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► To cite this version:

Pascal Auscher, Frédéric Bernicot, Jiman Zhao. Maximal regularity and Hardy spaces. *Collectanea Mathematica*, 2008, 59 (1), pp.103-127. hal-00179552

HAL Id: hal-00179552

<https://hal.science/hal-00179552>

Submitted on 16 Oct 2007

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Maximal regularity and Hardy spaces

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October 10, 2007

Abstract

In this work, we consider the Cauchy problem for $u' - Au = f$ with A the Laplacian operator on some Riemannian manifolds or a sublapacian on some Lie groups or some second order elliptic operators on a domain. We show the boundedness of the operator of maximal regularity $f \mapsto Au$ and its adjoint on appropriate Hardy spaces which we define and study for this purpose. As a consequence we reobtain the maximal L^q regularity on L^p spaces for $1 < p, q < \infty$.

Key words: maximal regularity, Laplace-Beltrami operator, heat kernel, Hardy spaces, atomic decomposition.

AMS2000 Classification: 34G10, 35K90, 42B30, 42B20, 47D06

1 Introduction

Let B be a Banach space, A the infinitesimal generator of a bounded analytic semigroup of operators on B and $I = (0, +\infty)$ (I bounded is also of interest). We shall consider the problem

$$\begin{cases} u'(t) - Au(t) = f(t), & t \in I, \\ u(0) = 0, \end{cases} \quad (1.1)$$

where $f : I \rightarrow B$ is given. If T_t is the semigroup generated by A , then u is given by

$$u(t) = \int_0^t T_{t-s} f(s) ds.$$

For fixed $q \in (1, +\infty)$, one says that there is maximal L^q regularity if for every $f \in L^q(I, B)$, u' and Au belongs to $L^q(I, B)$ with

$$\|u'\|_{L^q(I, B)} + \|Au\|_{L^q(I, B)} \leq C\|f\|_{L^q(I, B)}. \quad (1.2)$$

When B is an L^p space ($1 < p < \infty$) we refer to [5], [8], [10], [11], [15], [17] etc. When B is an UMD space, [23] gives a necessary and sufficient condition for maximal L^q regularity in terms of a notion called R-boundedness. An excellent survey is [16].

Let us come back to $B = L^p(\Omega, d\mu)$. A classical method to prove maximal regularity (already in [11]) relies on singular integral theory. It is well-known that maximal L^q regularity for one $q \in (1, +\infty)$ implies the same for all $q \in (1, +\infty)$. Hence, it is enough to prove (1.2) for $q = p$, that is the boundedness of the map of maximal regularity $f \mapsto Au$ on $L^p(I \times \Omega, dtd\mu)$. Seen as a singular integral on $I \times \Omega$ equipped with a parabolic distance and product measure $dtd\mu$, one proves this map and its adjoint have weak type (1,1) and applies interpolation. This method has been refined so as to relax hypotheses as much as possible on Ω and A . Completing a theorem in [15], this method has been successful for all $1 < p < \infty$ in [10] with Ω a (subset of a) space of homogeneous type and pointwise Gaussian upper bounds for the kernel of the semigroup and even generalized to a restricted range of p in [3] under weaker generalized Gaussian estimates.

There exists a criterion for L^1 maximal regularity on L^1 (see [14]) but when A is a second order differential operator it does not apply. Of course, seen from the point of view of singular integrals, the map $f \mapsto Au$ is naturally not bounded on $L^1(I \times \Omega)$ and one should replace $L^1(I \times \Omega)$ by a Hardy space and prove boundedness into $L^1(I \times \Omega)$. (The notion of H^p -maximal regularity introduced in [4] is a different problem.) When A is the Laplacian on \mathbb{R}^n , kernel representation suggests to use a parabolic Hardy space of Coifman-Weiss type on $I \times \mathbb{R}^n$. Indeed, this is done implicitly in [20, Appendix] given $H^1 - BMO$ duality, where the authors prove boundedness from $L^\infty(\mathbb{R} \times \mathbb{R}^n)$ to parabolic BMO of a dual problem. They also do this for some time dependent generalization of (1.1) with the Laplacian on \mathbb{R}^n .

In the abstract setting of [10] and [3], it is not clear which Hardy space is appropriate¹. Here, we restrict ourselves to the Laplace-Beltrami operator Δ on some manifolds M or a sublaplacian on some Lie groups, or second order operators on \mathbb{R}^n with bounded measurable coefficients with additional assumptions. We prove that the operator of maximal regularity is not just bounded from a Hardy space into L^1 but into itself. We prove a similar phenomenon for the dual operator but with two different Hardy spaces as source and target. The Hardy spaces on which we prove boundedness are of course of parabolic type. Interpolation applies to reobtain maximal L^q regularity on L^p for $1 < p, q < \infty$.

The plan of the article is as follows. We give next the notation and then state the main result in the setting of a Riemannian manifold. We define and study the parabolic Hardy spaces in Section 3. Then the proof of the main result on

¹This will be done in a subsequent paper by the last two authors.

manifolds is presented in Section 4 and 5. The consequence for the maximal L^q regularity is in Section 6. We give in Section 7 the needed ingredients to prove the analogous results in a connected Lie group with polynomial growth. We also present the analogous results for a class of second order operators on domains of \mathbb{R}^n in Section 8.

2 Main result on manifolds

Let M be a complete non-compact connected Riemannian manifold, μ the Riemannian measure. Denote by $\|\cdot\|_p$ the norm in $L^p(M, \mu)$, $1 \leq p \leq +\infty$. Denote by ∇ the Riemannian gradient, $|\cdot|$ the length in the tangent space, Δ the (negative) Laplace-Beltrami operator corresponding to ∇ and d the geodesic distance. For all $x \in M$, all $r > 0$, $B(x, r)$ stands for the open geodesic ball with center x and radius r , and its measure is denoted $V(x, r)$.

Say that M satisfies the doubling property if and only if there exists $C > 0$ such that for all $r > 0$ and $x \in M$,

$$V(x, 2r) \leq CV(x, r). \quad (D)$$

A straightforward consequence of (D) is that there exist $C, \delta > 0$ such that for all $x \in M$, all $r > 0$ and all $\theta > 1$,

$$V(x, \theta r) \leq C\theta^\delta V(x, r). \quad (D')$$

The hypothesis (D) exactly means that M , equipped with its geodesic distance and its Riemannian measure, is a space of homogeneous type in the sense of Coifman and Weiss ([7]).

Say that M satisfies the (scaled) Poincaré inequalities if there exists $C > 0$ such that for every ball $B = B(x, r)$ and every f with $f, \nabla f$ locally square integrable,

$$\int_B |f - f_B|^2 d\mu \leq Cr^2 \int_B |\nabla f|^2 d\mu, \quad (P)$$

where f_B denotes the mean of f on B .

The assumptions (D) and (P) guarantee the needed properties on the heat kernel (size and regularity) in the sequel.

Consider the Cauchy problem

$$\begin{cases} \frac{\partial u}{\partial t}(t, x) - \Delta u(t, x) = f(t, x), & t > 0, \\ u(0, x) = 0. \end{cases} \quad (2.1)$$

Now define

$$Tf(t, x) = \int_0^t [\Delta e^{(t-s)\Delta} f(s, \cdot)](x) ds. \quad (2.2)$$

By [11], we know that T is bounded from $L^2(X)$ to $L^2(X)$, where $X = \mathbb{R}_+^* \times M$ and product measure $dt d\mu$. Its adjoint is given by

$$T^* f(t, x) = \int_t^\infty [\Delta e^{(s-t)\Delta} f(s, \cdot)](x) ds. \quad (2.3)$$

We introduce in the next section two Hardy spaces of parabolic type on X with $H_z^1(X) \subset H_r^1(X) \subset L^1(X)$. Our main result is the following.

Theorem 2.1. *Assume that M satisfies (D) and (P). Then T is $H_z^1(X)$ -bounded and T^* is bounded from $H_r^1(X)$ to $H_z^1(X)$.*

Corollary 2.2. *With the hypotheses above, one has*

$$\|u'\|_{H_z^1(X)} + \|\Delta u\|_{H_z^1(X)} \leq C \|f\|_{H_z^1(X)}.$$

Proof: If $f \in H_z^1(X) \cap L^2(X)$, we have $\Delta u = Tf$ and $u' = f + \Delta u$ by uniqueness of the Cauchy problem. Thus this estimate holds for those f 's and we conclude by density in $H_z^1(X)$ (See Section 3). \square

We shall also show that a better result cannot hold even in the simplest situation $X = \mathbb{R}_+^* \times \mathbb{R}^n$: T is not bounded from $H_r^1(X)$ to $L^1(X)$ and T^* is not bounded on $L^1(X)$.

3 Parabolic $H_r^1(X)$ and $H_z^1(X)$

In this section we study the parabolic Hardy spaces $H_r^1(X)$ and $H_z^1(X)$, following the ideas of [6] and [2] in an elliptic context.

We assume that (E, d, μ) is a space of homogeneous type. Let N be the product space $\mathbb{R} \times E$ endowed with the measure ν , product of the Lebesgue measure with μ , and with the quasi-distance \tilde{d} defined by

$$\tilde{d}(t, x; s, y)^2 = \sup(d(x, y)^2, |t - s|), \quad (t, x), (s, y) \in N.$$

Suppose $Q = \{(t, x) \in N : \tilde{d}(t, x; s, y) < r\}$ is the ball in N centered at (s, y) with radius r . Then the relation between the volumes of Q and $B(y, r)$ is $\nu(Q) \sim r^2 V(y, r)$, where $V(y, r) = \mu(B(y, r))$. One then easily checks that (N, \tilde{d}, ν) is a space of homogeneous type and

$$\nu(\theta Q) \leq C \theta^{\delta+2} \nu(Q) \quad (3.1)$$

for all balls Q and all $\theta > 1$ where δ is the doubling exponent for E as in (D'). Thus, one can consider the Hardy space $H^1(N) = H_{CW}^1(N)$ of [7]. Let us review some definitions and facts.

Definition 3.1. A classical $(1, \infty)$ -atom on N is a measurable function a supported on a ball $Q \subseteq N$ that satisfies

- (1) $\int_N a d\nu = 0$,
- (2) $\|a\|_{L^\infty(N)} \leq \nu(Q)^{-1}$.

Definition 3.2. An $(1, 2)$ -atom on N is a measurable function a supported on a ball $Q \subseteq N$ that satisfies

- (1) $\int_N a d\nu = 0$,
- (2) $\|a\|_{L^2(N)} \leq \nu(Q)^{-\frac{1}{2}}$.

An $L^1(N)$ function f belongs to $H^1(N)$ if it can be written as $\sum \lambda_Q a_Q$ where $\sum |\lambda_Q| < \infty$ and a_Q are $(1, 2)$ -atoms. The norm is given by the infimum of such $\sum |\lambda_Q|$ taken over all possible decompositions called atomic decompositions. One obtains the same space with an equivalent norm by replacing $(1, 2)$ -atoms by $(1, \infty)$ -atoms.

We now introduce two Hardy spaces on $X = \mathbb{R}_+^* \times E$.

Definition 3.3. Set $H_r^1(X) = \{f \in L^1(X) : \exists F \in H^1(N), F|_X = f\}$. The norm of an element in this space is the quotient norm : the infimum of the $H^1(N)$ norms of all possible extensions to N .

Definition 3.4. Set $H_z^1(X) = \{f|_X : f \in H^1(N), f = 0 \text{ on } \mathbb{R}_- \times E\}$. The norm of an element in this space is the norm of its zero extension in $H^1(N)$.

It is clear that $H_z^1(X) \subset H_r^1(X)$, but the converse is not true as shown by the following example.

Example 1. Let

$$f(t, x) = \begin{cases} 1, & (t, x) \in (0, 1) \times B(0, 1), \\ 0, & \text{elsewhere,} \end{cases}$$

then f is the restriction to X of

$$F(t, x) = \begin{cases} 1, & (t, x) \in (0, 1) \times B(0, 1), \\ -1, & (t, x) \in (-1, 0) \times B(0, 1), \\ 0, & \text{elsewhere.} \end{cases}$$

It is clear that F is a $(1, \infty)$ -atom on N , so $f \in H_r^1(X)$, but $f \notin H_z^1(X)$ because f does not have vanishing moment.

Definition 3.5. Let a be a measurable function supported in a ball $Q \subset X$, and satisfying $\|a\|_{L^2(N)} \leq \nu(Q)^{-\frac{1}{2}}$.

- (1) a is a type (a) atom if, moreover, $4Q \subseteq X$ and $\int_X a d\nu = 0$.
- (2) a is a type (b) atom if, moreover, $2Q \subseteq X$ and $4Q \not\subseteq X$ (a does not necessarily have vanishing moment).

We have an atomic characterization of $H_r^1(X)$. The proof, analogous to the one in [6], is included for completeness.

Theorem 3.6. *Let f be an integrable function on X . The followings are equivalent:*

- (1) $f \in H_r^1(X)$;
- (2) f has an atomic decomposition

$$f = \sum_{\text{type (a) atoms}} \lambda_Q a_Q + \sum_{\text{type (b) atoms}} \mu_Q b_Q \quad (3.2)$$

with

$$\sum_{\text{type (a) atoms}} |\lambda_Q| + \sum_{\text{type (b) atoms}} |\mu_Q| < \infty. \quad (3.3)$$

Further, the norm on $H_r^1(X)$ is equivalent to the infimum of (3.3) taken over all decompositions (3.2).

Proof: (2) \Rightarrow (1). Let f satisfy (3.2) and (3.3). As a_Q is a type (a) atom, then A_Q defined by

$$A_Q(t, x) = \begin{cases} a_Q(t, x), & (t, x) \in X, \\ 0, & (t, x) \notin X, \end{cases}$$

is a (1,2)-atom of $H^1(N)$. By definition we conclude that $a_Q \in H_r^1(X)$.

As b_Q is a type (b) atom, then define

$$B_{\tilde{Q}}(t, x) = \begin{cases} b_Q(t, x), & (t, x) \in Q, \\ -b_Q(-t, x), & (t, x) \in Q_-, \end{cases}$$

where Q_- is the reflection of Q across $\{0\} \times M$ given by $(t, x) \mapsto (-t, x)$, and \tilde{Q} is a ball such that $Q \cup Q_- \subseteq \tilde{Q}$ and $\nu(\tilde{Q}) \leq C\nu(Q)$. For instance, Q is centered at (s, y) with radius r , the ball \tilde{Q} centered at $(0, y)$ with radius $5r$ works.

Then we have

$$\begin{aligned} \int_N B_{\tilde{Q}} d\nu &= \iint_X B_{\tilde{Q}}(t, x) dt d\mu(x) + \iint_{\mathbb{R}_- \times M} B_{\tilde{Q}}(t, x) dt d\mu(x) \\ &= \iint_Q b_Q(t, x) dt d\mu(x) - \iint_Q b_Q(t, x) dt d\mu(x) = 0, \end{aligned}$$

and

$$\|B_{\tilde{Q}}\|_{L^2(N)} \leq \|b_Q\|_{L^2(N)} + \|-b_Q\|_{L^2(N)} \leq \nu(Q)^{-\frac{1}{2}} + \nu(Q_-)^{-\frac{1}{2}} \leq c\nu(\tilde{Q})^{-\frac{1}{2}}.$$

This means that $c^{-1}B_{\tilde{Q}}$ is an (1,2)-atom for $H^1(N)$. Now let

$$F = \sum_{\text{type (a) atoms}} \lambda_Q A_Q + \sum_{\text{type (b) atoms}} \mu_Q B_{\tilde{Q}},$$

then $F \in H^1(N)$ and $F|_X = f$, therefore $f \in H_r^1(X)$.

(1) \Rightarrow (2). If $f \in H_r^1(X)$, then there exists $F \in H^1(N)$ such that $F|_X = f$. So F has an atomic decomposition $\sum \lambda_Q A_Q$ [7]. It suffices therefore to concentrate on the restriction of one single atom A_Q . Restricting A_Q to X , then we consider only those balls Q which intersect X and we have the following situations:

If $4Q \subseteq X$, then $a_Q = A_Q$ is a type (a)-atom.

If $4Q \not\subseteq X$, but $2Q \subseteq X$, then $b_Q = A_Q$ is a type (b)-atom.

If $2Q \not\subseteq X$, then we can decompose $A_Q|_X$ into type (b)-atoms as follows. Using Whitney decomposition on X with respect to $\mathbb{R}_- \times E$ (see [12]), one can find a family of balls $Q_j \subset X$ with $Q \cap X \subseteq \cup Q_j$ and bounded overlap. Hence,

$\sum \nu(Q_j) \leq c\nu(Q \cap X)$ for some constant $c > 0$ depending on the metric and the doubling property of ν . Write

$$A = \sum_j \chi_{Q_j} A = \sum_j \frac{\|\chi_{Q_j} A\|_2}{\nu(Q_j)^{-\frac{1}{2}}} \cdot \frac{\nu(Q_j)^{-\frac{1}{2}}}{\|\chi_{Q_j} A\|_2} \chi_{Q_j} A = \sum \lambda_{Q_j} a_{Q_j},$$

where the indices j are those for which $\chi_{Q_j} A$ are not identically 0 and

$$\lambda_{Q_j} = \frac{\|\chi_{Q_j} A\|_2}{\nu(Q_j)^{-\frac{1}{2}}}, \quad a_{Q_j} = \frac{\nu(Q_j)^{-\frac{1}{2}}}{\|\chi_{Q_j} A\|_2} \chi_{Q_j} A.$$

Then we have $\text{supp } a_{Q_j} \subseteq Q_j$ and

$$\|a_{Q_j}\|_{L^2(N)} \leq \frac{\nu(Q_j)^{-\frac{1}{2}}}{\|\chi_{Q_j} A\|_2} \cdot \|\chi_{Q_j} A\|_2 \leq \nu(Q_j)^{-\frac{1}{2}}.$$

It follows that a_{Q_j} is a type (b)-atom and

$$\sum |\lambda_{Q_j}| = \sum_j \frac{\|\chi_{Q_j} A\|_2}{\nu(Q_j)^{-\frac{1}{2}}} \leq \left(\sum_j \|\chi_{Q_j} A\|_2^2 \right)^{\frac{1}{2}} \cdot \left(\sum_j \nu(Q_j) \right)^{\frac{1}{2}} \lesssim \|A\|_{L^2(N)} \nu(Q \cap X)^{\frac{1}{2}}.$$

As $\|A\|_{L^2(N)} \leq \nu(Q)^{-\frac{1}{2}}$, the series is bounded independently of A and this atomic decomposition converges boundedly in $H_r^1(X)$. \square

For the space $H_z^1(X)$, we have the following characterization, whose proof is again inspired by [6].

Theorem 3.7. *Let f be an integrable function on X . The followings are equivalent:*

- (1) $f \in H_z^1(X)$;
- (2) f has an atomic decomposition $f = \sum_{Q \subset X} \lambda_Q a_Q$ with $\sum |\lambda_Q| < +\infty$, where a_Q are (1,2)-atoms with support entirely contained in X .

Proof: (2) \Rightarrow (1) is obvious and we turn to (1) \Rightarrow (2). Let $f \in H_z^1(X)$. Then the function F defined by

$$F(t, x) = \begin{cases} f(t, x), & (t, x) \in X, \\ 0, & (t, x) \notin X. \end{cases}$$

belongs to $H^1(N)$. Set $f_e(t, x) = F(t, x) + F(-t, x)$ with $(t, x) \in N$ then f_e is even in the t variable, $f_e \in H^1(N)$, $f_e|_X = f$. Pick an atomic decomposition $f_e = \sum_Q \lambda_Q A_Q$, where A_Q are (1,2)-atoms for $H^1(N)$. One can rewrite f_e as follows:

$$f_e(t, x) = \frac{f_e(t, x) + f_e(-t, x)}{2} = \sum_Q \lambda_Q \frac{A_Q(t, x) + A_Q(-t, x)}{2} = \sum_Q \lambda_Q \tilde{A}_Q(t, x),$$

where $\tilde{A}_Q(t, x) = \frac{A_Q(t, x) + A_Q(-t, x)}{2}$.

1° If $\text{supp } A_Q \subseteq X$, then $\tilde{A}_Q|_X = \frac{A_Q}{2}$, and we define

$$a_Q(t, x) = \frac{A_Q(t, x)}{2}, \quad (t, x) \in X.$$

2° If $\text{supp } A_Q \subseteq \mathbb{R}_- \times M$, $\tilde{A}_Q|_X(t, x) = \frac{A_Q(-t, x)}{2}$, and we define

$$a_Q(t, x) = \frac{A_Q(-t, x)}{2}, \quad (t, x) \in X.$$

3° If neither case occurs, then let $Q_0 = (\text{supp } A_Q) \setminus X$, and the "reflection" of Q_0 is defined by $Q_r = \{(t, x) \in X : (-t, x) \in Q_0\}$. Define $a_Q = \tilde{A}_Q|_X$. Clearly, if Q is centered at (s, y) has radius r , we have $|s| \leq r^2$ and $\text{supp } a_Q \subset Q|_X \cup Q_r$. Hence $\text{supp } a_Q$ is contained in the ball \tilde{Q} centered at (r^2, y) with radius r which is contained in X . Next,

$$\|a_Q\|_{L^2(N)} \leq \nu(Q)^{-\frac{1}{2}} = \nu(\tilde{Q})^{-\frac{1}{2}}.$$

Finally

$$\begin{aligned} \int_X a_Q d\nu &= \frac{1}{2} \iint_{Q \cap X} A_Q(t, x) dt d\mu(x) + \frac{1}{2} \iint_{Q_r} A_Q(-t, x) dt d\mu(x) \\ &= \frac{1}{2} \iint_{Q \cap X} A_Q(t, x) dt d\mu(x) + \frac{1}{2} \iint_{Q \setminus X} A_Q(t, x) dt d\mu(x) \\ &= \iint_Q A_Q(t, x) dt d\mu(x) = 0. \end{aligned}$$

Thus a_Q are $(1, 2)$ -atoms with $\text{supp } a_Q \subset X$. So

$$f = f_e|_X = \sum_Q \lambda_Q \tilde{A}_Q|_X = \sum_Q \lambda_Q a_Q.$$

This completes the proof of the theorem. \square

Let us make some further remarks. If one uses the reflection $(t, x) \mapsto (-t, x)$ in N , then we say that a function on N is odd (resp. even) if it is changed to its opposite (resp. itself).

Proposition 3.8. $H_r^1(X)$ coincides with the space of restrictions to X of odd functions in $H^1(N)$. Similarly, $H_z^1(X)$ coincides with the space of restrictions to X of even functions in $H^1(N)$.

Proof: The argument is essentially contained in the proofs of the last two results and is skipped (See also [6]). \square

Remark 3.9. Density of $H_z^1(X) \cap L^2(X)$ (resp. $H_r^1(X) \cap L^2(X)$) in $H_z^1(X)$ (resp. $H_r^1(X)$) follows from the atomic decomposition.

Since (X, \tilde{d}, ν) is also a space of homogeneous type, let $H^1(X) = H_{CW}^1(X)$ be the Hardy space of Coifman and Weiss as defined in [7]. Each $f \in H^1(X)$ can be written as follows:

$$f = \sum_Q \lambda_Q a_Q,$$

where $\text{supp } a_Q \subset Q \cap X$, Q is a ball centered in X , $\int_{Q \cap X} a d\nu = 0$, and $\|a_Q\|_{L^2(N)} \leq \nu(Q \cap X)^{-\frac{1}{2}}$ and where $\sum |\lambda_Q| < \infty$. Observe that since Q is centered in X , one has $\nu(Q) \lesssim \nu(Q \cap X)$ with implicit constant independent of Q .

Proposition 3.10. $H^1(X) = H_z^1(X)$.

Proof: That $H_z^1(X) \subseteq H^1(X)$ is obvious. To prove the converse inclusion let a be an atom in $H^1(X)$: there exists a ball Q in N centered at a point in X such that $\text{supp } a \subseteq Q \cap X$. If $Q \subset X$, then $a \in H_z^1(X)$. Otherwise, let (s, y) be the center of Q and r its radius and observe that $0 < s$. Let \tilde{Q} be the ball centered at (r^2, y) with radius r . Then

$$\nu(\tilde{Q}) \sim r^2 V(y, r) = \frac{r^2}{s + r^2} (s + r^2) V(y, r) \leq 2\nu(Q \cap X).$$

Clearly $\text{supp } a \subseteq \tilde{Q}$ and a is up to a fixed multiplicative constant, an $(1, 2)$ -atom supported inside X . This proves that $a \in H_z^1(X)$ and completes the proof of the proposition. \square

Remark 3.11. One obtains all the above atomic decompositions replacing atoms with L^2 estimates by atoms with L^∞ estimates. This will be used in the sequel.

We now recall the notion of molecules.

Definition 3.12. Let Q be a ball contained in X .

A function m is called a molecule associated to Q if there exists $\alpha > 0$ such that

- (1) $\int_X m d\nu = 0$,
 - (2) for all $j \geq 1$, $\|m\|_{L^2(B_j(Q))} \leq 2^{-j\alpha} \nu(2^{j+1}Q \cap X)^{-\frac{1}{2}}$,
- where

$$B_j(Q) = \begin{cases} 4Q \cap X, & j = 1, \\ (2^{j+1}Q \setminus 2^jQ) \cap X, & j \geq 2. \end{cases}$$

A function m is called a type (a) molecule associated to Q if $4Q \subseteq X$ and there exists $\alpha > 0$ such that (1) and (2) hold.

A function m is called a type (b) molecule associated to Q if $2Q \subseteq X$, $4Q \not\subseteq X$ and there exists $\alpha > 0$ such that (2) holds.

Proposition 3.13. *Molecules form a bounded set in $H_z^1(X)$. Molecules of type (a) and (b) form a bounded set in $H_r^1(X)$.*

Proofs are somewhat analogous to the ones in [7] but we include them for convenience.

Proof: Let us begin with the case of $H_z^1(X)$. Let m be a molecule associated to a cube Q . By Proposition 3.10, it suffices to show that m belongs to $H^1(X)$ with norm independent of Q . Set χ_A the indicator of a set A and $m_{B_j(Q)}$ the mean value of m over $B_j(Q)$: $\frac{1}{\nu(B_j(Q))} \int_{B_j(Q)} m d\nu$. Write

$$m = \sum_{j=1}^{\infty} m \chi_{B_j(Q)} = \sum_{j=1}^{\infty} (m - m_{B_j(Q)}) \chi_{B_j(Q)} + \sum_{j=1}^{\infty} m_{B_j(Q)} \chi_{B_j(Q)}.$$

Using (2) in Definition 3.12, we obtain that

$$\int_{2^{j+1}Q \cap X} |(m - m_{B_j(Q)}) \chi_{B_j(Q)}|^2 d\nu \leq 4 \int_{B_j(Q)} |m|^2 d\nu \lesssim \frac{1}{4^{j\alpha}} \cdot \frac{1}{\nu(2^{j+1}Q \cap X)},$$

where we have used the doubling property $\nu(2^{j+1}Q) \sim \nu(2^jQ)$ and also the fact that $\nu(Q \cap X) \sim \nu(Q)$ for all balls centered at a point in X .

Set

$$\lambda_j = \frac{1}{2^{j\alpha}}, \quad a_j = \lambda_j^{-1} (m - m_{B_j(Q)}) \chi_{B_j(Q)}.$$

We have $\int_{2^{j+1}Q \cap X} a_j d\nu = 0$, $\text{supp } a_j \subset 2^{j+1}Q \cap X$, and $\|a_j\|_{L^2(2^{j+1}Q \cap X)} \leq C \nu(2^{j+1}Q \cap X)^{-\frac{1}{2}}$, with C independent of j and Q . So $a_j \in H^1(X)$ and $\|a_j\|_{H^1(X)}$ is bounded. Since $\sum_{j=1}^{\infty} \lambda_j \leq C(\alpha) < \infty$,

$$\sum_{j=1}^{\infty} (m - m_{B_j(Q)}) \chi_{B_j(Q)} = \sum_{j=1}^{\infty} \lambda_j a_j \in H^1(X).$$

For the second term, we have

$$\begin{aligned} \sum_{j=1}^{\infty} m_{B_j(Q)} \chi_{B_j(Q)} &= \sum_{j=1}^{\infty} \int_{B_j(Q)} m d\nu \frac{\chi_{B_j(Q)}}{\nu(B_j(Q))} \\ &= \sum_{j=2}^{\infty} \left(\int_{2^{j+1}Q \cap X} m d\nu - \int_{2^jQ \cap X} m d\nu \right) \frac{\chi_{B_j(Q)}}{\nu(B_j(Q))} + \left(\int_{4Q \cap X} m d\nu \right) \frac{\chi_{B_1(Q)}}{\nu(B_1(Q))} \\ &= \sum_{j=2}^{\infty} \left(\int_{2^jQ \cap X} m d\nu \right) \left[\frac{\chi_{B_{j-1}(Q)}}{\nu(B_{j-1}(Q))} - \frac{\chi_{B_j(Q)}}{\nu(B_j(Q))} \right]. \end{aligned}$$

Set

$$\mu_j = \int_{2^jQ \cap X} m d\nu, \quad b_j = \frac{\chi_{B_{j-1}(Q)}}{\nu(B_{j-1}(Q))} - \frac{\chi_{B_j(Q)}}{\nu(B_j(Q))}.$$

Then by (1) in Definition 3.12,

$$\mu_j = \int_{2^jQ \cap X} m d\nu = - \left(\int_X m d\nu - \int_{2^jQ \cap X} m d\nu \right) = - \sum_{i \geq j} \int_{B_i(Q)} m d\nu,$$

so

$$|\mu_j| \leq \sum_{i \geq j} \int_{B_i(Q)} |m| d\nu \lesssim \sum_{i \geq j} 2^{-i\alpha} \sim 2^{-j\alpha},$$

thus

$$\sum_{j=1}^{\infty} |\mu_j| \lesssim \sum_{j=1}^{\infty} 2^{-j\alpha} < +\infty.$$

On the other hand, $\text{supp } b_j \subseteq 2^{j+1}Q \cap X$,

$$\int_{2^{j+1}Q \cap X} b_j d\nu = 0,$$

and

$$\int_{2^{j+1}Q \cap X} |b_j|^2 d\nu \leq \frac{C^2}{\nu(2^{j+1}Q \cap X)},$$

for some $C > 0$ independent of j and Q . To prove the last inequality, we remark that $\nu(B_j(Q))$ is comparable to $\nu(2^{j+1}Q \cap X)$ for all $j \geq 1$. This is obvious for $j = 1$ and we turn to $j \geq 2$. As $B_j(Q)$ is contained in $2^{j+1}Q \cap X$, we obtain one inequality. Next, write $Q = I(t_0, r^2) \times B(x_0, r) = I \times B$ and remark that $B_j(Q)$ contains the set $[t_0 + (2^j r)^2, t_0 + (2^{j+1} r)^2] \times (2^{j+1} B)$. Hence

$$\nu(B_j(Q)) \geq (2^j r)^2 \mu(2^{j+1} B) \sim \nu(2^{j+1} Q) \sim \nu(2^{j+1} Q \cap X)$$

where the last equivalence comes from the fact that $2^{j+1}Q$ is centered in X .

Thus $C^{-1}b_j$ are atoms for $H^1(X)$, and therefore $\sum_{j=1}^{\infty} \mu_j b_j \in H^1(X)$.

Clearly the same argument applies for type (a) molecules so they belong to $H_z^1(X) \subset H_r^1(X)$. It remains to consider a type (b) molecule m associated to a cube Q . According to Proposition 3.8, the odd extension m_{odd} is, up to a constant that depends only on N , a molecule for $H^1(N)$ associated to a cube \tilde{Q} of size comparable to Q (since $4Q \not\subseteq X$), containing Q and its reflection across $\{0\} \times E$. The same argument as above shows that m_{odd} belongs to $H^1(N)$. Thus $m \in H_r^1(X)$ and the control of its norm follows from an examination of the argument. \square

We finish with a remark on how to prove boundedness on Hardy spaces for an operator, bypassing any other knowledge such as weak type (1,1).

Proposition 3.14. *Let H_1, H_2 be any Hardy spaces obtained via an atomic decomposition. Any linear operator mapping the atoms of H_1 into a bounded set in H_2 has a bounded extension from H_1 into H_2 .*

Proof: Let S be the subset of H_1 consisting of finite linear combinations $f = \sum_{i=1}^n \lambda_i a_i$, $n \geq 1$, $\lambda_i \in \mathbb{C}$, a_i atoms of H_1 with $\|f\|_{H_1} \geq \frac{1}{10} \sum_{i=1}^n |\lambda_i|$. We claim that S is dense in H_1 . Indeed, if $f \in H_1$, there is a (possibly) infinite representation $f = \sum_{i=1}^{\infty} \lambda_i a_i$ with $\sum_{i=1}^{\infty} |\lambda_i| \leq 2\|f\|_{H_1}$. As the partial sums f_n converge to f in H_1 , we have $\|f\|_{H_1} \leq 5\|f_n\|_{H_1}$ for n large enough and

$$\sum_{i=1}^n |\lambda_i| \leq 2\|f\|_{H_1} \leq 10\|f_n\|_{H_1}.$$

This proves that $f_n \in S$ for all n large enough.

By assumption, we have a linear operator T satisfying $\|Ta\|_{H_2} \leq C$ for all atoms $a \in H_1$. By linearity and definition of S , we have $\|Tf\|_{H_2} \leq 10C\|f\|_{H_1}$ for all $f \in S$ and we conclude by density. \square

4 Boundedness of T

Let us come back to the situation where M is a Riemannian manifold and $X = \mathbb{R}_+^* \times M$. Before proving the regularity of the operator T , that is its boundedness on $H_z^1(X)$, let us first give some lemmas about the heat kernel $p_t(x, y)$. Indeed, using the definition of T , we can write

$$Tf(t, x) = \int_0^t \Delta e^{(t-s)\Delta} f(s, x) ds = \iint_X p'_{t-s}(x, y) \chi_{\{t>s\}} f(s, y) ds d\mu(y),$$

where $p'_t(x, y)$ stands for $\frac{\partial}{\partial t} p_t(x, y)$. The last equality is a formal one as the integral may not converge.

Lemma 4.1. ([19], Proposition 3.3) *If M satisfies the doubling property (D) and the Poincaré inequality (P), then the heat kernel satisfies the Gaussian upper estimate:*

$$p_t(x, y) \leq \frac{C_1}{V(x, \sqrt{t})} \cdot e^{-c_1 \frac{d(x, y)^2}{t}} \quad (G)$$

for all $t > 0$ and $x, y \in M$.

There is also a lower estimate but we do not need it.

Lemma 4.2. ([10], Lemma 2.5) *If the heat kernel satisfies (G), then the time derivatives of p_t satisfy:*

$$\left| \frac{\partial^k}{\partial t^k} p_t(x, y) \right| \lesssim \frac{1}{t^k V(x, \sqrt{t})} e^{-\frac{Cd(x, y)^2}{t}}, \quad (4.1)$$

for all $t > 0$, $x, y \in M$ and all integer $k \geq 1$.

Lemma 4.3. *If (D) and (P) hold, for all $t > 0$, $x, y, x_0 \in M$, if $d(y, x_0) \leq \sqrt{t}$, then*

$$\left| \frac{\partial}{\partial t} p_t(x, y) - \frac{\partial}{\partial t} p_t(x, x_0) \right| \lesssim \frac{1}{tV(x, \sqrt{t})} \left(\frac{d(y, x_0)}{\sqrt{t}} \right)^\gamma e^{-\frac{Cd(x, y)^2}{t}},$$

where $\gamma \in (0, 1)$, $C > 0$ are independent of x, y, x_0, t .

Proof: Following [13] (see also [18]), one can write

$$\frac{\partial}{\partial t} (p_t(x, y) - p_t(x, x_0)) = \int_M \frac{\partial}{\partial t} p_{t/2}(x, z) (p_{t/2}(z, y) - p_{t/2}(z, x_0)) dz$$

and use the estimate (4.1) together with the Hölder estimate of $p_t(x, y)$, as $p_t(x, y)$ is a solution of heat equation, from Harnack inequality (see [19], Proposition 3.2). This is where we use (indirectly) the assumption on Poincaré's inequality. \square

Lemma 4.4. Assume (D) and (G). One has for all $y \in M$ and $s > 0$.

$$\int_M |\nabla_x p_s(x, y)| d\mu(x) \lesssim \frac{1}{\sqrt{s}}.$$

Proof: It is proved in [9, Lemma 2.3] that there is $\gamma > 0$ such that

$$\int_M |\nabla_x p_s(x, y)|^2 e^{\gamma \frac{d^2(x, y)}{s}} d\mu(x) \lesssim \frac{1}{sV(y, \sqrt{s})},$$

for all $y \in M$ and $s > 0$. It suffices to use Cauchy-Schwarz inequality. \square

Now we can prove the first part of Theorem 2.1, namely that T is bounded from $H_z^1(X)$ to itself. By Proposition 3.10, we prove that T is bounded from $H_z^1(X)$ to $H^1(X)$. To this end, it is enough by Proposition 3.14 to check that T maps atoms to molecules defined as above, up to a multiplicative factor.

Suppose $a \in H_z^1(X)$ is an $(1, \infty)$ -atom : there is a ball $Q = \{(t, x) \in X : \tilde{d}(t, x; t_0, x_0) < r\}$ contained in X such that $\text{supp } a \subset Q$, $\|a\|_\infty \leq \nu(Q)^{-1}$ and $\int_Q a d\nu = 0$. Since T is bounded from $L^2(X)$ to $L^2(X)$,

$$\left(\int_{4Q \cap X} |Ta|^2 d\nu \right)^{\frac{1}{2}} \leq C \left(\int_X |a|^2 d\nu \right)^{\frac{1}{2}} \leq C\nu(Q)^{-\frac{1}{2}} \leq c\nu(4Q \cap X)^{-\frac{1}{2}}.$$

Thus $c^{-1}Ta$ satisfies (2) with $j = 1$ in Definition 3.12.

Now to prove that Ta satisfies (2) with $j \geq 2$ in Definition 3.12, up to some multiplicative constant, it suffices to prove that there exists $\alpha > 0$ such that for every $j \geq 2$ and $(t, x) \in B_j(Q)$,

$$|Ta(t, x)| \lesssim 2^{-j\alpha} \nu(2^{j+1}Q \cap X)^{-1}.$$

Observe that $Ta(t, x) = 0$ if $t \leq t_0 - r^2$. Hence we assume $t > t_0 - r^2$ in the sequel and we have

$$Ta(t, x) = \int_{t_0 - r^2}^t \int_{d(y, x_0) < r} p'_{t-s}(x, y) a(s, y) ds d\mu(y).$$

If $(t, x) \in B_j(Q)$, then $2^j r \leq \tilde{d}(t, x; t_0, x_0) \leq 2^{j+1} r$, and we have the following two cases:

Case I: $2^j r \leq d(x, x_0) \leq 2^{j+1} r, |t - t_0| \leq (2^{j+1} r)^2$.

Case II: $d(x, x_0) \leq 2^{j+1} r, (2^j r)^2 \leq |t - t_0| \leq (2^{j+1} r)^2$.

Case I: we distinguish two subcases.

1°: If $2^j r \leq d(x, x_0) \leq 2^{j+1} r$ and $|t - t_0| \leq (2r)^2$, then we have that $2^{j-2} r \leq d(x, y) \leq 2^{j+2} r$, whenever $(s, y) \in Q$. As a consequence,

$$\begin{aligned}
|Ta(t, x)| &\leq \int_{t_0-r^2}^t \int_{d(y, x_0) < r} \frac{1}{(t-s)V(x, \sqrt{t-s})} e^{-\frac{Cd(x, y)^2}{t-s}} |a(s, y)| ds d\mu(y) \\
&\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_0^{t-t_0+r^2} \frac{1}{sV(x, \sqrt{s})} e^{-\frac{C(2^j r)^2}{s}} ds \\
&\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_{\frac{4^j}{5}}^{+\infty} \frac{1}{uV(x, \frac{2^j r}{\sqrt{u}})} e^{-Cu} du \\
&\lesssim \frac{2^{j(\delta+2)}}{\nu(2^{j+1}Q)} \int_{\frac{4^j}{5}}^{+\infty} \frac{V(x_0, r)}{uV(x, \frac{2^j r}{\sqrt{u}})} e^{-Cu} du \\
&\lesssim \frac{2^{j(\delta+2)}}{\nu(2^{j+1}Q)} \int_{\frac{4^j}{5}}^{+\infty} \frac{1}{u} \left(\frac{\sqrt{u}}{2^j} \right)^\delta e^{-Cu} du \\
&\lesssim \frac{1}{\nu(2^{j+1}Q)} e^{-C'4^j} \lesssim \frac{1}{\nu(2^{j+1}Q \cap X)} e^{-C''4^j}.
\end{aligned} \tag{4.2}$$

The first inequality is obtained by Lemma 4.2, the second follows from the size condition of a , the third is by changing variable: $u = \frac{(2^j r)^2}{s}$, the fourth is by using doubling property (3.1) for ν and the fifth is by the doubling property of μ :

$$V(x_0, r) \lesssim V(x, r) \left(1 + \frac{d(x_0, x)}{r} \right)^\delta \lesssim V\left(x, \frac{2^j r}{\sqrt{5u}}\right) \left(1 + \frac{2^j r}{r} \right)^\delta \left(\frac{\sqrt{u}}{2^j} \right)^\delta.$$

2°: If $2^j r \leq d(x, x_0) \leq 2^{j+1} r$ and $(2r)^2 \leq |t - t_0| \leq (2^{j+1} r)^2$, then $d(x, y) \sim 2^j r$ and $t - s \sim t - t_0$ when $(s, y) \in Q$. In this case, using $\iint a(s, y) ds d\mu(y) = 0$, write $Ta(t, x) = I_1 + I_2$ where

$$I_1 = \int_{t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) < r} (p'_{t-s}(x, y) - p'_{t-s}(x, x_0)) a(s, y) ds d\mu(y),$$

and

$$I_2 = \int_{t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) < r} (p'_{t-s}(x, x_0) - p'_{t-t_0-r^2}(x, x_0)) a(s, y) ds d\mu(y).$$

For I_1 , we have the following estimate:

$$\begin{aligned}
|I_1| &\leq \iint_Q |p'_{t-s}(x, y) - p'_{t-s}(x, x_0)| |a(s, y)| ds d\mu(y) \\
&\lesssim \frac{1}{\nu(Q)} \iint_Q \frac{1}{(t-s)V(x, \sqrt{t-s})} \left(\frac{d(y, x_0)}{\sqrt{t-s}} \right)^\gamma e^{-\frac{Cd(x, y)^2}{t-s}} ds d\mu(y) \\
&\lesssim \left(\frac{r}{\sqrt{t-t_0}} \right)^\gamma \frac{1}{|t-t_0|V(x, \sqrt{|t-t_0|})} e^{-\frac{C(2^j r)^2}{t-t_0}} \\
&\lesssim \frac{1}{2^{j\gamma}\nu(2^{j+1}Q)} \left(\frac{2^j r}{\sqrt{t-t_0}} \right)^\gamma \frac{(2^j r)^2 V(x_0, 2^j r)}{|t-t_0|V(x, \sqrt{|t-t_0|})} e^{-\frac{C(2^j r)^2}{t-t_0}} \\
&\lesssim \frac{1}{2^{j\gamma}\nu(2^{j+1}Q)} \\
&\lesssim \frac{1}{2^{j\gamma}\nu(2^{j+1}Q \cap X)}.
\end{aligned}$$

The second inequality is obtained by Lemma 4.3, and the fifth inequality is by doubling property $V(x_0, 2^j r) \lesssim V(x, 2^{j+1}r) \lesssim (\frac{2^j r}{\sqrt{t-t_0}})^\delta V(x, \sqrt{t-t_0})$, and that the function $u \rightarrow u^\beta e^{-Cu}$ is bounded on \mathbb{R}_+ .

Similarly, for I_2 , we have

$$\begin{aligned}
|I_2| &\leq \left| \int_{s=t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) < r} \int_{w=s}^{t_0+r^2} p''_{t-w}(x, x_0) dw a(s, y) ds d\mu(y) \right| \\
&\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_{s=t_0-r^2}^{t_0+r^2} \int_{w=s}^{t_0+r^2} |p''_{t-w}(x, x_0)| dw ds \\
&\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_{w=t_0-r^2}^{t_0+r^2} \int_{s=t_0-r^2}^w \frac{1}{(t-w)^2 V(x, \sqrt{t-w})} e^{-\frac{Cd(x, x_0)^2}{(t-w)}} dw ds \\
&\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_{t_0-r^2}^{t_0+r^2} (w-t_0+r^2) \frac{1}{(t-w)^2 V(x, \sqrt{t-w})} e^{-C\frac{(2^j r)^2}{t-w}} dw \\
&\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_0^{2r^2} w dw \frac{1}{(t-t_0)^2 V(x, \sqrt{t-t_0})} e^{-C\frac{4^j r^2}{t-t_0}} \\
&\lesssim \frac{r^2 V(x_0, r)}{\nu(Q)} \cdot \frac{1}{(2^j r)^2 V(x_0, 2^j r)} \cdot \frac{r^2 (2^j r)^2 V(x, 2^j r)}{(t-t_0)^2 V(x, \sqrt{t-t_0})} e^{-C\frac{4^j r^2}{t-t_0}} \\
&\lesssim \frac{1}{4^j \nu(2^{j+1}Q)} \lesssim \frac{1}{4^j \nu(2^{j+1}Q \cap X)},
\end{aligned}$$

where the third inequality is obtained by Lemma 4.2, the fifth inequality by substitution and the fact $t-w \sim t-t_0$, the sixth inequality by $V(x, 2^j r) \sim V(x_0, 2^j r)$, the seventh by $r^2 V(x_0, r) \sim \nu(Q)$, and that the function $u \rightarrow u^\beta e^{-Cu}$ is bounded on \mathbb{R}_+ .

Case II: $(2^j r)^2 \leq |t-t_0| \leq (2^{j+1}r)^2$ and $d(x, x_0) \leq 2^{j+1}r$. Then we have $t-s \sim$

$t - t_0$. Write again $Ta = I_1 + I_2$, where I_1 and I_2 are as in Case I. Thus

$$\begin{aligned}
|I_1| &\leq \frac{1}{\nu(Q)} \int_{t_0-r^2}^{t_0+r^2} \int_{d(y,x_0) < r} \left(\frac{r}{\sqrt{t-s}} \right)^\gamma \frac{1}{(t-s)V(x, \sqrt{t-s})} e^{-\frac{Cd(x,y)^2}{t-s}} ds d\mu(y) \\
&\lesssim \frac{1}{\nu(Q)} \int_{t_0-r^2}^{t_0+r^2} \int_{d(y,x_0) < r} \left(\frac{r}{\sqrt{t-s}} \right)^\gamma \frac{1}{(t-s)V(x, \sqrt{t-s})} ds d\mu(y) \\
&\lesssim \frac{1}{2^{j\gamma}} \frac{1}{(2^j r)^2 V(x, 2^j r)} \lesssim \frac{1}{2^{j\gamma} \nu(2^{j+1} Q)} \lesssim \frac{1}{2^{j\gamma} \nu(2^{j+1} Q \cap X)},
\end{aligned}$$

where the first inequality follows from Lemma 4.3, the third inequality follows from the fact that $t - s \sim (2^j r)^2$, the fourth from $V(x, 2^j r) \sim V(x_0, 2^j r)$.

For I_2 , we have

$$\begin{aligned}
|I_2| &\leq \int_{s=t_0-r^2}^{t_0+r^2} \int_{d(y,x_0) < r} \left\{ \int_{\theta=s}^{t_0+r^2} |p''_{t-\theta}(x, x_0)| d\theta \right\} |a(s, y)| ds d\mu(y) \\
&\lesssim \frac{1}{\nu(Q)} \int_{t_0-r^2}^{t_0+r^2} \int_{d(y,x_0) < r} \int_s^{t_0+r^2} \frac{1}{(t-\theta)^2 V(x, \sqrt{t-\theta})} d\theta ds d\mu(y) \\
&\lesssim \frac{r^2}{(2^j r)^4 V(x_0, 2^j r)} \lesssim \frac{1}{4^j \nu(2^{j+1} Q)} \lesssim \frac{1}{4^j \nu(2^{j+1} Q \cap X)},
\end{aligned}$$

where the second inequality is obtained by Lemma 4.2, the third one by $t - \theta \sim (2^j r)^2$ and $V(x, 2^j r) \sim V(x_0, 2^j r)$.

Thus we proved that Ta satisfies (2) with $\alpha = \gamma$ in Definition 3.12.

We claim that $\int_X Ta d\nu = 0$. Assuming this claim for a moment, we conclude from Proposition 3.13 that $Ta \in H^1(X)$ with $\|Ta\|_{H^1(X)} \leq C$ uniformly with respect to a . \square

Lemma 4.5. $\int_X Ta d\nu = 0$ when $a \in H_z^1(X)$ is an $(1, \infty)$ -atom.

Proof: We give two proofs, each having advantages. The first one is as follows. Let u be the solution (since $a \in L^2(X)$) of (2.1) with data a . Then $Ta = \Delta u = u' - a$. Furthermore $a \in L^1(X)$ and $Ta, u' \in L^1(X)$ as we just showed. Fix h a bounded function with compact support on M . Then integrate the equation against h on X . Note that $u \in C([0, \infty), L^2(M))$ with $u(0, \cdot) = 0$ and $\lim_{t \rightarrow \infty} u(t, \cdot) = 0$ in $L^2(M)$ (since a is 0 for t larger than some τ , hence $u(t, \cdot) = e^{(t-\tau)\Delta} u(\tau, \cdot)$ for $t > \tau$) and $u' \in L^2(\mathbb{R}_+^*, L^2(M))$. Integrating by parts, one finds

$$\iint_X u'(t, x) h(x) dt d\mu(x) = \int_0^\infty \langle u'(t, \cdot), h \rangle dt = 0,$$

hence,

$$\iint_X Ta(t, x) h(x) dt d\mu(x) = - \iint_X a(t, x) h(x) dt d\mu(x).$$

Since a and Ta both belong to $L^1(X)$, letting h converge to 1 everywhere with $\|h\|_\infty = 1$ implies

$$\iint_X Ta(t, x) dt d\mu(x) = - \iint_X a(t, x) dt d\mu(x) = 0.$$

The second argument is more precise: by the previous calculations, we know that

$$\iint_X |Ta(t, x)| dt d\mu(x) \lesssim \sum_{j=1}^{\infty} 2^{-j\alpha} < +\infty,$$

thus $Ta \in L^1(X)$ and, therefore, for a.e. $t > 0$, $Ta(t, \cdot) \in L^1(M)$. It is enough to show that for such t 's, $\int_M Ta(t, x) d\mu(x) = 0$.

Now let $\varphi \in C_0^\infty([0, +\infty))$, with $\varphi = 1$ on $[0, 1]$, $\varphi = 0$ on $[2, +\infty)$ and $0 \leq \varphi \leq 1$. Fix $x_0 \in M$ and set $\varphi_R(x) = \varphi(R^{-1}\eta(x))$ for $R > 0$ where $\eta \in C^\infty(M \setminus \{x_0\})$ with $\eta(x) \sim d(x_0, x)$ and $\nabla\eta$ bounded. Recall that $Ta = \Delta u$ where $u(t) = \int_0^t e^{(t-s)\Delta} a(s) ds$. Hence

$$\int_M Ta(t, x) \varphi_R(x) d\mu(x) = - \int_M \nabla u(t, x) \cdot \nabla \varphi_R(x) d\mu(x). \quad (4.3)$$

Since $\|\nabla \varphi_R\|_\infty \leq \frac{C}{R}$, we obtain

$$\left| \int_M Ta(t, x) \varphi_R(x) d\mu(x) \right| \lesssim \frac{1}{R} \int_M |\nabla u(t, x)| d\mu(x). \quad (4.4)$$

It follows immediately from Lemma 4.4 that

$$\|\nabla e^{s\Delta} g\|_{L^1(M)} \lesssim \frac{1}{\sqrt{s}} \|g\|_{L^1(M)},$$

for all $g \in L^1(M)$ and $s > 0$, therefore

$$\int_M |\nabla u(t, x)| d\mu(x) \leq \int_0^t \frac{C}{\sqrt{t-s}} \|a(s, \cdot)\|_{L^1(M)} ds \leq 2C\sqrt{t} \|a\|_\infty r^{-2}, \quad (4.5)$$

where the ball supporting a has radius r . Thus by (4.3), (4.4), (4.5), we have

$$\lim_{R \rightarrow \infty} \left| \int_M Ta(t, x) \varphi_R(x) d\mu(x) \right| = 0.$$

Hence, as $t > 0$ is such that $Ta(t, \cdot) \in L^1(M)$, dominated convergence yields

$$\int_M Ta(t, x) d\mu(x) = \lim_{R \rightarrow \infty} \int_M Ta(t, x) \varphi_R(x) d\mu(x) = 0.$$

□

Remark 4.6. The second argument shows that for almost all $t > 0$, $Ta(t, \cdot)$ is in $L^1(M)$ and has a vanishing integral. Actually, an examination of the proof (left to the reader) shows that for almost all $t > 0$, $Ta(t, \cdot)$ is up to a constant depending on t, t_0, r a molecule in M , hence belongs to $H_{CW}^1(M)$. Calling $c_{t_0, r}(t)$ its norm in this space, one can check that $\int_0^\infty c_{t_0, r}(t) dt \lesssim 1$ uniformly with respect to t_0, r , that is uniformly with respect to the atom a . Thus, the maximal regularity map T extends boundedly from $H_z^1(X)$ into $L^1(\mathbb{R}_+^*, H^1(M))$. It is easy to show that $H_z^1(\mathbb{R}_+^* \times M)$ and $L^1(\mathbb{R}_+^*, H^1(M))$ are not comparable spaces. This extra regularity of the map T is thus not shared by the map $f \mapsto u'$.

We conclude this section by observing that the source space $H_z^1(X)$ is best possible.

Proposition 4.7. *If $M = \mathbb{R}^n$, T is not bounded from $H_r^1(X)$ to $L^1(X)$.*

Proof: Let $f(t, x) = 1$ if $0 < t < 1$ and $|x| \leq 1$, and 0 otherwise. We have seen that $f \in H_r^1$. We claim that $Tf \notin L^1(X)$. Note that $X = \mathbb{R}_+^* \times \mathbb{R}^n$.

Recall that if $p_t(z)$ denotes the heat kernel, then the time derivative is given by

$$p'_t(z) = \frac{1}{(4\pi t)^{n/2}} \left(\frac{|z|^2}{4t^2} - \frac{n}{2t} \right) e^{-\frac{|z|^2}{4t}},$$

hence

$$\frac{|z|^2}{4t} \leq \frac{n}{4} \implies p'_t(z) \leq -\frac{n}{4t(4\pi t)^{n/2}}.$$

Thus, if $|x| \leq \frac{1}{2}(nt)^{1/2}$, $|y| \leq 1$, $t \geq \sup(4/n, 2) = \tau$ and $0 < s < 1$ we have

$$p'_{t-s}(x-y) \leq -\frac{n}{4t(4\pi t)^{n/2}}.$$

It follows that there is a constant $c_n > 0$ such that

$$\begin{aligned} \iint |Tf(t, x)| dt dx &\geq \int_\tau^\infty \int_{|x| \leq \frac{1}{2}(nt)^{1/2}} \left| \int_{|y| \leq 1} \int_0^1 p'_{t-s}(x-y) ds dy \right| dt dx \\ &\geq c_n \int_\tau^\infty \frac{dt}{t} \\ &= +\infty. \end{aligned}$$

5 Boundedness of T^*

In this section, we study the boundedness of T^* given by (2.3). If we let

$$v(t, x) = \int_t^{+\infty} [e^{(s-t)\Delta} f(s, \cdot)](x) ds, t > 0,$$

then $\Delta v = T^*f$ and satisfies:

$$\frac{\partial v}{\partial t}(t, x) + \Delta v(t, x) = -f(t, x).$$

We now prove the part of Theorem 2.1 concerning the boundedness of T^* .

Proof: As before, it is enough to prove the boundedness from $H_r^1(X)$ into $H^1(X)$ and for that purpose that H_r^1 -atoms are mapped to H^1 -molecules of Definition 3.12 up to a multiplicative constant. There are two kinds of atoms in $H_r^1(X)$: type (a) atoms and type (b) atoms.

Suppose a is a type (a) atom, with $\text{supp } a \subseteq Q = \{(t, x) : \tilde{d}(t, x; t_0, x_0) < r\}$ and $4Q \subseteq X$. To see that T^*a is a molecule we prove (1) and (2) in Definition 3.12. We begin with (2). First, since T is $L^2(X)$ -bounded, so is T^* and

$$\|T^*a\|_{L^2(4Q \cap X)} \lesssim \nu(Q \cap X)^{-\frac{1}{2}} \lesssim \nu(4Q \cap X)^{-\frac{1}{2}}. \quad (5.1)$$

Next, we show that there exists $\alpha > 0$ such that for each $j \geq 2$ and $(t, x) \in B_j(Q)$,

$$|T^*a(t, x)| \lesssim 2^{-j\alpha} \nu(2^{j+1}Q \cap X)^{-1}. \quad (5.2)$$

We observe that $T^*a(t, x) = 0$ if $t > t_0 + r^2$. Hence we assume $t \leq t_0 + r^2$. For $(t, x) \in B_j(Q)$, we have that $2^j r \leq \tilde{d}(t, x; t_0, x_0) \leq 2^{j+1}r$. So we distinguish in addition two cases:

Case I: $2^j r \leq d(x, x_0) \leq 2^{j+1}r, |t - t_0| \leq (2^{j+1}r)^2$.

Case II: $(2^j r)^2 \leq |t - t_0| \leq (2^{j+1}r)^2, d(x, x_0) \leq 2^j r$.

Case I: $2^j r \leq d(x, x_0) \leq 2^{j+1}r$ and $|t - t_0| \leq (2^{j+1}r)^2$. If, in addition $|t - t_0| \leq (2r)^2$ then

$$\begin{aligned} |T^*a(t, x)| &= \left| \int_t^{+\infty} \int_{d(y, x_0) < r} p'_{s-t}(x, y) a(s, y) ds d\mu(y) \right| \\ &\lesssim \frac{1}{\nu(Q)} \int_t^{t_0+r^2} \int_{d(y, x_0) < r} \frac{1}{(s-t)V(x, \sqrt{s-t})} e^{-\frac{Cd(y, x)^2}{s-t}} ds d\mu(y) \\ &\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_t^{t_0+r^2} \frac{1}{(s-t)V(x, \sqrt{s-t})} e^{-\frac{C4^j r^2}{s-t}} ds \lesssim e^{-C4^j} \frac{1}{\nu(2^j Q \cap X)}, \end{aligned}$$

where we used Lemma 4.2 and that a is supported in Q . The last inequality is obtained repeating the proof of (4.2).

If $t \leq t_0 - (2r)^2$, as $t \leq t_0 + r^2$, we have $t \leq t_0 - 4r^2$, then

$$\begin{aligned} |T^*a(t, x)| &\leq \frac{1}{r^2} \int_{t_0-4r^2}^{t_0+r^2} \frac{1}{(s-t)V(x, \sqrt{s-t})} e^{-\frac{C4^j r^2}{s-t}} ds \\ &\leq \frac{V(x_0, r)}{\nu(Q)} \int_{\frac{4^j r^2}{t_0+4r^2-t}}^{\frac{4^j r^2}{t_0-r^2-t}} \frac{e^{-Cu} du}{uV(x, \frac{2^j r}{\sqrt{u}})} \\ &\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_{\frac{4^j}{2}}^{+\infty} \frac{e^{-Cu} du}{uV(x, \frac{2^j r}{\sqrt{u}})} \\ &\lesssim e^{-c4^j} \frac{1}{\nu(2^{j+1}Q \cap X)}, \end{aligned}$$

where the second inequality is obtained by substitution, the third from the fact that $t_0 - t + r^2 \geq r^2$, and the last one repeating the proof of (4.2).

Case II: $(2^j r)^2 \leq |t - t_0| \leq (2^{j+1} r)^2$, $d(x, x_0) \leq 2^j r$. In this case, using the vanishing moment condition of a , we can write $T^*a(t, x)$ as follows:

$$\begin{aligned} T^*a(t, x) &= \int_{t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) \leq r} p'_{s-t}(x, y) a(s, y) ds d\mu(y) \\ &= \int_{t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) \leq r} (p'_{s-t}(x, y) - p'_{s-t}(x, x_0)) a(s, y) ds d\mu(y) \\ &\quad + \int_{t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) \leq r} (p'_{s-t}(x, x_0) - p'_{r^2+t_0-t}(x, x_0)) a(s, y) ds d\mu(y) = J_1 + J_2. \end{aligned}$$

Following the proof of (4.2), we have the estimation of J_1 :

$$|J_1| \lesssim \frac{1}{\nu(Q)} \int_{t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) \leq r} \frac{1}{(s-t)V(x, \sqrt{s-t})} e^{-\frac{Cd(y, x)^2}{s-t}} ds d\mu(y) \lesssim \frac{1}{2^{j\gamma} \nu(2^{j+1}Q \cap X)},$$

For J_2 , we have the estimation:

$$\begin{aligned} |J_2| &\leq \int_{t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) \leq r} \int_{\theta=s}^{t_0+r^2} |p''_{\theta-t}(x, x_0)| d\theta |a(s, y)| ds d\mu(y) \\ &\lesssim \frac{1}{\nu(Q)} \int_{s=t_0-r^2}^{t_0+r^2} \int_{d(y, x_0) \leq r} \int_{\theta=s}^{t_0+r^2} \frac{d\theta ds dy}{(\theta-t)^2 V(x, \sqrt{\theta-t})} \lesssim \frac{1}{4^j} \cdot \frac{1}{\nu(2^j Q \cap X)}, \end{aligned}$$

where we use Lemma 4.2 and the fact that $\theta - t \sim (2^j r)^2$ to get the last inequality.

Now the argument to show $\int_X T^*a(t, x) dt d\mu(x) = 0$ is the same as in Lemma 4.5. This proves that T^*a is a molecule.

Now we prove that for type (b)-atoms b , T^*b is a molecule. Suppose that b is a type (b) atom, supported in $Q = \{(t, x) \in X, \tilde{d}(t, x; t_0, x_0) < r\}$ with $2Q \subseteq X$ but $4Q \not\subseteq X$, and $\|b\|_\infty \leq \frac{1}{\nu(Q)}$.

As for type (a) atoms, we have

$$\|T^*b\|_{L^2(4Q \cap X)} \lesssim \nu(Q)^{-\frac{1}{2}} \lesssim \nu(4Q \cap X)^{\frac{1}{2}}. \quad (5.3)$$

Next, we establish that for all $j \geq 2$ and $(t, x) \in B_j(Q)$,

$$|T^*b(t, x)| \lesssim \nu(2^j Q \cap X)^{-\frac{1}{2}} e^{-C4^j}. \quad (5.4)$$

Fix $j \geq 2$ and suppose $(t, x) \in B_j(Q)$. Observe that $T^*b(t, x) = 0$ if $t > t_0 + r^2$. As $t_0 \leq (4r)^2$, this implies that $|t - t_0| \leq 16r^2$. Thus the condition $2^j r \leq \tilde{d}(t, x; t_0, x_0) \leq 2^{j+1} r$ implies $d(x, x_0) \sim 2^j r$ for $j \geq 5$ (we skip the details for

$j||e4$ which are not difficult). Like the proof of (4.2), we have

$$\begin{aligned}
|T^*b(t, x)| &\leq \int_t^{t_0+r^2} \int_{d(y, x_0) \leq r} |p'_{s-t}(x, y)| |a(s, y)| ds d\mu(y) \\
&\lesssim \frac{1}{\nu(Q)} \int_t^{t_0+r^2} \int_{d(y, x_0) \leq r} \frac{1}{(s-t)V(x, \sqrt{s-t})} e^{-\frac{Cd(x, y)^2}{s-t}} ds d\mu(y) \\
&\lesssim \frac{V(x_0, r)}{\nu(Q)} \int_t^{t_0+r^2} \frac{1}{(s-t)V(x, \sqrt{s-t})} e^{-\frac{C4^j}{s-t}} ds \\
&\lesssim e^{-C4^j} \frac{1}{\nu(2^j Q)}.
\end{aligned}$$

The second argument in the proof of Lemma 4.5 implies that $\iint_X T^*b d\nu = 0$. Hence, T^*b is, up to a multiplicative constant, a molecule in $H^1(X)$. The theorem is proved. \square

Remark 5.1. As for T , one can check that the proof yields that T^* is bounded from $H_r^1(X)$ into $L^1(\mathbb{R}_+^*, H^1(M))$.

Remark 5.2. If $X = \mathbb{R}_+^* \times \mathbb{R}^n$, T^* is not bounded on $L^1(X)$. Indeed, this would mean that $\iint_X \chi_{s>t} |p'_{s-t}(x-y)| ds dy$ is bounded uniformly with respect to $(t, x) \in X$. But this integral equals $c \int_0^\infty \frac{dt}{t} = \infty$ as $c = \int_{\mathbb{R}^n} |p'_1(x)| dx > 0$.

6 Maximal L^q regularity on L^p

The well-known maximal L^q regularity of the Cauchy problem in our setting (see [11]) states as follows

Theorem 6.1. *Assume that M satisfies (D) and (P). There is maximal L^q regularity property for the Cauchy problem (2.1):*

$$\|u'\|_{L^q(\mathbb{R}_+, Y)} + \|Tf\|_{L^q(\mathbb{R}_+, Y)} \lesssim \|f\|_{L^q(\mathbb{R}_+, Y)},$$

where $Y = L^p(M)$, $1 < p, q < \infty$.

Let us see a proof as a corollary of Theorem 2.1 (as said in the introduction, [10] gives a proof under weaker assumptions).

Proof: It suffices to prove that T is $L^p(X)$ -bounded, with $1 < p < \infty$. To see this, we construct an operator \tilde{T} as follows:

$$\begin{cases} \tilde{T}f(t, x) = T(\chi_X f)(t, x), & t > 0, \\ \tilde{T}f(t, x) = 0, & t \leq 0. \end{cases}$$

Then by Theorem 2.1, we know that T is bounded from $H_z^1(X)$ to $H_z^1(X)$, therefore \tilde{T} is bounded from $H^1(N)$ to $H^1(N)$. On the other hand T is $L^2(X)$ -bounded, so \tilde{T} is $L^2(N)$ -bounded too. Then using the interpolation theorem in [7], we obtain

that \tilde{T} is $L^p(N)$ -bounded, where $1 < p \leq 2$. Therefore T is $L^p(X)$ -bounded, $1 < p \leq 2$.

Similarly, we can repeat the argument for the operator \tilde{T}^* defined by

$$\begin{cases} \tilde{T}^* f(t, x) = T^*(\chi_X f)(t, x), & t > 0, \\ \tilde{T}^* f(t, x) = 0, & t \leq 0. \end{cases}$$

and obtain that T^* is $L^p(X)$ -bounded, $1 < p \leq 2$.

By duality, we obtain that T is $L^p(X)$ -bounded, $2 < p < \infty$. \square

7 Sublaplacians of connected Lie groups of polynomial growth

First let us recall some properties of connected Lie groups of polynomial growth. For further informations, see [22], [1], and the references given therein.

Let G be a connected Lie group. Assume G is unimodular and fix a (left-invariant) Haar measure μ on G . If A is a measurable subset of G , we will denote the measure of A by $|A|$.

Let X_1, \dots, X_n be left-invariant vector fields on G which satisfy the Hörmander's condition, i.e. they generate, together with their successive Lie brackets $[X_{i_1}, [\dots, X_{i_\alpha}] \dots]$, the Lie algebra of G .

Associated with X_1, \dots, X_n , in a canonical way, there is a control distance ρ , which is left-invariant and compatible with the topology on G .

If we put $V(r) = |B(e, r)|$ for $r > 0$, where $B(e, r) = \{x \in G : \rho(x, e) < r\}$, there exists $d \in \mathbb{N}$, which depends on the vector fields, such that

$$V(r) \sim r^d, \quad \text{for } 0 < r \leq 1.$$

We assume that G has polynomial volume growth, namely that

$$V(r) \sim r^D, \quad \text{for } 1 < r < +\infty.$$

Otherwise, by a theorem of Guivarc'h, G would have exponential growth. By direct calculation, we have the "doubling property" : $V(sr) \lesssim s^{\max\{d, D\}} V(r)$, where $s > 1$.

For example every connected nilpotent Lie group has polynomial volume growth.

On connected Lie groups of polynomial growth, associated with the sublaplacian $\Delta = \sum_{j=1}^n X_j^2$ and corresponding heat kernel, Poincaré inequalities also hold (see [21], [19]). The heat kernel satisfies the Gaussian upper estimate (G) (see [22], Theorem VIII.2.9). In this setting, the Harnack inequality holds too, see [1], Theorem 3.1. Gaussian estimates (4.1) of time derivatives of the heat kernel are obtained in [22], Theorem VIII.2.4. Next, as Harnack inequality holds, Lemma 4.3 is valid.

Eventually Lemma 4.4 is easy. The simple proof is a consequence of the following Gaussian estimates of first order space derivatives of the heat kernel.

Lemma 7.1. ([22], Theorem VIII.2.7) *If G is a connected Lie group with polynomial growth, then for $i \in \{1, \dots, n\}$, we have*

$$|X_i p_t(x)| \lesssim t^{-1/2} V(\sqrt{t})^{-1} e^{-c \frac{\rho(x,e)^2}{t}},$$

for all $t > 0$ and $x \in G$.

So far we have recovered all the basic tools used in Section 4. The “parabolic” Hardy spaces on $N = \mathbb{R} \times G$ and the z and r Hardy spaces on $X = \mathbb{R}_+^* \times G$ can be defined analogously. Consider the Cauchy problem associated with Δ and define the operators T and T^* as before.

Theorem 7.2. *Let G be a connected Lie group with polynomial growth. Then T is bounded on $H_z^1(X)$ and T^* is bounded from $H_r^1(X)$ to $H_z^1(X)$.*

The proof is the same as before. Details are left to the reader.

8 Second order operators

Let $A \in L^\infty(\mathbb{R}^n, M_n(\mathbb{C}))$. Assume that A is uniformly elliptic, which means that there exists $\delta > 0$ such that, for almost all $x \in \mathbb{R}^n$ and all $\xi \in \mathbb{C}^n$,

$$\operatorname{Re} A(x) \xi \cdot \bar{\xi} \geq \delta |\xi|^2.$$

Let Ω be an open subset of \mathbb{R}^n . If $V(\Omega)$ is a closed subset of $W^{1,2}(\Omega)$ containing $W_0^{1,2}(\Omega)$, there is a unique operator L which is the maximal accretive operator associated with the accretive sesquilinear form

$$Q(f, g) = \int_{\Omega} A(x) \nabla f(x) \cdot \overline{\nabla g(x)} dx$$

for all $(f, g) \in V \times V$. We are interested in the Dirichlet boundary condition ($V = W_0^{1,2}(\Omega)$) and the Neumann boundary condition ($V = W^{1,2}(\Omega)$).

Say that L satisfies (G) is the following three conditions hold:

The kernel of e^{-tL} , denoted by $K_t(x, y)$, is a measurable function on $\Omega \times \Omega$ and there exist $C_G, \alpha > 0$ such that, for all $0 < t < +\infty$ and almost every $x, y \in \Omega$,

$$|K_t(x, y)| \leq \frac{C_G}{t^{n/2}} e^{-\alpha \frac{|x-y|^2}{t}}. \quad (8.1)$$

For all $y \in \Omega$ and all $0 < t < +\infty$, the function $x \mapsto K_t(x, y)$ is Hölder continuous in Ω and there exist $C_H, \mu > 0$ such that, for all $0 < t < +\infty$ and all $x, x', y \in \Omega$,

$$|K_t(x, y) - K_t(x', y)| \leq \frac{C_H}{t^{n/2}} \frac{|x - x'|^\mu}{t^{\mu/2}}. \quad (8.2)$$

For all $x \in \Omega$ and all $0 < t < +\infty$, the function $y \mapsto K_t(x, y)$ is Hölder continuous in Ω and there exist $C_H, \mu > 0$ such that, for all $0 < t < +\infty$ and all $y, y', x \in \Omega$,

$$|K_t(x, y) - K_t(x, y')| \leq \frac{C_H}{t^{n/2}} \frac{|y - y'|^\mu}{t^{\mu/2}}. \quad (8.3)$$

We assume from now on that Ω is a strongly Lipschitz domain (see [2] for a definition). Hence Ω equipped with Euclidean distance and Lebesgue measure is a space of homogeneous type. We set $N = \mathbb{R} \times \Omega$, $X = \mathbb{R}_+^* \times \Omega$ and we can define the parabolic Hardy spaces $H^1(N)$, $H_r^1(X)$ and $H_z^1(X) = H^1(X)$ following Section 3. Note that the subscripts r and z are with respect to extension with respect to the variable t across $t = 0$, the spatial variable remaining in Ω .

Consider the Cauchy problem associated with $-L$ on X and call again T the map of maximal regularity and T^* its adjoint in $L^2(X)$. The result is as follows.

Theorem 8.1. *Suppose Ω is a strongly Lipschitz domain and that L satisfies condition (G).*

1. *In the case of Dirichlet boundary condition, T is bounded from $H_z^1(X)$ into $H_r^1(X)$ and T^* from $H_r^1(X)$ into $H_z^1(X)$.*
2. *In the case of Neumann boundary condition, T is bounded from $H_z^1(X)$ into $H_z^1(X)$ and T^* from $H_r^1(X)$ into $H_z^1(X)$.*

For example, this theorem applies when A has real-valued entries.

Let us sketch the proof as it is essentially similar to the ones in the case of manifolds.

First, the four lemmas on $p_t(x, y)$ have analogs for $K_t(x, y)$. The upper bound in (G) correspond to the Gaussian upper bound of Lemma 4.1. By analyticity of the semigroup generated by $-L$, it easily implies Lemma 4.2 and Lemma 4.3. Finally, the estimate on the gradient of Lemma 4.4 is a consequence of [2, Proposition A.4].

This yields that the action of T and T^* on appropriate atoms are functions with same estimates as molecules.

For both boundary conditions the source spaces are the same. What changes are the target spaces, let us explain why. In the case of Neumann boundary conditions, one has $\int_\Omega \partial_t K_t(x, y) dx = 0$ so one expects Ta and T^*a to have mean value zero. Indeed, this is obtained by mimicking the second argument for Lemma 4.5. In the case of Dirichlet boundary condition, the first argument for Lemma 4.5 applies. Hence, when the corresponding atom has mean value zero, then the same is true for its image under T or T^* . One concludes with Proposition 3.13. Further details are left to the reader.

Remark 8.2. One could also consider elliptic operators on manifolds with (D) and (P) or subelliptic operators on connected Lie groups of polynomial growth with coefficients and appropriate estimates, and develop a similar theory. This is left to the reader.

Remark 8.3. All results have local analogs under local doubling, local Poincaré and local Gaussian bounds in which t is restricted to a finite interval $(0, \tau)$. In this case, \mathbb{R}_+^* and \mathbb{R} are replaced in X and N by $(0, \tau)$ and $(-\tau, \tau)$ and the Hardy spaces by their local analogs in the sense of Goldberg (where one relaxes the mean value zero property for atoms or molecules associated to large balls). This is left to the interested reader.

Acknowledgements: The authors thank Thierry Coulhon for indicating useful bibliographical references and Emmanuel Russ for reading carefully the manuscript and suggesting some improvements. The third author would like to express her thanks to professor Lizhong Peng and professor Heping Liu for their encouragements, and thanks the Department of Mathematics of the University Paris-Sud for its hospitality.

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