

Grain size distributions of fault rocks: a comparison between experimentally and naturally deformed granitoids.

Nynke Keulen, Renée Heilbronner, Holger Stünitz, Anne-Marie Boullier, Hisao

Ito

To cite this version:

Nynke Keulen, Renée Heilbronner, Holger Stünitz, Anne-Marie Boullier, Hisao Ito. Grain size distributions of fault rocks: a comparison between experimentally and naturally deformed granitoids.. Journal of Structural Geology, 2007, 29, pp.1282-1300. 10.1016/j.jsg.2007.04.003. hal-00194259

HAL Id: hal-00194259 <https://hal.science/hal-00194259v1>

Submitted on 6 Dec 2007

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Manuscript

37 **Grain size distributions of fault rocks: a comparison between experimentally and**

38 **naturally deformed granitoids.**

39

40 Nynke Keulen, Renée Heilbronner, Holger Stünitz, Anne-Marie Boullier, Hisao Ito

41

42 **1 Abstract**

43 We have investigated the grain size distribution (GSD) of granitoid fault rock by comparing 44 experimentally produced gouge with fault rock from the Nojima Fault Zone. Triaxial experiments 45 were carried out on wet and dry intact samples of Verzasca Gneiss at $T = 300$ and 500° C, $P_c = 500$ and 1030 MPa, $ε = 0.013 - 1.6 \times 10^{-4} s^{-1}$. The GSD has been determined from SEM-micrographs and is 47 characterized by the slope, D, of its log(frequency)-log(radius) histogram. The GSD is not fractal; we 48 observe two slopes for all GSDs. The larger grains in experimentally deformed samples have a D-49 value, D_>, of 2.04 and 2.26 for feldspar and quartz gouge. Cracked grains yield values of D = 1.5-1.6.
50 Increasing the confining pressure or temperature decreases the D-value. For grains smaller than \sim 2 Increasing the confining pressure or temperature decreases the D-value. For grains smaller than \sim 2 51 µm another D-value, $D_{\leq x}$ of 0.9-1.1 is observed. The grain radius at the slope change, r_K , corresponds 52 to the grinding limit of quartz, so that r_K probably represents a change in the dominant comminution 53 mechanism from grinding to attrition-processes. The GSD obtained from deformation experiments 54 agrees well with results for the Nojima Fault Zone: $D_>= 2.02$ for gouge and 1.64 for cracked grains; 55 D \leq 0.97. Grain size reduction in fault zones develops by a two-stage process: rupturing creates 56 cracked grains; further displacement of fragments causes further comminution by wear and attrition. 57 Cracked grains have been used to calculate the surface energy associated with faulting; it follows that 58 this energy forms a small fraction in the total energy-budget of earthquakes.

59

60 **Keywords: faulting, granite, gouge, grinding, surface density, power-law distribution**

61 **2. Introduction**

62 **2.1 Overview**

63

64 Major fault zones and earthquakes in the upper crust typically develop highly localized zones 65 of deformation. At the macroscopic scale, faults usually appear as narrow, highly deformed 66 zones or networks embedded in wide damage zones of relatively low deformation. 67 Observations on natural fault zones have shown that this localization of deformation along 68 faults occurs both in space and time (e.g. Mitra, 1993; Faulkner et al., 2003; Wilson et al., 69 2003)*.* Most of the displacement takes place in very narrow zones within the gouge of 70 typically a few to 100 µm wide, leaving the rest of the gouge inactive (e.g. Wibberley and 71 Shimamoto, 2005; Chester et al., 2005; Ma et al., 2006). Bos et al. (2000) have made this 72 same observation for laboratory experiments on halite. 73 74 Within these narrow deformation zones gouge with a wide range of grain sizes has been 75 produced (e.g. Storti et al., 2003, Boullier et al., 2004a). At the microscopic scale, the fault 76 rock consists of cataclasite or gouge with properties that change as a function of time and 77 space. Displacement on a fault zone causes the evolution from a thin fracture to a wider gouge 78 zone. The microstructures of the fault rocks influence and reflect the mechanics of faulting, 79 i.e., earthquakes. For example, evidence for coseismic hydraulic fracturing (Boullier et al., 80 2004b) or for a change in the frictional properties of faults resulting from their grain shape 81 (Mair et al., 2002) has been obtained through microstructural studies. Despite experimental 82 data and field observations on gouge, the properties of fault rock and the evolution of cracked 83 grains to gouge still remain rather poorly understood. 84 85 86 **2.2 Fault rocks** 87 88 A number of definitions for fault rocks can be found in the literature (e.g. Sibson, 1977; 89 White, 1982; Wise et al., 1984; Chester et al., 1985; Schmid and Handy, 1991). Usually, 90 cohesive and noncohesive fault rocks are attributed to cataclasis, a process involving

- 91 fracturing, frictional sliding, dilatancy, and rigid body rotation between grain fragments,
- 92 grains, or clusters of grains. Sibson (1977) has defined gouge as a cataclastic rock without
- 93 cohesion; the cohesive rock is called cataclasite. Many cohesive cataclasites in nature are

94 thought to be the product of syn- or post-tectonic healing processes (Schmid and Handy,

95 1991).

96

97 In this study, however, we will distinguish different fault rocks on the basis of microstructures 98 rather than inferred mechanical properties (see Heilbronner & Keulen, 2006). We will use the 99 term "cracked grains" for angular fragments that have fractured without much displacement 100 (such that the original geometrical relationship of the fragments can still be recognized). The 101 term "gouge" will be used for more highly deformed fault rocks where the grains are more 102 rounded and the original spatial relationship between the fragments is lost. The term "fault 103 rocks" thus applies to both cracked grains and gouge.

- 104
-

105 **2.3 Grain size distributions**

106

107 A large number of studies has been dedicated to measuring and analyzing grain size 108 distributions of fault rocks (e.g. Mitra, 1993; Blanpied et al., 1995; Lin, 1999; Bos et al., 109 2000; Stewart et al., 2000; Kato et al., 2003). The grain size distribution (GSD) is considered 110 as an important parameter for the characterization of the gouge (e.g. Sammis et al., 1987; 111 Marone and Scholz, 1989; Hadizadeh and Johnson, 2003; Wilson et al., 2005; Chester et al., 112 2005). Within a certain size range the GSD may be self-similar (scale invariant) and in this 113 case can be described by a power-law relationship (Mandelbrot, 1982). The slope of some 114 power-law distributions, D, is called the fractal dimension (Turcotte, 1992; Korvin, 1992, and 115 references in both). This D-value has often been used to compare the GSD of natural 116 earthquake and fault zones, like the San Andreas Fault zone (e.g. Anderson et al., 1980; 117 Blenkinsop, 1991), the Moine Thrust (Blenkinsop and Rutter, 1986) or fault systems in Japan 118 (Monzawa and Otsuki, 2003; Otsuki et al., 2003) with fault rocks derived from experiments 119 (e.g. Marone and Scholz, 1989; Morrow and Byerlee, 1989; Amitrano and Schmittbuhl, 120 2002). 121 122 For the measurements of the GSD, several different methods have been used (see Table 1). 123 The most classical one, originating from sedimentology, is sieving. By this method a

124 distribution of mass against size is obtained (Anderson et al., 1980; Sammis et al., 1986).

- 125 Other studies have used the Coulter counter (An and Sammis, 1994; Amitrano and
- 126 Schmittbuhl, 2002; Wilson et al., 2005) or optical microscopy on thin sections (Biegel et al.,
- 127 1989; Monzawa and Otsuki, 2003); these methods provide a description of the GSD in terms
- 128 of frequency versus grain size. A disadvantage of sieving and the use of light microscopes is
- 129 that the size of smallest detectable grains is \sim 2 μ m (Table 1).
- 130
- 131 An easy way to measure the sizes of the grains smaller than 2 μ m is scanning electron
- 132 microscopy (SEM), as was done by Sammis et al. (1987) and Shao and Zou (1996), but in
- 133 their samples no small grains were observed. Olgaard and Brace (1983), Yund et al. (1990),
- 134 and Chester et al. (2005) found grains as small as 15-50 nm with the TEM, but the TEM is
- 135 unsuitable for measuring large numbers of grains necessary to obtain a size distribution.
- 136 Wilson et al. (2005) used a Coulter counter and measured grains of 40 nm to 300 µm.
- 137 Nevertheless, the frequency distribution of grains sized between 30 nm and 2 µm is not very
- 138 well known. We find that using scanning electron micrographs of impregnated and polished
- 139 sections of gouge, reliable GSD in the size range larger than 60 nm can be obtained.
- 140

141 In a previous study (Heilbronner & Keulen, 2006) two types of gouge can be distinguished:

- 142 cracked grains and gouge. Cracked grains and gouge can be discriminated on the basis of
- 143 particle shape and grain size distributions. The D-values of gouge are always higher than
- 144 those of cracked grains. Composition (quartz versus feldspar) influences various aspects of
- 145 the particle shape of cracked grains and gouge (elongation, paris factor, etc.).
- 146

147 Table 1 shows a compilation of published results on the GSD of fault rocks: values of the

148 fractal dimension vary between 0.8 and 2.6, but values of D between 1.5 and 2.1 are the most

- 149 common for granitoid rocks. High D-values are generally obtained for more recently active
- 150 fault zones, though a wide spread of D-values within a single fault zone can be observed.
- 151

152 As there appears to be a link between fault rock microstructure (shape and GSD) and the 153 mechanical behavior of faults, and since fossil fault rocks are important clues to seismic 154 activity in the past, this study focuses on the microstructures of naturally and experimentally 155 produced faults. In this paper we examine the influence of experimental conditions on the 156 fault rock microstructure and we compare the results to a natural fault rock cored from a 157 recently active segment of the Nojima Fault (Japan). The focus of the paper is on the grain 158 size distributions (D-values cracked grains and gouge) and on the geometry and spatial 159 distribution of the fault rock in the samples (D-mapping). The aim is to study the influence of 160 temperature, confining pressure, axial shortening rate, total axial shortening, and fluid content 161 on the grain size development of gouge material in deformed granitoids.

196 **3.2 Rock deformation experiments**

197

198 Samples of Verzasca gneiss measuring 6.39 mm in diameter and 10.1 to 12.9 mm in length 199 (L_0) are cored parallel to the mesoscopic lineation. Table 2 lists the used abbreviations. At this 200 size, the sample width corresponds to approximately 25 grains. After polishing their flat ends 201 the samples are dried in an oven at 110°C and atmospheric pressure for at least 24 hours. The 202 samples are wrapped in Ni-foil and placed in a gold jacket. In all cases except for experiment 203 63nk, a "room-humidity" sample, approximately 0.2 wt% H_2O (1.68 to 2.14 µl) is added and 204 the jacket is weld-sealed. During welding, the jacket and sample are cooled in order to avoid 205 water loss. The exact water content of the "wet" and "room-humidity" samples has not been 206 determined, and we assume that also in room-humidity samples some water is present at grain 207 boundaries.

208

209 Deformation is carried out in a Grigg's solid medium deformation apparatus as described by 210 Tullis & Tullis (1986) and De Ronde (2004), using straight furnaces, alumina pistons and 211 sodium chloride as the confining medium. The pistons are slightly bevelled, thus reducing 212 their diameter from 6.39 mm to 6.2 mm at the sample-piston interface. Temperature is 213 measured with Pt/Pt-10%Rh or Cromel/Alumel thermocouples. The standard or reference 214 deformation conditions are "wet" samples, 300°C temperature, 500 MPa confining pressure, 215 10^{-4} s⁻¹ shortening rate, and a total axial displacement of 3.5 - 5.0 mm (samples 38nk and 216 64nk). In separate runs, one or more parameters are varied: 500°C temperature, 1030 MPa 217 confining pressure, up to 5.8 mm total displacement, 10^{-6} s⁻¹ shortening rate or no H₂O is 218 added (see Table 3). During the experiment, pressure, load and displacement are recorded on 219 a strip chart. Using the rigC4 program (www.unibas.ch/earth/micro/software/), with a 220 correction for rig stiffness, the force-displacement and stress-strain curves are calculated. 221 222 After the experiment, the samples are quenched. To avoid loss or damage of the non-cohesive 223 fault rocks, the jackets are punctured and the samples are impregnated with epoxy. The 224 samples are cut along the compressional axis and the deformed length of the sample is re-

225 measured. Thin sections are prepared (polished and carbon coated) for observation in the light 226 and scanning electron microscope.

- 227
- 228 **3.3 Grain size analysis**
- 229

230 For the grain size analysis, images with a wide range of magnifications are used. The 231 micrographs of the experimental microstructures are collected using a Philips XL30 ESEM 232 scanning electron microscope (20kV acceleration voltage; backscatter contrast); those of the 233 Nojima Fault Zone are acquired with a Leica Stereoscan 440 SEM. For each sample a cascade 234 of micrographs is prepared, starting with overviews at low magnifications (100x) and 235 zooming in (to 20,000x) on selected areas with a factor of 2 to 2.5 from one magnification to 236 the next. At each magnification, three or four images are taken of each kind of fault rock 237 (quartz, feldspar, cracked grains and gouge). The total range of grain sizes (radius, r) 238 observable at these magnifications is from <20 nm to >200 µm, however, because of the 239 scarcity of measurable grains at either end of the distribution, the statistically reliable range

- 240 does not exceed three orders of magnitude.
- 241

242 We use Image SXM (http://reg.ssci.liv.ac.uk/) and ImageJ (http://rsb.info.nih.gov/ij/) and 243 special Macros and Plug-Ins (http://www.unibas.ch/earth/micro/software) for segmentation 244 and image analysis. First, the images are pre-processed (removal of noise) and scaled to 245 pixels (Figure 2a). Bitmaps of quartz and feldspar grains are obtained by density slicing at the 246 corresponding grey levels and subsequent separation of grains (Figure 2b). The cross 247 sectional areas, A (pixel²), and the perimeters, P (pixel), of the grains are measured using the 248 • Analyze menu of the software. After removing grains with areas smaller than 20 pixel², 300 249 to 1200 grains per image could be used. Some of the images taken at the highest 250 magnification could not be processed automatically; in this case the grain boundaries are 251 traced manually.

252

253 The resulting files of areas and perimeters were transferred to a spreadsheet program. The 254 areas were restored by adding the pixel values of P, i.e., the band of pixels outlining the grain, 255 to the area A within the outline. The equivalent radii, r, were calculated as

256

257 $r = \sqrt{(A+P)/\pi}$ (1)

258

259 and collected in a histogram (Figure 2c). On a log(frequency)-log(size) plot with 20 bins per

- 260 order of magnitude of grain size, the slope (D) of the power-law fit is determined (Figure 2d).
- 261 The results from all magnifications are then combined into a single plot (Figure 2e),
- 262 multiplying the frequency in each bin by a factor that reflects the relative magnification of the

297 by audible cracking. Sample 60nk (run at the highest confining pressure) is different from all 298 others in that it has produced distinct acoustic events, which accompany load drops on the 299 order of 20 kN (corresponding to approximately 500 MPa).

300

301 At 45kN (1000-1200 MPa), sample 70nk, deformed at relatively high temperature (500°C) 302 and low axial shortening rate $(10^{-6} s^{-1})$, is the weakest. At 60-75 kN (1750 MPa) sample 63nk

303 (room humidity sample) is the strongest. Samples deformed at standard conditions (64nk and

304 38nk) support loads of 55-65 kN, corresponding to strengths of 1300-1500 MPa.

305

306 For comparison with published data, we also show a stress-strain diagram (Figure 3b). The

307 differential stress is calculated with a correction for an increase in cross-sectional area,

308 assuming constant volume. The Verzasca Gneiss has approximately the same strength as

309 Westerly Granite (Fig. 3b) deformed by Tullis and Yund (1977). The main differences with

310 respect to Westerly Granite experiments are the higher shortening at yield stress (with typical

311 values ~10% versus ~5%), the higher shortening rate (mainly at 10^{-4} s⁻¹, see Table 3), and the

312 amount of added water with which the Verzasca Gneiss has been deformed. In both sets of

313 experiments, a temperature increase from 300°C to 500°C has caused a strength reduction of

314 approximately 200 MPa. Note that Tullis and Yund (1977) have used pyrophyllite or talc as

315 confining medium, which means that the true confining pressures probably are \sim 2/3 of the

- 316 published values (J.Tullis, pers.comm., 2004).
- 317

318 We have noted a difference between the measured lengths of the deformed samples and the

319 deformed lengths calculated from the total axial displacements derived from the force-

320 displacement records (compare Figure 3a and Table 3). The calculated displacements exceed

321 the measured displacements by up to 1 mm. In the rest of the paper, we will therefore use the

322 more conservative values of the measured displacements, not the displacements shown in

323 Figure 3. This difference in measured length is not critical for the results, as the displacement 324 is used semi-quantitatively in this study.

325

- 326 **4.2 Faulted samples**
- 327

328 In most samples deformed at $T = 300^{\circ}$ C and $P_c = 500$ MPa, through-going fault zones are 329 formed accommodating the axial displacement within a single gouge zone. At $T = 500^{\circ}C$, P_c

330 = 500 MPa and occasionally at T = 300 \degree C, P_c = 500 MPa, conjugate fault zones have

331 developed and partitioned the axial displacement over at least two fault zones, leading to a

332 smaller displacement on any given fault. Sample 60nk deformed at $P_c = 1030$ MPa shows a

333 set of parallel faults. The corresponding load-displacement curve (Figure 3) shows an

334 equivalent number of load drops. It is likely that each drop in the differential stress is related

335 to the creation of a new fault plane.

336

337 The samples are deformed along relatively wide fault zones of 1.0-1.4 mm width, which 338 become narrower with higher confining pressure or higher temperature; with increasing

339 displacement, they tend to broaden. Most of the gouge zones are formed around mica grains

340 or the connection of two mica grains. At a temperature of 300 $^{\circ}$ C and shortening rates of 10⁻⁴

 341 s⁻¹ biotite deforms by gliding, kinking and occasionally by fracturing.

342

343 Using the D-mapping technique described elsewhere (Heilbronner & Keulen, 2006), maps of 344 local variations of the grain size distribution are prepared for three of the experimental 345 samples (Figure 4). Cracking and gouge formation occur throughout the samples. In samples 346 38nk and 102nk gouge regions coalesce; in sample 60nk there is less coalescence. The 347 volume proportion of gouge formed during the deformation is determined by measuring the 348 black area (indicating $D > 1.75$) in Figure 4 and dividing this proportion by the total sample 349 area. For increasing axial displacement, d, increasing amounts of gouge are formed. After $d =$ 350 4 mm, sample 38nk (reference sample) has 24.5% of gouge; after d = 5.8 mm, sample 102nk 351 (deformed at the same conditions) has 54.8%. At higher confining pressures, there is less 352 gouge: after $d = 3.5$ mm, sample 60nk ($P_c = 1000$ MPa) has only 15.5% gouge.

353

354

355 **4.3 Microstructures**

356

357 In the experiments at $P_c = 500 \text{ MPa}$ and T = 300°C, the original grains in the gouge zone form 358 broad and elongated arrays of fractured particles (Figure 5a); at higher confining pressures, 359 these arrays are much narrower (Figure 5b). Quartz and feldspar are both strongly fractured; 360 at higher confining pressures, quartz is less intensively fractured than feldspar. A variation in 361 grain shapes can be observed: quartz grains are more elongated, feldspar grains are more 362 irregularly shaped (Figure 5c & 5d); the same variation in shape has been observed for a 363 sample deformed at low PT-conditions (Heilbronner and Keulen, 2006).

400

- 401 The GSD of the experimentally deformed granitoid samples are shown in Figure 6 & 7, the 402 derived values of D_{<,} D_>, and r_K are shown in Tables 4 & 5. The d_K grain diameter is two 403 times the r_K grain radius indicated in Figure 6, and has been used to compare our data to other 404 studies. We consider the samples 64nk and 38nk as reference samples. The experiments are 405 performed at T = 300°C, confining pressure P_c = 500 MPa, shortening rate $\cdot \varepsilon = 1.5 \cdot 10^{-4} \text{ s}^{-1}$ 406 and 0.2 %wt H2O was added to the samples. The total axial displacement, ∆L, of 64nk and 407 38nk is 2.2 and 3 mm, respectively. In Figure 6, each of the plots shows the GSDs resulting 408 from the change of one of the experimental conditions with respect to 64nk.
- 409

410 The reference values derived from the GSDs of sample 64nk and 38nk are the following: D_{\leq}

- 411 ranges from 0.95 to 1.07 (D_< of cracked feldspar = 1.00; D_< of feldspar gouge = 1.02; D_< of
- 412 cracked quartz = 0.95 ; D_< of quartz gouge = 1.07), cracked grains have somewhat lower
- 413 values than gouge. r_K is 1.1 1.8 µm (cracked feldspar = 1.8 µm; r_K of feldspar gouge = 1.6
- 414 μ m; r_K of cracked quartz = 1.6 μ m; r_K of quartz gouge = 1.1 μ m), quartz yields lower values
- 415 than feldspar and cracked grains have higher values than gouge. $D_>\iota$ is 1.5 1.6 for cracked
- 416 grains (D_> of cracked feldspar = 1.6; D_> of cracked quartz = 1.5) and 2.04 2.26 for gouge
- 417 (D, of feldspar gouge = 2.04; D, of quartz gouge = 2.26), quartz has a lower D,-value than
- 418 feldspar for cracked grains and a higher D>-value for gouge.
- 419

420 Increasing the confining pressure from 500 MPa to 1030 MPa leads to a marked reduction of 421 both D< and D>, both for cracked grains and gouge. The grain sizes, r_K , at the slope break of 422 cracked grains and gouge do not change significantly, with the possible exceptions of r_K of 423 feldspar gouge and cracked quartz: these values increase from 1.6 to 1.8 µm and from 1.5 to 424 1.8 µm, respectively. This variation is within the error of the measurements.

425

426 Reducing the shortening rate from $10^{-4}s^{-1}$ to $10^{-6}s^{-1}$ does not affect the values of D>; of the D< 427 values only those of cracked feldspar and of quartz gouge appear to be lowered: from 1.00 to 428 0.86 and 1.07 to 0.97, respectively. The grain size r_K of feldspar remains unchanged, while 429 that of cracked quartz and quartz gouge increases from 1.5 to 1.8 μ m and from 1.1 to 1.4 μ m 430 respectively.

499 between 10^{-4} and 10^{-6} s⁻¹ does not show any mechanical effect, consistent with the results of 500 Tullis & Yund (1980, 1992).

501

502

503 The effect of a high confining pressure $(\sim 1 \text{ GPa})$ is comparable to experiments by Tullis & 504 Yund (1977, 1980). Very narrow slip zones accommodate the axial shortening. According to 505 Tullis and Yund (1977) and Hirth and Tullis (1994), the formation of these narrow zones and 506 the decrease in volume of fault rock is related to the transition from mode I to mode II 507 cracking above $P_c \approx 600$ -800 MPa. The number of these narrow zones corresponds to the 508 number of audible cracking events in the experiment 60nk. This observation suggests that the 509 mechanical behavior is not a case of stick-slip as described in the literature (e.g. Byerlee and 510 Brace, 1968; Stesky et al., 1974; Lockner et al., 1986). Instead, the apparent stick slip 511 character of the mechanical data (Figure 3) is produced by multiple rupture events, 512 demonstrating the great strength of already ruptured material at high confining pressure. 513 514 **5.2 Grain size distributions obtained in experiments** 515 516 Changing the physical parameters of experimental deformation also affects the grain size 517 distribution. The effects are discussed with reference to sample 38nk and 64nk, which have 518 been deformed under identical temperature, pressure and shortening rate conditions. 519 520 The largest effect is observed for increasing the confining pressure to \sim 1 GPa. It causes a 521 reduction in the D₂-values (Figure 6, 7, Table 4) and a decrease in the total amount of gouge 522 and cracked grains in the fault zone (see Figure 4). This smaller amount of gouge is related to 523 the transition from mode I to mode II cracking: mode I fracturing is dilatant; in mode II 524 fracturing sharp shear fractures are formed when dilatancy is suppressed. Dilatancy 525 suppression appears to inhibit the formation of large amounts of surface (small grains) and 526 thus may reduce the D_z-value. 527 528 At higher temperatures, there is a decrease in the D_{\ge} values of the gouge (Figure 7, Table 4). 529 In feldspar deformed at $T = 500^{\circ}$ C we observe undulatory extinction and deformation bands, 530 which are developed by micro-cracking at a very small scale, as described by Tullis and Yund

531 (1987). At a temperature of 500°C, Hirth & Tullis (1994) have found an increased dislocation

532 activity in quartz. The less efficient grain size reduction at elevated temperature may be 533 caused by a semi-brittle behavior.

534

535 The effects of not adding water and slower shortening rate are either unsystematic (higher

536 D_svalues for feldspar and lower D_s values for quartz in gouge of non-water added sample

537 63nk; Table 4) or not existent (samples 38nk, 64nk, 102nk; Table 4). The D-values are the

538 same regardless of the amount of shortening of the sample (Figure 7, Table 3).

539

540 After 4 mm axial shortening a gouge of 1.0 mm thickness has been formed. Once a mature 541 zone with gouge is established, this zone is growing wider at greater displacement, but $D_{>}$ of 542 the gouge is not changing. Marone and Scholz (1989) have also observed similar D-values for 543 different shear strains. In general, it can be concluded that the effects of changing the physical 544 conditions of deformation on the grain size distribution are of minor importance except for 545 the increase in confining pressure, and, to a lesser extent, temperature.

- 546
-

547 **5.3 Cracked material and gouge**

548

549 All experimental and Nojima fault rock samples show the same type of grain size distribution 550 (GSD). For small grain sizes ($d < 2 - 3 \mu m$) the GSD is characterized by D-values of ~ 1 , 551 whereas for the larger sizes the D-values range between \sim 1.5 and \sim 2.3 (Table 4). The gouge 552 always shows higher D-values than the cracked material (Figure 6, 8, Table 4). It is 553 interesting to note that the values of D_z in cracked material and gouge, in quartz and feldspar, 554 and in experimental and in Nojima fault rocks all have very similar values and the same range 555 of variation (Table 4). Furthermore, this grain size range is not affected by changing the 556 experimental deformation conditions. In contrast, D> of gouge is the parameter that has been 557 measured in several other studies before (Table 1). It is the most suitable of the D-values to 558 compare natural fault zones in different regions or to compare natural faults with 559 experimentally obtained values. 560

561 Higher D-values indicate a more efficient grain size comminution producing a larger number 562 of smaller grains. The difference between gouge and cracked grains clearly demonstrates that 563 the initial rupturing process does not produce the full amount and size fraction of small grains

564 observed in gouge (Figure 6, 8). Instead, slip on the fault zone causes further comminution of

565 grains in the gouge and an increase in D_{\ge} . Thus, it can be concluded that after the initial

- 566 fragmentation of the rock by rupturing there is a post-rupture processing during slip which
- 567 takes place in the gouge to produce the higher values of D_{\geq} .
- 568
- 569 Further comminution of grains in gouge has been observed by Sammis et al. (1986, 1987) and
- 570 Hadizadeh & Johnson (2003). However, any quantitative characterization of fault
- 571 displacement by GSD measurements appears impossible because an increase in sample
- 572 shortening from 2.2 to 4.8 mm in our experiments (64nk and102nk) does not change any of
- 573 the D-values systematically (Figure 6, 9, Table 5). Thus, the more efficient grain
- 574 comminution in gouge must take place in the first stages of the fault displacement after
- 575 rupturing.
- 576

577 In the experiments performed under identical conditions (38nk, 64nk, 102nk) the variation of 578 D-values of the cracked material is always larger than in the gouge (Table 4). The only 579 systematic difference appears to be the consistently higher D_> values of quartz gouge 580 compared to feldspar gouge, i.e., for the coarser grain size fraction $(> 2 \mu m)$ quartz shows a 581 more efficient comminution of grains than feldspar. The D>-values for quartz and feldspar 582 gouge are almost identical in these three experiments and indicate that this D-value is very 583 well reproducible and therefore forms a reliable parameter to compare the GSD of fault zones. 584 Another difference is the lower r_K -value for quartz compared to feldspar. Thus, the change in 585 the GSD occurs at \sim 1 µm in quartz, while it occurs at \sim 2 µm in feldspar (Table 4).

586

587 **5.4 Natural fault rock**

588

589 The range of values in the natural samples from the Nojima fault is approximately the same as 590 those of the feldspars in experiments (Figure 6, 8, Table 4) so that the natural fault rock 591 appears to behave in a similar way as feldspar. The granodiorite deformed by the Nojima fault 592 consists of ∼53% feldspar (plagioclase and K-feldspar); therefore, feldspar is expected to play 593 a dominant role in the GSD of a mixed gouge.

594

595 The recently formed fault rock of the Nojima Fault zone shows the same D-values as the

- 596 experimentally deformed granitoid rock, especially for feldspar minerals. If we compare the
- 597 D-values of the recent deformation episode with feldspar in experiment 64nk and 38nk,
- 598 performed at T = 300°C, P_c = 500 MPa, ϵ = 1.5·10⁻⁴ s⁻¹, with < 4 mm shortening and with 0.2

600 Temperature and pressure conditions during Nojima deformation were lower than in the

- 601 deformation experiments.
- 602

603 D_> of the Nojima gouge (D_> = 2.02) is in fairly good agreement with data from coarser grain

604 size fractions by Monzawa and Otsuki (2003); they obtain $D_>= 2.192$ to 2.559, for five

605 different samples of the surface outcrop of the Nojima Fault Zone, with a mean of 2.347.

606 Monzawa and Otsuki's (2003) data is more similar to the values of $D_>$ for the experimentally

- 607 deformed quartz so that all data appears to be within the variation of the experimental data. 608
- 609 Our results are consistent with studies of other fault zones (Table 1). Blenkinsop (1991) and

610 An and Sammis (1994) describe both low ($D_0 < 1.6$) and high ($D_0 > 2.0$) values for analyses

611 of gouge originating from the same fault zone. Several other studies describe high D_{\ge} values

612 (e.g. Blenkinsop, 1991; Hadizadeh and Johnson, 2003). In some gouge the lower limit of D>

613 is 1.4 to 1.6, except for the San Andreas Fault Cajon Pass Drillhole granite ($D = 0.8$;

614 Blenkinsop, 1991). We speculate that these gouge with $D_>$ below 1.6 might be monomineralic

615 gouge where fluids and temperature caused healing (e.g. Sammis et al., 1987, Keulen et al.,

616 submitted) or were only cracked (e.g. Marone and Scholz, 1989).

- 617
-

618 **5.5. Surface densities of fault rocks**

619

620 The surface energy in gouge can be used to calculate a part of the energy release during an 621 earthquake event (Kanamori, 1994; Wilson et al., 2005; Chester et al., 2005; Ma et al., 2006).

622 The grain size distribution allows the calculation of the total surface area of the fault rock

623 (Chester et al. 2005, Wilson et al. 2005). The estimates of Chester et al. (2005) for the

624 Punchbowl fault are based on GSDs with an assumed minimum grain size, r_{min} , of 1.6 nm and

625 a constant D-value of $3 (= 2 \text{ in } 2 \text{ dimensions})$ for the entire grain size range. In contrast, we

626 find that in all of our experimental and natural samples the D-value below $r \approx 1$ to 2 µm is D

 $627 \approx 1$.

628

629 If the evolution of the GSD is a two-stage process involving initial rupturing and further grain

630 comminution in the gouge, probably by attrition, shearing, grinding, etc.; it is questionable

631 whether the surface energy calculated from the GSD can directly be related to the seismic

632 energy release as proposed by Wilson et al. (2005) and Chester et al. (2005). Neither is the

- 633 surface energy obtained from the GSD of fault rocks identical to the specific surface energy
- 634 as applied in earthquake energy budget calculations (e.g. Ventakaraman & Kanamori, 2004).
- 635 For the following calculation we use the surface density of cracked grains, but not of gouge,
- 636 to compare to the data of Wilson et al. (2005) and Chester et al. (2005), because we assume
- 637 that this is the material, which is created by the rupture process during the seismic event,
- 638 while the higher surface density of gouge is due to the combined effect of fracturing and
- 639 subsequent deformation during slip.
- 640
- 641 To estimate the surface density for our cracked minerals and gouge we use the concept of the
- 642 fractal fragmentation of the cube and a simple spread sheet program. We set up a histogram of
- 643 13 (logarithmic) bins, extending from 0.03 μ m to 128 μ m and from 0.015 μ m to 64 μ m,
- 644 respectively. Between r_{min} and r_K , the 3-dimensional fractal dimension is set to 2 in both
- 645 cases, corresponding to a value of D_< = 1.0. Between r_K and r_{max} , the fractal dimension is set
- 646 to 2.5 for cracked material and to 3 for gouge, corresponding to $D_>= 1.50$ and $D_>= 2.00$,
- 647 respectively. A value of $r_K = 1 \mu m$ is assumed in both cases. For each of the bins, the number,
- 648 total volume and total surface of the cubes is calculated. Summing over all bins, the surface to
- 649 volume ratio of cracked material and gouge is 1.3 and 6.0 μ m²/ μ m³ i.e., 1.3·10⁶ m⁻¹ and
- 650 6.0 \cdot 10⁶ m⁻¹ for cracked material and gouge.
- 651
- 652 Using the same approach to calculate surface densities for the Punchbowl Fault using the data 653 $(r_{min} = 1.6 \text{ nm} \text{ and } r_{max} = 100 \text{ µm}, D = 2.0)$ given by Chester et al. (2005), we obtain 6.8·10⁸ 654 m⁻¹. If we now compare our value of cracked grains of $1.3 \cdot 10^6$ m⁻¹ to a value of $4.2 \cdot 10^8$ m⁻¹ 655 (corresponding to $4.2 \cdot 10^5$ m² in the 1mm thick layer of the fault, given by Chester et al., 656 2005) or 2.7 \cdot 10⁸ m⁻¹ (corresponding to 80 m²g⁻¹ given by Wilson et al., 2005) we note that in 657 our case, the surface density, and accordingly, the energy required to generate it, is smaller 658 than the values assumed by Wilson et al. (2005) and Chester et al. (2005) by a factor of 200 to 659 300. Finding that the surface densities in cracked rocks and even in gouge are so low suggests 660 to us that the creation of those surfaces cannot play a major role in the energy budget of a 661 seismic event.
- 662
- 663 The discrepancy between our data and the data published by Chester et al. (2005) can be
- 664 explained by the fact that they have measured their small grain size fraction (down to $r = 25$)
- 665 nm) in the TEM, which is a notoriously difficult method for obtaining good statistics of any
- 666 parameter observed. Thus, they did not detect the change of D-values at $r \approx 1 \mu m$. Using D =
- 667 2.0 over the entire range of grain sizes results in a total surface area which is 2 orders of
- 668 magnitude too high. Furthermore, they assume $r_{min} = 1.6$ nm and thus many tiny grains with a
- 669 high surface area density. Note that finding lower surface densities rather supports and
- 670 strengthens their general conclusion that the surface energies do not contribute in any major
- 671 way to the energy dissipated during earthquakes. Pittarello et al. (2006) come to the same
- 672 conclusion, using GSD measurements of pseudotachylyte-related granitoid cataclasites.
- 673

674 **5.6 Grinding limit**

675

676 Grains with very small radii formed by comminution have been found before in deformation 677 experiments (d = 10 nm; Yund et al., 1990), in mining induced fault zones (\sim 40 nm; Olgaard 678 and Brace, 1983; Wilson et al., 2005) and in natural fault zones (40 nm; Wilson et al., 2005, 679 Chester et al. 2005). We observe grains of 30 nm in both experimental and natural fault rock 680 (Figure 9). r_K is 1.2 \pm 0.3 µm for quartz and 1.8 \pm 0.3 µm for feldspar grains. The slope 681 change is observed for both cracked grains and gouge (Figure 6, 8, Table 5). Very small 682 grains are already formed upon cracking of the grains; they are not only a product of gouge 683 formation. D_{\leq} is less dependent on deformation conditions, the type of mineral (quartz or

684 feldspar) and is more similar for cracked grains and gouge than D>.

685

686 Kendall (1978) has demonstrated that a critical radius exists below which particles cannot be 687 comminuted further by grinding. This critical radius, called the grinding limit, is 0.3 µm for

688 SiO2 glass, 0.9 µm for quartz and 2.2 µm for calcite (Steier and Schönert, 1972 in Prasher,

689 1987). Small particles have a different internal stress distribution and are stronger due to

690 smaller flaw sizes and flaw size density (Prasher, 1987) so that very high stresses are required

691 to break them (Mitra, 1978). This is consistent with our observation that no intragranular

692 cracks within grains smaller than 1 µm in radius are observed.

693

694 The r_K-values for quartz are slightly higher than the grinding limit (given as radius) for quartz

695 (Figure 10). There is no data for the feldspar grinding limit available but the higher r_K -values

- 696 for feldspar (Figure 10) are consistent with the observation that the critical grain size for
- 697 grinding is linearly dependent on the elastic modulus E of the mineral (Kendall, 1978),
- 698 because the bulk E-modulus of feldspar is approximately 2 times that of quartz (Bass, 1995).
- 699

700 Thus, it appears likely that r_K represents approximately the grinding limit in the comminution 701 of fault rock. For smaller grain sizes, where grinding is inhibited, communition is less 702 efficient, leading to a lower D<. Grinding is performed by compressional fracturing of grains: 703 a network of cracks will develop and the grain will break into several smaller ones (Jaeger,

704 1967). Other processes, like shearing or attrition also produce comminution. Reches and

705 Dewers (2005) and Dor et al. (2006) describe comminution as dynamic pulverization

706 (explosive granulation followed by dynamic contraction) at earthquake crack tips. For grains

707 smaller than r_K compressional fracturing appears inhibited, but shearing or attrition remains 708 significant. From the difference between D-values in cracked grains and gouge, it is clear that

709 grinding, shearing and attrition during slip cause further comminution in gouge.

- 710
-

711 **5.7 Fractal dimension**

712

713 Fractal dimensions are used to describe geometries that are scale invariant (Turcotte, 1992). 714 To describe a GSD as fractal, this distribution should lie on a straight line over several orders 715 of magnitude in a log(frequency)-log(size) distribution. The GSD of both the experiments and 716 the Nojima Fault Zone gouge does not show a fractal distribution. The data cover 3.5 orders 717 of magnitude, but cannot be fitted with a single straight line. Furthermore, there are statistical

718 problems because most fault rock grain size measurements only cover 1 to 4 orders of

719 magnitude.

720

721 The advantage of having a fractal distribution is the scale invariance of the measurements.

722 Analyses performed on laboratory experiments or on hand specimens could easily be

723 extrapolated to large fault systems. However, Turcotte (1992) emphasizes that not every

724 power-law distribution can be described by a fractal dimension. Fractal dimensions are

725 limited to D-values between 0 and 2 in 2D (Turcotte, 1992), so that gouge with D_{\geq} larger than

726 2.0 cannot be described as fractal.

727

728 Several authors have commented on the ideal nature of a fractal dimension or D-value of 1.6:

729 a GSD with a D-value of 1.6 appears the most ideal for close packing (Monzawa and Otsuki,

730 2003). Hoffmann and Schönert (1971) and Sammis et al. (1987) describe a D-value of 1.6 (for

731 constrained comminution) as ideal, because the probability that grains of the same size are

732 neighboring each other and the amount of pore space for a close-packed array of spherical

733 grains are minimized. Morgan (1999) and Morgan and Boettcher (1999) have shown with

- 734 numerical simulations of gouge that a D of 1.6 marks the onset of shear localization, inter-
- 735 particle rolling, and a decrease in the general comminution rate. Biegel et al. (1989) and
- 736 Marone and Scholz (1989) find a D-value of 1.6 for the onset of shear localization in their
- 737 experimentally formed gouge. It seems that the onset of gouge evolution from cracked grains,
- 738 leading to an increase of D_5 above 1.6, causes dilatancy (less close packing), which will
- 739 increase the possibility for grinding, creating more small grains and thus a higher D.
- 740 Dilatancy may allow for inter-particle rolling; in this way smaller and rounder grains are
- 741 formed, therefore D> is increased and more grains with $r \le r_K$ are formed. The smaller D-
- 742 values of gouge in high confining pressure experiments support such an interpretation,
- 743 because of the effective suppression of dilatancy at high pressures.
- 744

745 Storti et al. (2003) report measurements in carbonate fault zones, in which the D-values at the

746 boundary of the shear zone are between 0.9 and 1.4 and evolve to D-values between 1.6 and

747 2.5 for the interior shear bands. These observations are consistent with our observations:

748 gouge ($D > 1.6$) is localized in slip zones, whereas cracked grains ($D < 1.6$) make up a larger 749 part of the sample without major slip.

- 750
- 751

752 **6. Conclusions**

753

754 The analyzed fault rocks derived from experimentally faulted granite and from the Nojima 755 fault zone consist of two different types of material, cracked grains and gouge. These can be 756 distinguished on the basis of their microstructures and their grain size distribution (GSD). On

757 log-log plots the GSDs display two distinct power-law fits: D_{\leq} for grain sizes smaller than r_{K}

758 and D> for grain sizes larger than r_K . D> is the value that corresponds to other published D-

- 759 values (so-called "fractal dimensions").
- 760

761 D< is 0.9-1.0 for cracked grains and 0.9-1.1 for gouge and more or less independent of the

762 deformation conditions or the type of mineral. The lower limit of grain sizes in

763 experimentally and naturally produced gouge is $r = 15$ nm.

- 765 D> is 1.5-1.6 for cracked grains and 2.0-2.3 for gouge and depends on the deformation
- 766 conditions or the type of mineral. The upper limit of grain sizes evaluated in this study, r_{max} , is

767 100 µm. D> for gouge is a good parameter to compare natural and experimental fault rock of 768 grain sizes smaller than ∼100 µm.

769

770 Cracked grains result from initial fragmentation by rupturing. They develop into gouge by

771 subsequent grain comminution, grinding, attrition, or shear during slip along the fault zone.

772 These processes produce larger D>-values and therefore represent a more efficient grain size

773 reduction for quartz and feldspar grains in the size range greater than 1-2 µm.

774

775 The grain size, r_K , at the slope break of the log(frequency)-log(size) histogram of the GSD 776 occurs at 1.2 ± 0.3 µm for quartz and at 1.8 ± 0.3 µm for feldspar. The grain size r_K coincides

777 approximately with the grinding limit in quartz and probably corresponds to a change in the 778 physical process of grain comminution. Attrition and shear may dominate below the transition 779 value r_K .

780

781 Most of the experimental conditions during deformation (temperature, confining pressure,

782 H2O-content, displacement rate) have a minor effect on the GSD. Most noticeable is an

783 increase in confining pressure (~1 GPa), which reduces the efficiency of grain comminution. 784

785 The gouge of the Nojima Fault Zone shows $D_>= 2.02$ for gouge and 1.64 for cracked grains. 786 These values are the same as for the experimentally deformed granitoids.

787

788 Experiments and natural seismic fault rocks from the Nojima fault zone show the same GSD, 789 although the experiments have been carried out at higher PT-conditions and at considerably 790 slower displacement rates than in nature.

791

792 The development of the GSD is a two-stage process. First rupturing of the rock causes 793 cracked grains that evolve to a gouge as a result of further movement on the fault zone. The 794 surface densities calculated from the GSD of cracked minerals and gouge are 200-300 times 795 lower than published data, confirming that the creation of surface plays a minor role in the 796 energy budget of earthquakes.

- 797
- 798
- 799

800 **Acknowledgements**

801

802 The authors thank Jan Tullis and Chris Spiers for discussion and Jan for providing thin

803 sections of her Westerly Granite samples. We thank H.R. Rüegg and C. Schneider for

804 technical and A. de Ronde for scientific and general support in the laboratory. W. Tschudin

805 provided excellent thin sections. The Geological Survey of Japan is thanked for the use of

806 their thin sections of the Nojima Fault Zone. M. Düggelin has provided photographs for

807 Figures 9a and b. ZMB Basel and O. Romeyer at Savoie University in Chambéry are thanked

808 for the use of their electron microscope facilities. We are grateful to M. Boettcher and J.

809 Hazidadeh for very constructive comments, which have improved the paper. Financial

810 support for Nynke Keulen by Swiss National Fond grants nrs. 200020-100616 and 200020-

811 108082 is gratefully acknowledged.

- 813 References
- 814 Amitrano, D., Schmittbuhl, J., 2002. Fracture roughness and gouge distribution of a granite shear band. Journal 815 of Geophysical Research 107, 2375–2390.
- 816

817 An, L.-J., Sammis, C., 1994. Particle size distribution of cataclastic fault materials from southern California: A 818 3-D study. Pure & Applied Geophysics 143, 203–227.

819
820 820 Anderson, J., Osborne, R., Palmer, D., 1980. Petrogenesis of cataclastic rocks within the San Andreas Fault Zone of southern California, U.S.A. Tectonophysics 67, 221–249. Zone of southern California, U.S.A. Tectonophysics 67, 221–249.

822
823

823 Anderson, J., Osborne, R., Palmer, D., 1983. Cataclastic rocks of the San Gabriel Fault- An expression of deformation at deeper crustal levels in the San Andreas Fault Zone. Tectonophysics 98, 209–251. 824 deformation at deeper crustal levels in the San Andreas Fault Zone. Tectonophysics 98, 209–251.

825
826

826 Bass, J., 1995. Elasticity of Minerals, Glasses, and Melts. In: Ahrens, T. (Ed.), Mineral Physics & Crystallography - A handbook of Physical Constants, AGU Reference Shelf 2, 45–63. 827 Crystallography - A handbook of Physical Constants, AGU Reference Shelf 2, 45–63. 828

829 Biegel, R., Sammis, C., Dietrich, J., 1989. The frictional properties of a simulated gouge having a fractal 830 particle distribution. Journal of Structural Geology 11, 827–846. 831

832 Blanpied, M., Lockner, D., Byerlee, J., 1995. Frictional slip of granite at hydrothermal conditions. Journal of Geophysical Research 100, 13045–13064. Geophysical Research 100, 13045–13064.

- 834
835 835 Blenkinsop, T., 1991. Cataclasis and Processes of Particle Size Reduction. Pure & Applied Geophysics 136, 836 60–86. 837
- 838 Blenkinsop, T., Rutter, E., 1986. Cataclastic deformation of quartzite in the Moine thrust zone. Journal of 839 Structural Geology 8, 669–681.
- 840
841

841 Bos, B., Peach, C., Spiers, C., 2000. Slip behaviour of simulated gouge-bearing faults under conditions 842 favoring pressure solution. Journal of Geophysical Research 105, 16699–16717.

843

844 Boullier, A.-M., Fujimoto, K., Ito, H., Ohtani, T., Keulen, N., Fabbri, O., Amitrano, D., Dubois, M., Pezard, P., 845 2004a. Structural evolution of the Nojima Fault (Awaji Island, Japan) revisited from the GSJ drillhole at Hirabayashi. Earth Planets Space 56, 1233–1240. Hirabayashi. Earth Planets Space 56, 1233–1240.

847
848

848 Boullier, A.-M., Fujimoto, K., Ohtani, T, Roman-Ross, G., Lewin, E., Ito, H., Pezard,P. Ildefonse, B., 2004b. 849 Textural evidence for recent co-seismic circulation of fluids in the Nojima fault zone, Awaji island, Japan. 850 Tectonophysics 378, 165-181.

851

852 Byerlee, J., Brace, W., 1968. Stick slip, stable sliding, and earthquakes - effect of rock type, pressure, strain rate, and stiffness. Journal of Geophysical Research 73, 6031–6037. rate, and stiffness. Journal of Geophysical Research 73, 6031–6037.

854
855

Carter, N., Anderson, D., Hansen, F., Kranz, R., 1981. Creep and creep rupture of granitic rocks. In: Carter, N., 856 Friedman, M., Logan, J., Stearns, D. (Eds.), Geophysical Monograph: Mechanical behavior of crustal rocks; the 857 Handin volume, 24, American Geophysical Union, Washington DC, 61–62. 858

- 859 Chester, F., Friedman, M., Logan, J., 1985. Foliated cataclasites. Tectonophysics 111, 139–146.
- 860

861 Chester, J., Chester, F., Kronenberg, A., 2005. Fracture surface energy of the Punchbowl fault, San Andreas 862 system. Nature 437, 133–136.

863

864 De Ronde, A., 2004. Mineral Reaction and Deformation in Plagioclase-Olivine Composites: An Experimental 865 Study. PhD thesis, Basel University.

- 867 Dor, O., Ben-Zion, Y., Rockwell, T.K. Brune, J., 2006. Pulverized rocks in the Mojave section of the San 868 Andreas Fault Zone. Earth and Planetary Science Letters 245, 642-654.
- 869

870 Fabbri, O., Iwamura, K., Matsunaga, S., Coromina, G., Kanaori, Y., 2004. Distributed strike-slip faufting, 871 block rotation, and possible intracrustal vertical decoupling in the convergent zone. In Grocott, J., Tikoff, B., McCaffrey, K., Taylor, G. (Eds.), Vertical coupling and decoupling in the lithosphere. Geological Societ

872 McCaffrey, K., Taylor, G. (Eds.), Vertical coupling and decoupling in the lithosphere. Geological Society 873 London Special Publications, London. 227, 141–166.

- 874
875 Faulkner, D., Lewis, A., Rutter, E., 2003. On the internal structure and mechanics of large strike-slip fault 876 zones: field observations of the Carboneras fault in southeastern Spain. Tectonophysics 367, 235–251. 877
- 878 Griggs, D., Turner, F., Heard, H., 1960. Deformation of rocks at $T = 500$ °C to $T = 800$ °C. In Griggs, D. Handin, J. (Eds.), Geological Society of America Memoir, 79, Waverly Press, Baltimore MD, 39–104. 879 Handin, J. (Eds.), Geological Society of America Memoir, 79, Waverly Press, Baltimore MD, 39–104.
- 880
881 881 Hadizadeh, J., Johnson, W., 2003. Estimating local strain due to comminution in experimental cataclastic 882 textures. Journal of Structural Geology 25, 1973–1979.
- 883
884

884 Heilbronner, R., Keulen, N., 2006. Grain size and grain shape analysis of fault rocks. Tectonophysics 427, 199-
885 216 doi:10.1016/j.tecto.2006.05.020. 885 216 doi:10.1016/j.tecto.2006.05.020.

-
- 886
887 887 Hirth, G., Tullis, J., 1994. The brittle-plastic transition in experimentally deformed quartz aggregates. Journal 888 of Geophysical Research 99, 11731–11747.
- 889
890 890 Hoffmann, N., Schönert, K., 1971. Bruchanteil von Glaskugeln in Packungen von Fraktionen und binären Mischungen. Aufbereitungstechnik 9, 513–518. Mischungen. Aufbereitungstechnik 9, 513–518. 892
- 893 Ito, H., Fujimoto, K., Tanaka, H., Lockner, D., 1999. Proceedings of the international workshop on the Nojima
894 fault core and borehole data analysis, Geological Survey of Japan Interim report No. EO/00/1, USGS Openfault core and borehole data analysis, Geological Survey of Japan Interim report No. EQ/00/1, USGS Open-File 895 Report 000-129.
- 896
897 Jaeger, J., 1967. Failure of rocks under tensile conditions. International Journal of Rock Mechanics and Mining 898 Sciences 4, 219–237.
- 899
900 900 Kanamori, H., 1994. Mechanics of Earthquakes. Annual Review of Earth and Planetary Sciences 22, 207-237.
- 901
902 Kato, A., Ohnaka, M., Mochizuki, H., 2003. Constitutive properties for the shear failure of intact granite in 903 seismogenic environments. Journal of Geophysical Research 108, doi:10.1029/2001JB000791. 904
- 905 Kendall, K., 1978. The impossibility of comminuting small particles by compression. Nature 272, 710–711.
- 907 Korvin, G., 1992. Fractal Models in the Earth Sciences. Elsevier, Amsterdam. 908
- 909 Lin, A., 1999. S-C cataclasite in granitic rock. Tectonophysics 304, 57–273.
- 910

- 911 Lockner, D., Summers, R., Byerlee, J., 1986. Effects of temperature and sliding rate on frictional strength of granite. Pure & Applied Geophysics 124, 445–469. granite. Pure $&$ Applied Geophysics 124, 445–469.
- 913
914 914 Ma, K.-F., Tanaka, H., Song, S.-R., Wang, C.-Y., Hung, J.-H., Tsai, Y.-B., Mori, J., Song, Y.-F., Yeh, E.-C.,
- 915 Soh, W., Sone, H., Kuo, L.-W., Wu, H.-Y., 2006. Slip zone and energetics of a large earthquake from the
- 916 Taiwan Chelungpu-fault Drilling Project. Nature 444, doi:10.1038/nature05253.
- 917
918 Mair, K., Frye, K., Marone, C., 2002. Influence of grain characteristics on the friction of granular shear zones.
- 919 Journal of Geophysical Research 107, ECV 4-1, doi:2001JB000516.
- 920
- 921 Mandelbrot, B., 1982. The fractal Geometry of Nature. Freeman and Co., New York.
- 922
923
- Marone, C., Scholz, C., 1989. Particle-size distribution and microstructures within simulated gouge. Journal of 924 Structural Geology 11, 799–814.
- 925
926 926 Mitra, G., 1978. Ductile deformation zones and mylonites. the mechanical processes involved in the deformation of crystalline basement rocks. American Journal of Science 278, 1057–1084. deformation of crystalline basement rocks. American Journal of Science 278, 1057–1084.
- 928
- 929 Mitra, G., 1993. Deformation processes in brittle deformation zones in granitic basement rocks: A case study
930 from the Torrey Creek area, Wind River Mountains. In: Schmith, C., Chase, R., Erslev, E. (Eds.), Laramide from the Torrey Creek area, Wind River Mountains. In: Schmith, C., Chase, R., Erslev, E. (Eds.), Laramide
- 931 Basement Deformation in the Rocky Mountains Foreland of the Western United States. Geological Society of America Special Paper 280, 177–195. America Special Paper 280, 177–195.
- 933
934 934 Monzawa, N., Otsuki, K., 2003. Comminution and fluidization of granular fault materials: implications for fault slip behavior. Tectonophysics 367, 127–143. fault slip behavior. Tectonophysics 367, 127–143.
- 936
937 Morgan, J., 1999. Numerical simulations of granular shear zones using the distinct element method. 2. Effects 938 of particle size distribution and interparticle friction on mechanical behavior. Journal of Geophysical Research
939 104, 2721–2732. 939 104, 2721–2732.
- 940
- 941 Morgan, J., Boettcher, M., 1999. Numerical simulations of granular shear zones using the distict element
942 method. 1. Shear zone kinematics and the micromechanics of localization. Journal of Geophysical Researmethod. 1. Shear zone kinematics and the micromechanics of localization. Journal of Geophysical Research 943 104, 2703–2719.
- 944
945 945 Morrow, C., Byerlee, J., 1989. Experimental studies of compaction and dilatancy during frictional sliding on faults containing gouge. Journal of Structural Geology 11, 815–826. faults containing gouge. Journal of Structural Geology 11, 815–826. 947
- 948 Murakami, M., Tagami, T., 2004. Dating pseudotachylyte of the Nojima fault using the zircon fission-track
949 method. Geophysical Research Letters 31, doi:10.1029/2004GL020211. method. Geophysical Research Letters 31, doi:10.1029/2004GL020211. 950
- 951 Murakami, M., Tagami, T., Hasebe, N., 2002. Ancient thermal anomaly of an active fault system: zircon fission-track evidence from Nojima GSJ 750m borehole samples. Geophysical Research Letters 29, 952 fission-track evidence from Nojima GSJ 750m borehole samples. Geophysical Research Letters 29, 953 doi:10.1029/2002GL015679.
- 954
955 955 Murata, A., Takemura, K., Miyata, T., Lin, A., 2001. Quarternary vertical offset and average slip rate of the 956 Nojima Fault on Awaii Island, Japan. The Island Arc 10, 360–367. Nojima Fault on Awaji Island, Japan. The Island Arc 10, 360–367. 957
- 958 Ohtani, T., Fujimoto, K., Ito, H., Tanaka, H., Tomida, N., Higuchi T., 2000. Fault rocks and past to recent fluid
959 characteristics from the borehole survey of the Nojima fault ruptured in the 1995 Kobe earthquake, s 959 characteristics from the borehole survey of the Nojima fault ruptured in the 1995 Kobe earthquake, southwest 960 Japan. Journal of Geophysical Research 105, 16161-16172. 961
- 962 Olgaard, D., Brace, W., 1983. The microstructure of gouge from a mining-induced seismic shear zone.
963 International Journal of Rock Mechanics and Mining Science & Geomechanical Abstracts 20, 11–19. 963 International Journal of Rock Mechanics and Mining Science & Geomechanical Abstracts 20, 11–19. 964
- 965 Otsuki, K., Monzawa, N., Nagase, T., 2003. Fluidization and melting of gouge during seismic slip:
966 Identification in the Nojima fault zone and implications for focal earthquake mechanisms. Journal o
- Identification in the Nojima fault zone and implications for focal earthquake mechanisms. Journal of 967 Geophysical Research 108, ESE 4–1 – 4–18.
- 968
- 969 Pittarello, L., Di Toro, G., Bizzarri, A., Hadizadeh, J., Pennacchioni, G., 2006. Seismic energy partitioning
970 inferred from pseudotachylyte-bearing faults (Gole Larghe Fault, Adamello batholith, Italy). EGU, Vienna 970 inferred from pseudotachylyte-bearing faults (Gole Larghe Fault, Adamello batholith, Italy). EGU, Vienna. 971 Geophysical Research Abstracts 8, EGU06–A–01883.
- 972
973
- Prasher, C., (1987). Crushing and Grinding Process Handbook. Wiley, New York.
- 974
975 975 Reches, Z., Dewers, T., 2005. Gouge formation by dynamic pulverization during earthquake rupture. Earth and 976 Planetary Science Letters 235, 361–374.
-
- 977
978 978 Sammis, C., Biegel, R., 1989. Fractals, fault-gouge, and friction. Pure & Applied Geophysics 131, 255–271.
- 979
980
- 980 Sammis, C., King, G., Biegel, R., 1987. The kinematics of gouge deformation. Pure & Applied Geophysics 125, 777–812. 125, 777–812.
- 982

983 Sammis, C., Osborne, R., Anderson, J., Banerdt, M., White, P., 1986. Self-similar cataclasis in the formation of 984 gouge. Pure & Applied Geophysics 124, 53–78.

- 985
- 986 Schmid, S., Handy, M., 1991. Towards a genetic classification of fault rocks: Geological usage and tectonophysical implications. In: Müller, D., McKenzie, J., Weissert, H. (Eds.), Controversies in Mo
- 987 tectonophysical implications. In: Müller, D., McKenzie, J., Weissert, H. (Eds.), Controversies in Modern 988 Geology, Academic Press, London, 339–361.
- 989
990
- Shao, S.-M., Zou, J.-C., 1996. Fractal research of gouge. Acta Seismologica Sinica 9, 485–491.
- 991
992 Sibson, R., 1977. Fault rocks and fault mechanisms. Journal of the Geological Society of London 133, 191– 993 213.
- 994
995
- Stesky, R., Brace, W., Riley, D., Robin, P.-Y., 1974. Friction in faulted rock at high temperature and pressure. 996 Tectonophysics 23, 177–203.
- 997
998 998 Stewart, M., Holdsworth, R., Strachan, R., 2000. Deformation processes and weakening mechansms within the frictional-viscous transition zone of maior crustal-scale faults: insights from the Great Glen fault zone. 999 frictional-viscous transition zone of major crustal-scale faults: insights from the Great Glen fault zone,
1000 Scotland. Journal of Structural Geology 22, 543–560. Scotland. Journal of Structural Geology 22, 543-560.
- 1001 1002 Storti, F., Billi, A., Salvini, F., 2003. Particle size distributions in natural carbonate fault rocks; insights for 1003 non-self-similar cataclasis. Earth and Planetary Science Letters 206, 173-186. 1004
- 1005 Tanaka, H., Fujimoto, K., Ohtani, T.,Ito H., 2000. Structural and chemical characterization of shear zones in 1006 the freshly activated Nojima fault, Awaji Island, southwest Japan. Journal of Geophysical Research 106, 8789-
1007 8810, 10.1029/2000JB900444. 8810, 10.1029/2000JB900444. 1008
- 1009 Tullis, J., Yund, R., 1977. Experimental deformation of dry Westerly Granite. Journal of Geophysical Research 82, 5705–5718. 1011
- 1012 Tullis, J., Yund, R., 1980. Hydrolytic weakening of experimentally deformed Westerly Granite and Hale Albite
1013 rock. Journal of Structural Geology 2, 439–451. rock. Journal of Structural Geology 2, 439–451.
- 1014
1015 1015 Tullis, J., Yund, R., 1987. Transition from cataclastic flow to dislocation creep of feldspar: Mechanisms and nicrostructures. Geology 15, 606–609. microstructures. Geology 15, 606–609.
- $\frac{1017}{1018}$ 1018 Tullis, J., Yund, R., 1992. The brittle-ductile transion in feldspar aggregates: An experimental study. In: Evans, 1019 B., Wong, T.-f. (Eds.), Fault Mechanics and transport properties of Rocks, Academic Press, San Diego, 89– 1020 117.
- 1021
1022
- 1022 Tullis, T., Tullis, J., 1986. Experimental rock deformation techniques. In: Hobbs, B. Heard, H. (Eds.), AGU
1023 Geophysical Monograph, Mineral and rock deformation: Laboratory Studies The Paterson Volume 36, 297 1023 Geophysical Monograph, Mineral and rock deformation: Laboratory Studies - The Paterson Volume 36, 297– 1024 324.
- 1025
1026 Turcotte, D., (1992). Fractals and chaos in geology and geophysics. University Press, Cambridge.
- $\frac{1027}{1028}$ 1028 Ventakaraman, A., Kanamori, H., 1992. Observational constraints on the fracture energy of subduction zone
1029 earthquakes. Journal of Geophysical Research 109, B05302, doi:10.1029/2003JB002549. 1029 earthquakes. Journal of Geophysical Research 109, B05302, doi:10.1029/2003JB002549.
- 1030
- 1031 White, S., 1982. Fault rocks of the Moine thrust zone: a guide to their nomenclature. Textures and
- 1032 microstructures 4, 211–221.
- 1033
1034
- Wibberley, C., Shimamoto, T., 2005. Earthquake slip weakening and asperities explained by thermal 1035 pressurisation. Nature 436, 689–692.
- 1036
1037
- Wilson, B., Dewers, T., Reches, Z., Brune, J., 2005. Particle size and energetics of gouge from earthquake 1038 rupture zones. Nature 434, 749–752.
- 1039
1040 Wilson, J., Chester, J., Chester, F., 2003. Microfracture analysis of fault growth and wear processes, 1041 Punchbowl Fault, San Andreas system, California. Journal of Structural Geology 25, 1855–1873.
-
- 1042
1043 1043 Wise, D., Dunn, D., Engelder, J., Geiser, P., Hatcher, R., Kish, S., Odom, A., Schamel, S., 1984. Fault-related rocks: Suggestions for terminology. Geology 12, 391–394. rocks: Suggestions for terminology. Geology 12, 391–394.
- 1045
1046 1046 Yund, R., Blanpied, M., Tullis, T., Weeks, J., 1990. Amorphous material in high strain experimental gouge.
1047 Journal of Geophysical Research 95, 15589–15602. Journal of Geophysical Research 95, 15589–15602.
- 1048 1049
- 1050
- 1051 **Figure captions & Table captions**
- 1052
- 1053 Figure 1:
- 1054 Starting material for experiments and natural fault rocks used in study.
- 1055 (a) Undeformed sample material from the Val Verzasca, Switzerland; scanning electron micrographs,
- 1056 back scatter contrast. This gneiss is isotropic at the scale of a thin section and has a granite
- 1057 composition. L = lineation. Minerals, in order of increasing brightness: quartz (Qtz), Na-rich
- 1058 plagioclase (Plg), muscovite (Musc), K-feldspar (Kfs), biotite (Bt).
- 1059 (b) Map of the northern part of Awaji Island, Japan, showing the locality of the Hirabayashi drill site
- 1060 in the Nojima Fault [after Ohtani et al. (2000)]. Inset shows the cities Tokyo and Kobe (stars) and the
- 1061 enlarged area (black square).
- 1062 (c, d) Fault rocks of the Nojima Fault Zone (recent event); scanning electron micrographs, back
- 1063 scatter contrast. (c) Gouge with rounded siderite fragments (white arrows) and pseudotachylyte
- 1064 fragment (black arrow). 623.50 m depth. (d) Cracked grains: individual parts can still be fitted
- 1065 together. The white arrow indicates siderite grains. 633.08 m depth.
- 1066
- 1067 Figure 2:
- 1068 Schematic of procedure for grain size analysis.
- 1069 (a) Starting image: gouge of Nojima Fault Zone; scanning electron micrographs, backscatter contrast.
- 1070 (b) Binary image. (c) Histogram of equivalent radii (1285 grains). (d) Log-log plot of frequency
- 1071 versus equivalent radius (20 bins per order of magnitude). The slope of the line fit yields the D-value.
- 1072 (e) Combined log-log plot for entire range of magnification, showing two distinct average D-values
- 1073 (D_< and D_>) and r_K (see text). (f) D-values versus magnification. Note that intermediate magnifications
- 1074 yield both D-values (D< and D>).
- 1075
- 1076 Figure 3:
- 1077 Mechanical data.
- 1078 (a) Force versus axial shortening of samples of Verzasca Gneiss. For experimental conditions see
- 1079 Table 3, r-h = room-humidity. (b) Differential stress versus axial shortening, calculated using
- 1080 program rigC4 (see text). For comparison, 3 experiments on Westerly Granite by Tullis and Yund
- 1081 (1977) are included. Since pyrophyllite or talc was used as confining medium, the true P_c -values
- 1082 (indicated in brackets) are approx. 2/3 of the values given in the literature (see text).
- 1083
- 1084 Figure 4:
- 1085 Maps of local variations of grain size distribution.
- 1086 The radius of the Gauss-filter was 5 pixels (= 40 μ m); for D-mapping technique, see Heilbronner &
- 1087 Keulen (2006). D-values > 1.75 indicate gouge, D-values < 1.75 indicate cracked grains. Holes and
- 1088 biotite in samples are masked. Three samples are shown: 38nk (reference sample): d = 3 mm; 102nk:
- 1089 $d = 4.8$ mm; 60nk: $d = 2.5$ mm. Fault surfaces are indicated. Scale applies to all samples.
- 1090
- 1091 Figure 5:
- 1092 Microstructures of experimentally deformed samples of Verzasca Gneiss.
- 1093 Scanning electron micrographs, backscatter contrast. Minerals, in order of increasing brightness:
- 1094 quartz, Na-rich plagioclase, muscovite, K-feldspar, biotite.
- 1095 (a) Fault zone consisting of gouge with adjacent cracked grains (reference sample 38nk). (b)
- 1096 Relatively narrow gouge zone (sample 60nk, high confining pressure). (c) Reference sample 64nk. (d)
- 1097 Reference sample 38nk. (e) Reference sample 38nk. (f) Reference sample 38nk. (g) Cracked quartz
- 1098 grains at high magnification (high displacement sample 102nk). (h) Quartz gouge at high
- 1099 magnification (reference sample 38nk).
- 1100

1101 Figure 6:

- 1102 Grain size distribution of experimentally deformed Verzasca Gneiss plotted in a log(frequency)-
- 1103 log(equivalent radius) histogram. Results for quartz and feldspar are shown in separate graphs, results
- 1104 for cracked grains and gouge are shown with separate symbols. The frequency is normalized to 100
- 1105 grains with a radius of 10 μ m for cracked grains and to 1000 grains with a radius of 10 μ m for gouge.
- 1106 Note the slope change in most of the curves at $r \approx 2 \mu m$. The D-values (D_< and D_>) and the grain size
- 1107 at the slope change, r_K , are listed in the Table 4 and 5, respectively.
- 1108
- 1109 Figure 7:
- 1110 D-values of the gouge plotted against magnification of analyzed images.
- 1111 Average D-values are indicated by the horizontal lines. Feldspar and quartz are shown separately.
- 1112 Analyses at 1000x and 2000x magnification yield a value for $D₅$ and $D₅$, because r_K lies in the range
- 1113 of analyzed grain sizes .
- 1114
- 1115 Figure 8:
- 1116 Grain size distribution and D-values of the Nojima Fault Zone.
- 1117 (a) Grain size distribution of cracked grains and gouge (all minerals). The frequency is normalized to
- 1118 100 grains with a radius of 10 µm for cracked grains and to 1000 grains with a radius of 10 µm for
- 1119 gouge. (b) D-values versus magnification, compiled from 9 cascades (see text). The average values
- 1120 per magnification are indicated with solid symbols. Overall average values for D_{\le} and D_{\ge} are
- 1121 indicated by solid lines. The thin-sections used for the analyses originate from the following locations
- 1122 in the cores: a) HR4-1 (522.79 m), HR4-11 (633.08 m) b) HR4-11, HR2-52 (626.90 m), HR256
- 1123 (643.10 m) and HR3-10 (623.46 m).
- 1124
- 1125 Figure 9:
- 1126 Scanning electron micrographs of the smallest grains found in gouge.
- 1127 (a,c) Nojima Fault Zone. (b,d) Experiment 63nk (room-humidity sample). (a,b) Secondary electron
- 1128 mode. (c,d) backscatter contrast mode. The pockmark structure in (a) is an effect of the carbon coating
- 1129 and does not show individual atoms or another substructure.
- 1130
- 1131 Figure 10:
- 1132 Grain size, r_k , at the transition from D< to D> (slope change in GSD plots) and grinding limit for
- 1133 quartz, after Steier & Schönert (1972) in Prasher (1987). The uncertainty in the position of r_K is ± 0.3
- 1134 µm.
- 1135
- 1136
- 1137 Table 1: 1138 Published data on grain size analyses of fault rocks. 1139 Method of analysis: LM = light microscope image of thin section, SEM = scanning electron 1140 microscope image, manual = analogue box counting, digitizer = instrument for digital analysis of 1141 photomicrographs, digital = analyzed with computer, Coult.count. = laser diffraction particle size 1142 analysis. The indicated fractal dimension (D-value) is for 2-dimensional analysis. S.A.F. = San 1143 Andreas Fault, R.S.A. = Republic of South Africa. 1144 1145 Table 2: 1146 Abbreviations and definitions of measurements as used in this paper. 1147 1148 1149 Table 3: 1150 Deformation conditions of the experiments on Verzasca gneiss. 1151 T = temperature, P_c = confining pressure, \dot{e} = axial shortening rate rate, L_0 = starting length of 1152 sample; $d_{\text{chart}} =$ axial displacement from chart strip, $d_{\text{chart} - 1 \text{ mm}} =$ corrected axial displacement, $H_2O =$ 1153 amount of added water. Samples were quenched immediately after deformation. 1154 1155 Table 4: 1156 D-values (D_z and D_z) of cracked grains and gouge. 1157 In experiments, quartz and feldspar were measured separately; in natural samples (Nojima), all 1158 mineral grains were combined. 1159 1160 Table 5: 1161 r_K values determined from grain size distributions. 1162 In experiments, the values for quartz and feldspar were determined separately; in natural samples 1163 (Nojima), all mineral grains were combined. 1164 1165 1166
- 1167

Figure03 [Click here to download high resolution image](http://ees.elsevier.com/sg/download.aspx?id=32788&guid=d395aada-71e9-420a-8bf8-499920335361&scheme=1)

Figure08 [Click here to download high resolution image](http://ees.elsevier.com/sg/download.aspx?id=32792&guid=31312912-7155-4fe6-a172-b4ef0f421a95&scheme=1)

Authors	year	rock type	location	nat / exp	method of		$(2D)$ -fractal smallest diameter
					analysis	dimension	measured, μm
Sammis et al	1987 1989	\rm{granic}	Lopez Canyon	natural	$LM + SEM$, manual		
Sammis & Biegel		granite		natural	manual		
Marone & Scholz	1989	quartz sand (porous)	Ottawa	experimentsel ^a	SEM, manual	1.6	ГJ
Blenkinsop	1991	granite, gneiss	$\operatorname{SAF},$ Cajon Pass Drillhole	natural ^b	LM, digitizer LM, digitizer		
		arkose		$\operatorname{natural}^b$			
An & Sammis	1994	granite	SAF, Tejon Pass	natural	sieve + Coult. Count	$\begin{array}{c} 0.8 - 2.1 \\ 1.6 - 2.6 \\ 1.4 - 2.6 \\ 1.6 - 1.9 \\ \end{array}$	
		gneiss	San Gabriel	natural	sieve + Coult. Count. sieve + Coult. Count.		
		tonalite		natural			
Shao $&$ Zou	1996	schist and gneiss	$\begin{array}{c} \text{Lopez } \text{Canyon}\\ \text{Qining } \text{Mountain} \end{array}$	natural	$LM + SEM$, manual		
Monzawa & Otsuki	2003	granite	$\begin{array}{c} \text{Tanakura} \\ \text{It.-Shimostutaki} \end{array}$	natural	LM, manual LM, manual LM, manual LM, manual	1.7	
		granite		natural		$\frac{1}{2}$	
		$grant$ e	Koi	natural			
		$grant$ e $% \left(\left(\mathcal{N}\right) \right)$	Nojima	natural		2.3	
Hadizadeh & Johnson		quartz sandstone		experimental ^c		$1.9 - 2.5$	
Wilson et al.; $\,$		granite	$\begin{array}{c} \rm Masillon, Ohio \\ \rm SAF, Tejon \; Pass \\ \rm Bosman \; fault, \; R.S.A. \end{array}$	natural	LM, digitizer Coult. Count.		
Reches $\&$ Dewers		granite		natural	Coult. Count		
Chester et al.	2003 2005 2006 2006	granite		$\operatorname{natural}$		2.0	
Heilbronner & Keulen		granitoids	$\text{SAF}, \text{Punchbowl}$ Verzasca	\exp erimental ^d	$LM + TEM, manual$ $SEM, digital$		
This study		granitoids	Nojima	natural	SEM, digital	$1.4 - 2.3$ $1.6 - 2.3$ $1.4 - 2.3$	
		granitoids	Verzasca	\exp erimental ^e	SEM, digital		0.03

Table 1:

Symbol Definition

\sim , \sim .	\sim 0.11111.010.11
D_{\leq}	slope of power law fit for small grain sizes $(r < r_K)$
$D_{>}$	slope of power law fit for large grain sizes $(r > r_K)$
r_K	radius of grains at slope break of GSD
L_0	original length of sample
Δ_L	axial shortening of sample
r_{min}	minimum grain size used for calculation of GSD
r_{max}	maximum grain size used for calculation of GSD
A	cross sectional area of grain $(pixel^2)$
	perimeter of cross section of grain (pixel)

Table 2:

sample nr.	T	P_c	$ax.\dot{\epsilon}$	L_0	d_{chart}	$d_{(chart-1mm)}$	H_2O			
	$\rm ^{\circ}C$	MPa	$\rm x\; 10^{-4} s^{-1}$	$\,mm$	$\,mm$	$\mathscr{K}wt$				
38nk	300	510	1.5	10.44	4.01	3.0	0.2			
64nk	300	500	1.3	12.33	3.18	2.2	0.2			
102nk	300	510	1.3	10.17	5.78	4.8	$0.2\,$			
60nk	300	1030	1.5	10.13	3.47	2.5	0.2			
68nk	500	520	1.4	10.32	4.28	3.3	0.2			
63nk	300	530	1.4	12.11	4.25	3.3	0.0			
66nk	300	490	0.015	12.37	4.64	3.6	0.2			
70nk	500	500	0.013	12.95	4.97	4.0	0.2			

Deformation conditions

Table 3:

D-values for cracked grains and gouge								
	feldspar				quartz			
	cracked			gouge	cracked		gouge	
sample nr.	D_{\leq}	$D_{>}$	D_{\leq}	$D_{>}$	$D_{<}$	$D_{>}$	$D_{<}$	$D_{>}$
38nk	1.02	1.68	1.00	2.03	1.09	1.56	1.10	2.26
64nk	0.95	1.52	1.05	2.06	0.74	1.44	1.05	2.26
102nk	$0.87\,$	1.52	1.01	2.07	0.92	1.55	1.00	2.27
60nk	0.74	1.37	0.91	1.85	0.96	1.72	0.78	1.94
68nk	0.92	1.67	1.01	1.95	0.92	1.47	1.01	2.07
63nk	0.72	1.62	1.12	2.12	0.78	1.55	0.83	2.32
66nk	0.86	1.37	1.05	2.09	0.96	1.46	0.97	2.26
			Nojima Fault Zone					
	cracked gouge							
	D_{\leq}	D_{\geq}	D_{\leq}	$D_{>}$				
	0.97	1.64	0.97	2.02				

D-values for cracked grains and gouge

Table 4:

	feldspar					quartz			
		cracked		gouge	cracked		gouge		
sample nr.	$r_{\cal K}$	range	$r_K\;$	range	r_K	range	$r_K\;$	range	
38nk 64nk 102nk 60nk 68nk 63nk 66nk	1.6 $1.4 - 1.7$ 1.9 $1.7 - 2.0$ 1.8 $1.7 - 2.0$ $1.6 - 1.7$ 1.7 1.8 $1.6 - 1.9$ 1.8 $1.7 - 1.9$ 1.6 $1.3 - 1.6$		1.5 $1.4 - 1.7$ 1.6 $1.5 - 1.7$ 1.5 $1.3 - 1.7$ 1.8 $1.6 - 1.9$ 1.9 $1.9 - 2.0$ 1.5 $1.5 - 1.6$ 1.3 $1.3 - 1.6$		1.5 1.3 1.4 1.8 1.5 1.1 1.8	$1.3 - 1.6$ $1.2 - 1.4$ $1.4 - 1.5$ $1.8 - 2.0$ $1.4 - 1.6$ $1.1 - 1.3$ $1.6 - 1.9$	1.1 1.1 1.1 1.1 1.3 1.1 1.4	$0.9 - 1.1$ $1.0 - 1.3$ $0.9 - 1.1$ $1.0 - 1.2$ $1.3 - 1.5$ $1.0 - 1.3$ $1.4 - 1.6$	
	Nojima Fault Zone cracked gouge range range r_K r_K 1.9 $1.6 - 2.0$ 1.5 $1.2 - 1.6$								

 r_K : Grain radius at slope change [μ m]

Table 5: