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# An alternative expression for the Black-Scholes formula in terms of Brownian first and last passage times 

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## 1 Introduction and main results

- a) Let $\left(B_{t}, t \geq 0\right)$ denote a one-dimensional Brownian motion starting from 0 , and let

$$
\mathcal{E}_{t}=\exp \left(B_{t}-\frac{t}{2}\right),(t \geq 0)
$$

A reduced form of the celebrated Black-Scholes formula is the following.
$\left(\mathbf{1}_{K}\right)$

$$
E\left[\left(\mathcal{E}_{t}-K\right)^{+}\right]=\mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}+\frac{\sqrt{t}}{2}\right)-K \mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}-\frac{\sqrt{t}}{2}\right)
$$

where $K \geq 0$, and as usual:

$$
\mathcal{N}(x)=\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{x} d y \exp \left(-\frac{y^{2}}{2}\right)
$$

In fact, formula $\left(\mathbf{1}_{\mathbf{K}}\right)$ may be split into two parts:
$\left(\mathbf{1}_{\mathbf{K}}^{+}\right)$

$$
E\left[\mathcal{E}_{t} 1_{\mathcal{E}_{t}>K}\right]=\mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}+\frac{\sqrt{t}}{2}\right)
$$

$\left(\mathbf{1}_{\mathbf{K}}^{-}\right)$

$$
K P\left(\mathcal{E}_{t}>K\right)=K \mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}-\frac{\sqrt{t}}{2}\right)
$$

As noted in Section 5 of [1], formula ( $\mathbf{1}_{\mathbf{K}}^{ \pm}$) is obtained in an elementary manner, after performing the change of probability:

$$
P_{\mathcal{F}_{t}}^{\prime}=\mathcal{E}_{t} \cdot P_{\left.\right|_{\mathcal{F}_{t}}}
$$

which transforms $\left(B_{t}\right)$ in $\left(B_{t}+t\right)$ hence $\left(B_{t}-\frac{t}{2}\right)$ in $\left(B_{t}+\frac{t}{2}\right)$.

- b) In this Note, we give and discuss a different representation of $\left(\mathbf{1}_{\mathbf{K}}^{ \pm}\right)$.

Theorem 1. For any $K \geq 0$, there are the representations:

$$
\begin{array}{ll}
\left(\mathbf{2}_{\mathbf{K}}^{-}\right) & E\left[\mathcal{E}_{t} 1_{\left(\mathcal{E}_{t}>K\right)}\right]-K P\left(\mathcal{E}_{t}>K\right)=P\left(G_{(\ell n K)}^{(1 / 2)} \leq t\right) \\
\left(2_{\mathbf{K} \geq 1}^{+}\right) & E\left[\mathcal{E}_{t} 1_{\mathcal{E}_{t}>K}\right]+K P\left(\mathcal{E}_{t}>K\right)=P\left(T_{(\ell n K)}^{(1 / 2)} \leq t\right) \\
\left(\mathbf{2}_{\mathbf{K} \leq 1}^{+}\right) & E\left[\mathcal{E}_{t} 1_{\mathcal{E}_{t}<K}\right]+K P\left(\mathcal{E}_{t}<K\right)=P\left(T_{(\ell n K)}^{(1 / 2)} \leq t\right)
\end{array}(0 \leq K \leq 1) .
$$

or, equivalently:

$$
\left(2_{\mathbf{K} \leq 1}^{++}\right) \quad E\left[\mathcal{E}_{t} 1_{\mathcal{E}_{t}>K}\right]+K P\left(\mathcal{E}_{t}>K\right)=K+P\left(T_{(\ell n K)}^{(1 / 2)}>t\right) \quad(0 \leq K \leq 1)
$$

where, for $\nu \in \mathbb{R}$, and $B_{t}^{(\nu)} \equiv B_{t}+\nu t$, we write:

$$
T_{a}^{(\nu)}=\inf \left\{t: B_{t}^{(\nu)}=a\right\} ; G_{a}^{(\nu)}=\sup \left\{t: B_{t}^{(\nu)}=a\right\}
$$

Comment and Complements about Theorem 1:
(i) Our motivation to prove formulae such as $\left(\mathbf{2}_{\mathbf{K}}^{-}\right)$was our desire to obtain an expression on the RHS showing in a clear manner that the LHS $\left(=E\left(\left(\mathcal{E}_{t}-K\right)_{+}\right)\right)$is an increasing function of $t$.
This is not clear from $\left(\mathbf{1}_{\mathbf{K}}\right)$, although this property of increase is a consequence of the submartingale property of $\left(\mathcal{E}_{t}-K\right)^{+}$; see Section 4 for a more extended discussion.
(ii) Obviously, an equivalent presentation of the "system" $\left(\mathbf{2}_{\mathbf{K}}^{ \pm}\right)$is, for $K \geq 1$ :

$$
\left(3_{\mathbf{K} \geq 1}^{+}\right) \quad E\left[\mathcal{E}_{t} 1_{\left(\mathcal{E}_{t}>K\right)}\right]=\frac{1}{2}\left\{P\left(T_{(\ell n K)}^{(1 / 2)} \leq t\right)+P\left(G_{(\ell n K)}^{(1 / 2)} \leq t\right)\right\}
$$

$\left(\mathbf{3}_{\mathbf{K} \geq \mathbf{1}}^{-}\right) \quad K P\left(\mathcal{E}_{t}>K\right)=\frac{1}{2}\left\{P\left(T_{(\ell n K)}^{(1 / 2)} \leq t\right)-P\left(G_{(\ell n K)}^{(1 / 2)} \leq t\right)\right\}$

$$
=\frac{1}{2} P\left(T_{(\ell n K)}^{(1 / 2)} \leq t \leq G_{(\ell n K)}^{(1 / 2)}\right)
$$

and, for $0 \leq K \leq 1$ :
$\left(3_{\mathbf{K} \leq 1}^{+}\right) \quad E\left[\mathcal{E}_{t} 1_{\left(\mathcal{E}_{t}>K\right)}\right]=\frac{1}{2}\left\{1+K-P\left(T_{(\ell n K)}^{(1 / 2)} \leq t \leq G_{(\ell n K)}^{(1 / 2)}\right)\right\}$
$\left(\mathbf{3}_{\mathbf{K} \leq 1}^{-}\right) K P\left(\mathcal{E}_{t}>K\right)=\frac{1}{2}\left\{1+K-\left[P\left(T_{(\ell n K)}^{(1 / 2)} \leq t\right)+P\left(G_{(\ell n K)}^{(1 / 2)} \leq t\right)\right]\right\}$
(iii) In order to give formulae ( $\mathbf{2}_{\mathbf{K}}^{ \pm}$) an "explicit" character, we now recall the distributions of

$$
T_{a}^{(\nu)}=\inf \left\{t: B_{t}^{(\nu)}=a\right\} \text { and } G_{a}^{(\nu)}=\sup \left\{t: B_{t}^{(\nu)}=a\right\},
$$

for $\nu>0$, and $a>0$ (these formulae will then be used with $\nu=\frac{1}{2}$ and $a=\ln K)$ : denoting by $p_{t}^{(\nu)}(a)$ the density of $B_{t}^{(\nu)}$, we have:

$$
\begin{equation*}
P\left(T_{a}^{(\nu)} \in d t\right)=\left(\frac{a}{t}\right) p_{t}^{(\nu)}(a) d t \equiv \frac{a}{\sqrt{2 \pi t^{3}}} \exp \left(-\frac{1}{2 t}(a-\nu t)^{2}\right) d t \tag{4}
\end{equation*}
$$

whereas

$$
\begin{equation*}
P\left(G_{a}^{(\nu)} \in d t\right)=\nu p_{t}^{(\nu)}(a) d t \equiv \frac{\nu}{\sqrt{2 \pi t}} \exp \left(-\frac{1}{2 t}(a-\nu t)^{2}\right) d t \tag{5}
\end{equation*}
$$

Formula (4) may be obtained from the combination of (4) for $\nu=0$, which is very well known, followed by an application of the Cameron-Martin relationship between the laws of $B^{(\nu)}$ and $B$. Formula (5) is a particular case of a more general formula for last passage times of transient diffusions, obtained in Pitman-Yor [3].
(iv) Although this is not strictly necessary at this point (but will be useful in our proof of Theorem 1), we also present the distributions of $T_{a}^{(-\nu)}$ and $G_{a}^{(-\nu)}$, for $a \geq 0$.
In fact, they may be obtained easily from those of $T_{a}^{(\nu)}$ and $G_{a}^{(\nu)}$ thanks to the Cameron-Martin absolute continuity relationships:

$$
\begin{equation*}
W_{\mid \mathcal{F}_{T_{a}} \cap\left(T_{a}<\infty\right)}^{(-\nu)}=\exp (-2 \nu a) \cdot W_{\mid \mathcal{F}_{T_{a}}}^{(\nu)} \tag{6}
\end{equation*}
$$

$$
W_{\mid \mathcal{F}_{G_{a}} \cap\left(G_{a}>0\right)}^{(-\nu)}=\exp (-2 \nu a) \cdot W_{\mid \mathcal{F}_{G_{a}}}^{(\nu)}
$$

where, here, $W^{(\mu)}$ denotes the law of $\left(B_{t}+\mu t, t \geq 0\right)$ on canonical space, and $T_{a}$, resp. $G_{a}$, is the first, resp. last, hitting time of $a$ by the coordinate process.
Thus, we deduce from formulae (6), (4) and (5) that:

$$
P\left(G_{a}^{(-\nu)}>0\right)=P\left(T_{a}^{(-\nu)}<\infty\right)=\exp (-2 \nu a)
$$

whereas:

$$
\begin{equation*}
P\left(T_{a}^{(-\nu)} \in d t\right)=\left(\frac{a}{t}\right) p_{t}^{(-\nu)}(a) d t=\frac{a}{\sqrt{2 \pi t^{3}}} \exp \left(-\frac{1}{2 t}(a+\nu t)^{2}\right) d t \tag{7}
\end{equation*}
$$

$$
P\left(G_{a}^{(-\nu)} \in d t\right)=\nu p_{t}^{(-\nu)}(a) d t=\frac{\nu}{\sqrt{2 \pi t}} \exp \left(-\frac{1}{2 t}(a+\nu t)^{2}\right) d t
$$

- c) Organisation of the remainder of the paper:
- In Section 2, we prove Theorem 1, independently from formulae ( $\mathbf{1}_{\mathbf{K}}^{ \pm}$)
- In Section 3, we give an elementary proof of the agreement between formulae ( $\mathbf{1}_{\mathbf{K}}^{ \pm}$) and ( $\mathbf{3}_{\mathbf{K}}^{ \pm}$)
- Section 4 concludes, by setting the matter in a broader context.


## 2 Proof of Theorem 1

Clearly, in order to prove Theorem 1, it now suffices to prove $\left(\mathbf{2}_{\mathbf{K}}^{-}\right),\left(\mathbf{3}_{\mathbf{K} \geq \mathbf{1}}^{-}\right)$ and $\left(\mathbf{2}_{\mathbf{K} \leq 1}^{+}\right)$.
a) Proof of $\left(\mathbf{2}_{\mathbf{K}}^{-}\right)$(for any $\left.K \geq 0\right)$

We shall show that

$$
E\left[\left(\mathcal{E}_{t}-K\right)^{+}\right] \underset{(i)}{=} K P\left(0<G_{(\ln K)}^{(-1 / 2)} \leq t\right) \underset{(i i)}{=} P\left(G_{(\ln K)}^{(+1 / 2)} \leq t\right)
$$

The proof of (ii) follows from the relationship between the laws of $G_{a}^{(-\nu)}$ and $G_{a}^{(\nu)}$ as discussed in Section 1; namely:

$$
P\left(G_{a}^{(-\nu)}>0\right)=\exp (-2 \nu a) \text { and } P\left(G_{a}^{(-\nu)} \in d t \mid G_{a}^{(-\nu)}>0\right)=P\left(G_{a}^{(\nu)} \in d t\right)
$$

For the proof of $(i)$, we rely upon the following formula

$$
\begin{equation*}
P\left(G_{a}^{(\mu)} \geq t \mid \mathcal{F}_{t}\right)=\left(\frac{\exp (2 \mu a)}{\exp \left(2 \mu B_{t}^{(\mu)}\right)}\right) \wedge 1 \tag{8}
\end{equation*}
$$

which is valid for all $\mu \in \mathbb{R}$; this is a particular case of the results for last passage times of a transient real-valued diffusion, as discussed in Pitman-Yor [3].
In particular, for $\mu=-\nu, \nu>0$, we get:

$$
\exp (2 \nu a) P\left(0<G_{a}^{(-\nu)} \leq t \mid \mathcal{F}_{t}\right)=\left(\exp \left(2 \nu B_{t}^{(-\nu)}\right)-\exp (2 \nu a)\right)^{+}
$$

This obviously proves $(i)$, by taking $a=(\ln K), \nu=1 / 2$.
b) Proof of $\left(\mathbf{3}_{\mathbf{K}>\mathbf{1}}^{-}\right)$(for $K \geq 1$ )

Using again formula (8), we see that $\left(\mathbf{3}_{\mathbf{K} \geq \mathbf{1}}^{-}\right)$is equivalent to:
$\left(\mathbf{9}_{\mathbf{K}}\right) K P\left(B_{t}-\frac{t}{2}>(\ell n K)\right)=\frac{1}{2} E\left(1_{\left(T_{(\ell n K)}^{(1 / 2)} \leq t\right)} \cdot\left(\frac{K}{\exp \left(B_{t}+\frac{1}{2}\right)} \wedge 1\right)\right)$
We now use the Cameron-Martin relationship on both sides to reduce the statement of $\left(\mathbf{9}_{\mathbf{K}}\right)$ to a statement about standard Brownian motion $\left(B_{t}\right)$, for which we denote: $M_{t}=\sup _{s \leq t} B_{s}$. Thus, we find that $\left(\mathbf{9}_{\mathbf{K}}\right)$ is equivalent to:

$$
\begin{equation*}
K E\left(1_{\left(B_{t}>(\ell n K)\right)} \cdot \exp \left(-\frac{B_{t}}{2}\right)\right) \tag{K}
\end{equation*}
$$

$$
=\frac{1}{2} E\left(1_{\left(M_{t}>(\ell n K)\right)} \cdot\left(\frac{K}{\exp \left(B_{t}\right)} \wedge 1\right) \cdot \exp \left(\frac{B_{t}}{2}\right)\right)
$$

We now decompose the RHS of $\left(\mathbf{1 0}_{\mathbf{K}}\right)$ in a sum of two quantities:

$$
\frac{1}{2}\left\{\begin{array}{l}
E\left[1_{\left(M_{t}>(\ln K)\right)} \cdot 1_{\left(B_{t}>\ln K\right)} \cdot K \exp \left(-\frac{B_{t}}{2}\right)\right] \\
+E\left[1_{\left(M_{t}>\ln K\right)} \cdot 1_{\left(B_{t}<\ln K\right)} \cdot \exp \left(+\frac{B_{t}}{2}\right)\right]
\end{array}\right\}
$$

Thus, $\left(\mathbf{1 0}_{\mathbf{K}}\right)$ now gets simplified to the equivalent form:
(11 ${ }_{\mathrm{K}}$ )

$$
\begin{aligned}
& \frac{K}{2} E\left(1_{\left(B_{t}>(\ln K)\right)} \cdot \exp \left(-\frac{B_{t}}{2}\right)\right) \\
= & \frac{1}{2} E\left(1_{\left(M_{t}>(\ln K)\right)} \cdot 1_{\left(B_{t}<\ln K\right)} \cdot \exp \left(\frac{B_{t}}{2}\right)\right)
\end{aligned}
$$

which, taking $x=(\ell n K)$, may be written as:
$\left(\mathbf{1 2}_{x}\right) \quad E\left(1_{\left(B_{t}>x\right)} \cdot \exp \left(x-\frac{B_{t}}{2}\right)\right)=E\left(1_{\left(M_{t}>x>B_{t}\right)} \cdot \exp \left(\frac{B_{t}}{2}\right)\right)$
We now show $\left(\mathbf{1 2}_{x}\right)$, from the right to the left, as a consequence of the reflection principle:
conditionally on $\mathcal{F}_{T_{x}}$, and $T_{x}<t$, we have:

$$
B_{t}-x=\widehat{B}_{\left(t-T_{x}\right)}, \text { with } \widehat{B} \text { independent from } \mathcal{F}_{T_{x}}
$$

hence, under this condition, the reflection principle boils down to:

$$
B_{t}-x \stackrel{(\text { law })}{=}-\left(B_{t}-x\right)
$$

Thus, the RHS of $\left(\mathbf{1 2}_{x}\right)$ is:

$$
\begin{array}{ll} 
& E\left(1_{\left(T_{x}<t\right)} \cdot 1_{\left(B_{t}-x<0\right)} \cdot \exp \left(\frac{1}{2}\left\{x+\left(B_{t}-x\right)\right\}\right)\right) \\
\underset{\text { (from (w)) }}{=} & E\left(1_{\left(T_{x}<t\right)} \cdot 1_{\left(B_{t}-x<0\right)} \cdot \exp \left(\frac{1}{2}\left\{x-\left(B_{t}-x\right)\right\}\right)\right) \\
= & E\left(1_{\left(B_{t}>x\right)} \cdot \exp \left(x-\frac{B_{t}}{2}\right)\right), \text { which is the LHS of }\left(\mathbf{1 2}_{x}\right)
\end{array}
$$

This proves $\left(\mathbf{3}_{K \geq 1}\right)$, and, with $\left(\mathbf{2}_{K}^{-}\right),\left(\mathbf{2}_{K \geq 1}^{+}\right)$.
c) We now prove that $\left(\mathbf{2}_{K \geq 1}^{+}\right)$implies $\left(\mathbf{2}_{K \leq 1}^{+}\right)$(for $\left.0 \leq K \leq 1\right)$ :

We introduce the probability $P^{\prime}$ such that:

$$
P_{\mathcal{F}_{t}}^{\prime}=\mathcal{E}_{t} \cdot P_{\mathcal{F}_{t}}
$$

We note that, under $P^{\prime}, \frac{1}{\mathcal{E}_{t}}:=\widehat{\mathcal{E}}_{t}=\exp \left(\widehat{B}_{t}-\frac{t}{2}\right)$, for a new Brownian motion $\left(\widehat{B}_{t}, t \geq 0\right)$. Thus, the LHS of $\left(\mathbf{2}_{K \leq 1}^{+}\right)$writes:

$$
\begin{aligned}
& P^{\prime}\left(\mathcal{E}_{t}<K\right)+K P^{\prime}\left(\widehat{\mathcal{E}}_{t}-1_{\left(\mathcal{E}_{t}<K\right)}\right) \\
= & P^{\prime}\left(\widehat{\mathcal{E}}_{t}>\frac{1}{K}\right)+K E^{\prime}\left(1_{\left(\widehat{\mathcal{E}}_{t}>\frac{1}{K}\right)}\right) \\
= & K E^{\prime}\left(1_{\left(\widehat{\mathcal{E}_{t}>\frac{1}{K}}\right)}\right)+\frac{1}{K} P^{\prime}\left(\widehat{\mathcal{E}}_{t}>\frac{1}{K}\right) \\
= & K P\left(T_{(\ell n K)}^{(1 / 2)} \leq t\right) \quad\left(\operatorname{from}\left(2_{\frac{1}{K}}^{+}\right)\right) \\
= & K P\left(T_{(\ell n K)}^{(-1 / 2)} \leq t\right) \quad(\text { by symmetry }) \\
= & P\left(T_{(\ell n K)}^{(1 / 2)} \leq t\right) \quad(\text { from }(6))
\end{aligned}
$$

d) Finally, we observe that $\left(\mathbf{2}_{K \leq 1}^{+}\right)$is equivalent to $\left(2_{K \leq 1}^{++}\right)$, since:

$$
\begin{aligned}
& E\left(\mathcal{E}_{t} 1_{\left(\mathcal{E}_{t}<K\right)}\right)+K P\left(\mathcal{E}_{t}<K\right) \\
= & 1-E\left(\mathcal{E}_{t} 1_{\left(\mathcal{E}_{t}>K\right)}\right)+K\left(1-P\left(\mathcal{E}_{t}>K\right)\right)\left(\text { since } E\left(\mathcal{E}_{t}\right)=1\right) \\
= & 1+K-\left\{E\left(\mathcal{E}_{t} 1_{\left(\mathcal{E}_{t}>K\right)}\right)+K P\left(\mathcal{E}_{t}>K\right)\right\}
\end{aligned}
$$

## 3 On the agreement between the classical BlackScholes formula ( $\mathbf{1}_{\mathrm{K}}^{ \pm}$) and our main result

a) We now check in an elementary manner formulae ( $\mathbf{2}_{\mathbf{K}}^{ \pm}$) by comparing their LHS, as given from the "traditional" Black-Scholes formulae ( $\mathbf{1}_{\mathbf{K}}^{ \pm}$), with their RHS, as given by (4) and (5).
b) The case $K \geq 1$. Since both sides of $\left(\mathbf{2}_{\mathbf{K}}^{-}\right)$and $\left(\mathbf{2}_{K \geq 1}^{+}\right)$are equal to 0 for $t=0$, we need only check that the derivatives in $t$ are equal; thus, our task is to show:
$\left(\mathbf{1 3} \mathbf{K}_{\mathbf{K}}^{-}\right) \quad \frac{\partial}{\partial t}\left\{\mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}+\frac{\sqrt{t}}{2}\right)-K \mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}-\frac{\sqrt{t}}{2}\right)\right\}$

$$
\begin{aligned}
= & \left(\frac{K}{2}\right) \frac{1}{\sqrt{2 \pi t}} \exp \left(-\frac{1}{2 t}\left((\ln K)+\frac{t}{2}\right)^{2}\right) \\
& \frac{\partial}{\partial t}\left\{\mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}+\frac{\sqrt{t}}{2}\right)+K \mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}-\frac{\sqrt{t}}{2}\right)\right\} \\
= & K(\ln K) \frac{1}{\sqrt{2 \pi t^{3}}} \exp \left(-\frac{1}{2 t}\left((\ln K)+\frac{t}{2}\right)^{2}\right)
\end{aligned}
$$

$\left(13_{\mathrm{K} \geq 1}^{+}\right)$
(We also see on these expressions the relationships between $p_{t}^{(-1 / 2)}(x)$ and $\left.p_{t}^{(+1 / 2)}(x)\right)$.
To prove $\left(\mathbf{1 3}_{\mathbf{K}}^{-}\right)$and $\left(\mathbf{1 3}_{\mathbf{K} \geq \mathbf{1}}^{+}\right)$, we compute:

- $\quad \frac{\partial}{\partial t}\left(\mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}+\frac{\sqrt{t}}{2}\right)\right)$

$$
=\frac{\sqrt{K}}{\sqrt{2 \pi}} \exp \left(-\frac{1}{2}\left(\frac{(\ln K)^{2}}{t}+\frac{t}{4}\right)\right)\left(\frac{\partial}{\partial t}\left\{-\frac{(\ln K)}{\sqrt{t}}+\frac{\sqrt{t}}{2}\right\}\right)
$$

- $K \frac{\partial}{\partial t}\left(\mathcal{N}\left(-\frac{\ln K}{\sqrt{t}}-\frac{\sqrt{t}}{2}\right)\right)$

$$
=\frac{\sqrt{K}}{\sqrt{2 \pi}} \exp \left(-\frac{1}{2}\left(\frac{(\ln K)^{2}}{t}+\frac{t}{4}\right)\right)\left(\frac{\partial}{\partial t}\left\{-\frac{(\ln K)}{\sqrt{t}}-\frac{\sqrt{t}}{2}\right\}\right)
$$

and $\left(13_{\mathrm{K}}^{-}\right)$and $\left(\mathbf{1 3}_{\mathrm{K} \geq 1}^{+}\right)$are then obtained by elementary algebraic manipulations. (In fact, it is these very manipulations which led us to believe in the truth of Theorem 1!!).
c) The case $0 \leq K \leq 1$. Since both sides of $\left(\mathbf{2}_{K \leq 1}^{++}\right)$are equal to $1+K$ for $t=0$, it suffices to prove, for $K \leq 1$ :

$$
\begin{aligned}
& \frac{\partial}{\partial t}\left\{\mathcal{N}\left(-\frac{\log K}{\sqrt{t}}+\frac{\sqrt{t}}{2}\right)+K \mathcal{N}\left(-\frac{\log K}{\sqrt{t}}-\frac{\sqrt{t}}{2}\right)\right\} \\
= & \frac{\partial}{\partial t} P\left(T_{(\ln K)}^{(1 / 2)}>t\right)
\end{aligned}
$$

and this latter relation follows immediately from the computations done in point b) above and from (7).
Note also that $\left(\mathbf{2}_{K \geq 1}^{+}\right)$and $\left(\mathbf{2}_{K \leq 1}^{+}\right)$coincide for $K=1$.

## 4 Further remarks and conclusion

a) Easy variants of formulae ( $2_{\mathrm{K}}^{ \pm}$) may be written, e.g. by using the scaling property of Brownian motion, so that $T_{a}^{(\nu)}$ and $G_{a}^{(\nu)}$ appear on the RHS of $\left(2_{\mathbf{K}}^{ \pm}\right)$; however, writing down these variants would only complicate unnecessarily these formulae.
b) We recall that, as a consequence of the time inversion property of Brownian motion, there are the relations:

$$
\left(T_{a}^{(\nu)}, G_{a}^{(\nu)}\right) \stackrel{(\mathrm{law})}{=}\left(\frac{1}{G_{\nu}^{(a)}}, \frac{1}{T_{\nu}^{(a)}}\right)
$$

(see e.g., Pitman-Yor [3] for a more general discussion).
In particular, for $a=0$, one gets:

$$
G_{0}^{(\nu)} \stackrel{(\text { law })}{=} \frac{1}{T_{\nu}^{(0)}} \stackrel{(\text { law })}{=} \frac{B_{1}^{2}}{\nu^{2}}
$$

In particular, formula ( $\mathbf{2}_{\mathbf{K}=\mathbf{1}}^{-}$) becomes:

$$
\begin{equation*}
E\left[\left(\mathcal{E}_{t}-1\right)^{+}\right]=E\left[\left(\mathcal{E}_{t}-1\right)^{-}\right]=P\left(4 B_{1}^{2} \leq t\right) \tag{1}
\end{equation*}
$$

which allowed us to answer M. Qian's question [4]:
is there a simple formula for:

$$
\int_{0}^{\infty} \theta(d t) E\left[\left(\mathcal{E}_{t}-1\right)^{ \pm}\right]
$$

where $\theta(d t)$ is a probability on $\mathbb{R}_{+}$?
From ( $\mathbf{1 4}_{1}$ ), we easily obtain:

$$
\begin{equation*}
\int_{0}^{\infty} \theta(d t) E\left[\left(\mathcal{E}_{t}-1\right)^{ \pm}\right]=E\left[\bar{\theta}\left(4 B_{1}^{2}\right)\right] \tag{15}
\end{equation*}
$$

where $\bar{\theta}(x)=\theta([x, \infty))$ is the tail of $\theta$.
To particularise even more, we give the explicit form of the Laplace transform:

$$
\begin{align*}
\int_{0}^{\infty} d t e^{-\lambda t} E\left[\left(\mathcal{E}_{t}-1\right)^{ \pm}\right] & =\frac{1}{\lambda} E\left[\exp \left(-\lambda\left(4 B_{1}^{2}\right)\right)\right]  \tag{16}\\
& =\frac{1}{\lambda} \frac{1}{\sqrt{1+8 \lambda}}
\end{align*}
$$

It is this question which set us on the general quest for a representation of

$$
E\left[\left(\mathcal{E}_{t}-K\right)^{+}\right]
$$

as a cumulative distribution function in $t$.

- c) We come back to the time inversion property of BM, in order to throw another light upon our main result $\left(\mathbf{2}_{K}^{-}\right)$, which relates the European call price with the cumulative function of last Brownian passage times. (This paragraph has been partly inspired by unpublished notes by Peter Carr [2].) Indeed, a variant of $\left(\mathbf{2}_{K}^{-}\right)$is the following: for every $t \geq 0, K \geq 0$, and $\phi: C([0, t]) \rightarrow \mathbb{R}_{+}$, measurable,

$$
\begin{equation*}
E\left[\phi\left(B_{u}, u \leq t\right)\left(K-\mathcal{E}_{t}\right)^{+}\right]=K E\left[\phi\left(B_{u}, u \leq t\right) 1_{\left(\mathcal{G}_{K} \leq t\right)}\right] \tag{17}
\end{equation*}
$$

where $\mathcal{G}_{K}=\sup \left\{u: \mathcal{E}_{u}=K\right\}$.
Writing (17) in terms of the Brownian motion ( $\widehat{B}_{v}, v \geq 0$ ) such that: $B_{u}=$ $u \widehat{B}_{(1 / u)}$, and setting $s=1 / t$, it is clearly seen that $(\mathbf{1 7})$ is equivalent to:

$$
K P\left(\widehat{T}_{1 / 2}^{(-\ell n K)} \geq s \mid \widehat{B}_{s}\right)=\left(K-\exp \left(\frac{1}{s} \widehat{B}_{s}-\frac{1}{2 s}\right)\right)^{+}
$$

where: $\widehat{T}_{a}^{(\nu)}=\inf \left\{u: \widehat{B}_{u}+\nu u=a\right\}$.
Since hats are no longer necessary for our purpose, we drop them, and we now look for an independent proof of:

$$
\begin{equation*}
P\left(T_{1 / 2}^{(-\ell n K)} \geq s \mid B_{s}=x\right)=\left(1-\frac{1}{K} \exp \left(\frac{x}{s}-\frac{1}{2 s}\right)\right)^{+} \tag{18}
\end{equation*}
$$

On the LHS of (18), we may replace $\left(B_{s}=x\right)$ by $\left(B_{s}-s(\ln K)=x-s(\ell n K)\right)$. Now, as a consequence of the Cameron-Martin relationship, the conditional expectation:

$$
E\left[F\left(B_{u}-\nu u, u \leq s\right) \mid B_{s}-\nu s=y\right]
$$

does not depend on $\nu$; hence, (18) is equivalent to:

$$
P\left(T_{1 / 2}^{(-\ell n K)} \geq s \mid B_{s}=x-s(\ell n K)\right)=\left(1-\frac{1}{K} \exp \left(\frac{x-\frac{1}{2}}{s}\right)\right)^{+}
$$

which simplifies to:

$$
P\left(\left.\sup _{u \leq s} B_{u}<\frac{1}{2} \right\rvert\, B_{s}=y\right)=\left(1-\exp \left(\frac{y-\frac{1}{2}}{s}\right)\right)^{+}
$$

or, by scaling:

$$
P\left(\left.\sup _{u \leq 1} B_{u}<\frac{1}{2 \sqrt{s}} \right\rvert\, B_{1}=\frac{y}{\sqrt{s}}\right)=\left(1-\exp \left(\frac{1}{\sqrt{s}}\left(\frac{y}{\sqrt{s}}-\frac{1}{2 \sqrt{s}}\right)\right)\right)^{+}
$$

This is equivalent to:

$$
\begin{equation*}
P\left(\sup _{u \leq 1} B_{u}<\sigma \mid B_{1}=y\right)=\left(1-\exp (2 \sigma(y-\sigma))^{+}\right. \tag{19}
\end{equation*}
$$

for $\sigma \geq 0$, and $y \in \mathbb{R}$.
This formula is trivial for $\sigma<y$, and, for $\sigma \geq y$, it follows from the classical formula:

$$
\begin{equation*}
P\left(\sup _{u \leq 1} B_{u} \in d \sigma, B_{1} \in d a\right)=\frac{d a d \sigma}{\sqrt{2 \pi}} 2(2 \sigma-a) e^{-\frac{(2 \sigma-a)^{2}}{2}} 1_{\{a<\sigma ; \sigma \geq 0\}} \tag{20}
\end{equation*}
$$

d) In a future work, we plan to study more generally how quantities such as the calls and puts:

$$
E\left[\left(S_{t}-K\right)^{+}\right] \text {and } E\left[\left(S_{t}-K\right)^{-}\right]
$$

associated with a general $\mathbb{R}_{+}$-valued continuous local martingale $\left(S_{t}, t \geq 0\right)$ may be written in terms of cumulative functions.
e) In [1], the authors present eight different approaches to the BlackScholes formula, among which the change of numéraire approach (Section 5 of [1]), and the local time approach (Section 6 of [1]). This local time approach, together with $\left(\mathbf{2}_{\mathbf{K}}^{-}\right)$yields the relationship:

$$
\begin{equation*}
P\left(G_{(\ell n K)}^{(1 / 2)} \leq t\right)=\left(\frac{1}{2}\right) E\left[\mathcal{L}_{t}^{K}\right] \tag{21}
\end{equation*}
$$

where $\left(\mathcal{L}_{t}^{K}, t \geq 0\right)$ denotes the local time at level $K$ of $\left(\mathcal{E}_{t}, t \geq 0\right)$. It is this kind of relationship (21) which is central in the obtention in [3] of a general expression for the law of a last passage time of a transient diffusion. However, to our knowledge, despite the remarkable survey [1] of methods leading to the Black-Scholes formula, no interpretation of this formula seems to have been made in terms of last passage times distributions.

## References

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