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4	naturally deformed granitoids.
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- Grain size distributions of fault rocks: a comparison between experimentally and naturally deformed granitoids.
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- 40 Nynke Keulen, Renée Heilbronner, Holger Stünitz, Anne-Marie Boullier, Hisao Ito
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- 42 1 Abstract
- We have investigated the grain size distribution (GSD) of granitoid fault rock by comparing
- 44 experimentally produced gouge with fault rock from the Nojima Fault Zone. Triaxial experiments
- 45 were carried out on wet and dry intact samples of Verzasca Gneiss at T = 300 and 500° C, $P_c = 500$
- and 1030 MPa, $\varepsilon = 0.013 1.6 \times 10^{-4} \text{s}^{-1}$. The GSD has been determined from SEM-micrographs and is
- 47 characterized by the slope, D, of its log(frequency)-log(radius) histogram. The GSD is not fractal; we
- 48 observe two slopes for all GSDs. The larger grains in experimentally deformed samples have a D-
- value, $D_>$, of 2.04 and 2.26 for feldspar and quartz gouge. Cracked grains yield values of D = 1.5-1.6.
- 50 Increasing the confining pressure or temperature decreases the D-value. For grains smaller than ~2
- 51 μm another D-value, D_<, of 0.9-1.1 is observed. The grain radius at the slope change, r_K, corresponds
- 52 to the grinding limit of quartz, so that r_K probably represents a change in the dominant comminution
- mechanism from grinding to attrition-processes. The GSD obtained from deformation experiments
- agrees well with results for the Nojima Fault Zone: $D_> = 2.02$ for gouge and 1.64 for cracked grains;
- $D_{<}$ = 0.97. Grain size reduction in fault zones develops by a two-stage process: rupturing creates
- cracked grains; further displacement of fragments causes further comminution by wear and attrition.
- 57 Cracked grains have been used to calculate the surface energy associated with faulting; it follows that
- 58 this energy forms a small fraction in the total energy-budget of earthquakes.

60 Keywords: faulting, granite, gouge, grinding, surface density, power-law distribution

2. Introduction

62	2.1	Overview
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Major fault zones and earthquakes in the upper crust typically develop highly localized zones of deformation. At the macroscopic scale, faults usually appear as narrow, highly deformed zones or networks embedded in wide damage zones of relatively low deformation.

Observations on natural fault zones have shown that this localization of deformation along faults occurs both in space and time (e.g. Mitra, 1993; Faulkner et al., 2003; Wilson et al., 2003). Most of the displacement takes place in very narrow zones within the gouge of

typically a few to 100 µm wide, leaving the rest of the gouge inactive (e.g. Wibberley and

Shimamoto, 2005; Chester et al., 2005; Ma et al., 2006). Bos et al. (2000) have made this

same observation for laboratory experiments on halite.

Within these narrow deformation zones gouge with a wide range of grain sizes has been produced (e.g. Storti et al., 2003, Boullier et al., 2004a). At the microscopic scale, the fault rock consists of cataclasite or gouge with properties that change as a function of time and space. Displacement on a fault zone causes the evolution from a thin fracture to a wider gouge zone. The microstructures of the fault rocks influence and reflect the mechanics of faulting, i.e., earthquakes. For example, evidence for coseismic hydraulic fracturing (Boullier et al., 2004b) or for a change in the frictional properties of faults resulting from their grain shape (Mair et al., 2002) has been obtained through microstructural studies. Despite experimental data and field observations on gouge, the properties of fault rock and the evolution of cracked grains to gouge still remain rather poorly understood.

2.2 Fault rocks

A number of definitions for fault rocks can be found in the literature (e.g. Sibson, 1977; White, 1982; Wise et al., 1984; Chester et al., 1985; Schmid and Handy, 1991). Usually, cohesive and noncohesive fault rocks are attributed to cataclasis, a process involving fracturing, frictional sliding, dilatancy, and rigid body rotation between grain fragments, grains, or clusters of grains. Sibson (1977) has defined gouge as a cataclastic rock without cohesion; the cohesive rock is called cataclasite. Many cohesive cataclasites in nature are

94 thought to be the product of syn- or post-tectonic healing processes (Schmid and Handy, 95 1991).

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In this study, however, we will distinguish different fault rocks on the basis of microstructures rather than inferred mechanical properties (see Heilbronner & Keulen, 2006). We will use the term "cracked grains" for angular fragments that have fractured without much displacement (such that the original geometrical relationship of the fragments can still be recognized). The term "gouge" will be used for more highly deformed fault rocks where the grains are more rounded and the original spatial relationship between the fragments is lost. The term "fault rocks" thus applies to both cracked grains and gouge.

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2.3 Grain size distributions

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107 A large number of studies has been dedicated to measuring and analyzing grain size 108 distributions of fault rocks (e.g. Mitra, 1993; Blanpied et al., 1995; Lin, 1999; Bos et al., 109 2000; Stewart et al., 2000; Kato et al., 2003). The grain size distribution (GSD) is considered 110 as an important parameter for the characterization of the gouge (e.g. Sammis et al., 1987; 111 Marone and Scholz, 1989; Hadizadeh and Johnson, 2003; Wilson et al., 2005; Chester et al., 112 2005). Within a certain size range the GSD may be self-similar (scale invariant) and in this 113 case can be described by a power-law relationship (Mandelbrot, 1982). The slope of some 114 power-law distributions, D, is called the fractal dimension (Turcotte, 1992; Korvin, 1992, and 115 references in both). This D-value has often been used to compare the GSD of natural 116 earthquake and fault zones, like the San Andreas Fault zone (e.g. Anderson et al., 1980; 117 Blenkinsop, 1991), the Moine Thrust (Blenkinsop and Rutter, 1986) or fault systems in Japan 118 (Monzawa and Otsuki, 2003; Otsuki et al., 2003) with fault rocks derived from experiments 119 (e.g. Marone and Scholz, 1989; Morrow and Byerlee, 1989; Amitrano and Schmittbuhl, 120 2002).

- 122 For the measurements of the GSD, several different methods have been used (see Table 1).
- 123 The most classical one, originating from sedimentology, is sieving. By this method a
- 124 distribution of mass against size is obtained (Anderson et al., 1980; Sammis et al., 1986).
- 125 Other studies have used the Coulter counter (An and Sammis, 1994; Amitrano and
- 126 Schmittbuhl, 2002; Wilson et al., 2005) or optical microscopy on thin sections (Biegel et al.,
- 127 1989; Monzawa and Otsuki, 2003); these methods provide a description of the GSD in terms

128 of frequency versus grain size. A disadvantage of sieving and the use of light microscopes is 129 that the size of smallest detectable grains is $\sim 2 \mu m$ (Table 1). 130 131 An easy way to measure the sizes of the grains smaller than 2 µm is scanning electron 132 microscopy (SEM), as was done by Sammis et al. (1987) and Shao and Zou (1996), but in 133 their samples no small grains were observed. Olgaard and Brace (1983), Yund et al. (1990), 134 and Chester et al. (2005) found grains as small as 15-50 nm with the TEM, but the TEM is 135 unsuitable for measuring large numbers of grains necessary to obtain a size distribution. 136 Wilson et al. (2005) used a Coulter counter and measured grains of 40 nm to 300 μm. 137 Nevertheless, the frequency distribution of grains sized between 30 nm and 2 µm is not very 138 well known. We find that using scanning electron micrographs of impregnated and polished 139 sections of gouge, reliable GSD in the size range larger than 60 nm can be obtained. 140 141 In a previous study (Heilbronner & Keulen, 2006) two types of gouge can be distinguished: 142 cracked grains and gouge. Cracked grains and gouge can be discriminated on the basis of 143 particle shape and grain size distributions. The D-values of gouge are always higher than 144 those of cracked grains. Composition (quartz versus feldspar) influences various aspects of 145 the particle shape of cracked grains and gouge (elongation, paris factor, etc.). 146 147 Table 1 shows a compilation of published results on the GSD of fault rocks: values of the 148 fractal dimension vary between 0.8 and 2.6, but values of D between 1.5 and 2.1 are the most 149 common for granitoid rocks. High D-values are generally obtained for more recently active 150 fault zones, though a wide spread of D-values within a single fault zone can be observed. 151 152 As there appears to be a link between fault rock microstructure (shape and GSD) and the 153 mechanical behavior of faults, and since fossil fault rocks are important clues to seismic 154 activity in the past, this study focuses on the microstructures of naturally and experimentally 155 produced faults. In this paper we examine the influence of experimental conditions on the 156 fault rock microstructure and we compare the results to a natural fault rock cored from a 157 recently active segment of the Nojima Fault (Japan). The focus of the paper is on the grain 158 size distributions (D-values cracked grains and gouge) and on the geometry and spatial 159 distribution of the fault rock in the samples (D-mapping). The aim is to study the influence of 160 temperature, confining pressure, axial shortening rate, total axial shortening, and fluid content 161 on the grain size development of gouge material in deformed granitoids.

162 163 164 3. Materials and methods 165 166 167 3.1 Description of samples 168 169 3.1.1 Verzasca gneiss: 170 The starting material for the study of experimental fault rocks is cored from a gneiss that was 171 collected in the Val Verzasca (location: Swiss coordinates 704.65 - 126.30, Figure 1a). The 172 composition is granitic, consisting of ~29% plagioclase, ~27% K-feldspar, ~35% quartz and 173 ~7% mica, mainly biotite. The average grain size is ~280 µm and the rock exhibits a faint 174 foliation and lineation while appearing isotropic at the thin section scale. 175 176 3.1.2 Nojima fault rock: 177 The samples for the study of natural fault rocks are taken from the Nojima fault, which 178 transects a Cretaceous granodiorite and the overlying sediments of the Kobe Group and Osaka 179 Group (Figure 1b). The modal composition of the granodiorite is ~33% quartz, ~34% plagioclase, ~19% K-feldspar, ~11% biotite and ~4% hornblende. One year after the Hyogo-180 181 ken Nambu Earthquake (Kobe Earthquake, 17 January 1995, Mw = 7.2) three bore holes have 182 been drilled through the Nojima fault at different depths (Ito et al., 1999; Ohtani et al., 2000; 183 Tanaka et al., 2001; Murakami et al., 2002). The fault rock samples studied here (Figure 1c 184 and 1d) are from the Hirabayashi drill core, which has reached the fault zone at a depth 185 between 623.1 to 625.3 m (the thin sections are provided by the Geological Survey of Japan). 186 187 The history, as derived from the cores, can be divided into two major deformation episodes 188 (Boullier et al., 2004a). An earlier deformation episode is related to a left-lateral movement 189 on the Nojima Fault (Fabbri et al., 2004) prior to the deposition of the Middle to Late Eocene 190 Kobe Group. Pseudotachylytes associated with this episode have been dated at 56.4 Ma by 191 fission-tracks on zircons (Murakami and Tagami, 2004). A more recent Quaternary 192 deformation episode is associated with a right-lateral movement along the Nojima fault and 193 offsets the Plio-Pleistocene unconformity at the base of the Kobe Group (Murata et al., 2001). 194 The studied thin sections are from the Ouaternary deformation episode.

3.2 Rock deformation experiments

Samples of Verzasca gneiss measuring 6.39 mm in diameter and 10.1 to 12.9 mm in length (L_0) are cored parallel to the mesoscopic lineation. Table 2 lists the used abbreviations. At this size, the sample width corresponds to approximately 25 grains. After polishing their flat ends the samples are dried in an oven at 110°C and atmospheric pressure for at least 24 hours. The samples are wrapped in Ni-foil and placed in a gold jacket. In all cases except for experiment 63nk, a "room-humidity" sample, approximately 0.2 wt% H_2O (1.68 to 2.14 μ l) is added and the jacket is weld-sealed. During welding, the jacket and sample are cooled in order to avoid water loss. The exact water content of the "wet" and "room-humidity" samples has not been determined, and we assume that also in room-humidity samples some water is present at grain boundaries.

Deformation is carried out in a Grigg's solid medium deformation apparatus as described by Tullis & Tullis (1986) and De Ronde (2004), using straight furnaces, alumina pistons and sodium chloride as the confining medium. The pistons are slightly bevelled, thus reducing their diameter from 6.39 mm to 6.2 mm at the sample-piston interface. Temperature is measured with Pt/Pt-10%Rh or Cromel/Alumel thermocouples. The standard or reference deformation conditions are "wet" samples, 300°C temperature, 500 MPa confining pressure, $10^{-4} \, \text{s}^{-1}$ shortening rate, and a total axial displacement of 3.5 - 5.0 mm (samples 38nk and 64nk). In separate runs, one or more parameters are varied: 500°C temperature, 1030 MPa confining pressure, up to 5.8 mm total displacement, $10^{-6} \, \text{s}^{-1}$ shortening rate or no H₂O is added (see Table 3). During the experiment, pressure, load and displacement are recorded on a strip chart. Using the rigC4 program (www.unibas.ch/earth/micro/software/), with a correction for rig stiffness, the force-displacement and stress-strain curves are calculated.

After the experiment, the samples are quenched. To avoid loss or damage of the non-cohesive fault rocks, the jackets are punctured and the samples are impregnated with epoxy. The samples are cut along the compressional axis and the deformed length of the sample is remeasured. Thin sections are prepared (polished and carbon coated) for observation in the light and scanning electron microscope.

3.3 Grain size analysis

230 For the grain size analysis, images with a wide range of magnifications are used. The 231 micrographs of the experimental microstructures are collected using a Philips XL30 ESEM 232 scanning electron microscope (20kV acceleration voltage; backscatter contrast); those of the 233 Nojima Fault Zone are acquired with a Leica Stereoscan 440 SEM. For each sample a cascade 234 of micrographs is prepared, starting with overviews at low magnifications (100x) and 235 zooming in (to 20,000x) on selected areas with a factor of 2 to 2.5 from one magnification to 236 the next. At each magnification, three or four images are taken of each kind of fault rock 237 (quartz, feldspar, cracked grains and gouge). The total range of grain sizes (radius, r) 238 observable at these magnifications is from <20 nm to >200 µm, however, because of the 239 scarcity of measurable grains at either end of the distribution, the statistically reliable range 240 does not exceed three orders of magnitude. 241 242 We use Image SXM (http://reg.ssci.liv.ac.uk/) and ImageJ (http://rsb.info.nih.gov/ij/) and 243

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special Macros and Plug-Ins (http://www.unibas.ch/earth/micro/software) for segmentation and image analysis. First, the images are pre-processed (removal of noise) and scaled to pixels (Figure 2a). Bitmaps of quartz and feldspar grains are obtained by density slicing at the corresponding grey levels and subsequent separation of grains (Figure 2b). The cross sectional areas, A (pixel²), and the perimeters, P (pixel), of the grains are measured using the Analyze menu of the software. After removing grains with areas smaller than 20 pixel², 300 to 1200 grains per image could be used. Some of the images taken at the highest magnification could not be processed automatically; in this case the grain boundaries are traced manually.

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The resulting files of areas and perimeters were transferred to a spreadsheet program. The areas were restored by adding the pixel values of P, i.e., the band of pixels outlining the grain, to the area A within the outline. The equivalent radii, r, were calculated as

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$$r = \sqrt{((A+P)/\pi)}$$
 (1)

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259 and collected in a histogram (Figure 2c). On a log(frequency)-log(size) plot with 20 bins per 260 order of magnitude of grain size, the slope (D) of the power-law fit is determined (Figure 2d). 261 The results from all magnifications are then combined into a single plot (Figure 2e), 262 multiplying the frequency in each bin by a factor that reflects the relative magnification of the 263 images and corrects for areas in the micrograph that were not investigated (e.g. large grains at 264 edge of the micrograph, cracks). Overlapping bins of grain radii have been analyzed to check 265 the correct up-scaling; the average frequency of the overlapping bins is used to draw the 266 combined GSD plot. Figure 2f shows the individual D-values as a function of magnification. 267 268 In all samples, two distinct D-values are obtained: a shallower slope (D_<) at smaller grain 269 sizes and a steeper slope (D_>) at larger grain sizes. The change from D_< to D_> occurs at r_K. 270 The point r_K is determined in the following manner: On grain size plots derived from the 271 magnifications around r_K, the curves for D_< and D_> are drawn. The intersection of the 272 corresponding line yields r_K. 273 274 From the experimentally deformed Verzasca Gneiss, only limited fault rock material is 275 available. We therefore use only one combined log-log plot and one (average) D_> and D_< 276 value per magnification. In the case of the Nojima fault zone, we have more fault rock 277 samples. We have been able to evaluate a number of cascades from different samples of the 278 drill cores and derive 4 to 5 D_< and D_> values per magnification. 279 280 For the evaluation of the largest grain sizes, light micrographs (taken with a ZEISS Axioplan 281 Polarisation microscope) are included. The difference between the slopes (D_>) determined 282 from light and scanning electron micrographs is less than 0.05. We have also tested the 283 reliability of the semi-automatic segmentation by comparing the derived D-values to D-values 284 obtained from manual tracings of the outlines. On two sets of test images taken at 285 magnifications of 500x and 2000x, the difference in D_> is less than 0.04. 286 287 288 4. Results 289 290 4.1 Mechanical data 291 292 The results of the triaxial compression experiments are shown as force-displacement diagrams 293 (Figure 3a). With the exception of sample 60nk all samples display the same general 294 behavior. Yield strength is reached after approximately 1 mm displacement, after that the load 295 increases steadily, the differential stress decreases slowly after an axial shortening of 8-10%. 296 In some of the experiments deformed at 300°C, reaching the yield strength was accompanied

297 by audible cracking. Sample 60nk (run at the highest confining pressure) is different from all 298 others in that it has produced distinct acoustic events, which accompany load drops on the order of 20 kN (corresponding to approximately 500 MPa). 299 300 301 At 45kN (1000-1200 MPa), sample 70nk, deformed at relatively high temperature (500°C) and low axial shortening rate (10⁻⁶ s⁻¹), is the weakest. At 60-75 kN (1750 MPa) sample 63nk 302 303 (room humidity sample) is the strongest. Samples deformed at standard conditions (64nk and 304 38nk) support loads of 55-65 kN, corresponding to strengths of 1300-1500 MPa. 305 306 For comparison with published data, we also show a stress-strain diagram (Figure 3b). The 307 differential stress is calculated with a correction for an increase in cross-sectional area, 308 assuming constant volume. The Verzasca Gneiss has approximately the same strength as 309 Westerly Granite (Fig. 3b) deformed by Tullis and Yund (1977). The main differences with 310 respect to Westerly Granite experiments are the higher shortening at yield stress (with typical values $\sim 10\%$ versus $\sim 5\%$), the higher shortening rate (mainly at 10^{-4} s⁻¹, see Table 3), and the 311 312 amount of added water with which the Verzasca Gneiss has been deformed. In both sets of 313 experiments, a temperature increase from 300°C to 500°C has caused a strength reduction of 314 approximately 200 MPa. Note that Tullis and Yund (1977) have used pyrophyllite or talc as 315 confining medium, which means that the true confining pressures probably are $\sim 2/3$ of the 316 published values (J. Tullis, pers. comm., 2004). 317 318 We have noted a difference between the measured lengths of the deformed samples and the 319 deformed lengths calculated from the total axial displacements derived from the force-320 displacement records (compare Figure 3a and Table 3). The calculated displacements exceed 321 the measured displacements by up to 1 mm. In the rest of the paper, we will therefore use the 322 more conservative values of the measured displacements, not the displacements shown in 323 Figure 3. This difference in measured length is not critical for the results, as the displacement 324 is used semi-quantitatively in this study. 325 326 4.2 Faulted samples 327 328 In most samples deformed at T = 300°C and $P_c = 500$ MPa, through-going fault zones are 329 formed accommodating the axial displacement within a single gouge zone. At T = 500°C, P_c

= 500 MPa and occasionally at T = 300° C, $P_c = 500$ MPa, conjugate fault zones have

developed and partitioned the axial displacement over at least two fault zones, leading to a smaller displacement on any given fault. Sample 60nk deformed at P_c = 1030 MPa shows a set of parallel faults. The corresponding load-displacement curve (Figure 3) shows an equivalent number of load drops. It is likely that each drop in the differential stress is related to the creation of a new fault plane.

The samples are deformed along relatively wide fault zones of 1.0-1.4 mm width, which become narrower with higher confining pressure or higher temperature; with increasing displacement, they tend to broaden. Most of the gouge zones are formed around mica grains or the connection of two mica grains. At a temperature of 300°C and shortening rates of 10⁻⁴ s⁻¹ biotite deforms by gliding, kinking and occasionally by fracturing.

Using the D-mapping technique described elsewhere (Heilbronner & Keulen, 2006), maps of local variations of the grain size distribution are prepared for three of the experimental samples (Figure 4). Cracking and gouge formation occur throughout the samples. In samples 38nk and 102nk gouge regions coalesce; in sample 60nk there is less coalescence. The volume proportion of gouge formed during the deformation is determined by measuring the black area (indicating D > 1.75) in Figure 4 and dividing this proportion by the total sample area. For increasing axial displacement, d, increasing amounts of gouge are formed. After d = 4 mm, sample 38nk (reference sample) has 24.5% of gouge; after d = 5.8 mm, sample 102nk (deformed at the same conditions) has 54.8%. At higher confining pressures, there is less gouge: after d = 3.5 mm, sample 60nk ($P_c = 1000 \text{ MPa}$) has only 15.5% gouge.

4.3 Microstructures

In the experiments at $P_c = 500$ MPa and $T = 300^{\circ}$ C, the original grains in the gouge zone form broad and elongated arrays of fractured particles (Figure 5a); at higher confining pressures, these arrays are much narrower (Figure 5b). Quartz and feldspar are both strongly fractured; at higher confining pressures, quartz is less intensively fractured than feldspar. A variation in grain shapes can be observed: quartz grains are more elongated, feldspar grains are more irregularly shaped (Figure 5c & 5d); the same variation in shape has been observed for a sample deformed at low PT-conditions (Heilbronner and Keulen, 2006).

Within the gouge zones, quartz and feldspar are fractured to very small grain sizes in all experiments (Figure 5e & 5f). The gouge is equidimensional and subangular in shape. In the cracked grains and the gouge, the grains smaller than 2 µm diameter do not show any intragranular fractures, while grains larger than 10 µm are usually fractured into a spectrum of smaller grains (Figure 5g & 5h). Also, grains smaller than 1 µm are equiaxial and appear weakly indented. Within the gouge these small grains show more rounded edges.

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4.4 Grain size analysis of experimental fault rocks

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4.4.1. Dependence of grain size distribution on type and composition of fault rock

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The grain size distributions of all of our experiments display the same general characteristics

(Figure 6). On log-log plots of frequency versus equivalent radius, two distinct parts can be

discriminated (compare also Heilbronner & Keulen, 2006). Both for cracked grains and

gouge, two D-values are obtained: a lower value, D<, for small grain sizes and a higher value,

 $D_{>}$, for larger grain sizes. The grain size, r_{K} , at the intersection of the two curve fits is on the

order of 1 - 2 μ m. In gouge, the lower range of grain sizes extends from $r_{min} \approx 20$ nm to $r_K \approx 2$

382 µm with $D_{<} \approx 1.0$ and the upper range extends from r_{K} to $r_{max} \approx 20$ µm with $D_{>} \approx 2.0$. In

cracked grains, the lower range of grain sizes extends from $r_{min} \approx 50$ nm to $r_K \approx 1.5 \mu m$ with

 $D_{<} \approx 1.0$ and the upper range extends from r_{K} to $r_{max} \approx 100 \,\mu m$ with $D_{>} \approx 1.5$.

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Comparing feldspar and quartz we note that the D_< values are approximately identical under

almost all experimental conditions ranging from 0.72 to 1.02 in cracked feldspar, 0.74 to 1.09

in cracked quartz, from 0.91 to 1.12 in feldspar gouge and 0.78 to 1.10 in quartz gouge (Table

389 4). The D_> values, in contrast, are usually higher for quartz than for feldspar ranging from

1.37 to 1.68 in cracked feldspar, 1.44 to 1.72 in cracked quartz, from 1.85 to 2.12 in feldspar

gouge and 1.94 to 2.32 in quartz gouge (Table 4).

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On average, the grain size, r_K , of cracked grains is slightly larger than in gouge (Table 5).

However, this effect is more pronounced in quartz than in feldspar. The grain size, r_K, of

cracked quartz ranges from $1.1 - 1.8 \mu m$ and is smaller than r_K of cracked feldspar which

ranges from $1.6 - 2.0 \mu m$. Similarly, in gouge, r_K of quartz is $1.1 - 1.4 \mu m$ while r_K of feldspar

397 is $1.4 - 2.0 \, \mu m$.

399 4.4.2. Dependence of grain size distribution on experimental conditions 400 401 The GSD of the experimentally deformed granitoid samples are shown in Figure 6 & 7, the 402 derived values of D_<, D_>, and r_K are shown in Tables 4 & 5. The d_K grain diameter is two 403 times the r_K grain radius indicated in Figure 6, and has been used to compare our data to other 404 studies. We consider the samples 64nk and 38nk as reference samples. The experiments are performed at T = 300°C, confining pressure $P_c = 500$ MPa, shortening rate $\varepsilon = 1.5 \cdot 10^{-4}$ s⁻¹ 405 and 0.2 %wt H_2O was added to the samples. The total axial displacement, ΔL , of 64nk and 406 407 38nk is 2.2 and 3 mm, respectively. In Figure 6, each of the plots shows the GSDs resulting 408 from the change of one of the experimental conditions with respect to 64nk. 409 410 The reference values derived from the GSDs of sample 64nk and 38nk are the following: D< 411 ranges from 0.95 to 1.07 ($D_{<}$ of cracked feldspar = 1.00; $D_{<}$ of feldspar gouge = 1.02; $D_{<}$ of 412 cracked quartz = 0.95; D_< of quartz gouge = 1.07), cracked grains have somewhat lower 413 values than gouge. r_K is $1.1 - 1.8 \mu m$ (cracked feldspar = 1.8 μm ; r_K of feldspar gouge = 1.6 414 μm ; r_K of cracked quartz = 1.6 μm ; r_K of quartz gouge = 1.1 μm), quartz yields lower values 415 than feldspar and cracked grains have higher values than gouge. $D_{>}$ is 1.5 - 1.6 for cracked 416 grains (D_> of cracked feldspar = 1.6; D_> of cracked quartz = 1.5) and 2.04 - 2.26 for gouge 417 $(D_{>})$ of feldspar gouge = 2.04; $D_{>}$ of quartz gouge = 2.26), quartz has a lower $D_{>}$ -value than 418 feldspar for cracked grains and a higher D>-value for gouge. 419 420 Increasing the confining pressure from 500 MPa to 1030 MPa leads to a marked reduction of 421 both D< and D>, both for cracked grains and gouge. The grain sizes, r_K, at the slope break of 422 cracked grains and gouge do not change significantly, with the possible exceptions of r_K of 423 feld spar gouge and cracked quartz: these values increase from 1.6 to 1.8 µm and from 1.5 to 424 1.8 µm, respectively. This variation is within the error of the measurements. 425 Reducing the shortening rate from 10^{-4} s⁻¹ to 10^{-6} s⁻¹ does not affect the values of D_>; of the D_< 426 427 values only those of cracked feldspar and of quartz gouge appear to be lowered: from 1.00 to 428 0.86 and 1.07 to 0.97, respectively. The grain size r_K of feldspar remains unchanged, while 429 that of cracked quartz and quartz gouge increases from 1.5 to 1.8 µm and from 1.1 to 1.4 µm 430 respectively. 431

432 Increasing the axial shortening from 2.2 mm to 4.8 mm has no significant effect on any of the 433 measured values. As has been mentioned before, the only difference between samples 64nk 434 $(\Delta L = 2.2 \text{ mm})$, 38nk $(\Delta L = 3.0 \text{ mm})$ and 102nk $(\Delta L = 4.8 \text{ mm})$ is the width of the fault 435 zone(s), i.e., the amount of fault rock (cracked grains and gouge) created during deformation. 436 437 Deforming the sample without adding water (under room-humidity conditions) has a marked 438 effect on the GSD. With the exception of feldspar gouge, all D_< values are lowered 439 significantly to values as low as 0.72 to 0.83. At the same time, all values of D> increase. The 440 grain size at the slope break is not affected except for r_K of cracked quartz, which is lowered 441 from 1.5 to $1.2 \mu m$. 442 443 Increasing the temperature from 300°C to 500°C affects the D_> values. In particular, the 444 values for feldspar gouge and quartz gouge are lowered from 2.04 to 1.95 and from 2.26 to 2.07, respectively. The major effect on r_K is to increase the values for feldspar and quartz 445 446 gouge from 1.6 to 2.0 µm and from 1.1 to 1.4 µm, respectively. 447 448 4.5 Grain size analysis of natural fault rocks 449 450 For the Nojima samples the GSDs are given as bulk results, i.e., the different mineral phases, 451 in particular, quartz and feldspar, were not separated. Comparing the GSDs of cracked 452 minerals and gouge to the experimental results, we note that the values of D_{\leq} (1.64 and 2.02), 453 $D_{>}$ (0.97) and r_{K} (1.9 µm and 1.5 µm) are very close to the corresponding values obtained for 454 feldspar cracked grains and gouge (see Table 4, Figure 8). 455 456 457 4.6 Characteristics of GSDs 458 459 The GSDs of both the experimental and natural fault rocks display a number of characteristics 460 that we want to highlight (compare Figure 6, 7 & 8). The lower limit of grain sizes that could 461 be used for the calculation of statistically reliable distributions is $r_{min} = 30$ to 60 nm for 462 experimentally produced and $r_{min} = 100$ to 200 nm for naturally produced gouge and cracked 463 grains, respectively. Similarly, the upper limit is $r_{max} = 50 - 100 \mu m$ for experimentally 464 produced and $r_{max} = 100 - 200 \mu m$ for naturally produced gouge and cracked grains,

respectively. In other words r_{min} and r_{max} of gouge are always smaller than r_{min} and r_{max} of 465 466 cracked grains by a factor of ~ 2 . Note, however, that r_{min} is not the absolutely smallest grain 467 size that can be detected in the SEM (Figure 9). Grains as small as $r \approx 5$ nm have been 468 detected both in experimental and in natural fault rocks, but are at the detection limit of the 469 SEM so that their number is not large enough for a grain size analysis. 470 471 The grain size r_K at the slope break of the power law fit to the GSD has a remarkably constant 472 value. In contrast to the marked decrease of minimum and maximum grain size from cracked 473 grains to gouge, the r_K values remain approximately constant or decrease by < 30% from an 474 average value of 1.7 to 1.4 μm in experimental and from 1.9 to 1.5 μm in natural fault rocks 475 (Figure 10). 476 477 478 5 Discussion 479 480 5.1 Faulting experiments 481 482 Brittle deformation experiments on granitoid rock at elevated pressure-temperature conditions 483 have been performed by Griggs et al. (1960); Stesky et al. (1974); Carter et al. (1981); 484 Blanpied et al. (1995); Kato et al. (2003) and others. Our experiments on Verzasca Gneiss 485 compare best with the data on Westerly Granite published by Griggs et al. (1960) and Tullis 486 and Yund (1977, 1980). The geometry and distribution of strain in our samples is variable 487 (Figure 4). Therefore, we prefer to present our mechanical data in terms of force versus axial 488 displacement rather than stress and strain (Figure 3a). For comparison with published data, we 489 have converted our force data to stress data (Figure 3b). 490 491 The experiments on Westerly Granite at $P_c = 500$ and 750 MPa have been preformed without 492 added water; samples have been pre-dried at a higher temperature (300°C) than ours (Tullis & 493 Yund 1977). Since pyrophyllite or talc has been used as outer confining medium, the 494 confining pressure has been reduced to approximately 350 and 500 MPa. Their experiments at 495 750 MPa are therefore best compared to ours at 500 MPa, because we use NaCl as outer

confining medium. The strengths of Westerly Granite and Verzasca Gneiss are comparable

showing strengths of 1200 to 1600 MPa under similar conditions, although no experiments

have been performed under identical conditions. The variation of the bulk shortening rate

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499 between 10⁻⁴ and 10⁻⁶ s⁻¹ does not show any mechanical effect, consistent with the results of 500 Tullis & Yund (1980, 1992). 501 502 503 The effect of a high confining pressure (~1 GPa) is comparable to experiments by Tullis & 504 Yund (1977, 1980). Very narrow slip zones accommodate the axial shortening. According to 505 Tullis and Yund (1977) and Hirth and Tullis (1994), the formation of these narrow zones and 506 the decrease in volume of fault rock is related to the transition from mode I to mode II 507 cracking above $P_c \approx 600-800$ MPa. The number of these narrow zones corresponds to the 508 number of audible cracking events in the experiment 60nk. This observation suggests that the 509 mechanical behavior is not a case of stick-slip as described in the literature (e.g. Byerlee and 510 Brace, 1968; Stesky et al., 1974; Lockner et al., 1986). Instead, the apparent stick slip 511 character of the mechanical data (Figure 3) is produced by multiple rupture events, 512 demonstrating the great strength of already ruptured material at high confining pressure. 513 514 5.2 Grain size distributions obtained in experiments 515 516 Changing the physical parameters of experimental deformation also affects the grain size 517 distribution. The effects are discussed with reference to sample 38nk and 64nk, which have 518 been deformed under identical temperature, pressure and shortening rate conditions. 519 520 The largest effect is observed for increasing the confining pressure to ~1 GPa. It causes a 521 reduction in the D_>-values (Figure 6, 7, Table 4) and a decrease in the total amount of gouge 522 and cracked grains in the fault zone (see Figure 4). This smaller amount of gouge is related to 523 the transition from mode I to mode II cracking: mode I fracturing is dilatant; in mode II 524 fracturing sharp shear fractures are formed when dilatancy is suppressed. Dilatancy 525 suppression appears to inhibit the formation of large amounts of surface (small grains) and 526 thus may reduce the D>-value. 527 528 At higher temperatures, there is a decrease in the D_> values of the gouge (Figure 7, Table 4). 529 In feldspar deformed at $T = 500^{\circ}$ C we observe undulatory extinction and deformation bands,

which are developed by micro-cracking at a very small scale, as described by Tullis and Yund

(1987). At a temperature of 500°C, Hirth & Tullis (1994) have found an increased dislocation

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activity in quartz. The less efficient grain size reduction at elevated temperature may be caused by a semi-brittle behavior.

The effects of not adding water and slower shortening rate are either unsystematic (higher D_<values for feldspar and lower D_< values for quartz in gouge of non-water added sample 63nk; Table 4) or not existent (samples 38nk, 64nk, 102nk; Table 4). The D-values are the same regardless of the amount of shortening of the sample (Figure 7, Table 3).

After 4 mm axial shortening a gouge of 1.0 mm thickness has been formed. Once a mature zone with gouge is established, this zone is growing wider at greater displacement, but $D_>$ of the gouge is not changing. Marone and Scholz (1989) have also observed similar D-values for different shear strains. In general, it can be concluded that the effects of changing the physical conditions of deformation on the grain size distribution are of minor importance except for the increase in confining pressure, and, to a lesser extent, temperature.

5.3 Cracked material and gouge

All experimental and Nojima fault rock samples show the same type of grain size distribution (GSD). For small grain sizes (d < 2 - 3 μ m) the GSD is characterized by D-values of ~1, whereas for the larger sizes the D-values range between ~1.5 and ~2.3 (Table 4). The gouge always shows higher D-values than the cracked material (Figure 6, 8, Table 4). It is interesting to note that the values of D< in cracked material and gouge, in quartz and feldspar, and in experimental and in Nojima fault rocks all have very similar values and the same range of variation (Table 4). Furthermore, this grain size range is not affected by changing the experimental deformation conditions. In contrast, D> of gouge is the parameter that has been measured in several other studies before (Table 1). It is the most suitable of the D-values to compare natural fault zones in different regions or to compare natural faults with experimentally obtained values.

Higher D-values indicate a more efficient grain size comminution producing a larger number of smaller grains. The difference between gouge and cracked grains clearly demonstrates that the initial rupturing process does not produce the full amount and size fraction of small grains observed in gouge (Figure 6, 8). Instead, slip on the fault zone causes further comminution of grains in the gouge and an increase in D_>. Thus, it can be concluded that after the initial

fragmentation of the rock by rupturing there is a post-rupture processing during slip which takes place in the gouge to produce the higher values of D_>.

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Further comminution of grains in gouge has been observed by Sammis et al. (1986, 1987) and

Hadizadeh & Johnson (2003). However, any quantitative characterization of fault

displacement by GSD measurements appears impossible because an increase in sample

shortening from 2.2 to 4.8 mm in our experiments (64nk and 102nk) does not change any of

the D-values systematically (Figure 6, 9, Table 5). Thus, the more efficient grain

comminution in gouge must take place in the first stages of the fault displacement after

575 rupturing.

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In the experiments performed under identical conditions (38nk, 64nk, 102nk) the variation of

D-values of the cracked material is always larger than in the gouge (Table 4). The only

systematic difference appears to be the consistently higher D_> values of quartz gouge

compared to feldspar gouge, i.e., for the coarser grain size fraction (> 2 µm) quartz shows a

more efficient comminution of grains than feldspar. The D_>-values for quartz and feldspar

gouge are almost identical in these three experiments and indicate that this D-value is very

well reproducible and therefore forms a reliable parameter to compare the GSD of fault zones.

Another difference is the lower r_K-value for quartz compared to feldspar. Thus, the change in

585 the GSD occurs at $\sim 1 \mu m$ in quartz, while it occurs at $\sim 2 \mu m$ in feldspar (Table 4).

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5.4 Natural fault rock

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The range of values in the natural samples from the Nojima fault is approximately the same as

those of the feldspars in experiments (Figure 6, 8, Table 4) so that the natural fault rock

appears to behave in a similar way as feldspar. The granodiorite deformed by the Nojima fault

consists of ~53% feldspar (plagioclase and K-feldspar); therefore, feldspar is expected to play

a dominant role in the GSD of a mixed gouge.

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The recently formed fault rock of the Nojima Fault zone shows the same D-values as the

experimentally deformed granitoid rock, especially for feldspar minerals. If we compare the

D-values of the recent deformation episode with feldspar in experiment 64nk and 38nk,

performed at T = 300°C, $P_c = 500 \text{ MPa}$, $\varepsilon = 1.5 \cdot 10^{-4} \text{ s}^{-1}$, with < 4 mm shortening and with 0.2

599 wt% H₂O added to the sample, all D-values are the same within the error of 0.05. 600 Temperature and pressure conditions during Nojima deformation were lower than in the 601 deformation experiments. 602 603 $D_{>}$ of the Nojima gouge ($D_{>} = 2.02$) is in fairly good agreement with data from coarser grain 604 size fractions by Monzawa and Otsuki (2003); they obtain $D_> = 2.192$ to 2.559, for five 605 different samples of the surface outcrop of the Nojima Fault Zone, with a mean of 2.347. 606 Monzawa and Otsuki's (2003) data is more similar to the values of D_> for the experimentally 607 deformed quartz so that all data appears to be within the variation of the experimental data. 608 609 Our results are consistent with studies of other fault zones (Table 1). Blenkinsop (1991) and 610 An and Sammis (1994) describe both low ($D_> < 1.6$) and high ($D_> > 2.0$) values for analyses 611 of gouge originating from the same fault zone. Several other studies describe high D_> values 612 (e.g. Blenkinsop, 1991; Hadizadeh and Johnson, 2003). In some gouge the lower limit of D 613 is 1.4 to 1.6, except for the San Andreas Fault Cajon Pass Drillhole granite (D = 0.8; 614 Blenkinsop, 1991). We speculate that these gouge with D_> below 1.6 might be monomineralic 615 gouge where fluids and temperature caused healing (e.g. Sammis et al., 1987, Keulen et al., 616 submitted) or were only cracked (e.g. Marone and Scholz, 1989). 617 618 5.5. Surface densities of fault rocks 619 620 The surface energy in gouge can be used to calculate a part of the energy release during an 621 earthquake event (Kanamori, 1994; Wilson et al., 2005; Chester et al., 2005; Ma et al., 2006). 622 The grain size distribution allows the calculation of the total surface area of the fault rock 623 (Chester et al. 2005, Wilson et al. 2005). The estimates of Chester et al. (2005) for the 624 Punchbowl fault are based on GSDs with an assumed minimum grain size, r_{min}, of 1.6 nm and 625 a constant D-value of 3 (= 2 in 2 dimensions) for the entire grain size range. In contrast, we 626 find that in all of our experimental and natural samples the D-value below $r \approx 1$ to 2 μ m is D_< 627 $\approx 1.$ 628 629 If the evolution of the GSD is a two-stage process involving initial rupturing and further grain 630 comminution in the gouge, probably by attrition, shearing, grinding, etc.; it is questionable

whether the surface energy calculated from the GSD can directly be related to the seismic

energy release as proposed by Wilson et al. (2005) and Chester et al. (2005). Neither is the

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633 surface energy obtained from the GSD of fault rocks identical to the specific surface energy 634 as applied in earthquake energy budget calculations (e.g. Ventakaraman & Kanamori, 2004). 635 For the following calculation we use the surface density of cracked grains, but not of gouge, 636 to compare to the data of Wilson et al. (2005) and Chester et al. (2005), because we assume 637 that this is the material, which is created by the rupture process during the seismic event, while the higher surface density of gouge is due to the combined effect of fracturing and 638 639 subsequent deformation during slip. 640 641 To estimate the surface density for our cracked minerals and gouge we use the concept of the 642 fractal fragmentation of the cube and a simple spread sheet program. We set up a histogram of 643 13 (logarithmic) bins, extending from 0.03 μm to 128 μm and from 0.015μm to 64 μm, 644 respectively. Between r_{min} and r_K, the 3-dimensional fractal dimension is set to 2 in both 645 cases, corresponding to a value of $D_{<} = 1.0$. Between r_{K} and r_{max} , the fractal dimension is set 646 to 2.5 for cracked material and to 3 for gouge, corresponding to $D_> = 1.50$ and $D_> = 2.00$, 647 respectively. A value of $r_K = 1 \mu m$ is assumed in both cases. For each of the bins, the number, 648 total volume and total surface of the cubes is calculated. Summing over all bins, the surface to volume ratio of cracked material and gouge is 1.3 and 6.0 $\mu m^2/\mu m^3$ i.e., $1.3\cdot 10^6~m^{-1}$ and 649 650 6.0·10⁶ m⁻¹ for cracked material and gouge. 651 652 Using the same approach to calculate surface densities for the Punchbowl Fault using the data $(r_{min} = 1.6 \text{ nm} \text{ and } r_{max} = 100 \text{ } \mu\text{m}, D = 2.0)$ given by Chester et al. (2005), we obtain $6.8 \cdot 10^8$ 653 m⁻¹. If we now compare our value of cracked grains of 1.3·10⁶ m⁻¹ to a value of 4.2·10⁸ m⁻¹ 654 (corresponding to 4.2·10⁵ m² in the 1mm thick layer of the fault, given by Chester et al., 655 2005) or 2.7·10⁸ m⁻¹ (corresponding to 80 m²g⁻¹ given by Wilson et al., 2005) we note that in 656 657 our case, the surface density, and accordingly, the energy required to generate it, is smaller 658 than the values assumed by Wilson et al. (2005) and Chester et al. (2005) by a factor of 200 to 659 300. Finding that the surface densities in cracked rocks and even in gouge are so low suggests 660 to us that the creation of those surfaces cannot play a major role in the energy budget of a 661 seismic event. 662 663 The discrepancy between our data and the data published by Chester et al. (2005) can be 664 explained by the fact that they have measured their small grain size fraction (down to r = 25665

nm) in the TEM, which is a notoriously difficult method for obtaining good statistics of any

parameter observed. Thus, they did not detect the change of D-values at $r \approx 1 \mu m$. Using D =

2.0 over the entire range of grain sizes results in a total surface area which is 2 orders of magnitude too high. Furthermore, they assume $r_{min} = 1.6$ nm and thus many tiny grains with a high surface area density. Note that finding lower surface densities rather supports and strengthens their general conclusion that the surface energies do not contribute in any major way to the energy dissipated during earthquakes. Pittarello et al. (2006) come to the same conclusion, using GSD measurements of pseudotachylyte-related granitoid cataclasites.

5.6 Grinding limit

Grains with very small radii formed by comminution have been found before in deformation experiments (d = 10 nm; Yund et al., 1990), in mining induced fault zones (~40 nm; Olgaard and Brace, 1983; Wilson et al., 2005) and in natural fault zones (40 nm; Wilson et al., 2005, Chester et al. 2005). We observe grains of 30 nm in both experimental and natural fault rock (Figure 9). r_K is 1.2 ± 0.3 μ m for quartz and 1.8 ± 0.3 μ m for feldspar grains. The slope change is observed for both cracked grains and gouge (Figure 6, 8, Table 5). Very small grains are already formed upon cracking of the grains; they are not only a product of gouge formation. $D_{<}$ is less dependent on deformation conditions, the type of mineral (quartz or feldspar) and is more similar for cracked grains and gouge than $D_{>}$.

Kendall (1978) has demonstrated that a critical radius exists below which particles cannot be comminuted further by grinding. This critical radius, called the grinding limit, is 0.3 μm for SiO₂ glass, 0.9 μm for quartz and 2.2 μm for calcite (Steier and Schönert, 1972 in Prasher, 1987). Small particles have a different internal stress distribution and are stronger due to smaller flaw sizes and flaw size density (Prasher, 1987) so that very high stresses are required to break them (Mitra, 1978). This is consistent with our observation that no intragranular cracks within grains smaller than 1 μm in radius are observed.

The r_K -values for quartz are slightly higher than the grinding limit (given as radius) for quartz (Figure 10). There is no data for the feldspar grinding limit available but the higher r_K -values for feldspar (Figure 10) are consistent with the observation that the critical grain size for grinding is linearly dependent on the elastic modulus E of the mineral (Kendall, 1978), because the bulk E-modulus of feldspar is approximately 2 times that of quartz (Bass, 1995).

Thus, it appears likely that r_K represents approximately the grinding limit in the comminution of fault rock. For smaller grain sizes, where grinding is inhibited, communition is less efficient, leading to a lower $D_{<}$. Grinding is performed by compressional fracturing of grains: a network of cracks will develop and the grain will break into several smaller ones (Jaeger, 1967). Other processes, like shearing or attrition also produce comminution. Reches and Dewers (2005) and Dor et al. (2006) describe comminution as dynamic pulverization (explosive granulation followed by dynamic contraction) at earthquake crack tips. For grains smaller than r_K compressional fracturing appears inhibited, but shearing or attrition remains significant. From the difference between D-values in cracked grains and gouge, it is clear that grinding, shearing and attrition during slip cause further comminution in gouge.

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5.7 Fractal dimension

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Fractal dimensions are used to describe geometries that are scale invariant (Turcotte, 1992). To describe a GSD as fractal, this distribution should lie on a straight line over several orders of magnitude in a log(frequency)-log(size) distribution. The GSD of both the experiments and the Nojima Fault Zone gouge does not show a fractal distribution. The data cover 3.5 orders of magnitude, but cannot be fitted with a single straight line. Furthermore, there are statistical problems because most fault rock grain size measurements only cover 1 to 4 orders of magnitude.

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- The advantage of having a fractal distribution is the scale invariance of the measurements.
- Analyses performed on laboratory experiments or on hand specimens could easily be
- extrapolated to large fault systems. However, Turcotte (1992) emphasizes that not every
- power-law distribution can be described by a fractal dimension. Fractal dimensions are
- limited to D-values between 0 and 2 in 2D (Turcotte, 1992), so that gouge with $D_{>}$ larger than
- 726 2.0 cannot be described as fractal.

- Several authors have commented on the ideal nature of a fractal dimension or D-value of 1.6:
- a GSD with a D-value of 1.6 appears the most ideal for close packing (Monzawa and Otsuki,
- 730 2003). Hoffmann and Schönert (1971) and Sammis et al. (1987) describe a D-value of 1.6 (for
- 731 constrained comminution) as ideal, because the probability that grains of the same size are
- 732 neighboring each other and the amount of pore space for a close-packed array of spherical
- grains are minimized. Morgan (1999) and Morgan and Boettcher (1999) have shown with

734 numerical simulations of gouge that a D of 1.6 marks the onset of shear localization, inter-735 particle rolling, and a decrease in the general comminution rate. Biegel et al. (1989) and 736 Marone and Scholz (1989) find a D-value of 1.6 for the onset of shear localization in their 737 experimentally formed gouge. It seems that the onset of gouge evolution from cracked grains, 738 leading to an increase of D_> above 1.6, causes dilatancy (less close packing), which will 739 increase the possibility for grinding, creating more small grains and thus a higher D>. 740 Dilatancy may allow for inter-particle rolling; in this way smaller and rounder grains are 741 formed, therefore D_r is increased and more grains with $r \le r_K$ are formed. The smaller D_r 742 values of gouge in high confining pressure experiments support such an interpretation, 743 because of the effective suppression of dilatancy at high pressures. 744 745 Storti et al. (2003) report measurements in carbonate fault zones, in which the D-values at the 746 boundary of the shear zone are between 0.9 and 1.4 and evolve to D-values between 1.6 and 747 2.5 for the interior shear bands. These observations are consistent with our observations: 748 gouge (D > 1.6) is localized in slip zones, whereas cracked grains (D < 1.6) make up a larger 749 part of the sample without major slip. 750 751 752 6. Conclusions 753 754 The analyzed fault rocks derived from experimentally faulted granite and from the Nojima 755 fault zone consist of two different types of material, cracked grains and gouge. These can be 756 distinguished on the basis of their microstructures and their grain size distribution (GSD). On 757 log-log plots the GSDs display two distinct power-law fits: D_< for grain sizes smaller than r_K 758 and D_> for grain sizes larger than r_K. D_> is the value that corresponds to other published D-759 values (so-called "fractal dimensions"). 760 761 D_< is 0.9-1.0 for cracked grains and 0.9-1.1 for gouge and more or less independent of the 762 deformation conditions or the type of mineral. The lower limit of grain sizes in 763 experimentally and naturally produced gouge is r = 15 nm. 764 765 D_> is 1.5-1.6 for cracked grains and 2.0-2.3 for gouge and depends on the deformation 766 conditions or the type of mineral. The upper limit of grain sizes evaluated in this study, r_{max}, is 767 100 µm. D_> for gouge is a good parameter to compare natural and experimental fault rock of 768 grain sizes smaller than ~100 µm. 769 770 Cracked grains result from initial fragmentation by rupturing. They develop into gouge by 771 subsequent grain comminution, grinding, attrition, or shear during slip along the fault zone. 772 These processes produce larger D>-values and therefore represent a more efficient grain size 773 reduction for quartz and feldspar grains in the size range greater than 1-2 µm. 774 775 The grain size, r_K , at the slope break of the log(frequency)-log(size) histogram of the GSD 776 occurs at 1.2 ± 0.3 µm for quartz and at 1.8 ± 0.3 µm for feldspar. The grain size r_K coincides 777 approximately with the grinding limit in quartz and probably corresponds to a change in the 778 physical process of grain comminution. Attrition and shear may dominate below the transition 779 value r_K. 780 781 Most of the experimental conditions during deformation (temperature, confining pressure, 782 H₂O-content, displacement rate) have a minor effect on the GSD. Most noticeable is an 783 increase in confining pressure (~1 GPa), which reduces the efficiency of grain comminution. 784 785 The gouge of the Nojima Fault Zone shows $D_> = 2.02$ for gouge and 1.64 for cracked grains. 786 These values are the same as for the experimentally deformed granitoids. 787 788 Experiments and natural seismic fault rocks from the Nojima fault zone show the same GSD, 789 although the experiments have been carried out at higher PT-conditions and at considerably 790 slower displacement rates than in nature. 791 792 The development of the GSD is a two-stage process. First rupturing of the rock causes 793 cracked grains that evolve to a gouge as a result of further movement on the fault zone. The 794 surface densities calculated from the GSD of cracked minerals and gouge are 200-300 times 795 lower than published data, confirming that the creation of surface plays a minor role in the 796 energy budget of earthquakes. 797 798 799

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1050				
1051	Figure captions & Table captions			
1052				
1053	Figure 1:			
1054	<u>C</u>			
1055	(a) Undeformed sample material from the Val Verzasca, Switzerland; scanning electron micrographs			
1056	back scatter contrast. This gneiss is isotropic at the scale of a thin section and has a granite			
1057				
1058				
1059	(b) Map of the northern part of Awaji Island, Japan, showing the locality of the Hirabayashi drill site			
1060				
1061	enlarged area (black square).			
1062	(c, d) Fault rocks of the Nojima Fault Zone (recent event); scanning electron micrographs, back			
1063	scatter contrast. (c) Gouge with rounded siderite fragments (white arrows) and pseudotachylyte			
1064	fragment (black arrow). 623.50 m depth. (d) Cracked grains: individual parts can still be fitted			
1065	together. The white arrow indicates siderite grains. 633.08 m depth.			
1066				
1067	Figure 2:			
1068	Schematic of procedure for grain size analysis.			
1069	(a) Starting image: gouge of Nojima Fault Zone; scanning electron micrographs, backscatter contrast			

(b) Binary image. (c) Histogram of equivalent radii (1285 grains). (d) Log-log plot of frequency

versus equivalent radius (20 bins per order of magnitude). The slope of the line fit yields the D-value.

1070

- 1072 (e) Combined log-log plot for entire range of magnification, showing two distinct average D-values
- 1073 (D_< and D_>) and r_K (see text). (f) D-values versus magnification. Note that intermediate magnifications
- 1074 yield both D-values (D< and D>).

- 1076 Figure 3:
- 1077 Mechanical data.
- 1078 (a) Force versus axial shortening of samples of Verzasca Gneiss. For experimental conditions see
- Table 3, r-h = room-humidity. (b) Differential stress versus axial shortening, calculated using
- program rigC4 (see text). For comparison, 3 experiments on Westerly Granite by Tullis and Yund
- 1081 (1977) are included. Since pyrophyllite or talc was used as confining medium, the true P_c-values
- 1082 (indicated in brackets) are approx. 2/3 of the values given in the literature (see text).

1083

- 1084 Figure 4:
- 1085 Maps of local variations of grain size distribution.
- 1086 The radius of the Gauss-filter was 5 pixels (= 40 μm); for D-mapping technique, see Heilbronner &
- Keulen (2006). D-values > 1.75 indicate gouge, D-values < 1.75 indicate cracked grains. Holes and
- biotite in samples are masked. Three samples are shown: 38nk (reference sample): d = 3 mm; 102nk:
- 1089 d = 4.8 mm; 60nk; d = 2.5 mm. Fault surfaces are indicated. Scale applies to all samples.

1090

- 1091 Figure 5:
- 1092 Microstructures of experimentally deformed samples of Verzasca Gneiss.
- Scanning electron micrographs, backscatter contrast. Minerals, in order of increasing brightness:
- quartz, Na-rich plagioclase, muscovite, K-feldspar, biotite.
- 1095 (a) Fault zone consisting of gouge with adjacent cracked grains (reference sample 38nk). (b)
- Relatively narrow gouge zone (sample 60nk, high confining pressure). (c) Reference sample 64nk. (d)
- Reference sample 38nk. (e) Reference sample 38nk. (f) Reference sample 38nk. (g) Cracked quartz
- grains at high magnification (high displacement sample 102nk). (h) Quartz gouge at high
- magnification (reference sample 38nk).

- 1101 Figure 6:
- Grain size distribution of experimentally deformed Verzasca Gneiss plotted in a log(frequency)-
- log(equivalent radius) histogram. Results for quartz and feldspar are shown in separate graphs, results
- for cracked grains and gouge are shown with separate symbols. The frequency is normalized to 100
- grains with a radius of 10 µm for cracked grains and to 1000 grains with a radius of 10 µm for gouge.

- Note the slope change in most of the curves at $r \approx 2 \mu m$. The D-values (D< and D>) and the grain size
- at the slope change, r_K , are listed in the Table 4 and 5, respectively.

- 1109 Figure 7:
- D-values of the gouge plotted against magnification of analyzed images.
- 1111 Average D-values are indicated by the horizontal lines. Feldspar and quartz are shown separately.
- Analyses at 1000x and 2000x magnification yield a value for $D_{<}$ and $D_{>}$, because r_{K} lies in the range
- of analyzed grain sizes .

1114

- 1115 Figure 8:
- Grain size distribution and D-values of the Nojima Fault Zone.
- 1117 (a) Grain size distribution of cracked grains and gouge (all minerals). The frequency is normalized to
- 1118 100 grains with a radius of 10 μm for cracked grains and to 1000 grains with a radius of 10 μm for
- gouge. (b) D-values versus magnification, compiled from 9 cascades (see text). The average values
- per magnification are indicated with solid symbols. Overall average values for D< and D> are
- indicated by solid lines. The thin-sections used for the analyses originate from the following locations
- in the cores: a) HR4-1 (522.79 m), HR4-11 (633.08 m) b) HR4-11, HR2-52 (626.90 m), HR256
- 1123 (643.10 m) and HR3-10 (623.46 m).

1124

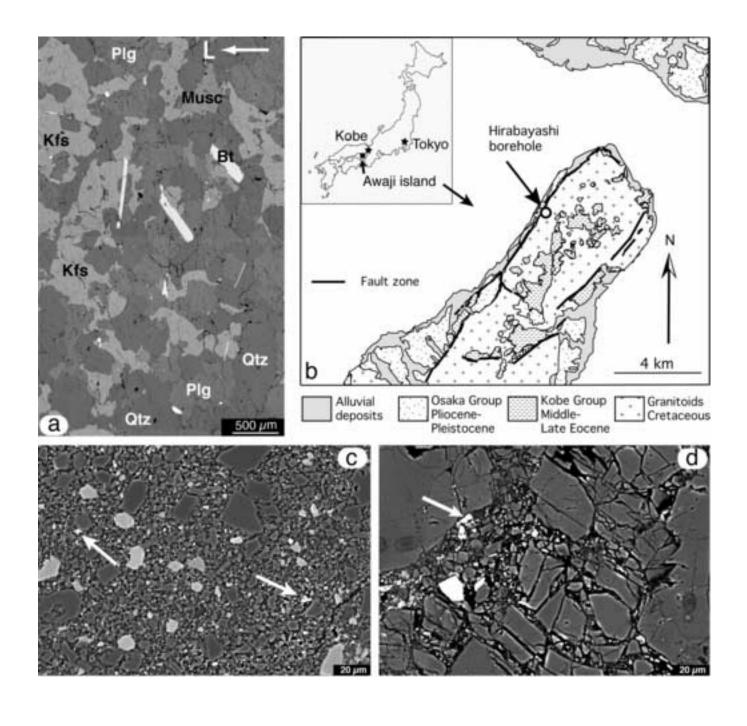
- 1125 Figure 9:
- 1126 Scanning electron micrographs of the smallest grains found in gouge.
- 1127 (a,c) Nojima Fault Zone. (b,d) Experiment 63nk (room-humidity sample). (a,b) Secondary electron
- mode. (c,d) backscatter contrast mode. The pockmark structure in (a) is an effect of the carbon coating
- and does not show individual atoms or another substructure.

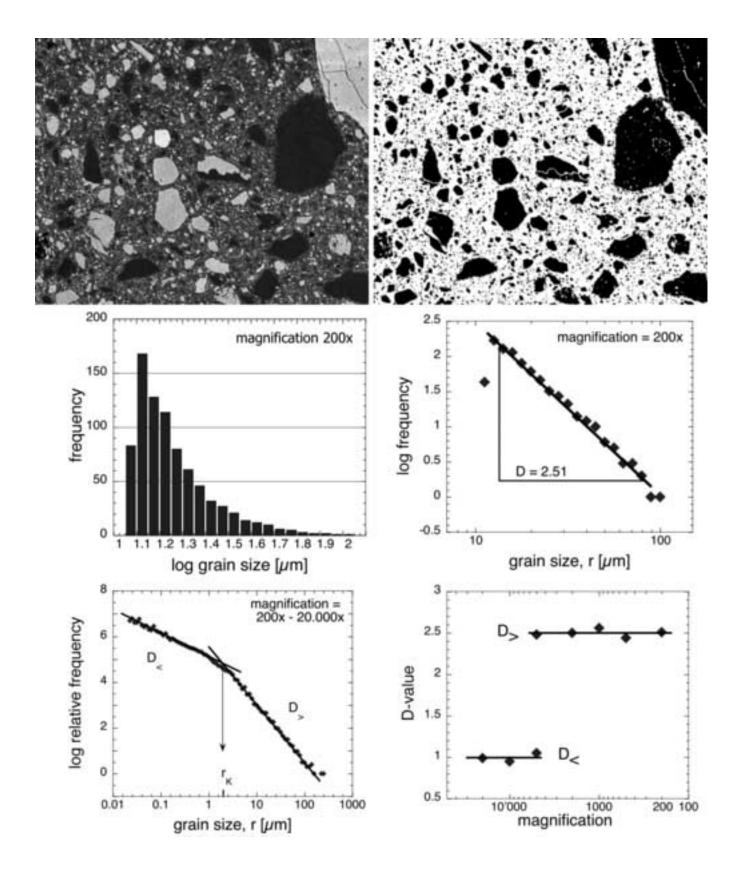
1130

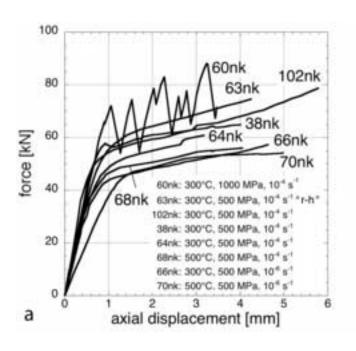
- 1131 Figure 10:
- Grain size, r_k, at the transition from D< to D> (slope change in GSD plots) and grinding limit for
- quartz, after Steier & Schönert (1972) in Prasher (1987). The uncertainty in the position of r_K is ± 0.3
- 1134 µm.

1135

1137 Table 1: 1138 Published data on grain size analyses of fault rocks. 1139 Method of analysis: LM = light microscope image of thin section, SEM = scanning electron 1140 microscope image, manual = analogue box counting, digitizer = instrument for digital analysis of 1141 photomicrographs, digital = analyzed with computer, Coult count. = laser diffraction particle size 1142 analysis. The indicated fractal dimension (D-value) is for 2-dimensional analysis. S.A.F. = San 1143 Andreas Fault, R.S.A. = Republic of South Africa. 1144 1145 Table 2: 1146 Abbreviations and definitions of measurements as used in this paper. 1147 1148 Table 3: 1149 1150 Deformation conditions of the experiments on Verzasca gneiss. 1151 T = temperature, P_c = confining pressure, \dot{e} = axial shortening rate rate, L_0 = starting length of 1152 sample; $d_{chart} = axial displacement from chart strip, <math>d_{chart-1 mm} = corrected axial displacement, H₂O =$ 1153 amount of added water. Samples were quenched immediately after deformation. 1154 1155 Table 4: 1156 D-values (D< and D>) of cracked grains and gouge. 1157 In experiments, quartz and feldspar were measured separately; in natural samples (Nojima), all 1158 mineral grains were combined. 1159 1160 Table 5: 1161 r_K values determined from grain size distributions. 1162 In experiments, the values for quartz and feldspar were determined separately; in natural samples 1163 (Nojima), all mineral grains were combined.







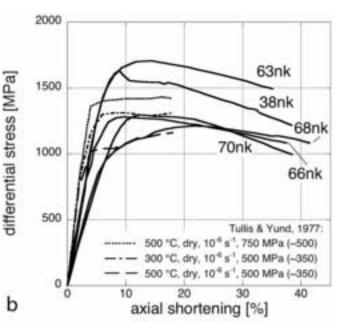


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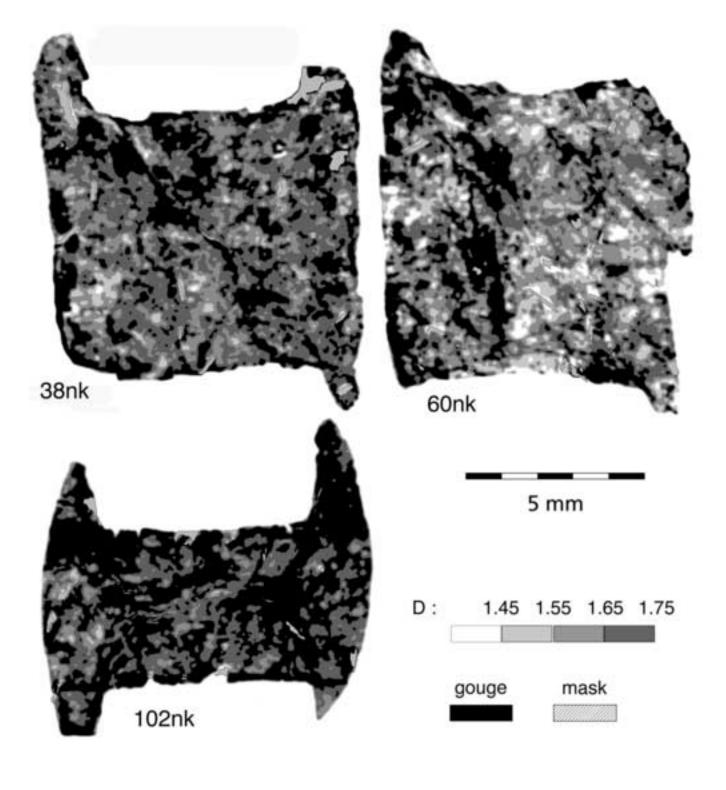
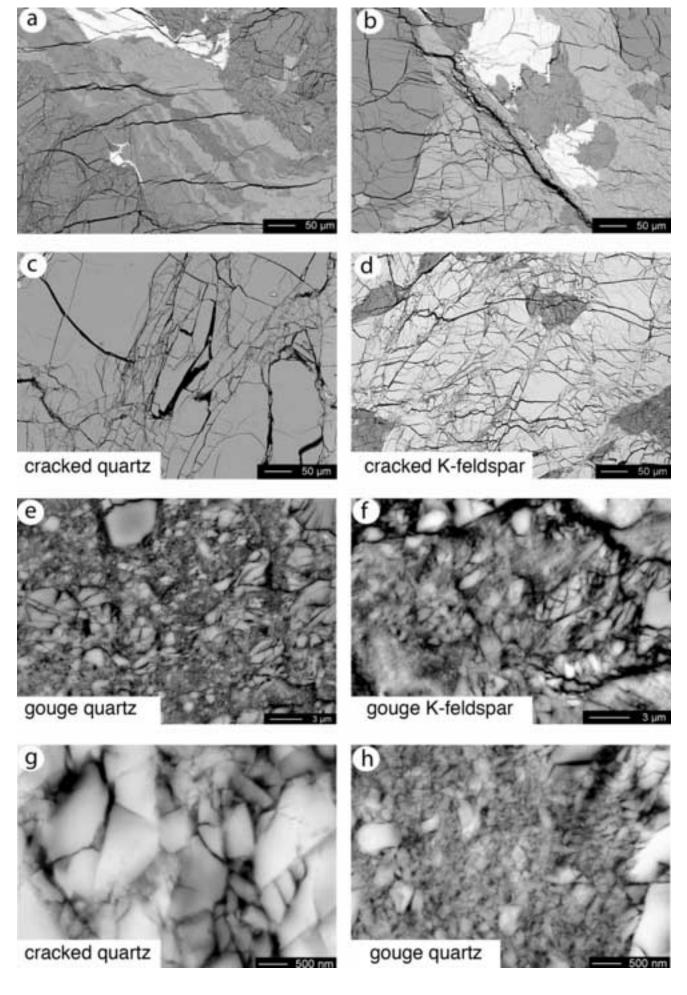
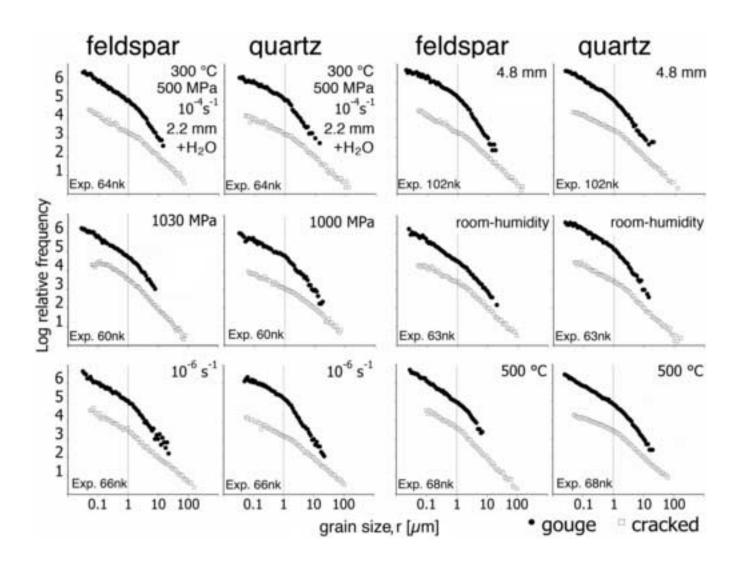
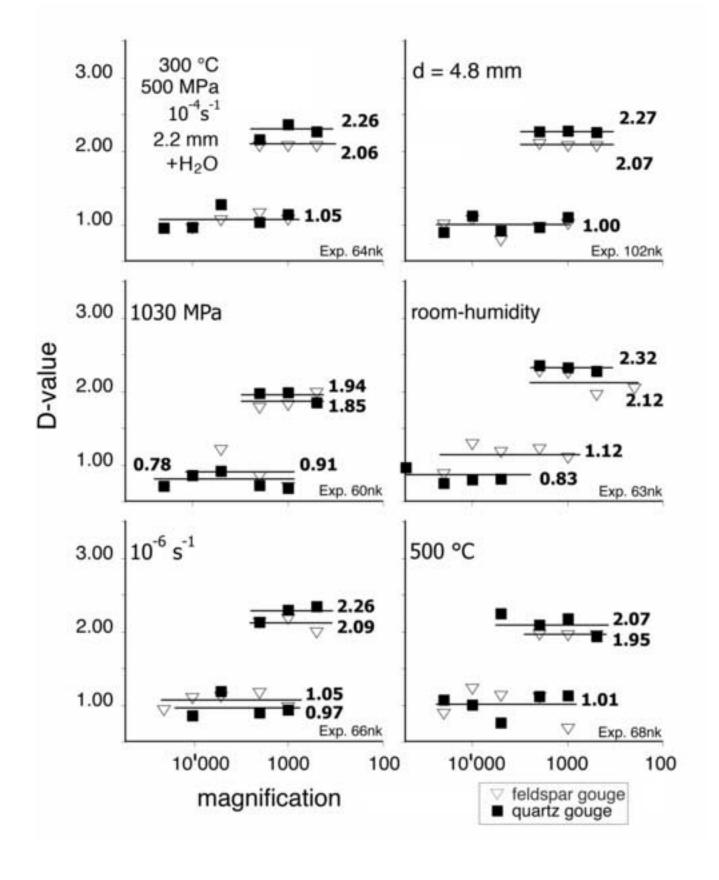


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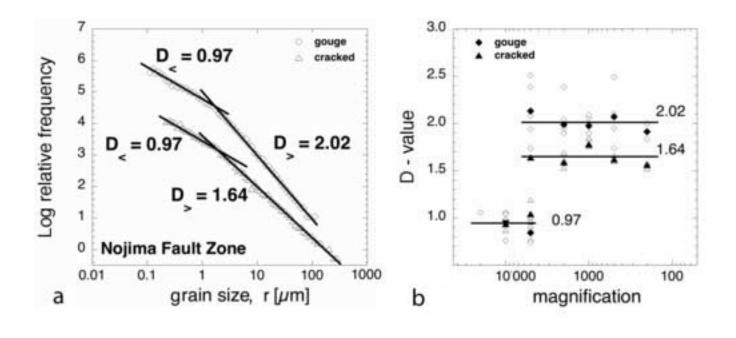
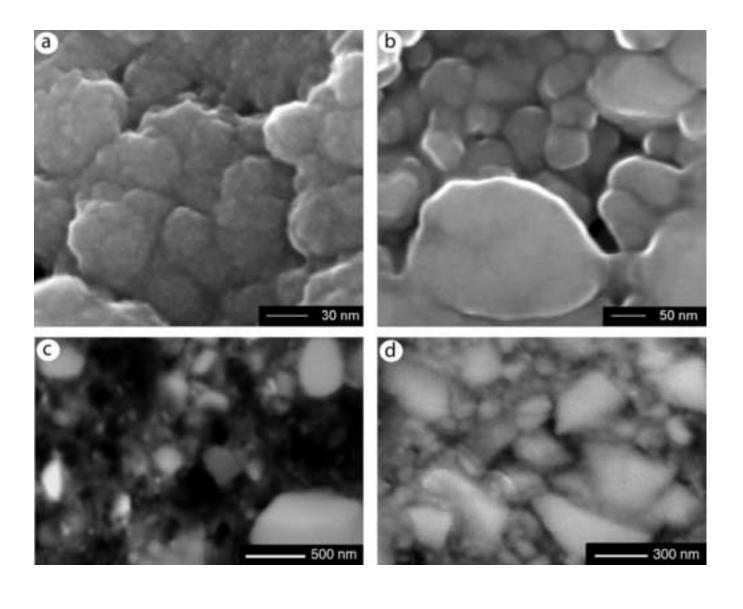
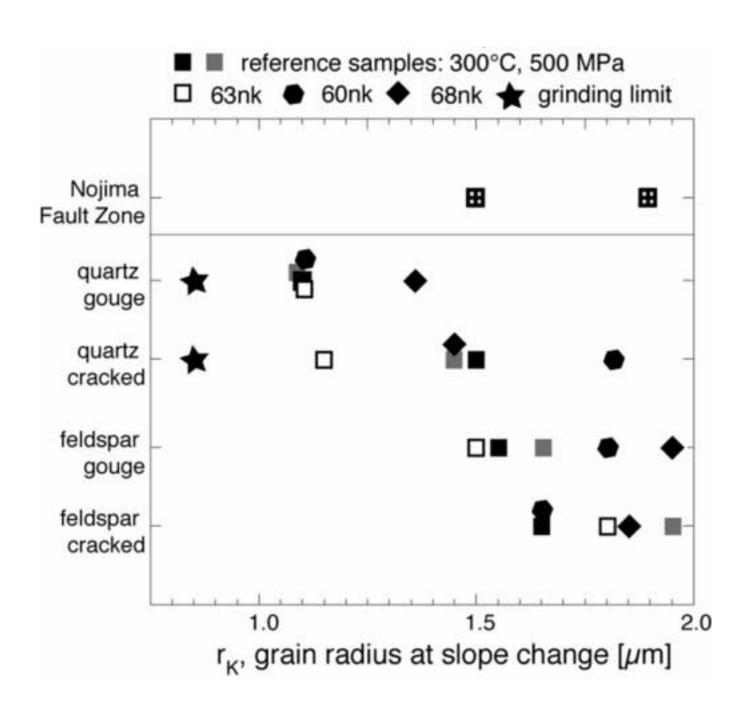


Figure09 Click here to download high resolution image





Authors	year	rock type	location	nat/exp	method of	(2D)-fractal	(2D)-fractal smallest diameter	
					analysis	dimension	measured, μm	
Sammis et al.	1987	granite	Lopez Canyon	natural	LM + SEM, manual	1.6	2	
Sammis & Biegel	1989	granite		natural	mannal	1.6		
Marone & Scholz	1989	quartz sand (porous)	Ottawa	experimental ^a	SEM, manual	1.6	1.5	
Blenkinsop	1991	granite, gneiss	SAF,	$\operatorname{natural}^b$	LM, digitizer	0.8 - 2.1	3	
		arkose	Cajon Pass Drillhole	$\operatorname{natural}^{b}$	LM, digitizer	1.6 - 2.0	ಬ	
An & Sammis	1994	granite	SAF, Tejon Pass	natural	sieve $+$ Coult. Count.	1.4 - 2.6	2	
		gneiss	San Gabriel	natural	sieve $+$ Coult. Count.	1.6 - 1.9	2	
		tonalite	Lopez Canyon	natural	sieve $+$ Coult. Count.	1.5 - 1.9	2	
Shao & Zou	1996	schist and gneiss	Qinling Mountain	natural	LM + SEM, manual	1.6	9	
Monzawa & Otsuki	2003	granite	$\operatorname{Tanakura}$	natural	LM, manual	1.7	10	
		granite	ItShimotsutaki	natural	LM, manual	2.1	11	
		granite	Koi	natural	LM, manual	2.1	8.1	
		granite	Nojima	natural	LM, manual	2.3	7.7	
Hadizadeh & Johnson	2003	quartz sandstone	Masillon, Ohio	experimental ^c	LM, digitizer	1.9-2.5	3.5	
Wilson et al.;	2002	granite	SAF, Tejon Pass	natural	Coult. Count.		0.04	
Reches & Dewers	2002	granite	Bosman fault, R.S.A.	natural	Coult. Count.		0.04	
Chester et al.	2002	granite	SAF, Punchbowl	natural	LM + TEM, manual	2.0	0.05	
Heilbronner & Keulen	2006	granitoids	Verzasca	$experimental^d$	SEM, digital	1.4 - 2.3	0.03	
This study		granitoids	Nojima	natural	SEM, digital	1.6 - 2.4	0.03	
		granitoids	Verzasca	experimental ^e	SEM, digital	1.4 - 2.3	0.03	

Table 1:

Symbol	Definition
$\overline{\mathrm{D}_{<}}$	slope of power law fit for small grain sizes $(r < r_K)$
$D_{>}$	slope of power law fit for large grain sizes $(r > r_K)$
\mathbf{r}_K	radius of grains at slope break of GSD
L_0	original length of sample
Δ_L	axial shortening of sample
\mathbf{r}_{min}	minimum grain size used for calculation of GSD
\mathbf{r}_{max}	maximum grain size used for calculation of GSD
A	cross sectional area of grain (pixel ²)
P	perimeter of cross section of grain (pixel)

Table 2:

Deformation conditions

sample nr.	T	P_c	$\mathrm{ax}.\dot{\epsilon}$	L_0	d_{chart}	$d_{(chart-1mm)}$	H_2O
	$^{\circ}\mathrm{C}$	MPa	$\ge 10^{-4} s^{-1}$	mm	mm	%wt	
38nk	300	510	1.5	10.44	4.01	3.0	0.2
$64\mathrm{nk}$	300	500	1.3	12.33	3.18	2.2	0.2
102nk	300	510	1.3	10.17	5.78	4.8	0.2
$60\mathrm{nk}$	300	1030	1.5	10.13	3.47	2.5	0.2
$68\mathrm{nk}$	500	520	1.4	10.32	4.28	3.3	0.2
$63\mathrm{nk}$	300	530	1.4	12.11	4.25	3.3	0.0
$66\mathrm{nk}$	300	490	0.015	12.37	4.64	3.6	0.2
$70\mathrm{nk}$	500	500	0.013	12.95	4.97	4.0	0.2

Table 3:

D-values for cracked grains and gouge

		feld	spar		quartz				
	crac	cked	go	uge	cracked		gouge		
sample nr.	$D_{<}$	$D_{>}$	$D_{<}$	$D_{>}$	$D_{<}$	$D_{>}$	$D_{<}$	$D_{>}$	
38nk 64nk 102nk 60nk 68nk 63nk 66nk	1.02 0.95 0.87 0.74 0.92 0.72 0.86	1.68 1.52 1.52 1.37 1.67 1.62 1.37	1.00 1.05 1.01 0.91 1.01 1.12 1.05	2.03 2.06 2.07 1.85 1.95 2.12 2.09	1.09 0.74 0.92 0.96 0.92 0.78 0.96	1.56 1.44 1.55 1.72 1.47 1.55 1.46	1.10 1.05 1.00 0.78 1.01 0.83 0.97	2.26 2.26 2.27 1.94 2.07 2.32 2.26	
	"	ima Feked $\begin{array}{c} D_{>} \\ 1.64 \end{array}$	$\begin{array}{c} \text{Fault } D_{<} \\ D_{<} \\ 0.97 \end{array}$	Zone uge $D_{>}$ 2.02					

Table 4:

 r_K : Grain radius at slope change $[\mu m]$

	feldspar			quartz				
	C	racked		gouge	cracked go		gouge	
sample nr.	r_K	range	r_K	range	r_K	range	r_K	range
38nk 64nk 102nk 60nk 68nk 63nk 66nk	1.6 1.9 1.8 1.7 1.8 1.8 1.6	1.4 - 1.7 1.7 - 2.0 1.7 - 2.0 1.6 - 1.7 1.6 - 1.9 1.7 - 1.9 1.3 - 1.6	1.5 1.6 1.5 1.8 1.9 1.5 1.3	1.4 - 1.7 1.5 - 1.7 1.3 - 1.7 1.6 - 1.9 1.9 - 2.0 1.5 - 1.6 1.3 - 1.6	1.5 1.3 1.4 1.8 1.5 1.1 1.8	1.3 - 1.6 1.2 - 1.4 1.4 - 1.5 1.8 - 2.0 1.4 - 1.6 1.1 - 1.3 1.6 - 1.9	1.1 1.1 1.1 1.1 1.3 1.1 1.4	0.9 - 1.1 1.0 - 1.3 0.9 - 1.1 1.0 - 1.2 1.3 - 1.5 1.0 - 1.3 1.4 - 1.6
	m Nojima~Fa $ m cracked$ $ m range$ $ m 1.9~1.6-2.0$			Zone gouge range 1.2 - 1.6				

Table 5: