

Structure of the Synthetic K-rich Phyllomanganate Birnessite Obtained by High-Temperature Decomposition of KMnO4. Substructures of K-rich Birnessite from 1000°C Experiment

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1	Structure of the Synthetic K-rich Phyllomanganate Birnessite Obtained
2	by High-Temperature Decomposition of KMnO ₄ .
3	Substructures of K-rich Birnessite from 1000°C Experiment.
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Abstract

1	9	

The structure of a synthetic potassium-rich birnessite prepared from the thermal 20 decomposition of KMnO₄ at 1000°C in air has been refined by Rietveld analysis of the 21 powder X-ray diffraction (XRD) data, and the structure model shown to be consistent with 22 extended X-ray absorption fine structure data. K-rich birnessite structure is a two-layer 23 orthorhombic polytype (20) with unit-cell parameters a = 5.1554(3) Å, b = 2.8460(1) Å, c =24 14.088(1) Å, $\alpha = \beta = \gamma = 90^{\circ}$, $a/b = \sqrt{3.281}$, and was refined in the *Ccmm* space group. The 25 structure is characterized by the regular alternation of octahedral layers rotated with respect to 26 each other by 180°. Octahedral layers are essentially devoid of vacant sites, the presence of 27 $0.25 \text{ Mn}^{3+}_{\text{laver}}$ cations within these layers being the main source of their deficit of charge, 28 which is compensated for by interlayer K^+ cations. Mn^{3+} octahedra, which are distorted by the 29 Jahn-Teller effect, are systematically elongated along the a axis (cooperative Jahn-Teller 30 effect) to minimize steric strains, thus yielding an orthogonal layer symmetry. In addition, 31 Mn^{3+} octahedra are segregated in Mn^{3+} -rich rows parallel to the **b** axis that alternate with two 32 Mn^{4+} rows according to the sequence ...- $Mn^{3+}-Mn^{4+}-Mn^{3+}-...$ along the **a** direction, 33 thus leading to a A = 3a super-periodicity. At 350°C, the structure partially collapses due to 34 the departure of interlayer H₂O molecules and undergoes a reversible 2O-to-2H phase 35 36 transition. This transition results from the relaxation of the cooperative Jahn-Teller effect, that is from the random orientation of elongated Mn^{3+} octahedra. 37

Introduction

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Birnessite is a phyllomanganate, that is a manganese oxide containing predominantly 40 Mn⁴⁺ cations assembled in layers of edge-sharing MnO₆ octahedra. A layer charge deficit 41 arises from the presence within layers of Mn³⁺ cations and/or vacant layer octahedra and is 42 compensated for by the presence of interlayer cations which are typically hydrolysable cations 43 [1-8]. As such phyllomanganates are very similar to expandable 2:1 phyllosilicates (smectites) 44 and can be considered as $(2 \times \infty)$ microporous solids [9,10]. As smectites, phyllomanganates 45 can be intercalated with a variety of organic and inorganic compounds to form multilayer 46 nanocomposites [11-17], or pillared structures [18-24]. In addition, phyllomanganates have 47 been widely used as templates for the formation of octahedral molecular sieves with variable 48 tunnel sizes that have demonstrated excellent potential in heterogeneous catalysis, hazardous 49 waste remediation, and rechargeable battery technology [25-42]. More recently, mesoporous 50 hollow shells with birnessite walls have also been synthesized [43]. 51 52 Besides, birnessites play a pivotal role in the fate of heavy metals and other pollutants in contaminated water systems and soils because they possess unique crystal-chemical 53 characteristics which confer them extensive redox and sorption properties [3,6,44-49]. A large 54 55 range of birnessite varieties with different structural and chemical characteristics can be synthesized in the laboratory. Synthetic birnessite can be obtained from the oxidation of Mn^{2+} 56 in highly alkaline medium [1,2,50,51], from the thermal decomposition of MnO-containing 57 mixtures [52], and from the reduction of Mn^{7+} [53-59]. In most cases, syntheses are 58 performed at low temperature to obtain surrogate materials representative of natural 59 60 specimens [1,2,60,61]. Low-temperature birnessites are usually finely dispersed, and often present a low degree of structural order. 61

Over the last decade an interest for high temperature birnessite (from 170°C up to
 1000°C) has emerged as the result of their promising potential as electrodes in secondary

64	lithium batteries [52,56-58,62-65]. Kim et al. showed that birnessites with high structural
65	perfection can be obtained at high temperature, these varieties being more stable than those
66	commonly used in electrochemistry [58]. In particular the transition to a spinel structure was
67	not observed after cell cycling. The protocol is based on the thermal decomposition of
68	$KMnO_4$, large K^+ cations pillaring the synthesis products, thus stabilizing the layered
69	framework. In addition, the production of highly oxidizing species during the decomposition
70	prevents the formation of suboxides. As a result, a new variety of synthetic K-rich birnessites
71	with a two-layer periodicity along the c axis was obtained [58]. This variety will be hereafter
72	referred to as KBi. These authors performed experiments from 200-1000°C, but structural and
73	crystal chemical study was limited to high-temperature varieties. The Rietveld technique was
74	used to determine the crystal structure of KBi obtained at 800°C. In the refined model
75	adjacent vacancy-free layers are rotated with respect to each other by 180° around the c axis,
76	which is perpendicular to the layers, and interlayer K^+ is located in a prism defined by
77	tridentate cavities of adjacent layers [58]. Kim et al. also synthesized at 1000°C a KBi sample
78	that presents a different chemistry but did not refine its structure [58].
79	
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81 82 83 84 85 86	characterization of high-temperature KBi synthesized according to the protocol of Kim et al. [58]. This series of articles was initiated by the structure refinement of a KBi sample obtained at 800°C from single-crystal and powder X-ray diffraction (XRD) and extended X-ray absorption fine structure (EXAFS) spectroscopy experiments [67]. XRD structure refinement has pointed out several significant differences with the model proposed by Kim et al. [58,67]. First, the cation composition of KBi interlayers is heterogeneous and includes both K ⁺ and Mn ³⁺ cations. Second, layers of KBi synthesized at 800°C contain significant amount of

the interlayer mid plane. The determination of these fine structure details should provide a 90 better understanding of the relationship between KBi structure and its properties. The second 91 paper of this series describes the phase and structural heterogeneities of KBi samples 92 93 synthesized at 200-1000°C, together with the factors responsible for these heterogeneities [68]. The present work is devoted to the structural characterization of a homogeneous KBi 94 95 sample obtained at 1000°C and of this KBi variety when heated to 350°C. The present article will focus on the average distribution of layer and interlayer cations within the subcell and on 96 the origin of the layer charge, whereas KBi superstructure and its origin will be discussed in 97 the last part of this series of articles devoted to high-temperature KBi structure [66]. One 98 99 effective tool to study superstructures is to use, along with XRD, selected area electron diffraction (SAED) [5,6,46,69-71]. However, comparison of XRD and SAED data may 100 appear irrelevant because of the possible modification of interlayer species distribution as the 101 result of dehydration under vacuum [57]. XRD experiments were thus performed under deep 102 vacuum conditions to check the consequences of such dehydration on the layer structure and 103 104 on the interlayer cation distribution. 105 **Experimental Section** 106 107 Experimental Methods. K-rich phyllomanganate birnessite (hereafter referred to as 108 KBi) was prepared by thermal decomposition of fine grained KMnO₄ powder (particle size < 109 50 µm) at 1000°C in air according to the procedure of Kim et al. [58,67]. As for the 800°C 110 synthesis, structural homogeneity of synthetic products was maximized by using flat crucibles 111 112 containing a thin layer of KMnO₄ powder. KBi samples synthesized at 1000°C (this study) and 800°C will be hereafter referred to as KBi₁₀ and KBi₈ [67], respectively. 113 The morphology of KBi₁₀ particles was observed on a JEOL JSM 6320F high-resolution 114 scanning electron microscope equipped with a field emission electron source. No significant 115

difference in size or shape was observed between KBi_8 and KBi_{10} crystals, and KBi_{10} particles consist of large micro-crystals with well-defined crystallographic faces [67].

Thermal and Chemical Analyses. Thermal analyses were carried out with a NETZSCH 118 119 Simultan STA 409 EP analyzer. DT-TG data were recorded in air using a 6°C/min heating rate over the 20-1100°C temperature range to determine the amount of structural water. Total 120 K and Mn contents were determined using a Perkin-Elmer Optima 3000 ICP-AES after 121 digestion of about 8 mg of KBi powder in 200 mL of 1% HNO3 / 0.1% NH3OHCl matrix. 122 The mean oxidation degree of manganese in birnessite was determined by potentiometric 123 titration using $(NH_4)_2$ Fe (SO_4) Mohr salt and $Na_4P_2O_7$ [72,73]. The mean oxidation state, 2x, 124 in combination with the atomic ratio y = K/Mn determined from the chemical analyses, allows 125 calculating the structural formula of KBi using the relation [67]: 126

127
$$K^{+}_{2y/w}(Mn^{4+}_{(4x-6)/w}Mn^{3+}_{(8-4x)/w}\square_{1-2/w})O_2$$
 (1)

128 where $w = \frac{2x + y}{2}$ and \Box represents vacant layer sites. This structural formula can be refined 129 by taking into account the amount of interlayer H₂O molecules (H₂O_{interlayer}) deduced from the 130 DT-TG curves.

X-Ray Diffraction Data Collection. Powder XRD patterns were recorded using a 131 Bruker D5000 powder diffractometer equipped with a Kevex Si(Li) solid detector and 132 CuK α_{1+2} radiation. Intensities were recorded from 5 to 90° with a 0.04° 20 interval, using a 133 40 sec counting time per step. A rotating sample holder was used for room temperature 134 measurements to minimize the effect of preferential orientation. A Anton Paar TTK450 135 chamber was used to record XRD patterns from samples heated in situ. XRD patterns of 136 KBi₁₀ obtained at 100°C, 150°C, 250°C and 350°C will be hereafter referred to as 100-KBi₁₀, 137 150-KBi₁₀, 250-KBi₁₀ and 350-KBi₁₀, respectively, while the notations KBi₁₀ and AfterT-138 KBi₁₀ will be used to differentiate XRD patterns recorded at room-temperature before and 139 after temperature experiments, respectively. All heated KBi diffractograms were recorded 140 successively in a single experiment. The sample was heated *in situ* at a 6°C/min rate until the 141

142 target temperature was reached and maintained at a temperature plateau for two hours before collecting the diffractogram using the same conditions than at room temperature. After ~24 143 hours of data collection, temperature was further increased to the next two-hour plateau. XRD 144 145 data was finally collected again after cooling the sample back to room temperature (sample AfterT-KBi₁₀). Finally, the same Anton Paar TTK450 chamber allowed studying the structure 146 of KBi₁₀ under deep vacuum conditions using a Varian V70 turbo-molecular pump connected 147 to the sample chamber to reach $\sim 10^{-5}$ hPa. After waiting for 15 hours to ensure a complete 148 dehydration of the sample, XRD data was collected using the same experimental conditions. 149 This sample will be referred to as Vacuum-KBi₁₀, while the XRD pattern recorded once back 150 151 at the atmospheric pressure to allow rehydration of this sample will be referred to as AfterVac-KBi₁₀. Note that indexing was performed systematically assuming an orthorhombic 152 *C*-centered unit cell whatever the actual layer symmetry (orthogonal or hexagonal). 153

Simulation of Powder XRD Patterns. The Rietveld technique was used to refine the 154 structure of sample KBi₁₀ using the XND code [74]. To de-correlate instrumental broadening, 155 profile shape function (PSF) parameters and spectral distribution were refined first using a 156 reference quartz powder sample [8]. During the subsequent structure refinement process these 157 158 instrumental parameters were not further refined. A spline function was used to interpolate a background defined by 6-to-8 points. Scale factor and unit-cell parameters were refined first 159 before site coordinates and occupancies. Atomic thermal motion was modeled using isotropic 160 Debye-Waller factors (B). In the final stage of the refinement, orientation parameters were 161 introduced using first degree spherical functions. Fit quality was estimated with the standard 162 $R_{\rm wp}$ and $R_{\rm B}$ factors. More details on the XND code can be found at 163 164 http://www.ccp14.ac.uk/ccp/ccp14/ftp-mirror/xnd/pub/xnd/html/xnd.html. **EXAFS Experiments.** Mn EXAFS spectra were recorded at the LURE synchrotron 165

radiation laboratory (Orsay – France) on the D42 spectrometer of the EXAFS I station.

167 Experimental conditions and data analysis were the same as those reported for KBi₈ [67]. The

168	EXAFS spectra were apodized by a Kaiser-Bessel function ($\beta = 2.5$) and then Fourier
169	transformed to real space to generate radial structure functions (RSFs). RSFs were not
170	corrected for phase shifts, causing peaks to appear at shorter distances ($R + \Delta$, with $\Delta \approx$ -
171	0.4 Å) relative to the true interatomic distances (R). Structural parameters from the nearest
172	oxygen and manganese shells (Mn-O and Mn-Mn atomic pairs) were determined by back-
173	transforming to k space the first two RSF peaks, and fitting the data with phase shift and
174	amplitude functions calculated by FEFF7 using λ -MnO ₂ as structure model [75,76]. The
175	short-distance Mn-O and Mn-Mn contributions were fitted together because of incomplete
176	separation in the Fourier-filtering process.
177	
178	Results
179	
180	Chemical Analyses and Structural Formulae.
101	Amount of Interlance Water The DTA surve of KDi contains two endetherms at 170°C
181	Amount of Interlayer Water. The DTA curve of KBi ₁₀ contains two endotherms at 178°C
181	and 926°C, respectively (Fig. 1). The low-temperature endotherm (178°C) corresponds to the
182	and 926°C, respectively (Fig. 1). The low-temperature endotherm (178°C) corresponds to the
182 183	and 926°C, respectively (Fig. 1). The low-temperature endotherm (178°C) corresponds to the loss of weight due to interlayer water, whose amount equals 7.5 wt% (between 120-210°C)
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182 183 184 185 186	and 926°C, respectively (Fig. 1). The low-temperature endotherm (178°C) corresponds to the loss of weight due to interlayer water, whose amount equals 7.5 wt% (between 120-210°C) while the high-temperature endotherm (926°C) is likely related to KBi melting. The additional loss of weight within the 60-120°C range is most likely related to water adsorbed on grain surfaces. In contrast to what was observed for sample KBi ₈ and hydrothermal
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182 183 184 185 186 187 188	and 926°C, respectively (Fig. 1). The low-temperature endotherm (178°C) corresponds to the loss of weight due to interlayer water, whose amount equals 7.5 wt% (between 120-210°C) while the high-temperature endotherm (926°C) is likely related to KBi melting. The additional loss of weight within the 60-120°C range is most likely related to water adsorbed on grain surfaces. In contrast to what was observed for sample KBi ₈ and hydrothermal birnessite [67,77], no DT peak or change of slope in TG curve were observed over the 250-400°C interval.
182 183 184 185 186 187 188 189	and 926°C, respectively (Fig. 1). The low-temperature endotherm (178°C) corresponds to the loss of weight due to interlayer water, whose amount equals 7.5 wt% (between 120-210°C) while the high-temperature endotherm (926°C) is likely related to KBi melting. The additional loss of weight within the 60-120°C range is most likely related to water adsorbed on grain surfaces. In contrast to what was observed for sample KBi ₈ and hydrothermal birnessite [67,77], no DT peak or change of slope in TG curve were observed over the 250- 400°C interval. <i>Structural Formula</i> . The K/Mn atomic ratio [$y = 0.320(5)$] was calculated from K and Mn
182 183 184 185 186 187 188 189 190	and 926°C, respectively (Fig. 1). The low-temperature endotherm (178°C) corresponds to the loss of weight due to interlayer water, whose amount equals 7.5 wt% (between 120-210°C) while the high-temperature endotherm (926°C) is likely related to KBi melting. The additional loss of weight within the 60-120°C range is most likely related to water adsorbed on grain surfaces. In contrast to what was observed for sample KBi ₈ and hydrothermal birnessite [67,77], no DT peak or change of slope in TG curve were observed over the 250- 400°C interval. <i>Structural Formula</i> . The K/Mn atomic ratio [$y = 0.320(5)$] was calculated from K and Mn ICP concentrations values measured from the same solution whereas the H ₂ O/Mn mole ratio

formula, in which \Box correspond to vacant layer sites, can be deduced for KBi₁₀ crystals at room temperature and pressure:

196
$$K_{0.314}^{+}(Mn_{0.246}^{3+}Mn_{0.737}^{+}\square_{0.017})O_2 \cdot 0.50H_2O$$
 (2)

197 Indexing of the experimental XRD patterns.

Sample KBi₁₀. The powder XRD pattern (Fig. 2A) contains a rational series of basal 198 reflections with d(00l) values corresponding to a minimum periodicity along the c* axis equal 199 to 7.044 Å. The non-basal reflections are intense and sharp. The positions of *hkl* reflections 200 are compatible with an orthorhombic two-layer (20) C-centered cell with a = 5.1554(3) Å, 201 b = 2.8460(1) Å, c = 14.088(1) Å, $\alpha = \beta = \gamma = 90^{\circ}$, $a/b = \sqrt{3.281}$ (Table 1). With such an 202 indexing, XRD patterns of birnessites contain 201/111 reflections over the 34-64° 20 CuKa 203 angular range, 31l/02l maxima being visible over the 64-74° 20 CuK α range and 40l/22l over 204 the 74-90° 20 range. Along with the described *hkl* reflections, the XRD pattern contains 205 additional weaker reflections, mainly over the 5-34° 20 region (inset on Fig. 2A), which can 206 be indexed using a C-centered supercell with A = 3a, B = b, C = c, $\alpha = \beta = \gamma = 90^{\circ}$. This 207 super-structure will be described and its origin discussed in the final part of this four paper 208 series [66]. 209

Sample 100-KBi₁₀. The XRD pattern of 100-KBi₁₀ (Fig. 2B) is very similar to that of KBi₁₀ and can be indexed with a 2O unit cell (Table 1). Previously described super-reflections are still visible.

Sample 150-KBi₁₀. Although TD-TG curves indicate that KBi₁₀ dehydrates essentially at ~180°C, the XRD pattern of 150-KBi₁₀ (Fig. 2C) shows a significant decrease of the d(002)value at ~6.44 Å, and related shifts of *hkl* reflections. The higher temperature deduced from TD-TG curves for the departure of interlayer H₂O molecules likely results from the high heating rate used for the thermal analysis, whereas the 2-hour plateau at 150°C likely allowed reaching complete dehydration before XRD data collection. Apart from the reduction of the unit-cell dimension along the **c*** axis, the birnessite layer structure remains essentially stable,

the main reflections being indexed with a 20 unit cell (Table 1). A close look at the

221 experimental XRD pattern reveals however that the profile of intense non-basal reflections is

altered with shoulders visible on their low-angle sides (shown for the 114 reflection on the

inset – Fig. 2C). These additional maxima can be indexed using a two-layer hexagonal (2H)

unit cell with a = 5.010(3) Å, b = 2.891(1) Å, c = 12.871(4) Å, $\alpha = \beta = \gamma = 90^{\circ}$, $a/b = \sqrt{3.00}$

- (Table 1). Additional low-intensity peaks visible over the 5-30° 2θ range can be indexed
- using a *C*-centered supercell with A = 3a, B = 3b, C = c, $\alpha = \beta = \gamma = 90^{\circ}$.

Samples 250-KBi₁₀ and 350-KBi₁₀. At 250°C dehydration is complete and heating of 227 KBi₁₀ at this temperature dramatically modifies the experimental XRD pattern as the result of 228 important structural changes that persist at least up to 350°C, XRD patterns of both 350-KBi₁₀ 229 and 250-KBi₁₀ being similar (Figs 2D-E). The two patterns contain two basal reflections 230 characteristic of birnessite together with a set of intense and sharp non-basal reflections 231 showing that the structural perfection is preserved. However, the orthogonal symmetry of the 232 low-temperature KBi₁₀ layers is transformed into a hexagonal one for samples 250-KBi₁₀ and 233 350-KBi₁₀, *hkl* reflections being then indexed using a two-layer hexagonal unit cell (2H -234

Table 1). XRD patterns of samples 250-KBi₁₀ and 350-KBi₁₀ also contain weak super-

reflections in the 5-35° 2 θ range that are similar to those observed at 150°C and indexed using

237 a *C*-centered A = 3a, B = 3b, C = c, $\alpha = \beta = \gamma = 90^{\circ}$ supercell.

238 Sample AfterT-KBi₁₀. Reflections present in the experimental XRD patterns of KBi₁₀

specimens recorded at room temperature before and after heating at 350°C (Figs. 2A and 2F,

respectively) have similar positions and intensities. This means that the structural

241 modification undergone during the heating at 350°C is essentially reversible. The increase of

- the minimum periodicity along the **c*** axis from 6.42 Å back to 7.02 Å after cooling of the
- sample shows that birnessite layers re-hydrate. The orthogonal layer symmetry is also
- recovered, reflections of sample AfterT-KBi₁₀ being indexed with a two-layer orthorhombic
- unit cell (20 Table 1). Super-reflections corresponding to the initial A = 3a, B = b, C = c,

 $\alpha = \beta = \gamma = 90^{\circ}$ C-centered supercell are also recovered. However, differences are visible 246 between KBi₁₀ and AfterT-KBi₁₀ patterns. In particular, in the latter pattern 20*l* reflections are 247 broader and their intensity relative to the 11*l* reflections (with the same *l* index) significantly 248 dropped as compared to sample KBi₁₀ (Figs. 2A, 2F and 3). This effect was described in 249 details by Gaillot et al. for the heterogeneous KBi_{10h} variety obtained at 1000°C [68]. It was 250 shown that 201 reflection broadening and weakening are linked to the coexistence of several 251 populations of KBi particles. All populations are 20 modifications with slightly different a 252 and *b* parameters that are related by the equation: 253

254
$$\frac{1}{d(110)^2} = \frac{1}{a_i^2} + \frac{1}{b_i^2}$$
(3)

As a consequence, positions of 11*l* reflections are insensitive to the variation of unit-cell parameters, whereas 20*l* reflections of individual 2*O* modifications are slightly shifted. The resulting overlap of 20*l* reflections from different 2*O* modifications broadens the resulting diffraction maxima, peak broadening decreasing with increasing *l* values (solid circles, Fig. 3).

Sample Vacuum-KBi₁₀. The XRD pattern of Vacuum-KBi₁₀ is very similar to that of 150-260 KBi₁₀ and presents a shift of basal reflections from 7.044 Å down to 6.397 Å, suggesting the 261 departure of H₂O_{interlayer} (Fig. 2G). As for 150-KBi₁₀, hkl reflections are similar to those of 262 KBi₁₀ but shifted towards higher angles. In contrast to sample 150-KBi₁₀, no reflections 263 corresponding to a 2H modification are visible, all reflections being indexed with a unique 2O 264 unit cell (Table 1). Additional weak reflections correspond to a C-centered supercell with A = 265 $3a, B = 3b, C = c, \alpha = \beta = \gamma = 90^{\circ}$. Finally, two additional very weak maxima at 6.964 Å and 266 3.486 Å likely correspond to basal reflections corresponding to a few partially hydrated 267 crystals. 268

AfterVac-KBi₁₀. The XRD pattern is very similar to that of KBi₁₀ with the 002 basal
 reflection shifted back from 6.397 Å to 7.038 Å indicating the reversibility of the dehydration
 process. Reflections are indexed using a 2*O* unit cell with parameters very similar to those

obtained before dehydration (Table 1). Visible super-reflections correspond to the initial *C*-

273 centered supercell with A = 3a, B = b, C = c, $\alpha = \beta = \gamma = 90^{\circ}$.

- Unit-cell parameters for all KBi₁₀ samples can be found in Table 1, indexation of experimental diffraction lines being provided as supplementary material.
- 276

Structure Refinement of KBi₁₀ Subcells.

*Rietveld Refinement of KBi*₁₀ *Subcell*. On KBi₁₀ XRD pattern all *hkl* reflections 277 consistently have low full width at half maximum intensity (fwhm). In particular, after 278 correction for crystal-size broadening by a $\cos\theta$ factor [78], fwhm of 11*l* reflections only 279 slightly increases with increasing l values (from 0.125° to $0.165^{\circ} 2\theta$ – open triangles, Fig. 3) 280 whereas that of 20*l* reflections is independent of l (0.165° 2 θ – open circles, Fig. 3). This data 281 indicates the high structural perfection of KBi₁₀. Small angular regions containing weak 282 super-reflections (at $\sim 37.5^\circ$, $\sim 39.5^\circ$ and $\sim 51.0^\circ 2\theta$) were excluded from the Rietveld 283 refinement. Atomic positions of KBi₁₀ subcell were refined using space group *Ccmm*. Initial 284 occupancies of structural sites correspond to the structural formula of KBi_{10} (eq 2). In the 285 initial model schematized on Fig. 4A Mnlayer are located at the subcell origin, and Olaver have 286 coordinates (0.333, 0, 0.071) to match the ~2.00 Å layer thickness previously refined for 287 several synthetic birnessite structures [5,8,58,67,77,79]. K⁺ are located in the interlayer mid-288 plane, above and below the layer empty tridentate cavities, sharing three edges with octahedra 289 of adjacent layers (TE sites - Fig. 4A), whereas H₂O_{interlayer} are located in-between the nearest 290 O_{laver} from adjacent layers. At this stage $R_{wp} = 22.65\%$ and $R_B = 23.45\%$. All structural 291 parameters were then successively refined, starting from unit-cell parameters, PSFs (the width 292 of basal reflections 00l and non-basal hkl reflections being described by two sets of 293 parameters) and preferential orientation. Atomic positions and occupancies were then refined. 294 As compared to the initial structure model O_{laver} and H₂O_{interlaver} were shifted along the **a** axis 295 from the initial positions. In addition, the K site was split between three positions located 296 within the interlayer prismatic cavity ($R_{wp} = 12.52\%$, $R_B = 9.57\%$). At this stage, calculated 297

20*l* and 11*l* reflections are slightly more and less intense, respectively, than corresponding
experimental reflections. As 20*l* fwhm values are slightly higher than those of 11*l* ones (solid
circles and solid triangles, respectively – Fig. 3A) some KBi₁₀ particles likely have unit-cell
parameters differing from those of the major part of the sample, unit-cell parameters from the
different populations following eq 3.

To better fit the broader 20l reflections, two PSFs were thus considered for 20l/31l/40l 303 and 111/021/221 diffraction lines, respectively, and structural parameters were refined again. 304 The use of two PSFs allowed fitting better the width of all hkl reflections, but did not affect 305 significantly the details of the optimum structure model, whose main parameters are listed in 306 Table 2. The refined site atomic coordinates and occupancies given in Table 3 led to the 307 optimum fit shown in Fig. 5A ($R_{wp} = 11.21\%$, $R_B = 5.50\%$ – Table 2). Selected interatomic 308 distances are listed in Table 4. Consistently with the calculated structural formula (eq 1), 309 KBi₁₀ consists of vacancy-free layers, the refined Mn_{laver} occupancy being very close to 1.0 310 (0.985). Deviation of layer symmetry from the hexagonal one induces a significant shift of 311 O_{laver} from the ideal anion close-packing sites [(0.342, 0, 0.070) – Table 3]. This position is 312 similar to that of O_{laver} in Mn³⁺-bearing birnessite varieties exhibiting an orthogonal layer 313 symmetry [5,8,70]. The refined K^+ site is split into three ones each being shifted in the **ab** 314 plane from the center of the interlayer prism towards its faces [(-0.253, 0, 0.25)] and (-0.134, 0.25)315 $\pm 0.375, 0.25$) – Fig. 4C, Table 3]. Similar positions were obtained if K⁺ was initially located 316 at the center of prism faces. Similar K⁺ positions were found for KBi₈ and hydrothermal KBi 317 [67,77]. H₂O_{interlayer} are also slightly shifted in the **ab** plane from their initial positions to form 318 a 132.9(3)° angle with the two nearest O_{laver} from adjacent layers [(0.127, 0, 0.25) – Figs 4C-319 D]. All attempts to split the H₂O_{interlayer} site from the edge of the prismatic cavity along [100], 320 [110] and $[1\overline{1}0]$ failed, the site along the **a** direction being systematically the sole with a final 321 322 significant occupancy.

*Rietveld Refinement of 350-KBi*₁₀ *Subcell*. The initial model was deduced from the 323 refined KBi₁₀ model, but 350-KBi₁₀ layers were considered to have a hexagonal symmetry, 324 the structure being described in the $P6_3/mmc$ space group. O_{laver} were first assumed to be 325 326 located in (1/3, 0, 0.078), K⁺ positions and occupancies were kept as in KBi₁₀, and no H₂O_{interlayer} was introduced according to TD-TG results. Structural parameters leading to the 327 best fit between experimental and calculated XRD patterns (Fig. 5B – R_{wp} = 16.80%, R_{B} = 328 13.22%) are given in Table 3 whereas selected interatomic distances are listed in Table 4. As 329 for KBi₁₀, 350-KBi₁₀ layers are essentially devoid of vacancies as the occupancy refined for 330 Mn_{laver} is close to 1.0 (0.983), in agreement with the structural formula proposed for KBi_{10} (eq 331 2). Due to the hexagonal layer symmetry O_{layer} are no longer shifted from the ideal anion 332 close-packing sites [(0.333, 0, 0.079) – Table 3]. Note however that this position was highly 333 unstable when refined. As discussed below this instability may arise from the presence in the 334 layers of Mn³⁺-octahedra with their long Mn³⁺-O bonds randomly oriented with respect to the 335 **a** axis. As for KBi₁₀ at room temperature, the refined K^+ site is split into three positions, 336 equivalent to (-0.300, 0, 0.25), shifted in the **ab** plane from the prism's center toward one of 337 its faces. Compared to KBi₁₀ structure this position is closer to the prism's center likely to 338 provide reasonable K-O_{laver} distances (2.73 Å Vs 2.97 Å, for 350-KBi₁₀ and KBi₁₀, 339 respectively) despite the interlayer collapse [77]. 340 Structure of Heated KBi₁₀ Samples. Between 100°C and 150°C a 20-to-2H structural 341 transformation occurs as described in the indexation section. The structure of 100-KBi₁₀ is 342

assumed to be identical to the 2*O* modification of KBi_{10} whereas 150-KBi₁₀ can be considered

as a physical mixture of hydrated 2O and dehydrated 2H structures (similar to samples KBi₁₀

and 350-KBi₁₀, respectively). At 250°C, dehydration and 2*O*-to-2*H* structural changes are

346 complete, and the structure of sample 250-KBi₁₀ was assumed to be identical to that of 350-

347 KBi₁₀. No structural refinement was performed on these three samples.

348	Rietveld Refinement of Vacuum-KBi10 Subcell. As for KBi10, small angular regions
349	containing super-reflections (at ~43.5° and ~52.0° 2 θ) were excluded from the refinement.
350	The initial structure model was identical to that of KBi_{10} without $H_2O_{interlayer}$ in agreement
351	with the dehydration under deep vacuum conditions. As for KBi_{10} , two PSFs were used to
352	account for the increased width of $20l/31l/40l$ diffraction lines compared to $11l/02l/22l$ ones.
353	The best fit is shown on Fig. 5C ($R_{wp} = 10.91\%$, $R_B = 6.85\%$). Refined O _{layer} positions (0.341,
354	0, 0.076) are similar to those in the original model, while refined K^+ positions are moved
355	toward the prism's center [(-0.300, 0, 0.25) and (0.153, ± 0.459 , 0.25)], as found for sample
356	350-KBi ₁₀ . Refined structural parameters and selected inter-atomic distances are listed in
357	Tables 2 and 3, respectively.

Qualitative and Quantitative Analysis of the EXAFS Spectra of KBi₁₀ and Related 358 Synthetic Birnessites. The short range order of Mn in KBi₁₀ was determined by using 359 synthetic Na- and H-birnessite (hereafter referred to as NaBi and HBi, respectively) as 360 references. The crystal structures of NaBi and HBi were determined by X-ray and electron 361 diffraction and EXAFS spectroscopy [8,46,47,80]. Triclinic NaBi consists of vacancy-free 362 Mn layers and its structural formula is $Na_{0,31}^+(Mn_{0,69}^{4+}Mn_{0,31}^{3+})O_2 \cdot 0.40H_2O$. Mn³⁺ octahedra, 363 which are elongated by the Jahn-Teller effect, are ordered in rows along [010] that are 364 separated by two Mn⁴⁺ rows along [100]. Due to this cation ordering and the distortion of the 365 Mn³⁺ octahedra, the layer symmetry departs from hexagonal to orthogonal [8,46]. The 366 structural formula of HBi at pH4 is 367

 $H^{+}_{0.33}Mn^{2+}_{0.043}Mn^{3+}_{0.123}(Mn^{4+}_{0.722}Mn^{3+}_{0.111}\square_{0.167})O_{2}(OH)_{0.013}, where \square represents vacancies$ $[80]. Its layer contains much less Mn^{3+} than NaBi (13%$ *Vs.*31%) and, as a consequence, ithas a hexagonal symmetry. HBi structure also differs from that of NaBi by the presence of $cationic layer vacancies, which are capped by interlayer Mn^{3+} and Mn^{2+} cations. As a result,$ NaBi has only edge-sharing Mn_{layer} octahedra whereas HBi contains in addition cornersharing Mn_{interlayer}-Mn_{layer} pairs with a distance separation of about 3.52 Å. Another important distinction between these two references is that the Mn_{layer}-Mn_{layer} EXAFS distances are equal
in HBi (2.89 Å) and unequal in NaBi (2.89 and 3.01 Å) [67], owing to the hexagonal *Vs*.
orthogonal layer symmetry.

Figs. 6A-B show that the resonance at k = 8.05 Å⁻¹ in HBi is split in two maxima at 7.8 377 and 8.1-8.2 Å⁻¹ in NaBi and KBi₁₀. Previous studies have shown that this indicator region is 378 sensitive to the amount and ordering of Mn^{3+}_{laver} [67,81-83]. When the layer has no Mn^{3+} , the 379 resonance peaks at 8.05 Å⁻¹. When it contains $1/3 \text{ Mn}^{3+}$ and $2/3 \text{ Mn}^{4+}$, and the Mn³⁺ cations 380 are not orderly distributed in rows, as in lithiophorite, then the hexagonal symmetry of the 381 layer is preserved and the resonance is shifted to lower k values (higher interatomic 382 distances), peaking at 7.9 Å⁻¹ [83]. In contrast, the resonance is split when the Mn-Mn 383 distances have essentially a bimodal distribution, that is when Mn³⁺_{laver} are segregated in 384 rows. Using that spectroscopic feature as a structural fingerprint for Mn ordering in 385 phyllomanganates, KBi₁₀ and NaBi have the same Mn³⁺-Mn⁴⁺ distribution patterns. The 386 absence of interlayer Mn atoms follows also from the comparison of the RSFs for KBi₁₀, 387 NaBi and HBi (Figs. 6C-D). The corner-sharing Mn_{interlayer}-Mn_{layer} correlation gives a peak at 388 $R + \Delta = 3.1$ Å [47,84], observed in HBi but neither in NaBi nor in KBi₁₀. Although Mn 389 clearly has a similar environment in KBi₁₀ and NaBi, the average bond lengths at the Mn site 390 are not strictly identical. The EXAFS spectrum of KBi_{10} is shifted towards the high k values 391 relative to NaBi (Fig. 6A), and this frequency modification induces a left-shift of the 392 imaginary part of the Mn shells in the real space (Fig. 6E). Comparing the imaginary parts for 393 the three phyllomanganates, we find that the average Mn-Mn distances decreases from NaBi, 394 to KBi₁₀, to HBi (Figs. 6E-F). This trend is consistent with the evolution of the fractional 395 amounts of Mn³⁺_{laver} in the three compounds: 0.31 in NaBi, 0.25 in KBi₁₀, and 0.13 in HBi. 396 The O shell has lesser sensitivity than the Mn shell to the amount of Mn^{3+}_{laver} because the four 397 equatorial Mn³⁺-O distances (~1.93 Å) [85], are similar to the Mn⁴⁺-O distances (~1.91-398 1.92 Å) [47,76,86], and because the two distant Mn³⁺-O pairs at 2.2-2.3 Å make a low 399

400 contribution to the total EXAFS signal (33% of the Mn-O pairs for 25% of the Mn cations, that is 8% of the total Mn-O pairs) [47]. Still, close examination of Figs. 6E-F shows that the 401 imaginary part of the O shell for KBi₁₀ is superimposed to that for HBi, but slightly left-402 403 shifted relative to that for NaBi. Consistently, the best-fit Mn-O EXAFS distance for KBi₁₀ was equal to 1.91 Å as in HBi, 0.01 Å shorter than in NaBi (Table 5). The similarity of the O 404 shell parameters for KBi₁₀ and HBi is explained with reference to the fact that the first RSF 405 peaks are the weighted sum from all the Mn atoms present in the structure. When Mn_{laver} and 406 $Mn_{interlayer}$ atoms in HBi are added, then HBi and KBi_{10} have almost the same Mn^{3+}/Mn_{total} 407 ratio (0.24 Vs. 0.25). 408

Similarly to NaBi, a two-shell fit to the Mn-Mn contribution provided an optimal 409 simulation to the data (Table 5). The short-distance Mn-Mn correlation at 2.89 Å is attributed 410 to the Mn^{4+} - Mn^{4+} and Mn^{3+} - Mn^{3+} pairs, and the long-distance Mn-Mn correlation at ~3.01-411 3.02 Å to the Mn⁴⁺-Mn³⁺ pairs [47]. In NaBi, the first subshell contains about 3.6 atoms and 412 the second about 2.4, while they contain about 3.5 and 1.6 atoms in KBi₁₀, respectively. 413 Although the differences in coordination numbers between the two samples are within 414 accuracy, the apparent decrease of the number of the long-distance pairs in KBi₁₀ is consistent 415 with the decrease of its Mn^{3+} content. The fact that the double antinode at ~8 Å⁻¹ is less 416 pronounced in KBi₁₀ than in NaBi (Fig. 6A) is definite evidence of the attenuation of the Mn 417 shell splitting caused by lower amounts of Mn^{3+} . 418

Overall, the structure models proposed here for KBi_{10} , and previously for HBi and NaBi [47], are consistent with XRD and EXAFS data. Upon closer examination, however, it appears that EXAFS gives shorter Mn-O and longer Mn-Mn distances than XRD, and that the number of oxygens in the first shell often is lower than six (Tables 3, 4 and 5). The discrepancy between the first shell structural parameters derived by the two techniques has a simple explanation. In Mn³⁺-containing phyllomanganates, the distribution of Mn-O distances is broad (from ~1.91 Å to 2.2-2.3 Å) and highly asymmetrical due to the elongation of the

Mn³⁺ octahedra. When the interlayer contains Mn atoms, as in HBi, the interlayer Mn-O and 426 Mn-H₂O distances are longer than the average Mn-O distance in the layer (Table 6). In these 427 mixed-valency and multi Mn sites compounds, the EXAFS analysis of the O shell with a 428 429 standard Gaussian-shell model excludes the full distribution of interatomic distances [87]. Since the tail of the distance separation is stretched towards higher R values (e.g., 2.2-2.3 Å 430 for Mn³⁺ and 1.98-2.08 Å for interlayer Mn), analysis of the data using symmetrical 431 distributions leads to lower numbers of O atoms around each Mn atom, and the EXAFS-432 derived <Mn-O> distance is thus shorter than the weighted average crystallographic value. 433 Cumulant analysis was tested in the present study but proved unsuccessful to detect the "lost" 434 atoms, because the tail apparently does not die off fast enough for cumulants of high orders to 435 be neglected (i.e., the asymmetry is too high). 436

This complication is circumvented sometimes in the literature by fixing the total number 437 of oxygens (CN) to six in the spectral fits [88,89]. This modelling strategy is inappropriate 438 when the structural disorder is high because the effective number of oxygen atoms detected 439 by EXAFS in the harmonic approximation is lower than six. In a recent study, the O shell of a 440 series of phyllomanganates was fitted with two Gaussian functions centred at 1.85-1.89 Å and 441 1.90-1.94 Å, and by constraining CN = 6 (Table 6) [90]. However, the distance separation 442 between the two functions is not high enough to include the longer bond lengths and the total 443 number of oxygens which can be recovered by this alternative model is thus lower than six. 444 For example, in NaBi, the number of oxygens at short distance is $0.69 \times 6 + 0.31 \times 4 = 5.4$. In 445 this recent study, the analysis of the Mn shell is also problematic for three reasons (Table 6). 446 First, NaBi does not have corner-sharing Mn octahedra. Second, if 0.8 Mn-Mn_{corner} pairs were 447 present in this compound, as assumed from EXAFS, then the number of Mn-Mn_{edge} pairs 448 should be lower than six, since interlayer Mn octahedra do not share edges with other Mn 449 octahedra. The same reasoning applies to HBi. Although this compound actually contains 450 interlayer Mn, the number of Mn-Mn_{edge} pairs was arbitrarily fixed to six by the authors. The 451

452 average number of Mn-Mn_{edge} pairs in HBi can be calculated from its structure as the
453 weighted sum of the different environments:

454
$$CN_{\text{edge}} = \sum_{i} W_{i} CN_{i}$$
(4)

where i refers to Mn site, W_i to Mn site occupancy, and CN_i to the number of Mn neighbors 455 for site i. For HBi, $CN_{edge} = 0.666 \times 4 + 0.166 \times 5 = 3.5$, as there are 50% vacancies in the 456 Mn³⁺-rich rows, and each Mn in these rows is surrounded by five Mn [47]. Our EXAFS CN 457 value for KBi_{10} (4.8) agrees with theory (3.5) within accuracy, but not the prescribed value of 458 six in the alternative model. Third, the Mn-Mn edge-sharing distances are split in NaBi, due 459 to the orthogonal symmetry of the layer, but not in HBi, whose layer symmetry is hexagonal. 460 From XRD data, the distance separation in NaBi is 2.95 - 2.85 = 0.10 Å (Table 6). In the 461 alternative EXAFS model, the distance separation in NaBi is as small as 2.90 - 2.88 = 0.02 Å 462 and, thus, incompatible with the orthogonal layer symmetry. The results reported for HBi are 463 even more questionable: not only the Mn-Mn distances were assumed to be split but, in 464 addition, in an amount (2.89 - 2.84 = 0.05 Å) higher than in NaBi. 465

The Mn-Mn edge-sharing distances obtained here by EXAFS for HBi (2.89 Å) and NaBi 466 (2.89 + 3.01 Å) are ~0.04 Å longer than those obtained by XRD (2.85 Å and 2.85 + 2.95 Å, 467 respectively – Tables 4 and 5). This difference is regarded as significant as EXAFS and XRD 468 give the same Mn-Mn distance in λ -MnO₂. The discrepancy between the lattice parameter 469 values (XRD) and the cation-cation distances (EXAFS) can be reconciled by considering that 470 the octahedral layer is corrugated. In increasing the Mn-Mn distance, the buckling of the 471 octahedral layer minimizes the cationic repulsion and contributes to increase the layer 472 stability. This effect has been described first in hydrotalcite by combining XRD, anomalous 473 474 scattering and EXAFS spectroscopy [91], and was subsequently inferred to occur also in birnessite from the quantitative analysis of the long distance Mn-Mn-Mn multiple scattering 475 paths in EXAFS [90,92], The angular deviation from the layer planarity, as estimated from 476

477 the difference of the EXAFS and XRD distances, is $\alpha = \cos^{-1}(2.85/2.89) \sim \cos^{-1}(2.95/3.01) \sim$ 478 10.5(10)° in HBi and NaBi.

Discussion

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The structure of high-temperature KBi synthesized at 1000°C (KBi₁₀ – this study) differs 482 substantially from that of KBi obtained at 800°C (KBi₈), whose structure was described in the 483 first of this four paper series [67]. Differences in chemical composition, layer symmetry and 484 origin of the layer charge between KBi₈ and KBi₁₀ are discussed first. Then, the nature of 485 structural modifications induced by thermal treatment at 350°C of KBi₁₀ is examined. From 486 these considerations and the comparison with structural features of other birnessites and layer 487 Mn oxyhydroxides, a generalized relationship between the layer symmetry and the origin of 488 the layer charge is proposed. Finally, the possible modification of the layer and interlayer 489 structure of KBi₁₀ in vacuum is discussed, and supporting arguments for the relevance of 490 SAED to study KBi₁₀ structure are presented. 491

492

Structure Model for KBi₁₀ Sample.

Chemical Composition of KBi₁₀ Layers and Layer Symmetry. One of the major differences 493 between KBi₈ and KBi₁₀ is their layer symmetry, as reflected by their a/b ratios (1.732 = $\sqrt{3}$ 494 and $1.811 = \sqrt{3.28}$, respectively) [67]. Drits et al. showed that the symmetry of the 495 phyllomanganate layers departs from being hexagonal when Mn³⁺ and Mn⁴⁺ cations are 496 orderly distributed in rows parallel to the **b** axis. Jahn-Teller distorted Mn^{3+} octahedra being 497 systematically elongated along the a axis (cooperative Jahn-Teller effect) [46]. The ordering 498 of heterovalent Mn cations in birnessite layers is energetically favorable as it minimizes the 499 steric strain induced by the presence of high amounts of distorted Mn³⁺ octahedra. In this 500 case, the unit-cell b dimension (2.84-2.86 Å) is not significantly modified because the b 501 direction is perpendicular to the direction of elongation of the Mn³⁺ octahedra and the Mn⁴⁺-O 502

and the four equatorial Mn^{3+} -O distances are equivalent. In contrast, the unit-cell *a* dimension 503 is increased and, consequently, the a/b ratio varies as a function of the Mn³⁺/Mn⁴⁺ ratio in the 504 octahedral layer. For example, in crednerite (CuMnO₂) a/b = 1.936 (Mn³⁺ = 100 %) [85], in 505 Na-rich buserite a/b = 1.835 (Mn³⁺ = 33 %) [70], in Na-rich birnessite a/b = 1.817 (Mn³⁺ = 506 31 %) [5,8,46], and in Ca-rich birnessite a/b = 1.808 (Mn³⁺ = 22%) [70]. Thus, the a/b ratio is 507 an indicator of the unique orientation of the elongated Mn³⁺ octahedra (cooperative Jahn-508 Teller effect) sensitive to the Mn^{3+}/Mn_{total} ratio in the layer when Mn^{3+} cations are ordered in 509 rows. Note that in lithiophorite $[(Al^{3+}_{0.67}Li^{+}_{0.32})(Mn^{4+}_{0.68}Mn^{3+}_{0.32})O_2(OH)_2]$, which contains 510 32% of layer Mn³⁺, the hexagonal symmetry of the layer is preserved ($a/b = \sqrt{3}$) because the 511 Mn^{3+} octahedra are elongated along three directions in the **ab** plane [83]. Thus, both the *a* and 512 b dimensions are enlarged relative to Mn^{3+} -free layers (b = 2.925 Å Vs. 2.842-2.850 Å) [93]. 513 KBi₁₀ provides new evidence for the relationship between the a/b and the Mn³⁺/Mn_{total} 514 ratios when the layer symmetry is orthogonal, as both its a/b ratio (1.811) and chemical 515 analysis indicate that it contains less Mn^{3+} than NaBi (0.25 Vs. 0.31) and that they are ordered 516 and systematically elongated along one direction. This finding is supported also by the 517 shallowness of the "dunce's cap" feature at $\sim 8 \text{ Å}^{-1}$ and the lower frequency of the EXAFS 518 spectrum for KBi₁₀ compared to NaBi (Fig. 6A). Lanson et al. showed that the amount of 519 layer Mn^{3+} could be estimated from the <Mn-O> distance, taking as end-members $<d(Mn^{3+}-Mn^{3+})$ 520 O)> = 2.04 Å, as calculated from crednerite and Mn oxyhydroxides, and $<d(Mn^{4+}-O)> =$ 521 1.912 Å, as calculated for λ -MnO₂ and hydrothermal birnessite (KBi_{hvdr}) [8,76,77,85,94]. The 522 <Mn-O> distance calculated for the (Mn⁴⁺0.75Mn³⁺0.25) layer cation composition of KBi₁₀ 523 (1.950 Å) is close to the experimental value (1.949 Å – Table 4). Similarly, the long Mn-O 524 distance calculated as the weighted sum of the Mn⁴⁺-O distance (1.912 Å in λ -MnO₂) and of 525 the long Mn³⁺-O distance (2.26 Å in credenerite) is 1.999 Å, and agrees with the XRD value 526 (2.021 Å – Table 4). 527

Chemical Composition of KBi₁₀ Layers and Origin of the Layer Charge. Another 528 important difference between KBi₈ and KBi₁₀ is the origin of their layer charge. KBi₈ has a 529 layer charge, which arises uniquely from cationic vacancies since the layer has no Mn^{3+} 530 cations. This charge is compensated for by the binding of $Mn_{interlayer}^{3+}$ at vacancy sites and the 531 sorption of K^+ [67]. In contrast, the layer charge deficit in KBi₁₀ arises mostly from the 532 substitution of 0.25 Mn^{4+}_{laver} by the same mole fraction of Mn^{3+}_{laver} and is compensated by 533 0.31 K^+ in the interlayer. According to XRD, layers are almost vacancy free. Additional 534 support for the absence of layer vacancies in KBi₁₀ follows from EXAFS results. Quantitative 535 analysis of the double antinode at $\sim 8 \text{ Å}^{-1}$ has shown that its unique "dunce's cap" shape 536 resulted from multiple scattering paths of the photoelectron between four aligned Mn shells in 537 the layer plane [83]. This high-order pair correlation, up to an effective radial distance of 538 $\sim 2.90 \times 3 = 8.7$ Å, can only be observed if the three successive Mn positions around a central 539 Mn site are filled. 540

Laver Symmetry and Orientation of Mn³⁺ Octahedra in Sample 350-KBi₁₀. Structure 541 models of KBi₁₀ at room temperature and at 350°C differ by their layer symmetries and, 542 consequently, by their unit-cell parameters and interatomic distances. The in-plane layer cell 543 parameters for hexagonal 350-KBi₁₀ (b = 2.894 Å – Table 1) are significantly higher than 544 those for hexagonal (1*H* and 2*H*) birnessites (2.844-2.848 Å at room temperature) 545 [58,67,80,95]. As a result, 350-KBi₁₀ has a longer < d(Mn-O) > (1.956 Å - Table 4) than 1H 546 and 2*H* birnessites at room temperature (~1.91 Å) but, for steric reasons, it still coincides with 547 the average distance in KBi₁₀ (1.949 Å – Table 4), since the two samples likely have the same 548 Mn³⁺/Mn_{total} ratio. The slight increase of the <Mn-O> distance (from 1.949 to 1.956 Å, for 549 KBi₁₀ and 350-KBi₁₀, respectively) likely is likely related to thermal expansion. 550 One likely explanation for the layer symmetry modification with thermal treatment is a 551 change of the azimuthal orientation of the elongation axes of the Mn³⁺ octahedra. If these axes 552

are oriented with equal probability at $\pm n120^{\circ}$ (*n* integer) directions, the resulting layer

symmetry is hexagonal. This configuration likely distorts locally the O_{laver} lattice, and may 554 explain the instability of the O_{laver} position during the structure refinement of 350-KBi₁₀. 555 Random orientation of Mn³⁺ octahedra induces unfavorable lattice strains that are likely 556 compensated for by the thermal energy available at 350°C. Such random distribution of Mn³⁺ 557 octahedra azimuthal orientation induces a significant increase of the unit-cell b parameter as 558 compared to that measured for vacancy-free layers or for layers with a unique orientation of 559 the long Mn^{3+} -O bonds in the **ac** plane. As mentioned above, such large *b* unit-cell dimension 560 was reported for lithiophorite (b = 2.912 Å) [93] whose layers have a preserved hexagonal 561 symmetry due to the random orientation of the Mn^{3+} octahedra in the **ab** plane [83]. 562 Consistently, the *b* parameter determined for 350-KBi₁₀ is increased compared to KBi₁₀ 563 (2.895 and 2.846 Å, respectively). As described by Gaillot et al. for hydrothermal KBi 564 samples [77], this abrupt increase with increasing temperature (Table 1) cannot result solely 565 from thermal motion as the unit-cell dilatation usually depends linearly on the temperature if 566 the structure is not modified. The 2O-to-2H transition with an increase of the b parameter 567 (from 2.846 to 2.895 Å) observed above 150°C can thus be considered as a direct evidence for 568 the presence of a significant amount of Mn³⁺ octahedra in KBi₁₀ layers and for their 569 systematic elongation along the **a** axis at room temperature. 570

General Relationship Between Layer Symmetry and Origin of the Layer Charge in 571 Birnessites. The comparison between structural features of KBi₈, KBi₁₀ and 350-KBi₁₀ can be 572 extended to all birnessite varieties, allowing general relationships between layer symmetry, 573 layer cation composition, <Mn-O> distance and the origin of the layer charge to be drawn. 574 HBi, KBi₈, KBi_{hvdr} are characterized by low amounts of Mn³⁺_{laver}, high layer vacancy contents 575 being responsible for the layer charge deficit. As a result, these species exhibit short </Mn-O> 576 distances (1.91-1.92 Å) and hexagonal symmetry with small b unit-cell dimensions (2.84-577 2.85 Å). In contrast, NaBi and KBi₁₀ layers are almost devoid of vacancy and the layer deficit 578 arises essentially from the presence of layer Mn^{3+} cations (25-30%), whose presence induces 579

a higher <Mn-O> distance (>1.94 Å). In the latter case Jahn-Teller distorted Mn³⁺ octahedra 580 are ordered at room temperature and their systematic elongation along the **a** axis leads to an 581 orthogonal layer symmetry. When temperature is increased up to 250-350°C, the symmetry of 582 these Mn³⁺-rich layers becomes hexagonal because of the random orientation of elongated 583 Mn^{3+} octahedra. In this case, $\langle Mn-O \rangle$ distance remains high (>1.94 Å) whereas the *b* unit-584 cell dimension is increased significantly (>2.86 Å). Similar structural features are observed at 585 room temperature when Mn³⁺ octahedra are ordered along the three equivalent directions 586 [100], [110] and $[1\overline{1}0]$ of the **ab** plane, as in lithiophorite. 587

This global multi-parameter relationship between layer cation composition, <Mn-O> distances, *b* parameter values and layer symmetry allows deducing fundamental crystalstructure information on the origin of the layer charge and on the location of Mn³⁺ cations, when present (within the octahedral layer or in the interlayer), even for poorly crystallized birnessite varieties. Such inference was recently demonstrated by Villalobos et al. for turbostratic birnessite samples and further discussed by Drits et al. [96,97].

Interlayer Structure of KBi₁₀. In KBi₁₀, the cohesion between adjacent layers mainly 594 results from the electrostatic interaction between the interlayer K⁺ cations and O_{laver}, H-bonds 595 between H₂O_{interlaver} and O_{laver} providing additional ties between layers. K⁺ cations occupy one 596 of the three possible sites in the interlayer prisms (Fig. 4). In such position, four K-O_{laver} 597 distances are shorter than the other two (~2.95 and ~3.28 Å, respectively – Table 4). A similar 598 shift of the interlayer K⁺ position was reported for KBi₈ and hydrothermal KBi_{hvdr} birnessite 599 varieties as resulting from the strong undersaturation of O_{layer} coordinated to only two Mn^{4+} 600 [67,77]. In KBi₁₀, whose layers are essentially vacancy-free, the presence of Mn^{3+}_{laver} is 601 responsible for the undersaturation of O_{laver}. When a O_{laver} is coordinated to three Mn⁴⁺ 602 octahedra, it ideally receives $0.667 \times 3 = 2$ valence units (v.u.) and is fully saturated. If this 603 O_{laver} is coordinated to two Mn⁴⁺ and one Mn³⁺, or to one Mn⁴⁺ and two Mn³⁺, it becomes 604 partially undersaturated, ideally receiving only $0.667 \times 2 + 0.500 = 1.833$ v.u. or 0.667 +605

606 $0.500 \times 2 = 1.667$ v.u., respectively. Therefore, a specific distribution of Mn⁴⁺ and Mn³⁺ 607 cations within adjacent layers (Fig. 7), and the resulting distribution of O_{layer} with various 608 degrees of undersaturation, may account for the shift of the K⁺ cations within the interlayer 609 prisms, so as to better achieve local charge compensation of most undersaturated O_{layer}.

As found for samples KBi₈ and KBi_{hydr}, H₂O_{interlayer}, which are located in the interlayer mid-plane, are slightly shifted in the **ab** plane from the prism's edges defined by O_{layer} from adjacent layers towards the nearest Mn_{layer} (Fig. 4B) [67,77]. However, in KBi₁₀ H₂O_{interlayer} are shifted only along the **a** axis, refined occupancy of equivalent sites being nil. Resulting H₂O-O_{layer} distances and O_{layer}-H₂O-O_{layer} angle, equal to 2.77 Å and 132.9°, respectively (Table 4), are typical for H-bonds.

Layer and Interlayer Structures of Dehydrated Samples KBi₁₀. When heated up to 616 350°C sample KBi₁₀ dehydrates, and the basal spacing is decreased from ~7.05 Å to ~6.39 Å 617 (Table 1) due to the departure of H₂O_{interlaver}. Additional structural changes resulting from the 618 heating include layer symmetry increase (20-to-2H transition) as a result of the loss of 619 cooperative Jahn-Teller effect. Despite the symmetry change and the partial interlayer 620 collapse, position of interlayer K^+ cations is essentially unchanged. Such split of the K^+ site 621 has been reported for hexagonal KBi structures (KBi₈ and KBi_{hvdr}), both at room temperature 622 and at 350°C [67,77]. 623

624 For sample Vacuum-KBi₁₀, the cooperative Jahn-Teller effect is preserved as no additional energy is available. Apart from the departure of H₂O_{interlayer} and from the resulting 625 interlayer collapse (from \sim 7.05 Å to \sim 6.47 Å), the structure of KBi₁₀ and in particular its 626 interlayer structure is maintained. K^+ cations occupy split positions similar to those in KBi₁₀, 627 the six-fold coordination being provided by the nearest O_{laver} from adjacent layers with only 628 limited migration of K⁺ towards the center of the interlayer prism to allow for more 629 appropriate K-O_{laver} distances. This similarity between hydrated and dehydrated interlayer 630 structures of high-temperature KBi varieties thus allows using SAED to study the distribution 631

of interlayer K^+ cations in dehydrated KBi₁₀ structure, and extrapolating the obtained results to the hydrated structure. Results from such a study will be reported in the fourth paper of the series devoted to the study of high-temperature KBi varieties [66].

Origin of Sample AfterT-KBi₁₀ Structural Heterogeneity. The main difference 635 between KBi₁₀ specimens before and after the thermal treatment is a significant broadening, 636 and weakening, of the 20*l* reflections in sample After-KBi₁₀ (Figs. 2 and 3). In addition, the 637 fwhm of 201 reflections of sample After-KBi₁₀ decreases with increasing l index (solid circles 638 - Fig. 3). These effects are similar to those described for heterogeneous KBi_{10h} birnessite 639 [68]. Gaillot et al. demonstrated that these effects arise from the coexistence in the sample of 640 several populations of particles which all have a 20 structure but differ from each other by 641 their unit-cell parameters that are related by eq 3 [68]. Two hypotheses were explored to 642 account for this specific variation of the a and b parameters. First, local fluctuations of the 643 redox conditions during cooling of sample 350-KBi₁₀ may have led slightly different Mn³⁺ 644 contents in individual crystals. This hypothesis is incompatible with the observed variation of 645 the *a* and *b* unit-cell parameters, the a/b ratio determined for AfterT-KBi₁₀ being slightly 646 lower (1.805 Å) than that of KBi₁₀ (1.811 Å). If in both specimens layer Mn^{3+} octahedra were 647 systematically elongated along the **a** axis, the *a* parameter decrease would then be related with 648 the Mn³⁺-to-Mn⁴⁺ oxidation during the thermal treatment. However, at 350°C, reduction 649 rather than oxidation was observed for hydrothermal KBi_{hvdr} and pyrolusite (MnO₂) 650 [68,77,98]. 651

In the second hypothesis, the unit-cell dimension heterogeneity is related to contrasting orientation distribution of the long Mn^{3+} -O_{layer} bonds with respect to the **a** axis [77]. The maximum *a/b* ratio corresponds to KBi₁₀ particles in which most Mn^{3+} octahedra are oriented with their long bonds along the **a** axis whereas in particles with lower *a/b* ratios the long Mn^{3+} -O bonds of some Mn^{3+} octahedra are oriented at ±120° with respect to the **a** axis. Such different azimuthal orientation distributions of Mn^{3+} octahedra likely originates from the

658	cooling rate of 350-KBi ₁₀ which was too fast to allow a complete $2H$ -to- $2O$ structure
659	transformation. During this "air quenching" the random orientation of a small fraction of
660	elongated Mn ³⁺ octahedra was "frozen" in some KBi ₁₀ particles.

661

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824	KBi ₁₀ , respectively).(F) Pattern recorded after complete thermal treatment to 350°C
825	and subsequent cooling to room temperature (sample AfterT-KBi10). (G) Pattern
826	recorded under vacuum condition (sample Vacuum-KBi10). (H) Pattern recorded
827	after the sample was brought back to atmospheric pressure (sample AfterVac-
828	KBi_{10}). Dotted and dot-dashed lines indicate the position of $00l$ reflections for
829	hydrated and dehydrated samples, respectively. Intensity scale is enlarged over the
830	30-80° 20 CuK α angular range, by 9×, 18×, 14×, 13×, 7×, 10×, 14× and 17×, for
831	XRD pattern shown in (A) to (G), respectively. Reflections are indexed using
832	orthorhombic C-centered 20 (A-C, F-G) or 2H (D-E) unit cells (see supplementary
833	material). Super-reflections are indicated by star symbols.
834	Fig. 3. Evolution of the full width at half-maximum intensity (fwhm) of 20 <i>l</i> and 11 <i>l</i>
835	reflections (circles and triangles respectively) as a function of the Miller index l for
836	samples KBi_{10} (open symbols) and AfterT- KBi_{10} (solid symbols). Experimental
837	fwhm values are corrected by a $\cos\theta$ factor to account for crystal-size broadening
838	[78].
839	Fig. 4. Structure model for KBi ₁₀ . (A) Projection along the b axis. Open and solid symbols
840	indicate atoms at $y = 0$ and $y = \pm 1/2$, respectively. Large circles represent O _{layer}
840 841	indicate atoms at $y = 0$ and $y = \pm 1/2$, respectively. Large circles represent O _{layer} atoms, small circles represent Mn _{layer} atoms, shaded circles represent K _{interlayer} and
841	atoms, small circles represent Mn_{layer} atoms, shaded circles represent $K_{interlayer}$ and
841 842	atoms, small circles represent Mn_{layer} atoms, shaded circles represent $K_{interlayer}$ and open circles with dashed outline represent $H_2O_{interlayer}$. Dot-dashed lines outline the
841 842 843	atoms, small circles represent Mn_{layer} atoms, shaded circles represent $K_{interlayer}$ and open circles with dashed outline represent $H_2O_{interlayer}$. Dot-dashed lines outline the interlayer prisms defined by two empty tridentate layer cavities. The center of these
841 842 843 844	atoms, small circles represent Mn_{layer} atoms, shaded circles represent $K_{interlayer}$ and open circles with dashed outline represent $H_2O_{interlayer}$. Dot-dashed lines outline the interlayer prisms defined by two empty tridentate layer cavities. The center of these prisms is shown by regular dashed lines, and the arrows outline the shift of K ⁺
841 842 843 844 845	atoms, small circles represent Mn_{layer} atoms, shaded circles represent $K_{interlayer}$ and open circles with dashed outline represent $H_2O_{interlayer}$. Dot-dashed lines outline the interlayer prisms defined by two empty tridentate layer cavities. The center of these prisms is shown by regular dashed lines, and the arrows outline the shift of K ⁺ cations from this ideal position. Solid lines represent H-bonds between O_{layer} and
 841 842 843 844 845 846 	atoms, small circles represent Mn_{layer} atoms, shaded circles represent $K_{interlayer}$ and open circles with dashed outline represent $H_2O_{interlayer}$. Dot-dashed lines outline the interlayer prisms defined by two empty tridentate layer cavities. The center of these prisms is shown by regular dashed lines, and the arrows outline the shift of K ⁺ cations from this ideal position. Solid lines represent H-bonds between O_{layer} and $H_2O_{interlayer}$, and the arrows outline the shift of H ₂ O along the a axis from the edge
 841 842 843 844 845 846 847 	atoms, small circles represent Mn_{layer} atoms, shaded circles represent $K_{interlayer}$ and open circles with dashed outline represent $H_2O_{interlayer}$. Dot-dashed lines outline the interlayer prisms defined by two empty tridentate layer cavities. The center of these prisms is shown by regular dashed lines, and the arrows outline the shift of K ⁺ cations from this ideal position. Solid lines represent H-bonds between O_{layer} and $H_2O_{interlayer}$, and the arrows outline the shift of H_2O along the a axis from the edge of interlayer prisms. (B) Projection on the ab plane. The upper surface of the lower

- 850 represent interlayer potassium, and open circles with dashed outline represent water 851 molecules. The arrows outline the shift of H_2O along the **a** axis.
- Fig. 5. Comparison between experimental (crosses) and calculated (solid line) XRD patterns
 for KBi₁₀ samples. (A) Sample KBi₁₀, (B) sample 350-KBi₁₀ and (C) sample
 Vacuum-KBi₁₀. Atomic coordinates, site occupancies and other structural
 parameters used for the calculations are listed in Table 3. Difference plot is given
 for each case. Angular regions with significant contributions from super-reflections
 were excluded from the fits.
- Fig. 6. k^3 -weighted EXAFS spectra and Fourier transforms (modulus and imaginary parts) of *k*-weighted EXAFS spectra for KBi₁₀, triclinic 1*T* NaBi and hexagonal 1*H* HBi (pH4). The EXAFS spectrum of KBi₁₀ has the same shape as that of NaBi, but a lower frequency, which manifests itself in a shift to lower distance of the Mn-O and Mn-Mn pairs in the real space.
- Fig. 7. Idealized distribution in projection on the **ab** plane of undersaturated O_{laver} within an 863 interlayer region. 4+ represent Mn⁴⁺ cations located in the lower and/or in the upper 864 octahedral layer. 3+ on black and white backgrounds represent Mn^{3+} cations located 865 in the lower and upper octahedral layers, respectively. Circles represent Olaver from 866 both the upper and lower surfaces of adjacent octahedral layers whose positions 867 coincide in projection on the **ab** plane and define the edges of interlayer prisms. 868 Different circle symbols are used to represent the degree of under-saturation of the 869 O-O edge as a function of the total number of Mn^{3+}_{laver} coordinated to it (from 0-4 – 870 See figure for equivalence). (A) Mn^{3+} -rich rows from adjacent layers coincide in 871 projection on the **ab** plane. (B) Mn^{3+} -rich rows from adjacent layers do not coincide 872 in projection on the **ab** plane (shift along the **a** axis). 873

	a^a	b	С	a/b	Polytype ^b
KBi ₁₀	5.1554(3)	2.8460(1)	14.088(1)	1.811	20
100-KBi ₁₀	5.141(1)	2.850(0)	14.002(2)	1.804	20
150-KBi ₁₀	5.117(1)	2.845(0)	12.860(2)	1.798	20
150-KBi ₁₀	5.010(3)	2.891(1)	12.871(4)	1.733	2H
350-KBi ₁₀	5.010(2)	2.894(1)	12.954(7)	1.731	2H
AfterT-KBi ₁₀	5.149(3)	2.846(1)	14.042(7)	1.809	20
Vacuum-KBi ₁₀	5.114(2)	2.840(1)	12.787(3)	1.801	20
AfterVac-KBi ₁₀	5.160(2)	2.844(1)	14.072(5)	1.815	20

Table 1: Unit-Cell Parameters of KBi10 Samples

^{*a*} All parameters are expressed in Å. $\alpha = \beta = \gamma = 90^{\circ}$. ^{*b*} Indexations for 20 and 2*H*

polytypes are given as supplementary material.

Ideal formula	$K^{+}_{0.314}(Mn^{3+}_{0.246}Mn^{4+}_{0.737}\square_{0.017})O_{2}{\cdot}0.50H_{2}O$
Wavelength	1.54056 Å + 1.54439 Å (Rel. Int. 0.476)
Space group	Ccmm
	a = 5.1554(3) Å
Unit-cell dimensions	b = 2.8460(1) Å
Unit-cen dimensions	c = 14.088(1) Å
	$V = 206.70 \text{ Å}^3$
Angular range	34-90°
Data / Parameters	1319 / 13
Number of reflections	100
R_{Exp}	3.5%
R _{wp}	11.21%
R _{Bragg}	5.55%
Goodness of fit	3.16

Table 2: Crystal data and structure refinement for the KBi_{10} sample

	KBi ₁₀					Vacuum-KBi ₁₀			350-KBi10			
	\mathbf{x}^{a}	У	Z	Occ.	X	у	Z	Occ.	X	У	Z	Occ.
Mn _{layer}	0	0	0	0.985(6)	0	0	0	0.983	0	0	0	0.983
O _{layer}	0.342(1)	0	0.0700(6)	2.000	0.341	0	0.0762	2.000	0.333	0	0.0788	2.000
K	-0.253(24)	0	1/4	0.077(28)	-0.300	0	1/4	0.121	-0.300	0	1/4	0.104
K	0.134(8)	±0.375(23)	1/4	0.235(28)	0.153	±0.459	1/4	0.191	0.150	±0.450	1/4	0.208
H ₂ O	0.127(5)	0	1/4	0.500								

^a Atomic positions are given according to orthogonal axes, even for 2*H* 350-KBi₁₀. The values of Debye-Waller factors (B) are 0.5, 1.0, 2.0,
2.0 for Mn_{layer}, O_{layer}, K and H₂O, respectively.

		KBi ₁₀	Vacuum-KBi ₁₀	350-KBi ₁₀
Height of M	n layer	1.972 ^{<i>a</i>}	1.974	2.015
	short	2.846(0) ×2	2.839 ×2	
Mn _{layer} -Mn _{layer}	long	2.944(0) ×4	2.924 ×4	2.895
5	average	2.911	2.896	
	short	1.914(5) ×4	1.903 ×4	
Mn _{layer} -O _{layer}	long	2.018(6) ×2	1.998 ×2	1.958
	average	1.949	1.934	
КO	short	2.948(28)-2.973(23) ×4	2.735 ×4	2.731 ×4
K-O _{layer}	long	3.277(78)-3.287(38) ×2	2.868 ×2	2.880 ×2
	average	3.061-3.067	2.784	2.808
Kinterlayer-H2Ointerlayer		3.197(137) - 3.611(49)		
Olayer-H2Oi		2.767(12)		
O _{layer} -H ₂ O-O _{la}	ayer angle	132.88(25)°		

Table 4. Selected Inter-Atomic Distances in High-Temperature KBi₁₀ Samples

^{*a*} All distances in Å.

Sample	$R+\Delta R$ window (Å)	fit interval (Å ⁻¹)	shell	$R^{a}(A)$	CN^b	$\sigma^{2c}_{\cdot}(\text{\AA}^2)$	$\Delta E^{d} (\mathrm{eV})$	residual ^e
λ -MnO ₂	1.1 - 3.3	3.7 < <i>k</i> < 13.5	Mn-O	1.91	6	25×10^{-4}	1.5	12
			Mn-Mn _{edge}	2.85	6	28×10^{-4}		
			Mn-O ^f	3.52	6 ^{<i>g</i>}	88×10^{-4}		
HBi ^h	1.1 - 3.4	3.7 < <i>k</i> < 12	Mn-O	1.91	4.5	32×10^{-4}	1.6	7
			Mn-Mn _{edge}	2.89	4.8	$51 \times 10^{-4 i}$		
			Mn- Mn _{corner}	3.49	2.8	$51 \times 10^{-4 i}$		
			Mn-O ^f	3.61	6 ^{<i>g</i>}	49×10^{-4}		
NaBi ^{<i>h</i>}	1.1 - 3.7	3.7 < <i>k</i> < 13.5	Mn-O	1.92	5.6	30×10^{-4}	0.7	14
			Mn-Mn _{edge}	2.89	3.6	$40 \times 10^{-4 i}$		
			Mn-Mn _{edge}	3.01	2.4	$40 \times 10^{-4 i}$		
			Mn-O ^f	3.57	6 ^{<i>g</i>}	12×10^{-3}		
KBi ₁₀	1.1 - 3.7	3.7 < <i>k</i> < 12	Mn-O	1.91	4.8	29×10^{-4}	1.3	14
			Mn-Mn _{edge}	2.89	3.5	$29 \times 10^{-4 i}$		
			Mn-Mn _{edge}	3.02	1.6	$29 \times 10^{-4 i}$		
			Mn-O ^f	3.58	6 ^{<i>g</i>}	14×10^{-3}		

 Table 5. EXAFS Parameters for the Nearest Mn-O and Mn-Mn Pairs in Birnessite Samples

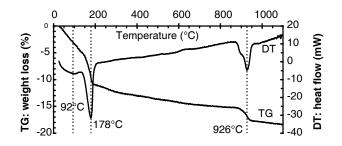
and in the Reference λ -MnO₂

^{*a*} Accuracy and precision in average distances are ±0.02 and ±0.01 Å, respectively. In λ -MnO₂, the crystallographic Mn-O and Mn-Mn distances are 1.91 Å and 2.84 Å, respectively [76]. ^{*b*} Coordination number. The scaling factor S_0^2 was calculated to obtain CN = 6 in the reference λ -MnO₂ : $S_0^2 = 0.73$ for the Mn-O pair and 0.80 for the Mn-Mn pair [96]. Typical accuracy on coordination numbers is ±1.5. ^{*c*} Debye-Waller factor. ^{*d*} Variation of the energy threshold treated as a single adjustable parameter for all subshells. ^{*e*} Residual is calculated from $R = [\Sigma|(k^3\chi_{exp}-k^3\chi_{cal})|/\Sigma|k^3\chi_{exp}|] \times 100$. ^{*f*} Second Mn-O shell. ^{*g*} Fixed value in the optimization procedure. ^{*h*} EXAFS results are, within accuracy, identical to those published by Silvester et al. [47]. ^{*i*} Parameter varied but constrained equal for the two subshells.

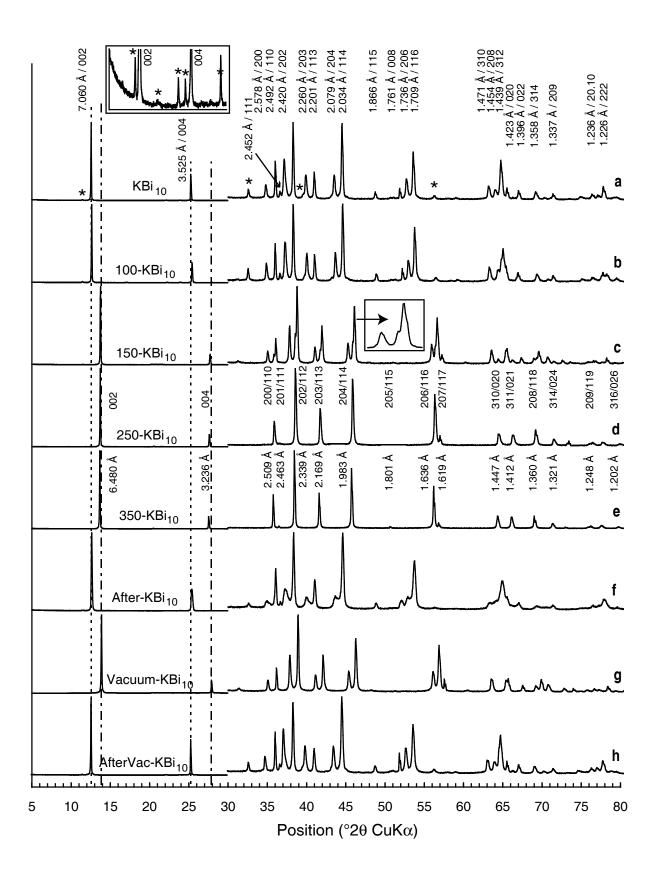
Table 6. EXAFS Parameters for the Nearest Mn-O and Mn-Mn Pairs in Birnessite Reportedby Webb et al. [90]

Sample	shell	XR	D^a	EXAFS				
		<i>R</i> (Å)	CN^b	<i>R</i> (Å)	CN	$\sigma^2(\text{\AA}^2)$		
NaBi	Mn-O	1.92	2	1.87 (2)	4^b	0.005 (2)		
		1.94	2	1.94 (1)	2^b	0.002(1)		
		2.00	2					
	Mn-Mn _{edge}	2.85	2	2.88 (3)	2^c	$0.006(1)^d$		
	0	2.95	4	2.90 (3)	4^c	$0.006(1)^d$		
	Mn-Mn _{corner}	-	-	3.36 (9)	0.8 (8)	0.006 (3)		
HBi	Mn _{laver} -O	1.92	5	1.89 (3)	4^c	0.008 (2)		
	Mn _{interlayer} -O	1.98	0.5	1.90 (2)	2^c	0.004 (2)		
	Mn-H ₂ Ó	2.08	0.5					
	Mn-Mn _{edge}	2.85	3.5	2.84(1)	2^c	0.005 (1) d		
	0			2.89(1)	4^c	$0.005(1)^d$		
	Mn-Mn _{corner}	3.54	1.8	3.48(1)	2.0 (1.1)	0.003 (3)		

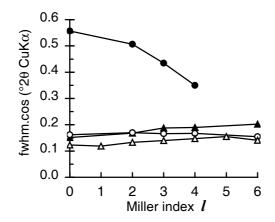
^{*a*} Data from Lanson et al. [8,80]. ^{*b*} Theoretical EXAFS CN values calculated from XRD data. ^{*c*} *CN* values were fixed during the fit. ^{*d*} Parameter varied but constrained equal for the two subshells.



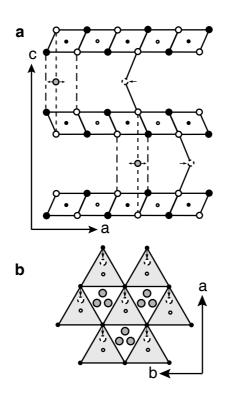
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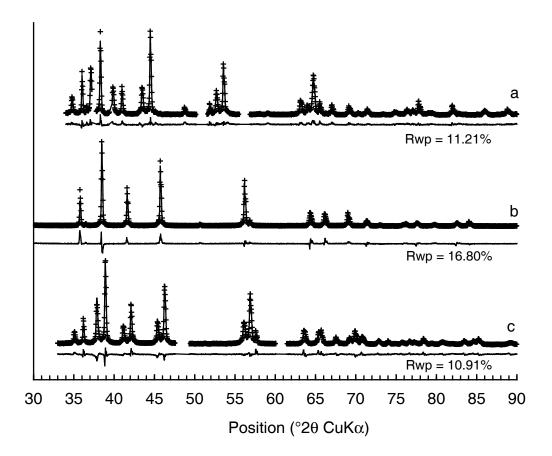
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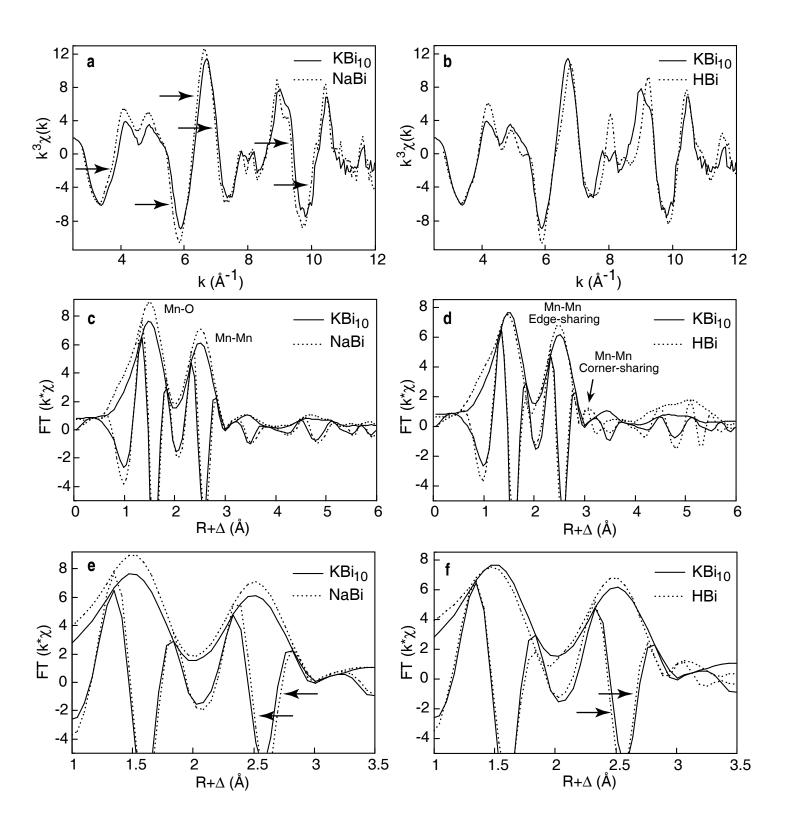
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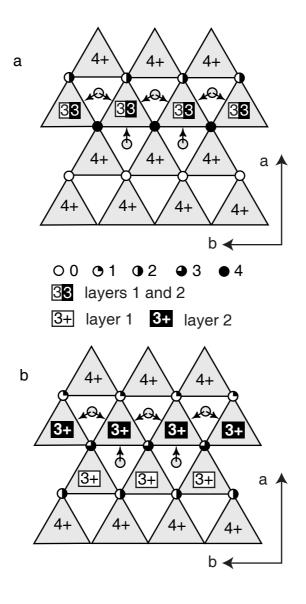




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Indexing in Terms of the 2*O* Polytype of the Experimental XRD Patterns of KBi₁₀, AfterT-KBi₁₀, Vacuum-KBi₁₀, 100-KBi₁₀, and 150-KBi₁₀ (Main Phase).

20	KBi ₁₀			-KBi ₁₀		n-KBi ₁₀	100-KBi ₁₀		150-KBi ₁₀	
hkl	d_{\exp}^{a}	$d_{\rm calc}{}^b$	d_{exp}	$d_{ m calc}$	d_{exp}	$d_{ m calc}$	$d_{ m exp}$	$d_{ m calc}$	$d_{ m exp}$	$d_{ m calc}$
	7.750°	7.733	7.691		7.653		7.722		7.706	
002^{d}					6.944					
002	7.060	7.050	7.021	7.022	6.394	6.390	6.995	7.000	6.435	6.430
	6.784	6.780					5.191			
	5.206	5.210			4.375		3.998		4.385	
	4.013	4.016	4.006		4.257		3.866		4.260	
	3.730	3.729	3.719		3.723		3.718		3.754	
004					3.486					
004	3.525	3.525	3.506	3.511	3.195	3.195	3.501	3.500	3.216	3.215
	3.202	3.207			3.111		3.193			
	2.745	2.745	2.744		2.850		2.748		2.857	
200	2.578	2.578	2.572	2.569	2.557	2.556	2.572	2.571	2.559	2.559
110	2.492	2.492	2.488	2.489	2.483	2.481	2.494	2.493	2.488	2.487
111	2.452	2.454		2.451		2.436	2.455	2.455		2.442
202	2.420	2.421	2.414	2.412	2.375	2.373	2.413	2.413	2.377	2.377
	2.403	2.405								
112	2.348	2.349	2.345	2.346	2.313	2.313	2.348	2.349	2.319	2.320
006		2.349		2.341		2.130		2.333		2.144
	2.276	2.278								
203	2.260	2.260	2.255	2.252	2.192	2.192	2.252	2.252	2.197	2.197
113	2.201	2.201	2.198	2.198	2.144	2.144	2.199	2.199	2.152	2.151
							2.091			
204	2.079	2.081	2.071	2.073	1.997	1.996	2.072	2.072	2.003	2.002
114	2.034	2.035	2.030	2.031	1.961	1.960	2.030	2.036	1.967	1.967
							2.007			1.814
205				1.895		1.807		1.894		
115	1.866	1.867	1.864	1.863		1.780	1.862	1.862	1.787	1.788
008	1.761	1.763	1.758	1.755	1.599	1.598	1.751	1.750	1.608	1.609
							1.797			
206	1.736	1.737	1.731	1.730	1.637	1.636	1.728	1.728	1.643	1.643
116	1.709	1.710	1.705	1.705	1.617	1.616	1.704	1.704	1.624	1.624
					1.553		1.691		1.536	
	1.634		1.632		1.528		1.630		1.501	
					1.504		1.560			
					1.487					
310	1.471	1.471	1.467	1.467	1.462	1.461	1.469	1.469	1.463	1.463
208	1.454	1.455	1.451	1.449	1.356	1.355	1.447	1.447	1.361	1.361
118	1.439	1.439	1.436	1.435	1.345	1.343	1.433	1.432	1.350	1.350
312	1.439	1.439	1.436	1.436	1.425	1.424	1.437	1.437	1.426	1.427
020	1.423	1.423	1.423	1.423	1.421	1.419	1.426	1.426		1.423
022	1.396	1.395	1.395	1.395	1.387	1.385	1.396	1.397	1.389	1.389
							1.387		1.380	
314	1.358	1.358	1.356	1.354	1.330	1.329	1.354	1.354	1.331	1.332
209	1.337	1.339	1.334	1.334	1 200	1.00-	1.331	1.331	1.001	
024	1.319	1.320	1.320	1.319	1.298	1.297	1.320	1.320	1.301	1.301
119	1.319	1.302	1.320	1.322			1.320	1.320		
00.10				1.404	1.050	1.278		1.400	1.286	1.286
400				1.284	1.279	1.278		1.285		1.280
					1.267					
400				1.0.52	1.260	1.050	1.0.4	1.0.4	1.054	1
402				1.263	1.254	1.253	1.264	1.264	1.256	1.255
000	1.045	1.045	1.0/5	1.045	1.249	1.0/1	1.01-	1.0.47	1.0.42	1
220	1.246	1.246	1.245	1.245	1.242	1.241	1.246	1.247	1.243	1.244

316			1.245	1.243			1.246	1.243	1.209	1.208
20.10	1.236	1.237		1.232				1.229	1.149	1.149
222	1.226	1.227	1.226	1.226	1.219	1.218	1.227	1.227	1.221	1.221
11.10	1.226	1.227	1.226	1.223	1.136	1.136	1.221	1.221	1.142	1.142
	1.209						1.239			
223			1.204	1.203		1.193	1.205	1.204	1.194	1.194
404			1.204	1.206		1.187	1.205	1.206	1.188	1.189
224	1.175	1.175	1.173	1.173	1.157	1.157	1.174	1.174	1.160	1.160
00.12	11170	11170	1.173	1.173	11107	11107	1.167	1.167	11100	11100
00.12			1.175	1.175			1.160	1.107		
					1.131		1.155			
405			1.169	1.168	1.145	1.143	1.155			
11.11			1.107	1.136	1.145	1.145	1.133	1.134		
225			1.130	1.130		1.116	1.133	1.134		
-			1.150							
318				1.126		1.078	1.125	1.125		
406				1.126			1.125	1.126	1.098	1.098
226				1.099			1.099	1.100		

^{*a*} $d_{exp}(hkl)$ are measured experimentally. ^{*b*} $d_{calc}(hkl)$ values are calculated using the unit-cell parameters listed in Table 1. ^{*c*} d(hkl) values in italics correspond to super-reflections (see companion article). ^{1 *d*} 002 and 004 reflections from remnant hydrated crystals.

2 <i>H</i>	150-1	KBi ₁₀	350-1	KBi ₁₀
$hk\ell$	d_{\exp}^{a}	$d_{\rm calc}^{\ \ b}$	d_{\exp}	$d_{ m calc}$
002	6.435	6.441	6.480	6.489
	4.385 ^c		4.341	
	4.260		4.130	
	3.754			
004	3.216	3.221	3.236	3.245
200/110	2.502	2.505	2.509	2.510
201/111	2.488	2.459	2.463	2.464
202/112	2.335	2.335	2.339	2.341
203/113	2.163	2.164	2.169	2.171
204/114	1.978	1.978	1.983	1.985
205/115		1.796	1.801	1.804
206/116	1.630	1.630	1.636	1.639
008	1.608	1.610	1.619	1.622
310/020	1.445	1.446	1.447	1.450
312/022	1.410	1.411	1.412	1.415
208/118	1.354	1.355	1.360	1.362
314/024	1.319	1.319	1.321	1.323
00.10	1.286	1.288	1.295	1.298
400/220	1.243	1.253	1.253	1.255
209/119	1.243	1.243	1.248	1.250
402/222	1.229	1.230		1.232
316/026	1.209	1.200	1.202	1.204
404/224	1.167	1.168	1.168	1.171
20.10/11.10	1.149	1.146	1.151	1.153
406/226	1.098	1.083		1.086

Indexing in Terms of 2H Polytypes of the Experimental XRD Patterns of 150-KBi₁₀ (Accessory Phase) and 350-KBi₁₀.

^{*a*} $d_{\exp}(hk\ell)$ are measured experimentally. ^{*b*} $d_{calc}(hk\ell)$ values are calculated using the unitcell parameters listed in Table 1. ^{*c*} d(hkl) values in italics correspond to super-reflections (see companion article).¹

1. A.-C. Gaillot, V. A. Drits and B. Lanson, in preparation.