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► **To cite this version:**

| Yves Colin de Verdière. Semi-classical analysis and passive imaging. 2009. hal-00264373v2

HAL Id: hal-00264373

<https://hal.science/hal-00264373v2>

Preprint submitted on 17 Feb 2009 (v2), last revised 17 Apr 2009 (v3)

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Semi-classical analysis and passive imaging

Yves Colin de Verdière *

February 17, 2009

Abstract

The propagation of elastic waves inside the earth provides us information's about the geological structure of the earth's interior. Since the beginning of seismology, people are using waves created by earthquakes or by artificial explosions. They record the waves as functions of the time using seismometers located at different stations on the earth's surface. Even without any earthquake or explosion, a weak signal is still recorded which has no evident structure: it is a "noise". How to use these noises? This is the goal of the method of "passive imaging". The main observation is the following one: the time correlation of the noisy fields, computed from the fields recorded at the points A and B , is "close" to the Green's function $G(\tau, A, B)$ of the wave propagation. The aim of this paper is to provide a mathematical context for this approach and to show, in particular, how the methods of semi-classical analysis can be used in order to find the asymptotic behavior of the correlations.

Introduction

The seismologists want to recover the physical parameters of the earth's interior from records (called seismogram's), at the earth's boundary, of the elastic waves propagating inside the earth. From the mathematical point of view, it is an example of a so-called "inverse problem": the most famous inverse problems are the Calderon problem (recovering the conductance of a domain from boundary measurements) and the Kac problem (recovering an Euclidian domain from the spectrum of its Dirichlet Laplacian). The use of elastic waves created by an earthquake (resp. an artificial explosion) has been quite successful in order to recover the large scale structure (resp. the small scale structure: oil detection for example). The geological structure of the earth crust up to depths of 30-50 kilometers is much more difficult to know.

On the other hand, only a small part of the seismogram's was used: this part corresponds to the propagation of well identified body and surface waves: typically the P-waves, the S-waves and the Rayleigh waves. The last part (called the "coda") of the seismogram's were not used. This was a long standing question asked by K. Aki, one of the founders of the modern seismology. Quite recently, it was observed by M. Campillo and A. Paul [5] that the time-correlation of coda waves is closely related to the Green's function.

Later, M. Campillo and his collaborators were showing that the correlations of ambient noises already gives the Green's functions [8, 24, 25]. Similar results were already known to work in acoustics [18, 31, 32, 9] and helio-seismology [12]. Theoretical results were derived by several people in [2, 21, 23]. Applications to clock

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synchronization [22] and to volcanic eruptions forecasting [4] have been developed. The paper [16] is a recent survey paper. So, let us assume that we can recover the Green's function: how does it help in order to solve our inverse problem? Here enters semi-classics (or ray theory): the semi-classical behavior of the Green's function contains the ray dynamics. Even more, the Green's function is not really needed for that, its phase is enough. So it is natural to study the problem of passive imaging using the tools of semi-classical analysis.

The goal of my paper is to present some mathematical models for this method of passive imaging. They will be applicable to any kind of wave, but I have in mind the seismic waves. Semi-classical analysis will be used as a way to get a geometric approximation to the wave propagation AND in the modeling of the source noise. I do not at all pretend that the assumptions of my models are realistic from the physical point of view. I just hope that they will help seismologists by providing another point of view and other words. On the other hand, these models put light onto several interesting mathematical problems: mode conversions, inverse spectral problems, modeling of random fields . . .

Things work as follows: let us assume that we have a medium X (a smooth manifold) and a smooth, deterministic (no randomness in it), linear wave equation in X . We hope to recover (part of) the geometry of X from the wave propagation. We assume that there is somewhere in X a source of noise $\mathbf{f}(x, t)$ which is a stationary random field. This source generates, by the (linear) wave propagation, a field $\mathbf{u}(x, t)$ (assumed to be scalar in this introduction). This field \mathbf{u} is recorded at different points A, B, \dots on long time intervals. We want to get some information on the propagation of waves from B to A in X from the correlation

$$C_{A,B}(\tau) = \lim_{T \rightarrow +\infty} \frac{1}{T} \int_0^T \mathbf{u}(A, t) \bar{\mathbf{u}}(B, t - \tau) dt$$

which can be computed numerically from the fields recorded at A and B . In seismology the medium X will be the earth, the waves are the elastic waves, the main source noise is due to the interaction of ocean and atmosphere with the earth crust and the field is recorded at the surface of the earth.

It turns out that $C_{A,B}(\tau)$ is closely related to the deterministic *Green's function* of the wave equation in X . It means that one can hope to recover, using Fourier analysis, the propagation speeds of waves between A and B as a function of the frequency, and, more precisely, the Hamiltonian or, more generally, the so-called *dispersion relation*.

If the wave dynamics is time reversal symmetric, the correlation admits also a symmetry by change of τ into $-\tau$; this observation has been used for clock synchronization, see [22].

The goal of this paper is to give precise formulas for $C_{A,B}(\tau)$ in the high frequency limit under some assumption on the source \mathbf{f} . We will assume that the source field is a stationary ergodic random field whose covariance $K(x, y, s - s') = \mathbb{E}(f(x, s) \bar{f}(y, s'))$ is given by the Schwartz kernel of a pseudo-differential operator. This is a strong assumption which implies in particular a rapid decay of the correlations of the source \mathbf{f} outside the diagonal.

More precisely, we have two small parameters, one of them entering into the decorrelation distance of the source noise, the other one in the high frequency propagation (the *ray* method). The fact that both are of the same order of magnitude is crucial for the method. The method works well in certain frequency ranges of the wave propagation. From the point of view of physics, the main result is that, with some assumptions on the support of the source noise, one can recover the dispersion relation (the ray dynamics) in some frequency interval. A mathematical statement will be that, for $\tau > 0$, $C_{A,B}(\tau)$ is close to the Schwartz kernel of $\Omega(\tau) \circ \Pi$ where Π

is a suitable *pseudo-differential operator* (a Ψ DO), whose principal symbol can be explicitly computed, and $\Omega(\tau)$ is the (semi-)group of the (damped) wave propagation. This closeness is in general only in the L^2 sense, but can be point-wise if the damping is strong enough. It implies that we can recover the dispersion relation, i.e. the classical dynamics, from the knowledge of all two-points correlations.

The waves are recorded at the boundary, it implies that the main part of the recorded noisy fields is coming from surface waves. So that we get another inverse problem: how do the dispersion relation of the elastic surface waves determines the structure of the earth's interior? We will discuss the effective (dispersive) ray dynamics of the guided surface waves. The corresponding ideal inverse spectral problem has been solved in the paper [7] of which we recall the main result.

I guess that our results can be interesting because they provide a quite general situation in which the source is not a white noise and for which we can still recover the most interesting part of the Green's function (the phase).

Let us also mention on the technical side that, rather than using mode decompositions, we prefer to work directly with the dynamics; in other words, we need really a *time dependent* rather than a *stationary approach*. Mode decompositions are often useful, but they are of no much help for general operators with no particular symmetry.

For clarity, we will first discuss the non-physical case of a first order wave equation like the Schrödinger equation, then the case of a more usual wave equations (acoustics, elasticity).

In order to make the paper readable by a large set of people, we have tried to make it self-contained by including sections on pseudo-differential operators and on random fields.

Contents

- In Section 1, we start with a quite general setting and discuss a very general formula for the correlation (Equation (2)). This formula will be made more precise in several cases in the following Sections.
- The goal of Section 2 is to introduce the basic examples which we will discuss in the paper.
- In Section 3, we introduce a large family of anisotropic random fields and show the relation between their power spectra and their Wigner measures. They are build from white noises using the Radonification process: our random fields will be the images of some white noises by some suitable linear operators.
- Section 4 is devoted to the case of a scalar field driven by a first order in time differential equation. This is quite academic, but will be used later for standard wave equations. We will discuss first the (non-realistic) case where the source is a white noise, then the semi-classical case assuming the absence of mode conversions.
- In Section 5, we discuss the case of multi-component wave equations: we will discuss first the (non-realistic) case where the source is a white noise, then the semi-classical case using a reduction to the case of scalar fields.
- In Section 6, going back to the beginning of the story, we discuss the question of the correlations of codas in the framework of “quantum chaos”.
- In Section 7, we focus on the case of seismology and discuss the remarkable fact that the correlation of the surface waves is enough to image the crust of the earth using a wave-guide model of the crust.

- Finally, there are four appendices:
 - Appendix A on pseudo-differential operators,
 - Appendix B on classical dynamics, dispersion relation and generating functions,
 - Appendix C on Egorov Theorem for long times,
 - Appendix D on time averages versus ensemble averages, i.e. ergodicity of the source field.

As can be seen in the paper, there are several interesting mathematical problems which come from our discussions; let us mention

- Mode conversions coming from an eigenvalue going into the continuous spectrum (Section 7.3).
- Extension of the large time Egorov Theorem with Fourier integral operators instead of Ψ DO's in order to be able to discuss properly the case of the source noise located on a sub-manifold.
- The large time behaviour of the solutions of a wave equation with initially localized Cauchy data, like an earthquake (Section 6).

Acknowledgment: I would like to thank Michel Campillo and his colleagues from LGIT for discussions and collaborations. I thank also the referee for his great help in order to improve the initial manuscript.

1 A general formula for the correlation

In this quite mathematical Section, we present the general context and a general formula for the correlation which makes quite explicit the relation with the Green's function: we will do that on the level of operators and on the level of more concrete functions, the integral (or Schwartz) kernel of the operators. Our model is quite general: it includes a linear wave equation with some attenuation and a source noise which lies on the right-hand side of the wave equation.

We will consider the following general damped wave equation:

$$\frac{d\mathbf{u}}{dt} + \frac{i}{\varepsilon} \hat{H} \mathbf{u} = \mathbf{f} \quad (1)$$

1.1 Operators level

We make the following assumptions:

- The parameter $\varepsilon > 0$ can be either fixed say $\varepsilon = 1$ or can be a semi-classical parameter going to 0^+ , like \hbar in the Schrödinger like case or the inverse of the frequency for wave equations.
- The operator $i\hat{H}/\varepsilon$ is the infinitesimal generator of a strongly continuous group (see [36], Chap. IX)

$$\Omega(t) = \text{"exp}(-it\hat{H}/\varepsilon)\text{"}$$

acting on some Hilbert space \mathcal{H} .

- $\Omega(t)$ is a contraction group: there exists $T_{\text{att}} > 0$, called the *attenuation time*, and $C > 0$ so that the operator norm of $\Omega(t)$ satisfies

$$\forall t \geq 0, \|\Omega(t)\| \leq C e^{-t/T_{\text{att}}}$$

- The source noise \mathbf{f} is a stationary (see Section 3.4) random field $\mathbf{f} : \mathbb{R} \rightarrow \mathcal{H}$. We assume that $\mathbb{E}(\|\mathbf{f}(0)\|^2) < \infty$ ⁽¹⁾ and the covariance of \mathbf{f} is given by⁽²⁾

$$\mathbb{E}(|\mathbf{f}(s) \rangle \langle \mathbf{f}(s')|) = K(s - s')$$

with $K(t) \in \mathcal{L}(\mathcal{H})$ ⁽³⁾, a continuous function of t with compact support $\{|t| \leq t_0\}$ ⁽⁴⁾.

The field \mathbf{u} is then stationary and given by

$$\mathbf{u}(t) = \int_0^\infty \Omega(s)\mathbf{f}(t - s)ds .$$

It satisfies $\mathbb{E}(\|\mathbf{u}(0)\|^2) < \infty$ and its covariance, defined by

$$C(\tau) = \mathbb{E}(|\mathbf{u}(0) \rangle \langle \mathbf{u}(-\tau)|) ,$$

is a continuous linear operator on \mathcal{H} .

Theorem 1.1 *The covariance $C(\tau)$ is given by*

- for $\tau > t_0$,

$$C(\tau) = \Omega(\tau)\Pi \text{ with } \Pi = \int_0^\infty \Omega(s)\mathcal{L}\Omega^*(s)ds \text{ and } \mathcal{L} = \int_{-t_0}^{t_0} \Omega(-u)K(u)du \quad (2)$$

- for $\tau < 0$,

$$C(\tau) = (C(-\tau))^* .$$

Assuming \mathbf{f} to be ergodic (see Section 3.4), the covariance $C(\tau)$ is also given almost surely by the *time average*:

$$C(\tau) = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T |\mathbf{u}(t) \rangle \langle \mathbf{u}(t - \tau)| dt .$$

In order to perform effective passive imaging, people use the previous formula for a large enough T and ... computers.

1.2 Kernels level

In order to have concrete functions and not only distributions, we will make some smoothness assumptions:

- X is a *smooth compact manifold* of dimension d with a smooth measure $|dx|$. For applications to seismology X will be the earth or (in Section 7) the surface of the earth.
- $\mathbf{u}(x, t)$, $x \in X$, $t \in \mathbb{R}$ is *the field* (scalar or vector valued) with values in \mathbb{C}^N (or \mathbb{R}^N).

¹ \mathbb{E} denotes the ensemble average or expectation value

²We use Dirac's bra-ket notation, $|f \rangle$ is a vector, $\langle f|$ is the Hermitian scalar product with f , so that $|f \rangle \langle g|$ is the rank one linear operator: $x \rightarrow \langle g|x \rangle f$

³ $\mathcal{L}(E, F)$ is the space of continuous linear maps from the Banach space E into the Banach space F with the norm topology and $\mathcal{L}(E) := \mathcal{L}(E, E)$

⁴The assumption on compact support could be replaced by some suitable decay in t . It means that the source noise is uncorrelated for large enough times

- The generator of the group $\Omega(t)$ is an elliptic differential (or pseudo-differential) linear operator \hat{H} acting with a dense domain of an Hilbert space $\mathcal{H} \subset \mathcal{D}'(X, \mathbb{C}^N)$.
- The *source* \mathbf{f} is a smooth *stationary random field* on $X \times \mathbb{R}$ with values in \mathbb{C}^N (or \mathbb{R}^N) whose matrix valued correlation kernel is the smooth function on $X \times X \times \mathbb{R}$ with values in $\text{End}(\mathbb{C}^N)$ given by

$$\mathbb{E}(\mathbf{f}(x, s) \otimes \mathbf{f}^*(y, s')) = K(s - s', x, y) . \quad (3)$$

Definition 1.1 *The propagator or Green's function P is the Schwartz kernel of $\Omega(t)$ defined by the formula:*

$$(\Omega(t)\mathbf{v})(x) = \int_X P(t, x, y)\mathbf{v}(y)|dy| .$$

The propagator P satisfies

$$\int_X P(t, x, y)P(s, y, z)|dy| = P(t + s, x, z)$$

which comes from: $\Omega(t + s) = \Omega(t) \circ \Omega(s)$. The causal solution of Equation (1) is then given by

$$\mathbf{u}(x, t) = \int_0^\infty ds \int_X P(s, x, y)\mathbf{f}(t - s, y)|dy| .$$

It is also a smooth function (ellipticity of \hat{H}).

The kernel of $C(\tau)$ is defined by:

$$C_{A,B}(\tau) = \mathbb{E}(\mathbf{u}(A, 0) \otimes \mathbf{u}(B, -\tau)^*) , \quad (4)$$

where $A, B \in X$ and $C_{A,B}(\tau) \in L(\mathbb{C}^N)$.

From Equation (2), we get an explicit integral formula for $C_{A,B}(\tau)$ in terms of $P(t, x, y)$ and $K(t, x, y)$; Equation (2) can be written for $\tau > t_0$ as follows⁽⁵⁾:

$$C_{A,B}(\tau) = [\Omega(\tau)\Pi](A, B) \quad (5)$$

2 Examples

In this Section, we will present the examples of “damped wave equations” which will be used in this paper. We start with concrete examples and go then to more general semi-classical examples in the scalar case as well as in the multi-component case.

2.1 Damped Schrödinger equation

Let X be a smooth compact boundary-less Riemannian manifold whose Laplace-Beltrami operator is denoted by Δ . Let us give $a : X \rightarrow \mathbb{R}$ a smooth non negative function, V a smooth real valued function on X , \hbar a non negative *constant*, and $\hat{K} = -\hbar^2\Delta + V(x)$, and take:

$$\frac{\hbar}{i}(u_t + a(x)u) + \hat{K}u = g .$$

It is a particular case of Equation (1) where $\varepsilon = \hbar$, $\hat{H} = \hat{K} + \frac{\hbar}{i}a(x)$, $f = \frac{i}{\hbar}g$ and the Hilbert space is $L^2(X)$.

⁵If R is an operator we will denote by $[R](x, y)$ its Schwartz kernel

Let us note for future use that, if $\hbar \rightarrow 0$, the principal symbol of our equation is $\omega + \|\xi\|^2 + V(x)$ where our (extended) phase space is $T^*(X \times \mathbb{R})$ with canonical coordinates (x, ξ, t, ω) and $a(x)$ is only a sub-principal term entering into the transport equation, but not in the classical dynamics (see Appendix A). The attenuation time can be taken as $T_{\text{att}} = 1/\inf_{x \in X} a(x)$. Let us note that this equation has no physical meaning, it is just a kind of toy model.

2.2 Damped wave equations

As in the previous Section, X is a compact Riemannian manifold. Let us start with

$$u_{tt} + au_t - \Delta u = f, \quad (6)$$

(a is a smooth function which satisfies $\inf_{x \in X} a(x) > 0$) which corresponds to Equation (1) with

$$\mathbf{u} = \begin{pmatrix} \varepsilon \sqrt{-\Delta} u \\ -i\varepsilon u_t \end{pmatrix}, \quad \mathbf{f} := \begin{pmatrix} 0 \\ -\varepsilon^2 f \end{pmatrix}.$$

and

$$\hat{H} = \begin{pmatrix} 0 & -\varepsilon \sqrt{-\Delta} \\ -\varepsilon \sqrt{-\Delta} & -i\varepsilon a \end{pmatrix}.$$

The Hilbert space is defined by the energy norm $\|(u, v)\|^2 = \int_X (|u|^2 + |v|^2) |dx|$.

Remark 2.1 *It follows from [19], that, if $a(x) \geq 0$ everywhere and there exists some $T_0 > 0$ so that the support of a crosses every geodesic arc of length $\geq T_0$, then $\Omega(t)$ is a contraction group.*

One can replace the very simple wave equation (6) by more sophisticated multi-component wave equations like the equation of elastic waves. The general form will be:

$$u_{tt} + Au_t - \Lambda u = 0 \quad (7)$$

where $u(x, t) \in \mathbb{C}^n$, Λ is an elliptic system of degree 2 and A is a differential operator of degree 0 representing the attenuation.

2.3 Pseudo-differential equations

This example will be the subject of Section 4.

This is a generalization of the damped Schrödinger equation (Section 2.1); we now introduce a small parameter ε which in the case of the damped Schrödinger equation is the Planck constant \hbar . The semi-classical regime will be $\varepsilon \rightarrow 0$,

We assume that the dynamics is generated by a pseudo-differential operator (a Ψ DO, see Appendix A). Our equation looks then like:

$$\mathbf{u}_t + \frac{i}{\varepsilon} \hat{H}_\varepsilon \mathbf{u} = \mathbf{f}$$

with

$$\hat{H}_\varepsilon = \text{Op}_\varepsilon(H_0 + \varepsilon H_1).$$

Here H_0 is real valued and elliptic, and is the Hamiltonian of the ray dynamics, while $\Im H_1 < 0$ provides the attenuation.

In the case of damped Schrödinger equation (Section 2.1), we have $\varepsilon = \hbar$, $H_0 = \|\xi\|^2 + V(x)$ and $H_1 = -ia(x)$.

2.4 Pseudo-differential systems

This example will be the subject of Section 5.

The system is of the following form

$$\mathbf{u}_t + \frac{i}{\varepsilon} \hat{H}_\varepsilon \mathbf{u} = \mathbf{f} ,$$

with

$$\hat{H}_\varepsilon = \text{Op}_\varepsilon(H_0 + \varepsilon H_1) .$$

Here

- $\mathbf{u}(x, t) \in \mathbb{C}^N$
- $H_0(x, \xi)$ is an Hermitian symmetric matrix. Each eigenvalue $\lambda(x, \xi)$ of $H_0(x, \xi)$ gives a ray dynamics associated to a polarized wave. The polarization bundle is the associated eigenbundle $E_\lambda(x, \xi)$
- $H_1(x, \xi)$ is a complex matrix. Let us introduce the quadratic forms $q_\lambda(x, \xi)$ which is the restriction of $\langle \Im H_1 x | x \rangle$ to $E_\lambda(x, \xi)$. The condition $q_\lambda < 0$ provides the attenuation.

This includes

- The high frequency limit of the wave equations (6): we have

$$H_0 = \begin{pmatrix} 0 & -\|\xi\| \\ -\|\xi\| & 0 \end{pmatrix}, \quad H_1 = \begin{pmatrix} 0 & 0 \\ 0 & -ia \end{pmatrix} .$$

A similar transform can be used for more general wave equations, like the elastic wave equation with $N = 6$.

- An effective pseudo-differential wave equation associated to stratified media. They are usually Ψ DO's with a non trivial dispersion relation (see Section 7).
- Phase space dependent damping included in H_1 : this is usually the case for seismic waves: the attenuation can be created by the small scales physics which is not uniform in general.

3 Random fields and semi-classics

Our aim in this section is to build quite general random fields with correlation distances given by a small parameter ε . It seems to be natural for that purpose to use ε -pseudo-differential operators. We will see how to compute the generalized power spectrum using Wigner measures. We will not restrict our-self to time-dependent random fields. The subsection 3.4 is devoted to the time-dependent case and we introduce there the notions of stationary and ergodicity.

3.1 Covariance of Random fields

Let \mathcal{H} be an Hilbert space. We will consider random fields $f = f(\bar{\omega}) \in \mathcal{H}$ where $\bar{\omega}$ belongs to Ω , the sample probability space. We can even take $\Omega = \mathcal{H}$ (and $f = \text{identity}$) and the random field is then a probability on \mathcal{H} . We will assume that $f \in L^2(\Omega, \mathcal{H})$ and, without loss of generality, that our fields have mean 0: $\mathbb{E}(f) = 0$.

The *covariance* of f is the operator $C = \mathbb{E}(|f\rangle\langle f|)$. This operator is *positive and Hermitian*. Often, \mathcal{H} will be a space of functions and we will consider the Schwartz kernel $[C](x, y)$ of C .

If \mathcal{H} is finite dimensional, there are random fields on it whose correlation is the Identity. We call them the *white noises*.

3.2 White noises in Hilbert spaces

For this section, see [27] Part II and [15] Chap IV

A white noise w on the Hilbert space \mathcal{H} is a noise of mean 0 and correlation the Identity. From a more formal mathematical point of view, w can be seen as a *cylindrical probability* on \mathcal{H} : the images of w by continuous linear mapping of finite rank are usual probabilities. In particular, bounded continuous functions of a *finite number* of linear forms are integrable. The fact that the covariance of w is the Identity reads as $\forall e, f \in \mathcal{H}, \mathbb{E}(\langle e|w\rangle\langle w|f\rangle) = \langle e|f\rangle$.

Definition 3.1 *The Fourier transform of a cylindrical probability w is defined as usual by $\hat{w}(\xi) = w(\exp(-i\langle \xi|x\rangle))$.*

Definition 3.2 *The Gaussian white noise is the white noise on \mathcal{H} whose Fourier transform is $\exp(-\|\xi\|^2)$.*

If the cylindrical probability comes from an usual probability, this probability is unique and we will say that our cylindrical probability is a probability.

We have the important

Theorem 3.1 (Sazonov [27]) *If the Fourier transform \hat{w} of the cylindrical white noise w is continuous and $A : \mathcal{H} \rightarrow \mathcal{K}$ is Hilbert-Schmidt, then $A(w)$ is a probability measure on \mathcal{K} . The covariance of $A(w)$ is the operator AA^* .*

The previous Theorem applies in particular to the *Gaussian white noise*.

If the Hilbert space \mathcal{H} is not finite dimensional, the white noises are “random distributions” in the following sense: if $\mathcal{H} = L^2(X)$, we can take for \mathcal{H} a Sobolev space H^s of negative order with $s < -\dim X/2$ and for A the injection from $L^2(X)$ into $H^s(X)$. So that w can be viewed as a random distribution in $H^s(X)$.

3.3 An important example

Let us give a self-adjoint operator \hat{H} on \mathcal{H} with compact resolvent. Let us denote by A_I the (finite rank) spectral projection onto the compact interval I . The random field $f = A_I w$ can also be defined in terms of the spectral decomposition (f_j, λ_j) of \hat{H} . $f = \sum_{\lambda_j \in I} a_j(\bar{\omega}) f_j$ where the a_j satisfies $\mathbb{E}(a_i) = 0$ and $\mathbb{E}(a_i \bar{a}_j) = \delta_{i,j}$.

3.4 Stationary Random fields

We will now consider *time dependent* random fields $f : t \rightarrow f(t) \in \mathcal{K}$ with $t \in \mathbb{R}$. Our Hilbert space will be $\mathcal{H} = L^2(\mathbb{R}, \mathcal{K})$. For $T \in \mathbb{R}$, we define $T.f$ by $(T.f)(t) = f(t-T)$. We will say that

Definition 3.3 *f is stationary if there exists a group of probability preserving transformations of Ω , denoted by X_T so that $T.(f(\bar{\omega})) = f(X_T(\bar{\omega}))$.*

Let us consider the time dependent fields $f : \mathbb{R} \rightarrow \mathcal{K}$ as elements of a topological vector space \mathcal{B} . A time dependent random field is a random element of \mathcal{B} . This can be viewed as a probability dP on $\Omega = \mathcal{B}$. The stationarity is then clearly equivalent to the fact that dP is invariant by the translations $X_T : f \rightarrow T.f$ of functions in \mathcal{B} .

Stationarity corresponds to the physical intuition that $T.f$ is another random choice of the function f (see also Wiener book [35]).

If f is a stationary random field, the co-variance kernel $\tilde{K}(t, t') = \mathbb{E}(|f(t)\rangle\langle f(t')|)$ is a function $K(t - t') \in \text{Herm}(\mathcal{K})$.

Definition 3.4 *A stationary random field f is ergodic if any measurable subset of Ω , invariant by the transformations X_T , is of probability 0 or 1.*

We have then the celebrated

Theorem 3.2 (Birkhoff) *If f is stationary and ergodic and $F : \Omega \rightarrow \mathbb{R}$ is integrable with respect to the probability measure, then, for almost every $\bar{\omega} \in \Omega$,*

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T F(X_t(\bar{\omega})) dt = \mathbb{E}(F) .$$

3.5 Examples of stationary fields

We will start with a Gaussian white noise w on $L^2(\mathbb{R}, \mathcal{K})$. Let us consider an Hilbert-Schmidt operator $A : L^2(\mathbb{R}, \mathcal{K}) \rightarrow \mathcal{K}$. Let us introduce the stationary random field f defined by

$$f(t) = A(t.w) .$$

If the Schwartz kernel of A is $a(s, x, y)$ we have, formally, $f(t, x) = \int_X |dy| \int_{\mathbb{R}} ds a(s+t, x, y) w(s, y)$.

Let us assume that $Aw = \int_{\mathbb{R}} \tilde{A}(s)w(s)ds$ with $\tilde{A}(s)$ a smoothly dependent of s Hilbert-Schmidt operator with compact support $\subset [-t_0/2, t_0/2]$. From the formula in Theorem 3.1, the correlation operator is

$$K(t) = \int_{\mathbb{R}} \tilde{A}(t+s)\tilde{A}(s)ds \quad (8)$$

which is supported in $[-t_0, t_0]$.

Example 3.1 Stationary scalar noise on the real line: *let us take $\mathcal{K} = \mathbb{R}$, $F \in C_0^\infty(\mathbb{R}, \mathbb{R})$ ⁽⁶⁾ and the operator A given by $w \rightarrow \int_{\mathbb{R}} F(-s)w(s)ds$. Then f is the convolution $F \star w$. On the level of Fourier transforms $\hat{f} = \hat{F}\hat{w}$, $\mathbb{E}(|\hat{f}(\omega)|^2) = |\hat{F}|^2(\omega)$ and the positive function $|\hat{F}|^2(\omega)$ is usually called the power spectrum of the stationary noise.*

Example 3.2 Brownian motions: *if $X = \mathbb{R}$ and w is the Gaussian white noise, the Brownian motion is defined by $b(t) = \int_0^t w(s)ds$, which is in $L^2([0, T])$ for all finite T . b is NOT a stationary field!*

3.6 Random fields from pseudo-differential operators

The main goal of the present section is to build natural random fields which are non homogeneous with small distances of correlation of the order of $\varepsilon \rightarrow 0$. The noise is non homogeneous in X , but could also be non isotropic w.r. to directions.

We will use as Hilbert-Schmidt operators some ε -pseudo-differential operators.

3.6.1 Random fields from pseudo-differential operators

Let us take for ε -dependent random fields on a manifold Z (which can be $X \times \mathbb{R}$) the image of a white noise on $L^2(Z)$ by an ε -pseudo-differential operator A of smooth compactly supported symbol $a(z, \zeta) \in S_{\text{class}}^{-\infty, 0}$ (A is Hilbert-Schmidt!).

The correlation $C(z, z')$ will then be given as the Schwartz kernel of AA^* which is a ΨDO of principal symbol $|a|^2$. We have

$$C(z, z') = (2\pi\varepsilon)^{-d} k \left(z, \frac{z' - z}{\varepsilon} \right) + O(\varepsilon^{-(d-1)}),$$

where $k(z, \cdot)$ is the Fourier transform with respect to ζ of $|a|^2(z, \zeta)$.

⁶If X is a smooth manifold $C_0^\infty(X, \mathbb{R})$ denotes the space of smooth functions with compact support

This construction gives smooth random fields which can be localized in some very small domains of the manifold Z , which are non isotropic (the covariance kernel $C(z, z')$ for z' close to z , depends on the direction and not only on the distance $|z - z'|$) and which have small distance of correlations of the order of ε . Moreover it will allow to use techniques of micro-local analysis with the small parameter given by ε .

3.6.2 Power spectrum and Wigner measures

Wigner measure represent the phase space energy densities of a family of functions depending on a small parameter ε . There is a scaling ε in the:

Definition 3.5 *If $f = (f_\varepsilon)$ is a family of functions on Z , bounded in $L^2_{\text{loc}}(Z)$, the Wigner measures W_f^ε of f are the measures on the phase space T^*Z defined, for $\varphi \in C_o^\infty(T^*Z)$, by*

$$\int \varphi dW_f^\varepsilon := \langle \text{Op}_\varepsilon^+(\varphi) f_\varepsilon | f_\varepsilon \rangle .$$

The measures dW_f^ε are the phase space densities of energy of the functions (f_ε) .

We now define:

Definition 3.6 *The power spectrum of the random field $f = (f_\varepsilon)$ is the phase space measure $\mathcal{P}_f^\varepsilon$ defined by:*

$$\mathcal{P}_f^\varepsilon = \mathbb{E}(W_f^\varepsilon) \quad :$$

the power spectrum of a random field is the average of its Wigner measures.

Proposition 3.1 *The power spectrum \mathcal{P} of $f_\varepsilon = \text{Op}(a)w$, satisfies:*

$$\mathcal{P}_f^\varepsilon = (2\pi\varepsilon)^{-d} |a|^2(x, \xi) |dx d\xi| + O(\varepsilon^{-(d-1)}) .$$

Proof.-

Let us put $A = \text{Op}_\varepsilon(a)$, we have

$$\langle \text{Op}(\varphi) Aw | Aw \rangle = \langle A^* \text{Op}^+(\varphi) Aw | w \rangle$$

and $\forall B, \mathbb{E}(\langle Bw | w \rangle) = \text{trace}(B)$. We get

$$\mathbb{E} \left(\int \varphi dW_f^\varepsilon \right) = \text{trace}(A^* \text{Op}^+(\varphi) A)$$

which can be evaluated using the Ψ DO calculus as

$$\mathbb{E} \left(\int \varphi dW_f^\varepsilon \right) = (2\pi\varepsilon)^{-d} \int \varphi |a|^2 dx d\xi + O(\varepsilon^{-(d-1)}) .$$

□

3.6.3 Time dependent random fields

If $Z = X \times \mathbb{R}$ is the space-time, we will denote by (x, ξ) (resp. (t, ω)) the canonical local coordinates in T^*X (resp. $T^*\mathbb{R}$). We will take our random fields as before: $f = L_\varepsilon w$; the symbol l of $L_\varepsilon = \text{Op}_\varepsilon(l)$ is assumed to be $l(x, \xi, \omega) \in S_{\text{class}}^{-\infty}(X \times \mathbb{R})$ (independent of t). The field f is stationary.

In this case, the correlation is given by:

$$K(x, y; t) = [LL^*](x, y; 0, t) \tag{9}$$

which is the Schwartz kernel of a Ψ DO of principal symbol $|l|^2(x, \xi; \omega)$.

Lemma 3.1 *The covariance operator $K(t)$ is a Ψ DO in the space $\Psi^{-\infty}$ of principal symbol*

$$\frac{1}{2\pi\varepsilon} \int_{\mathbb{R}} |l|^2(x, \xi; \omega) e^{it\omega/\varepsilon} d\omega . \tag{10}$$

4 Correlation for 1-component wave equation

In this section, we will study the case of a scalar field given as in Section 2.3. First we describe, as a warming up, a case where the correlation is exactly computable, then we compute the asymptotic of the correlation in the "semi-classical case".

4.1 An example of an exact formula for the correlation

Let us consider the 1-component "wave equation" (with $\varepsilon = 1$)

$$\frac{du}{dt} + ku + i\hat{H}_0 u = f$$

where $k > 0$ is a constant and \hat{H}_0 is self-adjoint on \mathcal{K} with a compact resolvent. In particular $\Omega(t) = e^{-kt}U(t)$ where $U(t)$ is the unitary group $U(t) = \exp(-it\hat{H}_0)$. Let us assume that $\mathbf{f} = \Phi \star A_I w$, i.e. $f(t) = A_I (\int \Phi(t-s)w(s)ds)$, with w a Gaussian white noise on $L^2(\mathbb{R}, \mathcal{K})$, $\Phi \in C_0^\infty(]-t_0/2, t_0/2[; \mathbb{R})$ and $A_I = \chi_I(\hat{H}_0)$ is the finite rank spectral projection of \hat{H}_0 associated to the compact interval I . The operator $A : w \rightarrow \int_{\mathbb{R}} \Phi(-s)A_I w(s)ds$ is clearly Hilbert-Schmidt. From Equation (8), the correlation is given by $K(t) = \Psi(t)A_I$ with $\Psi(t) = \int_{\mathbb{R}} \Phi(t+s)\Phi(s)ds$. Using Theorem 1.1, we get

Theorem 4.1 *We have the following simple formula for $\tau > t_0$:*

$$C(\tau) = \frac{1}{2k} F(\hat{H}_0) \Omega(\tau) ,$$

with $F(\omega) = \chi_I(\omega) \mathcal{F}_{t \rightarrow \omega}(\Psi(t)e^{kt})(\tau)$, which is an exact relation between the correlation and the propagator P_0 of the wave equation without attenuation.

Remark 4.1 *At least formally, we can take $\Phi = \delta(0)$, $I = \mathbb{R}$, then $f = w$ and $F \equiv 1$. The correlation is then equal to the Green's function up to the factor $1/2k$.*

Remark 4.2 *Assuming that \hat{H} is "real valued" (i.e. $\overline{\hat{H}f} = \hat{H}\bar{f}$), we have $C_{A,B}(-\tau) = \overline{C_{A,B}(\tau)}$.*

4.2 Semi-classical assumptions

Let us start with the semi-classical Schrödinger like equation

$$u_t + \frac{i}{\varepsilon} \hat{H} u = f$$

where

- \hat{H} is **admissible** (Definition 10.1 in Appendix C): \hat{H} is an ε -pseudo-differential operator:

$$\hat{H} := \text{Op}(H_0 + \varepsilon H_1)$$

with

- The **principal symbol** $H_0(x, \xi) : T^*X \rightarrow \mathbb{R}$, which gives the classical ("rays") dynamics, belongs to Σ_m , with $m > 0$, and is elliptic of degree m .

⁷ \mathcal{F} is the Fourier transform $\mathcal{F}(g)(u) = \int \exp(-ius)g(s)ds$.

- The **sub-principal symbol** $H_1(x, \xi)$ belongs to Σ_{m-1} and admits some positivity property which controls the attenuation: there exists $k > 0$, such that

$$h_1 := \Im H_1 \leq -k .$$

- The **random field** f is given by $f = \text{Op}(l(x, \xi, \omega))w$ with w a white noise on $X \times \mathbb{R}$ and with l smooth, compactly supported w.r. to (x, ξ) and whose Fourier transform w.r. to ω is compactly supported in $] -T_0, T_0[$. The power spectrum of f is $(2\pi\varepsilon)^{-(d+1)}|l|^2(x, \xi, \omega)$. The covariance operator $K(t)$ is a ΨDO given by Lemma 3.1 and supported in the interval $] -T_0\varepsilon, T_0\varepsilon[$.
- We introduce some compact energy interval I so that $\text{Support}(l) \subset H_0^{-1}(I)$.

Lemma 4.1 *Under the previous assumption on H_1 , we can take for T_{att} any number $> 1/k$ provided that ε is small enough.*

Proof.–

$$\frac{d}{dt} \langle v(t) | v(t) \rangle = 2\Re \langle v(t) | -i\hat{H}_1(t)v(t) \rangle$$

and we use Gårding inequality (see [10]): if $a \geq 0$, $\text{Op}_\varepsilon(a) \geq -C$ for any $C > 0$ and ε small enough.

□

4.3 Ehrenfest time

Definition 4.1 (Ehrenfest time) *Let X be a smooth compact manifold of dimension d . Let us consider a smooth proper Hamiltonian H_0 on T^*X . We fix also some compact energy interval I .*

- The Liapounov exponent $\Lambda_I = \Lambda \in]0, +\infty[$ is **any** real number⁽⁸⁾ for which the differential of the Hamiltonian flow Φ_t of H_0 satisfies the following uniform estimate:

$$\exists C > 0, \forall z \in T^*X \text{ with } H_0(z) \in I, \forall t \geq 0, \|d\Phi_t(z)\| \leq Ce^{\Lambda t} .$$

- Let us denote by $\varepsilon > 0$ the semi-classical parameter. The Ehrenfest time, $T_{\text{Ehrenfest}} \in]0, +\infty[$ is defined as

$$T_{\text{Ehrenfest}} := |\log \varepsilon| / \Lambda .$$

For any time smaller than $T_{\text{Ehrenfest}}$, any cell of diameter ε is not expanded to the whole phase space. We will need a smaller time:

Definition 4.2 *Let us give $0 < \gamma < \frac{1}{2}$. The time T_γ is defined as*

$$T_\gamma := \left(\frac{1}{2} - \gamma \right) T_{\text{Ehrenfest}} .$$

The previous definition will be used via the

Lemma 4.2 *Assume that we have a flow Φ_t on a compact manifold X with a Liapounov exponent Λ , then we have the following uniform estimate*

$$\forall \alpha, \forall \Lambda' > \Lambda, |D^\alpha \Phi_t(z)| = O(e^{\Lambda' |t| |\alpha|}) .$$

⁸For technical reasons, not taking the infimum of these Λ 's is more convenient

4.4 Main result ($N = 1$)

We get the main result:

Theorem 4.2 *With the assumptions and notations of Section 4.2 and $\gamma < \frac{1}{2}$, the correlation is given, for $\tau > 0$, by $C(\tau) = \Omega(\tau) \circ \text{Op}_\varepsilon(\pi) + R$ where*

•

$$\pi(x, \xi) = \int_0^{T_\gamma(I)} \exp\left(2 \int_{-t}^0 \Im(H_1)(\Phi_s(x, \xi)) ds\right) |l|^2(\Phi_{-t}(x, \xi), -H_0(x, \xi)) dt$$

belongs to $\Sigma_{\frac{1}{2}-\gamma}^{-\infty}$, and is supported in $H_0^{-1}(I)$.

• 1. [No assumption on the attenuation k] we take

$$\gamma = \gamma_1 = \frac{k}{2(\Lambda + k)},$$

and, for any $\alpha < 2\gamma_1$, we have the Hilbert-Schmidt estimate

$$\|R\|_{\text{H-S}} = O(\varepsilon^{\alpha-d/2})$$

(the operator $\Omega(\tau) \circ \text{Op}_\varepsilon(\pi)$ is Hilbert-Schmidt⁹ with an Hilbert-Schmidt norm of the order of $\varepsilon^{-d/2}$ if π is not $\equiv 0$);

2. [Large attenuation] if $k > 2d\Lambda$, we take

$$\gamma = \gamma_2 = \frac{k - 2d\Lambda}{2(\Lambda + k)},$$

and, for any $\beta < 2\gamma_2$, for (A, B, τ) generic (see Definition 9.1), we have the point-wise estimate

$$[R](A, B) = O(\varepsilon^{\beta-\frac{d}{2}}).$$

Proof.–

From Lemma 8.1, we know that

$$\mathcal{L} \in \Psi_{\text{class}}^{-\infty}$$

with the principal symbol $|l|^2(x, \xi, -H_0(x, \xi))$. The correlation operator $C(\tau)$ is given (Theorem 1.1) by $C(\tau) = \Omega(\tau) \circ \int_0^\infty \Omega(u) \mathcal{L} \Omega^*(u) du$. We split the integral $\int_0^\infty = \int_0^{T_\gamma} + \int_{T_\gamma}^\infty$.

- The first integral is estimated using the long time Egorov Theorem.
- The second one is estimated using the exponential decay of $\Omega(t)$

The main term comes from Theorem 10.2 with the formula for a_0 given in Theorem 10.1. We split the remainder R into $R = R_1 + R_2$ with

•

$$R_1 = \Omega(\tau) \int_0^{T_\gamma} r_1(s) ds$$

and $r_1(s) \in \varepsilon^{2\gamma} \Psi_{\frac{1}{2}-\gamma}^{-\infty}$ (using Theorem 10.2)

⁹An Hilbert-Schmidt operator A is an operator whose Schwartz kernel $[A](x, y)$ is in $L^2(X \times X)$ and the Hilbert-Schmidt norm $\|A\|_{\text{H-S}}$ of A is the L^2 norm of $[A]$.

•

$$R_2 = \int_{T_\gamma}^{\infty} \Omega(s + \tau) \mathcal{L} \Omega^*(s) ds$$

with $\mathcal{L} \in \Psi_{\text{class}}^{-\infty}$.

1. *The Hilbert-Schmidt estimate:* from Lemma 8.2, we get

$$\|R_1\|_{\text{H-S}} = O(\varepsilon^{-d/2} \varepsilon^{2\gamma} |\log \varepsilon|) .$$

Using the fact that the Hilbert-Schmidt norm of \mathcal{L} is of order $\varepsilon^{-d/2}$, we have

$$\|R_2\|_{\text{H-S}} = O\left(\varepsilon^{-d/2} \int_{T_\gamma}^{\infty} e^{-2ks} ds\right) .$$

Using the explicit value $\gamma = \gamma_1$, we get the result.

2. *The point-wise estimate:*

Lemma 4.3 *The Sobolev H^s norm of a distribution u in X is, uniformly in ε , equivalent to*

$$\varepsilon^{-s} \|\hat{H}^{s/m} u\|_{L^2} + \|u\|_{L^2} .$$

Taking $s > d$, using Sobolev's embedding Theorem on $X \times X$ and the fact that \hat{H} commutes with $\Omega(s)$, we get the point-wise estimates

$$[R_2](A, B) = O\left(\varepsilon^{-2s} \int_{T_\gamma}^{\infty} \|\Omega(s + \tau) \hat{H}^{s/m} \mathcal{L} \hat{H}^{s/m} \Omega^*(s) ds\|_{\text{H-S}} ds\right)$$

and

$$[R_2](A, B) = O\left(\varepsilon^{-2s} \varepsilon^{(1-2\gamma)k/\Lambda}\right) .$$

The result follows from the optimization $\gamma = \gamma_2$.

□

4.5 Applications: recovering the Hamiltonian from the Correlations

Theorem 4.3 *Let U be the set of $(B, \xi_B) \in T^*X$ so that there exists $t > 0$ with $l(\Phi_{-t}(B, \xi_B), H_0(B, \xi_B)) \neq 0$. For any $\tau_0 > 0$, the restriction of the Hamiltonian H_0 to U can be recovered from $C(\tau)$ for $0 < \tau < \tau_0$.*

Remark 4.3 *The physical meaning of the previous assumption is that the ray dynamics and its Hamiltonian generator H_0 can be recovered in the domain where the rays cross, in the past time, the support of the power spectrum of the source noise.*

Proof.–

Let us assume, for simplicity, that $\partial^2 H_0 / \partial \xi^2$ is non degenerate. For $\tau > 0$ small enough, there is no conjugate points on the unique trajectory from B to A in time τ contained in $H^{-1}(I)$. From Theorem 4.2 and Appendix B, we get the following expression:

$$C_{A,B}(\tau) = (2\pi i \varepsilon)^{-d/2} c(A, B, \tau) e^{iS(A,B,\tau)/\varepsilon} + R(A, B, \tau)$$

with

- $c(A, B, \tau) > 0$ if $(B, \xi_B) \in U$
- The $L^2(X \times X)$ norm of $R(\cdot, \cdot, \tau)$ is $O(\varepsilon^{\alpha-d/2})$ with $\alpha > 0$.

The Theorem follows: S is a generating function for Φ_τ , the flow Φ_τ determines H_0 up to a constant. The constant can be extracted from S which is given as the action integral $\int_\gamma \xi dx - H_0 dt$ where γ is the trajectory from B to A in time τ .

Without the simplifying assumption, the Theorem still holds: the H-S norm of the “Fourier integral operator” $\Omega(\tau) \circ \text{Op}_\varepsilon(\pi)$ is of order $\varepsilon^{-d/2}$.

□

5 The case of multi-component wave equations

5.1 Wave equations: source white noise

Let us take the case of a scalar wave equation with constant damping; the closest results were derived in [33, 21]. We will make a formal calculation in the spirit of Remark 4.1. We will consider the wave equation (6), $u_{tt} + 2au_t - \Delta u = f$, with

- $a > 0$ is a constant damping coefficient
- Δ a Riemannian laplacian in some Riemannian manifold X , possibly with boundary:

$$\Delta = g^{ij}(x)\partial_{ij} + b_i(x)\partial_i$$

which is self-adjoint with respect to $|dx|$ and appropriate boundary conditions; in fact we could replace the Laplacian by any self-adjoint operator on X !

- $f = f(x, t)$ the source of the noise which will be assumed to be a scalar white noise (homogeneous diffuse field):

$$\mathbb{E}(f(x, s)f(y, s')) = \delta(s - s')\delta(x - y)$$

Let us compute the “causal solution”, i.e. the solution given by $u = \mathbf{G}f$ with \mathbf{G} linear and satisfying $u(\cdot, t) = 0$ if $f(\cdot, s)$ vanishes for $s \leq t$. In order to get a readable expression, it is convenient to introduce $Q = \sqrt{-(\Delta + a^2)}$. We get then easily:

$$u(x, t) = \int_0^\infty ds \int_X e^{-as} \left[\frac{\sin sQ}{Q} \right] (x, y) f(y, t-s) |dy|$$

where $\sin sQ$ is defined from the spectral decomposition of Q (any choice of the square root gives the same result for $\frac{\sin sQ}{Q}$). We define

$$G_a(t, x, y) = Y(t) \left[e^{-at} \frac{\sin tQ}{Q} \right] (x, y)$$

with Y the Heaviside function. We will call G_a the (*causal*) *Green's function*. We can rewrite:

$$u(x, t) = \int_{\mathbb{R}} ds \int_X G_a(t-s, x, y) f(y, s) |dy| .$$

Let us assume now that f is an homogeneous white noise and compute formally the correlation $C_{A,B}(\tau)$.

We get quite easily, using

$$\sin \alpha \sin \beta = \frac{1}{2}(\cos(\alpha - \beta) - \cos(\alpha + \beta)) :$$

$$C_{A,B}(\tau) = \frac{e^{-a|\tau|}}{4a} \left[(Q^2 + a^2)^{-1} \left(\cos \tau Q + \frac{a \sin |\tau| Q}{Q} \right) \right] (A, B) . \quad (11)$$

Taking the τ derivative, we get the simpler formula:

$$\frac{d}{d\tau} C_{A,B}(\tau) = \begin{cases} \frac{-1}{4a} G_a(\tau, A, B) & \text{for } \tau > 0 \\ \frac{1}{4a} G_a(-\tau, A, B) & \text{for } \tau < 0 \end{cases} \quad (12)$$

This “exact” relation between the correlation and the Green’s function in the case of a white noise source field is the starting point of many works in seismology. It has been derived in various ways by several authors in particular R. Weaver and O. Lobkis [18, 31, 33, 21].

Remark 5.1 *The previous calculation can be extended to the more general wave equation (7):*

$$u_{tt} + ku_t - \Lambda u = f$$

with $k > 0$ constant, Λ self-adjoint and f a white noise. In particular, $C_{A,B}(\tau)$ is an **even** function of τ . This is the basis of

- *Clock synchronization [22]: if the clocks at the stations A and B are not synchronized, the correlation is even with respect to a shifted origine of times.*
- *Test of applicability of the theory: in the semi-classical regime studied later, approximate evenness holds. In fact the oscillating part is even because the corrections to the Green’s function do not affects the phases. In Theorem 4.2 the principal symbol π is ≥ 0 . This is basically due to the time reversibility of the ray dynamics.*

5.2 General wave equations: semi-classics

5.2.1 Assumptions

We want to derive results similar to those of Theorem 4.2 in the case of (multi-component) wave equations. For that, we will introduce a matrix version of what is done in Section 4 and use the results of Section 8.7. Let us start with the following case of Equation (1):

$$\frac{\varepsilon}{i} \mathbf{u}_t + \hat{H} \mathbf{u} = \frac{\varepsilon}{i} \mathbf{f} \quad (13)$$

where

- \hat{H} is a matrix of ε -pseudo-differential operators:

$$\hat{H} := \text{Op}_\varepsilon(H_0 + \varepsilon H_1)$$

with (see Appendix A, Section 8.7)

- The **principal symbol** $H_0(x, \xi) : T^*X \rightarrow \text{Herm}(\mathbb{C}^N)$ belongs to $\Sigma_m^{\mathbb{C}^N, \mathbb{C}^N}$ and is elliptic of degree m with the eigenvalues $\lambda_1(x, \xi) < \dots < \lambda_j(x, \xi) < \dots < \lambda_n(x, \xi)$ of constant multiplicities m_j ; the corresponding eigenspaces $E_j(x, \xi) \subset \mathbb{C}^N$, of dimension m_j , are called the *polarizations*. This assumption is satisfied for the elastic wave equation in a manifold of dimension 3 with $N = 6$: 2 eigenvalues $\pm \lambda_S$ of multiplicity 2 (the S -waves) and 2 eigenvalues $\pm \lambda_P$ of multiplicity 1 (the P -waves).
- The **sub-principal symbol** $H_1(x, \xi) \in \Sigma_m^{\mathbb{C}^N, \mathbb{C}^N}$ admits some positivity property which controls the attenuation: there exists $k > 0$, such that the sub-principal symbols $H_{1,j}$, $j = 1, \dots, n$, satisfy

$$\Im H_{1,j} \leq -k < 0 .$$

- The **random field** \mathbf{f} is given by $\mathbf{f} = \text{Op}(l(x, \xi, \omega))\mathbf{w}$:
 - The noises $\mathbf{w} = (w_1, \dots, w_N)$ are independent white noises on $X \times \mathbb{R}$
 - The matrix l is smooth, compactly supported w.r. to (x, ξ) and its Fourier transform w.r. to ω is compactly supported. The power spectrum of f is $(2\pi\varepsilon)^{-(d+1)}ll^*(x, \xi, \omega)$.

5.2.2 The main result

Using the gauge transform reducing \hat{H} to the block diagonal normal form of Section 8.7.2, we can decompose the correlation as blocks $C_{A,B}^{j,k}(\tau)$, $j, k = 1, \dots, n$.

Theorem 5.1 *With the assumptions of Section 5.2.1 the correlations $C_{A,B}^{j,k}(\tau)$ are given*

- If $j = k$, as in Theorem 4.2 with $H_0 = \lambda_j$
- If $j \neq k$ and $l(x, \xi, \omega)$ vanishes for x near B : $C_{A,B}^{j,k}(\tau) = O(\varepsilon^\infty)$

The proof makes a strong use of the normal form of Section 8.7.2: we use a block decomposition of $\mathcal{L} = (\mathcal{L}_{ij})$ in Equation (2) and the block diagonal reduction $\Omega(t) = (\Omega_i(t))$. We are reduced to the study of

$$\Omega_i(\tau) \int_0^\infty \Omega_i(t) \mathcal{L}_{ij} \Omega_j^*(t) dt .$$

If $i = j$, we are in a situation very close to that of a scalar Hamiltonian (Section 4.4), while if $i \neq j$, we can use Lemma 10.2.

In particular, for the elastic wave equation, if the source noise is supported away from B , the correlation $C_{A,B}^{P,S}$ between S-waves and P-waves is very small.

6 Coda's correlations

In [5], M. Campillo and A. Paul computed the correlation of *coda waves* and saw the emergence of the Green's function.

Let $u(x, t)$ be the solution of the wave equation with initially localized Cauchy data at time $t = 0$ (an "earthquake") The coda is the field $u(x, t)$ in some window $[T_1, T_2]$ with $T_1, T_2 \gg 0$. Let us look at the correlation

$$C_{A,B}(\tau) = \frac{1}{T_2 - T_1} \int_{T_1}^{T_2} u(A, t + \tau) \bar{u}(B, t) dt .$$

Introducing the propagator P , we get:

$$C_{A,B}(\tau) = \int_X |dx| P(\tau, A, x) \left(\frac{1}{T_2 - T_1} \int_{T_1}^{T_2} u(x, t) \bar{u}(B, t) dt \right) .$$

So that the question is reduced to the study of the limits of $\frac{1}{T_2 - T_1} \int_{T_1}^{T_2} u(A, t) \bar{u}(B, t) dt$ as $T_1, T_2, T_2 - T_1 \rightarrow \infty$. If we have some nice (smooth?) limit $K(A, B)$, then $C_{A,B}(\tau) = [\Omega(\tau) \circ \hat{K}](A, B)$.

Looking at some semi-classical limit, this is closely related to the following result by R. Schubert [26]: *if the classical dynamics is hyperbolic and t of the order of the Ehrenfest time, the Wigner measures of $u(t, \cdot)$ are smooth and uniform on each energy shell*. Does this uniformity of the repartition of energy extends to the wave functions themselves, in particular to the phases?

7 Using surface waves in order to image the inner crust

In real applications to seismology, the fields are recorded at the surface of the earth. Moreover the main contribution to the correlations comes from the surface waves:

- They have a smaller time decay than the body waves.
- The source of the noise itself (interaction with atmosphere and ocean) is located at the surface.

The earth's crust acts as a wave guide. This implies that we really have to look at a 2D version of the wave equation: the wave equation for surface wave is an effective wave equation on the surface. This equation is given in terms of eigenvalues of a vertical Sturm-Liouville equation. Solving an inverse spectral problem for these Sturm-Liouville equation allows (in principle!) to image the crust. The inverse spectral problem is solved in [7].

In this Section, we will discuss these effective Hamiltonians in the scalar model of the acoustic wave equation and show how one can reduce the problem of imaging earth's crust to an inverse spectral problem.

7.1 A mathematical model

We work locally in $X = \{(\mathbf{x}, z) \in \mathbb{R}^2 \times \mathbb{R} \mid z \leq 0\}$ so that $z = 0$ is the surface. We will consider the very simple case of an acoustic wave equation near the origin of X . Let us give a function

$$N(\mathbf{x}, Z) : \mathbb{R}^2 \times \mathbb{R}_- \rightarrow \mathbb{R}_+$$

which satisfies the following assumptions:

1. There exists $N_\infty > 0$ and $Z_0 < 0$, so that $N(\mathbf{x}, Z) = N_\infty$ for $Z \leq Z_0$.
2. For every $\mathbf{x} \in \mathbb{R}^2$, $\inf_Z N(\mathbf{x}, Z) = N_0(\mathbf{x}) > 0$ and $N_0(\mathbf{x}) < N_\infty$ for \mathbf{x} close to 0
3. N is C^1 w.r. to Z and the map $\mathbf{x} \rightarrow N(\mathbf{x}, \cdot)$, from \mathbb{R}^2 into the Banach space $C^1([Z_0, 0], \mathbb{R})$, is smooth.

Let us introduce a small parameter $\varepsilon > 0$ and the following acoustic wave equation where

$$n(\mathbf{x}, z) = N(\mathbf{x}, \frac{z}{\varepsilon}) : \tag{14}$$

$$\begin{cases} u_{tt} - \operatorname{div}(n \operatorname{grad} u) = 0 \\ \frac{\partial u}{\partial z}(\mathbf{x}, 0) = 0 \end{cases} \tag{15}$$

We plan to see that Equation (15) admits, as $\varepsilon \rightarrow 0$, asymptotic solutions of frequency of order ε^{-1} located near the surface. Moreover these solutions are determined by solving an effective pseudo-differential equation on the surface $\partial X = \mathbb{R}^2 \times \{0\}$. In order to show that we can apply Theorem 8.3 of Appendix A, we will first introduce the operator H_0 in Section 7.3. It will be a Sturm-Liouville operator in the variable Z .

7.2 Discussion of the assumptions on the model

The scalar acoustic wave equation (15) is a simplified model of the elastic wave equation. We could work with the elastic wave equation, but it will be more complicated and the main ideas would be the same.

The function n represent the square of the local propagation speed (group velocity). The formula (14) means that the speed admits fast variations in the vertical

directions and is smooth in the horizontal direction. It is a plausible assumption that the geological structure near the surface of the earth is “stratified”.

Assumption (1) on N means that the speed is constant far from the surface. It is an assumption which could be removed at the price of a technical work.

Assumption (2) implies that the speed is somewhere smaller than the speed at large depth. This kind of assumption is usually satisfied in seismology where the speed of elastic waves into sediments layers is smaller than the speed inside the rocks below the sediments.

Assumption (3) means that we need less regularity in the vertical direction. It could be relaxed using piecewise smoothness w.r. to Z , but at the price of a rather technical work: the domain of the Sturm-Liouville operator introduced below will then depend on \mathbf{x} .

7.3 A family of Sturm-Liouville operators

Let us consider, for each $(\mathbf{x}, \xi) \in T^*\partial X$, the self-adjoint differential operator $L_{\mathbf{x}, \xi}$ on the half line $Z \leq 0$, with Neumann boundary condition at $Z = 0$, defined by:

$$L_{\mathbf{x}, \xi} v := -\frac{d}{dZ} \left(N(\mathbf{x}, Z) \frac{dv}{dZ} \right) + N(\mathbf{x}, Z) |\xi|^2 v . \quad (16)$$

The domain $D \subset L^2(\mathbb{R}^-)$ of $L_{\mathbf{x}, \xi}$ is

$$D := H^2(\mathbb{R}^-) \cap \{v \mid v'(0) = 0\}$$

which is independent of (\mathbf{x}, ξ) .

The spectrum of $L_{\mathbf{x}, \xi}$ consists of a finite discrete spectrum and a continuous spectrum $[N_\infty |\xi|^2, +\infty[$. $L_{\mathbf{x}, \xi}$ admits, for ξ large enough, a non empty discrete spectrum of simple eigenvalues

$$N_0(\mathbf{x}) |\xi|^2 < \lambda_1(\mathbf{x}, \xi) < \dots < \lambda_j(\mathbf{x}, \xi) < \dots < \lambda_k(\mathbf{x}, \xi) < N_\infty |\xi|^2 ,$$

which depend smoothly of (\mathbf{x}, ξ) . In order to see that, we can interpret $|\xi|^{-2} L_{\mathbf{x}, \xi}$ as a semi-classical Schrödinger type operator with an effective Planck constant $|\xi|^{-1}$ and a principal symbol

$$p_{\mathbf{x}}(Z, \zeta) = N(\mathbf{x}, Z)(\zeta^2 + 1)$$

which admits a “well” because of Assumption 2 in Section 7.1 on N .

We should however take care of the fact that the number k of eigenvalues depends on (\mathbf{x}, ξ) and goes to ∞ as ξ does. This will give birth to mode conversions between surface waves and body waves. To my knowledge, such mode conversions have not yet been studied from a mathematical point of view.

7.4 Wave guides as an adiabatic limit

We will rewrite Equation (15) as an operator valued pseudo-differential equation as in Appendix A, Section 7.4, in the spirit of the book [29], Chapter 3 on “space-adiabatic perturbation theory”.

Putting

$$u(t, \mathbf{x}, z) = v(t, \mathbf{x}, \frac{z}{\varepsilon})$$

we get the following wave Equation for v :

$$\varepsilon^2 v_{tt} - (\partial_z (N \partial_z v) + \varepsilon^2 N \Delta_{\mathbf{x}} v + \varepsilon^2 \text{grad}_{\mathbf{x}} N \cdot \text{grad}_{\mathbf{x}} v) = 0 ,$$

which can be rewritten, using Weyl quantization in the variable \mathbf{x} :

$$\varepsilon^2 v_{tt} + \text{Op}_\varepsilon (L_{(\mathbf{x}, \xi)} - i\varepsilon \xi \cdot \text{grad}_{\mathbf{x}} N) v = 0 .$$

Putting $H_0((\mathbf{x}, \xi)) = L_{(\mathbf{x}, \xi)}$ and $H_1((\mathbf{x}, \xi)) = -i\xi \cdot \text{grad}_{\mathbf{x}} N$, we will to apply Theorem 8.3 in Appendix A, Let us choose, for (\mathbf{x}, ξ) close to 0, some eigenvalue $\lambda_j(x, \xi)$ of $L_{(\mathbf{x}, \xi)}$ with a normalized eigenfunction $\varphi_j(\mathbf{x}, \xi, Z)$. Then we have the existence of a propagation mode of the form

$$v(t, \mathbf{x}, z) = w(t, \mathbf{x}) (\varphi_j(\mathbf{x}, \xi, z/\varepsilon) + \varepsilon \cdots)$$

where w satisfying the **scalar** wave equation:

$$\varepsilon^2 w_{tt} + \text{Op}_{\varepsilon}(\lambda_j(x, \xi) + \cdots)w = 0 .$$

For each eigenvalue λ_j , we can then apply the result of Section 5.

From the passive imaging method, we can hope to recover the effective Hamiltonians $\pm\sqrt{\lambda_j}$. Solving an inverse spectral problem for the Sturm-Liouville equation allows in principle to recover the function n which encodes the physics of the propagation medium.

Let us put $\hbar = |\xi|^{-1}$ and, for $(\mathbf{x}, \xi) = (\mathbf{x}_0, \xi_0)$ fixed, $\mathcal{L}_{\hbar} = \hbar^2 L_{(\mathbf{x}_0, \xi_0)}$. This operator is a classical ΨDO on \mathbb{R}^+ of principal symbol $N(\mathbf{x}_0, \xi_0, Z)(\zeta^2 + 1)$.

From the paper [7], we know that the function N can be recovered from the $\hbar \rightarrow 0$ asymptotics of the λ_j 's.

8 Appendix A: a review about pseudo-differential operators

We review the basic definitions and properties of pseudo-differential operators (ΨDO 's). Pseudo-differential operators (first called singular integral operators by Calderon and Zygmund) were developed in order to get a class of operators generalizing linear differential operators and containing approximate inverse of elliptic operators. They became a basic tool in the theoretical study of linear differential operators with variable coefficients. A more general theory was then developed including a small parameter, here denoted by $\varepsilon \in]0, \varepsilon_0]$. Several references are now available [6, 10, 11, 30].

8.1 Basic calculus

We will define the pseudo-differential operators (ΨDO 's) on \mathbb{R}^d . ΨDO 's on manifolds are defined locally by the same formulas.

Definition 8.1 (Classical symbols) • *The space Σ_k of symbols of degree k is the space of smooth functions $p : T^*\mathbb{R}^d \rightarrow \mathbb{C}$ which satisfy*

$$\forall \alpha, \beta \in \mathbb{N}^d, |D_x^\alpha D_\xi^\beta p(x, \xi)| \leq C_{\alpha, \beta} \langle \xi \rangle^k ,$$

with $\langle \xi \rangle = 1 + \|\xi\|$.

- *A classical symbol of degree m is a family of functions*

$$p_\varepsilon : T^*\mathbb{R}^d \rightarrow \mathbb{C}$$

depending of some small parameter $\varepsilon > 0$ which admits, as $\varepsilon \rightarrow 0$, an asymptotic expansion

$$p_\varepsilon \equiv \sum_{j=0}^{\infty} \varepsilon^j p_j(x, \xi)$$

with $p_j \in \Sigma_m$. We will denote this space by S_{class}^m .

Definition 8.2 (Pseudo-differential operators) A classical ε -pseudo-differential operator P (a ΨDO) of degree m on \mathbb{R}^d is given by the (distribution) kernel

$$[P](x, x') = (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^d} e^{i(x-x'|\xi)/\varepsilon} p_\varepsilon\left(\frac{x+x'}{2}, \xi\right) |d\xi|$$

where $p_\varepsilon(x, \xi)$, the so-called (total) symbol of P , is in S_{class}^m .

We will denote $P = \text{Op}(p_\varepsilon)$ and by Ψ_{class}^m the space of classical ΨDO 's of degree m . The operator $P_\varepsilon = \text{Op}_\varepsilon(p_\varepsilon)$ is called the Weyl-quantization of p_ε .

The operators P_ε act from C_o^∞ into C^∞ . The kernel of P_ε is given by:

$$[P_\varepsilon](x, x') = (2\pi\varepsilon)^{-d} \tilde{p}_\varepsilon\left(\frac{x+x'}{2}, \frac{x'-x}{\varepsilon}\right) \quad (17)$$

with \tilde{p}_ε the partial Fourier transform of $p_\varepsilon(x, \xi)$ w.r. to ξ . Very often, one is only able to compute the symbol p_0 which is called the *principal symbol* of P . The term p_1 is called the *sub-principal symbol*.

The most basic fact about ΨDO 's is the fact they form a (non-commutative) algebra: if $P_\varepsilon = \text{Op}(p_\varepsilon)$ and $Q_\varepsilon = \text{Op}(q_\varepsilon)$, we have $P_\varepsilon Q_\varepsilon = \text{Op}(p_\varepsilon q_\varepsilon + O(\varepsilon))$. The composition formula for the total symbol is called the *Moyal \star -product* $\text{Op}(a) \circ \text{Op}(b) = \text{Op}(a \star b)$ and is given by

$$a \star b \equiv ab + \frac{\varepsilon}{2i} \{a, b\} + \sum_{j=2}^{\infty} \varepsilon^j P_j(a, b)$$

where

- $\{a, b\}$ is the Poisson bracket
- $P_j(a, b)$ is a bi-linear bi-differential operator, homogeneous of degree j with respect to a and b .

8.2 The symbols S_δ^m

In order to formulate the Egorov Theorem for long times, we will need more sophisticated classes of symbols.

Definition 8.3 (S_δ^m) For any real number m and any δ with $0 \leq \delta < \frac{1}{2}$, a symbol $a \in S_\delta^m$ is a smooth function on T^*X depending on ε which satisfies

$$\forall \alpha \in \mathbb{N}^{2d}, \exists C_\alpha > 0, \forall (x, \xi) \in T^*X, |D_{x, \xi}^\alpha a(x, \xi)| \leq C_\alpha \varepsilon^{-|\alpha|} \langle \xi \rangle^m.$$

We have $S_{\text{class}}^m \subset S_0^m$.

We can associate to $a \in S_\delta^m$ a pseudo-differential operator $\text{Op}(a) \in \Psi_\delta^m$ using formula (17). Such operators obey a nice calculus, see [10, 14]: if, for $j = 1, 2$, $a_j \in S_{\delta_j}^{m_j}$, $A_j := \text{Op}(a_j)$ and $\delta := \max(\delta_1, \delta_2)$, we have $A_1 \circ A_2 = \text{Op}(a_1 \star a_2)$ with $a_1 \star a_2 \in S_\delta^{m_1+m_2}$ and admits the asymptotic expansion given by the Moyal \star -product. We have $P_j(a_1, a_2) \in \varepsilon^{-j(\delta_1+\delta_2)} S_\delta^{m_1+m_2}$. In particular $[A_1, A_2] = \frac{\varepsilon}{i} \text{Op}(\{a_1, a_2\}) + \text{Op}(b)$ with $b \in \varepsilon^{2(1-\delta_1-\delta_2)} S_\delta^{m_1+m_2}$.

Theorem 8.1 (Calderon-Vaillancourt) The operators of Ψ_δ^0 extend to operators from L^2 to L^2 , uniformly continuous in ε . The operators in Ψ_δ^m extends to operators from L^2 into the Sobolev space H^{-m} with a norm of order $O(\varepsilon^m)$.

8.3 Asymptotic expansions

The basic tool in order to find a ΨDO is an inductive construction based on the realization of any asymptotic expansion: given a sequence $a_j \in S_\delta^m$ and $l_1 = 0 < l_2 < \dots < l_j < \dots$ and $l_j \rightarrow +\infty$ as $j \rightarrow \infty$, there exists $a \in S_\delta^m$ such that

$$a \equiv \sum_{j=1}^{\infty} \varepsilon^{l_j} a_j$$

meaning that

$$\forall n, a - \sum_{j=1}^{n-1} \varepsilon^{l_j} a_j \in \varepsilon^{l_n} S_\delta^m .$$

Such an a is uniquely defined modulo a remainder term r which lies in $\varepsilon^\infty S_\delta^m$. The corresponding $\Psi DO \text{ Op}_\varepsilon(a)$ is well defined modulo $O(\varepsilon^\infty)$. This implies, if $m \leq 0$ and a compactly supported in x , that the $L^2 \rightarrow L^2$ norm of $\text{Op}_\varepsilon(r)$ is $O(\varepsilon^\infty)$.

8.4 Micro-localization

A family u_ε of functions on X is said to be *admissible* if the L^2 norm on any compact set is at most of polynomial growth with respect to ε^{-1} .

The functions u_ε are said to be $O(\varepsilon^\infty)$ near a point $(x, \xi) \in T^*X$ if $\exists \varphi \in C_o^\infty(X)$, $\varphi(x_0) \neq 0$ and the ε -Fourier transform $\mathcal{F}_\varepsilon(\varphi u_\varepsilon)(\xi) := (2\pi\varepsilon)^{-d/2} \int e^{-i\langle x|\xi \rangle/\varepsilon} u_\varepsilon(x) dx$ satisfies $\mathcal{F}_\varepsilon(\varphi u_\varepsilon)(\xi) = O(\varepsilon^\infty)$ uniformly for ξ close to ξ_0 .

The micro-support $\text{MS}(u_\varepsilon)$ of u_ε is the set of all (x, ξ) 's at which u_ε is not $O(\varepsilon^\infty)$. For any $\Psi DO A$, we have $\text{MS}(Au_\varepsilon) \subset \text{MS}(u_\varepsilon)$.

If U is an open subset of T^*X and if u_ε and v_ε are admissible, we will shortly denote $u_\varepsilon = v_\varepsilon + O(\varepsilon^\infty)$ in U if $\text{MS}(u - v) \cap U = \emptyset$.

8.5 Functional calculus and and an useful Lemma

A simple case of the functional calculus for ΨDO 's is as follows:

Theorem 8.2 *Let $f \in C_o^\infty(\mathbb{R})$ and $A = \text{Op}_\varepsilon(a_0 + \varepsilon a_1 + \dots) \in \Psi_{\text{class}}^m$ with a_0 elliptic and real. Then one can define $f(A)$ using the spectral theory and $f(A) \in \Psi_{\text{class}}^{-\infty}$ with principal symbol $f(a_0)$.*

Remark 8.1 *The previous Theorem is usually stated for A self-adjoint (see [28] Chap. 8). It can be extended to our case with a very close proof.*

Lemma 8.1 *Let us consider the operator \mathcal{L} , given by*

$$\mathcal{L} := \frac{1}{2\pi\varepsilon} \int \int \Omega(-t) e^{it\omega/\varepsilon} \text{Op}_\varepsilon^x(L(x, \xi, \omega)) d\omega dt$$

with $L \in S_{\text{class}}^{-\infty}(X \times \mathbb{R})$ with $L(x, \xi, \omega)$ having an ω Fourier transform compactly supported as a function of (x, ξ, t) . Here $\Omega(t) = \exp(-it\hat{H}/\varepsilon)$ where \hat{H} satisfies the assumptions of Section 4.2. Then \mathcal{L} belongs to $\Psi_{\text{class}}^{-\infty}$ and his principal symbol is given by $l_0(x, \xi) = L(x, \xi, -H_0(x, \xi))$.

Using a change of variable $t = \varepsilon t'$ in the Definition of \mathcal{L} , the proof is a simple corollary of the functional calculus of ΨDO 's: if $A = \text{Op}(a)$ with a compactly supported, the operator $B = \exp(-it'\hat{H})A$ belongs to $\Psi_{\text{class}}^{-\infty}$ and the principal symbol of B is $b = \exp(-it'H_0)a$; because the function $\exp(-it')$ is not compactly supported, we need first to rewrite $A = F(\hat{H})A$ with $F \in C_o^\infty$ and then use the functional calculus with the function $f = F \exp(-it')$.

8.6 Hilbert-Schmidt estimate

Lemma 8.2 *If $P \in \Psi_\delta^{-\infty}(X)$ with X a compact manifold, P is Hilbert-Schmidt with $\|P\|_{\text{H-S}} = O(\varepsilon^{-d/2})$.*

Proof.–

$P^*P \in \Psi_\delta^{-\infty}$ is trace class (smooth kernel) and its trace is easily seen to be $O(\varepsilon^{-d})$. □

8.7 Multi-component operators

8.7.1 Systems of Ψ DO

Definition 8.4 *Let E and F be 2 Banach spaces $\Sigma_m^{E,F}$ is the space of symbols of degree m with values in the Banach space $\mathcal{L}(E, F)$ of continuous linear maps from E into F .*

Similarly, we can define the spaces $S_{\text{class}}^{m,E,F}, \Psi_{\text{class}}^{m,E,F}, \dots$. The operators of $\Psi_{\text{class}}^{m,E,F}$ send functions on X with values in E into functions on X with values in F .

Definition 8.5 $H_0 \in \Sigma_m^{E,F}$ is elliptic if $H_0(x, \xi)$ is invertible for large ξ 's and the inverse is in $\Sigma_{-m}^{F,E}$.

Matrix valued Ψ DO's have properties similar to scalar Ψ DO's.

8.7.2 Reduction

Let \mathcal{H} be an Hilbert space and $D \subset \mathcal{H}$ a dense sub-space.

Definition 8.6 *A matrix valued Hamiltonian $H_0 : T^*X \rightarrow \text{Herm}(\mathcal{H})$ is an element of $\Sigma_m^{D,\mathcal{H}}$: for every $(x, \xi) \in T^*X$, $H_0(x, \xi)$ is a self-adjoint operator with domain $D \subset \mathcal{H}$.*

Let us consider a classical symbol $H_\varepsilon \equiv H_0 + \varepsilon H_1 + \dots \in S_{\text{class}}^{m,D,\mathcal{H}}$ with H_0 a matrix values Hamiltonian. Let us consider near some point $z_0 \in T^*X$ an eigenvalue $\lambda_0(z)$ of constant multiplicity m and eigenspace $E(z)$ (the polarization space) and assume that $\lambda_0(z)$ is an isolated point in the spectrum of $H_0(z)$.

Then we have, inspired from [13, 29]:

Theorem 8.3 *In some neighborhood of z_0 , there exists*

- *A $\mathcal{L}(\mathcal{H})$ -valued Ψ DO $\hat{U} = \text{Op}_\varepsilon(U_0 + \varepsilon U_1 + \dots)$ with U_0 unitary.*
- *A classical symbol in $S_{\text{class}}^{m,\mathbb{C}^m,\mathbb{C}^m}$ $\lambda(z) = \lambda_0(z)\text{Id} + \varepsilon \lambda_1(z) + \dots$*
- *A symbol K with values in $\mathcal{L}(D \cap F, F)$ with $F = E^\perp$ with $\lambda_0(z) \notin \text{Spectrum}(K_0(z))$,*

so that

$$U^{-1} \text{Op}_\varepsilon(H) U \equiv \begin{pmatrix} \text{Op}_\varepsilon(\lambda) & 0 \\ 0 & \text{Op}_\varepsilon(K) \end{pmatrix}$$

near z_0 .

Moreover $\lambda_1(z) = \pi H_1 \pi + R_1$ with π the orthogonal projector onto E and R_1 self-adjoint.

The proof is done by induction on the powers of ε : at the first step, we choose an unitary gauge transform $U(z) \in \mathcal{L}(\mathcal{H})$ so that

$$U^{-1}H_0U = \begin{pmatrix} \lambda_0 \text{Id} & 0 \\ 0 & K_0 \end{pmatrix}$$

with $\lambda_0 \text{Id} - K_0$ invertible. At the next step, we try to find

$$B = \begin{pmatrix} 0 & b \\ c & 0 \end{pmatrix}$$

so that

$$(1 - \varepsilon B) \star \left(\begin{pmatrix} \lambda_0 \text{Id} & 0 \\ 0 & K_0 \end{pmatrix} + \varepsilon \tilde{H}_1 \right) \star (1 + \varepsilon B) = \begin{pmatrix} \lambda_0 \text{Id} + \varepsilon \lambda_1 & 0 \\ 0 & K_0 + \varepsilon K_1 \end{pmatrix} + o(\varepsilon^2).$$

Collecting terms in ε^1 , we get the following equations:

$$b(\lambda_0 \text{Id} - K_0) = \beta, \quad (\lambda_0 \text{Id} - K_0)c = -\gamma$$

with

$$\tilde{H}_1 = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}.$$

These equations can be solved with the assumption that $\lambda_0 \text{Id} - K_0$ is invertible.

Remark 8.2 *In the Theorem 8.3, we do not need that \hat{H} is self-adjoint (only H_0 is self-adjoint), neither that the symbol of \hat{H} takes values into bounded operators on \mathcal{H} .*

Remark 8.3 *The Theorem 8.3 will be used as well to reduce a multi-component wave equation to scalar equation in Section 5.2 and as tool for adiabatic limits in Section 7.4*

Corollary 8.1 *Assuming that all eigenvalues $\lambda_1 < \dots < \lambda_j < \dots < \lambda_n$ of $H_0(x, \xi)$ are of constant multiplicities m_j , $j = 1, \dots, n$, the Hamiltonian \hat{H} admits a micro-local normal form with diagonal blocks of sizes $m_j \times m_j$ and principal symbols of the j -th block $\lambda_j \text{Id}_{\mathbb{C}^{m_j}}$.*

The sub-principal symbol of the j -th block is $H_{1,j} := \pi_j H_1 \pi_j + R_j$ where R_j is self-adjoint while π_j is the orthogonal projection on E_j ; it means that, as quadratic forms, $H_{1,j}$ is the restriction of H_1 to E_j .

The formula for the sub-principal symbols is a direct consequence of the paper [13].

8.7.3 Wave equations

Let us consider the wave equation (7)

$$u_{tt} + Au_t - \Lambda u = 0$$

with $\varepsilon^2 \Lambda = \text{Op}_\varepsilon(H_0)$ (H_0 elliptic) and assume w.l.o.g. that $-\Lambda$ is > 0 so that $\Lambda_{\frac{1}{2}} = \varepsilon \sqrt{-\Lambda}$ is also a ΨDO . Putting

$$U := \begin{pmatrix} \Lambda_{\frac{1}{2}} u \\ -i\varepsilon u_t \end{pmatrix}$$

Equation (7) rewrites:

$$\frac{\varepsilon}{i} U_t + \begin{pmatrix} O & -\Lambda_{\frac{1}{2}} \\ -\Lambda_{\frac{1}{2}} & -i\varepsilon A \end{pmatrix} U = 0,$$

which is of the form

$$\frac{\varepsilon}{i} U_t + \hat{H}U = 0$$

with \hat{H} an elliptic ΨDO .

8.7.4 Extension of Lemma 8.1

Let us consider $a \in \mathcal{S}(\mathbb{R}_\omega, L(\mathbb{C}^N))$ and H_0 an $N \times N$ Hermitian matrix, we define $a(H_0)$ as

$$a(H_0) := \frac{1}{2\pi} \int e^{itH_0} \mathcal{F}a(t) dt .$$

The following Lemma is proved the same way as Lemma 8.1.

Lemma 8.3 *Let us consider the operator \mathcal{L} , given by*

$$\mathcal{L} := \frac{1}{2\pi\varepsilon} \iint \Omega(-t) e^{it\omega/\varepsilon} \text{Op}_\varepsilon^x(L(x, \xi, \omega)) d\omega dt$$

with $L \in S_{\text{class}}^{-\infty, \mathbb{C}^N, \mathbb{C}^N}(X \times \mathbb{R})$ with $L(x, \xi, \omega)$ having an ω Fourier transform compactly supported as a function of (x, ξ, t) . Here $\Omega(t) = \exp(-it\hat{H}/\varepsilon)$ where \hat{H} satisfies the assumptions of Section 5.2. Then \mathcal{L} belongs to $\Psi_0^{-\infty, \mathbb{C}^N, \mathbb{C}^N}$ and his principal symbol is given by $l_0(x, \xi) = L(x, \xi, -H_0(x, \xi))$.

9 Appendix B: asymptotics of the Green's function and classical dynamics

9.1 The scalar case

We want to describe quite explicitly the propagator $P(t, x, y) = [\Omega(t)](x, y)$ with $\Omega(t) = \exp(-it\hat{H}/\varepsilon)$, \hat{H} as in 4.2. We will do that in the energy interval I ; it means that we will only describe the action of $\Omega(t)$ on functions whose micro-supports are in $H_0^{-1}(I)$.

To the Hamiltonian $H_0 : T^*X \rightarrow \mathbb{R}$, we associate the *ray dynamics* defined by

$$\frac{dx_j}{dt} = \frac{\partial H_0}{\partial \xi_j}, \quad \frac{d\xi_j}{dt} = -\frac{\partial H_0}{\partial x_j} . \quad (18)$$

Let us denote by ϕ_t the flow of the Hamiltonian vector field X_{H_0} defined in Equation (18). The main result is that $P(t, \dots)$ is a *Fourier Integral Operator* associated to the flow ϕ_t . For simplicity, we will only describe the so-called *generic* case where P admits a WKB expansion.

If $x_0, y_0 \in X$ and $\phi_t(y_0, \eta_0) = (x_0, \xi_0)$, we say that x_0 and y_0 are *not conjugate* along the ray $\gamma(s) = \phi_s(y_0, \eta_0)$, $0 \leq s \leq t$, if the differential $\delta\eta \rightarrow d\phi_t(0, \delta\eta)$ is invertible at the point (y_0, η_0) . This implies that for (x, y) close enough to (x_0, y_0) there exists a unique (y, η) close to (y_0, η_0) with $\phi_t(y, \eta) = (x, \xi)$. In this case the action $S_\gamma(t, x, y)$ is defined by $S_\gamma(t, x, y) = \int_\gamma \xi dx - H_0(x, \xi) dt$. $S_\gamma(t, \dots)$ is a *generating function* of ϕ_t near (y_0, η_0) , meaning that $\phi_t(y, -\partial S_\gamma / \partial y) = (x, \partial S_\gamma / \partial x)$ for (y, η) close to (y_0, η_0) .

The action S_γ is useful in order to describe a WKB expansion of the propagator $P(t, x, y)$ for (x, y) close to (x_0, y_0) : $P(t, x, y)$ is a sum of contributions $P_\gamma(t, x, y)$ of rays going from y to x in the time t .

Assuming that x and y are not conjugate along γ , we have

$$P_\gamma(t, x, y) \equiv (2\pi\varepsilon)^{-d/2} \left(\sum_{j=0}^{\infty} \varepsilon^j a_{j,\gamma}(t, x, y) \right) e^{iS(t,x,y)/\varepsilon} .$$

Definition 9.1 *We will say that (x, y, t) is generic if x and y are not conjugate along any γ so that $\gamma(0) = x$, $\gamma(t) = y$ and $\gamma \subset H_0^{-1}(I)$.*

If (x, y, t) is generic, there is only a finite number of rays from y to x in time t included in $H_0^{-1}(I)$.

Dispersive rays: we say that the dynamics generated by H_0 is dispersive if $(H_0)_{\xi\xi}$ is invertible. This means that the speed $\partial H_0/\partial\xi$ depends on ξ ! If this is the case, (t, x, y) is generic for x close enough to y .

The dispersion relation: the principal symbol of the wave equation is $\omega + H_0(x, \xi) : T^*(X \times \mathbb{R}) \rightarrow \mathbb{R}$. It is clear enough that the ray dynamics is determined by the hyper-surface $\mathbb{D} := \{(x, \xi, \omega) | \omega + H_0(x, \xi) = 0\}$. This surface is called the *dispersion relation*.

9.2 The matrix case

Let us assume that H_0 is matrix valued as in Section 5.2. The dispersion relation is then the hyper-surface

$$\det(\omega \text{Id} - \mathbf{H}_0(x, \xi)) = \mathbb{D}(x, \xi, \omega) = 0 ,$$

we can (at least outside a *mode conversion set*), rewrite

$$\mathbb{D}(x, \xi, \omega) = \prod_{j=0}^M (\omega - \lambda_j(x, \xi))^{d_j}$$

which gives M ray dynamics λ_j associated to different polarizations

$$E_j = \ker(\mathbf{H}_0(x, \xi) - \lambda_j(x, \xi)\text{Id}) .$$

10 Appendix C: long time Egorov Theorem

The main reference for this part is [3] where the authors get only an operator norm estimates for the remainder. The proofs of reference [3] can be extended in order

- to get a point-wise result under some ellipticity assumption
- to cover the case of manifolds
- to use a local Liapounov exponent: the Liapounov exponent in the some energy interval I
- to allow an attenuation term.

10.1 Egorov theorem

Definition 10.1 *An admissible Hamiltonian of degree $m > 0$ is $\hat{H} := \text{Op}_\varepsilon(H_0 + \varepsilon H_1)$ with*

- $H_0 \in \Sigma_m$ a real valued function which is elliptic:

$$\exists C > 0, \forall (x, \xi) \in T^*X, H_0(x, \xi) \geq C \langle \xi \rangle^m .$$

- $H_1 \in \Sigma_{m-1}$ and $h_1 = \Im H_1 \leq -k \leq 0$ everywhere in T^*X .

Example 10.1 *A typical admissible Hamiltonian is the damped Schrödinger operator $\hat{H} = -\varepsilon^2 \Delta + V(x) - i\varepsilon a(x)$ with $V(x) \geq 0$ and $a(x) \geq 0$.*

We define the dynamics, for $t \geq 0$, by:

$$\Omega(t) = \exp(-it\hat{H}/\varepsilon) .$$

Theorem 10.1 (Egorov Theorem) *Let us give \hat{H} an admissible Hamiltonian and $A(0) = \text{Op}(a)$ with $a \in C_o^\infty(H_0^{-1}(I))$ with I a compact interval. Then the operator $A(t) = \Omega(t)A(0)\Omega^*(t)$ is, for t fixed, a pseudo-differential operator whose symbol $a(t) \equiv \sum_{j=0}^\infty \varepsilon^j a_j(t)$ belongs to $S_{\text{class}}^{-\infty}$ and satisfies*

$$a_0(t)(z) = \exp\left(2 \int_{-t}^0 h_1(\Phi_s(z)) ds\right) a(\Phi_{-t}(z)) . \quad (19)$$

Let us remark that the previous result is a priori *uniform* only in a fixed interval $t \in [0, T_0]$ independent of ε .

Let us sketch the proof. $A(t)$ is the value for $s = t$ of the solution $B(s)$ of the ODE

$$\frac{d}{ds} \Omega(t-s)B(s)\Omega^*(t-s) = 0, \quad B(0) = A(0) .$$

We have

$$\frac{d}{ds} \Omega(t-s)B(s)\Omega^*(t-s) = \Omega(t-s) \left(B'(s) + \frac{i}{\varepsilon} (\hat{H}B(s) - B(s)\hat{H}^*) \right) \Omega^*(t-s) .$$

Trying $B(s) = \text{Op}_\varepsilon(b_0(s) + \varepsilon b_1(s) + \dots)$ and using the Moyal product, we have to solve

$$b'(s) + \frac{i}{\varepsilon} ((H_0 + \varepsilon H_1) \star b(s) - b(s) \star (H_0 + \varepsilon H_1)) \equiv 0 .$$

We get the following equations by looking at the ε^j term:

$$\frac{d}{dt} b_j + \{H_0, b_j\}_1 - 2h_1 b_j = c_j ,$$

with $c_0 = 0$ and, for $j \geq 1$, $c_j = \sum_{k=0}^{j-1} P_{k,j}(b_k)$ where the $P_{k,j}$'s are linear differential operators constructed by using the Moyal product with H_0 and H_1 . If $\tilde{B}(s) = \text{Op}_\varepsilon(b_0(s) + \varepsilon b_1(s) + \dots)$, we have

$$\frac{d}{ds} \Omega(t-s)\tilde{B}(s)\Omega^*(t-s) = \Omega(t-s)R(s)\Omega^*(t-s)$$

with $R(s) \in \varepsilon^\infty \Psi_{\text{class}}^{-\infty}$.

It remains to prove that $\Omega(t-s)R(s)\Omega^*(t-s)$ is $O(\varepsilon^\infty)$ as an operator from between the Sobolev spaces H^{-s} into H^s for any s . From Calderon-Vaillancourt Theorem 8.1, it is enough to prove that $L^2 \rightarrow L^2$ norm of $\hat{H}^{s/m} \Omega(t-s)R(s)\Omega^*(t-s)(\hat{H}^*)^{s/m}$ is $O(\varepsilon^\infty)$. This is clear because the powers of \hat{H} (resp. of \hat{H}^*) commute with $\Omega(t-s)$ (resp. $\Omega^*(t-s)$).

10.2 Long time Egorov Theorem

We want to extend Egorov Theorem to times going to infinity with ε . From Lemma 4.2, we get (see [3]):

Lemma 10.1 *Uniformly, for $t \in [0, T_\gamma(\varepsilon)]$, $a_0(t)$ belongs to $S_\delta^{-\infty}$ with $\delta = \frac{1}{2} - \gamma < \frac{1}{2}$ and the constants in the estimates of the derivatives of the symbol are uniform in t . Similarly, $\varepsilon^j a_j$ belongs to $\varepsilon^{j(1-2\delta)} S_\delta^{-\infty}$.*

The main result is:

Theorem 10.2 (Long time Egorov Theorem) *Let us give \hat{H} admissible (Section 10.1) and $a \in C_o^\infty(T^*X)$, for any t with $t \in [0, T_\gamma]$ and with uniform symbol*

estimates w.r. to t , the operator $A(t) := \Omega(t)A\Omega^*(t)$ is a pseudo-differential operator whose total symbol is given as in the classical Egorov Theorem 10.1 by

$$A(t) \equiv \text{Op} \left(\sum_{j=0}^{\infty} \varepsilon^j a_j(t) \right)$$

where $\varepsilon^j a_j(t) \in \varepsilon^{j(1-2\delta)} S_{\delta}^{-\infty}$.

More precisely, for any M and k , there exists N so that if $R_N(t) = A(t) - \text{Op} \left(\sum_{j=0}^N \varepsilon^j a_j(t) \right)$, the C^k norm of the kernel of $R_N(t)$ satisfies, uniformly for $t \in [0, T_{\gamma}]$, the bound

$$\| [R_N(t)] \|_{C^k(X \times X)} = O(\varepsilon^M) .$$

Remark 10.1 *The long time Egorov Theorem was first proved by Bambusi, Graffi and Paul [1] and improved by Bouzouina and Robert [3].*

Our estimate of the remainder term is better than the one given in [3] which is only an L^2 operator norm. We need the ellipticity of \widehat{H}_0 .

Using Lemma 10.1, the proof follows the the proof for a fixed time interval in Section 10.1.

10.3 Integrals with two dynamics

Lemma 10.2 *Let us assume that we have two admissible Hamiltonians $\widehat{H}_{\pm} = \text{Op}(H_{0,\pm} + \varepsilon H_{1,\pm})$ with $k_{\pm} > 0$. Let us define the operator $J := \int_0^{\infty} \Omega_+(t + \tau) \text{Op}(a) \Omega_-^*(t) dt$, with $a \in C_0^{\infty}(T^*X)$. If we assume that, for $z \in \text{Support}(a)$, $H_{0,+}(z) \neq H_{0,-}(z)$ (no mode conversions in the support of a), then*

$$J = \varepsilon \Omega_+(\tau) K + R ,$$

with K a Ψ DO of principal symbol $k_0 = -ia/(H_{0,+} - H_{0,-})$ and $R \in \varepsilon^{\infty} \Psi^{-\infty}$.

In particular, if a vanishes in some neighborhood of B , $[J](A, B) = 0(\varepsilon^{\infty})$.

Proof.–

There exists some Ψ DO K so that $\text{Op}(a) = -i(\widehat{H}_+ K - K \widehat{H}_-)^* + r$ with $r \in \varepsilon^{\infty} \Psi^{-\infty}$. J can be rewritten as

$$J = \varepsilon \int_0^{\infty} \frac{d}{dt} (\Omega_+(t + \tau) K \Omega_-^*(t)) dt + \int_0^{\infty} (\Omega_+(t + \tau) r \Omega_-^*(t)) dt .$$

Integration by part gives the result. □

11 Appendix D: stationary Gaussian fields are ergodic

This section is inspired from [2]. Let us define

$$C_T(\tau, A, B) := \frac{1}{T} \int_0^T u(A, t) \otimes \overline{u(B, t - \tau)} dt ,$$

we have

Theorem 11.1 *Let us assume that $f(\cdot, t) = Kw(\cdot, t)$ where $w(x, t)$ is a (Gaussian) white noise on $X \times \mathbb{R}$ and $K : \mathcal{H} \rightarrow \mathcal{H}$. Then*

1. *If K is Hilbert-Schmidt, $C_T(\tau, A, B) \rightarrow \mathbb{E}(C(\tau, A, B))$, as $T \rightarrow \infty$, almost surely for almost every pair (A, B) .*
2. *If K is smoothing and \hat{H} elliptic, $C_T(\tau, A, B) \rightarrow \mathbb{E}(C(\tau, A, B))$, as $T \rightarrow \infty$, almost surely in the $C^\infty(X \times X)$ topology.*

Proof.–

1. We will assume w.l.o.g. that u is scalar valued and real and use the following classical identity: if $f = (f_1, f_2, f_3, f_4) : \Omega \rightarrow \mathbb{R}^4$ is a Gaussian random variable,

$$\mathbb{E}(f_1 f_2 f_3 f_4) = \mathbb{E}(f_1 f_2) \mathbb{E}(f_3 f_4) + \mathbb{E}(f_1 f_3) \mathbb{E}(f_2 f_4) + \mathbb{E}(f_1 f_4) \mathbb{E}(f_2 f_3). \quad (20)$$

We have

$$\mathbb{E}(\|C_T(\tau)\|_{\mathbb{H}-S}^2) = \frac{1}{T^2} \int_0^T dt \int_0^T dt' \int_{X \times X} dAdB \mathbb{E}(u(A, t)u(B, t-\tau)u(A, t')u(B, t'-\tau)).$$

We apply the identity (20) and get:

$$\mathbb{E}(\|C_T(\tau)\|_{\mathbb{H}-S}^2) = \|\mathbb{E}(C_T(\tau))\|_{\mathbb{H}-S}^2 + \frac{1}{T^2} \int_0^T dt \int_0^T dt' \int_{X \times X} dAdB (II + III),$$

with

$$II = \mathbb{E}(u(A, t)u(A, t')) \mathbb{E}(u(B, t-\tau)u(B, t'-\tau))$$

and

$$III = \mathbb{E}(u(A, t)u(B, t'-\tau)) \mathbb{E}(u(A, t-\tau)u(B, t')).$$

For example, we have

$$\mathbb{E}(u(A, t)u(B, t'-\tau)) = \left[\int_0^\infty ds \int_0^\infty ds' \Omega(s) K K^* \Omega^*(s') \delta(t-s = t'-s'-\tau) \right](A, B)$$

whose L^2 norm as a function of (A, B) is estimated as $O(\exp(-|t-t'+\tau|))$. Using Cauchy-Schwarz inequality and the fact that a product of two Hilbert-Schmidt operators is trace class, we get

$$\left| \frac{1}{T^2} \int_0^T dt \int_0^T dt' \int_{X \times X} dAdB (II + III) \right| = O(1/T).$$

Standards argument involving Tchebichev inequality allow to conclude.

2. The second assertion is proved in a similar way using the trick of Section 10.1.

□

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