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WEAK SOLUTIONS FOR DISLOCATION TYPE EQUATIONS

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Abstract. We describe recent results obtained by G. Barles, P. Cardaliaguet, R. Monneau and the author in [9, 7]. They are concerned with nonlocal Eikonal equations arising in the study of the dynamics of dislocation lines in crystals. These equations are nonlocal but also non monotone. We use a notion of weak solution to provide solutions for all time. Then, we discuss the link between these weak solutions and the classical viscosity solutions, and state some uniqueness results in particular cases. A counter-example to uniqueness is given.

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1 Introduction

It is a great honor to contribute to this proceedings of the Conference for the 25th Anniversary of Viscosity Solution and the Celebration of the 60th birthday of Professor Hitoshi Ishii.

In this proceedings, we describe recent results [9, 7] obtained by the author in collaboration with G. Barles, P. Cardaliaguet and R. Monneau for first-order nonlocal Hamilton-Jacobi modelling the dynamics of dislocations.

Dislocations are defects in crystals of typical length $10^{-6}m$ and the dynamics of dislocations is the main microscopic explanation of the macroscopic behaviour of metallic crystals. For details about the physics of dislocations, see for instance Nabarro [26] or Hirth and Lothe [23]. We are interested in a particular model introduced in Rodney, Le Bouar and Finel [30]; the dislocation line evolves in a plane called slip plane, with a normal velocity proportional to the Peach-Koehler force acting on this line. This Peach-Koehler force have two contributions. The first one is the self-force created by the elastic field generated by the dislocation line itself (i.e. this self-force is a nonlocal function of the shape of the dislocation line). The second one is due to exterior forces (like an exterior stress applied on the material for instance).

More precisely, we study the evolution of a dislocation line Γ_t which is, at any time $t \geq 0$, the boundary of an open bounded set $\Omega_t \subset \mathbb{R}^N$ (with $N = 2$ for the physical application). The normal velocity, at each point $x \in \Gamma_t = \partial\Omega_t$ of the dislocation line, is given by

$$V_n = c_0 \star \mathbb{1}_{\overline{\Omega}_t} + c_1 \quad (1)$$

where $\mathbb{1}_{\overline{\Omega}_t}(x)$ is the indicator function of the set $\overline{\Omega}_t$. The function $c_0(x, t)$ is a kernel which only depends on the physical properties of the crystal. In the special case of the study of dislocations, the kernel c_0 does not depend on time, but to keep a general setting we allow here a dependence on the time variable. Here \star denotes the convolution in space, namely

$$(c_0(\cdot, t) \star \mathbb{1}_{\overline{\Omega}_t})(x) = \int_{\mathbb{R}^N} c_0(x - y, t) \mathbb{1}_{\overline{\Omega}_t}(y) dy, \quad (2)$$

and this term appears to be the Peach-Koehler self-force created by the dislocation itself, while $c_1(x, t)$ is the exterior contribution to the velocity, created by everything exterior to the dislocation line. We refer to Alvarez, Hoch, Le Bouar and Monneau [4] for a detailed presentation and a derivation of this model.

Using the level-set approach to front propagation problems, we can derive a partial differential equation to represent the evolution of Γ_t . The level-set approach was introduced by Osher and Sethian [29], and then developed first by Chen, Giga and Goto [17], and Evans and Spruck [20]. This approach produced a lot of applications and now there is a huge literature; see the monograph of Giga [21] for details.

The level-set approach consists in replacing the evolution of the set Γ_t by the evolution of the zero level-set of an auxiliary function u . More precisely, given a set Γ_0 (the dislocation line at time $t = 0$) and a bounded uniformly continuous function $u_0 : \mathbb{R}^N \rightarrow \mathbb{R}$ such that

$$\{u_0 = 0\} = \Gamma_0 \quad \text{and} \quad \{u_0 > 0\} = \Omega_0 \quad (3)$$

(u_0 represents the initial dislocation line), we are looking for a function $u : \mathbb{R}^N \times [0, T] \rightarrow \mathbb{R}$ which satisfies

$$\{u(\cdot, t) = 0\} = \Gamma_t \quad \text{and} \quad \{u(\cdot, t) > 0\} = \Omega_t \quad \text{for all } t \geq 0. \quad (4)$$

The function u has to satisfy the level-set equation (see [21]) which reads here

$$\begin{cases} \frac{\partial u}{\partial t} = (c_0(\cdot, t) \star \mathbb{1}_{\{u(\cdot, t) \geq 0\}}(x) + c_1(x, t))|Du| & \text{in } \mathbb{R}^N \times (0, T) \\ u(\cdot, 0) = u_0 & \text{in } \mathbb{R}^N, \end{cases} \quad (5)$$

where $\frac{\partial u}{\partial t}$, Du and $|\cdot|$ denote respectively the time and the spatial derivative of u , and the Euclidean norm. Note that (2) now reads

$$c_0(\cdot, t) \star \mathbb{1}_{\{u(\cdot, t) \geq 0\}}(x) = \int_{\mathbb{R}^N} c_0(x - y, t) \mathbb{1}_{\{u(\cdot, t) \geq 0\}}(y) dy. \quad (6)$$

Note that (5) is not really a level-set equation because it is not invariant by increasing changes of functions. In order to have rigorously a level-set equation, the nonlocal term $\{u(\cdot, t) \geq 0\}$ should be replaced by $\{u(\cdot, t) \geq u(x, t)\}$ (see Slepčev [31]). But here, (5) is the equation we are interested in.

The study of Equation (5) raises three main difficulties: the first one is the presence of the nonlocal term (6).

The second difficulty is the weak regularity in time of the equation. Indeed, as soon as $\{u(\cdot, t) = 0\}$ develops an interior (fattening phenomenon), the map $t \mapsto c_0(\cdot, t) \star \mathbb{1}_{\{u(\cdot, t) \geq 0\}}(x)$ is no longer continuous and we have to deal with (5) which is an equation with measurable-in-time coefficients. The study of such equations was initiated by Ishii [24] (see the Appendix).

The third difficulty, which is more involved, is a lack of monotonicity for (5). In many cases, proofs of existence and uniqueness for such geometrical equations rely on the preservation of inclusion property which can be stated as follows. Consider a front propagation problem (see (4)) with a given normal velocity. Let Γ_0 and $\tilde{\Gamma}_0$ be two different initial fronts evolving independently. Then,

$$\overline{\Omega}_0 \subset \text{int}(\tilde{\Omega}_0) \quad \implies \quad \overline{\Omega}_t \subset \text{int}(\tilde{\Omega}_t) \quad \text{for all time } t \geq 0. \quad (7)$$

Such a property is the key point to use the classical viscosity solutions' theory. For instance, it is satisfied for local evolution problems as propagation by constant normal velocities, mean curvature flow (see [21]) or for some nonlocal problems as in Cardaliaguet [14, 15], Dalio, Kim and Slepčev [19], Srour [33], etc. But, for dislocation dynamics, the kernel c_0 has a zero mean which implies that it changes sign. Therefore, the preservation of inclusion property is not true in general. It follows that we cannot expect a principle of comparison (that is: the subsolutions of (5) are below the supersolutions).

For geometrical evolutions without preservation of inclusion, few results are known, see however Giga, Goto and Ishii [22], Soravia and Souganidis [32] and Alibaud [1]. In the case of (5), under suitable assumptions on c_0, c_1 (see **(H1)**-**(H2)**) and on the initial data, the existence and the uniqueness of the solution were proved first for short time in [3, 4]. In [2, 9, 16], such results were proved for all time under the additional assumption that $V_n \geq 0$, which is for instance always satisfied for c_1 satisfying $c_1(x, t) \geq |c_0(\cdot, t)|_{L^1(\mathbb{R}^N)}$. In the general case, a notion of weak solutions was introduced in [7].

The aim of this paper is to describe global-in-time results obtained in [2, 9, 7]. In Section 2, we define the weak solutions and prove an existence theorem. In Section 3, we state some uniqueness results. Section 4 is devoted to the study of a counter-example to uniqueness. Finally, we recall the definition of L^1 -viscosity solutions and a new stability result proved by Barles [6] in the Appendix.

2 Definition and existence of weak solutions

We introduce the following notion of weak solutions for (5):

Definition 2.1 (Classical and weak solutions) [7]

For any $T > 0$, we say that a Lipschitz continuous function $u : \mathbb{R}^N \times [0, T] \rightarrow \mathbb{R}$ is a weak solution of equation (5) on the time interval $[0, T]$, if there is some measurable map $\chi : \mathbb{R}^N \times (0, T) \rightarrow [0, 1]$ such that u is a L^1 -viscosity solution of

$$\begin{cases} \frac{\partial u}{\partial t} = \bar{c}(x, t)|Du| & \text{in } \mathbb{R}^N \times (0, T) \\ u(\cdot, 0) = u_0 & \text{in } \mathbb{R}^N, \end{cases} \quad (8)$$

where

$$\bar{c}(x, t) = c_0(\cdot, t) \star \chi(\cdot, t)(x) + c_1(x, t) \quad (9)$$

and

$$\mathbb{1}_{\{u(\cdot, t) > 0\}}(x) \leq \chi(x, t) \leq \mathbb{1}_{\{u(\cdot, t) \geq 0\}}(x), \quad (10)$$

for almost all $(x, t) \in \mathbb{R}^N \times [0, T]$. We say that u is a classical solution of equation (5) if u is a weak solution to (8) and if

$$\mathbb{1}_{\{u(\cdot, t) > 0\}}(x) = \mathbb{1}_{\{u(\cdot, t) \geq 0\}}(x) \quad (11)$$

for almost all $(x, t) \in \mathbb{R}^N \times [0, T]$.

We recall that L^1 -viscosity solutions were introduced by Ishii [24], see the appendix for details. Note that, for classical solutions, we have $\chi(x, t) = \mathbb{1}_{\{u(\cdot, t) > 0\}}(x) = \mathbb{1}_{\{u(\cdot, t) \geq 0\}}(x)$ for almost all $(x, t) \in \mathbb{R}^N \times [0, T]$.

To state our first existence result, we introduce the following assumptions

(H0) $u_0 : \mathbb{R}^N \rightarrow [-1, 1]$ is Lipschitz continuous and there exists $R_0 > 0$ such that $u_0(x) \equiv -1$ for $|x| \geq R_0$,

(H1) $c_0 \in C([0, T]; L^1(\mathbb{R}^N))$, $D_x c_0 \in L^\infty([0, T]; L^1(\mathbb{R}^N))$, $c_1 \in C(\mathbb{R}^N \times [0, T])$ and there exists constants M_1, L_1 such that, for any $x, y \in \mathbb{R}^N$ and $t \in [0, T]$

$$|c_1(x, t)| \leq M_1 \quad \text{and} \quad |c_1(x, t) - c_1(y, t)| \leq L_1|x - y|. \quad (12)$$

Let us make some comments about these assumptions. The role of u_0 is to represent the initial dislocation Γ_0 which lies in a bounded region (see (3)). In general, we choose u_0 as a truncation of the signed distance to Γ_0 (positive in Ω_0). Such a function u_0 is Lipschitz continuous and satisfies **(H0)**. Note that we do not impose any sign condition on c_0 in

(H1). In the sequel, we denote by M_0, L_0 some constants such that, for any (or almost every) $t \in [0, T)$, we have

$$|c_0(\cdot, t)|_{L^1(\mathbb{R}^N)} \leq M_0 \quad \text{and} \quad |D_x c_0(\cdot, t)|_{L^1(\mathbb{R}^N)} \leq L_0. \quad (13)$$

Our first main result is the following.

Theorem 2.2 (Existence of weak solutions) [7]

Under assumptions **(H0)**-**(H1)**, for any $T > 0$ and for any initial data u_0 , there exists a weak solution of equation (5) on the time interval $[0, T]$.

We give only the main ideas of the proof of Theorem 2.2. The whole proof can be found in [7] and an alternative proof is presented in [8].

Sketch of proof of Theorem 2.2.

1. *Introduction of a perturbed equation.* We consider the equation

$$\frac{\partial u^\varepsilon}{\partial t} = c_\varepsilon[u^\varepsilon](x, t)|Du^\varepsilon| \quad \text{in } \mathbb{R}^N \times (0, T), \quad (14)$$

where the unknown is u^ε ,

$$c_\varepsilon[u] = (c_0(\cdot, t) \star \psi_\varepsilon(u(\cdot, t)))(x) + c_1(x, t) \quad \text{for any } u : \mathbb{R}^N \times [0, T] \rightarrow \mathbb{R},$$

and $\psi_\varepsilon : \mathbb{R} \rightarrow \mathbb{R}$ is a sequence of continuous functions such that $\psi_\varepsilon(r) \equiv 0$ for $r \leq -\varepsilon$, $\psi_\varepsilon(r) \equiv 1$ for $r \geq 0$ and ψ_ε is an affine function on $[-\varepsilon, 0]$.

2. *Definition of a map \mathcal{T} .* Let

$$X = \{u \in C(\mathbb{R}^N \times [0, T]) : u \equiv -1 \text{ in } \mathbb{R}^N \setminus B(0, R_0 + MT), \\ |Du|, \left| \frac{\partial u}{\partial t} \right| / M \leq |Du_0|_{L^\infty(\mathbb{R}^N)} e^{LT}\},$$

where $M = M_0 + M_1$ and $L = L_0 + L_1$ (see (12) and (13) for the definition of M_0, M_1, L_0, L_1). By Ascoli's Theorem, X is a compact and convex subset of $(C(\mathbb{R}^N \times [0, T]), |\cdot|_\infty)$. We define the map $\mathcal{T} : X \rightarrow X$ by : if $u \in C(\mathbb{R}^N \times [0, T])$, then $u^\varepsilon := \mathcal{T}(u)$ is the unique solution of (14) with $c_\varepsilon[u]$ (instead of $c_\varepsilon[u^\varepsilon]$). The existence and uniqueness of u^ε come from classical results for Eikonal equations with finite speed propagation property (see [7, Theorem 2.1], Crandall & Lions [18], [25] and [9]) since, under assumption **(H1)** on c_1 and c_0 , $c_\varepsilon[u]$ satisfies **(H1)** with fixed constants M and L .

3. *Application of Schauder's fixed point theorem to \mathcal{T} .* The map \mathcal{T} is continuous since ψ_ε is continuous, by using the classical stability result for viscosity solutions (see Barles [5]). Therefore, \mathcal{T} has a fixed point u_ε which is bounded in $W^{1,\infty}(\mathbb{R}^N \times [0, T])$ uniformly with respect to ε (since M and L are independent of ε).

4. *Convergence of the fixed point when $\varepsilon \rightarrow 0$.* From Ascoli's Theorem, we extract a subsequence $(u_{\varepsilon'})_{\varepsilon'}$ which converges locally uniformly to a function denoted by u . The functions $\chi_{\varepsilon'} := \psi_{\varepsilon'}(u_{\varepsilon'})$ satisfy $0 \leq \chi_{\varepsilon'} \leq 1$. Therefore, we can extract a subsequence—still denoted $(\chi_{\varepsilon'})$ —which converges weakly-* in $L^\infty_{\text{loc}}(\mathbb{R}^N \times [0, T])$ to some function χ :

$\mathbb{R}^N \times (0, T) \rightarrow [0, 1]$. Furthermore, setting $c_{\varepsilon'} = c_0 \star \chi_{\varepsilon'} + c_1$, we have, for all $(x, t) \in \mathbb{R}^N \times [0, T]$,

$$\begin{aligned} \int_0^t c_{\varepsilon'}(x, s) ds &= \int_0^t \int_{\mathbb{R}^N} c_0(x - y, s) \chi_{\varepsilon'}(y, s) dy ds + \int_0^t c_1(x, s) ds \\ &\rightarrow \int_0^t \bar{c}(x, s) ds, \end{aligned}$$

where $\bar{c}(x, t) = c_0(\cdot, t) \star \chi(\cdot, t)(x) + c_1(x, t)$. The above convergence is pointwise but, noticing that $c_{\varepsilon'}$ is bounded Lipschitz continuous in space uniformly in time and measurable in time, we can apply the stability Theorem 4.3 of Barles [6] for weak convergence in time. We conclude that u is L^1 -viscosity solution to (8) with \bar{c} satisfying (9)-(10). \square

3 Classical solutions and uniqueness results

Our second main result gives a sufficient condition for a weak solution to be a classical one.

Theorem 3.1 (Links between weak solutions and classical continuous viscosity solutions) [7]

Assume **(H0)**-**(H1)** and suppose that there is some $\delta \geq 0$ such that, for all measurable map $\chi : \mathbb{R}^N \times (0, T) \rightarrow [0, 1]$,

$$\text{for all } (x, t) \in \mathbb{R}^N \times [0, T], \quad c_0(\cdot, t) \star \chi(\cdot, t)(x) + c_1(x, t) \geq \delta, \quad (15)$$

and that the initial data u_0 satisfies (in the viscosity sense)

$$-|u_0| - |Du_0| \leq -\eta_0 \quad \text{in } \mathbb{R}^N \quad (16)$$

for some $\eta_0 > 0$. Then any weak solution u of (5) in the sense of Definition 2.1, is a classical continuous viscosity solution of (5).

Assumption (15) ensures that the velocity V_n in (1) is nonnegative, i.e. the dislocation line is expanding. Of course, we can state similar results in the case of negative velocity for shrinking dislocation lines. Assumption (16) comes from [25]. It means that u_0 is a viscosity subsolution of $-|v(x)| - |Dv(x)| + \eta_0 \leq 0$. It can be seen as a nonsmooth generalization of the following situation: if u_0 is C^1 , (16) implies that the gradient of u_0 does not vanish on the set $\{u_0 = 0\}$ and therefore this latter set is a C^1 hypersurface.

Sketch of proof of Theorem 3.1. At first, if u is a weak solution and \bar{c} is associated with u , then, from (15), for any $x \in \mathbb{R}^N$ and for almost all $t \in [0, T]$, we have

$$\bar{c}(x, t) \geq \delta \geq 0$$

and therefore the Hamiltonian $\bar{c}(x, t)|Du|$ of (5) is convex 1-homogeneous in the gradient variable. Then, the conclusion is a consequence of a preservation of the lower-bound gradient estimate (16) proved in [25, Theorem 4.2] for equations with convex Hamiltonians H (such that $H(x, t, \lambda p) = \lambda H(x, t, p)$ for all $\lambda \geq 0$): there exists $\eta(T) > 0$ such that

$$-|u(\cdot, t)| - |Du(\cdot, t)| \leq -\eta(T) \quad \text{on } \mathbb{R}^N \times (0, T). \quad (17)$$

It follows that for every $t \in (0, T)$, the 0-level-set of $u(\cdot, t)$ has a zero Lebesgue measure and therefore (11) holds. Moreover $t \mapsto \mathbb{1}_{\{u(\cdot, t) \geq 0\}}$ is also continuous in L^1 , and then \bar{c} is continuous. \square

Let us turn to uniqueness results. If the evolving set has positive velocity or if the velocity is nonnegative and the following additional condition is fulfilled, then we can prove uniqueness results.

(H2) c_1 and c_0 satisfy **(H1)** and there exists constants m_0, N_1 and a positive function $N_0 \in L^1(\mathbb{R}^N)$ such that, for any $x, h \in \mathbb{R}^N, t \in [0, T)$, we have

$$\begin{aligned} |c_0(x, t)| &\leq m_0, \\ c_1(x + h, t) + c_1(x - h, t) - 2c_1(x, t) &\geq -N_1|h|^2, \\ c_0(x + h, t) + c_0(x - h, t) - 2c_0(x, t) &\geq -N_0(x)|h|^2. \end{aligned}$$

Second and third conditions means that c_0 and c_1 are semiconvex in space.

Theorem 3.2 (Uniqueness results) [2, 9, 7]

Assume **(H0)**-**(H1)**-**(H2)** and suppose that (15) and (16) hold. The solution of (5) is unique if

(i) either $\delta = 0$ and u_0 is semiconvex, i.e. satisfies for some constant $C > 0$:

$$u_0(x + h) + u_0(x - h) - 2u_0(x) \geq -C|h|^2, \quad \forall x, h \in \mathbb{R}^N;$$

(ii) or $\delta > 0$.

Even if it has no physical meaning in the theory of dislocations, an important particular case of application of Theorem 3.1 is the uniqueness for (5) when $c_0 \geq 0$ and $c_1 \equiv 0$ (this implies (16)). In this case, the preservation of inclusion property (7) holds and some classical results apply, see Cardaliaguet [14] and [7, Theorem 1.5]. But let us point out that a nonnegative kernel c_0 does not ensure uniqueness in general, see the counter-example in Section 4.

Point (i) of the theorem is the main result of [2, 9]. Let us compare the two articles. In [2], it is proved that we have uniqueness for (5) if we start with an initial dislocation $\Gamma_0 = \partial\Omega_0$ such that Ω_0 has the interior ball property of radius $r > 0$ that is: for any $x \in \bar{\Omega}_0$, there exists $p \in \mathbb{R}^N \setminus \{0\}$ such that $\bar{B}(x - r\frac{p}{|p|}, r) \subset \bar{\Omega}_0$. In [9], uniqueness is proved under the assumption that u_0 is semiconvex and satisfies the lower-bound gradient (16). This latter set of assumptions is equivalent to the interior ball property for $\{u_0 \geq 0\}$ (see [9, Lemma A.1]).

Sketch of proof of Theorem 3.2.

1. Part (i) *Definition of a map \mathcal{F} .* We follow the ideas of [9] and refer to this paper for details. The proof relies on the Banach contraction fixed point theorem. Let

$$\begin{aligned} Y = \{ \chi \in C([0, T], L^1(\mathbb{R}^N)) : \\ 0 \leq \chi \leq 1, |\chi(\cdot, t)|_{L^1(\mathbb{R}^N)} \leq \mathcal{L}^N(B(0, R_0 + MT)) \}, \end{aligned}$$

where $M = M_0 + M_1$ (see (12) and (13) for the definition of M_0, M_1), \mathcal{L}^N is the Lebesgue measure in \mathbb{R}^N and $B(0, R)$ is the open ball of center 0 and radius $R > 0$. For $\tau > 0$ fixed, the set Y is endowed with the norm

$$|\chi|_{Y, \tau} = \sup_{t \in [0, \tau]} |\chi(\cdot, t)|_{L^1(\mathbb{R}^N)}$$

Define $\mathcal{F} : Y \rightarrow Y$ by: for all $\chi \in Y$, $\mathcal{F}(\chi) = \mathbb{1}_{u(\cdot, t) \geq 0}$ where u is the unique continuous viscosity solution of

$$\begin{cases} \frac{\partial u}{\partial t} = c[\chi](x, t)|Du| & \text{in } \mathbb{R}^N \times (0, T) \\ u(\cdot, 0) = u_0 & \text{in } \mathbb{R}^N, \end{cases} \quad (18)$$

where $c[\chi] = c_0(\cdot, t) \star \chi(\cdot, t)(x) + c_1(x, t)$. We have to check that \mathcal{F} is well-defined.

2. *Part (i) The map \mathcal{F} is well defined.* From **(H1)**-**(H2)**, for all $\chi \in Y$, the map $(x, t) \in \mathbb{R}^N \times [0, T] \mapsto c[\chi](x, t)$ is bounded continuous in $\mathbb{R}^N \times [0, T]$, Lipschitz continuous and semiconvex in x (uniformly with respect to t) with some constants which depends only on the given data $M_0, M_1, L_0, L_1, R_0, \eta(T), C$. It follows that for all Lipschitz continuous u_0 , (18) has a unique Lipschitz continuous viscosity solution u . Next, if u_0 satisfies **(H0)**, then, by the finite speed of propagation property, for all $t \geq 0$, $\{u(\cdot, t) \geq 0\} \subset B(0, R_0 + Mt)$. Let us give a geometrical interpretation of this latter property: by (15), Equation (18) is monotone and the preservation inclusion principle (7) holds. Noticing that M is an upper bound for the speed of propagation of the 0-level-set of u and that $B(0, R_0 + Mt)$ is the propagation of the ball $B(0, R_0)$ with normal velocity M , by preservation of inclusion, the property follows.

3. *Part (i) The map \mathcal{F} is continuous.* It comes from the continuity of the map $t \in [0, T] \mapsto \int_{\mathbb{R}^N} \mathbb{1}_{u(\cdot, t) \geq 0}(x) dx$. The proof of this result is an immediate consequence of the preservation of the lower-bound gradient estimate (16)-(17) (see the proof of Theorem 3.1).

4. *Part (i) Contraction property for \mathcal{F} (beginning of the calculation).* Let $\chi_1, \chi_2 \in Y$ and u_1, u_2 be the solution of (18) with $c[\chi_1]$ and $c[\chi_2]$ respectively. Set

$$\rho := \sup_{t \in [0, \tau]} |(u_1 - u_2)(\cdot, t)|_{L^\infty(\mathbb{R}^N)}. \quad (19)$$

(note that $\rho \rightarrow 0$ as $\tau \rightarrow 0$ since $u_1(\cdot, 0) = u_2(\cdot, 0) = u_0$). For all $t \in [0, T]$, a straightforward computation leads to

$$\begin{aligned} & |(\mathcal{F}(\chi_1) - \mathcal{F}(\chi_2))(\cdot, t)|_{L^1(\mathbb{R}^N)} \\ &= |\mathbb{1}_{u_1(\cdot, t) \geq 0} - \mathbb{1}_{u_2(\cdot, t) \geq 0}|_{L^1(\mathbb{R}^N)} \\ &\leq \mathcal{L}^N(\{u_1(\cdot, t) \geq 0, u_2(\cdot, t) < 0\}) + \mathcal{L}^N(\{u_2(\cdot, t) \geq 0, u_1(\cdot, t) < 0\}) \\ &\leq \mathcal{L}^N(\{-\rho \leq u_2(\cdot, t) < 0\}) + \mathcal{L}^N(\{-\rho \leq u_1(\cdot, t) < 0\}). \end{aligned} \quad (20)$$

5. *Part (i) Contraction property for \mathcal{F} (L^1 -estimates).* The estimate of the last two terms in (20) are based on some fundamental L^1 -estimates obtained in [9]: let φ_ε be a smooth

approximation of $\mathbb{1}_{[-\rho,0]}$ (with $\mathbb{1}_{[-\rho,0]} \leq \varphi_\varepsilon \leq \mathbb{1}_{[-\rho-\varepsilon,\varepsilon]}$) and $0 < \rho < \eta(T)/2$ (where $\eta(T)$ is given by (17)). Then, there exists $K > 0$ such that

$$\int_{\mathbb{R}^N} \varphi_\varepsilon(u_2(x, t)) dx \leq e^{Kt} \int_{\mathbb{R}^N} \varphi_\varepsilon(u_0(x)) dx \quad (21)$$

which implies by sending $\varepsilon \rightarrow 0$,

$$\mathcal{L}^N(\{-\rho \leq u_2(\cdot, t) < 0\}) \leq e^{Kt} \mathcal{L}^N(\{-\rho \leq u_0 < 0\})$$

(we have the same formula for u_1). We provide a formal calculation which emphasizes the main ideas (see [9, Proposition 3.1] for a rigorous computation). We have

$$\frac{d}{dt} \left(\int_{\mathbb{R}^N} \varphi_\varepsilon(u_2(x, t)) dx \right) = \int_{\mathbb{R}^N} \varphi'_\varepsilon(u_2(x, t)) \frac{\partial u_2}{\partial t}(x, t) dx$$

for a.e. $t \in [0, T]$. Using Equation (18), it follows

$$\begin{aligned} \int_{\mathbb{R}^N} \varphi'_\varepsilon(u_2) \frac{\partial u_2}{\partial t} dx &= \int_{\mathbb{R}^N} \varphi'_\varepsilon(u_2) c[\chi_2] |Du_2| dx \\ &= \int_{\mathbb{R}^N} \left\langle \varphi'_\varepsilon(u_2) Du_2, \frac{c[\chi_2] Du_2}{|Du_2|} \right\rangle dx \\ &= \int_{\mathbb{R}^N} \left\langle D\varphi_\varepsilon(u_2), \frac{c[\chi_2] Du_2}{|Du_2|} \right\rangle dx \end{aligned}$$

since, from $0 < \rho < \eta(T)/2$, and (17), we have $|Du_2| > \eta(T)/2$, for almost every (x, t) such that $\varphi(u_2(x, t)) \neq 0$. By an integration by parts, we obtain

$$\int_{\mathbb{R}^N} \left\langle D\varphi_\varepsilon(u_2), \frac{c[\chi_2] Du_2}{|Du_2|} \right\rangle dx = - \int_{\mathbb{R}^N} \varphi_\varepsilon(u_2) \operatorname{div} \left(c[\chi_2] \frac{Du_2}{|Du_2|} \right) dx.$$

Since u_0 is semiconvex and **(H1)**-**(H2)** hold, by [25, Theorem 5.2], the solutions u_1, u_2 of (18) are still semiconvex in space, i.e. there exists \bar{C} such that

$$D^2 u_1, D^2 u_2 \geq -\bar{C} Id \text{ for a.e. } (x, t) \in \mathbb{R}^N \times [0, T].$$

Using this estimate and the lower-bound gradient estimate (17) again, we have, for almost every $(x, t) \in \mathbb{R}^N \times [0, T]$ such that $\varphi(u_2(x, t)) \neq 0$,

$$-\operatorname{div} \left(\frac{Du_2}{|Du_2|} \right) = -\frac{1}{|Du_2|} \operatorname{trace} \left[\left(Id - \frac{Du_2 \otimes Du_2}{|Du_2|^2} \right) D^2 u_2 \right] \leq \frac{2\bar{C}}{\eta(T)}.$$

It gives

$$\begin{aligned} -\operatorname{div} \left(c[\chi_2] \frac{Du_2}{|Du_2|} \right) &= -\left\langle Dc[\chi_2], \frac{Du_2}{|Du_2|} \right\rangle - c[\chi_2] \operatorname{div} \left(\frac{Du_2}{|Du_2|} \right) \\ &\leq |Dc[\chi_2]|_{L^\infty(\mathbb{R}^N)} + \frac{2\bar{C}|c[\chi_2]|_{L^\infty(\mathbb{R}^N)}}{\eta(T)} := K \end{aligned}$$

since $c[\chi_2]$ is nonnegative bounded Lipschitz continuous by (15) and Step 1. Finally, we obtain, for a.e. $t \in [0, T]$,

$$\frac{d}{dt} \left(\int_{\mathbb{R}^N} \varphi_\varepsilon(u_2(x, t)) dx \right) \leq K \int_{\mathbb{R}^N} \varphi_\varepsilon(u_2(x, t)) dx$$

which yields (21) through a classical Gronwall's argument.

By the same kind of arguments, we can estimate $\mathcal{L}^N(\{-\rho \leq u_0 < 0\})$ to obtain

$$\mathcal{L}^N(\{-\rho \leq u_2(\cdot, t) < 0\}) + \mathcal{L}^N(\{-\rho \leq u_1(\cdot, t) < 0\}) \leq \frac{2C}{\eta_0} \mathcal{L}^N(B(0, R_0 + 1)) e^{Kt} \rho. \quad (22)$$

6. *Part (i) Contraction property for \mathcal{F} (stability estimates with respect to variations of the velocity).* Since u_1 and u_2 are the solutions of (18) with $c[\chi_1]$ and $c[\chi_2]$ respectively, we have the ‘‘continuous dependence’’ type result: for all $t \in [0, T]$,

$$|(u_1 - u_2)(\cdot, t)|_{L^\infty(\mathbb{R}^N)} \leq |Du_0|_{L^\infty(\mathbb{R}^N)} e^{\Lambda t} \int_0^t |(c[\chi_1] - c[\chi_2])(\cdot, s)|_{L^\infty(\mathbb{R}^N)} ds, \quad (23)$$

where $\Lambda = \max\{|Dc[\chi_1]|_{L^\infty(\mathbb{R}^N)}, |Dc[\chi_2]|_{L^\infty(\mathbb{R}^N)}\}$.

7. *Part (i) Contraction property for \mathcal{F} (end of the proof).* From (19), (20), (22) and (23), we get

$$\begin{aligned} & |\mathcal{F}(\chi_1) - \mathcal{F}(\chi_2)|_{Y, \tau} \\ & \leq \frac{2C}{\eta_0} \mathcal{L}^N(B(0, R_0 + 1)) e^{Kt} \sup_{t \in [0, \tau]} |(u_1 - u_2)(\cdot, t)|_{L^\infty(\mathbb{R}^N)} \\ & \leq \bar{L} \sup_{t \in [0, \tau]} \int_0^t |(c[\chi_1] - c[\chi_2])(\cdot, s)|_{L^\infty(\mathbb{R}^N)} ds \\ & \leq \bar{L} \tau |\chi_1 - \chi_2|_{Y, \tau} \end{aligned}$$

for some constant \bar{L} . Therefore, we have contraction for τ small enough. This implies the uniqueness of a classical solution to (5) on the time interval $[0, \tau]$. Noticing that all the constants depend only on the given data, we conclude by a step-by-step argument to obtain the uniqueness on the whole interval $[0, T]$.

8. *Part (ii).* The additional difficulty comparing to the proof of (i) is the fact that u_0 is not supposed to be semiconvex anymore and then $u(\cdot, t)$ is not semiconvex. Nevertheless, we assume that $\delta > 0$, i.e. the velocity is positive. Such a property implies the creation of the interior ball property of radius γt for $\{u(\cdot, t) \geq 0\}$ for every $t > 0$ (see Cannarsa and Frankowska [13] and [7, Lemma 2.3]). Roughly speaking, we recover this way the semiconvexity property for $u(\cdot, t)$ (see the comment after the statement of Theorem 3.2).

Using arguments similar to those in the proof of Part (i) and the interior ball regularization, we prove the following Gronwall type inequality

$$\begin{aligned} & |\mathbb{1}_{\{u_1(\cdot, t) \geq 0\}} - \mathbb{1}_{\{u_2(\cdot, t) \geq 0\}}|_{L^1(\mathbb{R}^N)} \\ & \leq C [\text{per}(\{u_1(\cdot, t) \geq 0\}) + \text{per}(\{u_2(\cdot, t) \geq 0\})] \\ & \quad \int_0^t |\mathbb{1}_{\{u_1(\cdot, s) \geq 0\}} - \mathbb{1}_{\{u_2(\cdot, s) \geq 0\}}|_{L^1(\mathbb{R}^N)} ds \end{aligned}$$

where $u_i, i = 1, 2$ are two weak solutions of (5), C is a constant depending on the constants of the problem and $\text{per}(\{u_i(\cdot, t) \geq 0\})$ is the \mathcal{H}^{N-1} measure (the perimeter) of the set $\partial\{u_i(\cdot, t) \geq 0\}$. In order to apply Gronwall's Lemma it is sufficient to know that the functions $t \mapsto \text{per}(\{u_i(\cdot, t) \geq 0\})$ belong to L^1 . This fact is proved by applying the co-area formula. Finally, it follows $\mathbb{1}_{\{u_1(\cdot, t) \geq 0\}} = \mathbb{1}_{\{u_2(\cdot, t) \geq 0\}}$ for all $t \in [0, T]$ and therefore $u_1 = u_2$ since they are solution of the same equation. \square

4 A counter-example to uniqueness [7]

The following example is inspired from [10].

Let us consider, in dimension $N = 1$, the following equation of type (5),

$$\begin{cases} \frac{\partial U}{\partial t} = (1 \star \mathbb{1}_{\{U(\cdot, t) \geq 0\}}(x) + c_1(t)) |DU| & \text{in } \mathbb{R} \times (0, 2] \\ U(\cdot, 0) = u_0 & \text{in } \mathbb{R}, \end{cases} \quad (24)$$

where we set $c_0(x, t) := 1$, $c_1(x, t) := c_1(t) = 2(t-1)(2-t)$ and $u_0(x) = 1 - |x|$. Note that $1 \star \mathbb{1}_A = \mathcal{L}^1(A)$ for any measurable set $A \subset \mathbb{R}$.

Note that $c_0 \equiv 1$ does not satisfies exactly **(H1)** but this is not the point here: because of the finite speed of propagation property, it is possible to modify c_0 such that **(H1)** and the construction below holds.

We start by solving auxiliary problems for time in $[0, 1]$ and $[1, 2]$ in order to produce a family of solutions for the original problem in $[0, 2]$.

1. *Construction of a solution for $0 \leq t \leq 1$.* The function $x_1(t) = (t-1)^2$ is the solution of the ordinary differential equation (ode in short)

$$\dot{x}_1(t) = c_1(t) + 2x_1(t) \text{ for } 0 \leq t \leq 1, \quad \text{and } x_1(0) = 1,$$

(note that $\dot{x}_1 \leq 0$ in $[0, 1]$). Consider

$$\begin{cases} \frac{\partial u}{\partial t} = \dot{x}_1(t) \left| \frac{\partial u}{\partial x} \right| & \text{in } \mathbb{R} \times (0, 1], \\ u(\cdot, 0) = u_0 & \text{in } \mathbb{R}. \end{cases} \quad (25)$$

There exists a unique continuous viscosity solution u of (25). Looking for u under the form $u(x, t) = v(x, \Gamma(t))$ with $\Gamma(0) = 0$, we obtain that v satisfies

$$\frac{\partial v}{\partial t} \dot{\Gamma}(t) = \dot{x}_1(t) \left| \frac{\partial v}{\partial x} \right|.$$

Choosing $\Gamma(t) = -x_1(t) + 1$, we get that v is the solution of

$$\begin{cases} \frac{\partial v}{\partial t} = - \left| \frac{\partial v}{\partial x} \right| & \text{in } \mathbb{R} \times (0, 1], \\ v(\cdot, 0) = u_0 & \text{in } \mathbb{R}. \end{cases}$$

By the Oleinik-Lax formula, $v(x, t) = \inf_{|x-y| \leq t} u_0(y)$. Since u_0 is even, we have, for all $(x, t) \in \mathbb{R} \times [0, 1]$,

$$u(x, t) = \inf_{|x-y| \leq \Gamma(t)} u_0(y) = u_0(|x| + \Gamma(t)) = u_0(|x| - x_1(t) + 1).$$

Therefore, for $0 \leq t \leq 1$,

$$\{u(\cdot, t) > 0\} = (-x_1(t), x_1(t)) \quad \text{and} \quad \{u(\cdot, t) \geq 0\} = [-x_1(t), x_1(t)]. \quad (26)$$

We will see in Step 3 that u is a solution of (24) in $[0, 1]$.

2. *Construction of solutions for $1 \leq t \leq 2$.* Consider now, for any measurable function $0 \leq \gamma(t) \leq 1$, the unique solution y_γ of the ode

$$\dot{y}_\gamma(t) = c_1(t) + 2\gamma(t)y_\gamma(t) \text{ for } 1 \leq t \leq 2, \quad \text{and } y_\gamma(1) = 0. \quad (27)$$

By comparison, we have $0 \leq y_0(t) \leq y_\gamma(t) \leq y_1(t)$ for $1 \leq t \leq 2$, where y_0, y_1 are the solutions of (27) obtained with $\gamma(t) \equiv 0, 1$. In particular, it follows that $\dot{y}_\gamma \geq 0$ in $[1, 2]$. Consider

$$\begin{cases} \frac{\partial u_\gamma}{\partial t} = \dot{y}_\gamma(t) \left| \frac{\partial u_\gamma}{\partial x} \right| & \text{in } \mathbb{R} \times (1, 2], \\ u_\gamma(\cdot, 1) = u(\cdot, 1) & \text{in } \mathbb{R}, \end{cases}$$

where u is the solution of (25). Again, this problem has a unique continuous viscosity solution u_γ and setting $\Gamma_\gamma(t) = y_\gamma(t) \geq 0$ for $1 \leq t \leq 2$, we obtain that v_γ defined by $v_\gamma(x, \Gamma_\gamma(t)) = u_\gamma(x, t)$ is the unique continuous viscosity solution of

$$\begin{cases} \frac{\partial v_\gamma}{\partial t} = \left| \frac{\partial v_\gamma}{\partial x} \right| & \text{in } \mathbb{R} \times (0, \Gamma_\gamma(2)], \\ v_\gamma(\cdot, 0) = u(\cdot, 1) & \text{in } \mathbb{R}. \end{cases}$$

Therefore, for all $(x, t) \in \mathbb{R} \times [1, 2]$, we have

$$u_\gamma(x, t) = \sup_{|x-y| \leq y_\gamma(t)} u(y, 1) = \begin{cases} 0 & \text{if } |x| \leq y_\gamma(t), \\ u(|x| - y_\gamma(t), 1) & \text{otherwise.} \end{cases}$$

(Note that $u(-x, t) = u(x, t)$ since u_0 is even and, since $u(\cdot, 1) \leq 0$, by the maximum principle, we have $u_\gamma \leq 0$ in $\mathbb{R} \times [1, 2]$.) It follows that, for all $1 \leq t \leq 2$,

$$\{u_\gamma(\cdot, t) > 0\} = \emptyset \quad \text{and} \quad \{u_\gamma(\cdot, t) \geq 0\} = \{u_\gamma(\cdot, t) = 0\} = [-y_\gamma(t), y_\gamma(t)]. \quad (28)$$

3. *There are several weak solutions of (24).* Set, for $0 \leq \gamma(t) \leq 1$,

$$\begin{aligned} c_\gamma(t) &= c_1(t) + 2x_1(t), & U_\gamma(x, t) &= u(x, t) & \text{if } (x, t) \in \mathbb{R} \times [0, 1], \\ c_\gamma(t) &= c_1(t) + 2\gamma(t)y_\gamma(t), & U_\gamma(x, t) &= u_\gamma(x, t) & \text{if } (x, t) \in \mathbb{R} \times [1, 2]. \end{aligned}$$

Then, from Steps 1 and 2, U_γ is the unique continuous viscosity solution of

$$\begin{cases} \frac{\partial U_\gamma}{\partial t} = c_\gamma(t) \left| \frac{\partial U_\gamma}{\partial x} \right| & \text{in } \mathbb{R} \times (0, 2], \\ U_\gamma(\cdot, 0) = u_0 & \text{in } \mathbb{R}. \end{cases}$$

Taking $\chi_\gamma(\cdot, t) = \gamma(t) \mathbb{1}_{[-y_\gamma(t), y_\gamma(t)]}$ for $1 \leq t \leq 2$, from (26) and (28), we have

$$\mathbb{1}_{\{U_\gamma(\cdot, t) > 0\}} \leq \chi_\gamma(\cdot, t) \leq \mathbb{1}_{\{U_\gamma(\cdot, t) \geq 0\}},$$

(see Figure 1). It follows that all the U_γ 's, for measurable $0 \leq \gamma(t) \leq 1$, are weak solutions of (24) so we do not have uniqueness and the set of solutions is quite large.

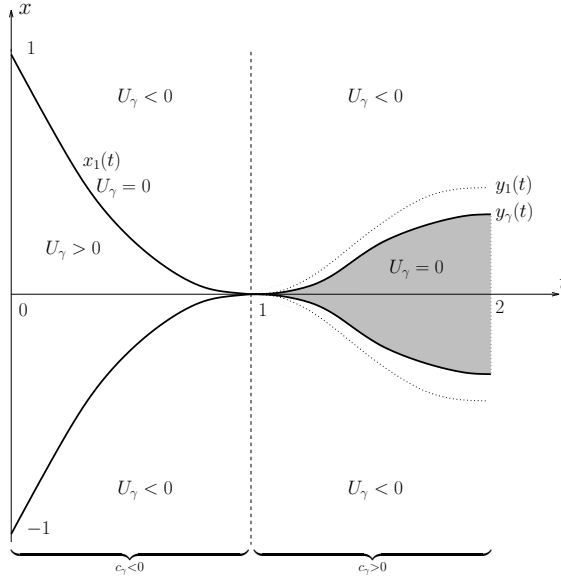


Figure 1: *Fattening phenomenon for the functions U_γ .*

Appendix: L^1 -viscosity solutions and a stability result for weak convergence in time

We recall that the definition of L^1 -viscosity solutions was introduced in Ishii's paper [24]. We refer also to Nunziante [27, 28] and Bourgoing [11, 12] for a complete presentation of the theory.

Consider the equation

$$\begin{cases} \frac{\partial v}{\partial t} = \bar{c}(x, t) |Dv| & \text{in } \mathbb{R}^N \times (0, T) \\ v(\cdot, 0) = u_0 & \text{in } \mathbb{R}^N, \end{cases} \quad (29)$$

where the velocity $\bar{c} : \mathbb{R}^N \times (0, T) \rightarrow \mathbb{R}$ is defined for almost every $t \in (0, T)$. We also assume that \bar{c} satisfies

(H3) The function \bar{c} is continuous with respect to $x \in \mathbb{R}^N$ and measurable in t . For all $x, y \in \mathbb{R}^N$ and almost all $t \in [0, T]$,

$$|\bar{c}(x, t)| \leq M \quad \text{and} \quad |\bar{c}(x, t) - \bar{c}(y, t)| \leq L|x - y|.$$

Definition 4.1 (L^1 -viscosity solutions)

An upper-semicontinuous (respectively lower-semicontinuous) function v on $\mathbb{R}^N \times [0, T]$ is a L^1 -viscosity subsolution (respectively supersolution) of (29), if

$$v(0, \cdot) \leq u_0 \quad (\text{respectively } v(0, \cdot) \geq u_0),$$

and if for every $(x_0, t_0) \in \mathbb{R}^N \times [0, T]$, $b \in L^1(0, T)$, $\varphi \in C^\infty(\mathbb{R}^N \times (0, T))$ and continuous function $G : \mathbb{R}^N \times (0, T) \times \mathbb{R}^N \rightarrow \mathbb{R}$ such that

(i) the function

$$(x, t) \mapsto v(x, t) - \int_0^t b(s) ds - \varphi(x, t)$$

has a local maximum (respectively minimum) at (x_0, t_0) over $\mathbb{R}^N \times (0, T)$ and such that (ii) for almost every $t \in (0, T)$ in some neighborhood of t_0 and for every (x, p) in some neighborhood of (x_0, p_0) with $p_0 = D\varphi(x_0, t_0)$, we have

$$\bar{c}(x, t)|p| - b(t) \leq G(x, t, p) \quad (\text{respectively } \bar{c}(x, t)|p| - b(t) \geq G(x, t, p))$$

then

$$\frac{\partial \varphi}{\partial t}(x_0, t_0) \leq G(x_0, t_0, p_0) \quad (\text{respectively } \frac{\partial \varphi}{\partial t}(x_0, t_0) \geq G(x_0, t_0, p_0)).$$

Finally we say that a locally bounded function v defined on $\mathbb{R}^N \times [0, T]$ is a L^1 -viscosity solution of (29), if its upper-semicontinuous (respectively lower-semicontinuous) envelope is a L^1 -viscosity subsolution (respectively supersolution).

Theorem 4.2 (Existence and uniqueness in the L^1 sense)

For any $T > 0$, under assumptions **(H0)** and **(H3)**, there exists a unique L^1 -viscosity solution to (29).

Finally, let us consider the solutions v^ε to the following equation

$$\begin{cases} \frac{\partial v^\varepsilon}{\partial t} = \bar{c}^\varepsilon(x, t)|Dv^\varepsilon| & \text{in } \mathbb{R}^N \times (0, T), \\ v^\varepsilon(\cdot, 0) = u_0 & \text{in } \mathbb{R}^N. \end{cases} \quad (30)$$

The following stability result is a particular case of a general stability result proved by Barles in [6].

Theorem 4.3 (L^1 -stability) [6]

Under assumption **(H0)**, let us assume that the velocity \bar{c}^ε satisfies **(H3)** (with some constants M, L independent of ε). Let us consider the L^1 -viscosity solution v^ε to (30). Assume that v^ε converges locally uniformly to a function v and, for all $x \in \mathbb{R}^N$,

$$\int_0^t \bar{c}^\varepsilon(x, s) ds \rightarrow \int_0^t \bar{c}(x, s) ds \quad \text{locally uniformly in } (0, T).$$

Then v is a L^1 -viscosity solution of (29).

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