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Sonochemical Approach to the Synthesis of Fe₃O₄@SiO₂ Core—Shell Nanoparticles with Tunable Properties

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agnetic iron oxide nanoparticles (NPs) have attracted much research interest over the recent years because of their unique physicochemical properties and great potential for various biomedical applications. In several pioneering works magnetic NPs were claimed as an effective tool for magnetically assisted biomolecule separation, biochemical sensing, NMR imaging, and targeted drug delivery, and magnetic hyperthermia. And magnetic NPs for tissue engineering, termed magnetic force-based tissue engineering (Mag-TE) has been proposed recently.

The requirements for any biomedical application of magnetic colloids include the chemical stability, biocompatibility, strong magnetization, and low coercivity of the dispersed magnetic NPs. In principal, silicacoated magnetite NPs, Fe₃O₄@SiO₂, comply with these requirements. The outer shell of silica not only protects the inner magnetite core from oxidation but also provides sites for surface functionalization with poly(ethylene glycol) and various biomolecules. For in vivo applications magnetic core-shell NPs should have a limiting size in the order of 20-50 nm to be able to migrate across the reticuloendothelial system and to be stealthy toward the mononuclear phagocyte system.3,7

There are essentially two methods for the coating of Fe_3O_4 NPs with silica: by acidic hydrolysis of silicate in aqueous solutions and by a modified Stöber process¹⁰ which consists of the alkaline hydrolysis of tetraethyl orthosilicate (TEOS) in ethanol/water mixture in the presence of Fe_3O_4 NPs. Deposition of silica from silicate solution

ABSTRACT In this study, we report a rapid sonochemical synthesis of monodisperse nonaggregated Fe₃O₄@SiO₂ magnetic nanoparticles (NPs). We found that coprecipitation of Fe(II) and Fe(III) in aqueous solutions under the effect of power ultrasound yields smaller Fe₃O₄ NPs with a narrow size distribution (4—8 nm) compared to the silent reaction. Moreover, the coating of Fe₃O₄ NPs with silica using an alkaline hydrolysis of tetraethyl orthosilicate in ethanol—water mixture is accelerated many-fold in the presence of a 20 kHz ultrasonic field. The thickness of the silica shell can be easily controlled in the range of several nanometers during sonication. Mössbauer spectra revealed that nonsuperparamagnetic behavior of obtained core—shell NPs is mostly related to the dipole—dipole interactions of magnetic cores and not to the particle size effect. Core—shell Fe₃O₄@SiO₂ NPs prepared with sonochemistry exhibit a higher magnetization value than that for NPs obtained under silent conditions owing to better control of the deposited silica quantities as well as to the high speed of sonochemical coating, which prevents the magnetite from oxidizing.

KEYWORDS: magnetite · silica · core—shell nanoparticles · sonochemistry · ultrasound · Mössbauer spectroscopy · FT-IR · XPS

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usually yields relatively big spherical magnetic particles with a mean size in the range of 80-200 nm. 11-13 These particles are useful for biomolecule separation or sensing in vitro but cannot be delivered intravenously. In the seminal work of Philipse and coworkers¹⁴ it was shown that Fe₃O₄@SiO₂ NPs with the mean diameter of 60-120 nm can be obtained by the Stöber process from initial 5-15 nm size Fe₃O₄ precursors only after the bare magnetite particles were pretreated with a small amount of silicate in an aqueous solution. Without this procedure, coating with TEOS in ethanol - agueous ammonia mixture yields magnetite clusters embedded in large silica aggregates. The thin silica layer at the Fe₃O₄ NPs surface reduces the isoelectric point of magnetite and thus maintains the colloidal stability during the second coating. However, the silicate precoating is not a sufficient condition to obtain the uniform

core-shell NPs. It was shown 14 that the concentration of pretreated magnetite particles should be less than 12 $mg \cdot L^{-1}$, otherwise large aggregates of magnetite are formed in TEOS solutions. The Fe₃O₄@SiO₂ NPs prepared by the Stöber process usually contain about six Fe₃O₄ cores even under optimal conditions. 14 Further detailed studies of Fe₃O₄ NPs coating using TEOS^{15–17} revealed the high sensitivity of this process to experimental conditions such as ethanol/water ratio, concentration of ammonia and TEOS, and temperature, etc. Moreover, the coating of Fe₃O₄ NPs with silica using TEOS was found to be a very slow process: from 12 to 48 h of mechanical stirring at room temperature is necessary to obtain uniform a silica shell on the surface of magnetite NPs. Heating of the reacting mixture causes formation of big magnetite-silica aggregates with irregular morphology. To improve the homogeneity of silica-coated iron oxide NPs and to increase the limiting concentration of NPs in reaction medium several authors used a modified Stöber process in reverse microemulsion conditions. 18,19 However, this method also requires at least 20-24 h of aging to obtain the final product and much effort to separate core-shell NPs from the large amount of surfactants associated with the microemulsion system.

In this work, we report the use of sonochemistry for the rapid synthesis of monodispersed Fe₃O₄@SiO₂ NPs with tunable properties using the Stöber process. The term sonochemistry is used to describe chemical processes occurring in liquids under the effect of power ultrasound. For the ultrasonic frequencies usually employed to influence chemical processes (20-60 kHz), the wavelengths produced in the liquid medium are in the range of 7.5 to 3.0 cm or considerably longer than the chemical bond length values. Therefore, the chemical effects of ultrasound are not the result of direct interactions between the molecules and the sonic waves. The origin of sonochemistry derives from acoustic cavitation, that is, formation, growth, and implosive collapse of gas-filled microbubbles in a liquid subjected to ultrasonic irradiation.^{20,21} Transient collapse generates high peak temperatures (~5000 K) within the cavitation bubbles and overheating of the interfacial region around the bubbles (~200 nm) until several hundred degrees, as well as the intense shock waves emanated upon collapse. Since the collapse occurs in less than a microsecond, very high cooling rates, in excess of 10¹⁰ K, are obtained. For these reasons, treatment of solutions with power ultrasound not only offers the possibility to disperse friable solids in the liquids but also provides a unique set of conditions to produce unusual materials from dissolved volatile and nonvolatile precursors. After the pioneering work of Suslick and co-workers on amorphous iron preparation by Fe(CO)₅ sonolysis²² a large number of nanosized metals, metal alloys, oxides, carbides, and sulfides were prepared using sonochemical reactions with volatile

metal—carbonyls in appropriate organic solvents. 20,23 On the other hand, sonochemistry was proven to be effective for NPs insertion into mesoporous matrixes and NPs deposition on ceramic and polymeric supports. 20,23 Magnetite NPs were prepared by sonolysis of Fe(CO)₅ in aqueous medium in the presence of sodium dodecyl sulfate²⁴ and by sonochemical oxidation of Fe(II) in aqueous solutions. 25,26 Very recently it was shown that ultrasonic treatment enhances the uniformity of SiO₂ NPs obtained by TEOS hydrolysis²⁷ and improves the coating of SiO₂ microspheres with FePt alloy²⁸ and FeNi fine particles with silica. 29

This paper reports that the sonochemical approach provides a resolution of several important problems in the synthesis of $\mathrm{Fe_3O_4@SiO_2}$ NPs by the Stöber process related to the necessity of silicate pretreatment, the slow kinetics, and the risk of aggregates formation. Moreover, the thickness of the silica shell over $\mathrm{Fe_3O_4}$ NPs can be easily tuned during sonication which is important in order to avoid magnetite nanoparticles oxidation and agglomeration. It is known that the superparamagnetism of uncoated single domain magnetic particles can vanish because of the intergrain magnetic dipole interaction. In the absence of an external magnetic field the energy of dipole—dipole monodomain attraction is given by 31

$$V(r) = -V_{\text{max}}^2(r)/6kt \tag{1}$$

where $V_{\rm max}^2(r)$ is a maximal attraction energy in a head-to-tail configuration and r is a distance between the centers of two monodomains. If a magnetic core is coated by an insulating shell the dipole—dipole attractions become negligible for some critical diameter of a insulating sphere $\sigma_{\rm B}$ (Boyle diameter):¹⁴

$$\sigma_{\rm B}^3 = \alpha d^6 / (6kt)^{1/2} \tag{2}$$

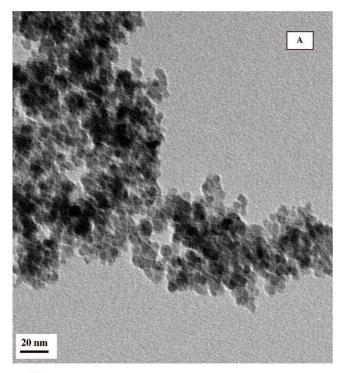
where d is a magnetic core diameter, $\alpha=M^2\pi/72\mu_0$, M is a volume saturation magnetization, and μ_0 is the permeability of a vacuum. From eq 2 it follows, for example, that for d=10 nm the dispersion of Fe₃O₄@SiO₂ NPs should be stable and superparamagnetic if $\sigma \geq 25$ nm. Here, we present the evidence that the superparamagnetic behavior of Fe₃O₄@SiO₂ NPs can be tailored by the silica shell thickness.

RESULTS AND DISCUSSION

Preparation of Fe₃0₄ NPs. The bare Fe_3O_4 NPs were prepared by the well-known Massart's method³² which consists of Fe(III) and Fe(III) coprecipitation in alkaline solutions:

$$2Fe(III) + Fe(II) + 8OH^{-} \rightarrow Fe_3O_4 + 8H_2O$$
 (3)

The only difference in our method was the application of ultrasonic irradiation at 20 kHz instead of mechanical



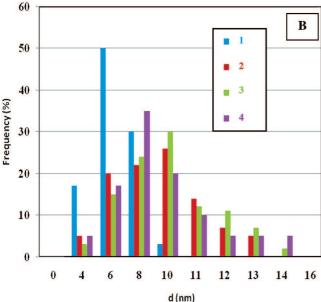


Figure 1. (a) TEM image of Fe $_3$ O $_4$ NPs obtained by ultrasonically assisted coprecipitation. An Fe(II)/Fe(III) mixture was injected into the active cavitation zone. Scale bar = 20 nm. (b) The size histograms of Fe $_3$ O $_4$ NPs obtained under ultrasound ([Fe $_3$ O $_4$] = 1.7 × 10 $^{-2}$ M, [NH $_3$ ·H $_2$ O] = 2 M, [N $_2$ H $_4$ ·H $_2$ O] = 0.01 M, T = 30-32 °C, f = 20 kHz, I = 30 W·cm $^{-2}$, P_{ac} = 0.67 W·mL $^{-1}$) by Fe(II)/Fe(III) mixture injection into the active cavitation zone (1) or in the upper part of the sonoreactor outside of the active cavitation zone (2) and mechanical stirring ([Fe $_3$ O $_4$] = 5 × 10 $^{-3}$ M, [NH $_3$ ·H $_2$ O] = 1 M, room temperature, ref 14 (3); [Fe $_3$ O $_4$] = 4 × 10 $^{-2}$ M, [NH $_3$ ·H $_2$ O] = 1 M, room temperature, ref 33 (4).

agitation during the synthesis. Figure 1 demonstrates the transmission electron microscopy (TEM) image and the size histograms for Fe₃O₄ NPs prepared with and without ultrasound at similar conditions. It can be seen that precipitation under power ultrasound yields par-

ticles that are smaller and have more narrow size distribution as compared to those prepared by mechanical stirring. Particle size histograms reveal that the effect of ultrasound on the magnetite particles size is observed much more clearly if the mixture of Fe(II) and Fe(III) is injected directly into the active cavitation zone of the sonoreactor. A TEM image of polymodal Fe₃O₄ NPs obtained with ultrasound but when the Fe(II)/Fe(III) mixture was injected outside of the active cavitation zone is presented in the Supporting Information. The particle size histogram for this case is very similar to that obtained under mechanical stirring.

The active zone of the sonoreactor is shown in Figure 2. A clearly visible cloudlike region with the maximal concentration of cavitation bubbles has a circular cone shape which is "stuck" to the radiating surface of the titanium horn. Obviously, the maximal effect of ultrasound should be observed in this "conebubble structure".21 Moreover, Fe(II)/Fe(III) coprecipitation is known to be a quasiimmediate process.33,34 Therefore, to observe the effect of ultrasound on magnetite precipitation and the hydrolysis of iron ions and magnetite particles, nucleation should occur in the active zone of acoustic cavitation to avoid interference with coprecipitation in the bulk solution in the absence of cavitation. Recently it was assumed that three effects of sonication contribute to the phenomenon of sonocrystallization:35,36 (i) the local transient heating of a liquid after bubble collapse, (ii) the shock waves generated during bubbles implosion hinder agglomeration, and (iii) the excellent mixing conditions created by acoustic cavitation. All these phenomena allow a reduction of particle size and increase particle size homogeneity owing to control of the local nucleus population.

The Fe2p high resolution XPS spectrum of Fe $_3$ O $_4$ NPs precipitated under ultrasound (Figure 3) reveals that, in contrast to γ -Fe $_2$ O $_3$ (Aldrich) XPS spectrum, it does not contain the charge transfer satellite of Fe2p $_{3/2}$ at 720 eV indicating formation of a mixed oxide of Fe(III) and Fe(III), such as Fe $_3$ O $_4$ NPs surface oxidation during precipitation under ultrasound. The binding energy (E_B) for Fe2p $_{3/2}$ photoelectrons (711.2 eV) of the synthesized NPs is very close to that published in the literature (710.6 eV).

Sonochemical Coating of Fe₃0₄ NPs with Silica. Sonication of Fe₃O₄ NPs suspension in alkaline ethanol—water solutions of TEOS causes rapid coating of the magnetic

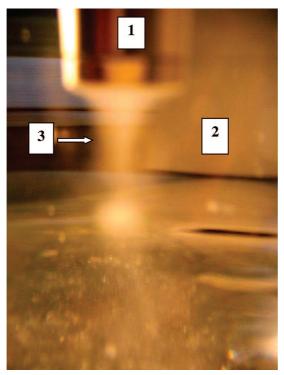


Figure 2. Image of the active cavitation zone of the sonoreactor (f = 20 kHz, $I = 30 \text{ W} \cdot \text{cm}^{-2}$, $P_{\text{ac}} = 0.67 \text{ W} \cdot \text{mL}^{-1}$). (1) Titanium ultrasonic horn, (2) sonicated water in the thermostatted cell under argon, (3) "cone bubble structure" of the active cavitation zone.

core with a silica shell. FT-IR spectra of the solid products demonstrate the increase of absorbance at 1080 cm $^{-1}$ attributed to $\nu_{\rm as}({\rm Si-O-Si})$ vibrations 38 during sonication. The calibration curve for FT-IR quantitative silica analyses is shown in Supporting Information. Figure 4 reveals that the concentration of deposited silica increases linearly with the time of sonochemical treatment indicating zero-order kinetics of SiO $_2$ formation.

Moreover, sonication causes at least 5-fold acceleration of TEOS hydrolysis compared to the silent reaction at similar experimental conditions.

The XPS spectra of Fe $2p_{3/2}$ ($E_B\approx711$ eV) and Si 2p ($E_B\approx103$ eV) core—electrons gives a further proof for the rapid Fe $_3O_4$ NPs coating with silica under ultrasound. Data in Table 1 clearly show the increase of Si 2p peak intensity during sonication. At the same time the peak intensity of Fe $2p_{3/2}$ sharply decreases until its practical disappearance during 3 h of sonication. The XPS surveys of Fe $_3O_4$ @Si O_2 NPs are shown in Supporting Information.

TEM images of Fe₃O₄@SiO₂ NPs (Figure 5) demonstrate the increase of silica shell thickness with sonication time in a good agreement with FT-IR and XPS measurements. According to TEM measurements the silica shell thickness is increased from 1.0-1.5 nm after 1 h of sonication to 3.0-3.5 nm after 3 h of ultrasonic treatment. Also TEM images reveal that the magnetic cores are not collapsed under sonication. Silica particles without magnetic cores are not observed indicating the heterogeneous mechanism of silica nucleation and growth at the magnetite surface. It is worth noting that practically all core-shell particles obtained with ultrasound contain a single magnetic core. By contrast, the TEM image presented in Supporting Information indicates that TEOS hydrolysis in the presence of Fe₃O₄ NPs under mechanical stirring and without silicate pretreatment causes formation of big irregular aggregates with the extremely heterogeneous distribution of silica in full agreement with published data.¹⁴ It should be emphasized that the overheating of the reaction mixture yields big aggregates even under ultrasound. The TEM image presented in the Supporting Information shows the typical cluster of Fe₃O₄ NPs embedded in silica ob-

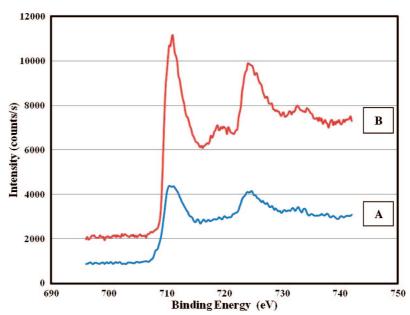


Figure 3. Fe 2p high resolution XPS spectra for (A) Fe₃O₄ NPs obtained under ultrasound and (B) γ-Fe₃O₃ provided by Aldrich.

TABLE 1. Surface Concentration of Fe and Si for Fe₃O₄@SiO₂ NPs Obtained by XPS Technique

	atomic %	
particles	Si2p	Fe 2p _{3/2}
Fe_3O_4	0	15.3
Fe ₃ O ₄ @SiO ₂ , 1 h of ultrasound	13.3	11.8
Fe ₃ O ₄ @SiO ₂ , 3 h of ultrasound	26.5	0.6

TABLE 2. Hydrodynamic Diameter of ${\rm Fe_3O_4@SiO_2}$ NPs as a Function of the Sonication Time Measured Using Dynamic Light Scattering in EtOH

particles	hydrodynamic diameter (nm)	polydispersity index (r.u.)
Fe ₃ 0 ₄ @Si0 ₂ , 1 h of ultrasound	49	0.01
Fe ₃ O ₄ @SiO ₂ , 3 h of ultrasound	53	0.12

tained after 2 h of sonication at 45 °C. It appears that the increase of temperature of the bulk solution causes acceleration of the undesirable homogeneous silica nucleation followed by the aggregation of ${\rm Fe_3O_4@SiO_2}$ and ${\rm SiO_2~NPs.}$

The values of hydrodynamic diameter for ${\rm Fe_3O_4@SiO_2}$ NPs measured with dynamic light scattering (DLS) in ethanol are summarized in Table 2. Particle size distribution curves exhibit only one peak with a small polydispersity index indicating the low extent of ${\rm Fe_3O_4@SiO_2}$ NPs aggregation in solutions. The hydrodynamic diameter of the particles is approximately three times larger than that measured with TEM. This difference may be due to the particles solvation and dynamic association in the liquid. Roughly, such associates contain on average three single ${\rm Fe_3O_4@SiO_2}$ NPs. The slight increase of hydrodynamic diameter with the time of sonication is correlated with the increase of silica shell thickness.

Probable Mechanism of TEOS Sonochemical Hydrolysis. The significant acceleration of TEOS alkaline hydrolysis under the effect of power ultrasound could be understood in light of the two-site model of the sonochemical reactions:³⁹ the dominant site is the cavitation bubble interior gas-phase while the secondary site is a liquid shell surrounding the cavitation bubble. The resonance size of a bubble in the solvents with low viscosity, like water, alkanes, or ethanol, sonicated at 20 kHz is approximately $150-200\mu$ m²¹ and the liquid reaction zone extends approximately 200 nm (500 molecules)³⁹ from the bubble surface. The lifetime of the overheated liquid zone is less than 2μ s after collapse. When the acoustic pressure from the ultrasonic waves decreases, the gas bubble is expanded, and gas or vapors flow into the bubble from the liquid in its neighborhood. At our experimental conditions the cavitation bubbles are filled with argon, ethanol, water, ammonia, and TEOS vapors. The vapor pressure

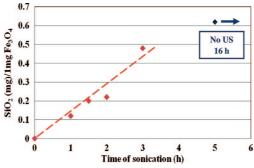


Figure 4. Kinetics of SiO₂ sonochemical deposition on Fe₃O₄ NPs in EtOH/NH₃ · H₂O solutions. [Fe₃O₄] = 7.8 × 10⁻³ M, [TEOS] = 4.0 × 10⁻² M, [NH₃ · H₂O] = 0.225 M, [H₂O] = 12.35 M, f = 20 kHz, I = 30 W · cm⁻², P_{ac} = 0.67 W · mL⁻¹, T = 32 °C, Ar. The blue diamond corresponds to the SiO₂ concentration deposited under silent conditions at intensive mechanical nonmagnetic stirring during 16 h under similar conditions.

of ethanol (bp 78 °C) at 32 °C is much higher than that of TEOS (bp 158 °C). Consequently, the concentration of TEOS molecules within the bubble is negligibly low and the sonochemically driven hydrolysis of TEOS, most probably, occurs in the liquid reaction zone and not in the gas-phase of the cavitation bubble. The Stöber process can be cast as hydrolysis eq 4 and condensation polymerization eq 5 steps:

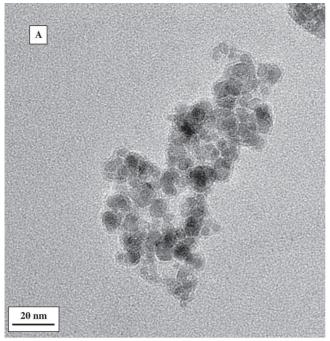
$$Si(OR)_4 + 4H_2O \rightarrow Si(OH)_4 + 4ROH$$
 (4)

$$nSi(OH)_4 \rightarrow n(SiO_2) + 2nH_2O$$
 (5)

In the presence of catalysts, hydrolysis of TEOS is generally faster than the condensation reaction. However, usually they cannot be separated. The overall TEOS hydrolysis in ethanol—water mixture is known to be temperature sensitive:²⁷ heating to 70 °C causes a significant increase in the reaction rate. As a consequence, transient temperature jumps in the liquid shell surrounding the cavitation bubble would lead to strong local acceleration of silica nucleation.

Furthermore, ultrasonically forced oscillations of the cavitation bubble can push ultrasmall particles to the bubble surface during the expansion half-cycle by a process similar to rectified diffusion described above for gases and vapors. This suggests that the liquid reaction zone surrounding the cavitation bubble is enriched with ${\rm Fe_3O_4}$ NPs which play a role of scavenger for silica nucleus. The presumable mechanism of sonochemical coating is shown schematically in Figure 6. Otherwise cold bulk solution avoids the secondary irregular nucleation of silica. Moreover, the significant acceleration of mass transport created by acoustic cavitation improves the homogeneity and monodispersity of core—shell particles.

Tailoring the Magnetic Properties of ${\rm Fe_3O_4@SiO_2}$ NPs. Mössbauer spectra of bare ${\rm Fe_3O_4}$ NPs at room temperature



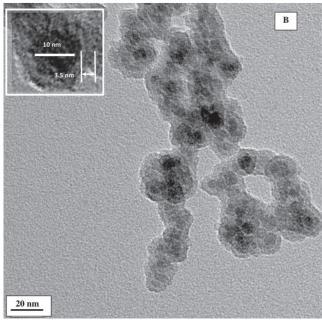


Figure 5. TEM images of Fe_3O_4 @SiO $_2$ NPs after 1 h (A) and 3 h (B) of sonication. The scale bar is 20 nm. Zoomed image of a single core—shell particle is shown in the inset of panel b.

and at 4.2 K (Figure 7) may be interpreted as a superposition of doublet (superparamagnetic phase) and broad sextet (magnetic phase) peaks similar to that observed in other studies. 40,41 These spectra were fitted in two stages: (i) decomposition into Lorentzian lines presuming doublet for the superparamagnetic phase and two sextets for the magnetic phase and (ii) fitting of obtained broad lines using the Hesse—Rubartsch approach. 42 The spectra of Fe₃O₄@SiO₂ NPs at room temperature were treated by the same procedure. Calculated hyperfine parameters of Fe₃O₄ and Fe₃O₄@SiO₂ NPs are listed in Table 3. Then these parameters of the same procedure.

eters were used to estimate the percentage of superparamagnetic and magnetic phase in the prepared materials (Table 4).

Undoubtedly, the superparamagnetic behavior is due to the particle size effect. The large majority of prepared Fe₃O₄ NPs have a diameter in the range of 4-10 nm (Figure 1) which is much less than the theoretical value for the monodomain size of magnetite (~25 nm).⁴¹ It seems to be reasonable to assume that the significant presence of magnetic phase in Mössbauer spectra of bare Fe₃O₄ NPs is related to the dipole – dipole interactions of magnetic nanoparticles according to eq 1. The sextet of magnetic phase which dominates at 4.2 K due to the blocking of Fe(III) – Fe(II) electron exchange does not exhibit a clear separation into tetrahedral (Td) and octahedral (Oh) sites typical for pure magnetite. 43 This effect has been recently interpreted as a formation of nonstoichiometric mixed oxide $(Fe^{3+})_{Td}[(Fe^{2+})_{1-3x}(Fe^{3+})_{1+2x}(\Box)_x]_{Oh}O_4$, where \Box is a vacancy formed due to the partial oxidation of Fe(II).41,44 Although high-resolution XPS (Figure 3) clearly indicates the presence of Fe(II) at the surface of NPs some surface oxidation of magnetite cannot be completely excluded. Mössbauer spectrum will be, in this case, similar to that of nonoxidized Fe₃O₄ NPs, however, with a different area ratio for the subspectra. In practice, that leads to line broadening.

Figure 8 reveals that Mössbauer spectra of Fe_3O_4 NPs are strongly influenced by silica coating. Data in Table 4 clearly indicate that the coating of Fe_3O_4 NPs with silica causes the increase in the yield of superparamagnetic phase in an accordance with equations 1 and 2. It can be thus concluded that nonsuperparamagnetic behavior of Fe_3O_4 @SiO $_2$ NPs observed with Mössbauer spectroscopy is mostly related to the dipole—dipole interactions of the magnetic cores and not to the particle size effect. This conclusion is of practical importance since it presumes principally superparamagnetic behavior of Fe_3O_4 @SiO $_2$ NPs in colloids, where dipole—dipole attractions are much weaker than those in magnetic solids.

Magnetization **M** and coercivity **H**_c values obtained with SQUID technique at room temperature are summarized in Table 5. Corresponding magnetization curves for Fe₃O₄ and Fe₃O₄@SiO₂ NPs are shown in Supporting Information. It is seen that as the silica shell thickness increases, the observed **M** value is decreased, but its value per unit of magnetite is practically not influenced by coating. This confirms that sonochemical coating does not change the magnetic cores of NPs. The magnetization value of bare Fe₃O₄ NPs obtained in this work is close to that published in the literature 16 and is almost 2-fold less than the M value for micrometer-size range magnetite.⁴³ It is well-known that such decrease in magnetization is due to the effect of increased thermal fluctuation near the particle surface or to the magnetically disordered surface

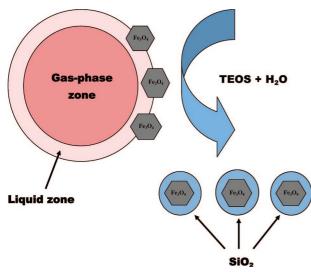


Figure 6. Schematic illustration of Fe₃O₄ NPs sonochemical coating with silica.

TABLE 3. Hyperfine Parameters of Mössbauer Spectra for Fe_3O_4 and Fe_3O_4 @SiO₂ NPs at Room Temperature

- 3 - 4	3-42				
	$^{\delta}_{(mm\cdot s^{-1})}$	$\Gamma \ \ (\text{mm} \cdot \text{s}^{-\text{1}})$	ϵ (mm·s ⁻¹)	Δ^a (mm·s $^{-1}$)	<i>Н</i> ^b (Т)
Magnetic Td [Fe ³⁺]					
Fe ₃ 0 ₄	0.29	0.35	0.001		7
Fe ₃ 0 ₄ @Si0 ₂ , 1 h ultrasound	0.29	0.50	-0.06		34
Fe ₃ O ₄ @SiO ₂ , 3 h ultrasound	0.29	0.50	0.001	-23.4	
	Magi	netic Oh [Fe³+	, Fe ²⁺]		
Fe ₃ O ₄	0.66	0.50	0.001		13
Fe ₃ 0 ₄ @SiO ₂ , 1 h ultrasound	0.66	0.50	0.001		36
Fe ₃ O ₄ @SiO ₂ , 3 h ultrasound	0.66	0.50	0.001		24.6
Superparamagnetic Td [Fe ³⁺]					
Fe ₃ 0 ₄	0.29	0.30		0.85	
Fe ₃ 0 ₄ @Si0 ₂ , 1 h ultrasound	0.29	0.40		1.43	
Fe ₃ O ₄ @SiO ₂ , 3 h ultrasound	0.29	0.30		0.93	
Superparamagnetic Oh [Fe ³⁺ , Fe ²⁺]					
Fe ₃ 0 ₄	0.66	0.30		0.90	
Fe ₃ 0 ₄ @Si0 ₂ , 1 h ultrasound	0.66	0.40		1.97	
Fe ₃ O ₄ @SiO ₂ , 3 h ultrasound	0.66	0.30		1.19	

^aMean value of quadrupole splitting. ^bMean value of hyperfine splitting.

TABLE 4. Magnetic and Superparamagnetic Phase Concentrations for ${\rm Fe_3O_4}$ and ${\rm Fe_3O_4}$ @SiO $_2$ Nanopowders Calculated from Mössbauer Spectra

particles	magnetic (%)	superparamagnetic (%)
Fe_3O_4	61	39
Fe ₃ O ₄ @SiO ₂ (1 h ultrasound)	48	52
Fe ₃ O ₄ @SiO ₂ (3 h ultrasound)	37	63

formed as a result of the large surface-to-volume ratio associated with the fine particle size. $^{\rm 33}$ Table 5 shows that ${\rm Fe_3O_4@SiO_2}$ NPs obtained with sonochemistry exhibit a much higher magnetization value than that obtained under mechanical stirring. This could be assigned to better control of the silica quantities deposited under ultrasound as well as to the high speed of sonochemical coating, which prevents the magnetite from oxidizing.

All particles obtained in this work demonstrate a very low coercivity close to that of the theoretical value for superparamagnetic particles ($\mathbf{H_c} \leq 5 \text{ mT}$)⁴⁵ and much lower than that for the micrometer-size range magnetite. In contrast to Mössbauer spectra, the coercivity remains nearly constant upon silica coating. This phenomenon could be attributed to the high magnetocrystal-

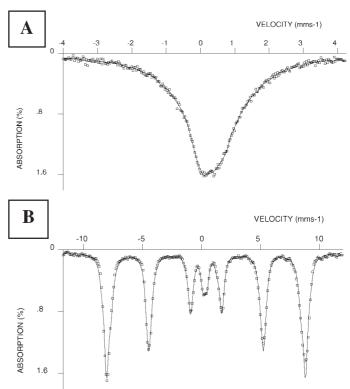


Figure 7. Mössbauer spectra of ${\rm Fe_3O_4}$ NPs at room temperature (A) and at 4.2 K (B).

TABLE 5. Magnetic Properties of ${\rm Fe_3O_4@SiO_2}$ NPs at Room Temperature

particles	M (emu · g ⁻¹)	M^a (emu · g ⁻¹ Fe ₃ 0 ₄)	H _c (mT)
Fe ₃ 0 ₄ this work	51	51	8
Fe ₃ 0 ₄ , 16 (d~15 nm)	55	55	
Fe ₃ O ₄ bulk ⁴³	92	92	120 - 400
Fe ₃ 0 ₄ @Si0 ₂ , 1 h of ultrasound	48	54	6
Fe ₃ 0 ₄ @Si0 ₂ , 3 h of ultrasound	33	50	6
Fe ₃ O ₄ @SiO ₂ ¹⁶	2-7		

 $^{^{}o}\text{Calculated}$ using FT-IR data and presuming insignificant influence of magnetite oxidation.

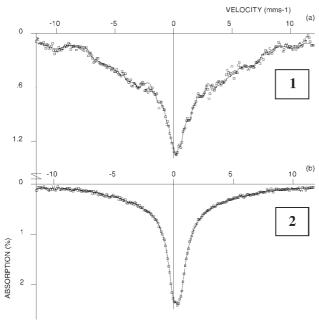


Figure 8. Mössbauer spectra of Fe_3O_4 @Si O_2 NPs at room temperature obtained after 1 (A) and 3 h (B) of sonochemical coating.

line anisotropy of magnetite. The coercivity is defined as a measure of the magnetic field strength that is required to achieve changes of magnetization direction in a material. The observed $\mathbf{H_C}$ value is a combination of many anisotropy mechanisms, such as magnetocrystalline anisotropy, surface anisotropy, and interparticle interactions. It was reported that for small magnetite particles magnetocrystalline anisotropy is the dominant form of anisotropy.⁴³ Silica coating does not lead to any

change in the shape and the size of magnetite cores and, therefore, would have little influence on the coercive force of the core—shell particles. A similar effect was observed recently for CoFe₂O₄@SiO₂ NPs also with high magnetocrystalline anisotropy.⁴⁶ Apparently, Mössbauer spectra exhibit higher sensitivity to the interparticle dipole—dipole magnetic interactions than SQUID measurements.

CONCLUSION

In this article, we reported a sonochemical approach to the synthesis of Fe₃O₄@SiO₂ core-shell NPs. First of all, it should be emphasized that coprecipitation of Fe(II) and Fe(III) in the ultrasonic field allows us to obtain smaller Fe₃O₄ NPs with narrow size distribution compared to those precipitated under mechanical stirring. Furthermore, we revealed that power ultrasound significantly accelerates alkaline TEOS hydrolysis in an alcohol-water mixture and in the presence of Fe₃O₄ NPs, hampers agglomeration, and improves the homogeneity of Fe₃O₄@SiO₂ NPs. It was shown that the thickness of the silica shell can be easily tailored with ultrasound in the range of several nanometers. Mössbauer spectroscopic study shows that silica coating improves the superparamagnetic behavior of magnetite nanopowder. From these measurements it follows that the presence of a magnetic component in Mössbauer spectra is attributed to dipole-dipole interactions between Fe₃O₄. Interestingly, the coercivity of magnetic NPs remains nearly constant upon silica coating which would be explained by the high magnetic magnetocrystalline anisotropy of magnetite.

EXPERIMENTAL SECTION

Sonochemical Setup. Ultrasonic treatment was performed in a thermostatically controlled tightly closed batch reactor with a volume of 50 mL. The sonoreactor was equipped with a titanium horn having 1 cm² of irradiating surface area and piezoelectric transducer supplied by a 20 kHz generator (Vibra Cell 600). The horn was immersed reproducibly below the surface of the sonicated liquid. The image of the experimental setup is available in Supporting Information. The temperature in the reactor during sonolysis was maintained using a Huber Unistat Tango thermocryostat. The experiments were performed at a steady-state temperature of 30-32 °C obtained after approximately 15 min of the sonication. Argon (<1 O₂ ppm) was bubbled at a rate of 100 $mL\cdot min^{-1}$ for about 30 min before sonication and during the ultrasonic treatment. The ultrasonic intensity, I (W·cm⁻²), and the specific absorbed acoustic power, P_{ac} (W·mL⁻¹), transmitted to the solution were measured by the thermal probe method as described recently for the sonoreactors with similar geometry. 47 Typically the values of I and $P_{\rm ac}$ were equal to 30 W \cdot cm $^{-2}$ and 0.67 $W \cdot mL^{-1}$, respectively.

Preparation of Fe₃0₄ NPs. All aqueous solutions were prepared using deionized water (Milli-Q 18 M Ω). Ethanol, hydrochloric acid, and ammonia solutions were of highest purity. A freshly prepared mixture of 1.5 mmol FeCl₃ · 6H₂O (97%, Aldrich) and 0.75 mmol FeCl₂ · 4H₂O (99%, Aldrich) in 5 mL of 0.05 M HCl was rapidly injected *via* a fine plastic tube to 40 mL of 2 M ammonia solution containing 0.01 M of hydrazine (98%, Aldrich) under power ultrasound at 30–32 °C in an argon flow. Hydrazine reduced the redox potential of solution preventing Fe(II) from

oxidizing. The injection was performed into the active zone of the sonoreactor *via* a fine plastic tube. Solids with a distinctive black color of magnetite were precipitated immediately after injection. The mixture was sonicated for 5 min. The sediment was transferred to a nitrogen filled glovebox, separated with a permanent magnet, and washed twice with argon purged 0.1 M NaCl solution and pure EtOH.

Preparation of Fe₃O₄@SiO₂ Core—Shell NPs. A suspension of 0.35 mmol freshly prepared Fe₃O₄ NPs in 30 mL of EtOH was sonicated for 15 min under argon, and then a chilled solution of 0.7 mL concentrated ammonia in 9.3 mL of water was added to the sonoreactor. This suspension was sonicated again for 15 min, and then a solution of 1.8 mmol TEOS (≥99%, Aldrich) in 5 mL of chilled EtOH was rapidly injected to the active zone of cavitation. The time of sonochemical treatment with TEOS was varied from 1 to 3 h at a steady-state temperature about 32 °C. A very stable black colored magnetic fluid was obtained even after 1 h of sonication. Silica-coated magnetite NPs were removed shortly after preparation by centrifugation at 15000 rpm for 30 min at a temperature of 10 °C. The precipitate was rapidly washed three times with chilled EtOH, and the final solids were redispersed in EtOH with power ultrasound ($I = 30 \text{ W} \cdot \text{cm}^{-2}$, $P_{ac} = 0.67$ $W \cdot mL^{-1}$, $\tau = 15$ min). Solid samples were obtained by EtOH evaporation at the reduced pressure in the entry-load chamber of an inert atmosphere glow box. A control experiment without ultrasound was performed under similar conditions with intensive mechanical nonmagnetic stirring in the presence of argon.

Characterization of Prepared Materials. Low-resolution TEM images were obtained with a Philips CM120 electron microscope

(120 keV) equipped with a USC1000 SSCCD $2k \times 2k$ Gatan camera. A drop of tested colloid in EtOH was deposited on each carbon-coated copper grid (300 mesh) and dried in air.

The hydrodynamic diameter of ${\rm Fe_3O_4@SiO_2}$ NPs in ethanol colloidal solutions was measured by using the quasi-elastic light scattering technique with a Malvern Zetasizer Nano ZS device.

Fourier transform infrared spectra (FT-IR) were recorded in KBr pellets with a Spectrum One Perkin-Elmer spectrometer. Calibration for the quantitative analysis of SiO $_2$ in the samples was performed at 1080 cm $^{-1}$ (ν_{as} Si-O-Si) using pure SiO $_2$ in a KBr matrix. This calibration is shown in Supporting Information. One can see that absorbance at 1080 cm $^{-1}$ is a linear function of SiO $_2$ content in the range of 0-0.6 mg of SiO $_2$ in the sample. Each value of absorbance was obtained as an average of three independent measurements.

XPS analysis was performed with an ESCALAB-220i-XL (THERMO-ELECTRON, VG Company) device. Photoemission was stimulated by a nonmonochromatized Mg $K\alpha$ source (1253.6 eV) for uncoated Fe $_3O_4$ NPs and by a monochromatized Al $K\alpha$ radiation (1486.6 eV) for more conductive Fe $_3O_4$ @SiO $_2$ samples. An area of about 150 μm diameter was analyzed. The XPS spectra of Fe2p and Si2p were collected at a pass energy of 20 eV. Binding energies of these spectra were referenced to the C 1s binding energy set at 284.6 eV. The samples were pressed on to an indium support in a nitrogen-filled glovebox and then were put into the entry-load chamber to pump for approximately 12 h at room temperature under a pressure of about 10^{-7} Pa to minimize surface contamination. Small amounts of activated carbon fine powder were added to the samples to improve their conductivity.

The hysteresis cycle of powdered materials was measured using a QPMS SQUID magnetometer in the -4.5 to 4.5 T **H** range at room temperature. Particles were stabilized in epoxy resin matrix prior to measurements.

Mössbauer spectroscopy (MS) studies were carried out using a conventional constant acceleration HALDER-type spectrometer equipped with a 57 Co:Rh source. Metallic α -Fe was used as a reference. Powdered samples were placed in tightly closed cells within a nitrogen-filled glovebox. MS spectra were collected at room temperature and at T = 4.2 K with a helium cryostat. The MS spectra were refined in two steps. In the first treatment, Lorentzian peaks were assumed, and the position (isomer shift, δ), amplitude, and width of each peak were refined. This preliminary calculation allowed the determination of experimental hyperfine parameters for the various iron sites present in the samples. The second computation allowed the analysis of spectra in terms of hyperfine field distributions $P(\mathbf{H})$ and quadrupolar splitting distribution $P(\Delta)$ using the method of Hesse and Rubartsch. 42 This method is often used for disordered compounds with a strong line broadening and non-Lorentzian line

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Supporting Information Available: Details of the sonochemical device, the TEM images of Fe_3O_4 and Fe_3O_4 @SiO $_2$ NPs, calibration curve for SiO_2 quantitative analysis with FT-IR, XPS spectra of Fe_3O_4 and Fe_3O_4 @SiO $_2$ NPs, and magnetization curves for Fe_3O_4 and Fe_3O_4 @SiO $_2$ NPs. This material is available free of charge via the Internet at http://pubs.acs.org.

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