(\Omega)))$ we have

$$\begin{aligned} \langle \mathbf{P}_n(m_n \wedge \Delta \frac{\partial m_n}{\partial t}) | \varphi \rangle &= - \langle \Delta \frac{\partial m_n}{\partial t} | m_n \wedge \mathbf{P}_n(\varphi) \rangle \\ &= \langle \nabla \frac{\partial m_n}{\partial t} | \nabla m_n \wedge \mathbf{P}_n(\varphi) \rangle + \langle \nabla \frac{\partial m_n}{\partial t} | m_n \wedge \nabla \mathbf{P}_n(\varphi) \rangle \end{aligned}$$

We integrate in time and we obtain that

$$\begin{aligned}
\left| \int_0^T \langle T_2 | \varphi \rangle \right| &\leq \|\mathbf{P}_n(\nabla m_n \wedge \nabla \frac{\partial m_n}{\partial t})\|_{L^{\frac{4}{3}}(0,T;H^{-1}(\Omega))} \|\varphi\|_{L^4(0,T;H_0^1(\Omega))} \\
&\quad + \|\nabla m_n \wedge \nabla \frac{\partial m_n}{\partial t}\|_{L^2(0,T;L^2(\Omega))} \|\nabla \mathbf{P}_n(\varphi)\|_{L^2(0,T;L^2(\Omega))} \\
&\leq \|\nabla m_n\|_{L^4(0,T;H^{\frac{3}{2}}(\Omega))} \|\nabla \frac{\partial m_n}{\partial t}\|_{L^2(0,T;L^2(\Omega))} \|\varphi\|_{L^4(0,T;H_0^1(\Omega))} \\
&\quad + \|m_n\|_{L^\infty(0,T \times \Omega)} \|\nabla \frac{\partial m_n}{\partial t}\|_{L^2(0,T;L^2(\Omega))} \|\nabla \varphi\|_{L^2(0,T;L^2(\Omega))}
\end{aligned}$$

Hence

$$\|T_2\|_{L^{\frac{4}{3}}(0,T;H^{-1}(\Omega))} \leq K$$

- we have

$$\begin{aligned}
\left\| \frac{\partial m_n}{\partial t} |\nabla m_n|^2 \right\|_{L^2(0,T;L^2(\Omega))} &\leq \left\| \frac{\partial m_n}{\partial t} \right\|_{L^2(0,T;L^6(\Omega))} \|\nabla m_n\|_{L^\infty(0,T;L^6(\Omega))} \\
&\leq \left\| \frac{\partial m_n}{\partial t} \right\|_{L^2(0,T;H^1(\Omega))} \|\nabla m_n\|_{L^\infty(0,T;H^1(\Omega))}.
\end{aligned}$$

In addition

$$\left\| \frac{\partial m_n}{\partial t} \wedge \Delta m_n \right\|_{L^2(0,T;L^{\frac{3}{2}}(\Omega))} \leq \left\| \frac{\partial m_n}{\partial t} \right\|_{L^2(0,T;L^6(\Omega))} \|m_n\|_{L^\infty(0,T;H^2(\Omega))}.$$

Hence

$$\|T_3\|_{L^2(0,T;H^{-1}(\Omega))} \leq K.$$

- We have $\|\nabla m_n\|_{L^4(0,T;H^{\frac{3}{2}}(\Omega))} \leq K$ by interpolation theorem. Hence, since for all $p < +\infty$, $L^4(0,T;H^{\frac{3}{2}}(\Omega)) \subset L^4(0,T;L^p(\Omega))$, we have that for all $\eta > 0$,

$$\begin{aligned}
\|T_4\|_{L^{\frac{4}{3}}(0,T;H^{-1}(\Omega))} &\leq \|T_4\|_{L^{\frac{4}{3}}(0,T;L^{2-\eta}(\Omega))} \\
&\leq \|m_n\|_{L^\infty(0,T \times \Omega)} \|\nabla m_n\|_{L^4(0,T;H^{\frac{3}{2}}(\Omega))} \left\| \nabla \frac{\partial m_n}{\partial t} \right\|_{L^2(0,T;L^2(\Omega))} \\
&\leq K.
\end{aligned}$$

- $\frac{\partial m_n}{\partial t}$ is bounded in $L^2(0,T;L^6(\Omega))$, $H(m_n) - v_n$ is bounded in $L^\infty(0,T;L^6(\Omega))$, and m_n is bounded in $L^\infty(0,T \times \Omega)$. Hence T_5 is bounded in $L^2(0,T;L^3(\Omega))$, so there exists a constant K such that

$$\|T_5\|_{L^2(0,T;H^{-1}(\Omega))} \leq K.$$

- $\frac{\partial m_n}{\partial t}$ is bounded in $L^2(0,T;L^2(\Omega))$ hence by property of the operator h_d (see Proposition 2.3), since m_n is bounded in $L^\infty(0,T \times \Omega)$,

$$\|T_6\|_{L^2(0,T;L^2(\Omega))} \leq K.$$

- $\frac{\partial v_n}{\partial t}$ is bounded in $L^2(0,T;L^6(\Omega))$, therefore since m_n is bounded in $L^\infty(0,T \times \Omega)$,

$$\|T_7\|_{L^2(0,T;L^2(\Omega))} \leq K.$$

Therefore we obtain that there exists a constant K independent of n such that

$$\left\| \frac{\partial^2 m_n}{\partial t^2} \right\|_{L^{\frac{4}{3}}(0,T;H^{-1}(\Omega))} \leq K.$$

Now $\frac{\partial m_n}{\partial t}$ is bounded in $L^2(0,T;H^1(\Omega))$. So by Simon's lemma,

$$\frac{\partial m_n}{\partial t} \longrightarrow \frac{\partial m}{\partial t} \text{ in } L^2(0,T;L^2(\Omega)) \text{ strong.}$$

We have $w_n = \beta\left(\frac{\partial m_n}{\partial t}\right) \rightharpoonup w$ in $L^2(0,T;L^2(\Omega))$. So

$$\langle w_n | \frac{\partial m_n}{\partial t} \rangle \longrightarrow \langle w | \frac{\partial m}{\partial t} \rangle .$$

Hence by Proposition 2.2, $w \in \beta\left(\frac{\partial m}{\partial t}\right)$, which concludes the proof of Theorem 1.1.

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