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Optical gain measurements in porous silicon planar waveguides codoped by erbium and ytterbium ions at 1.53 μm

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The on-off optical gain measurements as a function of the pump power were performed on porous silicon planar waveguides codoped by erbium and ytterbium ions. These measurements were obtained for different ratios of Yb concentration to Er concentration. The highest value of the gain was reached when the Yb concentration is three times higher than that of Er at a moderate 980 nm pump power value equal to 70 mW. Optical losses measurements have been performed on these waveguides and were equal to 2.1 dB/cm and an internal gain of about 6.4 dB/cm was obtained. © 2007 American Institute of Physics. [DOI: 10.1063/1.2789185]

Er doped materials are of great interest in optical communications technology as they can be used as erbium doped waveguide amplifiers to provide an optical gain over short propagation distances at the standard telecommunication wavelength of 1.5 μm . They can be used to compensate for coupling losses, waveguide losses, or the power division in optical splitter. All kinds of Er doped silicon nanostructures such as porous silicon,^{1,2} Si nanocrystals embedded in SiO₂ matrix,^{3,4} and Si/SiO₂ superlattices^{5,6} have been widely studied.

Among them, porous silicon (PS) was used as a host material for erbium (Er) atoms to obtain a high Er concentration [in the order of 10²⁰ cm⁻³ (Refs. 7 and 8)] because of its high specific surface. The electrochemical method is usually used to obtain a homogeneous Er doping into porous silicon layer.⁹ The same method can also be used to incorporate ytterbium (Yb). Thermal treatments are then necessary to diffuse and optically activate Er and Yb ions in order to observe rare earth luminescence. The Yb–Er codoping permits to obtain luminescent structures from silicon with a high radiative efficiency by direct energy transfer from the excited state of Yb to the state of Er.¹⁰ Optical gain measurements have still not been reported in the PS waveguides doped by Er or Er–Yb ions.

In this letter, the on-off optical gain was studied as a function of the pump power for different doping [Yb]/[Er] ratios performed on PS planar waveguides. The internal optical gain at 1.53 μm was then deduced. These results have permitted to optimize the Yb concentration incorporation in comparison with Er concentration to obtain a waveguide amplifier.

Porous silicon planar waveguides were prepared by electrochemical etching of a (100)-oriented heavily doped (4–6 m Ω cm resistivity) P⁺-type silicon substrate. The electrolyte was made by mixing HF (50%), ethanol, and H₂O (2:2:1). Two-layer (a guiding and a cladding layer) planar waveguides were formed by applying current densities of 50 and 80 mA/cm², respectively. The thickness of the guiding layer was set to 3 μm and that of the cladding layer was

5 μm to avoid the radiating mode from the silicon substrate.

Then, the Er and Yb ions were electrochemically deposited onto the porous silicon waveguides (PSWGs) at a constant current density of 0.1 mA/cm² for 10 min. The Er–Yb ions solutions were prepared by varying Yb concentrations whereas the Er concentration was kept constant. The term τ corresponds to the ratio of the Yb concentration to the one of Er in solution. Different values of τ were carried out ($\tau=1, 2, 3, \text{ or } 5$). Thermal treatments for ion diffusion and optical activation were performed on the codoped PSWG using a three-step process. First, the structure of PSWG was stabilized by a preoxidation at low temperature (300 °C) in wet O₂, and then the PSWGs were oxidized at 900 °C in wet O₂. Lastly, the temperature was raised to 1000 °C in N₂ gas to diffuse and to activate the rare earth ions. The samples were annealed in O₂ to oxidize the silicon skeleton and thus to convert them into porous silica.¹¹ The refractive indices of the guiding and cladding layers decreased after oxidation from 1.78 and 1.53 (before codoping and oxidation) down to 1.33 and 1.31, respectively.¹² Moreover, the Er and Yb concentrations in the layers were estimated by energy dispersive x-ray (EDX) analysis performed on the sample cross section. The concentration profiles were homogeneous in depth and their values are reported in Table I and confirm the ratio between Er and Yb concentrations incorporated in the different samples.

Gain measurements performed on these waveguides were characterized with a pump at 980 nm and a signal at 1530 nm. The source which was used for the pump is a 980 nm λ Highwave C band erbium doped fiber broadening source. A tunable laser system operating between 1470 and

TABLE I. Er and Yb concentrations determined by EDX analysis performed on the sample cross section as a function of the ratio (τ) of Yb concentration to Er concentration in solution.

Type of waveguides	$\tau=1$	$\tau=2$	$\tau=3$	$\tau=5$
[Er] 10 ²⁰ cm ⁻³	2.8	1.4	1.1	0.9
[Yb] 10 ²⁰ cm ⁻³	2.8	3	3.5	3.8

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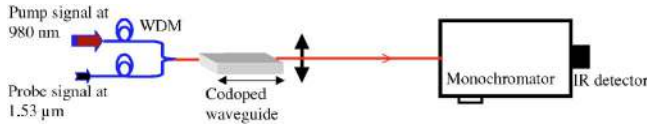


FIG. 1. (Color online) Experimental setup for gain measurements [wavelength-division multiplexer (WDM): 980/1550 nm].

1565 nm (Agilent 81680A) with a resolution of 0.001 nm was used as a probe signal at 1530 nm. To increase the fiber to waveguide coupling efficiency at the input face of the sample, a lens fiber with a 4 μm mode field diameter was used. Both pump and probe signals were coupled by using an output 980/1550 wavelength-division multiplexer coupler. The end-face guided light was collected with a microscope objective and was imaged into a monochromator interfaced with an IR digital camera which was allowed single spot measurements as a function of a fixed wavelength (Fig. 1). The IR digital camera response is constant in the used range of wavelength. The intensity of the light spot was recorded with a digital camera by integrating the values of pixel corresponding to the spot. All the spectra (output signals) were collected as a function of the wavelength and all the measurements were performed at a fixed wavelength with 11 nm resolution. Moreover, propagation losses were measured by studying the scattered light from the surface of the waveguide^{13,14} at near IR range.

On-off gain was determined from the ratio of the output signal as measured by the IR digital camera when both the pump and probe beams were coupled to the PS waveguide to the output signal in the absence of the pump. We had well checked well that this output signal is not due to amplified spontaneous emission (ASE) phenomena by measuring the intensity of the ASE that is not significant (without incoming signal).

The on-off gain measurements as a function of the pump power are reported in Fig. 2 for different τ values. For a given τ , the gain increases with the pump power and begins to saturate. Firstly, for the low pump powers ($P_{\text{pump}} < 50$ mW), on-off gain values increase slightly for the different codoping ratios. Then the gain values increase faster. This evolution is due to the variation of the absorption cross section of Yb ions at 980 nm. Lastly, for the high pump powers ($P_{\text{pump}} > 50$ mW), the gain values slightly increase until saturation. However, the saturation value (g_{sat}) depends on the τ value. The g_{sat} value increases when τ varies from 1 to 3 and strongly decreases for $\tau=5$. These values are in good agreement with the observed results by photoluminescence performed on codoped waveguides.¹⁵ This increase in on-off gain with the one of Yb concentration is related to the enhanced direct energy transfer from Yb excited states to the $I_{13/2}$ state of Er. The highest value was reached for $\tau=3$ at a power pump equal to 70 mW. The thermal oxidation of porous layers and the thermal treatment in N_2 at 1000 °C per-

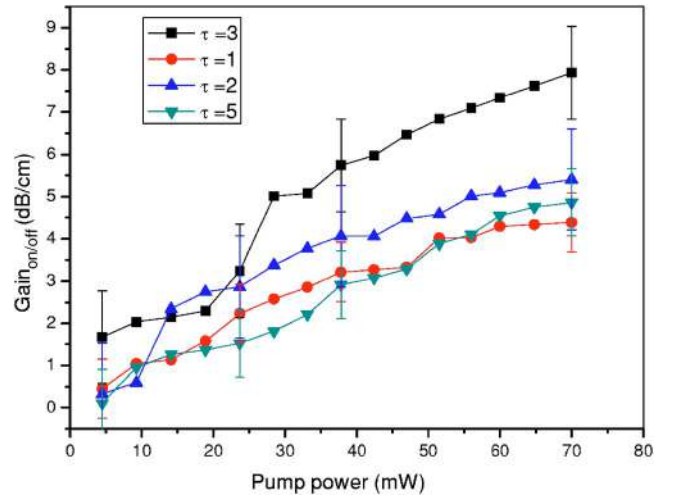


FIG. 2. (Color online) Experimental on-off gain at 1.53 μm as a function of the pump power at 980 nm for the different ratios (τ) of Yb concentration to Er concentration. The input signal power is equal to 100 μW .

mit to diffuse and to activate the Er ions by forming active complexes such as Si–Er–O. We think that up to $\tau=3$, the quantity of active complexes increases, increasing the gain by energy transfer from Yb ion to Er ion. For $\tau=5$, the gain decreases because the high concentration of Yb induces a strong interaction between neighbor ions of Yb. So, a low radiative efficiency by direct energy transfer at 1.53 μm is obtained.

From these on-off gains and from loss measurements, the internal gain for the codoped waveguides was inferred. The measurements of on-off gain have been performed for several lengths of sample for $\tau=3$. The mean value is equal to 8.5 dB/cm. The propagation losses were measured for the undoped waveguides and for the different codoped waveguides. These values are reported in Table II. The origins of the losses are principally due to surface scattering.¹⁶ The increase in losses after codoping is related to the presence of Er and Yb aggregates on the sample surface¹⁷ that increases the roughness interface, inducing additional losses and at Er–Yb absorption.

The highest internal gain reached in this study was equal to 6.4 dB/cm for $\tau=3$ obtained at a moderate pump power value of 70 mW. Larger gain could be expected by achieving single mode buried waveguides at 980 and 1530 nm to obtain a better overlap rate.

In conclusion, we have manufactured erbium-ytterbium codoped PS waveguides and evidenced their ability to be used as a light amplifier around the wavelength equal to 1.5 μm . Higher gain values can be expected by achieving single mode buried waveguides and by achieving waveguides at low temperature to obtain good quality interfaces and so reduce surface scattering losses due to the interface roughness.

TABLE II. Loss values of undoped waveguides and Er–Yb codoped waveguides as a function of the ratio (τ) of Yb concentration to Er concentration at 1.55 μm .

Type of waveguides	Codoped				
	Undoped	$\tau=1$	$\tau=2$	$\tau=3$	$\tau=5$
Losses (dB/cm)	0.7 ± 0.3	1.7 ± 0.5	1.9 ± 0.7	2.1 ± 0.8	1.9 ± 0.8

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