

Hitting Probabilities and the Hausdorff Dimension of the Inverse Images of Anisotropic Gaussian Random Fields

Hermine Biermé, Céline Lacaux, Yimin Xiao

▶ To cite this version:

Hermine Biermé, Céline Lacaux, Yimin Xiao. Hitting Probabilities and the Hausdorff Dimension of the Inverse Images of Anisotropic Gaussian Random Fields. Bulletin of the London Mathematical Society, 2009, 41 (2), pp.253-273. 10.1112/blms/bdn122 . hal-00264315

HAL Id: hal-00264315 https://hal.science/hal-00264315v1

Submitted on 15 Mar 2008

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Hitting Probabilities and the Hausdorff Dimension of the Inverse Images of Anisotropic Gaussian Random Fields

Hermine Biermé Céline Lacaux Université Paris Descartes Institut Élie Cartan, Nancy

> Yimin Xiao Michigan State University

> > March 11, 2008

Abstract

Let $X = \{X(t), t \in \mathbb{R}^N\}$ be a Gaussian random field with values in \mathbb{R}^d defined by

$$X(t) = (X_1(t), \dots, X_d(t)),$$

where X_1, \ldots, X_d are independent copies of a centered Gaussian random field X_0 . Under certain general conditions on X_0 , we study the hitting probabilities of X and determine the Hausdorff dimension of the inverse image $X^{-1}(F)$, where $F \subseteq \mathbb{R}^d$ is a non-random Borel set.

The class of Gaussian random fields that satisfy our conditions includes not only fractional Brownian motion, the Brownian sheet, but also such anisotropic fields as fractional Brownian sheets, solutions to stochastic heat equation driven by space-time white noise and the operator-scaling Gaussian random fields with stationary increments constructed in [4].

RUNNING HEAD: Inverse Images of Anisotropic Gaussian Random Fields

2000 AMS CLASSIFICATION NUMBERS: 60G60; 60G15; 60G17; 28A80.

KEY WORDS: Gaussian random fields; anisotropy; fractional Brownian sheet; stochastic heat equation; hitting probability; Hausdorff dimension; inverse image.

1 Introduction

In recent years, several classes of anisotropic Gaussian random fields have arisen naturally in the studies of random fields, stochastic partial differential equations as well as in many applied areas including image processing, hydrology, geostatistics and spatial statistics (see [23] and the references therein for more information).

Typical examples of anisotropic Gaussian random fields are fractional Brownian sheets and the solution to the stochastic heat equation. It has been known that the sample path properties such as fractal dimensions of these anisotropic Gaussian random fields can be very different from those of isotropic ones such as Lévy's fractional Brownian motion; see, for example, [1, 5, 6, 16, 21]. Recently, Xiao [23] studied systematically the analytic and geometric properties of anisotropic Gaussian random fields under certain general conditions, with the aim of characterizing their anisotropic nature by a multiparameter index $H = (H_1, \ldots, H_N) \in (0, 1)^N$. This index is often related to the operator-self-similarity or multi-self-similarity of the Gaussian random field under study.

Let $X = \{X(t), t \in \mathbb{R}^N\}$ be a Gaussian random field with values in \mathbb{R}^d defined on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$ by

$$X(t) = (X_1(t), \dots, X_d(t)).$$
(1.1)

We will call X an (N, d)-Gaussian random field. For any Borel set $F \subseteq \mathbb{R}^d$, its inverse image under X is defined by

$$X^{-1}(F) = \{ t \in \mathbb{R}^N : X(t) \in F \}.$$

In particular, if $F = \{x\}$ $(x \in \mathbb{R}^d)$, then $X^{-1}(\{x\})$ is the level set of X.

In the studies of random fields, it is interesting to consider the following questions:

- (i) Given a non-random Borel set $F \subseteq \mathbb{R}^d$, when is it *non-polar* for X in the sense that $\mathbb{P}\{X^{-1}(F) \neq \emptyset\} > 0$?
- (ii) If $\mathbb{P}\{X^{-1}(F) \neq \emptyset\} > 0$, what is the Hausdorff dimension of the inverse image $X^{-1}(F)$?

In this paper, we will use \dim_{H} to denote Hausdorff dimension, and refer to [8] and [11] for its definition and basic properties.

Question (i) is an important question in potential theory of random fields and complete answer has only been known for a few types of random fields with certain Markovian structures. Of particular significance is the result due to Khoshnevisan and Shi [14], who proved that if $W = \{W(t), t \in \mathbb{R}^N_+\}$ is the (N, d)-Brownian sheet, then for every Borel set $F \subseteq \mathbb{R}^d$,

$$\mathbb{P}\left\{W^{-1}(F)\cap(0,\infty)^N\neq\varnothing\right\}>0\iff\mathcal{C}_{d-2N}(F)>0.$$
(1.2)

Here and in the sequel, \mathcal{C}_{α} denotes the Bessel-Riesz capacity of order α defined by

$$\mathcal{C}_{\alpha}(F) = \left[\inf_{\mu \in \mathcal{P}(F)} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} f_{\alpha}(\|x - y\|) \mu(dx) \mu(dy)\right]^{-1},\tag{1.3}$$

where $\mathcal{P}(F)$ is the family of probability measures carried by F and the function $f_{\alpha}: (0, \infty) \to (0, \infty)$ is defined by

$$f_{\alpha}(r) = \begin{cases} r^{-\alpha} & \text{if } \alpha > 0, \\ \log\left(\frac{e}{r \wedge 1}\right) & \text{if } \alpha = 0, \\ 1 & \text{if } \alpha < 0. \end{cases}$$
(1.4)

Dalang and Nualart [7] extended the methods of Khoshnevisan and Shi [14] and proved similar results for the solution of a system of d nonlinear hyperbolic stochastic partial differential equations with two variables.

For random fields with general dependence structures, it is more difficult to establish results of the type (1.2). Some sufficient conditions and necessary conditions for $\mathbb{P}\left\{X^{-1}(F)\neq \varnothing\right\} > 0$ have been established by Testard [18] and Xiao [22] for fractional Brownian motion, by Dalang, Khoshnevisan and Nualart [5, 6] for the solutions to non-linear stochastic heat equations with additive and multiplicative noises, and by Xiao [23] for a large class of Gaussian random fields satisfying Conditions (C1) and (C2) (see Section 2 for the precise assumptions). Our Theorem 2.1 unifies and, in some cases, strengthens the results in the aforementioned works. The main objective of this paper is to determine the Hausdorff dimension of the inverse image $X^{-1}(F)$ for Gaussian random fields satisfying Conditions (C1) and (C2). The analogous problem for Lévy processes can be solved by applying potential theory of Lévy processes and a subordination technique, see Hawkes [9], Khoshnevisan and Xiao [13] for more details. However, when X is a non-Markovian process or a Gaussian random field and $F \subseteq \mathbb{R}^d$ is a general Borel set, few results on the geometric properties of $X^{-1}(F)$ are available. Testard [18] obtained an upper bound for $\dim_{\mathrm{H}} X^{-1}(F)$ when $X = \{X(t), t \in \mathbb{R}^N\}$ is a fractional Brownian motion of index $\alpha \in (0, 1)$, which is a centered Gaussian random field with values in \mathbb{R}^d and covariance function given by

$$\mathbb{E}[X_i(s)X_j(t)] = \frac{1}{2}\delta_{ij}\Big(\|s\|^{2\alpha} + \|t\|^{2\alpha} - \|s - t\|^{2\alpha}\Big), \quad \forall s, t \in \mathbb{R}^N,$$
(1.5)

where $\|\cdot\|$ denotes the Euclidean norm in \mathbb{R}^N and $\delta_{ij} = 1$ if i = j and 0 otherwise. Monrad and Pitt [17] proved that, if $N > \alpha d$, then almost surely

$$\dim_{\mathrm{H}} X^{-1}(F) = N - \alpha d + \alpha \dim_{\mathrm{H}} F \quad \text{for all Borel sets} \ F \subseteq \mathbb{R}^d.$$
(1.6)

Note that the exceptional null event on which (1.6) does not hold is independent of F, so such a result is called an *uniform Hausdorff dimension result*. The method of Monrad and Pitt [17] relies on rather strong conditions such as $N > \alpha d$ and the strong local nondeterminism of fractional Brownian motion. In Theorems 2.3 and 2.4 below, we determine $\dim_{\rm H} X^{-1}(F)$ for much more general Gaussian random fields. We should also point out that, compared with the isotropic case, the anisotropic nature of X induces far richer fractal structure into $X^{-1}(F)$.

The rest of this paper is organized as follows. In Section 2, we provide the general assumptions on the Gaussian random fields under investigation and state our main results (i.e., Theorems 2.1 and 2.3). Their proofs are given in Section 3. In Section 4 we apply our results to solutions of nonlinear stochastic heat equations considered by Dalang, Khoshnevisan and Nualart [5].

Throughout this paper we will use c to denote an unspecified positive and finite constant which may not be the same in each occurrence. More specific constants in Section i are numbered as $c_{i,1}, c_{i,2}, \ldots$

Acknowledgment The authors are grateful to F. Baudoin and A. Estrade since this work was partially supported by ANR grants ANR-05-BLAN-0017 and ANR-06-BLAN-0289. The research of Y. Xiao is also supported partial by NSF grant DMS-0706728.

This paper was finished while Y. Xiao was visiting the Statistical & Applied Mathematical Sciences Institute (SAMSI). He thanks the staff of SAMSI for their support and the good working conditions. The authors are indebted to Kenneth Falconer for stimulating discussions.

2 Main results

Let $(H_1, \ldots, H_N) \in (0, 1)^N$ be a fixed vector and, for $a, b \in \mathbb{R}^N$ with $a_j < b_j$ $(j = 1, \ldots, N)$, let $I = [a, b] := \prod_{j=1}^N [a_j, b_j] \subseteq \mathbb{R}^N$ denote an interval (or a rectangle). For simplicity, we set $I = [\varepsilon_0, 1]^N$, where $\varepsilon_0 \in (0, 1)$ is a fixed constant. By doing so, there is no loss of generality because, for a given concrete Gaussian random field, the σ -stability of Hausdorff dimension (cf. [8]) will make it clear when one can replace $[\varepsilon_0, 1]^N$ by $[0, 1]^N$ in Theorems 2.3 and 2.4.

Let $X = \{X(t), t \in \mathbb{R}^N\}$ be an (N, d)-Gaussian random field defined by (1.1). Throughout this paper we assume that the coordinate processes X_1, \ldots, X_d are independent copies of a real-valued, centered Gaussian random field $X_0 = \{X_0(t), t \in \mathbb{R}^N\}$. We assume that X_0 satisfies the following general conditions:

(C1) There exist positive and finite constants $c_{2,1}$, $c_{2,2}$ and $c_{2,3}$ such that $\mathbb{E}[X_0(t)^2] \ge c_{2,1}$ for all $t \in I$ and

$$c_{2,2} \sum_{j=1}^{N} |s_j - t_j|^{2H_j} \le \mathbb{E}\Big[\big(X_0(s) - X_0(t) \big)^2 \Big] \le c_{2,3} \sum_{j=1}^{N} |s_j - t_j|^{2H_j}, \quad \forall s, t \in I.$$
 (2.1)

(C2) There exists a constant $c_{2,4} > 0$ such that for all $s, t \in I$,

$$\operatorname{Var}(X_0(t)|X_0(s)) \ge c_{2,4} \sum_{j=1}^N |s_j - t_j|^{2H_j}.$$
(2.2)

Here $\operatorname{Var}(X_0(t)|X_0(s))$ denotes the conditional variance of $X_0(t)$ given $X_0(s)$.

It will be helpful to note that both (2.1) and (2.2) can be expressed in terms of the following metric ρ on \mathbb{R}^N :

$$\rho(s,t) = \sum_{j=1}^{N} |s_j - t_j|^{H_j}, \qquad \forall s, t \in \mathbb{R}^N,$$
(2.3)

since there exist positive constants $c_{\rm 2,5}$ and $c_{\rm 2,6}$ such that

$$c_{2,5} \sum_{j=1}^{N} |s_j - t_j|^{2H_j} \le \rho(s,t)^2 \le c_{2,6} \sum_{j=1}^{N} |s_j - t_j|^{2H_j}, \qquad \forall s,t \in \mathbb{R}^N.$$

The following are some remarks about Conditions (C1) and (C2).

- Condition " $\mathbb{E}[X_0(t)^2] \ge c_{2,1}$ for all $t \in I$ " in (C1) assumes that X is non-degenerate on I. This condition is needed to avoid some trivial cases in studying hitting probability of X and $\dim_{\mathrm{H}}(X^{-1}(F) \cap I)$.
- Under Condition (C1), X has a version which has continuous sample functions on I almost surely. Henceforth we will assume without loss of generality that the Gaussian random field X has continuous sample paths.
- Conditions (C1) and (C2) are closely related. It is easy to see that (C1) implies that $\operatorname{Var}(X_0(t)|X_0(s)) \leq c_{2,3} \sum_{j=1}^N |s_j t_j|^{2H_j}$ for all $s, t \in I$ and, on the other hand, (C2) implies $\sigma^2(s,t) \geq c_{2,4} \sum_{j=1}^N |s_j t_j|^{2H_j}$, where $\sigma^2(s,t)$ denotes $\mathbb{E}[(X_0(s) X_0(t))^2]$. Moreover, if the function $\sigma(0,t)$ satisfies certain smoothness condition, say, it has continuous first order partial derivatives on I, then one can show that (C1) implies (C2) by

using the following fact (which can be easily verified): If (U, V) is a centered Gaussian vector, then

$$\operatorname{Var}(U|V) = \frac{\left(\rho_{U,V}^2 - (\sigma_U - \sigma_V)^2\right)\left((\sigma_U + \sigma_V)^2 - \rho_{U,V}^2\right)}{4\sigma_V^2}, \qquad (2.4)$$

where $\rho_{U,V}^2 = \mathbb{E}[(U-V)^2], \sigma_U^2 = \mathbb{E}(U^2)$ and $\sigma_V^2 = \mathbb{E}(V^2)$.

• Condition (C2) can be referred to as "two-point local nondeterminism" with index $H = (H_1, \ldots, H_N)$. See [23] for more information on properties of local nondeterminism and their applications.

We point out that the class of Gaussian random fields satisfying Conditions (C1) and (C2) is large. It includes not only the well-known fractional Brownian motion and the Brownian sheet, but also such anisotropic random fields as fractional Brownian sheets ([1]), solutions to stochastic heat equation driven by space-time white noise ([5], [6], [16]) and many more. The results in this paper can be further extended so that they can be applied to the operator-scaling Gaussian fields with stationary increments constructed in [4]; see Remark 2.7 below.

Now we state the main results of this paper. We will always assume $F \subseteq \mathbb{R}^d$ to be nonempty and, except maybe in Remark 2.8, non-random. Theorem 2.1 is concerned with the hitting probabilities of X and provides a necessary condition and a sufficient condition for $\mathbb{P}\{X^{-1}(F) \cap I \neq \emptyset\} > 0$. Most part of it has been proved in [23].

Theorem 2.1 Let $X = \{X(t), t \in \mathbb{R}^N\}$ be an (N, d)-Gaussian random field defined by (1.1) and assume that X_0 satisfies Conditions (C1) and (C2). If $F \subseteq \mathbb{R}^d$ is a Borel set, then

$$c_{2,7} \mathcal{C}_{d-Q}(F) \le \mathbb{P}\left\{X^{-1}(F) \cap I \neq \varnothing\right\} \le c_{2,8} \mathcal{H}_{d-Q}(F),$$
(2.5)

where $c_{2,7}$ and $c_{2,8}$ are positive constants depending on I, F and H only. In the above, $Q := \sum_{j=1}^{N} \frac{1}{H_j}$ and $\mathcal{H}_q(F)$ is defined as the q-dimensional Hausdorff measure of F when q > 0, and $\mathcal{H}_q(F) = 1$ whenever $q \le 0$.

Remark 2.2

- Note that if $0 \in F$, then the second inequality in (2.5) may not hold without the condition " $\mathbb{E}[X_0(t)^2] \ge c_{2,1}$ for all $t \in I$ " in (C1).
- Let $F \subseteq \mathbb{R}^d$ be a Borel set and $Q = \sum_{j=1}^N \frac{1}{H_j}$. Theorem 2.1 implies the following statements:
 - (i) If $\dim_{\mathrm{H}} F < d Q$, then $X^{-1}(F) \cap I = \emptyset$ a.s.
 - (ii) If $\dim_{\mathrm{H}} F > d Q$, then $\mathbb{P}\{X^{-1}(F) \cap I \neq \emptyset\} > 0$.
- A natural conjecture is that, at least for a large subclass of Gaussian random fields satisfying the conditions of Theorem 2.1, $\mathbb{P}\{X^{-1}(F) \cap I \neq \emptyset\} > 0 \iff \mathcal{C}_{d-Q}(F) > 0$.

Next we consider the Hausdorff dimension of the inverse image $X^{-1}(F) \cap I$. In the special case of fractional Brownian motion, this problem was considered by Testard [18] and Monrad and Pitt [17]. The following two theorems determine $\dim_{\mathrm{H}} (X^{-1}(F) \cap I)$ in terms of $\dim_{\mathrm{H}} F$ and the parameters N, d and $H \in (0, 1)^N$ of X. Compared with the results for fractional Brownian motion, we see that the fractal structure of the inverse images of an anisotropic random field is much richer.

For convenience we will always assume

$$0 < H_1 \le H_2 \le \dots \le H_N < 1.$$
 (2.6)

Theorem 2.3 Let $X = \{X(t), t \in \mathbb{R}^N\}$ be an (N, d)-Gaussian random field defined by (1.1) with X_0 satisfying Conditions (C1) and (C2). Let $F \subseteq \mathbb{R}^d$ be a Borel set such that $\dim_{\mathrm{H}} F \geq d - Q$. Then the following statements hold:

(i) Almost surely

$$\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) \le \min_{1 \le k \le N} \left\{ \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k \left(d - \dim_{\mathrm{H}} F \right) \right\}.$$
(2.7)

In particular, if $\dim_{\mathrm{H}} F = d - Q$, then $\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) = 0$ a.s.

(ii) If $\dim_{\mathrm{H}} F > d - Q$, then for every $\varepsilon > 0$,

$$\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) \ge \min_{1 \le k \le N} \left\{ \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k \left(d - \dim_{\mathrm{H}} F \right) \right\} - \varepsilon$$
(2.8)

on an event of positive probability (which may depend on ε). In the special case when $\dim_{\mathrm{H}} F = 0$ and d < Q, we have

$$\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) = \min_{1 \le k \le N} \left\{ \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k d \right\}$$
(2.9)

on an event of positive probability.

As in [9] and [13], we can combine (2.7) and (2.8) in Theorem 2.3 and write

$$\left\|\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) \right\|_{L^{\infty}(\mathbb{P})} = \min_{1 \le k \le N} \left\{ \sum_{j=1}^{k} \frac{H_{k}}{H_{j}} + N - k - H_{k} \left(d - \dim_{\mathrm{H}} F \right) \right\}.$$
(2.10)

Here and in the sequel, for any function $Y: \Omega \to \mathbb{R}_+$, $||Y||_{L^{\infty}(\mathbb{P})}$ is defined as

 $||Y||_{L^{\infty}(\mathbb{P})} = \sup \left\{ \theta : Y \ge \theta \text{ on an event } E \text{ with } \mathbb{P}(E) > 0 \right\}.$

If one has more information on the geometry of F, then the conclusion in Part (ii) of Theorem 2.3 can be strengthened so that the positive probability is independent of ε . In the following we provide a sufficient (but not necessary) condition for this to be true. Because of (2.7) and (2.9), it suffices to consider Borel sets $F \subseteq \mathbb{R}^d$ satisfying $\dim_{\mathrm{H}} F > \max\{d - Q, 0\}$. **Theorem 2.4** Let $X = \{X(t), t \in \mathbb{R}^N\}$ be as in Theorem 2.3 and let $F \subseteq \mathbb{R}^d$ be a bounded Borel set such that $\dim_{\mathrm{H}} F > \max\{d - Q, 0\}$. If F satisfies the Condition (S): There is a finite constant $c_{2,9} \ge 1$ such that, for every $\gamma \in (0, \dim_{\mathrm{H}} F)$, there exists a probability measure $\mu_{0,\gamma}$ (which may depend on γ) with compact support in F such that

$$\mu_{0,\gamma}\big(B(x,r)\big) \le c_{2,9} r^{\gamma} \qquad for \ all \ x \in \mathbb{R}^d \ and \ r > 0.$$

$$(2.11)$$

Then with positive probability,

$$\dim_{H} \left(X^{-1}(F) \cap I \right) = \min_{1 \le k \le N} \left\{ \sum_{j=1}^{k} \frac{H_{k}}{H_{j}} + N - k - H_{k}(d - \dim_{H}F) \right\}$$

$$= \sum_{j=1}^{k} \frac{H_{k}}{H_{j}} + N - k - H_{k}(d - \dim_{H}F), \qquad if \quad \sum_{j=1}^{k-1} \frac{1}{H_{j}} \le d - \dim_{H}F < \sum_{j=1}^{k} \frac{1}{H_{j}}.$$

$$(2.12)$$

In the above we use the convention $\sum_{j=1}^{6} \frac{1}{H_j} = 0.$

We end this section with the following remarks about Theorems 2.3 and 2.4.

Remark 2.5 The key feature of Condition (S) is that the constant $c_{2,9}$ is independent of $\gamma < \dim_{\mathrm{H}} F$, even though the probability measure $\mu_{0,\gamma}$ may depend on γ . This assumption is stronger than the conclusion of Frostman's lemma (where both $c_{2,9}$ and $\mu_{0,\gamma}$ may depend on γ), and is verified by a large class of Borel sets including all self-similar sets satisfying the open set condition ([8]) and, more generally, the *d*-sets ([10]). By examining the proof of Frostman's lemma in [11, p. 130–131], one sees that Condition (S) is satisfied provided F has the following property: There is a constant $c_{2,10} > 0$ such that, for every $\gamma \in (0, \dim_{\mathrm{H}} F)$, one has

$$\inf\left\{\sum_{n=1}^{\infty} (\operatorname{diam} B_n)^{\gamma}: F \subseteq \bigcup_{n=1}^{\infty} B_n\right\} \ge c_{2,10},$$
(2.13)

where the infimum is taken over all coverings of F by balls $\{B_n, n \ge 1\}$. However, we have not been able to fully characterise the family of Borel sets satisfying Condition (S).

Remark 2.6 In the special case of $F = \{x\}, X^{-1}(F) \cap I$ becomes the level set of X on I and the above Hausdorff dimension formula (2.9) was proved by Ayache and Xiao [1] for fractional Brownian sheets, by Wu and Xiao [20] for the solutions of stochastic heat equation driven by a space-time white noise, and by Xiao [23] for Gaussian random fields satisfying Conditions (C1) and (C2). Theorems 2.3 and 2.4 extend their results and, as we will see in the next section, their proofs rely on some new ingredients.

Remark 2.7 From the proof of Theorem 2.3 given in Section 3 one can see that its conclusions still hold if (2.1) and (2.2) in Conditions (C1) and (C2) are replaced respectively by the following weaker conditions: For all $\varepsilon > 0$, there exist constants $c'_{2,i}$ (i = 2, 3, 4) which may depend on ε such that

$$c_{2,2}' \sum_{j=1}^{N} |s_j - t_j|^{2H_j + \varepsilon} \le \mathbb{E} \Big[\big(X_0(s) - X_0(t) \big)^2 \Big] \le c_{2,3}' \sum_{j=1}^{N} |s_j - t_j|^{2H_j - \varepsilon}, \quad \forall s, t \in I.$$
(2.14)

and

$$\operatorname{Var}(X_0(t)|X_0(s)) \ge c'_{2,4} \sum_{j=1}^N |s_j - t_j|^{2H_j + \varepsilon}, \quad \forall s, t \in I.$$
(2.15)

It can be proven that the operator-scaling Gaussian fields constructed in [4] satisfy conditions (2.14) and (2.15), while Conditions (C1) and (C2) are only satisfied by those random fields with diagonal scaling matrices. We refer to [2] and [3] for further information.

Remark 2.8 Note that the event on which (2.12) holds depends on F. In light of the result (1.6) of Monrad and Pitt [17], we may ask the following natural question: When $\sum_{j=1}^{N} \frac{1}{H_j} > d$, is it possible to have a single event $\Omega_1 \subseteq \Omega$ of positive probability such that on Ω_1 (2.12) holds for all Borel sets $F \subseteq \mathbb{R}^d$? This is simply referred to as a uniform Hausdorff dimension problem for the inverse images of X. An affirmative answer will make the formula (2.12) applicable to random Borel sets F as well.

If, in addition to Conditions (C1) and (C2), we also assume that Condition (C3) or (C3') in [23] holds and $H_1 = H_2 = \cdots = H_N$, then one can modify the proofs in [17] to prove that the answer to the above question is affirmative when $N > H_1 d$. In general, it can be proved that the upper bound (2.7) holds almost surely for all Borel sets $F \subseteq \mathbb{R}^d$. However it is possible that the lower bound may not hold uniformly due to the anisotropy of X.

Remark 2.9 Another useful dimension for studying fractals is packing dimension, denoted by $\dim_{\mathbf{P}}$, see [8] for its definition and properties. We mention that the problem of finding the packing dimension of $X^{-1}(F)$ remains widely open. Even though one can apply Lemma 3.1 and a covering argument to derive that under the conditions of Theorem 2.3,

$$\dim_{\mathbf{P}} \left(X^{-1}(F) \cap I \right) \le \min_{1 \le k \le N} \left\{ \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k (d - \dim_{\mathbf{P}} F) \right\}, \quad \text{a.s.}$$
(2.16)

It is possible that, for a general Borel set $F \subseteq \mathbb{R}^d$, the inequality in (2.16) may be strict. New ideas may be needed for solving this problem.

3 Proofs

In this section we prove the theorems stated in Section 2. The proofs of Theorem 2.1 and Theorem 2.3 are divided into proving the upper and lower bounds for $\mathbb{P}\{X^{-1}(F)\cap I\neq\emptyset\}$ and $\dim_{\mathrm{H}}(X^{-1}(F)\cap I)$ separately. The upper bounds will be proved by using a covering argument, and the lower bounds will be proved by constructing a random measure on $X^{-1}(F)\cap I$. These methods are similar to those in [11] and [18] for studying the level sets and hitting probabilities of fractional Brownian motion. However, in order to deal with anisotropic Gaussian random fields in general, we have to rely also on some new ingredients such as Lemma 3.1 and Lemma 3.2 below.

We will apply the following lemma on the hitting probability of X to prove the upper bounds in (2.5) and (2.7). A slightly stronger version of this lemma was provided in [23] with a longer proof. For the sake of completeness, we will provide a simpler proof by modifying the argument in the proof of Proposition 4.4 in Dalang, Khoshnevisan and Nualart [5]. **Lemma 3.1** Let $X = \{X(t), t \in \mathbb{R}^N\}$ be an (N, d)-Gaussian random field defined by (1.1) with X_0 satisfying Conditions (C1) and (C2). For any constant M > 0, there exist positive constants $c_{3,1}$ and δ_0 such that for all $r \in (0, \delta_0)$, $t \in I$ and all $x \in [-M, M]^d$,

$$\mathbb{P}\left\{\inf_{s\in B_{\rho}(t,r)\cap I}\left\|X(s)-x\right\|\leq r\right\}\leq c_{3,1}\,r^{d}.$$
(3.1)

In the above $B_{\rho}(t,r) = \{s \in \mathbb{R}^N : \rho(s,t) \leq r\}$ denotes the closed ball of radius r in the metric ρ in \mathbb{R}^N defined by (2.3).

Proof Since the coordinate processes of X are independent copies of X_0 , it is sufficient to prove (3.1) for X_0 . Note that for any $s \in I$, we have

$$\mathbb{E}(X_0(s)|X_0(t)) = \frac{\mathbb{E}(X_0(s)X_0(t))}{\mathbb{E}(X_0(t)^2)} X_0(t) := c(s,t)X_0(t).$$
(3.2)

Then the Gaussian random variables $X_0(s) - c(s,t)X_0(t)$ ($s \in I$) and $X_0(t)$ are independent. Denote

$$Z(t,r) = \sup_{s \in B_{\rho}(t,r) \cap I} |X_0(s) - c(s,t)X_0(t)|.$$

Then Z(t,r) is independent of $X_0(t)$. It follows from triangle's inequality that

$$\mathbb{P}\left\{\inf_{s\in B_{\rho}(t,r)\cap I} \left|X_{0}(s)-x\right| \leq r\right\} \\
\leq \mathbb{P}\left\{\inf_{s\in B_{\rho}(t,r)\cap I} \left|c(s,t)\left(X_{0}(t)-x\right)\right| \leq r+Z(t,r)+\sup_{s\in B_{\rho}(t,r)\cap I} \left|\left(1-c(s,t)\right)x\right|\right\}.$$
(3.3)

The Cauchy-Schwarz inequality and Condition (C1) imply that for all $s, t \in I$,

$$\left|1 - c(s,t)\right| = \frac{\left|\mathbb{E}\left[X_0(t)\left(X_0(t) - X_0(s)\right)\right]\right|}{\mathbb{E}\left(X_0(t)^2\right)} \le \left(\frac{c_{2,3}}{c_{2,1}}\sum_{j=1}^N |s_j - t_j|^{2H_j}\right)^{1/2}.$$
 (3.4)

Hence there exists a constant $\delta_0 > 0$ such that for all $r \in (0, \delta_0)$ and $s \in B_\rho(t, r) \cap I$, we have $1/2 \leq c(s, t) \leq 1$. Moreover, for all $r \in (0, \delta_0)$, all $s \in B_\rho(t, r) \cap I$ and all $x \in [-M, M]$,

$$\left| \left(1 - c(s,t) \right) x \right| \le \frac{M r}{2}.$$

Combining the above with (3.3) it follows that for all $x \in [-M, M]$ and $r \in (0, \delta_0)$,

$$\mathbb{P}\bigg\{\inf_{s\in B_{\rho}(t,r)\cap I} |X_0(s)-x| \le r\bigg\} \le \mathbb{P}\bigg\{x - (M+2)r - 2Z(t,r) \le X_0(t) \le x + (M+2)r + 2Z(t,r)\bigg\}.$$

Then using the independence of $X_0(t)$ and Z(t, r) and Condition (C1) we obtain that for all $x \in [-M, M]$ and $r \in (0, \delta_0)$,

$$\mathbb{P}\bigg\{\inf_{s\in B_{\rho}(t,r)\cap I} |X_0(s) - x| \le r\bigg\} \le c_{3,2} \,\mathbb{E}\big(r + Z(t,r)\big) = c_{3,2} \,r + c_{3,2} \,\mathbb{E}\big(Z(t,r)\big), \tag{3.5}$$

where $c_{3,2}$ is a finite positive constant which depends on M.

The last term in (3.5) can be estimated by applying Dudley's entropy theorem (see, e.g., [11], p. 219) to the Gaussian random field $Y(s) = X_0(s) - c(s,t)X_0(t)$ ($s \in B_\rho(t,r) \cap I$). To this end, note that Y(t) = 0 and the canonical metric

$$d(s, s') := \left\{ \mathbb{E} (Y(s) - Y(s'))^2 \right\}^{1/2} \le c \,\rho(s, s')$$

for all $s, s' \in B_{\rho}(t,r) \cap I$. Denote by $D := \sup_{s,s' \in B_{\rho}(t,r) \cap I} d(s,s')$ and $N_d(B_{\rho}(t,r),\varepsilon)$ the *d*-diameter and the metric entropy number of $B_{\rho}(t,r)$, respectively. Then $D \leq cr$ and

$$N_d(B_{\rho}(t,r),\varepsilon) \le c \left(\frac{r}{\varepsilon}\right)^Q$$
.

It follows from Dudley's theorem that

$$\mathbb{E}(Z(t,r)) \le c \int_0^D \sqrt{\log N_d(B_\rho(t,r),\varepsilon)} \, d\varepsilon \le c_{3,3} \, r.$$
(3.6)

This and (3.5) together imply (3.1) for d = 1. The proof of Lemma 3.1 is finished.

For proving the lower bounds in Theorems 2.1, 2.3 and Theorem 2.4 we will make use of the following lemma which extends a result of Testard [18] for fractional Brownian motion.

Lemma 3.2 Let $X = \{X(t), t \in \mathbb{R}^N\}$ be an (N, d)-Gaussian random field defined by (1.1) with X_0 satisfying Conditions (C1) and (C2). Then there exists a positive and finite constant $c_{3,4}$ such that for all $\varepsilon \in (0, 1)$, $s, t \in I$ and $x, y \in \mathbb{R}^d$, we have

$$\int_{\mathbb{R}^{2d}} e^{-i(\langle \xi, x \rangle + \langle \eta, y \rangle)} \exp\left(-\frac{1}{2}(\xi, \eta) \left(\varepsilon I_{2d} + \operatorname{Cov}(X(s), X(t))\right)(\xi, \eta)'\right) d\xi d\eta \\
\leq \frac{c_{3,4}}{\max\{\rho^d(s, t), \|x - y\|^d\}}.$$
(3.7)

In the above, I_{2d} denotes the identity matrix of order 2d, Cov(X(s), X(t)) denotes the covariance matrix of the random vector (X(s), X(t)), and $(\xi, \eta)'$ is the transpose of the row vector (ξ, η) .

Proof Since the coordinate processes X_1, \ldots, X_d of X are independent copies of X_0 , we first consider the case when d = 1.

For any $\varepsilon \in (0, 1)$, let

$$\Phi_{\varepsilon}^{s,t}(\xi_{1},\eta_{1}) = \varepsilon(\xi_{1}^{2}+\eta_{1}^{2}) + \mathbb{E}(\xi_{1}X_{1}(s)+\eta_{1}X_{1}(t))^{2}$$

= $\left(\varepsilon + \mathbb{E}(X_{1}(s)^{2})\right)\xi_{1}^{2} + 2\xi_{1}\eta_{1}\mathbb{E}(X_{1}(s)X_{1}(t)) + \left(\varepsilon + \mathbb{E}(X_{1}(t)^{2})\right)\eta_{1}^{2}.$ (3.8)

Denote by $\Gamma_{\varepsilon}(s,t)$ the symmetric matrix corresponding to the quadratic form (3.8). Then $\Gamma_{\varepsilon}(s,t) = \varepsilon I_2 + \text{Cov}(X_1(s), X_1(t))$ and its inverse is given by

$$\Gamma_{\varepsilon}^{-1}(s,t) = \frac{1}{\det(\Gamma_{\varepsilon}(s,t))} \begin{pmatrix} \varepsilon + \mathbb{E}(X_1(t)^2) & -\mathbb{E}(X_1(s)X_1(t)) \\ -\mathbb{E}(X_1(s)X_1(t)) & \varepsilon + \mathbb{E}(X_1(s)^2) \end{pmatrix},$$
(3.9)

where $\det(\Gamma_{\varepsilon}(s,t))$ denotes the determinant of $\Gamma_{\varepsilon}(s,t)$. Since the symmetric matrix $\Gamma_{\varepsilon}(s,t)$ is positive definite for all $\varepsilon > 0$, we have

$$\int_{\mathbb{R}^2} e^{-i(\xi_1 x_1 + \eta_1 y_1)} \exp\left(-\frac{1}{2} \Phi_{\varepsilon}^{s,t}(\xi_1, \eta_1)\right) d\xi_1 d\eta_1$$

$$= \frac{2\pi}{\sqrt{\det\left(\Gamma_{\varepsilon}(s,t)\right)}} \exp\left(-\frac{1}{2}(x_1, y_1) \Gamma_{\varepsilon}^{-1}(s,t) (x_1, y_1)'\right).$$
(3.10)

Let us remark that

$$(x_{1}, y_{1}) \Gamma_{\varepsilon}^{-1}(s, t) (x_{1}, y_{1})' = \frac{1}{\det(\Gamma_{\varepsilon}(s, t))} \left(\varepsilon(x_{1}^{2} + y_{1}^{2}) + \mathbb{E}\left[\left(x_{1}X_{1}(t) - y_{1}X_{1}(s) \right)^{2} \right] \right) \\ \geq \frac{1}{\det(\Gamma_{\varepsilon}(s, t))} \mathbb{E}\left[\left(x_{1}X_{1}(t) - y_{1}X_{1}(s) \right)^{2} \right].$$
(3.11)

We claim that for all $s, t \in I$ and all $x_1, y_1 \in \mathbb{R}$,

$$\mathbb{E}\Big[\big(x_1X_1(t) - y_1X_1(s)\big)^2\Big] \ge c_{3,5} (x_1 - y_1)^2, \tag{3.12}$$

where $c_{3,5} > 0$ is a constant depending only on $c_{2,1}, \ldots, c_{2,4}$ in Conditions (C1) and (C2).

In order to verify (3.12), we note that it is equivalent to prove the following: There exists a constant c > 0 such that, for all $s, t \in I$, the quadratic form

$$\mathbb{E}\left[\left(x_1X_1(s) - y_1X_1(t)\right)^2\right] - c\left(x_1 - y_1\right)^2$$
(3.13)

is positive semi-definite. Hence it is sufficient to show the existence of a constant c > 0 such that

- (i) $\mathbb{E}(X_1(s)^2) \ge c$ for all $s \in I$.
- (ii) for all $s, t \in I$ we have

$$\left(\mathbb{E}\left(X_1(s)^2\right) - c\right) \left(\mathbb{E}\left(X_1(t)^2\right) - c\right) - \left(\mathbb{E}\left(X_1(s)X_1(t)\right) - c\right)^2 \ge 0.$$
(3.14)

Since Condition (C1) implies (i) holds for any $c \in (0, c_{2,1}]$, it remains to verify (ii) for c > 0 small enough.

It can be seen that (3.14) holds for some c > 0 if and only if

$$c \leq \frac{\mathbb{E}(X_1(s)^2) \mathbb{E}(X_1(t)^2) - (\mathbb{E}(X_1(s)X_1(t)))^2}{\mathbb{E}[(X_1(s) - X_1(t))^2]}$$
(3.15)

for all $s, t \in I$ with $s \neq t$. Observe that the numerator in (3.15) is detCov $(X_1(s), X_1(t))$. It follows from (C1), (C2) and the identity that for any Gaussian vector (Z_1, Z_2) ,

$$\det \operatorname{Cov}(Z_1, Z_2) = \operatorname{Var}(Z_1) \operatorname{Var}(Z_2 | Z_1)$$
(3.16)

that

$$\det \operatorname{Cov}(X_1(s), X_1(t)) \ge c_{2,1} c_{2,4} \sum_{j=1}^N |s_j - t_j|^{2H_j}.$$
(3.17)

This and (2.1) imply that the right-hand side of (3.15) is bounded from below by $c_{2,1}c_{2,4}/c_{2,3}$. By taking $0 < c \le \min\{c_{2,1}, c_{2,1}c_{2,4}/c_{2,3}\}$, we have verified (3.12). Combining (3.10), (3.11) and (3.12) together, we obtain

$$\int_{\mathbb{R}^2} e^{-i(\xi_1 x_1 + \eta_1 y_1)} \exp\left(-\frac{1}{2} \Phi_{\varepsilon}^{s,t}(\xi_1, \eta_1)\right) d\xi_1 d\eta_1$$

$$\leq \frac{2\pi}{\sqrt{\det\left(\Gamma_{\varepsilon}(s,t)\right)}} \exp\left(-\frac{c_{3,5}}{2} \frac{(x_1 - y_1)^2}{\det\left(\Gamma_{\varepsilon}(s,t)\right)}\right). \tag{3.18}$$

Going back to \mathbb{R}^d , we derive from (3.18) that for all $\varepsilon > 0$,

$$\int_{\mathbb{R}^{2d}} e^{-i(\langle \xi, x \rangle + \langle \eta, y \rangle)} \exp\left(-\frac{1}{2}(\xi, \eta) \left(\varepsilon I_{2d} + \operatorname{Cov}(X(s), X(t))\right)(\xi, \eta)'\right) d\xi d\eta \\
\leq \frac{(2\pi)^d}{\left[\det\left(\Gamma_{\varepsilon}(s, t)\right)\right]^{d/2}} \exp\left(-\frac{c_{3,5}}{2} \frac{\|x - y\|^2}{\det\left(\Gamma_{\varepsilon}(s, t)\right)}\right).$$
(3.19)

Note that, if $\det(\Gamma_{\varepsilon}(s,t)) \geq ||x-y||^2$, then

$$\frac{(2\pi)^d}{\left[\det\left(\Gamma_{\varepsilon}(s,t)\right)\right]^{d/2}} \exp\left(-\frac{c_{3,5}}{2} \frac{\|x-y\|^2}{\det\left(\Gamma_{\varepsilon}(s,t)\right)}\right) \le \frac{(2\pi)^d}{\left[\det\left(\Gamma_{\varepsilon}(s,t)\right)\right]^{d/2}}.$$
(3.20)

On the other hand, if det $(\Gamma_{\varepsilon}(s,t)) < ||x-y||^2$, then we use the elementary inequality $x^{d/2}e^{-c\,x} \leq c_{\scriptscriptstyle 3.6}~(\forall\,x>0)$ to obtain

$$\frac{(2\pi)^d}{\left[\det\left(\Gamma_{\varepsilon}(s,t)\right)\right]^{d/2}} \exp\left(-\frac{c_{3,5}}{2} \frac{\|x-y\|^2}{\det\left(\Gamma_{\varepsilon}(s,t)\right)}\right) \le \frac{c_{3,7}}{\|x-y\|^d}.$$
(3.21)

By (3.20) and (3.21) we obtain

$$\frac{(2\pi)^d}{\left[\det\left(\Gamma_{\varepsilon}(s,t)\right)\right]^{d/2}} \exp\left(-\frac{c_{3,5}}{2} \frac{\|x-y\|^2}{\det\left(\Gamma_{\varepsilon}(s,t)\right)}\right) \le \frac{c_{3,8}}{\max\left\{\left[\det\left(\Gamma_{\varepsilon}(s,t)\right)\right]^{d/2}, \|x-y\|^d\right\}}.$$
 (3.22)

Finally, we note that for all $s, t \in I$,

$$\det(\Gamma_{\varepsilon}(s,t)) \ge \det\operatorname{Cov}(X_1(s), X_1(t)) \ge c_{2,1} c_{2,4} \rho(s,t)^2.$$
(3.23)

Therefore (3.7) follows from (3.19), (3.22) and (3.23).

We will also make use of the following elementary lemma from [1, Lemma 10].

Lemma 3.3 Let α , β and η be positive constants. For A > 0 and B > 0, let

$$J := J(A, B) = \int_0^1 \frac{dt}{(A + t^{\alpha})^{\eta} (B + t)^{\beta}}.$$
 (3.24)

Then there exist finite constants $c_{3,9}$ and $c_{3,10}$, depending on α , β , η only, such that the following hold for all real numbers A, B > 0 satisfying $A^{1/\alpha} \leq c_{3,9} B$:

(i) If $\alpha \eta > 1$, then

$$J \le c_{3,10} \, \frac{1}{A^{\eta - \alpha^{-1}} B^{\beta}}; \tag{3.25}$$

(ii) If $0 < \alpha \eta < 1$ and $\alpha \eta + \beta \neq 1$, then

$$J \le c_{3,10} \left(\frac{1}{B^{\alpha \eta + \beta - 1}} + 1 \right). \tag{3.26}$$

Now we are ready to prove Theorems 2.1, 2.3 and 2.4.

Proof of Theorem 2.1 We distinguish three cases: (i) d < Q, (ii) d = Q and (iii) d > Q. When d < Q, the second inequality in (2.5) holds automatically with $c_{2,8} = 1$. Moreover, Theorem 7.1 in [23] implies that X hits points, hence $\mathbb{P}\{X^{-1}(F) \cap I \neq \emptyset\} > 0$ since $F \neq \emptyset$.

When d > Q, (2.5) is proved in Theorem 7.6 in [23]. Hence it only remains to consider the case (ii). Again, the second inequality in (2.5) holds automatically with $c_{2.8} = 1$.

The proof of the lower bound in (2.5) is similar to that of Theorem 7.6 in [23]. Without loss of generality, we assume $C_0(F) > 0$ otherwise there is nothing to prove. By Choquet's capacity theorem (cf. [12]), we may and will assume F is compact and let M > 0 be a constant such that $F \subseteq [-M, M]^d$.

By definition (1.3), there is a Borel probability measure ν_0 on F such that

$$\mathcal{E}_{0}(\nu_{0}) := \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \log\left(\frac{e}{\|x-y\| \wedge 1}\right) \nu_{0}(dx)\nu_{0}(dy) \le \frac{2}{\mathcal{C}_{0}(F)}.$$
(3.27)

For all integers $n \ge 1$, we consider a family of random measures ν_n on I defined by

$$\int_{I} g(t) \nu_{n}(dt) = \int_{I} \int_{\mathbb{R}^{d}} (2\pi n)^{d/2} \exp\left(-\frac{n \left\|X(t) - x\right\|^{2}}{2}\right) g(t) \nu_{0}(dx) dt$$

$$= \int_{I} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \exp\left(-\frac{\|\xi\|^{2}}{2n} + i\langle\xi, X(t) - x\rangle\right) g(t) d\xi \nu_{0}(dx) dt,$$
(3.28)

where g is an arbitrary measurable, nonnegative function on I.

Denote the total mass of ν_n by $\|\nu_n\| := \nu_n(I)$. We claim that the following two inequalities hold:

$$\mathbb{E}(\|\nu_n\|) \ge c_{3,11} \quad \text{and} \quad \mathbb{E}(\|\nu_n\|^2) \le c_{3,12}\mathcal{E}_0(\nu_0), \tag{3.29}$$

where the constants $c_{3,11}$ and $c_{3,12}$ are independent of ν_0 and n.

Some simple calculation shows

$$\mathbb{E}(\|\nu_{n}\|) = \int_{I} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \exp\left(-\frac{1}{2}\left(\frac{1}{n} + \sigma^{2}(t)\right)\|\xi\|^{2} - i\langle\xi, x\rangle\right) d\xi \,\nu_{0}(dx) \,dt \\
\geq \int_{I} \int_{\mathbb{R}^{d}} \frac{(2\pi)^{d/2}}{(1 + \sigma^{2}(t))^{d/2}} \exp\left(-\frac{\|x\|^{2}}{2\sigma^{2}(t)}\right) \nu_{0}(dx) \,dt \\
\geq \int_{I} \frac{(2\pi)^{d/2}}{(1 + \sigma^{2}(t))^{d/2}} \exp\left(-\frac{dM^{2}}{2\sigma^{2}(t)}\right) dt := c_{3,11},$$
(3.30)

where $\sigma^2(t) = \mathbb{E}(X_0(t)^2)$ and the last inequality follows from the boundedness of F. Since $t \mapsto \sigma^2(t)$ is a positive continuous function on I, $c_{3,11} > 0$. This gives the first inequality in (3.29).

Denote by I_{2d} the identity matrix of order 2d and let $\Gamma_n(s,t) = n^{-1}I_{2d} + \text{Cov}(X(s), X(t))$. Since $\Gamma_n(s,t)$ is positive definite, we have

$$\mathbb{E}(\|\nu_{n}\|^{2}) = \int_{I} \int_{I} \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} e^{-i(\langle\xi, x\rangle + \langle\eta, y\rangle)} \\
\times \exp\left(-\frac{1}{2}(\xi, \eta) \Gamma_{n}(s, t) (\xi, \eta)'\right) d\xi d\eta \nu_{0}(dx) \nu_{0}(dy) ds dt \\
= \int_{I} \int_{I} \int_{\mathbb{R}^{2d}} \frac{(2\pi)^{d}}{\sqrt{\det(\Gamma_{n}(s, t))}} \exp\left(-\frac{1}{2}(x, y) \Gamma_{n}^{-1}(s, t) (x, y)'\right) \nu_{0}(dx) \nu_{0}(dy) ds dt \\
\leq c_{3,13} \int_{I} \int_{I} \int_{\mathbb{R}^{2d}} \frac{1}{\max\{\rho^{d}(s, t), \|x - y\|^{d}\}} \nu_{0}(dx) \nu_{0}(dy) ds dt,$$
(3.31)

where the last inequality follows from Lemma 3.2 and the constant $c_{3,13}$ is independent of ν_0 and n.

We claim that for all $x, y \in \mathbb{R}^d$,

$$\int_{I} \int_{I} \frac{dsdt}{\max\{\rho^{d}(s,t), \|x-y\|^{d}\}} \le c_{3,14} \log\left(\frac{e}{\|x-y\|\wedge 1}\right),$$
(3.32)

where $c_{3,14} > 0$ is a constant. To verify this, we break the integral in (3.32) over the regions $\{(s,t) \in I \times I : \rho(s,t) \leq ||x-y||\}$ and $\{(s,t) \in I \times I : \rho(s,t) > ||x-y||\}$ and denote them by \mathcal{I}_1 and \mathcal{I}_2 , respectively. Since d = Q, we derive that

$$\mathcal{I}_{1} = \int_{I} \int_{\{t \in I: \rho(s,t) \le \|x-y\|\}} \frac{dsdt}{\|x-y\|^{d}} \le \frac{c_{3,15}}{\|x-y\|^{d}} \|x-y\|^{Q} = c_{3,15}.$$
(3.33)

In the above, we have used the fact that, for every $s \in I$, the set $\{t \in I : \rho(s,t) \leq ||x-y||\}$ is contained in an rectangle of side-lengths $||x-y||^{1/H_j}$ (j = 1, ..., N).

On the other hand,

$$\mathcal{I}_2 = \int_I \int_{\{t \in I: \rho(s,t) > ||x-y||\}} \frac{dsdt}{\rho^d(s,t)} \le \int_{\{t \in [-1,1]^N: \rho(0,t) > ||x-y||\}} \frac{dt}{\rho^d(0,t)}$$

since $\rho(s,t) = \rho(0,t-s)$. Let us consider the diagonal matrix $E = \text{diag}(1/H_1, \ldots, 1/H_N)$. Then, $t \mapsto \rho(0,t)$ is E-homogeneous function in the sense of Definition 2.6 of [4], that is

$$\forall r > 0, \quad \rho\left(0, r^{E}t\right) = r\rho\left(0, t\right)$$

with $r^E = \exp(\log(r)E)$. By using the formula of integration in the polar coordinates with respect to E (see Proposition 2.3 in [15]) and Q = d, we have

$$\mathcal{I}_2 \leq \int_0^{+\infty} \int_{\mathcal{S}_E} \mathbf{1}_{\{r^E \theta \in [-1,1], r\rho(0,\theta) > \|x-y\|\}} \frac{\sigma_E(d\theta) dr}{r \rho^d(0,\theta)},$$

where S_E is a compact set which does not contain 0 and $\sigma_E(d\theta)$ is a finite positive Radon measure on S_E . Since

$$0 < \min_{\theta \in \mathcal{S}_E} \rho\left(0, \theta\right) \le \max_{\theta \in \mathcal{S}_E} \rho\left(0, \theta\right) < \infty,$$

there exist some finite positive constants $m, M, c, c_{3,16}$ such that

$$\mathcal{I}_{2} \le c \int_{0}^{M} \mathbf{1}_{\{r > m \| x - y \|\}} \frac{dr}{r} \le c_{3,16} \log\left(\frac{e}{\|x - y\| \wedge 1}\right).$$
(3.34)

Combining (3.33) and (3.34) one verifies (3.32). It is clear that (3.31), (3.32) and Fubini's theorem imply the second inequality in (3.29).

By using (3.29) and an argument in [18] (see Lemma 2.9) or in ([11], pp.204–206), one can verify that there is an event Ω_0 with probability at least $c_{3,11}^2/(2c_{3,12} \mathcal{E}_0(\nu_0))$ such that for every $\omega \in \Omega_0$, $\{\nu_n(\omega), n \ge 1\}$ has a subsequence that converges weakly to a finite positive measure ν which is supported on $X^{-1}(F) \cap I$. Then, we have

$$\mathbb{P}\Big\{X(I) \cap F \neq \varnothing\Big\} \ge \frac{c_{3,11}^2}{2c_{3,12}\,\mathcal{E}_0(\nu_0)}.$$
(3.35)

Combining this with (3.27) yields the lower bound in (2.5). The proof of Theorem 2.1 is completed.

Proof of Theorem 2.3

Proof of Part (i) It is clear that the second statement in Part (i) follows from (2.7) directly. Hence we proceed to prove (2.7).

Since dim_H is σ -stable ([8]), without loss of generality, we will assume that $F \subseteq [-M, M]^d$ for some constant M > 0 to prove that (2.7) holds almost surely.

Let us choose and fix an arbitrary constant $\gamma > \dim_{\mathrm{H}} F$. Then, for δ_0 in Lemma 3.1, there exists a constant $\delta \in (0, \delta_0)$ and a sequence of balls $\{B(x_j, r_j), j \geq 1\}$ in \mathbb{R}^d (in Euclidean metric) such that $r_j \leq \delta$ for all $j \geq 1$,

$$F \subseteq \bigcup_{j=1}^{\infty} B(x_j, r_j) \quad \text{and} \quad \sum_{j=1}^{\infty} (2r_j)^{\gamma} \le 1.$$
(3.36)

Moreover, we require $x_j \in [-M, M]^d$ for all $j \ge 1$.

For every integer $j \ge 1$, let n_j be the integer such that $2^{-n_j-1} \le r_j < 2^{-n_j}$. We divide I into sub-rectangles $R_{j,i}$ of side-lengths $2^{-n_j/H_\ell}$ ($\ell = 1, \ldots, N$). Note that there are at most $c 2^{n_j Q}$ such rectangles (recall that $Q = \sum_{j=1}^N \frac{1}{H_j}$), and each $R_{j,i}$ is equivalent to a ball of radius 2^{-n_j} in the metric ρ . Let N_j denote the number of $R_{j,i}$'s such that $R_{j,i} \cap X^{-1}(B(x_j, r_j)) \ne \emptyset$. Hence, by using Lemma 3.1, we derive

$$\mathbb{E}(N_j) \le c_{3,17} \, 2^{n_j Q} \, 2^{-n_j d} = c_{3,17} \, 2^{n_j (Q-d)}. \tag{3.37}$$

Since $X^{-1}(F) \subseteq \bigcup_{j=1}^{\infty} X^{-1}(B(x_j, r_j))$, we have obtained a covering of $X^{-1}(F) \cap I$ by a subsequence of the rectangles $\{R_{j,i}\}$ (i.e., those satisfying $R_{j,i} \cap X^{-1}(B(x_j, r_j)) \neq \emptyset$).

For every $k \in \{1, \ldots, N\}$, each rectangle $R_{j,i}$ can be covered by at most $\prod_{\ell=k}^{N} \left(2^{n_j \left(\frac{1}{H_k} - \frac{1}{H_\ell}\right)} + 1\right)$ cubes of side-length $2^{-n_j/H_k}$. In this way, we obtain a covering of $X^{-1}(F) \cap I$ by cubes of side-length $2^{-n_j/H_k}$ which can be used to bound the Hausdorff measure of $X^{-1}(F) \cap I$. Let

$$\beta_k = \sum_{j=1}^k \frac{H_k}{H_j} + N - k - H_k(d - \gamma).$$

It follows from (3.36) and (3.37) that

$$\mathbb{E}\left[\sum_{j=1}^{\infty} N_j \prod_{\ell=k}^{N} \left(2^{n_j\left(\frac{1}{H_k} - \frac{1}{H_\ell}\right)} + 1\right) 2^{-\frac{n_j}{H_k}\beta_k}\right] \le c_{3,18} \sum_{j=1}^{\infty} 2^{-n_j\gamma} \le c_{3,18}.$$
(3.38)

This and Fatou's lemma together imply that $\mathcal{H}_{\beta_k}(X^{-1}(F) \cap I) < \infty$ almost surely. Hence we have $\dim_{\mathrm{H}}(X^{-1}(F) \cap I) \leq \beta_k$ a.s. for every $k \in \{1, \ldots, N\}$. Letting $\gamma \downarrow \dim_{\mathrm{H}} F$ along rational numbers yields (2.7).

Proof of Part (ii) Assume that F is a non-empty set such that $\dim_{\mathrm{H}} F > d - Q$. Note that if $\dim_{\mathrm{H}} F = 0$, then we have $\sum_{j=1}^{N} \frac{1}{H_j} > d$. For any fixed $x \in F$, we have $X^{-1}(F) \supseteq X^{-1}(\{x\})$ and (2.9) follows from Theorem 7.1 in [23] and (2.7). It remains to consider those Borel sets $F \subseteq \mathbb{R}^d$ with $\dim_{\mathrm{H}} F > 0$.

For this purpose, let us note that (2.6) implies

$$\min_{1 \le k \le N} \left\{ \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k (d - \dim_{\mathrm{H}} F) \right\}$$

$$= \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k (d - \dim_{\mathrm{H}} F), \quad \text{if} \quad \sum_{j=1}^{k-1} \frac{1}{H_j} \le d - \dim_{\mathrm{H}} F < \sum_{j=1}^{k} \frac{1}{H_j}.$$
(3.39)

The verification of (3.39) is elementary (cf. Lemma 7 in [1]) and is omitted.

Let $k \in \{1, ..., N\}$ be the unique integer satisfying

$$\sum_{j=1}^{k-1} \frac{1}{H_j} \le d - \dim_{\mathrm{H}} F < \sum_{j=1}^k \frac{1}{H_j}.$$
(3.40)

By (3.39), it is sufficient to prove that for every $\varepsilon > 0$, the inequality

$$\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) \ge \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k(d - \dim_{\mathrm{H}} F) - \varepsilon$$
(3.41)

holds on an event with positive probability (which may depend on ε).

Let us choose and fix an arbitrary constant $\delta \in (0, \dim_{\mathrm{H}} F)$ such that

$$\sum_{j=1}^{k-1} \frac{1}{H_j} < d - \dim_{\mathrm{H}} F + \delta < \sum_{j=1}^k \frac{1}{H_j}$$
(3.42)

(this is possible because of (3.40)) and define

$$\beta_{\delta} := \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k (d - \dim_{\mathrm{H}} F) - H_k \delta.$$
(3.43)

Then it can be verified that $\beta_{\delta} \in (N-k, N-k+1)$, and (3.41) follows from

$$\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) \ge \beta_{\delta} \tag{3.44}$$

on an event with positive probability (which may depend on δ).

Now we prove (3.44). Let $\gamma \in (\dim_{\mathrm{H}} F - \delta, \dim_{\mathrm{H}} F)$ be a constant. Then (3.42) still holds when $\dim_{\mathrm{H}} F - \delta$ is replaced by γ . Then there is a compact set $F_{\gamma} \subseteq F$ such that $\dim_{\mathrm{H}} F_{\gamma} > \gamma$ (see [8]). It follows from Frostman's lemma (cf. [8] or [11]) that there exists a Borel probability measure $\mu_{0,\gamma}$ supported on F_{γ} such that

$$\mu_{0,\gamma}\big(B(x,r)\big) \le c_{3,19} r^{\gamma} \qquad \text{for all } x \in \mathbb{R}^d \text{ and } r > 0, \tag{3.45}$$

where $c_{3,19}$ may depend on γ .

We construct a sequence $\{\mu_{n,\gamma}, n \ge 1\}$ of random Borel measures on I as in (3.28) with ν_0 replaced by $\mu_{0,\gamma}$. We will prove the following statements:

(i) There exist finite positive constants $c_{3,20}$ and $c_{3,21}$ (which may depend on γ) such that

$$\mathbb{E}(\|\mu_{n,\gamma}\|) \ge c_{3,20} \qquad \text{and} \qquad \mathbb{E}(\|\mu_{n,\gamma}\|^2) \le c_{3,21} \tag{3.46}$$

hold for all integers $n \ge 1$.

(ii) For the constant β_{δ} defined by (3.43), there exists a finite positive constant $c_{3,22}$ such that for all $n \ge 1$

$$\mathbb{E} \int_{I} \int_{I} \frac{1}{\|s - t\|^{\beta_{\delta}}} \,\mu_{n,\gamma}(ds) \mu_{n,\gamma}(dt) \le c_{3,22}.$$
(3.47)

The first inequality in (3.46) is derived in the same way as (3.30), because the probability measure $\mu_{0,\gamma}$ is supported on the compact set F_{γ} . Now we prove the second inequality in (3.46). Similar to (3.31), we use Lemma 3.2 to obtain

$$\mathbb{E}(\|\mu_{n,\gamma}\|^2) \le c_{3,23} \int_I \int_I \int_{\mathbb{R}^{2d}} \frac{1}{\max\{\rho^d(s,t), \|x-y\|^d\}} \,\mu_{0,\gamma}(dx)\mu_{0,\gamma}(dy) \,dsdt.$$
(3.48)

By property (3.45) of $\mu_{0,\gamma}$, it can be shown that

$$\int_{\mathbb{R}^{2d}} \frac{1}{\max\{\rho^d(s,t), \|x-y\|^d\}} \,\mu_{0,\gamma}(dx) \mu_{0,\gamma}(dy) \le \frac{c_{3,24}}{\rho(s,t)^{d-\gamma}}.$$
(3.49)

To this end, we break the above integral over $\{(x, y) \in \mathbb{R}^{2d} : ||x - y|| \ge \rho(s, t)\}$ and $\{(x, y) \in \mathbb{R}^{2d} : ||x - y|| < \rho(s, t)\}$ and write them as \mathcal{J}_1 and \mathcal{J}_2 , respectively. For each fixed $x \in \mathbb{R}^d$, let κ_x be the image measure of $\mu_{0,\gamma}$ under the mapping $T : y \mapsto ||x - y||$ from \mathbb{R}^d to \mathbb{R}_+ .

$$\mathcal{J}_{1} = \int_{\mathbb{R}^{d}} \mu_{0,\gamma}(dx) \int_{\{\|x-y\| \ge \rho(s,t)\}} \frac{\mu_{0,\gamma}(dy)}{\|x-y\|^{d}} = \int_{\mathbb{R}^{d}} \mu_{0,\gamma}(dx) \int_{\rho(s,t)}^{+\infty} \frac{1}{r^{d}} \kappa_{x}(dr) \\
\leq d c_{3,19} \int_{\mathbb{R}^{d}} \mu_{0,\gamma}(dx) \int_{\rho(s,t)}^{+\infty} r^{\gamma-d-1} dr = \frac{c_{3,25}}{\rho(s,t)^{d-\gamma}},$$
(3.50)

where the above inequality follows from an integration-by-parts formula and (3.45).

On the other hand, (3.45) implies that

$$\mathcal{J}_{2} = \int_{\mathbb{R}^{d}} \mu_{0,\gamma}(dx) \int_{\{\|x-y\| < \rho(s,t)\}} \frac{\mu_{0,\gamma}(dy)}{\rho(s,t)^{d}} \le \frac{c_{3,19}}{\rho(s,t)^{d-\gamma}}.$$
(3.51)

Hence (3.49) follows from (3.50) and (3.51).

It follows from (3.48), (3.49), Condition (C2) and some elementary computation that

$$\mathbb{E}(\|\mu_{n,\gamma}\|^2) \le c_{3,26} \int_I \int_I \frac{1}{\left[\sum_{j=1}^N |s_j - t_j|^{H_j}\right]^{d-\gamma}} \, ds \, dt := c_{3,21} < \infty.$$
(3.52)

In the above, we have used (3.42) to show the last integral is convergent (cf. [1], p. 432). Thus we have verified the second inequality in (3.46).

Similar to (3.48)–(3.52) and recall that $I = [\varepsilon_0, 1]^N$, we derive

$$\mathbb{E} \int_{I} \int_{I} \frac{1}{\|s-t\|^{\beta_{\delta}}} \mu_{n,\gamma}(ds) \mu_{n,\gamma}(dt) \\
\leq c_{3,27} \int_{I} \int_{I} \frac{1}{\left[\sum_{j=1}^{N} |s_{j} - t_{j}|^{H_{j}}\right]^{d-\gamma} \left(\sum_{j=1}^{N} |s_{j} - t_{j}|\right)^{\beta_{\delta}}} ds dt \qquad (3.53) \\
\leq c_{3,28} \int_{[0,1]^{N}} \frac{1}{\left[\sum_{j=1}^{N} t_{j}^{H_{j}}\right]^{d-\gamma} \left(\sum_{j=1}^{N} t_{j}\right)^{\beta_{\delta}}} dt.$$

Let us denote the last integral in (3.53) by \mathcal{G} and apply Lemma 3.3 to verify that it is convergent. Without loss of generality, we assume $k \geq 2$, otherwise we jump to (3.56) directly. We integrate $[dt_1]$ first. By (3.42) we see that $H_1(d - \gamma) > 1$. Hence we can use (3.25) of Lemma 3.3 with $\alpha = H_1$, $\eta = d - \gamma$, $A = \sum_{j=2}^{N} t_j^{H_j}$ and $B = \sum_{j=2}^{N} t_j$ to get

$$\mathcal{G} \le c_{3,29} \int_0^1 dt_N \cdots \int_0^1 \frac{1}{\left(\sum_{j=2}^N t_j^{H_j}\right)^{d-\gamma-1/H_1} \left(\sum_{j=2}^N t_j\right)^{\beta_{\delta}}} dt_2.$$
(3.54)

We repeat this procedure for integrating dt_2, \ldots, dt_{k-1} to obtain

$$\mathcal{G} \le c_{3,30} \, \int_0^1 dt_N \cdots \int_0^1 \frac{dt_k}{\left(\sum_{j=k}^N t_j^{H_j}\right)^{d-\gamma - \sum_{j=1}^{k-1} 1/H_j} \left(\sum_{j=k}^N t_j\right)^{\beta_{\delta}}}.$$
 (3.55)

To integrate $[dt_k]$ in (3.55), note that $H_k(d - \gamma - \sum_{j=1}^{k-1} 1/H_j) < 1$. It follows from (3.26) that

$$\mathcal{G} \leq c_{3,31} \left[\int_{0}^{1} dt_{N} \cdots \int_{0}^{1} \frac{dt_{k+1}}{\left(\sum_{j=k+1}^{N} t_{j}\right)^{H_{k}(d-\gamma-\sum_{j=1}^{k-1} 1/H_{j})+\beta_{\delta}-1}} + 1 \right]$$

$$= c_{3,31} \left[\int_{0}^{1} dt_{N} \cdots \int_{0}^{1} \frac{dt_{k+1}}{\left(\sum_{j=k+1}^{N} t_{j}\right)^{N-k+H_{k}(\dim_{\mathrm{H}} F-\delta-\gamma)}} + 1 \right].$$
(3.56)

Since $\dim_{\mathrm{H}} F - \delta < \gamma$, the last integral is convergent. Hence (3.47) holds.

By using (3.46) and the argument in [18] (see Lemma 2.9) or in ([11], pp.204–206) again, one can verify that there is an event Ω_{γ} with probability at least $c_{3,20}^2/(2c_{3,21})$ such that for every $\omega \in \Omega_{\gamma}$, $\{\mu_{n,\gamma}(\omega), n \ge 1\}$ has a subsequence that converges weakly to a positive finite measure μ_{γ} , which is supported on $X^{-1}(F) \cap I$.

On the other hand, (3.47) and the monotone convergence theorem together imply that, on the event Ω_{γ} , the β_{δ} -energy of μ_{γ} is finite. Therefore by Frostman's theorem (cf. [8] or [11]) we derive that (3.41) holds on an event with probability at least $c_{3,20}^2/(2c_{3,21})$. This finishes the proof of Theorem 2.3. Finally we show that, under the extra condition (S), the proof of Theorem 2.3 can be strengthened to prove Theorem 2.4.

Proof of Theorem 2.4 The second equality in (2.12) is implied by (3.39). Hence it only remains to prove that, if $F \subseteq \mathbb{R}^d$ is a Borel set satisfying

$$\sum_{j=1}^{k-1} \frac{1}{H_j} \le d - \dim_{\mathrm{H}} F < \sum_{j=1}^k \frac{1}{H_j}$$
(3.57)

for some integer $k \in \{1, \ldots, N\}$, then

$$\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) \ge \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k (d - \dim_{\mathrm{H}} F)$$
(3.58)

holds with positive probability. We will prove the following statement: there exist constants $\delta_0 \in (0, \dim_{\mathrm{H}} F)$ and $c_{3,32} > 0$ such that for all $\delta \in (0, \delta_0)$ we have

$$\sum_{j=1}^{k-1} \frac{1}{H_j} < d - \dim_{\mathrm{H}} F - \delta < \sum_{j=1}^k \frac{1}{H_j}$$
(3.59)

and

$$\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) \ge \sum_{j=1}^{k} \frac{H_k}{H_j} + N - k - H_k (d - \dim_{\mathrm{H}} F) - H_k \delta$$
(3.60)

holds with probability greater than $c_{3,32}$ (which does not depend on δ). Letting $\delta \downarrow 0$, this shows that (3.58) holds with probability greater than $c_{3,32}$.

The proof of the above statement is almost the same as the proof of (3.44). Since (3.57) is fulfilled, we can choose $\delta_0 \in (0, \dim_{\mathrm{H}} F)$ such that (3.59) holds for all $\delta \in (0, \delta_0)$. Let $\delta \in (0, \delta_0)$, $\gamma \in (\dim_{\mathrm{H}} F - \delta, \dim_{\mathrm{H}} F)$ and define β_{δ} by (3.43). Let $\mu_{0,\gamma}$ be the probability measure with compact support in F given by Condition (S), and let { $\mu_{n,\gamma}, n \geq 1$ } be the sequence of random Borel measures on I constructed in the proof of Theorem 2.3. The only difference is that we will show the constants $c_{3,20}$ and $c_{3,21}$ in (3.46) are independent of δ and γ .

Since F is bounded, we derive as in (3.30) that $c_{3,20}$ is independent of $\mu_{0,\gamma}$, γ and δ . Next we prove that $\mathbb{E}(\|\mu_{n,\gamma}\|^2) \leq c_{3,21}$ for all $\gamma \in (\dim_{\mathrm{H}} F - \delta, \dim_{\mathrm{H}} F)$, and the constant $c_{3,21}$ only depends on F, δ_0 , $c_{2,9}$, I and H.

Recall from (3.48) that

$$\mathbb{E}(\|\mu_{n,\gamma}\|^2) \le c_{3,23} \int_I \int_I \int_{\mathbb{R}^{2d}} \frac{1}{\max\{\rho^d(s,t), \|x-y\|^d\}} \,\mu_{0,\gamma}(dx)\mu_{0,\gamma}(dy) \,dsdt.$$
(3.61)

Since $\gamma > \dim_{\rm H} F - \delta_0$, the probability measure $\mu_{0,\gamma}$ also satisfies

$$\mu_{0,\gamma}\big(B(x,r)\big) \le c_{2,9} r^{\dim_{\mathrm{H}} F - \delta_0} \qquad \text{for all } x \in \mathbb{R}^d \text{ and } r \in (0,1).$$

We can assume without loss of generality that $c_{2,9} \ge 1$. Then since $\mu_{0,\gamma}$ is probability measure and $\dim_{\mathrm{H}} F - \delta_0 > 0$, we have

$$\mu_{0,\gamma}\big(B(x,r)\big) \le c_{2,9} r^{\dim_{\mathrm{H}} F - \delta_0} \qquad \text{for all } x \in \mathbb{R}^d \text{ and } r \in (0,\infty).$$
(3.62)

Similar to the verification of (3.49), we derive from (3.62) that

$$\int_{\mathbb{R}^{2d}} \frac{1}{\max\{\rho^d(s,t), \|x-y\|^d\}} \,\mu_{0,\gamma}(dx) \mu_{0,\gamma}(dy) \le \frac{c_{3,33}}{\rho(s,t)^{d-\dim_{\mathrm{H}}F+\delta_0}}.$$
(3.63)

Here the constant $c_{3,33}$ depends only on d, F, $c_{2,9}$ and δ_0 . It follows from (3.61) and (3.63) that

$$\mathbb{E}(\|\mu_{n,\gamma}\|^2) \le c_{3,34} \int_I \int_I \frac{1}{\left[\sum_{j=1}^N |s_j - t_j|^{H_j}\right]^{d - \dim_{\mathrm{H}}F + \delta_0}} \, ds \, dt := c_{3,21} < \infty.$$
(3.64)

This verified the second inequality in (3.46) and the constant $c_{3,21}$ is *independent* of γ and δ .

As shown in the proof of Theorem 2.3, (3.46) implies the existence of a subsequence of $\{\mu_{n,\gamma}, n \geq 1\}$ that converges weakly to a positive finite measure μ_{γ} , which is supported on $X^{-1}(F) \cap I$ on an event Ω_{γ} with probability $c_{3,20}^2/(2c_{3,21}) := c_{3,32}$. The constant $c_{3,32}$ does not depend on γ and δ .

Observe that (3.47) remains valid for $\mu_{n,\gamma}$. On the event Ω_{γ} the β_{δ} -energy of μ is finite almost everwhere. Thus by Frostman's theorem, we derive

$$\dim_{\mathrm{H}} \left(X^{-1}(F) \cap I \right) \ge \beta_{\delta} \tag{3.65}$$

on an event with probability greater than $c_{3,32}$ which does not depend on γ and δ . This finishes the proof of Theorem 2.4.

4 Applications to stochastic partial differential equations

Theorems 2.1, 2.3 and 2.4 are applicable to solutions of stochastic partial differential equations (SPED) such as linear hyperbolic SPDE considered by Dalang and Nulart [7] and nonlinear stochastic heat equations considered by Dalang, Khoshnevisan and Nualart [5]. In this section we only consider the Hausdorff dimension of the inverse images of nonlinear stochastic heat equations in [5].

Let $\dot{W} = (\dot{W}_1, \ldots, \dot{W}_d)$ be a space-time white noise in \mathbb{R}^d . That is, the components $\dot{W}_1(x,t), \ldots, \dot{W}_d(x,t)$ of $\dot{W}(x,t)$ are independent space-time white noises, which are generalized Gaussian processes with covariance given by

$$\mathbb{E}[\dot{W}_i(t,x)\dot{W}_i(y,s)] = \delta(x-y)\delta(t-s), \quad (i=1,\ldots,d),$$

where $\delta(\cdot)$ is the Dirac delta function. For all $1 \leq j \leq d$, let $b_j : \mathbb{R}^d \to \mathbb{R}$ be globally Lipschitz and bounded functions, and $\sigma := (\sigma_{i,j})$ be a deterministic $d \times d$ invertible matrix.

Consider the system of stochastic partial differential equations

$$\frac{\partial u_i}{\partial t}(t,x) = \frac{\partial^2 u_i}{\partial x^2}(t,x) + \sum_{j=1}^d \sigma_{i,j} \dot{W}_j(t,x) + b_i \big(u(t,x) \big)$$
(4.1)

for $1 \le i \le d$, $t \in [0, T]$ and $x \in [0, 1]$, with the initial conditions u(0, x) = 0 for all $x \in [0, 1]$, and the Neumann boundary conditions

$$\frac{\partial u_i}{\partial x}(t,0) = \frac{\partial u_i}{\partial x}(t,1) = 0, \qquad 0 \le t \le T.$$
(4.2)

In the above, $u(t, x) = (u_1(t, x), \dots, u_d(t, x)).$

Following Dalang, Khoshnevisan and Nualart [5], Equation (4.1) can be interpreted rigorously as follows. Let $W_i = \{W_i(s, x), s \in \mathbb{R}_+, x \in [0, 1]\}$ $(1 \le i \le d)$ be independent Brownian sheets, defined on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$, and set $W = (W_1, \ldots, W_d)$. For $t \ge 0$, let $\mathcal{F}_t = \sigma(W(s, x) : s \in [0, t], x \in [0, 1])$. We say that a random field $u = \{u(t, x), t \in [0, T], x \in [0, 1]\}$ is adapted to (\mathcal{F}_t) if u(t, x) is \mathcal{F}_t -measurable for every $(t, x) \in [0, T] \times [0, 1]$. A random field $u = \{u(t, x), t \in [0, T], x \in [0, 1]\}$ is a solution of (4.1) if u is adapted to (\mathcal{F}_t) and if for every $i \in \{1, \ldots, d\}, t \in [0, T]$ and $x \in [0, 1]$,

$$u_i(t,x) = \int_0^t \int_0^1 G_{t-r}(x,v) \sum_{j=1}^d \sigma_{i,j} W_j(dr,\,dv) + \int_0^t \int_0^1 G_{t-r}(x,v) \,b_i(u(t,x)) \,drdv, \quad (4.3)$$

where $G_t(x, y)$ is the Green kernel for the heat equation with Neumann boundary conditions (see Walsh [19]).

For the linear form of (4.1) [i.e., $b \equiv 0$ and $\sigma \equiv I_d$ (the $d \times d$ identity matrix)], Mueller and Tribe [16] found necessary and sufficient conditions (in terms of the dimension d) for its solution u to hit points or to have double points of various types. Wu and Xiao [20] further studied the fractal properties of the sample paths of u. In particular, they obtained the Hausdorff dimensions of the level sets and the set of double times of u.

More generally, Dalang, Khoshnevisan and Nualart [5] studied hitting probabilities for the non-linear equation (4.1) by providing sufficient conditions and necessary conditions for a Borel set $F \subseteq \mathbb{R}^d$ to be polar for the processes $\{u(t,x), t \in [0,T], x \in [0,1]\}, \{u(t,x), x \in [0,1]\}$ $(t \in (0,T]$ is fixed) and $\{u(t,x), t \in [0,T]\}$ $(x \in [0,1]$ is fixed). They also determined the Hausdorff dimensions of the range and level sets of these processes.

In the following, we apply Theorem 2.3 to further determine the Hausdorff dimension of inverse images of the processes $u = \{u(t, x), t \in [0, T], x \in [0, 1]\}, u_t = \{u(t, x), x \in [0, 1]\}$ $(t \in (0, T] \text{ is fixed}) \text{ and } u_x = \{u(t, x), t \in [0, T]\} (x \in [0, 1] \text{ is fixed}).$ As shown by Proposition 5.1 in [5], it is sufficient to consider these problems for the solution of equation (4.1) in the following drift-free case [i.e., $b \equiv 0$]:

$$\frac{\partial u}{\partial t}(t,x) = \frac{\partial^2 u}{\partial x^2}(t,x) + \sigma \,\dot{W}.\tag{4.4}$$

The solution of (4.4) is the mean zero Gaussian random field $u = \{u(t, x), t \in [0, T], x \in [0, 1]\}$ with values in \mathbb{R}^d defined by

$$u(t,x) = \int_0^t \int_0^1 G_{t-r}(x,y) \, \sigma W(dr,dy), \qquad \forall \ t \in [0,T], \ x \in [0,1].$$
(4.5)

Moreover, since the matrix σ is invertible, a change of variables shows (see the proof of Proposition 4.1 in [5]) that $v := \sigma^{-1}u$ solves the following uncoupled system of SPDE

$$\frac{\partial v}{\partial t}(t,x) = \frac{\partial^2 v}{\partial x^2}(t,x) + \dot{W}.$$
(4.6)

Note that both processes u and v have the same hitting probability and Hausdorff dimension properties. Therefore, without loss of generality, we will assume that $\sigma = I_d$ in (4.4).

The following is a consequence of Lemmas 4.2 and 4.3 of Dalang, Khoshnevisan and Nualart [5].

Lemma 4.1 Let $u = \{u(t,x), t \in [0,T], x \in [0,1]\}$ be the solution of (4.4). Then for every $t_0 \in (0,T)$, there exist positive and finite constants $c_{4,1}, \ldots, c_{4,5}$ such that the following hold:

(i) For all
$$t \in I = [t_0, T] \times [0, 1]$$
, we have $c_{4,1} \leq \mathbb{E}[u(t, x)^2] \leq c_{4,2}$ and
 $c_{4,3} \left(|t - s|^{1/4} + |x - y|^{1/2} \right)^2 \leq \mathbb{E} \left[\left(u(t, x) - u(s, y) \right)^2 \right]$
 $\leq c_{4,4} \left(|t - s|^{1/4} + |x - y|^{1/2} \right)^2.$
(4.7)

(ii) For all $(t, x), (s, y) \in I$,

$$\operatorname{Var}\left(u(t,x)\big|u(s,y)\right) \ge c_{4,5}\left(|t-s|^{1/4} + |x-y|^{1/2}\right)^2.$$
(4.8)

In other words, Lemma 4.1 states that the Gaussian random field u satisfies Conditions (C1) and (C2) with $H_1 = 1/4$ and $H_2 = 1/2$.

Proof It follows from (4.5) that

$$\mathbb{E}\left[u(t,x)^2\right] = \int_0^t dr \int_0^1 \left(G_{t-r}(x,v)\right)^2 dv.$$

It can be verified that the function $\sigma^2(t, x) := \mathbb{E}[u(t, x)^2]$ is continuous in (t, x) and positive on *I*. This implies the first conclusion of the lemma. The inequality (4.7) is the same as (4.11) in Lemma 4.2 of Dalang, Khoshnevisan and Nualart [5].

To prove (4.8), we note that Lemma 4.3 of Dalang, Khoshnevisan and Nualart [5] shows that

$$\det \operatorname{Cov}(u(t,x), u(s,y)) \ge c_{4,6} \left(|t-s|^{1/4} + |x-y|^{1/2} \right)^2.$$
(4.9)

Hence (4.8) follows from (3.16), (4.9) and the fact that $\mathbb{E}[u(s,y)^2] \leq c_{4,2}$ for all $t \in I$. This proves Lemma 4.1.

Remark 4.2 Lemma 4.1 implies that, for every fixed $t \in [t_0, T]$, the Gaussian process $u_t = \{u(t, x), x \in [0, 1]\}$ has the following properties:

$$c_{4,1} \le \mathbb{E}[u_t(x)^2] \le c_{4,2}, \quad \forall x \in [0,1],$$
(4.10)

$$c_{4,3} |x - y| \le \mathbb{E} \Big[\big(u_t(x) - u_t(y) \big)^2 \Big] \le c_{4,4} |x - y|, \qquad \forall x, y \in [0, 1]$$
(4.11)

and

$$\operatorname{Var}(u_t(x)|u_t(y)) \ge c_{4,5} |x-y|, \qquad \forall x, y \in [0,1].$$
(4.12)

This shows that the Gaussian process u_t satisfies the *two-point local nondeterminism* with index $H_2 = 1/2$. Similarly, for every fixed $x \in (0, 1]$, the Gaussian process $u_x = \{u(t, x), t \in [t_0, T]\}$ satisfies the two-point local nondeterminism with $H_1 = 1/4$.

Hence we can apply Theorem 2.1 to recover the results of Dalang, Khoshnevisan and Nualart [5] on the hitting probabilities of the solution to (4.1). The following theorem extends the Hausdorff dimension results in [5].

Theorem 4.3 Let $u = \{u(t, x), t \in [0, T], x \in [0, 1]\}$ denote the solution to (4.1) and let $F \subseteq \mathbb{R}^d$ be a Borel set. The following conclusions hold:

(i) If $\dim_{\mathbf{H}} F \ge d - 6$, then

$$\left\| \dim_{\mathrm{H}} \left(u^{-1}(F) \cap I \right) \right\|_{L^{\infty}(\mathbb{P})} = \begin{cases} 2 - \frac{1}{4} (d - \dim_{\mathrm{H}} F) & \text{if } 0 \le d - \dim_{\mathrm{H}} F < 4, \\ 3 - \frac{1}{2} (d - \dim_{\mathrm{H}} F) & \text{if } 4 \le d - \dim_{\mathrm{H}} F \le 6. \end{cases}$$
(4.13)

(ii) If $\dim_{\mathrm{H}} F \geq d-2$, then for every fixed $t \geq t_0$ we have

$$\left\| \dim_{\mathrm{H}} \left(u_t^{-1}(F) \cap [0,1] \right) \right\|_{L^{\infty}(\mathbb{P})} = 1 - \frac{1}{2} \left(d - \dim_{\mathrm{H}} F \right).$$
(4.14)

(iii) If $\dim_{H} F \ge d-4$, then for every fixed $x \in [0,1]$ we have

$$\left\| \dim_{\mathrm{H}} \left(u_x^{-1}(F) \cap [0, T] \right) \right\|_{L^{\infty}(\mathbb{P})} = 1 - \frac{1}{4} \left(d - \dim_{\mathrm{H}} F \right).$$
(4.15)

Proof The conclusions follow from Theorem 2.3, Lemma 4.1 and Remark 4.2. \Box

References

- A. Ayache and Y. Xiao (2005), Asymptotic properties and Hausdorff dimensions of fractional Brownian sheets. J. Fourier Anal. Appl. 11, 407–439.
- [2] H. Biermé, A. Estrade, M. M. Meerschaert and Y. Xiao (2007), Sample path properties of operator scaling Gaussian random fields. *In Preparation*.
- [3] H. Biermé and C. Lacaux (2007), Hölder regularity for operator scaling stable random fields. Submitted.
- [4] H. Biermé, M. M. Meerschaert and H.-P. Scheffler (2007), Operator scaling stable random fields. Stoch. Process. Appl. 117, 312–332.
- [5] R. C. Dalang, D. Khoshnevisan and E. Nualart (2007a), Hitting probabilities for systems of non-linear stochastic heat equations with additive noise. *Latin Amer. J. Probab. Statist. (Alea)* 3, 231–271.
- [6] R. C. Dalang, D. Khoshnevisan and E. Nualart (2007b), Hitting probabilities for the non-linear stochastic heat equation with multiplicative noise. *Submitted*.
- [7] R. C. Dalang and E. Nualart (2004), Potential theory for hyperbolic SPDEs. Ann. Probab. 32, 2099–2148.
- [8] K. J. Falconer (1990), Fractal Geometry Mathematical Foundations and Applications. John Wiley & Sons Ltd., Chichester.
- [9] J. Hawkes (1971), On the Hausdorff dimension of the intersection of the range of a stable process with a Borel set. Z. Wahrsch. Verw. Geb. 19, 90–102.
- [10] A. Jonsson and H. Wallin (1984), Function Spaces on Subsets of \mathbb{R}^n . Harwood Academic Publishers, London.
- [11] J.-P. Kahane (1985), Some Random Series of Functions. 2nd edition, Cambridge University Press, Cambridge.

- [12] D. Khoshnevisan (2002), Multiparameter Processes: An Introduction to Random Fields. Springer, New York.
- [13] D. Khoshnevisan and Y. Xiao (2005), Lévy processes: capacity and Hausdorff dimension. Ann. Probab. 33, 841–878.
- [14] D. Khoshnevisan and Z. Shi (1999), Brownian sheet and capacity. Ann. Probab. 27, 1135–1159.
- [15] M. M. Meerschaert and H. P. Scheffler (2001), Limit Distributions for Sums of Independent Random Vectors. John Wiley & Sons Inc., New York.
- [16] C. Mueller and R. Tribe (2002), Hitting probabilities of a random string. *Electron. J. Probab.* 7, Paper No. 10, 1–29.
- [17] D. Monrad and L. D. Pitt (1987), Local nondeterminism and Hausdorff dimension. Progress in Probability and Statistics. Seminar on Stochastic Processes 1986, (E, Cinlar, K. L. Chung, R. K. Getoor, Editors), pp. 163–189, Birkhäuser, Boston.
- [18] F. Testard (1986), Polarité, points multiples et géométrie de certain processus gaussiens. Publ. du Laboratoire de Statistique et Probabilités de l'U.P.S. Toulouse, 01–86.
- [19] J. B. Walsh (1986), An introduction to stochastic partial differential equations. École d'été de probabilités de Saint-Flour, XIV—1984, 265–439, Lecture Notes in Math., 1180, Springer, Berlin.
- [20] D. Wu and Y. Xiao (2006), Fractal properties of random string processes. IMS Lecture Notes-Monograph Series-High Dimensional Probability. 51, 128–147.
- [21] D. Wu and Y. Xiao (2007), Geometric properties of the images of fractional Brownian sheets. J. Fourier Anal. Appl. 13, 1–37.
- [22] Y. Xiao (1999), Hitting probabilities and polar sets for fractional Brownian motion. Stoch. Stoch. Rep. 66, 121–151.
- [23] Y. Xiao (2007), Sample path properties of anisotropic Gaussian random fields. Submitted.

HERMINE BIERMÉ: MAP5, Université Paris Descartes, 45 rue des Saints-Pères, 75006 Paris, France E-mail: hermine.bierme@math-info.univ-paris5.fr URL: http://www.math-info.univ-paris5.fr/~bierme/

CÉLINE LACAUX: Institut Elie Cartan, Nancy-Université, CNRS, INRIA, Boulevard des Aiguillettes, BP 239, F-54506 Vandoeuvre-Lès-Nancy, France E-mail: Celine.Lacaux@iecn.u-nancy.fr URL: http://www.iecn.u-nancy.fr/~lacaux/

YIMIN XIAO: Department of Statistics and Probability, A-413 Wells Hall, Michigan State University, East Lansing, MI 48824, U.S.A. E-mail: xiao@stt.msu.edu URL: http://www.stt.msu.edu/~xiaoyimi/