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# A NONLINEAR ELLIPTIC PDE WITH THE INVERSE SQUARE POTENTIAL 

Louis Dupaigne

## 0. Introduction

## Statement of the problem.

This paper is concerned with the following equation :
$\left(P_{t, p}\right) \quad\left\{\begin{aligned}-\Delta u-\frac{c}{|x|^{2}} u & =u^{p}+t f & & \text { in } \Omega \\ u & >0 & & \text { in } \Omega \\ u & =0 & & \text { on } \partial \Omega\end{aligned}\right.$
Here, $\Omega$ is a smooth bounded open set of $\mathbb{R}^{n}(n \geq 3)$ containing the origin, $c>0$, $p>1, t>0$ are constants and $f \not \equiv 0$ is a smooth, bounded, nonnegative function.

We assume from now on that

$$
\begin{equation*}
0<c \leq c_{0}:=\frac{(n-2)^{2}}{4} \tag{0.1}
\end{equation*}
$$

The relevance of the constant $c_{0}$ will appear after we clarify the notion of a solution of $\left(P_{t, p}\right)$.

Three types of solution are defined thereafter : weak solutions, which provide a good setting for non-existence proofs (see Theorem 1 and Proposition 2.1), $H_{0}^{1}(\Omega)$ solutions, for which uniqueness results can be established (see Theorem 2) and strong solutions, which set the optimal regularity one can hope for (see Theorem 1 and Lemma 1.5.)

We shall say that $u \in L^{1}(\Omega)$ is a weak solution of $\left(P_{t, p}\right)$ if $u \geq 0$ a.e. and if it satisfies the two following conditions :

$$
\left\{\begin{array}{l}
\int_{\Omega}\left(\frac{u}{|x|^{2}}+u^{p}\right) \operatorname{dist}(x, \partial \Omega) d x<\infty \\
\int_{\Omega} u\left(-\Delta \phi-\frac{c}{|x|^{2}} \phi\right)=\int_{\Omega}\left(u^{p}+t f\right) \phi \quad \text { for } \phi \in C^{2}(\bar{\Omega}),\left.\phi\right|_{\partial \Omega}=0
\end{array}\right.
$$

Observe that the first condition merely ensures that the integrals in the second equation make sense.

An $\mathbf{H}_{\mathbf{0}}^{\mathbf{1}}(\Omega)$ solution is a function $u \in H_{0}^{1}(\Omega)$ such that $u \geq 0$ a.e., $u^{p} \in L^{\frac{2 n}{n+2}}(\Omega)$ and

$$
\int_{\Omega} \nabla u \nabla \phi-\int_{\Omega} \frac{c}{|x|^{2}} u \phi=\int_{\Omega}\left(u^{p}+t f\right) \phi \quad \text { for all } \phi \in H_{0}^{1}(\Omega)
$$

All integrals are well defined because of Sobolev's and Hardy's inequalities (see (0.3) for the latter.)

Finally, a strong solution $u$ is a $C^{2}(\bar{\Omega} \backslash\{0\})$ function satisfying the system of equations ( $P_{t, p}$ ) everywhere except possibly at the origin, such that for some $C>0$,

$$
0 \leq u \leq C|x|^{-a}
$$

where

$$
\begin{equation*}
a:=\frac{n-2-\sqrt{(n-2)^{2}-4 c}}{2}>0 \tag{0.2}
\end{equation*}
$$

Observe that $-a$ is the larger root of $P(X)=X(X-1)+(n-1) X+c=0$. Also define $a^{\prime}$ by

$$
-a^{\prime} \text { is the smaller root of } \mathrm{P}(\mathrm{X})
$$

## - Why are definitions (0.1), (0.2) important ?

The constant $c_{0}$ defined in (0.1) is the best constant in Hardy's inequality :

$$
\begin{equation*}
\int_{\Omega}|\nabla u|^{2} \geq c_{0} \int_{\Omega} \frac{u^{2}}{|x|^{2}} \quad \text { for all } u \in H_{0}^{1}(\Omega) \tag{0.3}
\end{equation*}
$$

Consequently, when $c<c_{0}$, the operator $-\Delta-\frac{c}{|x|^{2}}$ is coercive in $H_{0}^{1}(\Omega)$. This turns out to be crucial since Theorem 2.2 in [BG] implies that if $c>c_{0}$, there is no nonnegative $u, u \neq 0$ such that $-\Delta u-\frac{c}{|x|^{2}} u \geq 0$ and hence no solution of $\left(P_{t, p}\right)$, even in the weak sense. We arrive at the same conclusion if $c>0$ is arbitrary and the space dimension $n$ is 1 or 2 , as can be deduced from the first lines of the proof of Theorem 1.2 in [BC]. We therefore restrict to $n \geq 3$.

The constant $a$ defined in (0.2) plays a central role, even in the linear theory. Indeed, if $f \not \equiv 0$ is say, a smooth nonnegative bounded function on $\Omega$ and $u \in H_{0}^{1}(\Omega)$ is the unique solution of

$$
\left\{\begin{align*}
-\Delta u-\frac{c}{|x|^{2}} u & =f \text { in } \Omega  \tag{0.4}\\
u & =0 \text { on } \partial \Omega
\end{align*}\right.
$$

then $u(x) \geq C|x|^{-a}$ near the origin, for some $C>0$ (see Lemma 1.5.) In particular, strong solutions are the nicest one can hope for. In addition, $\psi:=|x|^{-a}$ solves $-\Delta \psi-\frac{c}{|x|^{2}} \psi=0$ in $\mathbb{R}^{n} \backslash\{0\}$.

We introduce a third constant, the exponent

$$
\begin{equation*}
p_{0}:=1+\frac{n-2+\sqrt{(n-2)^{2}-4 c}}{c} \tag{0.5}
\end{equation*}
$$

which satisfies

$$
a+2=p_{0} a
$$

Roughly speaking, if $u$ behaves like $|x|^{-a}$, then $-\Delta u-\frac{c}{|x|^{2}} u \sim|x|^{-(a+2)}$ and $u^{p} \sim|x|^{-a p}$. Hence, $p_{0}$ sets the threshold beyond which the nonlinear term produces a stronger singularity at the origin than the differential operator. In fact, we will show that for $p \geq p_{0},\left(P_{t, p}\right)$ has no solution, no matter how small $t>0$ is. See Theorem 1 for details.

This fact is somewhat surprising : one would expect that working with the map $F(u):=-\Delta u-\frac{c}{|x|^{2}} u-u^{p}$, which is such that $F^{\prime}(0)=-\Delta-\frac{c}{|x|^{2}}$ is formally bijective and $F(0)=0$, the inverse function theorem would yield solutions for $t>0$ sufficiently small. Such an argument fails because there is no functional setting in which it may be applied. See section 7 of $[\mathrm{BV}]$ or the introduction of $[\mathrm{BC}]$ for a similar situation.

Another interesting property of $p_{0}$ is its variation as $c$ decreases from $c=c_{0}$ to $c=0:$ when $c=c_{0}, p_{0}=\frac{n+2}{n-2}$ is the Sobolev exponent whereas when $c \rightarrow 0$, $p_{0} \rightarrow \infty$. This is natural in view of the case $c=0$, for which $p>1$ can be chosen arbitrarily (see e.g. $[\mathrm{D}],[\mathrm{BCMR}],[\mathrm{CR}]$. )

- How do strong, $\mathbf{H}_{0}^{1}(\Omega)$ and weak solutions relate ?


## Proposition 0.1.

Suppose (0.1) holds and recall (0.2), (0.5). Suppose also that $1<p<p_{0}$.

- If $u$ is a strong solution of $\left(P_{t, p}\right)$, then $u$ is an $H_{0}^{1}(\Omega)$ solution of $\left(P_{t, p}\right)$.
- If $u$ is an $H_{0}^{1}(\Omega)$ solution of $\left(P_{t, p}\right)$, then $u$ is a weak solution of $\left(P_{t, p}\right)$.
- If $u$ is a weak solution of $\left(P_{t, p}\right)$ and $0 \leq u \leq C|x|^{-a}$ then $u$ is an $H_{0}^{1}(\Omega)$ solution of $\left(P_{t, p}\right)$.
- If $u$ is an $H_{0}^{1}(\Omega)$ solution of $\left(P_{t, p}\right)$ and $0 \leq u \leq C|x|^{-a}$ then $u$ is a strong solution of $\left(P_{t, p}\right)$.

This will be proved in Section 1.
Remark 0.1. In section 5, we provide examples of both strong and $H_{0}^{1}(\Omega)$ solutions. We do not know however if there exist weak solutions that are not $H_{0}^{1}(\Omega)$.

With these definitions in mind, we investigate the existence, uniqueness and regularity of solutions of $\left(P_{t, p}\right)$ :

## Main results.

## Theorem 1.

Suppose (0.1) holds and recall (0.5).

- If $1<p<p_{0}$, there exists $t_{0}>0$ depending on $n, c, p, f$ such that
if $t<t_{0}$ then $\left(P_{t, p}\right)$ has a minimal strong solution,
if $t=t_{0}$ then $\left(P_{t, p}\right)$ has a minimal weak solution,
if $t>t_{0}$ then $\left(P_{t, p}\right)$ has no solution, even in the weak sense and there is complete blow-up.
- If $p \geq p_{0}$ then, for any $t>0$,
$\left(P_{t, p}\right)$ has no solution, even in the weak sense, and there is complete blow-up.

This result requires the following definition :
Definition 0.1. Let $\left\{a_{n}(x)\right\}$ and $\left\{g_{n}(u)\right\}$ be increasing sequences of bounded smooth functions converging pointwise respectively to $\frac{c}{|x|^{2}}$ and $u \rightarrow u^{p}$ and let $u_{n}$ be the minimal nonnegative solution of
$\left(P_{n}\right)$

$$
\left\{\begin{aligned}
-\Delta \underline{u_{n}}-a_{n} \underline{u_{n}} & =g_{n}\left(\underline{u_{n}}\right)+t f & & \text { in } \Omega \\
\underline{u_{n}} & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

We say that there is complete blow-up in $\left(P_{t, p}\right)$ if, given any such $\left\{a_{n}(x)\right\}$, $\left\{g_{n}(u)\right\}$ and $\left\{\underline{u_{n}}\right\}$,

$$
\frac{u_{n}(x)}{\delta(x)} \rightarrow+\infty \text { uniformly on } \Omega
$$

where $\delta(x):=\operatorname{dist}(x, \partial \Omega)$.

## Theorem 2.

Suppose (0.1) holds and $1<p<p_{0}, 0<t<t_{0}$. Then if $u_{t}$ denotes the minimal strong solution of $\left(P_{t, p}\right)$,

- $u_{t}$ is stable
- $u_{t}$ is the only stable $H_{0}^{1}(\Omega)$ solution of $\left(P_{t, p}\right)$

If $u_{t_{0}}$ denotes the minimal weak solution of $\left(P_{t_{0}, p}\right)$ and $0<c<c_{0}$ and

- if $u_{t_{0}}$ solves the problem in the strong sense then $\lambda_{1}\left(u_{t_{0}}\right)=0$

Stability is defined as follows :
Definition 0.2. We say that $u$ is stable if the generalized first eigenvalue $\lambda_{1}(u)$ of the linearized operator of equation $\left(P_{t, p}\right)$ is positive, i.e., if

$$
\lambda_{1}(u):=\inf \left\{J(\phi): \phi \in C_{c}^{\infty}(\Omega) \backslash\{0\}\right\}>0
$$

where

$$
J(\phi)=\frac{\int_{\Omega}|\nabla \phi|^{2}-\int_{\Omega} \frac{c}{|x|^{2}} \phi^{2}-\int_{\Omega} p u^{p-1} \phi^{2}}{\int_{\Omega} \phi^{2}}
$$

The proof of Theorem 1 is presented in sections 2 and 3 , whereas Theorem 2 is proved in section 4.

In section 5 , we study the extremal case $t=t_{0}$ and provide examples of two distinct behaviors of the extremal solution of $\left(P_{t_{0}, p}\right)$.

Finally, in section 6, proofs of all previously announced results pertaining to the case $c=c_{0}$ are given.

## Notation and further definitions.

Dealing with linear equations of the form (0.4) with $f \in L^{1}(\Omega, \operatorname{dist}(x, \partial \Omega) d x)$, a weak solution $u$ is one that satisfies the equation $\int_{\Omega} u\left(-\Delta \phi-\frac{c}{|x|^{2}} \phi\right)=\int_{\Omega} f \phi$ with the integrability condition $\int_{\Omega} \frac{|u|}{|x|^{2}}<\infty$. Strong solutions are defined as in the nonlinear case.

Of course, Proposition 0.1 need not be true in this setting.
Sometimes we shall refer to inequalities holding in the weak sense or talk about (weak) supersolutions. This means that we integrate the equation with nonnegative test functions.

For example, $-\Delta u-\frac{c}{|x|^{2}} u \geq f$ holds in the weak sense,
given $f \in L^{1}(\Omega, \operatorname{dist}(x, \partial \Omega) d x)$, if $\frac{u}{|x|^{2}} \in L^{1}(\Omega)$ and if

$$
\int_{\Omega} u\left(-\Delta \phi-\frac{c}{|x|^{2}} \phi\right) \geq \int_{\Omega} f \phi \quad \text { for all } \phi \in C^{2}(\bar{\Omega}) \text { with } \phi \geq 0 \text { and }\left.\phi\right|_{\partial \Omega}=0
$$

The following $L^{q}$ weighted spaces will be used in the sequel :

$$
\begin{aligned}
& L_{\delta}^{q}=L^{q}(\Omega, \delta(x) d x) \\
& L_{m}^{q}=L^{q}\left(\Omega,|x|^{m} d x\right) \\
& L_{m, \delta}^{q}=L^{q}\left(\Omega,|x|^{m} \delta(x) d x\right) \text { and } \\
& L_{m}^{\infty}=\left\{u: u \cdot|x|^{-m} \in L^{\infty}(\Omega)\right\}
\end{aligned}
$$

where $1 \leq q<\infty, \delta(x)=\operatorname{dist}(x, \partial \Omega)$ and $m \in \mathbb{R}$.
Also, for $\rho>0, B_{\rho}$ denotes the open ball of radius $\rho$ centered at the origin. The letter $C$ denotes a generic positive constant.

## 1. Preliminary : Linear theory

We construct here a few basic tools to be used later on and start out with the $L^{2}$ theory.

## Lemma 1.1.

Suppose $0<c<c_{0}$ and let $f \in H^{-1}(\Omega)$. There exists a unique $u \in H_{0}^{1}(\Omega)$, weak solution of

$$
\left\{\begin{align*}
-\Delta u-\frac{c}{|x|^{2}} u & =f \text { in } \Omega  \tag{1.1}\\
u & =0 \text { on } \partial \Omega
\end{align*}\right.
$$

Furthermore,

$$
\begin{equation*}
f \geq 0 \text { in the sense of distributions } \Rightarrow u \geq 0 \text { a.e. } \tag{1.3}
\end{equation*}
$$

Proof. Hardy's inequality (0.3) implies that $-\Delta-\frac{c}{|x|^{2}}$ is coercive in $H_{0}^{1}(\Omega)$. (1.2) follows from Lax-Milgram's lemma. Observe that, using approximation in $H_{0}^{1}(\Omega)$ by smooth functions and integration by parts in $\bar{\Omega} \backslash B_{\epsilon}$ with $\epsilon \rightarrow 0$, our definition of a weak solution and that of Lax-Milgram's lemma coincide in this setting.

For $u \in H_{0}^{1}(\Omega)$, it is well known that $u^{-} \in H_{0}^{1}(\Omega)$. Testing the variational formulation of (1.1) against $u^{-}$yields (1.3).

Next, we consider the $L^{q}$ theory and restrict ourselves to the radial case.

## Lemma 1.2.

Suppose $0<c<c_{0}$ (with $c_{0}$ defined in (0.1)) and recall (0.2).
Let $q \in\left(\frac{n}{n-a}, \frac{n}{2+a}\right), E=W^{2, q}\left(B_{1}\right) \cap W_{0}^{1, q}\left(B_{1}\right) \cap\left\{u: \frac{u}{|x|^{2}} \in L^{q}\left(B_{1}\right)\right\}$.
For any radial $f \in L^{q}\left(B_{1}\right)$, there exists a unique radial weak solution $u \in E$ of

$$
\left\{\begin{align*}
-\Delta u-\frac{c}{|x|^{2}} u & =f \text { in } B_{1}  \tag{1.4}\\
u & =0 \text { on } \partial B_{1}
\end{align*}\right.
$$

Furthermore,

$$
\begin{gather*}
\|u\|_{E} \leq C\|f\|_{L^{q}}  \tag{1.5}\\
f \geq 0 \text { a.e. } \Rightarrow u \geq 0 \text { a.e. } \tag{1.5'}
\end{gather*}
$$

Remark 1.2.

- It can be shown that $u \in W^{2, q} \cap W_{0}^{1, q} \Rightarrow \frac{u}{\mid x^{2}} \in L^{q}$ for $1<q<n / 2$, so that the definition of $E$ can be slightly simplified.
- Observe that the interval $\left(\frac{n}{n-a}, \frac{n}{2+a}\right)$ is nonempty if and only if $c<c_{0}$.
- The restrictions on the range of $q$ are optimal. If $q \leq \frac{n}{n-a}$, uniqueness is lost (see Remark 1.4), whereas if the lemma were to hold for some $q \geq \frac{n}{2+a}$, one could construct solutions of $\left(P_{t, p}\right)$ for some $p, p \geq p_{0}$ by means of the inverse function theorem, contradicting Theorem 1 (see the methods of Proposition 4.1 .)
- It would be natural to extend Lemma 1.2 to the nonradial case. The problem remains open.

Proof.
Uniqueness will follow from the maximum principle (Lemma 1.4) proved in this section, provided we can show that $E \subset L_{-a-2}^{1}$.

If $u \in E, \frac{u}{|x|^{2}} \in L^{q}$ and using Hölder's inequality, $u \in L_{-a-2}^{1}$ if $|x|^{-a} \in L^{\frac{q}{q-1}}$, which is equivalent to asking $q>\frac{n}{n-a}$.

For existence, we suppose (without loss of generality in view of estimate 1.5) that $f \in C_{c}^{\infty}(0,1), f \geq 0$ and define

$$
u(r):=\Phi(f)(r)=\frac{r^{-a}}{\alpha} \int_{0}^{1} f(s) \cdot s^{\frac{n+\alpha}{2}}\left[\max (s, r)^{-\alpha}-1\right] d s
$$

where $\alpha=\sqrt{(n-2)^{2}-4 c}, r \in(0,1)$.
(1.5') follows from the definition of $u$.

Since $f$ is supported away from the origin, it is quite clear that $\frac{u}{r^{-a}}$ is smooth everywhere on $[0,1]$ so that $|u| \leq C r^{-a}$ and $\left|u^{\prime}\right| \leq C r^{-a-1}$. Also, $u(1)=0$. Differentiating $u$, we get

$$
\begin{equation*}
-u^{\prime \prime}-\frac{n-1}{r} u^{\prime}-\frac{c}{r^{2}} u=f \tag{1.6}
\end{equation*}
$$

This equality holds for every $r \neq 0$ and also in the weak sense, using integration by parts in $B_{1} \backslash B_{\epsilon}$ with $\epsilon \rightarrow 0$ and the above estimate on $u$ and $u^{\prime}$.

So, we just have to prove (1.5), which we shall do using Hardy-inequality-type arguments. Using the definition of $u$, we see that

$$
\begin{aligned}
0 \leq C \frac{u}{r^{2}} & \leq r^{-\left(1+\frac{n+\alpha}{2}\right)} \int_{0}^{r} f(s) \cdot s^{\frac{n+\alpha}{2}} d s+r^{-(1+n / 2)+\alpha / 2} \int_{r}^{1} f(s) s^{\frac{n-\alpha}{2}} d s \\
& \equiv A \quad+B
\end{aligned}
$$

Letting $g(s)=f(s) s^{\frac{n+\alpha}{2}}$ for $0 \leq s \leq 1$ and $G(r)=\int_{0}^{r} g(s) d s$ for $0 \leq r \leq 1$, integration by parts yields

$$
\begin{aligned}
& I:=\int_{0}^{1} r^{-\left(1+\frac{n+\alpha}{2}\right) q} G^{q}(r) r^{n-1} d r= \\
& \begin{aligned}
\frac{1}{n-\left(1+\frac{n+\alpha}{2}\right) q} G^{q}(1)-\frac{q}{n-\left(1+\frac{n+\alpha}{2}\right) q} & \int_{0}^{1} r^{n-\left(1+\frac{n+\alpha}{2}\right) q} G^{q-1}(r) g(r) d r \\
& \leq C \int_{0}^{1} r^{n-\left(1+\frac{n+\alpha}{2}\right) q} G^{q-1}(r) g(r) d r
\end{aligned}
\end{aligned}
$$

The last inequality results from the fact that when $q>\frac{n}{n-a}, \frac{1}{n-\left(1+\frac{n+\alpha}{2}\right) q}<0$.
Applying Hölder,

$$
I \leq I^{\frac{q-1}{q}}\left(\int_{0}^{1} r^{\gamma} g^{q}(r) d r\right)^{1 / q}
$$

where $\gamma=q\left(n-\left(1+\frac{n+\alpha}{2}\right)\right)$. But $r^{\gamma} g^{q}(r)=r^{q(n-1)} f^{q}(r) \leq r^{(n-1)} f^{q}(r)$ so

$$
\begin{equation*}
\left(\int_{B_{1}} A^{q}\right)^{1 / q}=C \cdot I^{1 / q} \leq C\|f\|_{L^{q}} \tag{1.7}
\end{equation*}
$$

To bound $B$, we introduce similarily $h(s)=s^{\frac{n-\alpha}{2}} f(s)$ and $H(r)=\int_{r}^{1} h(s) d s$. Then, since $H(1)=0$ and $(-a-2) q+n>0$, integration by parts yields

$$
\begin{aligned}
\int_{0}^{1} r^{-(a+2) q} H^{q}(r) r^{n-1} d r & \leq C \int_{0}^{1} r^{-(a+2) q+n} H^{q-1}(r) h(r) d r \\
& \leq C\left(\int_{0}^{1} r^{-(a+2) q+n-1} H^{q}(r) d r\right)^{\frac{q-1}{q}}\left(\int_{0}^{1} r^{\gamma} h^{q}(r) d r\right)^{\frac{1}{q}}
\end{aligned}
$$

where $\gamma=n-1-q(n+\alpha) / 2$. Now, $r^{\gamma} h^{q}(r)=r^{n-1} f^{q}(r)$ and it follows that

$$
\begin{equation*}
\left(\int_{B_{1}} B^{q}\right)^{1 / q} \leq C\|f\|_{L^{q}} \tag{1.8}
\end{equation*}
$$

Combining (1.7) and (1.8) gives $\left\|u / r^{2}\right\|_{L^{q}} \leq C\|f\|_{L^{q}}$.
To get (1.5), using equation (1.6), it suffices to show that $u^{\prime} / r \in L^{q}$. From the definition of $u=\Phi(f)$, we see that

$$
u^{\prime} / r=-a \cdot u / r^{2}-\alpha A
$$

and the estimate follows from our previous analysis.

Existence or uniqueness hold in other functional spaces, as the following two lemmas show :

## Lemma 1.3.

Recall (0.1), (0.2), (0.2').
Let $f$ be such that $\int_{\Omega}|f| \cdot|x|^{-a} \operatorname{dist}(x, \partial \Omega) d x<\infty$.
There exists at least one weak solution $u$ with $u \cdot|x|^{-2} \in L^{1}(\Omega)$, of

$$
\left\{\begin{align*}
-\Delta u-\frac{c}{|x|^{2}} u & =f \text { in } \Omega  \tag{1.9}\\
u & =0 \text { on } \partial \Omega
\end{align*}\right.
$$

Furthermore,

$$
\begin{align*}
&\|u\|_{L_{-2}^{1}} \leq C\|f\|_{L_{-a, \delta}^{1}}  \tag{1.10}\\
&\|u\|_{L_{-a}^{\infty}} \leq C\|f\|_{L^{\infty}}  \tag{1.11}\\
&\|u\|_{L_{-b}^{\infty}} \leq C\|f\|_{L_{-b-2}^{\infty}} \quad \text { for } a<b<a^{\prime}
\end{align*}
$$

Proof. (Case $0<c<c_{0}$ )

We assume, without loss of generality, that $f \geq 0$ (for the general case, apply the result to the positive and negative parts of $f$ ).

Let $f_{k}=\min (f, k)$ for $k \in \mathbb{N}$. Then, $f_{k} \nearrow f$ in $L_{-a, \delta}^{1}$.
By Lemma 1.1, there exists $u_{k}$, unique solution in $H_{0}^{1}(\Omega)$ of (1.9) with $f_{k}$ in place of $f$. Clearly, $\left\{u_{k}\right\}$ is monotone increasing.

Let $\zeta_{0}$ be the $H_{0}^{1}(\Omega)$ solution of

$$
\left\{\begin{align*}
-\Delta \zeta_{0}-\frac{c}{|x|^{2}} \zeta_{0} & =1 \text { in } \Omega  \tag{1.12}\\
\zeta_{0} & =0 \text { on } \partial \Omega
\end{align*}\right.
$$

When $\Omega=B_{1}, \zeta_{0}=\zeta_{0}^{1}:=C\left(|x|^{-a}-|x|^{2}\right)$, for some $C>0$. Otherwise, $\Omega \subset B_{R}$ for some $R>0$ and $C \cdot \zeta_{0}^{1}(x / R)$ is a supersolution of problem (1.12), for some $C>0$. So,

$$
\begin{equation*}
0 \leq \zeta_{0} \leq C|x|^{-a} \delta(x) \quad \text { in } \Omega \tag{1.13}
\end{equation*}
$$

Since $u_{k}$ and $\zeta_{0} \in H_{0}^{1}(\Omega)$, they are valid test functions in their respective equations and

$$
\int_{\Omega} \nabla u_{k} \nabla \zeta_{0}-\int_{\Omega} \frac{c}{|x|^{2}} u_{k} \zeta_{0}=\int_{\Omega} u_{k}=\int_{\Omega} f_{k} \zeta_{0}
$$

Since $f \geq 0$, so are $f_{k}$ and $u_{k}$ and

$$
\begin{equation*}
\left\|u_{k}\right\|_{L^{1}}=\int_{\Omega} f_{k} \zeta_{0} \leq C\left\|f_{k}\right\|_{L_{-a, \delta}^{1}} \tag{1.14}
\end{equation*}
$$

Let $\zeta_{1}$ be the smooth solution of

$$
\left\{\begin{align*}
-\Delta \zeta_{1} & =1 \text { in } \Omega  \tag{1.15}\\
\zeta_{1} & =0 \text { on } \partial \Omega
\end{align*}\right.
$$

and integrate in the equation satisfied by $u_{k}$ :

$$
\begin{equation*}
\int_{\Omega} u_{k}-\int_{\Omega} \frac{c}{|x|^{2}} u_{k} \zeta_{1}=\int_{\Omega} f_{k} \zeta_{1} \tag{1.16}
\end{equation*}
$$

Using (1.14) and (1.16) and the inequality $m \delta(x) \leq \zeta_{1} \leq M \delta(x)$, where $m, M$ are some positive constants, we get

$$
\left\|u_{k}\right\|_{L_{-2}^{1}} \leq C\left\|f_{k}\right\|_{L_{-a, \delta}^{1}}
$$

It is then easy to construct by monotonicity a solution of (1.9) satisfying (1.10).

For estimate (1.11), one should just check that if $f \in L^{\infty},\|f\|_{L^{\infty}} \zeta_{0}$ is a supersolution of (1.9) and apply the maximum principle (see e.g. Lemma 1.4). Hence,

$$
u \leq\|f\|_{L^{\infty}} \zeta_{0}
$$

Applying this estimate to $-u$ yields (1.11).
For estimate (1.11'), $\|f\|_{L_{-b-2}^{\infty}} \zeta_{2}$ provides a supersolution of (1.9) where

$$
\left\{\begin{align*}
-\Delta \zeta_{2}-\frac{c}{|x|^{2}} \zeta_{2} & =|x|^{-b-2} \text { in } \Omega  \tag{1.17}\\
\zeta_{2} & =0 \text { on } \partial \Omega
\end{align*}\right.
$$

Observe that in the radial case $\zeta_{2}=C\left(|x|^{-b}-|x|^{-a}\right)$ so that in general $0 \leq \zeta_{2} \leq$ $C|x|^{-b}$ and that Lemma 1.4 may be applied because $a<b<a^{\prime}$.

## Remark 1.3.

- In view of Lemma 1.5, for equation (1.9) to have a solution with $f \in L_{\delta}^{1}$, it may be necessary that $f \in L_{-a, \delta}^{1}$.
- In the case $0<c<c_{0}$, if $\int_{\Omega}|f| \cdot|x|^{-a} \cdot|\ln (x)| \cdot \delta(x) d x<\infty$, that is, if we ask a little more regularity on $f$, then $u \in L_{-a-2}^{1}$ and is therefore unique (using Lemma 1.4.) For a proof, use the methods of the lemma with $\zeta_{2}$ solving

$$
\left\{\begin{aligned}
-\Delta \zeta_{3}-\frac{c}{|x|^{2}} \zeta_{3} & =|x|^{-a-2} & & \text { in } \Omega \\
\zeta_{3} & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

When $\Omega=B_{1}, \zeta_{3}=C|x|^{-a} \ln (1 /|x|)$.

Proof of Proposition 0.1 (case $0<c<c_{0}$ ).
Suppose first that $u$ is a strong solution of $\left(P_{t, p}\right)$.
Let $\zeta_{n} \in C_{c}^{\infty}(\Omega \backslash\{0\})$ be such that $0 \leq \zeta_{n} \leq 1,\left|\nabla \zeta_{n}\right| \leq C n,\left|\Delta \zeta_{n}\right| \leq C n^{2}$ and

$$
\zeta_{n}=\left\{\begin{array}{lll}
0 & \text { if } & |x| \leq 1 / n \text { and } \delta(x)<1 / n \\
1 & \text { if } & |x| \geq 2 / n \text { and } \delta(x)>2 / n
\end{array}\right.
$$

Multiplying $\left(P_{t, p}\right)$ by $u \zeta_{n}$ and integrating by parts, it follows that

$$
\int_{\Omega}\left(\frac{c}{|x|^{2}} u+u^{p}+t f\right) u \zeta_{n}=-\int_{\Omega} \Delta u u \zeta_{n}=\int_{\Omega}|\nabla u|^{2} \zeta_{n}+\int_{\Omega} u \nabla u \nabla \zeta_{n}
$$

Since $u \leq C|x|^{-a}$ and $p<p_{0}, u^{p} \leq C|x|^{-a-2}$. Hence, on the one hand, $u^{p} \in$ $L^{\frac{2 n}{n+2}}(\Omega)$ and on the other hand, the left-hand-side integral in the above equation is bounded by $C \int_{\Omega}|x|^{-2 a-2} \leq C$, whereas $\left|\int_{\Omega} u \nabla u \nabla \zeta_{n}\right|=\left|\frac{1}{2} \int_{\Omega} u^{2} \Delta \zeta_{n}\right| \leq$ $C n^{2} \int_{1 / n<|x|<2 / n}|x|^{-2 a} \rightarrow 0$ as $n \rightarrow \infty$. Hence $\int_{\Omega}|\nabla u|^{2} \zeta_{n} \leq C$ and $u \in H_{0}^{1}(\Omega)$.

Multipying $\left(P_{t, p}\right)$ by $\phi \zeta_{n}$ for $\phi \in C_{c}^{\infty}(\Omega)$ yields

$$
\int_{\Omega}\left(\frac{c}{|x|^{2}} u+u^{p}+t f\right) \phi \zeta_{n}=-\int_{\Omega} \Delta u \phi \zeta_{n}=\int_{\Omega} \zeta_{n} \nabla u \nabla \phi+\int_{\Omega} \phi \nabla u \nabla \zeta_{n}
$$

The last term in the right-hand-side can be rewritten as

$$
\int_{\Omega} \phi \nabla u \nabla \zeta_{n}=\int_{\Omega} \nabla(u \phi) \nabla \zeta_{n}-\int_{\Omega} u \nabla \phi \nabla \zeta_{n}=-\int_{\Omega} u \phi \Delta \zeta_{n}-\int_{\Omega} u \nabla \phi \nabla \zeta_{n}
$$

and converges to zero as in the previous case when $n \rightarrow \infty$.
It follows that $u$ is an $H_{0}^{1}(\Omega)$ solution of $\left(P_{t, p}\right)$.
Approximating $u \in H_{0}^{1}(\Omega)$ by smooth functions and integrating by parts implies that $H_{0}^{1}(\Omega)$ solutions are weak solutions.

Suppose now that $u$ is a weak solution satisfying the estimate $u \leq C|x|^{-a}$.
Then as before, $u^{p} \leq C|x|^{-a-2} \in L^{\frac{2 n}{n+2}}(\Omega) \subset H^{-1}(\Omega)$.
Letting $g=u^{p}+t f$, it follows from Lemma 1.1 that there exists a weak solution $v \in H_{0}^{1}(\Omega)$ of (1.9) with $g$ in place of $f . u$ is also a weak solution of (1.9) and by Remark 1.3, we must have $u=v \in H_{0}^{1}(\Omega)$. Hence, $u$ is an $H_{0}^{1}(\Omega)$ solution.

Finally if $u$ is an $H_{0}^{1}(\Omega)$ solution satisfying the estimate $u \leq C|x|^{-a}$, using local elliptic regularity theorems in $\Omega \backslash B_{\epsilon}$ for an arbitrary $\epsilon>0$, we may conclude that $u \in C^{\infty}(\bar{\Omega} \backslash\{0\})$ and satisfies $\left(P_{t, p}\right)$ in the strong sense.

## Lemma 1.4 (Maximum Principle).

If $\int_{\Omega}|u| \cdot|x|^{-a-2}<\infty$ and if

$$
\begin{equation*}
-\Delta u-\frac{c}{|x|^{2}} u \geq 0 \quad \text { in the weak sense. } \tag{1.16}
\end{equation*}
$$

then

$$
u \geq 0 \text { a.e. }
$$

Proof (case $0<c<c_{0}$ ).
It is enough to show that $\int_{\Omega} u \phi \geq 0$ for $\phi \in C_{c}^{\infty}(\Omega \backslash\{0\}), \phi \geq 0$.
For such a $\phi$ and $\epsilon>0$, construct $v_{\epsilon} \in C^{2}(\bar{\Omega}), v_{\epsilon} \geq 0$, solving

$$
\left\{\begin{aligned}
-\Delta v_{\epsilon}-\frac{c}{|x|^{2}+\epsilon} v_{\epsilon} & =\phi \text { in } \Omega \\
v_{\epsilon} & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

Also let $v \in H_{0}^{1}(\Omega)$ be the solution of

$$
\left\{\begin{aligned}
-\Delta v-\frac{c}{|x|^{2}} v & =\phi \text { in } \Omega \\
v & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

Using Lemma 1.1, since $-\Delta\left(v_{\epsilon}-v\right)-\frac{c}{|x|^{2}}\left(v_{\epsilon}-v\right) \leq 0$,

$$
\begin{equation*}
0 \leq v_{\epsilon} \leq v \quad \text { a.e. in } \Omega \tag{1.17}
\end{equation*}
$$

Applying (1.11) in Lemma 1.3 to $v$,

$$
0 \leq v \leq C|x|^{-a} \quad \text { a.e. in } \Omega
$$

Combining (1.17) and (1.18),

$$
\begin{equation*}
0 \leq v_{\epsilon} \leq C|x|^{-a} \quad \text { a.e. in } \Omega \tag{1.19}
\end{equation*}
$$

Applying (1.16) with $\phi=v_{\epsilon}$,

$$
\begin{gathered}
\int_{\Omega} u\left(-\Delta v_{\epsilon}-\frac{c}{|x|^{2}} v_{\epsilon}\right) \geq 0 \\
\text { Since }-\Delta v_{\epsilon}-\frac{c}{|x|^{2}} v_{\epsilon}=\phi-c\left\{\frac{1}{|x|^{2}}-\frac{1}{|x|^{2}+\epsilon}\right\} \\
\int_{\Omega} u \phi \geq \int_{\Omega} c\left\{\frac{1}{|x|^{2}}-\frac{1}{|x|^{2}+\epsilon}\right\} u v_{\epsilon} .
\end{gathered}
$$

Clearly, $\left\{v_{\epsilon}\right\}$ is monotone increasing and converges pointwise to a finite value a.e. in $\Omega$ by (1.19).

So the integrand in the right hand side of the previous equation converges a.e. to 0 . Using (1.19) and $u \in L_{-a-2}^{1}$, this integrand is dominated by an $L^{1}$ function.

By Lebesgue's theorem, we conclude that

$$
\int_{\Omega} u \phi \geq 0
$$

Remark 1.4.

- This maximum principle is sharp in the following sense :
if $q>-a$ then there exists $u \in L_{q-2}^{1}$ such that $-\Delta u-\frac{c}{|x|^{2}} u=0$ yet $u \not \equiv 0$.

Just take $\Omega=B_{1}$ and $u:=|x|^{-a^{\prime}}-|x|^{-a}$, with $-a^{\prime}$ and $-a$ defined in (0.2), (0.2').

We conclude this section with a lemma giving necessary conditions for the existence of a solution to the linear problem.

## Lemma 1.5.

Suppose $f \geq 0$ a.e., $f \not \equiv 0, \int_{\Omega} f(x) \operatorname{dist}(x, \partial \Omega) d x<\infty$.
If $u$ is a nonnegative weak solution of

$$
\left\{\begin{align*}
-\Delta u-\frac{c}{|x|^{2}} u & =f \text { in } \Omega  \tag{1.20}\\
u & =0 \text { on } \partial \Omega
\end{align*}\right.
$$

Then there exists a constant $C>0$ depending only on $\Omega$ such that

$$
u \geq C\left(\int_{\Omega} f \zeta_{0}\right) \zeta_{0} \quad \text { a.e. in } \Omega
$$

with $\zeta_{0}$ defined by (1.12). In particular, for some $m>0$

$$
u \geq m|x|^{-a} \text { a.e. near the origin }
$$

Furthermore, for any $\epsilon>0$, if $u$ denotes the minimal solution of (1.20) then

$$
\int_{\Omega} u \cdot|x|^{-a-2+\epsilon} d x \leq C_{\epsilon} \int_{\Omega} f \cdot|x|^{-a} d i s t(x, \partial \Omega) d x<\infty
$$

Most of the results of this lemma are a direct consequence of a more general theorem on the associated evolution equation, established by Baras and J. Goldstein (see [BG] Th 2.2 page 124.) We give here a simpler proof for convenience of the reader.

Proof (case $0<c<c_{0}$ ).

## Step 1. $\mathbf{u} \geq \mathbf{m}|\mathbf{x}|^{-a}$ near the origin.

Let $f_{1}=\min (f, k)$ with $k>0$ such that $f_{1} \not \equiv 0$ and $u_{1} \geq 0$ be the minimal solution of

$$
\left\{\begin{aligned}
-\Delta u_{1}-\frac{c}{|x|^{2}} u_{1} & =f_{1} \text { in } \Omega \\
u_{1} & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

Since $u$ is a supersolution of the above problem, $u_{1}$ is well defined and $0 \leq u_{1} \leq u$ so it suffices to prove the result for $u_{1}$.

Since $f_{1} \in L^{\infty}(\Omega)$, on the one hand $0 \leq u_{1} \leq C|x|^{-a}$ by (1.11) and on the other hand the equation has a solution $v \in H_{0}^{1}(\Omega)$. By Lemma 1.4, we must have $u_{1}=v$.

Now, since $u_{1} \not \equiv 0, u_{1} \geq 0$ and $-\Delta u_{1} \geq 0$ in the connected set $\Omega$, we have for some $\epsilon>0$ and $\eta>0$,

$$
u_{1} \geq \epsilon \quad \text { a.e. in } B_{2 \eta}
$$

Choose $C>0$ so that $\epsilon \geq C r^{-a}$ for $r \geq \eta$ and let $z=\left(u_{1}-C|x|^{-a}\right)^{-}$. Observe that $z \in H_{0}^{1}\left(B_{\eta}\right)$.

Next, we multiply $u_{1}-C|x|^{-a}$ by $z$ and integrate by parts :

$$
\begin{aligned}
0 \geq-\int_{\Omega}|\nabla z|^{2}+\int_{\Omega} \frac{c}{|x|^{2}} z^{2} & =\int_{\Omega} \nabla\left(u_{1}-C|x|^{-a}\right) \nabla z-\int_{\Omega} \frac{c}{|x|^{2}}\left(u_{1}-C|x|^{-a}\right) z \\
& =\int_{\Omega} f z-C\left(\int_{B_{\eta}} \nabla|x|^{-a} \nabla z-\int_{B_{\eta}} \frac{c}{|x|^{2}}|x|^{-a} z\right) \\
& \geq-C \int_{\partial B_{\eta}} z \partial_{\nu}|x|^{-a} \geq 0
\end{aligned}
$$

And hence $z \equiv 0$ in $B_{\eta}$.
Step 2. $\mathbf{u} \geq \mathbf{C}(\mathbf{K}, \Omega) \int_{\Omega} \mathbf{f} \zeta_{\mathbf{0}}$ in $\mathbf{K} \subset \subset \Omega$ when $\mathbf{f} \in \mathbf{L}^{\infty}(\Omega)$
The proof is an adaptation of Lemma 3.2 in [BC]. Observe that up to replacing $u$ by the minimal nonnegative solution of the problem, we may assume $u$ to be an $H_{0}^{1}(\Omega)$ solution satisfting $0 \leq u \leq C|x|^{-a}$.

Let $\rho=\operatorname{dist}(K, \partial \Omega) / 2$ and take $m$ balls of radius $\rho$ such that

$$
K \subset B_{\rho}\left(x_{1}\right) \cup \cdots \cup B_{\rho}\left(x_{m}\right) \subset \Omega
$$

Let $\zeta_{1}, \ldots, \zeta_{m}$ be the solutions (given, say, by Lemma 1.1) of

$$
\left\{\begin{aligned}
-\Delta \zeta_{i}-\frac{c}{|x|^{2}} \zeta_{i} & =\chi_{B_{\rho}\left(x_{i}\right)} \text { in } \Omega \\
\zeta_{i} & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

where $\chi_{A}$ denotes the characteristic function of $A$. There is a constant $C>0$ such that

$$
\zeta_{i}(x) \geq C \zeta_{0}(x) \quad \text { in } \Omega \quad \text { for } 1 \leq i \leq m
$$

Indeed, by Step 1, this inequality must hold near the origin and by Hopf's boundary lemma, we also have $\zeta_{i} \geq c \delta \geq C \zeta_{0}$ away from the origin.

Let now $x \in K$, and take a ball $B_{\rho}\left(x_{i}\right)$ containing $x$. Then $B_{\rho}\left(x_{i}\right) \subset B_{2 \rho}(x) \subset \Omega$ and, since $-\Delta u \geq 0$ in $\Omega$, we conclude

$$
\begin{aligned}
u(x) & \geq f_{B_{2 \rho}(x)} u=C \int_{B_{2 \rho}(x)} u \geq C \int_{B_{\rho}\left(x_{i}\right)} u \\
& =C \int_{\Omega} u\left(-\Delta \zeta_{i}-\frac{c}{|x|^{2}} \zeta_{i}\right)=C \int_{\Omega} f \zeta_{i} \\
& \geq C \int_{\Omega} f \zeta_{0}
\end{aligned}
$$

Step 3. $\mathbf{u} \geq \mathbf{C}(\Omega)\left(\int_{\Omega} \mathbf{f} \zeta_{\mathbf{0}}\right) \zeta_{\mathbf{0}}$ in $\Omega$ when $\mathbf{f} \in \mathbf{L}^{\infty}(\Omega)$
Suppose without loss of generality that $B_{1} \subset \Omega$ and let $K=\bar{B}_{1} \backslash B_{1 / 2}$. By Step 2, it suffices to prove the inequality in $\Omega \backslash K$. Let $w$ be the solution of

$$
\left\{\begin{aligned}
-\Delta w-\frac{c}{|x|^{2}} w & =0 \text { in } \Omega \backslash B_{1} \\
w & =0 \text { on } \partial \Omega \\
w & =1 \text { on } \partial B_{1}
\end{aligned}\right.
$$

and extend $w$ by $w:=(2|x|)^{-a}$ in $B_{1 / 2}$, so that the above equation still holds in $\Omega \backslash K$ with $\left.w\right|_{\partial K} \equiv 1$. By Hopf's boundary Lemma applied in $\Omega \backslash B_{1}$, we conclude that

$$
w \geq C \zeta_{0} \quad \text { in } \Omega \backslash K
$$

$u$ is assumed to be dominated by $C|x|^{-a}$ so we can apply the maximum principle (Lemma 1.4) in $\Omega \backslash K$ to conclude that

$$
u \geq C\left(\int_{\Omega} f \zeta_{0}\right) w \geq C\left(\int_{\Omega} f \zeta_{0}\right) \zeta_{0} \quad \text { in } \Omega \backslash K
$$

Step 4. $\int_{\Omega}|\mathbf{x}|^{-\mathbf{a}} \mathbf{f} \delta(\mathbf{x})<\infty$.
We assume for now that $f \in L^{\infty}(\Omega)$ and that $u \geq 0$ is the minimal solution of (1.20).

We let $\left\{\phi_{n}\right\}$ be a sequence of smooth, nonnegative and bounded functions converging pointwise and monotonically to $c|x|^{-a-2}$ and construct $v_{n}$ as the (smooth) solution of

$$
\left\{\begin{aligned}
-\Delta v_{n} & =\phi_{n} \text { in } \Omega \\
v_{n} & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

Testing $v_{n}$ in (1.20) yields

$$
\begin{equation*}
\int_{\Omega} f v_{n}=\int_{\Omega} u\left(-\Delta v_{n}-\frac{c}{|x|^{2}} v_{n}\right)=\int_{\Omega} u\left(\phi_{n}-\frac{c}{|x|^{2}} v_{n}\right) \tag{1.21}
\end{equation*}
$$

Now $\phi_{n} \nearrow c|x|^{-a-2}$ pointwise and in $L^{1}$, so, by Lemma 2.1, $v_{n} \nearrow|x|^{-a}-w$ pointwise and in $L^{1}$, where $w$ solves

$$
\left\{\begin{aligned}
-\Delta w & =0 \text { in } \Omega \\
w & =|x|^{-a} \text { on } \partial \Omega
\end{aligned}\right.
$$

Since $u$ is minimal, $0 \leq u \leq C|x|^{-a}$ by (1.11) and we can safely pass to the limit in (1.21) to obtain

$$
\int_{\Omega}\left(|x|^{-a}-w\right) f=\int_{\Omega} u\left(c|x|^{-a-2}-\frac{c}{|x|^{2}}\left(|x|^{-a}-w\right)\right)=c \int_{\Omega} \frac{u}{|x|^{2}} w
$$

Observe that $w$ is bounded and that $|x|^{-a}-w \geq C|x|^{-a} \delta(x)$, hence

$$
\int_{\Omega}|x|^{-a} f \delta(x) \leq C \int_{\Omega}|x|^{-2} u
$$

This estimate holds when $f \in L^{\infty}$ and $u$ is minimal but also in the general case, as approximation of $f$ by $f_{n}=\min (n, f)$ shows.

Step 5. $\mathbf{u} \geq \mathbf{C}(\Omega)\left(\int_{\Omega} \mathbf{f} \zeta_{\mathbf{0}}\right) \zeta_{\mathbf{o}}$ in $\Omega$ when $\mathbf{f} \in \mathbf{L}_{-\mathbf{a}, \delta}^{\mathbf{1}}$
Let $k>0$ be so large that $f_{k}=\min (f, k) \not \equiv 0$. Then $u$ is a supersolution of (1.20) with $f_{k}$ in place of $f$ and by Step 3 , we have

$$
u \geq C(\Omega)\left(\int_{\Omega} f_{k} \zeta_{0}\right) \zeta_{0}
$$

Letting $k \rightarrow \infty$, Lebesgue's theorem yields the desired result.
Step 6. $\int_{\Omega}|\mathbf{x}|^{-\mathbf{a}-2+\epsilon} \mathbf{u}<\infty$.
We proceed as in Step 4, only this time we let $\phi_{n} \nearrow-P(-a+\epsilon)|x|^{-a-2+\epsilon}$, where $P(X)=X(X-1)+(n-1) X+c$ and construct $v_{n}$ solving

$$
\left\{\begin{aligned}
-\Delta v_{n}-\frac{c}{|x|^{2}+1 / n} v_{n} & =\phi_{n} \text { in } \Omega \\
v_{n} & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

Hence,

$$
\begin{equation*}
\int_{\Omega} f v_{n}=\int_{\Omega} u \phi_{n}+\int_{\Omega}\left(\frac{c}{|x|^{2}+1 / n}-\frac{c}{|x|^{2}}\right) v_{n} u \tag{1.22}
\end{equation*}
$$

If $\zeta$ solves

$$
\left\{\begin{aligned}
-\Delta \zeta-\frac{c}{|x|^{2}} \zeta & =-P(-a+\epsilon)|x|^{-a-2+\epsilon} \text { in } \Omega \\
\zeta & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

then we have $0 \leq v_{n} \leq \zeta \leq C|x|^{-a}$. Indeed, if $\Omega=B_{1}$, then $\zeta=\zeta^{1}:=$ $C\left(|x|^{-a}-|x|^{-a+\epsilon}\right)$. Otherwise, $\Omega \subset B_{R}$ for some $R>0$ and $C \zeta^{1}(x / R)$ is a supersolution of the problem, for some $C>0$.

By Step $4, \int_{\Omega} f v_{n} \leq \int_{\Omega} f \zeta<\infty$. Assuming first that $f$ is bounded (whence $u \leq$ $C|x|^{-a}$ ) and then working by approximation, it follows from Lebesgue's theorem and from (1.22) that

$$
\int_{\Omega}|x|^{-a-2+\epsilon} u \leq C_{\epsilon} \int_{\Omega} f \zeta<\infty
$$

Remark. More results about the linear theory of our operator, with $c \in \mathbb{R}$ arbitrary have been detailed by F. Pacard in unpublished work (see [P].)

## 2. Existence vs. COMPLETE BLOW-UP

In this section, we will prove existence or nonexistence of weak solutions of $\left(P_{t, p}\right)$, using the tools we have just constructed and monotonicity arguments.

### 2.1. Case $\mathbf{p}<\mathbf{p}_{\mathbf{0}}, \mathbf{c}<\mathrm{c}_{\mathbf{0}}$ : existence for small $\mathbf{t}>\mathbf{0}$.

$p_{0}$ has been defined so that $p_{0} a=a+2$. So, for $p<p_{0}, a p<a+2$ and for some $b \in\left(a, a^{\prime}\right)$, the inequality $b p<b+2$ still holds. We fix such a $b$ and prove that for an appropriate choice of $A>0$ and for $t>0$ small,

$$
w:=A|x|^{-b} \in H^{1}(\Omega) \text { is a supersolution of }\left(P_{t, p}\right)
$$

Observe that $w \in H^{1}(\Omega)$ as long as $b$ is close enough to $a$, which may be assumed. We have
$-\Delta w-\frac{c}{|x|^{2}} w=-A P(-b)|x|^{-b-2} \quad$ where $P(X)=X(X-1)+(n-1) X+c$
Observe that $P(-b)<0$ since $b \in\left(a, a^{\prime}\right)$ and $a^{\prime}$ and $a$ are the roots of $P(X)$.
We would like to have $-A P(-b)|x|^{-b-2} \geq A^{p}|x|^{-p b}+t f$ in $\Omega$. This will be true as soon as

$$
\left\{\begin{array}{l}
-\frac{1}{2} A P(-b)|x|^{-b-2} \geq A^{p}|x|^{-p b} \quad \text { and } \\
-\frac{1}{2} A P(-b)|x|^{-b-2} \geq t f
\end{array}\right.
$$

The first inequality amounts to

$$
A \leq\left[-\frac{1}{2} P(-b)|x|^{p b-b-2}\right]^{\frac{1}{p-1}}
$$

which will be satisfied, taking $R>0$ such that $\Omega \subset B_{R}$, if

$$
A \leq\left[-\frac{1}{2} P(-b) R^{p b-b-2}\right]^{\frac{1}{p-1}}
$$

since $p b-b-2<0$.
With such a choice of $A$, pick any $t>0$ such that

$$
-\frac{1}{2} A P(-b) R^{-b-2} \geq t\|f\|_{L^{\infty}}
$$

We have just constructed $w \in H^{1}(\Omega)$ such that

$$
\left\{\begin{array}{rlrl}
-\Delta w-\frac{c}{|x|^{2}} w & \geq w^{p}+t f & \text { in } \Omega \\
w & \geq 0 & & \text { on } \partial \Omega
\end{array}\right.
$$

Finally we construct an $H_{0}^{1}(\Omega)$ supersolution of $\left(P_{t, p}\right)$. We let $w_{1}$ be a smooth extension inside $\Omega$ of $\left.w\right|_{\partial \Omega}$ which is also supported away from the origin. Then $g=\Delta w_{1}+\frac{c}{|x|^{2}} w_{1}$ is smooth and bounded and using Lemma 1.1, there is a unique strong solution $z$ of

$$
\left\{\begin{align*}
-\Delta z-\frac{c}{|x|^{2}} z & =g \text { in } \Omega  \tag{2.1}\\
z & =0 \text { on } \partial \Omega
\end{align*}\right.
$$

Letting $w_{2}=z+w_{1}$, it follows that

$$
\left\{\begin{align*}
-\Delta w_{2}-\frac{c}{|x|^{2}} w_{2} & =0 \text { in } \Omega  \tag{2.2}\\
w_{2} & =w \text { on } \partial \Omega
\end{align*}\right.
$$

Multiplying by $w_{2}^{-}$, it follows that $w_{2} \geq 0$ a.e. in $\Omega$.
It is now clear that $\tilde{w}=w-w_{2}$ is an $H_{0}^{1}(\Omega)$ supersolution of $\left(P_{t, p}\right)$. For convenience, we drop the superscript $\sim$ thereafter.

Construction of a minimal solution $u$ of $\left(P_{t, p}\right)$ is now just a matter of monotone iteration. For this purpose we recall the following lemma, proved in [BCMR] :

Lemma 2.1. Suppose $\int_{\Omega}|f(x)| \operatorname{dist}(x, \partial \Omega)<\infty$. Then there exists a unique $v \in$ $L^{1}(\Omega)$ which is a weak solution of

$$
\left\{\begin{aligned}
-\Delta v & =f \text { in } \Omega \\
v & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

Moreover,

$$
\|v\|_{L^{1}} \leq C\|f\|_{L_{\delta}^{1}}
$$

Moreover if $v \in L^{1}(\Omega)$ and $-\Delta v \geq 0$ weakly, i.e. if

$$
\int_{\Omega}(-\Delta \phi) v \geq 0 \quad \text { for all } \phi \in C^{2}(\bar{\Omega}),\left.\phi\right|_{\partial \Omega} \equiv 0, \phi \geq 0 \text { in } \Omega
$$

then

$$
v \geq 0 \quad \text { a.e. in } \Omega
$$

Define $\left\{u_{k}\right\}$ by induction to be the $L^{1}$ weak solutions of

$$
\begin{gathered}
\left\{\begin{array}{rll}
-\Delta u_{0}=t f & \text { in } \Omega & \text { for } k=0 \\
u_{0}=0 & \text { on } \partial \Omega
\end{array}\right. \\
\left\{\begin{aligned}
-\Delta u_{k}=\frac{c}{|x|^{2}} u_{k-1}+u_{k-1}^{p}+t f & \text { in } \Omega \\
u_{k}=0 & \text { on } \partial \Omega
\end{aligned}\right. \\
\text { for } k \geq 1
\end{gathered}
$$

We now check that this definition makes sense and that $\left(u_{k}\right)$ is monotone and satisfies $0 \leq u_{k} \leq w$ a.e. in $\Omega$.

For $u_{0}$ there is nothing to prove. Suppose the result true up to order $k-1$. Then

$$
\begin{aligned}
0 \leq \frac{c}{|x|^{2}} u_{k-1}+u_{k-1}^{p}+t f & \leq \frac{c}{|x|^{2}} w+w^{p}+t f \\
& \leq C|x|^{-a-2} \in L^{1}(\Omega)
\end{aligned}
$$

So $u_{k}$ is well defined using the previous lemma, $u_{k} \geq 0$ a.e. and since
$-\Delta\left(u_{k}-u_{k-1}\right)=\frac{c}{|x|^{2}}\left(u_{k-1}-u_{k-2}\right)+u_{k-1}^{p}-u_{k-2}^{p} \geq 0$ by induction hypothesis
and similarly $-\Delta\left(w-u_{k}\right) \geq 0$, we conclude using Lemma 2.1 that

$$
0 \leq u_{k-1} \leq u_{k} \leq w \quad \text { a.e. in } \Omega
$$

By a standard monotone convergence argument, $\left\{u_{k}\right\}$ converges to a weak solution of $\left(P_{t, p}\right)$.

### 2.2. Pushing to to $\mathrm{t}_{0}$.

We let $t_{0}=\sup \left\{t:\left(P_{t, p}\right)\right.$ has a weak solution. $\}$ and adapt the methods of [BCMR].

If $\phi_{1}$ is a positive eigenvector of $-\Delta$ (with zero Dirichlet condition) associated to its first eigenvalue $\lambda_{1}$, in other words if $\phi_{1}>0$ in $\Omega$ and, for some $\lambda_{1}>0$,

$$
\left\{\begin{aligned}
-\Delta \phi_{1} & =\lambda_{1} \phi_{1} \text { in } \Omega \\
\phi_{1} & =0 \text { on } \partial \Omega
\end{aligned}\right.
$$

and if $u$ is a weak solution of $\left(P_{t, p}\right)$, testing against $\phi_{1}$ yields

$$
\int_{\Omega} \frac{c}{|x|^{2}} u \phi_{1}+\int_{\Omega} u^{p} \phi_{1}+t \int_{\Omega} f \phi_{1}=\lambda_{1} \int_{\Omega} u \phi_{1}
$$

and, by Young's inequality,

$$
\lambda_{1} \int_{\Omega} u \phi_{1} \leq \frac{1}{2} \int_{\Omega} u^{p} \phi_{1}+C \int_{\Omega} \phi_{1} .
$$

Thus,

$$
\begin{equation*}
t \int_{\Omega} f \phi_{1}+\int_{\Omega} \frac{c}{|x|^{2}} u \phi_{1}+\int_{\Omega} u^{p} \phi_{1} \leq C \tag{2.3}
\end{equation*}
$$

which implies $t_{0}<\infty$. In particular, there are no weak solutions of $\left(P_{t, p}\right)$ for $t>t_{0}$. This implies complete blow-up (see Definition 0.1 ), as the following proposition shows.

## Proposition 2.1.

Suppose (0.1) holds, $p>1$ and $t>0$.
If $\left(P_{t, p}\right)$ has no weak solution then there is complete blow-up.
Proof.
The proof is an easy adaptation of Theorem 3.1 in [BC].
Suppose indeed that $\left(P_{t, p}\right)$ has no weak solution and by contradiction that $\int_{\Omega} g_{n}\left(\underline{u_{n}}\right) \delta+\int_{\Omega} a_{n} \underline{u_{n}} \delta \leq C$, where $\left\{a_{n}\right\},\left\{g_{n}\right\},\left\{u_{n}\right\}$, are given in Definition 0.1.

Then, multiplying $\left(P_{n}\right)$ by $\zeta_{1}$, solution of (1.15) we get
$\int_{\Omega} \underline{u_{n}}\left(-\Delta \zeta_{1}\right)-\int_{\Omega} a_{n} \underline{u_{n}} \zeta_{1}=\int_{\Omega} g_{n}\left(\underline{u_{n}}\right) \zeta_{1}+\int_{\Omega} t f \zeta_{1}$.
Hence, $\int_{\Omega} \underline{u_{n}} \leq C$ and there exists a $u$ such that $\underline{u_{n}} \nearrow u$ in $L^{1}(\Omega)$, by monotone convergence.

Since $\left\{a_{n}\right\}$ and $\left\{g_{n}\right\}$ converge monotonically, we can pass to the limit in $\left(P_{n}\right)$, using monotone convergence again and obtain a solution $u$ of $\left(P_{t, p}\right)$, which is a contradiction.

We have just proved that $\int_{\Omega} g_{n}\left(\underline{u_{n}}\right) \delta+\int_{\Omega} a_{n} \underline{u_{n}} \delta \rightarrow \infty$. Now, using $\left(P_{n}\right)$ and Lemma 3.2 in [BC], it follows that

$$
\frac{u_{n}(x)}{\delta(x)} \geq C(\Omega)\left(\int_{\Omega} g_{n}\left(\underline{u_{n}}\right) \delta+\int_{\Omega} a_{n} \underline{u_{n}} \delta\right) \rightarrow \infty
$$

Next, we want to prove that if $\left(P_{\tau, p}\right)$ has a solution then so does $\left(P_{t, p}\right)$ for $0<t \leq \tau$. This is true because $u_{\tau}$ is a supersolution of $\left(P_{t, p}\right)$ in the sense that, weakly,

$$
-\Delta u_{\tau} \geq \frac{c}{|x|^{2}} u_{\tau}+u_{\tau}^{p}+t f
$$

and with the help of Lemma 2.1, we may construct a solution of $\left(P_{t, p}\right)$ by monotone iteration.

Finally, we prove that $\left(P_{t_{0}, p}\right)$ has a weak solution. Choose a nondecreasing sequence $\left\{t_{n}\right\}$ converging to $t_{0}$ and for each $n \in \mathbb{N}$, let $u_{n}$ be a (weak) solution of $\left(P_{t_{n}, p}\right)$. Since $\phi_{1} \geq m \delta(x)$ for some $m>0$, equation (2.3) implies that

$$
\int_{\Omega} \frac{c}{|x|^{2}} u_{n} \delta(x)+\int_{\Omega} u_{n}^{p} \delta(x) \leq C
$$

Multiplying by $\zeta_{1}$, solution of (1.15) then implies boundedness of $\left\{u_{n}\right\}$ in $L^{1}$ and hence monotone convergence to a solution of $\left(P_{t_{0}, p}\right)$ as $t_{n} \rightarrow t_{0}$.

### 2.3. Case $0<c<c_{0}, p \geq p_{0}$ : blow-up for all $t>0$.

By Proposition 2.1, we just need to prove that there are no weak solutions of ( $P_{t, p}$ ) for $p \geq p_{0}$. Assume by contradiction there exists one and call it $u$. If we apply Lemma 1.5 with $u^{p}+t f$ in place of $f$, it follows that

$$
\int_{\Omega} u^{p}|x|^{-a} \delta(x)<\infty \quad \text { and } \quad u \geq m|x|^{-a} \quad \text { a.e. near the origin. }
$$

Using Hölder's inequality,
$\int_{\Omega} u|x|^{-a-2} \delta(x) \leq\left(\int_{\Omega} u^{p}|x|^{-a} \delta(x)\right)^{1 / p} \cdot\left(\int_{\Omega}|x|^{-a-2 \frac{p}{p-1}}\right)^{\frac{p-1}{p}}$.
If $p \geq p_{0}$ and $c<c_{0}$ then $-a-2 \frac{p}{p-1}>-n$, hence, since $u \in L^{1}(\Omega)$,

$$
\begin{equation*}
\int_{\Omega} u|x|^{-a-2}<\infty \tag{2.4}
\end{equation*}
$$

Suppose without loss of generality, that $\Omega \subset B_{1}$ and define $w=A|x|^{-a} \ln \left(\frac{1}{|x|}\right)$ for some $A>0$.

Then $-\Delta w-\frac{c}{|x|^{2}} w=A \sqrt{(n-2)^{2}-4 c}|x|^{-a-2}$. Also,
$-\Delta u-\frac{c}{|x|^{2}} u \geq u^{p} \geq m|x|^{-a p} \geq m|x|^{-a-2} \quad$ in $B_{\eta}$, for a fixed small $\eta>0$.
Let $A=m\left(\sqrt{(n-2)^{2}-4 c}+c \ln \frac{1}{\eta}\right)^{-1}$ and $C=A \eta^{-a} \ln \frac{1}{\eta}$.
Finally define $z=u+C-w$. Using (2.4), $z \in L^{1}\left(B_{\eta},|x|^{-a-2} d x\right)$. Furthermore,

$$
\begin{aligned}
-\Delta z-\frac{c}{|x|^{2}} z & \geq u^{p}-\frac{c C}{|x|^{2}}-A \sqrt{(n-2)^{2}-4 c}|x|^{-a-2} \\
& \geq m|x|^{-a p}-c C|x|^{-2}-A \sqrt{(n-2)^{2}-4 c}|x|^{-a-2} \\
& \geq|x|^{-2}\left[m|x|^{-a}-c C-A \sqrt{(n-2)^{2}-4 c}|x|^{-a}\right] \\
& \geq|x|^{-2}\left[\left(m-A \sqrt{\left.(n-2)^{2}-4 c\right)} \eta^{-a}-c C\right]\right. \\
& \geq 0
\end{aligned}
$$

All these inequalities hold in the weak sense in $B_{\eta}$ (since our choice of constants implies $\left.z\right|_{\partial B_{\eta}} \geq C-\left.w\right|_{\partial B_{\eta}} \geq 0$.)

Applying Lemma 1.4, we conclude

$$
u \geq A|x|^{-a} \ln \frac{1}{|x|}-C \quad \text { a.e. in } B_{\eta}
$$

Choosing $A$ and $\eta$ smaller, we may assume that

$$
u \geq A|x|^{-a} \ln \frac{1}{|x|} \geq 1 \quad \text { a.e. in } B_{\eta}
$$

The next step is to consider the function $\Phi \in C^{1}(\mathbb{R})$ defined by

$$
\Phi(x)=\left\{\begin{aligned}
\ln x & \text { if } x \geq 1 \\
x-1 & \text { otherwise } .
\end{aligned}\right.
$$

and apply Lemma 1.7 in $[\mathrm{BC}]$ to conclude that in $B_{\eta}$

$$
\begin{aligned}
-\Delta(\ln u) & \geq \frac{-\Delta u}{u} \geq u^{p-1} \geq A^{p-1}|x|^{-a(p-1)}\left(\ln \frac{1}{|x|}\right)^{p-1} \\
& \geq A^{p-1}|x|^{-2}\left(\ln \frac{1}{|x|}\right)^{p-1}
\end{aligned}
$$

Now if $v=\left(\ln \frac{1}{|x|}\right)^{p}$, a computation yields

$$
-\Delta v \leq C|x|^{-2}\left(\ln \frac{1}{|x|}\right)^{p-1}
$$

And by the $L^{1}$ maximum principle (Lemma 2.1),

$$
\ln u \geq d\left(\ln \frac{1}{|x|}\right)^{p}-C \quad \text { for some } d>0 \text { and } C>0
$$

This clearly violates $u \in L_{l o c}^{1}(\Omega)$.

## 3. Regularity

We start out with a result in the spirit of Lemma 5.3 in [BC] :
Lemma 3.1. Let $f \in L_{-a, \delta}^{1}$ and $v=|x|^{-a}$. Then if $u \in L_{-2}^{1}$ is the solution given by Lemma 1.3 of

$$
\left\{\begin{aligned}
-\Delta u-\frac{c}{|x|^{2}} u & =f & & \text { in } \Omega \\
u & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

and if $\Phi \in C^{1}(\mathbb{R})$ is concave, $\Phi^{\prime} \in L^{\infty}$ and $\Phi(1)=0$, then $v \Phi\left(\frac{u}{v}\right) \in L_{-2}^{1}$ and

$$
-\Delta\left(v \Phi\left(\frac{u}{v}\right)\right)-\frac{c}{|x|^{2}}\left(v \Phi\left(\frac{u}{v}\right)\right) \geq \Phi^{\prime}\left(\frac{u}{v}\right) f \quad \text { in the weak sense. }
$$

Proof (case $0<c<c_{0}$ ). Suppose first $u, v \in C^{2}(\bar{\Omega}), v>0$ in $\Omega$ and $\Phi \in C^{2}(\mathbb{R})$ and write $L=-\Delta-a(x)$ where $a(x)$ is a smooth bounded function. Applying Lemma 5.3 in [BC], it follows that a.e. in $\Omega$,

$$
\begin{aligned}
L w & \geq \Phi^{\prime}(u / v)(-\Delta u)+\left[\Phi(u / v)-\Phi^{\prime}(u / v) u / v\right](-\Delta v)-a(x) \Phi(u / v) v \\
& \geq \Phi^{\prime}(u / v) L u+\left[\Phi(u / v)-\Phi^{\prime}(u / v) u / v\right] L v \\
& \geq \Phi^{\prime}(u / v)(L u-L v)+\left[\Phi(u / v)-\Phi^{\prime}(u / v) u / v+\Phi^{\prime}(u / v)\right](L v)
\end{aligned}
$$

Since $\Phi$ is concave,

$$
\Phi(s)+(1-s) \Phi^{\prime}(s) \geq \Phi(1) \quad \text { for all } s \in \mathbb{R}
$$

Hence, if $w=v \Phi(u / v)$,

$$
\begin{equation*}
L w \geq \Phi^{\prime}(u / v)(L u-L v) \quad \text { a.e. in } \Omega \tag{3.1}
\end{equation*}
$$

Since $\Phi^{\prime}$ is bounded, we see, as in $[\mathrm{BC}]$, that

$$
\begin{equation*}
|v \Phi(u / v)|=|v(\Phi(u / v)-\Phi(0))+\Phi(0) v| \leq C(u+v) \tag{3.2}
\end{equation*}
$$

Hence, $w$ vanishes on $\partial \Omega$ and integrating by parts, (3.1) holds in the weak sense. By approximation of $\Phi$, we can also say that (3.1) holds even when $\Phi$ is only $C^{1}$.

In the general case, let $a_{n}=c /(|x|+1 / n)^{2}$ and $f_{n}$ be a smooth bounded function increasing pointwise and respectively to $c /|x|^{2}, f$. Let $u_{n}$ solve the equation $L_{n} u_{n}=$ $f$ (with zero boundary condition), where $L_{n}=-\Delta-a_{n}(x)$. Also write $w_{n}=$ $v_{n} \Phi\left(u_{n} / v_{n}\right)$ where $v_{n}=(|x|+1 / n)^{-a}$.

We can then apply (3.1) to obtain

$$
-\Delta w_{n}-a_{n}(x) w_{n} \geq \Phi^{\prime}\left(u_{n} / v_{n}\right) f_{n} \quad \text { weakly }
$$

Clearly, $v \Phi(u / v)$ is well defined a.e. Moreover, it is clear that $u_{n} \nearrow u$ in $L^{1}$ and that $a_{n}(x) u_{n}(x) \nearrow \frac{c}{|x|^{2}} u(x)$ in $L_{\delta}^{1}$ and similarly for $v$. So that, using the above equation and Lebesgue's theorem

$$
w_{n} \rightarrow w \quad \text { in } L^{1} \text { and } \quad a_{n}(x) w_{n} \rightarrow \frac{c}{|x|^{2}} w \quad \text { in } L_{\delta}^{1}
$$

Since $\Phi^{\prime}$ is bounded, we can also easily pass to the limit in the right-hand side and obtain the desired result.

## Lemma 3.2.

Let $u$ be the minimal weak solution of $\left(P_{t, p}\right)$ for $t<t_{0}\left(\right.$ and $\left.p<p_{0}\right)$.
Then $u$ is a strong solution of $\left(P_{t, p}\right)$

## Remark 3.2.

- By Proposition 0.1, we only need to show that $0<u \leq C|x|^{-a}$
- By Lemma 1.5, we also have the lower bound $u \geq m|x|^{-a} \operatorname{dist}(x, \partial \Omega)$.

Proof.
Recall that $\zeta_{0}$ solving, for $f$ as in the definition of $\left(P_{t, p}\right)$,

$$
\left\{\begin{aligned}
-\Delta \zeta_{0}-\frac{c}{|x|^{2}} \zeta_{0} & =f & & \text { in } \Omega \\
\zeta_{0} & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

satisfies $0<\zeta_{0} \leq C|x|^{-a}$. For $u \in \mathbb{R}^{+}$, let

$$
g(u)=\left(u+t_{0}\left\|\zeta_{0} / v\right\|_{L^{\infty}}\right)^{p} \quad \text { and } \quad \tilde{g}(u)=\left(u+t\left\|\zeta_{0} / v\right\|_{L^{\infty}}\right)^{p}
$$

and construct $\Phi \in C^{1}(\mathbb{R})$ with $\Phi(0)=0$ and

$$
\begin{equation*}
\Phi^{\prime}(u)=\frac{\tilde{g}(\Phi(u))}{g(u)} \tag{3.3}
\end{equation*}
$$

as in Lemma 4 of [BCMR].
Next, if $u_{0}$ is the minimal solution of $\left(P_{t_{0}, p}\right)$ then $z:=u_{0}-t_{0} \zeta_{0}$ is the minimal solution of

$$
\left\{\begin{aligned}
-\Delta z-\frac{c}{|x|^{2}} z & =\left(z+t_{0} \zeta_{0}\right)^{p} & & \text { in } \Omega \\
z & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

Applying Lemma 3.1 to $z$ with the above function $\Phi$ and $v=|x|^{-a}$,

$$
\begin{aligned}
-\Delta\left(v \Phi\left(\frac{z}{v}\right)\right)- & \frac{c}{|x|^{2}}\left(v \Phi\left(\frac{z}{v}\right)\right) \geq \Phi^{\prime}\left(\frac{z}{v}\right)\left(z+t_{0} \zeta_{0}\right)^{p} \geq \\
& \left(\frac{\Phi\left(\frac{z}{v}\right)+t\left\|\zeta_{0} / v\right\|_{L^{\infty}}}{\frac{z}{v}+t_{0}\left\|\zeta_{0} / v\right\|_{L^{\infty}}}\right)^{p}\left(z+t_{0} \zeta_{0}\right)^{p}
\end{aligned}
$$

We need the following easy lemma:
Lemma 3.3. Let $A, B>0$ such that $A \leq \frac{t}{t_{0}} B$.
Then $F(C):=\frac{A+t C}{B+t_{0} C}$ is increasing with $C$.
Observe that, since $\Phi$ is concave and $\Phi^{\prime}$ is defined by $(3.3), \Phi^{\prime}(u) \leq \Phi^{\prime}(0)=$ $\left(\frac{t}{t_{0}}\right)^{p}<\frac{t}{t_{0}}$ for $u \in \mathbb{R}^{+}$. Hence, since $\Phi(0)=0, \Phi(u) \leq \frac{t}{t_{0}} u$ for $u \in \mathbb{R}^{+}$. Applying Lemma 3.3 with $A=\Phi\left(\frac{z}{v}\right)$ and $B=\frac{z}{v}$, we get

$$
\frac{\Phi\left(\frac{z}{v}\right)+t \frac{\zeta_{0}}{v}}{\frac{z}{v}+t_{0} \frac{\zeta_{0}}{v}} \leq \frac{\Phi\left(\frac{z}{v}\right)+t\left\|\frac{\zeta_{0}}{v}\right\|_{L^{\infty}}}{\frac{z}{v}+t_{0}\left\|\frac{\zeta_{0}}{v}\right\|_{L^{\infty}}}
$$

and

$$
\begin{aligned}
-\Delta\left(v \Phi\left(\frac{z}{v}\right)\right)-\frac{c}{|x|^{2}}\left(v \Phi\left(\frac{z}{v}\right)\right) & \geq\left(\frac{\Phi\left(\frac{z}{v}\right)+t \zeta_{0} / v}{\frac{z}{v}+t_{0} \zeta_{0} / v}\right)^{p}\left(z+t_{0} \zeta_{0}\right)^{p} \\
& \geq\left(v \Phi\left(\frac{z}{v}\right)+t \zeta_{0}\right)^{p}
\end{aligned}
$$

We finally define $w=v \Phi\left(\frac{z}{v}\right)+t \zeta_{0}$, which satisfies

$$
\left\{\begin{aligned}
-\Delta w-\frac{c}{|x|^{2}} w & \geq w^{p}+t f & & \text { in } \Omega \\
w & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

We have just constructed a supersolution of problem $\left(P_{t, p}\right)$ satisfying $0<w \leq$ $C|x|^{-a}$ (since $\Phi(\infty)<\infty$ by Lemma 4 in $[B C M R]$ ) and, of course, the same estimate holds for $u$, the minimal solution of $\left(P_{t, p}\right)$.

This completes the proof of Theorem 1 (in the case $0<c<c_{0}$.)

## 4. Stability

We show first that $\lambda_{1}\left(u_{t}\right)>-\infty$ (recall Definition 0.2 ) and study the corresponding eigenfunction $\phi_{1}$.

Indeed, if $u_{t}$ is the minimal solution of $\left(P_{t, p}\right)$ with $t<t_{0}$, then $0 \leq u_{t} \leq C|x|^{-a}$ and

$$
\begin{aligned}
\int_{\Omega} u_{t}^{p-1} \phi^{2} & \leq C \int_{\Omega}|x|^{-a(p-1)} \phi^{2} \leq C\left(\int_{\Omega}|x|^{-2} \phi^{2}\right)^{\frac{a(p-1)}{2}} \cdot\left(\int_{\Omega} \phi^{2}\right)^{1-\frac{a(p-1)}{2}} \\
& \leq C\|\phi\|_{H_{0}^{1}(\Omega)}^{a(p-1)}\|\phi\|_{L^{2}}^{2-a(p-1)}
\end{aligned}
$$

So $\lambda_{1}>-\infty$ and if $\left\{\phi_{n}\right\}$ is a minimizing sequence of $J$ (see Definition 0.2 ), $\left\{\phi_{n}\right\}$ is bounded in $H_{0}^{1}(\Omega)$ and converges (weakly and up to a subsequence) to $\phi_{1} \in H_{0}^{1}(\Omega)$ solving

$$
\left\{\begin{align*}
-\Delta \phi_{1}-\frac{c}{|x|^{2}} \phi_{1} & =p u_{t}^{p-1} \phi_{1}+\lambda_{1} \phi_{1} & & \text { in } \Omega  \tag{4.1}\\
\phi_{1} & =0 & & \text { on } \partial \Omega
\end{align*}\right.
$$

Claim. $0 \leq \phi_{1} \leq C|x|^{-a}$
Testing equation (4.1) against $\phi_{1}^{+}$, it follows that

$$
\int_{\Omega}\left|\nabla \phi_{1}^{+}\right|^{2}-\int_{\Omega} \frac{c}{|x|^{2}}{\phi_{1}^{+2}}^{2}-\int_{\Omega} p u^{p-1} \phi_{1}^{+2}=\lambda_{1} \int_{\Omega} \phi_{1}^{+^{2}}
$$

Hence $\phi_{1}^{+}$is also a minimizer of $J$ and up to replacing $\phi_{1}$ by $\phi_{1}^{+}$, we may assume that $\phi_{1} \geq 0$.

Next, using local elliptic regularity, $\phi_{1} \in C^{\infty}(\bar{\Omega} \backslash\{0\})$. Also, pick $\tilde{c} \in\left(c, c_{0}\right)$ and $\eta>0$ so small that

$$
\frac{\tilde{c}-c}{|x|^{2}} \geq p u_{t}^{p-1}+\lambda_{1} \quad \text { a.e. in } B_{\eta}
$$

Let $z=\phi_{1}-M|x|^{-\tilde{a}}$ and $M=\left\|\phi_{1}\right\|_{L^{\infty}\left(\partial B_{\eta}\right)} \eta^{\tilde{a}}$ ( $-\tilde{a}$ being the greater root of $P(X)=X(X-1)+(n-1) X+\tilde{c}=0)$. Then,

$$
\left\{\begin{align*}
-\Delta z-\frac{\tilde{c}}{|x|^{2}} z \leq 0 & \text { in } B_{\eta}  \tag{4.3}\\
z \leq 0 & \text { on } \partial B_{\eta}
\end{align*}\right.
$$

Testing (4.3) against $z^{+}$(which is permitted since $z^{+} \in H_{0}^{1}\left(B_{\eta}\right)$ ),

$$
\phi_{1} \leq M \cdot|x|^{-\tilde{a}} \quad \text { a.e. in } B_{\eta} .
$$

With $\tilde{c}$ close enough to $c$, it follows that $p u_{t}^{p-1} \phi_{1}+\lambda_{1} \phi_{1} \leq C|x|^{-a-2+\epsilon}$, for some $\epsilon>0$. Let $\zeta \in H_{0}^{1}(\Omega)$ be the solution of

$$
\left\{\begin{align*}
-\Delta \zeta-\frac{c}{|x|^{2}} \zeta & =|x|^{-a-2+\epsilon} & & \text { in } \Omega  \tag{4.4}\\
\zeta & =0 & & \text { on } \partial \Omega
\end{align*}\right.
$$

As in the proof of Lemma 1.5,

$$
\begin{equation*}
0 \leq \phi_{1} \leq C \zeta \leq C|x|^{-a} \quad \text { a.e. in } \Omega \tag{4.5}
\end{equation*}
$$

Next, we prove that there exists $0<t_{1} \leq t_{0}$ such that $u_{t}$ is stable for $t<t_{1}$.
Fix $b \in\left(a, a^{\prime}\right)$ such that $p b<b+2$ and $b+a(p-1)<a+2$, and define $F: X \times \mathbb{R} \rightarrow Y$, by

- $X$ is the space of functions $v \in C(\bar{\Omega} \backslash\{0\})$ such that there exist a constant $C>0$ and a function $g \in C(\bar{\Omega} \backslash\{0\})$ satisfying $|v| \leq C|x|^{-b},|g| \leq C|x|^{-b-2}$ and

$$
\left\{\begin{aligned}
-\Delta v-\frac{c}{|x|^{2}} v & =g & & \text { in } \Omega \\
v & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

in the weak sense.
$X$ is a Banach space for the norm $\|v\|_{X}=\left\||x|^{b} v\right\|_{L^{\infty}}+\left\||x|^{b+2} g\right\|_{L^{\infty}}$

- $Y=\left\{f \in C(\bar{\Omega} \backslash\{0\}):|x|^{b+2} f \in L^{\infty}(\Omega)\right\},\|f\|_{Y}=\left\||x|^{b+2} f\right\|_{L^{\infty}}$
- $F(v, t)=-\Delta v-\frac{c}{|x|^{2}} v-|v|^{p}-t f$

Observe that $F$ is well defined with our choice of $b$, that $F \in C^{1}$ and that $F\left(u_{t}, t\right)=0$. Also $L:=F_{u}(0,0)$ is an isomorphism between $X$ and $Y$. Indeed $L$ is injective by Lemma 1.4 and surjective with continuous inverse by Lemma 1.3. These facts and a global form of the implicit function theorem (see e.g. Cor. 3 in [BN]) imply the existence of a maximal $t_{1}>0$ such that $t \rightarrow u_{t}$ is a $C^{1}$ map from $\left(0, t_{1}\right)$ to $X$ and $F_{u}\left(u_{t}, t\right) \in I \operatorname{so}(X, Y)$.

In particular, since $\phi_{1} \in X, \lambda_{1}\left(u_{t}\right) \neq 0$ for $t<t_{1}$. It can also be shown that $t \rightarrow \lambda_{1}\left(u_{t}\right)$ is continuous : if $\tau_{n} \rightarrow \tau<t_{1}$ and $\lambda_{1}^{n}$ and $\phi_{1}^{n}$ are the corresponding eigenvalues and eigenfunctions with $\left\|\phi_{1}^{n}\right\|_{L^{2}}=1$, looking carefully at the previous claim, we obtain that $\phi_{1}^{n}$ is bounded in $H_{0}^{1}(\Omega)$ and that

$$
0 \leq \phi_{1}^{n} \leq C|x|^{-a}
$$

Passing to a subsequence, it is then easy to show that $\lambda_{1}^{n} \rightarrow \lambda_{1}\left(u_{\tau}\right)$ and therefore that $\lambda_{1}$ is continuous.

Hence, since $\lambda_{1}(0)>0$ and $\lambda_{1}$ cannot vanish, we have $\lambda_{1}>0$ for $t<t_{1}$.
We now prove that $t_{1}=t_{0}$. If not, we would have for $t_{1}<t<t_{0}$,

$$
\begin{aligned}
& -\Delta\left(u_{t}-u_{t_{1}}\right)-\frac{c}{|x|^{2}}\left(u_{t}-u_{t_{1}}\right)-p u_{t_{1}}^{p-1}\left(u_{t}-u_{t_{1}}\right)= \\
& \\
& \quad u_{t}^{p}-u_{t_{1}}^{p}-p u_{t_{1}}^{p-1}\left(u_{t}-u_{t_{1}}\right)+\left(t-t_{1}\right) f \geq\left(t-t_{1}\right) f
\end{aligned}
$$

And testing against $\phi_{1}$, solution of (4.1) with $t_{1}$ in place of $t$, we would obtain

$$
0 \geq\left(t-t_{1}\right) \int_{\Omega} f \phi_{1}
$$

which is impossible. Hence, $t_{1}=t_{0}$.
Next, we prove that if $v$ is another stable $H_{0}^{1}(\Omega)$ solution then it must coincide with $u_{t}$.

Suppose indeed $v$ is another $H_{0}^{1}(\Omega)$ solution such that $\lambda_{1}(v) \geq 0$. Then $v \geq u_{t}$ and

$$
\begin{aligned}
\int_{\Omega} p v^{p-1}\left(v-u_{t}\right)^{2} & \leq \int_{\Omega}\left|\nabla\left(v-u_{t}\right)\right|^{2}-\int_{\Omega} \frac{c}{|x|^{2}}\left(v-u_{t}\right)^{2} \\
& \leq \int_{\Omega}\left(v^{p}+t f-u_{t}^{p}-t f\right)\left(v-u_{t}\right)
\end{aligned}
$$

So that,

$$
\int_{\Omega}\left(v-u_{t}\right)\left(v^{p}-u_{t}^{p}-p v^{p-1}\left(v-u_{t}\right)\right) \geq 0
$$

Since $u \rightarrow u^{p}$ is strictly convex and $v \geq u_{t}$, we must have $v=u_{t}$.
Finally, stability of strong extremal solutions is determined through the following proposition :

## Proposition 4.1.

Suppose that $0<c<c_{0}$ and $1<p<p_{0}$.

If $u$, the minimal solution of $\left(P_{t_{0}, p}\right)$, solves the problem in the strong sense then

$$
\lambda_{1}(u)=0
$$

Proof.

Arguing by contradiction, our general strategy is to use the implicit function theorem to extend the curve $t \rightarrow u_{t}$ of minimal solutions of $\left(P_{t, p}\right)$ beyond $t_{0}$ if $\lambda_{1}(u)>0$.

More precisely assume that $\lambda_{1}(u)>0$ and define $F: X \times \mathbb{R} \rightarrow Y$ as before. If we can prove that $F_{u}\left(u, t_{0}\right) \in I s o(X, Y)$, the implicit function theorem will yield the desired contradiction.

We first claim that $F_{u}\left(u, t_{0}\right)$ is injective. If not, there would be a weak solution $\phi_{1} \in X$ of

$$
\left\{\begin{aligned}
-\Delta \phi_{1}-\frac{c}{|x|^{2}} \phi_{1} & =p u^{p-1} \phi_{1} & & \text { in } \Omega \\
\phi_{1} & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

Since $b+a(p-1)<a+2, u^{p-1} \phi_{1} \in L^{\frac{2 n}{n+2}}(\Omega)$ and, using the methods of Proposition $0.1, \phi_{1}$ is an $H_{0}^{1}(\Omega)$ solution. Testing the above equation against $\phi_{1}$ would then imply $J\left(\phi_{1}\right)=0$, which contradicts $\lambda_{1}(u)>0$. Thus $F_{u}\left(u, t_{0}\right)$ is injective.

Next we prove that $F_{u}\left(u, t_{0}\right)$ is surjective.
First observe that $L:=F_{u}(0,0)$ is an isomorphism between $X$ and $Y$. Indeed $L$ is injective by Lemma 1.4 and surjective with continuous inverse by Lemma 1.3.

Let $Z:=\left\{f:|x|^{-a(p-1)} f \in Y\right\}$ and define $K \in \mathcal{L}(Z)$ by

$$
K:\left\{\begin{array}{lll}
Z & \rightarrow \quad Y & \rightarrow Z \\
\phi & \mapsto p u^{p-1} \phi & \mapsto \quad L^{-1}\left(p u^{p-1} \phi\right)
\end{array}\right.
$$

$K$ is compact in $Z$. Indeed if $\left\{\phi_{n}\right\}$ is a bounded sequence in $Z$ then $u_{n}:=K \phi_{n}$ is bounded in $X$, by continuity of $L^{-1}$. It follows from standard elliptic theory that up to a subsequence, $u_{n} \rightarrow u$ uniformly on compacts of $\bar{\Omega} \backslash\{0\}$ for some $u \in X$. Also, letting $\gamma=2-a(p-1)>0$, we have for $\epsilon>0$ small

$$
\left\|u_{n}-u\right\|_{Z} \leq C\left\|u_{n}-u\right\|_{L^{\infty}\left(\Omega \backslash B_{\epsilon}\right)}+\epsilon^{\gamma}\left\|u_{n}-u\right\|_{L_{-b}^{\infty}} \leq C\left(\left\|u_{n}-u\right\|_{L^{\infty}\left(\Omega \backslash B_{\epsilon}\right)}+\epsilon^{\gamma}\right)
$$

so that

$$
\limsup _{n \rightarrow \infty}\left\|u_{n}-u\right\|_{Z} \leq C \epsilon^{\gamma}
$$

Letting $\epsilon \rightarrow 0$, we obtain that $K$ is compact in $Z$.
With these notations, our problem reduces to showing that $I d-K$ is surjective. By Fredholm's alternative, we just need to prove that $I d-K$ is injective. Now if for some $\phi \in Z, \phi=K \phi$ then $\phi \in X$ by definition of $K$, and $F_{u}\left(u, t_{0}\right) \phi=0$. But we just showed that $F_{u}\left(u, t_{0}\right)$ is injective so $\phi \equiv 0$.

## 5. What happens in the extremal case $t=t_{0}$ ?

In this section, we look at two specific sets of conditions on $c, p, f$ and $\Omega$.
In one case, the minimal solution $u$ of $\left(P_{t_{0}, p}\right)$ solves the problem in the strong sense. It then follows from Proposition 5.1 that $\lambda_{1}(u)=0$.

In the other case, the minimal solution $u$ is not a strong one and its singularity at the origin is worse than $|x|^{-a}$. Moreover, $u$ is stable, i.e., $\lambda_{1}(u)>0$.

## Situation 1.

Suppose $\Omega=B_{1}, c<c_{0}$ close to $c_{0}, f$ radial and $p>1$ close to 1 .
Then $u$, the minimal solution of $\left(P_{t_{0}, p}\right)$, solves the problem in the strong sense and $\lambda_{1}(u)=0$. Furthemore, $u=u(r)$ is radial and

$$
\begin{aligned}
u & =r^{-a} w & & \text { where } w \in C[0,1] \cap C^{\infty}(0,1] \\
w^{\prime} & \sim m r^{-a(p-1)+1} & & \text { for some } m<0 \text { as } r \rightarrow 0 \\
w^{\prime} & <0 & & \text { in }(0,1)
\end{aligned}
$$

Proof.
We suppose for simplicity that $f \equiv 1$.
First, we note that for any rotation of the space $A \in S O(n, \mathbb{R}), u \circ A$ is a solution of ( $P_{t_{0}, p}$ ) and since $u$ is minimal, we must have $u \leq u \circ A$. This inequality holds almost everywhere in $B_{1}$, hence for almost all $y=A^{-1} x$ with $x \in \mathbb{R}^{n}$ so that $u$ must be radial.

Next, define $\alpha:=\sqrt{(n-2)^{2}-4 c}$ and for $r \in(0,1)$,

$$
\begin{equation*}
\Phi(u)(r):=\frac{r^{-a}}{\alpha} \int_{0}^{1} s^{1+\alpha+a} u^{p}(s)\left[\max (s, r)^{-\alpha}-1\right] d s+\frac{t_{0}}{2 n+c}\left[r^{-a}-r^{2}\right] \tag{5.1}
\end{equation*}
$$

In view of Lemma 1.5, $\Phi(u)(r)$ is well defined for $r \neq 0$ and it follows from Lebesgue's theorem that $w:=r^{a} \Phi(u) \in C(0,1]$.

Using Lebesgue's theorem again, it is also true that $w \in C^{1}(0,1]$ and that for $r \in(0,1]$,

$$
w^{\prime}(r)=-r^{-1-\alpha}\left(\int_{0}^{r} s^{1+\alpha+a} u^{p}(s) d s\right)-(2+a) \frac{t_{0}}{2 n+c} r^{1+a}
$$

Using the fundamental theorem of calculus, $w$ is twice differentiable a.e. in $(0,1)$ and
$w^{\prime \prime}(r)=-r^{a} u^{p}(r)-(1+\alpha) \frac{1}{r}\left[w^{\prime}(r)+(2+a) \frac{t_{0}}{2 n+c} r^{1+a}\right]-(2+a)(1+a) \frac{t_{0}}{2 n+c} r^{a}$
So that

$$
\begin{equation*}
-\left(w^{\prime \prime}+(1+\alpha) \frac{1}{r} w^{\prime}\right)=r^{a} u^{p}(r)+t_{0} r^{a} \quad \text { a.e. in }(0,1) \tag{5.2}
\end{equation*}
$$

Using the fundamental theorem of calculus again, this equation also holds in the sense of distributions in $(0,1)$. Furthermore, since $u$ is a weak solution of $\left(P_{t_{0}, p}\right)$, it is not hard to see that $\tilde{w}:=r^{a} u$ solves (5.2) in $\mathcal{D}^{\prime}(0,1)$.

So if $z=\tilde{w}^{\prime}-w^{\prime}$, it follows from (5.2) and this last remark that

$$
z^{\prime}+(1+\alpha) \frac{1}{r} z=0 \quad \text { in } \mathcal{D}^{\prime}(0,1)
$$

And by a straightforward computation, we see that

$$
\left[r^{1+\alpha} z\right]^{\prime}=0 \quad \text { in } \mathcal{D}^{\prime}(0,1)
$$

Hence $z=A r^{-(1+\alpha)}$ for some $A \in \mathbb{R}$ and, for some $B \in \mathbb{R}$,

$$
\begin{equation*}
\tilde{w}=w+\frac{A}{\alpha} r^{-\alpha}+B \tag{5.3}
\end{equation*}
$$

Since $w$ is $C^{1}$ away from $r=0$ (and hence, so must be $\tilde{w}$ ), we must have, on the one hand, using the boundary condition of $\left(P_{t_{0}, p}\right)$ and equation (5.1), that $w(1)=\tilde{w}(1)=0$ and $B=0$ and on the other hand that $u$ is $C^{1}$ away from the origin. Bootstrapping this result with the help of (5.1) and (5.3), it follows that $w, \tilde{w} \in C^{\infty}(0,1]$.

Let us now prove that $A=0$. Suppose by contradiction that $A>0$ and let $u_{1}(x):=|x|^{-a} w(|x|)$ for $x \in B_{1}$. Then,

$$
\begin{aligned}
-\Delta u_{1}-\frac{c}{|x|^{2}} u_{1} & =u^{p}+t_{0}=\left[u_{1}+\frac{A}{\alpha}\left(|x|^{-a^{\prime}}-|x|^{-a}\right)\right]^{p}+t_{0} \\
& \geq u_{1}^{p}+t_{0}
\end{aligned}
$$

This equation holds at every $x \neq 0$ and also in the weak sense, as integration by parts on $B_{1} \backslash B_{\epsilon}$ with $\epsilon \rightarrow 0$ shows. But then $u_{1}$ would be a nonnegative supersolution of problem $\left(P_{t_{0}, p}\right)$, contradicting minimality of $u$.

We have just shown that $A \leq 0$. We now prove that $A=0$. Recall that

$$
\begin{equation*}
\tilde{w}^{\prime}(r)=-r^{-1-\alpha} \int_{0}^{r} s^{1+\alpha-a} u^{p}(s) d s-(2-a) \frac{t_{0}}{2 n+c} r^{1+a}-A r^{-1-\alpha} \tag{5.4}
\end{equation*}
$$

By Hopf's boundary lemma, $u^{\prime}(1)=\tilde{w}^{\prime}(1)<0$. We claim that

$$
\tilde{w}^{\prime}(r)<0 \quad \text { for all } r \in(0,1]
$$

Suppose not and let $r_{0}=\sup \left\{r \in(0,1): \tilde{w}^{\prime}(r)=0\right\}$. Then $\tilde{w}^{\prime}<0$ on $\left(r_{0}, 1\right]$ and (5.2) implies that

$$
\tilde{w}^{\prime \prime}\left(r_{0}\right)=-(1+\alpha) \frac{1}{r_{0}} \tilde{w}^{\prime}\left(r_{0}\right)-\left(r_{0}^{a} u^{p}\left(r_{0}\right)+t_{0} r_{0}^{a}\right)<0
$$

So $\tilde{w}$ has a local maximum at $r_{0}$. Suppose by contradiction that $\tilde{w}$ has another critical point and let $r_{1}<r_{0}$ so that

$$
\tilde{w}^{\prime}\left(r_{1}\right)=0 \quad \text { and } \quad \tilde{w}^{\prime}(r)>0 \quad \text { for } r \in\left(r_{1}, r_{0}\right)
$$

From (5.2), it follows as before that $\tilde{w}^{\prime \prime}\left(r_{1}\right)<0$ and $r_{1}$ would be a local maximum of $\tilde{w}$, contradicting $\tilde{w}^{\prime}>0$ on $\left(r_{1}, r_{0}\right)$.

Hence $\tilde{w}$ has an absolute maximum at $r_{0}$ and must therefore be bounded, which forces $A=0$.

But then, using (5.4), $\tilde{w}^{\prime}(r)<0$ in $(0,1]$, contradicting $\tilde{w}^{\prime}\left(r_{0}\right)=0$.
So, we have proved that $\tilde{w}^{\prime}<0$ in $(0,1]$.
From (5.4), it follows that if $A<0, \tilde{w}^{\prime}(r)=-A r^{-1-\alpha}(1+o(1))$ as $r \rightarrow 0$ and we cannot have at the same time $\tilde{w}^{\prime}<0$ and $A<0$. Hence $A=0$.

So far we know that :

$$
\begin{gathered}
\tilde{w}=w \\
w^{\prime}<0 \quad \text { in }(0,1]
\end{gathered}
$$

We now prove that $u \leq C r^{-a}$. From equation (5.1), we already know that $u=\Phi(u) \leq C r^{-a-\alpha}$. Plugging this result into (5.1) again, we only need to show that the right-hand-side integral is bounded as $r \rightarrow 0$, which holds as soon as

$$
1-a(p-1)-\alpha p>-1
$$

This last condition is satisfied for $\alpha p$ small and in particular when $c$ is close to $c_{0}$ and $p$ close to 1 . This result, combined with (5.4) yields the asymptotic behaviour of $w^{\prime}$ at the origin.

Finally, by Proposition 4.1, we have that $\lambda_{1}(u)=0$.

When $p$ is chosen close to the critical exponent $p_{0}$, the minimal solution $u$ of $\left(P_{t_{0}, p}\right)$ may become more singular than when $t<t_{0}$, in such a way that $u^{p-1}$ has a singularity at the origin of same order as $\frac{1}{|x|^{2}}$ :

## Situation 2.

Suppose $0<c<c_{0}$ and $p$ close to $p_{0}$. Then there exists a smooth nonnegative nonzero data $f$ such that $u$, the minimal solution of $\left(P_{t_{0}, p}\right)$, is stable and such that, near the origin,

$$
u=m|x|^{-\gamma} \quad \text {, where } m>0 \quad \text { and } \quad \gamma=\frac{2}{p-1}>a>0
$$

Proof.
We adapt a proof given in [D].
Let $v=|P(-\gamma)|^{\frac{1}{p-1}}|x|^{-\gamma}$, where $P(X)=X(X-1)+(n-1) X+c$.
Then, $-\Delta v-\frac{c}{|x|^{2}} v=v^{p}$ in $\mathbb{R}^{n}$ and, since when $p \rightarrow p_{0}, \gamma \rightarrow a$, we may assume that $v \in H^{1}$.

Lemma 5 in [D] constructs a function $\psi \in C^{\infty}(\bar{\Omega})$ with the following properties:
(1) $\psi \geq 0$ in $\bar{\Omega}$
(2) $\Delta \psi+\frac{c}{|x|^{2}} \psi \geq 0$ in $\Omega$
(3) $\psi \equiv 0$ in a neighbourhood of 0 , and
(4) $\psi=v$ on $\partial \Omega$

We then let $u=v-\psi$ and see that

$$
\begin{aligned}
-\Delta u-\frac{c}{|x|^{2}} u & =-\Delta v-\frac{c}{|x|^{2}} v+\Delta \psi+\frac{c}{|x|^{2}} \psi \\
& =v^{p}+\Delta \psi+\frac{c}{|x|^{2}} \psi \\
& \geq 0
\end{aligned}
$$

and $u=0$ on $\partial \Omega$, so, by Lemma 1.1 say, $u \geq 0$.
Taking $f=\Delta \psi+\frac{c}{|x|^{2}} \psi+v^{p}-u^{p}$, we then have

$$
-\Delta u-\frac{c}{|x|^{2}} u=u^{p}+f
$$

Observe that $f \geq 0$ and is smooth since $u \leq v$ and $u \equiv v$ near the origin.
Next, we prove that $\lambda_{1}(u)>0$. Given $\phi \in H_{0}^{1}(\Omega)$,

$$
\begin{aligned}
\int_{\Omega} p u^{p-1} \phi^{2} & \leq \int_{\Omega} p v^{p-1} \phi^{2} \\
& =p|P(-\gamma)| \int_{\Omega} \frac{\phi^{2}}{|x|^{2}} \\
& \leq \int_{\Omega}|\nabla \phi|^{2}-\int_{\Omega} \frac{c}{|x|^{2}} \phi^{2}-\epsilon \int_{\Omega} \phi^{2}
\end{aligned}
$$

The last inequality holds, using Hardy's inequality (0.3), provided $c+p|P(-\gamma)|<$ $c_{0}$ and $\epsilon>0$ small. This condition is readily satisfied since as $p \rightarrow p_{0}, \gamma \rightarrow a$ and $P(-\gamma) \rightarrow P(-a)=0$. Hence, we get that $\lambda_{1}(u) \geq \epsilon>0$.

We still need to prove that, for our choice of $f, t_{0}=1$ and $u$ is the minimal solution of $\left(P_{t_{0}, p}\right)$.

If $u_{1}$ denotes the minimal solution of $\left(P_{1, p}\right)$, it is clear that $0 \leq u_{1} \leq u$, hence $u_{1}^{p} \leq \frac{C}{|x|^{2}}$ and using this inequality and $\left(P_{1, p}\right), u_{1} \in H_{0}^{1}(\Omega)$.

Since $\lambda_{1}(u) \geq 0$, it follows that

$$
\begin{aligned}
\int_{\Omega} p u^{p-1}\left(u-u_{1}\right)^{2} & \leq \int_{\Omega}\left|\nabla\left(u-u_{1}\right)\right|^{2}-\int_{\Omega} \frac{c}{|x|^{2}}\left(u-u_{1}\right)^{2} \\
& \leq \int_{\Omega}\left(u^{p}+f-u_{1}^{p}-f\right)\left(u-u_{1}\right)
\end{aligned}
$$

So that,

$$
\int_{\Omega}\left(u-u_{1}\right)\left(u^{p}-u_{1}^{p}-p u^{p-1}\left(u-u_{1}\right)\right) \geq 0
$$

Since $u \rightarrow u^{p}$ is strictly convex and $u \geq u_{1}$, we must have $u=u_{1}$. And since $u$ is not a strong solution of $\left(P_{1, p}\right)$, we must have $1=t_{0}$.

## 6. The case $c=c_{0}$

When $c=c_{0}$, the operator $-\Delta-\frac{c}{|x|^{2}}$ is no longer coercive in $H_{0}^{1}(\Omega)$. However, one can still make use of the improved Hardy inequality (see [BV] or [VZ])

$$
\begin{equation*}
\int_{\Omega}|\nabla u|^{2}-c_{0} \int_{\Omega} \frac{u^{2}}{|x|^{2}} \geq C(\Omega) \int_{\Omega} u^{2} \quad \text { for all } u \in C_{c}^{\infty}(\Omega) \tag{6.1}
\end{equation*}
$$

to define a new Hilbert space $H$ in which the operator is coercive, even when $c=c_{0}$.

## Definition.

$H$ is the space obtained by completing $C_{c}^{\infty}(\Omega)$ with respect to the norm

$$
\|u\|_{H}^{2}:=\int_{\Omega}|\nabla u|^{2}-c_{0} \int_{\Omega} \frac{u^{2}}{|x|^{2}}
$$

By analogy with the case $c<c_{0}$, an $H$ solution $u$ will be one such that $u^{p} \in H^{*}$ and such that the equation holds in the sense of Lax-Milgrams lemma in $H$.

We now list the modifications needed to prove Theorem 1 when $c=c_{0}$. When no proof is given, just replace $H_{0}^{1}(\Omega)$ by $H$ in the original demonstration.

Lemma 1.1'. Lemma 1.1 still holds if $c=c_{0}, H_{0}^{1}(\Omega)$ is replaced by $H$ and $H^{-1}$ by $H^{*}$, the dual of $H$.
Proof. Only the proof of (1.3) needs to be clarified in this setting.
Let $f \in H^{*}, f \geq 0$ and $u \in H$ be the corresponding solution of (1.1).
By definition of $H$, there exists a sequence $\left\{u_{n}\right\}$ in $C_{c}^{\infty}(\Omega)$ converging to $u$ in $H$. Letting $f_{n}=-\Delta u_{n}-\frac{c}{|x|^{2}} u_{n}$, it follows that $f_{n} \in H^{-1}(\Omega)$ and $f_{n} \rightarrow f$ in $H^{*}$.

Now, $u_{n} \in H_{0}^{1}(\Omega) \Rightarrow u_{n}^{-} \in H_{0}^{1}(\Omega)$ and integrating the equation satisfied by $u_{n}$ against $u_{n}^{-}$yields

$$
-\left\|u_{n}^{-}\right\|_{H}^{2}=\left\langle f_{n}, u_{n}^{-}\right\rangle_{H *, H}
$$

To pass to the limit in this last equation, we just need to prove that $\left\{u_{n}^{-}\right\}$remains bounded in $H$. But

$$
\begin{align*}
\left\|u_{n}^{-}\right\|_{H}^{2} & =\int_{\Omega}\left|\nabla u_{n}^{-}\right|^{2}-c_{0} \int_{\Omega} \frac{\left(u_{n}^{-}\right)^{2}}{|x|^{2}} \\
& =\int_{\Omega}\left|\nabla u_{n}^{-}\right|^{2}-c_{0} \int_{\Omega} \frac{u_{n}^{2}}{|x|^{2}}+\int_{\Omega} \frac{c_{0}}{|x|^{2}}\left(u_{n}^{+}\right)^{2}  \tag{6.2}\\
& \leq \int_{\Omega}\left|\nabla u_{n}^{-}\right|^{2}-c_{0} \int_{\Omega} \frac{u_{n}^{2}}{|x|^{2}}+\int_{\Omega}\left|\nabla u_{n}^{+}\right|^{2}=\int_{\Omega}\left|\nabla u_{n}\right|^{2}-c_{0} \int_{\Omega} \frac{u_{n}^{2}}{|x|^{2}} \\
& =\left\|u_{n}\right\|_{H}^{2}
\end{align*}
$$

where we've used (0.3) in the inequality.

Proposition 0.1'. Proposition 0.1 still holds when $c=c_{0}$ and $H_{0}^{1}(\Omega)$ solutions are replaced by $H$ solutions.

Proof. Suppose first that $u$ is a strong solution of $\left(P_{t, p}\right)$.
Let $\zeta_{n} \in C_{c}^{\infty}(\Omega \backslash\{0\})$ be such that $0 \leq \zeta_{n} \leq 1,\left|\Delta \zeta_{n}\right| \leq C n^{2}$ and

$$
\zeta_{n}=\left\{\begin{array}{lll}
0 & \text { if } & |x| \leq 1 / n \text { and } \delta(x) \leq 1 / n \\
1 & \text { if } & |x| \geq 2 / n \text { and } \delta(x) \geq 2 / n
\end{array}\right.
$$

Multiplying $\left(P_{t, p}\right)$ by $u \zeta_{n}$ and integrating by parts, it follows that

$$
\begin{aligned}
\int_{\Omega}\left(u^{p}+t f\right) u \zeta_{n} & =-\int_{\Omega} \Delta u u \zeta_{n}-\int_{\Omega} \frac{c}{|x|^{2}} u^{2} \zeta_{n} \\
& =\int_{\Omega}|\nabla u|^{2} \zeta_{n}-\int_{\Omega} \frac{c}{|x|^{2}} u^{2} \zeta_{n}+\int_{\Omega} u \nabla u \nabla \zeta_{n}
\end{aligned}
$$

Since $u \leq C|x|^{-a}$ and $p<p_{0}, u^{p} \leq C|x|^{-a-2+\epsilon}$, for some $\epsilon>0$, so that the first integral in the above equation is bounded by $C \int_{\Omega}|x|^{-2 a-2+\epsilon} \leq C$ whereas $\left|\int_{\Omega} u \nabla u \nabla \zeta_{n}\right|=\left|\frac{1}{2} \int_{\Omega} u^{2} \Delta \zeta_{n}\right| \leq C n^{2} \int_{1 / n<|x|<2 / n}|x|^{-2 a} \leq C$ as $n \rightarrow \infty$.

Hence $\int_{\Omega}|\nabla u|^{2} \zeta_{n}-\int_{\Omega} \frac{c}{|x|^{2}} u^{2} \zeta_{n} \leq C$ and $u \in H$. Approximating $u \in H$ by smooth functions and integrating by parts in $\Omega \backslash B_{\epsilon}$ with $\epsilon \rightarrow 0$, it follows that $u$ is a weak solution of $\left(P_{t, p}\right)$. For $u$ to be an $H$ solution, we only need to prove the following :

Claim. Suppose $u$ is a weak solution satisfying the estimate $u \leq C|x|^{-a}$. Then

$$
u^{p} \in H^{*}
$$

For $\phi \in C_{c}^{\infty}(\Omega), 1<q<2$, it follows from Hölder's inequality that

$$
\left.\left|\int_{\Omega}\right| x\right|^{-a p} \phi \left\lvert\, \leq\left(\int_{\Omega}|x|^{-(a p-1) \frac{q}{q-1}}\right)^{\frac{q-1}{q}}\left\|\frac{\phi}{|x|}\right\|_{L^{q}}\right.
$$

On the one hand, since $p<p_{0}$, the integral in the right hand side will be finite if $q$ is chosen close enough to 2 .

On the other hand, using Hardy's inequality in $L^{q}$ and the inclusion $H \hookrightarrow W_{0}^{1, q}$ (see section 4 of [VZ]),

$$
\left\|\frac{\phi}{|x|}\right\|_{L^{q}} \leq C\|\phi\|_{W_{0}^{1, q}} \leq C\|\phi\|_{H}
$$

and $u^{p} \in H^{*}$.
Hence, strong solutions are also $H$ solutions.
Showing that $H$ solutions are weak solutions is similar to the case $c<c_{0}$, whereas, starting from a weak solution $u$, we observe as above that $u^{p} \in H^{*}$ and define $u_{n} \geq 0$ to be the minimal weak solution of

$$
\left\{\begin{aligned}
-\Delta u_{n}-\frac{c-1 / n}{|x|^{2}} u_{n} & =u_{n}^{p}+t f & & \text { in } \Omega \\
u_{n} & =0 & & \text { on } \partial \Omega
\end{aligned}\right.
$$

$u$ is a supersolution of this equation so $u_{n}$ is well defined and $0 \leq u_{n} \leq u \leq$ $C|x|^{-a}$. By Proposition 0.1 (case $c<c_{0}$ ), it follows that $u_{n} \in H_{0}^{1}(\Omega)$ and testing in the above equation against $u_{n}$,

$$
\left\|u_{n}\right\|_{H}^{2} \leq\left(\left\|u_{n}^{p}\right\|_{H^{*}}+C\right)\left\|u_{n}\right\|_{H}
$$

Letting $n \rightarrow \infty$, we get $u \in H$. Since $H \hookrightarrow W_{0}^{1, q}$ for $1 \leq q<2$, elliptic regularity can be applied to complete the proof.

Lemma 1.3'. Lemma 1.3 still holds when $c=c_{0}$
Lemma 1.4'. Lemma 1.4 still holds when $c=c_{0}$

Lemma 1.5'. Lemma 1.5 still holds when $c=c_{0}$
Proof. We assume first that $u$ is the minimal (weak) solution of (1.20) and can therefore be written as the pointwise limit of an increasing sequence $\left\{u_{\epsilon}\right\}$, where $u_{\epsilon}$ solves (1.20) with $c-\epsilon$ in place of $c$. Since Lemma 1.5 can be applied to $u_{\epsilon}$, an argument of monotone convergence yields the result.

If $u$ isn't minimal, the above discussion yields all the results up to the conclusion of Step 6 in the original proof. That step can be applied -as is- in our context, which finishes the proof.

## 2.1' : Existence of a solution of $\left(\mathbf{P}_{\mathbf{t}, \mathrm{p}}\right)$ for $\mathbf{1}<\mathbf{p}<\mathbf{p}_{\mathbf{0}}$, small $\mathbf{t}>\mathbf{0}$.

For simplicity, we assume without loss of generality that $\Omega \subset B_{\frac{1}{2}}$. Consider $w=A|x|^{-a}\left(\ln \frac{1}{|x|}\right)^{\frac{1}{4 p}}$, with $A>0$ to be fixed later. Then $-\Delta w-\frac{c}{|x|^{2}} w=$ $\frac{4 p-1}{16 p^{2}} A|x|^{-a-2}\left(\ln \frac{1}{|x|}\right)^{\frac{1}{4 p}-2}$ and $w$ will be a supersolution of $\left(P_{t, p}\right)$ as soon as

$$
\left\{\begin{array}{l}
\frac{4 p-1}{32 p^{2}} A|x|^{-a-2}\left(\ln \frac{1}{|x|}\right)^{\frac{1}{4 p}-2} \geq A^{p}|x|^{-a p}\left(\ln \frac{1}{|x|}\right)^{p / 4} \\
\frac{4 p-1}{32 p^{2}} A|x|^{-a-2}\left(\ln \frac{1}{|x|}\right)^{\frac{1}{4 p}-2} \geq t f
\end{array}\right.
$$

The first inequality amounts to

$$
A \leq C \min _{r \in(0,1 / 2]}\left\{r^{-a-2+p a}\left(\ln \frac{1}{r}\right)^{\frac{1}{4 p}-2-p / 4}\right\}^{\frac{1}{p-1}}
$$

and the second to

$$
t \leq C \cdot A \min _{r \in(0,1 / 2]}\left\{r^{-a-2}\left(\ln \frac{1}{r}\right)^{\frac{1}{4 p}-2}\right\}
$$

Under these conditions, $w$ is a supersolution of $\left(P_{t, p}\right)$.
We now just have to construct a supersolution in $H$.
Let $w_{1}$ be a smooth extension inside $\Omega$ of $\left.w\right|_{\partial \Omega}$ such that $w_{1}=w$ in $\Omega \backslash B_{1 / 4}$ and $w_{1}=0$ in $B_{1 / 8}$. Next, we let $g=\Delta w_{1}+\frac{c}{|x|^{2}} w_{1}$ and construct $z \in H$ solving (2.1) and $w_{2}=z+w_{1}$ solving (2.2).

We would like to show that $w_{2} \geq 0$ and remark that $w_{2}^{-} \in H$. Indeed, let $\phi_{k} \in C_{c}^{\infty}(\Omega) \rightarrow z$ in $H$. Then $\left(\phi_{k}+w_{1}\right)^{-} \in H_{0}^{1}(\Omega) \subset H$ and

$$
\begin{aligned}
\left\|\left(\phi_{k}+w_{1}\right)^{-}\right\|_{H}^{2} & =\int_{\left\{\phi_{k}+w_{1}<0\right\}}\left(\left|\nabla\left(\phi_{k}+w_{1}\right)\right|^{2}-\frac{c}{|x|^{2}}\left(\phi_{k}+w_{1}\right)^{2}\right) \\
& \leq\left\|\phi_{k}\right\|_{H}^{2}+C+2 \int_{\left\{\phi_{k}+w_{1}<0\right\}}\left(\nabla \phi_{k} \cdot \nabla w_{1}-\frac{c}{|x|^{2}} \phi_{k} w_{1}\right) \\
& \leq\left\|\phi_{k}\right\|_{H}^{2}+C+\frac{1}{2}\left\|\left(\phi_{k}+w_{1}\right)^{-}\right\|_{H}^{2}
\end{aligned}
$$

Hence $\left\|\left(\phi_{k}+w_{1}\right)^{-}\right\|_{H} \leq C$ and passing to the limit (in the weak topology and for a subsequence), it follows that $w_{2}^{-} \in H$.

Letting $\psi_{k} \in C_{c}^{\infty}(\Omega) \rightarrow w_{2}^{-}$in $H$, integration by parts then yields

$$
\left(w_{2} \mid \psi_{k}\right)_{H}=\int_{\Omega}\left(\nabla z \nabla \psi_{k}-\frac{c}{|x|^{2}} z \psi_{k}\right)+\int_{\Omega}\left(\nabla w_{1} \nabla \psi_{k}-\frac{c}{|x|^{2}} w_{1} \psi_{k}\right)=\int_{\partial \Omega} \psi_{k} \partial_{\nu} w
$$

and letting $k \rightarrow \infty$ in $H$,

$$
\left\|w_{2}^{-}\right\|_{H}^{2}=\left(w_{2} \mid w_{2}^{-}\right)_{H}=\int_{\partial \Omega} \partial_{\nu} w w^{-}=0
$$

Hence $w_{2} \geq 0$.
Finally, letting $\tilde{w}=w-w_{2}=w-z-w_{1}$, we only need to prove that $\tilde{w} \in H$ and the rest of the proof will remain unchanged. Since $z \in H$, it is enough to show that $w-w_{1} \in H$.

If $H(\omega)$ denotes the space $H$ relative to the open set $\omega$ of $\mathbb{R}^{n}$, it has been shown in [VZ] (see 5.2) that $f$ defined for $0<r<r_{0}<1$ by

$$
f(r)=r^{-a}(\ln (1 / r))^{\alpha}
$$

and continued smoothly up to the boundary of the ball $B_{1}$, where $f=0$, belongs to $H\left(B_{1}\right)$ as long as $\alpha<1 / 2$.
$\left.\left(w-w_{1}\right)\right|_{B_{1 / 4}}$ precisely satisfies these conditions, hence belongs to $H\left(B_{1 / 4}\right)$. Since $w-w_{1} \equiv 0$ in $\Omega \backslash B_{1 / 4}$, it follows that $w-w_{1} \in H(\Omega)$.

## $2.3^{\prime}:$ Case $p \geq p_{0}:$ blow-up.

By Proposition 2.1, we just need to prove that $\left(P_{t, p}\right)$ has no weak solution if $p \geq p_{0}$. Assume by contradiction there exists one and call it $u$. If we apply Lemma 1.5 with $u^{p}+t f$ in place of $f$, it follows that

$$
\int_{\Omega} u^{p}|x|^{-a} \delta(x)<\infty \quad \text { and } \quad u \geq m|x|^{-a} \quad \text { a.e. near the origin. }
$$

This is impossible since near the origin,

$$
|x|^{-a} u^{p} \geq m|x|^{-a(p+1)} \geq m|x|^{-n}
$$

Lemma 3.1'. Lemma 3.1 still holds when $c=c_{0}$
Lemma 3.2'. Lemma 3.2 still holds when $c=c_{0}$ and $t<t_{0}$

Theorem 2'. Theorem 2 still holds when $c=c_{0}$ with $H_{0}^{1}(\Omega)$ solutions replaced by $H$ solutions.

Proof. For $t<t_{0}, u_{t}$ the strong minimal solution of $\left(P_{t, p}\right)$ can be written as the monotone limit of $u_{n}$, where $u_{n}$ is the strong minimal solution of the same problem with $c$ replaced by $c_{n}=c-1 / n$.By our analysis in the case $c<c_{0}$, we know that $\lambda_{1}\left(u_{n}\right)>0$. Passing to the limit, we easily get that $\lambda_{1}\left(u_{t}\right) \geq 0$.

To obtain a strict inequality, it is enough to show that $t \rightarrow \lambda_{1}\left(u_{t}\right)$ is a strictly decreasing function. It should be clear from its definition that $t \rightarrow \lambda_{1}\left(u_{t}\right)$ is nonincreasing.

Suppose that $\lambda_{1}\left(u_{t}\right)=\lambda_{1}\left(u_{s}\right)$ for some $s \leq t$. Call $\phi_{1}^{t}$ and $\phi_{1}^{s}$ the corresponding eigenfunctions, which can be constructed as in the case $c<c_{0}$. Then,

$$
\begin{aligned}
\lambda_{1}\left(u_{t}\right) & =\int_{\Omega}\left|\nabla \phi_{1}^{t}\right|^{2}-\int_{\Omega} \frac{c}{|x|^{2}} \phi_{1}^{t^{2}}-\int_{\Omega} p u_{t}^{p-1} \phi_{1}^{t^{2}} \\
& \leq \int_{\Omega}\left|\nabla \phi_{1}^{s}\right|^{2}-\int_{\Omega} \frac{c}{|x|^{2}} \phi_{1}^{s 2}-\int_{\Omega} p u_{t}^{p-1} \phi_{1}^{s 2} \\
& \leq \int_{\Omega}\left|\nabla \phi_{1}^{s}\right|^{2}-\int_{\Omega} \frac{c}{|x|^{2}} \phi_{1}^{s 2}-\int_{\Omega} p u_{s}^{p-1} \phi_{1}^{s 2} \\
& =\lambda_{1}\left(u_{s}\right)
\end{aligned}
$$

Hence,

$$
\begin{aligned}
\int_{\Omega}\left|\nabla \phi_{1}^{s}\right|^{2}-\int_{\Omega} \frac{c}{|x|^{2}} \phi_{1}^{s 2}-\int_{\Omega} p u_{t}^{p-1} \phi_{1}^{s 2} & = \\
& \int_{\Omega}\left|\nabla \phi_{1}^{s}\right|^{2}-\int_{\Omega} \frac{c}{|x|^{2}} \phi_{1}^{s 2}-\int_{\Omega} p u_{s}^{p-1} \phi_{1}^{s 2}
\end{aligned}
$$

and

$$
u_{t}=u_{s} \quad \text {,which implies } \quad t=s
$$

Hence $u_{t}$ is a stable solution of $\left(P_{t, p}\right)$.
To prove that $u_{t}$ is the only stable $H$ solution, we can argue exactly as in the case $c<c_{0}$.

The results of section 5 extend in the following way (we skip the proof) :
Situation 1'. Suppose $c=c_{0}, \Omega=B_{1}$, $f$ radial and $1<p<p_{0}$.
Then $u$, the minimal solution of $\left(P_{t_{0}, p}\right)$, solves the problem in the strong sense.
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