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Integrated optical amplifiers and microspherical lasers based on erbium-doped oxide glasses

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Er³⁺-doped glasses have been a subject of great interest in the recent years for their application in the areas of guided wave optical amplifiers and lasers. Oxide-glass matrices, in particular, offer the advantage of relatively simple fabrication processes—both for bulk glasses and optical fibers—and have demonstrated quite good properties in the 1.5 μm wavelength band. Here we present some results we have obtained in the development of different silica-based glasses and in their application to the production of integrated optical amplifiers and microspherical lasers. As to the former application, two classes of silica glasses have been produced and investigated, namely a melted soda-lime-alumino-silicate glass and a sol-gel silica-hafnia glass. Both of them exhibit a wide emission bandwidth and seem quite promising for broadband optical amplifiers. We also demonstrated low-threshold microlasers, based on whispering gallery mode spherical resonators, fabricated in different oxide glasses.

Keywords: Silica glass; Er³⁺-doped glass; Optical amplifier; Microsphere; Microlaser; Integrated optics

1. Introduction

The possibility of confining light inside guiding films and structures has opened new perspectives in the very broad area of passive and active optical components for either telecom or sensing systems. The aim of this paper is to discuss some advances in guided-wave struc-

tures for light amplification and lasing, based on oxide glasses. In particular, reference will be made to (a) planar optical waveguides, fabricated in erbium-doped silica glasses, that appear quite promising for the development of broadband integrated optical amplifiers, and (b) dielectric spherical microresonators, where the light is guided through whispering-gallery modes (WGMs), that allowed us to demonstrate low-threshold microlasers.

For what concerns the development of Er³⁺-doped waveguide amplifiers (EDWAs), we have focused our

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attention on silica-based glasses in virtue of their good chemical durability and adaptability to different waveguide fabrication process. A net optical gain of 1.5 dB/cm was demonstrated in $\text{Er}^{3+}/\text{Yb}^{3+}$ co-doped soda-lime-silicate glasses [1,2]. Their performance, however, was limited by the relatively narrow fluorescence bandwidth, the full width half maximum (FWHM) value being around 17 nm. We have therefore searched for novel materials that could offer a larger bandwidth: as it is well known that introduction of aluminum in a silicate glass may lead to increase the emission bandwidth [3,4], a first approach has been that of co-doping our soda-lime-silicate glass matrix with different quantities of aluminum oxide (up to about 20 mol%) and to investigate the effect of aluminum content on the optical properties of the glass and in particular on the emission bandwidth of the ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{15/2}$ Er^{3+} ions transition at 1550 nm. A second approach has concerned the use of sol-gel-derived silica-hafnia films, that represent an innovative binary glass system with valuable optical, spectroscopic and structural characteristics [5]. Here we report on the spectroscopic properties of these two glass systems, and on the characteristics of the corresponding optical waveguides.

As to the second issue, it is well known that in dielectric spheres light can be guided through high-Q whispering-gallery-modes with an unique combination of strong temporal and spatial confinement of light. Glass microspheres are of interest for a large number of applications, such as cavity quantum electrodynamics [6], nonlinear optics [7], photonics [8] and chemical or biological sensing [9]. In particular, since the studies on morphology-dependent-resonances (MDRs) and lasing effects in droplets during the 1980s [10], rare earth doped glass microsphere lasers have become the subject of many investigations and have been demonstrated as potentially compact laser sources [11–13]. Here we briefly report both on the spectroscopic characterization of an Er^{3+} doped microsphere in a modified-silicate glass, and on lasing action in that glass and in an $\text{Er}^{3+}/\text{Yb}^{3+}$ co-doped phosphate glass.

2. Experimental

2.1. Sol-gel silica-hafnia films

By using the sol-gel process, we produced and subsequently tested several films of composition $(100 - x)\text{SiO}_2 - x\text{HfO}_2$ ($x = 10, 20, 30, 40$), with the aim of assessing the role of hafnia concentration on the structural and optical properties of these waveguides. Two different sets of films were produced, containing 0.01 and 0.3 mol% Er^{3+} ions, respectively. All the films were deposited on cleaned pure-silica substrates by dip-coating. After each layer's deposition, the sample was

annealed in air for 50 s at 900 °C. After a 10-dip cycle, the film was heated for 2 min at 900 °C. Finally, the waveguides were subject to a further annealing at 900 °C, whose duration was different for each waveguide; for instance, for the samples containing 0.3 mol% Er^{3+} the annealing time was varying from 30 h ($x = 10$) down to 5 min ($x = 40$). The necessary number of dips for the corresponding samples was 50 and 25. More details on the fabrication process are available in previous papers [5,14].

Modal characterization was performed by the well known dark-line spectroscopy; one TE and one TM mode at 632.8 nm were supported by each waveguide.

Photoluminescence measurements in the region of the ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{15/2}$ transition and decay curves from the ${}^4\text{I}_{13/2}$ level were obtained with the experimental set-up described elsewhere [15]. All the spectroscopic measurements were, thus, made in guided-wave configuration, using prism coupling and as excitation source either the 980 nm line of a Ti:sapphire laser or the 514.5 nm line of an argon laser.

2.2. Soda-lime-alumino-silicate glasses and ion-exchanged waveguides

These glasses were prepared by conventional melting process and had a basic composition of the type $\text{Na}_2\text{O}-\text{CaO}-\text{Al}_2\text{O}_3-\text{SiO}_2$ with small percentages of P_2O_5 and K_2O . The batches were prepared by using reagent grade SiO_2 , Na_2CO_3 , K_2CO_3 , Al_2O_3 , CaCO_3 , $(\text{NH}_4)_3\text{PO}_4$, and Er_2O_3 as starting materials. Melting occurred in an electrically heated furnace within Pt crucibles, following the same heating cycle: from 20 to 1000 °C at 10 °C/min, with a 24 h soaking time at 1000 °C, from 1000 to 1550 °C at 20 °C/min and finally 1 h of soaking time at the maximum temperature of 1550 °C. Eventually, the melt was quenched in a graphite mould to obtain small pieces of glass having a bar form; each bar was then cut to 1 mm thickness and optically polished on both faces. Room temperature absorption spectra were measured in the visible region using a Kontron UVICON 941 Plus double beam spectrophotometer (spectral bandwidth of 0.5 nm) and in the near infrared region with a Nicolet Magna 760 FTIR spectrometer (spectral bandwidth of 2 cm^{-1}) or a double beam Perkin-Elmer spectrophotometer with a resolution of 1 nm. The fluorescence spectra of Er^{3+} ions were detected using a 976 nm laser diode as excitation source: luminescence was dispersed using a monochromator with a resolution of 1 nm and detected with an InGaAs photodiode and lock-in technique. Lifetime τ of the ${}^4\text{I}_{13/2}$ level was also measured after excitation at 976 nm.

As to the waveguide fabrication process in these glasses, our attention was focused on ion-exchange; indeed, we had designed a soda-lime-silicate glass just in

order to take advantage of a large content of Na^+ ions. Thus, planar optical waveguides were produced by using the $\text{Ag}^+ \rightleftharpoons \text{Na}^+$ exchange process in a dilute AgNO_3 solution ($\text{AgNO}_3:\text{NaNO}_3 = 0.5:99.5$ mol%) at 325°C . Exchange times were chosen in such a way to obtain multimode waveguides, in order to use with good accuracy the inverse WKB (IWKB) method for the determination of the refractive index profile resulting from the exchange process. Modal characterization at various wavelengths was performed by dark line spectroscopy using a semi-automatic system (*COMPASSO*) developed at IFAC [16]. The same device had also been used to measure the refractive indices of the bulk glasses.

2.3. Er-doped glass microspherical resonators

Different kinds of glasses were used to produce microspheres: three experimental modified-silica glass samples, doped with 0.2, 0.5 and 1.5 mol% of erbium, respectively [17], as well as a commercial $\text{Er}^{3+}/\text{Yb}^{3+}$ co-doped phosphate glass (Schott IOG-2), which contains 2 wt% of Er_2O_3 and 3 wt% of Yb_2O_3 . The refractive indices at 632.8 nm of the former three glasses were equal to 1.5604, 1.5607, and 1.5871, respectively.

Each glass was ground, and microspheres were prepared from the powders by using a microwave plasma torch. Powders were axially injected and melted when passing through the plasma flame, superficial tension forces giving them their spherical shape. Free spheres were collected a few ten centimeters down. The diameter of the spheres depends essentially on the powder size and may vary from 10 to 200 μm [18]. By visual analysis with a microscope, the microspheres with diameter in the range 50–100 μm were selected and then glued to the stretched tip of an optical fibre (~ 20 – 30 μm in diameter), which in turn was mounted on a submicrometer translation stage.

Photoluminescence spectroscopy in the 1.5 μm region was performed using the 980.8 nm line of a Ti:sapphire laser as excitation source. The luminescence was dispersed by a 320 mm single-grating monochromator with a resolution of 2 nm. The light was detected using a Si/InGaAs two-colors photodiode and standard lock-in technique. Lifetime decay curves were obtained recording the signal by a digital oscilloscope. For both the precursor bulk samples and the microspheres the standard bulk measurement configuration has been used.

A single half-fibre-taper was used to couple the pump light from a pigtailed laser diode into the WGMs of the microsphere and—at the same time—to allow us to couple the fluorescence or laser light out of the microsphere. The half-taper was obtained using a fusion splicer by heating and stretching a standard telecommunication fibre until breaking. The drawn length was typically

850 μm , and the taper end was reduced down to 1.5 μm in diameter.

3. Results and discussion

3.1. Sol-gel silica–hafnia waveguides for EDWAs

Table 1 summarizes some optical and spectroscopic parameters of the fabricated silica–hafnia planar waveguides, together with their compositions. In all samples, birefringence is not negligible, as indicated by the significant difference between the refractive indices measured for TE and TM polarizations. Propagation losses were also measured in various samples by using the technique based on the collection of the light scattered out of the plane of the waveguide, and an average value was ≈ 0.9 dB/cm at 632.8 nm.

The photoluminescence (PL) spectra relative to the ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{15/2}$ transition of the Er^{3+} ions for all the waveguides, obtained upon excitation at 514.5 nm and 980 nm, exhibit a main emission peak at 1.53 μm and a broad FWHM spectral width of about 50 ± 2 nm. Fig. 1 shows the emission spectra of the samples doped with 0.3 mol% Er^{3+} for excitation at 514.5 nm; the shape and the bandwidth of these spectra, however, do not change with the excitation wavelength, indicating that site selection is negligible. The similarity of PL spectra for all waveguides seems to indicate that a small inhomogeneous broadening is present independently on HfO_2 content.

The measured lifetime of the ${}^4\text{I}_{13/2}$ metastable state, also reported in Table 1, decreases with the increasing of the HfO_2 molar concentration. At very low Er^{3+} concentration, the energy transfer processes should be very improbable. Moreover, considering that in completely densified silicate glasses the multiphonon decay does not significantly affect the ${}^4\text{I}_{13/2}$ level lifetime, we can assume that the lifetimes measured at 0.01 mol% Er^{3+} doping level are very close to the radiative ones. Under this hypothesis, we can estimate that a quantum efficiency between 84% and 88% is obtained in all the samples doped with 0.3 mol% of erbium.

3.2. Soda-lime–alumino-silicate glasses and ion-exchanged waveguides

The main oxides' concentrations of the six produced soda-lime–alumino-silicate glasses are reported in Table 2, together with the emission bandwidth B , the measured τ_{meas} and the calculated radiative τ_{rad} lifetimes of the ${}^4\text{I}_{13/2}$ metastable state. The data in the table clearly indicate an increase of the fluorescence bandwidth with the increase of alumina content.

This is also shown in Fig. 2, where the FWHM bandwidth of both the absorption and the emission peak

Table 1

Compositional, optical and spectroscopic parameters of the Er^{3+} -doped $(100 - x) \text{SiO}_2 - x\text{HfO}_2$ planar waveguides

Waveguide label	$\text{SiO}_2:\text{HfO}_2$ molar ratio	Erbium content (mol%)	n at $632.8 \text{ nm} \pm 0.005$		Thickness (μm) ± 0.05	${}^4\text{I}_{13/2}$ lifetime (ms) ± 0.5
			TE	TM		
SH10	90:10	0.3	1.470	1.468	2.05	7.1
SH20	80:20	0.3	1.543	1.537	1.27	6.5
SH30	70:30	0.3	1.600	1.588	1.10	6.0
SH40	60:40	0.3	1.663	1.646	0.96	5.8
SH10a	90:10	0.01	1.494	1.494	1.16	8.5
SH20a	80:20	0.01	1.562	1.562	0.69	7.5
SH30a	70:30	0.01	1.605	1.601	0.69	7.0
SH40a	60:40	0.01	1.664	1.655	0.53	6.6

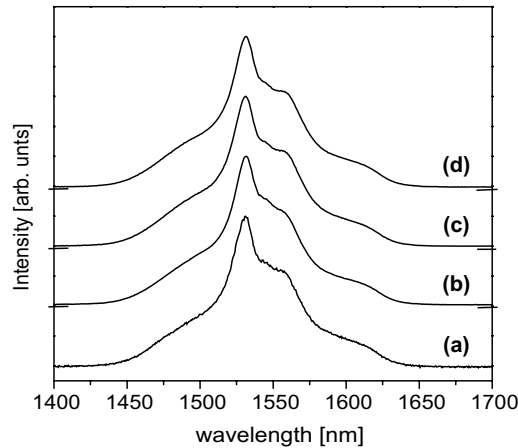


Fig. 1. Photoluminescence spectra relative to the ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{15/2}$ transition of the Er^{3+} ions for the (a) SH40, (b) SH30, (c) SH20, and (d) SH10 waveguides, obtained upon excitation at 514.5 nm.

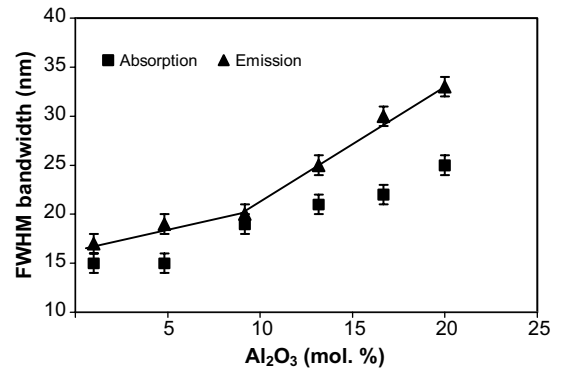


Fig. 2. Effective bandwidths (absorption and emission) as a function of Al_2O_3 concentration. The line is only traced as an help to eyes.

Table 2

Main chemical components, FWHM emission bandwidth, measured and calculated lifetimes of the ${}^4\text{I}_{13/2}$ metastable state of the six aluminosilicate glasses we had produced

Sample	Main oxides content (mol%)					B ($\pm 1 \text{ nm}$)	$\tau_{\text{meas.}}$ ($\pm 0.2 \text{ ms}$)	$\tau_{\text{rad.}}$ ($\pm 0.2 \text{ ms}$)
	SiO_2	Na_2O	Al_2O_3	CaO	Er_2O_3			
AL1	72.64	13.94	0.99	10.96	0.47	17	7.0	12.6
AL5	69.84	13.40	4.81	10.54	0.45	19	8.0	11.4
AL10	66.64	12.79	9.17	10.05	0.43	20	7.8	10.3
AL15	63.71	12.23	13.16	9.61	0.41	25	5.1	8.76
AL20	61.14	11.73	16.67	9.23	0.39	30	6.2	8.43
AL25	58.70	11.26	20	8.86	0.39	33	5.1	8.18

around $1.5 \mu\text{m}$ is plotted as a function of the alumina content. One can notice that the slope of increase of the emission bandwidth changes significantly from the region where the alumina content is lower than 9 mol% to the region where the content is equal to or higher than 13 mol%. Correspondingly, an abrupt change in the physical (density) and optical (bulk refractive index) properties of the glass was observed as well. These facts suggest that the behaviour of this class of

glasses has a drastic change when the concentration of aluminium ions becomes larger than that of sodium ions.

The measured lifetimes were compared with the radiative lifetimes calculated using Judd-Ofelt (J-O) analysis; in particular, the phenomenological intensity parameters Ω_i were estimated and are reported in Table 3. The values of the J-O parameters appear to be higher, especially for the samples at high alumina content, than

Table 3

Judd–Ofelt intensity parameters (Ω_i), with their uncertainties, calculated for the six soda-alumino-silicate glasses

Glass	Ω_2 (pm ²)	Ω_4 (pm ²)	Ω_6 (pm ²)
AL1	5.24 ± 0.18	1.11 ± 0.25	0.574 ± 0.088
AL5	5.43 ± 0.17	1.15 ± 0.24	0.691 ± 0.086
AL10	6.12 ± 0.19	1.38 ± 0.26	0.813 ± 0.093
AL15	6.47 ± 0.21	1.60 ± 0.28	0.98 ± 0.10
AL20	7.15 ± 0.25	1.95 ± 0.34	1.01 ± 0.12
AL25	7.89 ± 0.21	1.83 ± 0.28	1.06 ± 0.10

those found for other silicate glasses [19]. Moreover, a monotonic and remarkable increase of all the J–O parameters is observed while increasing the amount of alumina in the host glass; it is in agreement with the results reported by Tanabe [4]. These values were then used to calculate the emission probabilities between the energy levels of erbium in the soda-lime–alumino-silicate host matrix, and the radiative lifetime τ_{rad} of this emission. As an example, a quantum efficiency near 75% is estimated for the AL20 sample, which has a high Al₂O₃ concentration (16.67 mol%).

The presence of guided modes in the Ag⁺-exchanged glasses was checked by using an SF6 prism to couple light into the waveguides at the wavelength of 635 nm. The effective indices of TE modes were thus measured. It appeared that, under the same ion-exchange process conditions, the number of modes supported by each waveguide was depending on the alumina content. In order to derive accurate values for the ion-exchange parameters, namely the diffusion coefficient of incoming ions D_{Ag}^+ and the diffusion depth d , in the six glasses, we found out that the best fit for the measured effective indices was an analytical function given by a linear combination between erfc and exponential functions [20]:

$$n(x) = n_{\text{sub}} + \Delta n \left\{ H * \text{erfc} \left(\frac{x}{2\sqrt{Dt}} \right) + (1 - H) * \exp \left(- \left(\frac{x}{2\sqrt{2Dt}} \right)^3 \right) \right\} \quad (1)$$

where x is the depth coordinate, n_{sub} is the index of the glass substrate, $n_s = n_{\text{sub}} + \Delta n$ is the surface index, d is the effective depth of the waveguide, and $D = D_{\text{Ag}}$. In an analogous way, D_{Na} indicates the effective diffusion

coefficient for the Na⁺ ions. The value of D is related to the diffusion time t by the equation $d = (Dt)^{1/2}$; it also depends on the molten salt solution, the glass and the temperature. It is possible to relate the fitting value H to the Stewart coefficient $\alpha = 1 - D_{\text{Ag}}/D_{\text{Na}}$ [20]. Table 4 lists the number of TE modes and the values of D_{Ag} , Δn , H and α calculated for the six doped glasses after the ion-exchange. There is a marked decrease of Δn in this class of glasses with the increase of Al content, while the diffusion coefficient D_{Ag} does not present a monotonic behaviour, but exhibits a maximum for the AL10 sample, that has a sodium content slightly higher than aluminium content. Fig. 3 shows the behaviour of the diffusion coefficient and the index change versus the atomic concentrations of Al and Na (C_{Al} and C_{Na} , respectively). It confirms that there is a change in the diffusion process of the Ag⁺ ions when the value of C_{Al} approaches that of C_{Na} (‘breakpoint zone’). As a consequence, the reconstructed index profiles of the waveguides, which are shown in Fig. 4, have different shapes when $C_{\text{Al}} < 10$ mol%, while they are very similar for values of $C_{\text{Al}} > 10$ mol%.

3.3. Er-doped glass microspherical lasers

Spectroscopic characterization of Er³⁺-doped modified-silica glass microspheres indicated that both the

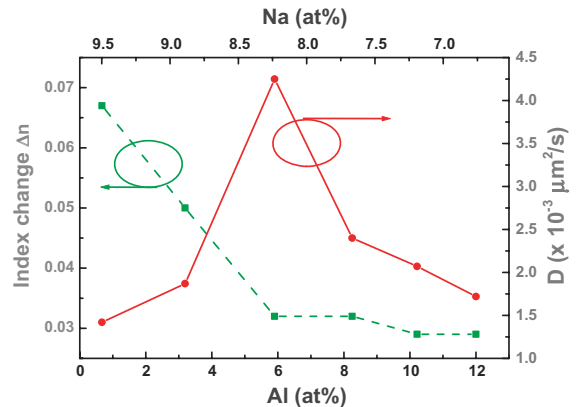


Fig. 3. Index change Δn after ion-exchange and diffusion coefficient of Ag⁺ ions versus the atomic concentrations of Al and Na. The lines are only traced as an help to eyes.

Table 4

Ion-exchange process parameters and fitted values of D_{Ag} , Δn , H and α obtained using Eq. (1) as index profile for the slab waveguides fabricated in six glasses

Sample	Melt salt and temperature	Exchange time (min)	Modes at 635 nm	Δn	D ($10^{-3} \mu\text{m}^2/\text{s}$)	H	α
AL1	AgNO ₃ :NaNO ₃ (0.5:99.5 mol%) 325 °C	90	11	0.067	1.42	0.25	0.94
AL5			11	0.050	1.87	0.24	0.93
AL10			13	0.032	4.25	0.36	0.89
AL15			10	0.034	2.57	0.50	0.83
AL20			8	0.029	2.07	0.53	0.81
AL25			7	0.029	1.72	0.58	0.78

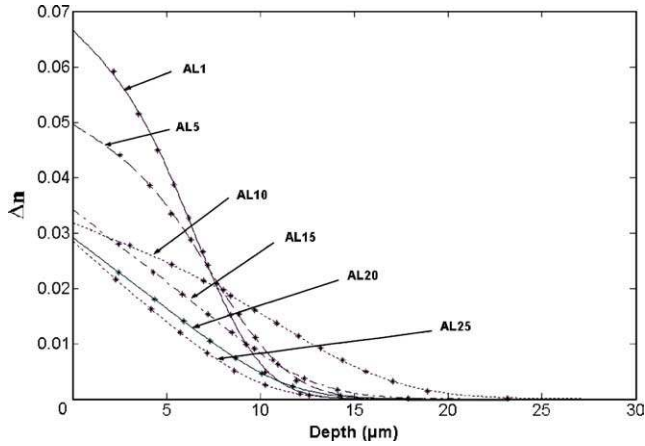


Fig. 4. Refractive index profiles of the waveguides produced by ion-exchange in the six aluminosilicate glasses. The profiles were numerically reconstructed by using the function given in Eq. (1).

luminescence spectrum and the lifetime changed when passing from the bulk glass to the microsphere [21]. In fact, the emission bandwidth is larger for microspheres than for the precursor glass, demonstrating that the site-to-site inhomogeneities are effective in the microspheres. This is probably due to the fabrication process, but, due to the small number of samples, so far it has been impossible to determine the specific reason of the induced inhomogeneity. Lifetimes measurements gave further information about the site-to-site inhomogeneities; the decay curves of the erbium in various microspheres all showed a faster relaxation than in the corresponding bulk glasses, indicating that energy transfer mechanism among active ions was not negligible. This behaviour is evident in Fig. 5, which compares the measured lifetimes of the bulk and microsphere glass doped with 0.5 mol% of erbium oxide.

Single line laser action upon pumping at 1480 nm was obtained by using a 0.5 mol% Er-doped microsphere of about 80 μm in diameter. The laser spectrum with emis-

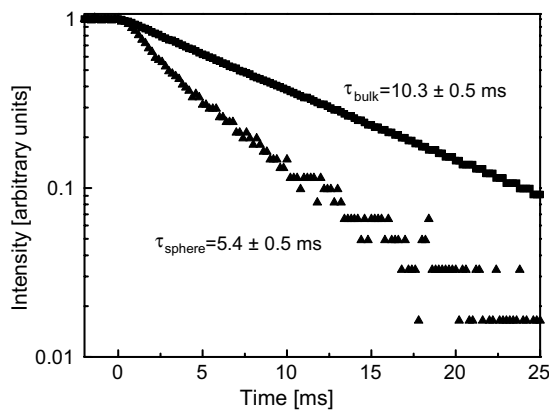


Fig. 5. Luminescence decay curve of the $^4I_{13/2}$ level of Er^{3+} ion in the bulk sample (squares) and in the microsphere (triangles) for the same modified silicate glass (0.5 mol% of erbium oxide).

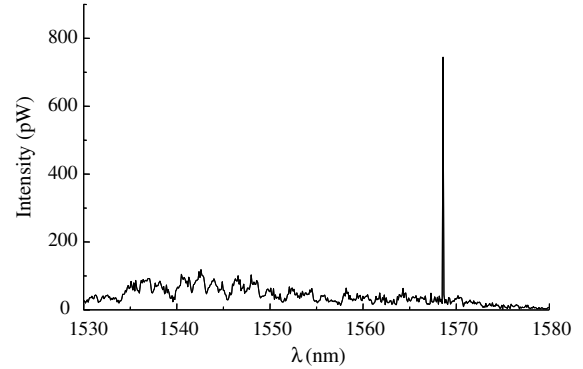


Fig. 6. Lasing spectrum of a microsphere of diameter 80 μm in the modified silicate glass.

sion at 1568.5 μm is shown in Fig. 6. The pump threshold was 15 mW (measured at the input of the tapered fiber).

As to the phosphate IOG-2 microspheres, lasing action was obtained quite easily for any sphere diameter on the range 50–100 μm , as reported in a previous paper [13].

4. Conclusions

The results we have obtained so far indicate that there are several classes of rare-earth-doped oxide glasses that are worth being investigated in greater detail with the aim of producing high-performance broadband integrated optical amplifiers. One of these glasses is represented by $\text{SiO}_2\text{-HfO}_2$ binary systems, where Er^{3+} -doped sol-gel planar waveguides have shown high quantum efficiency and ≈ 50 nm FWHM bandwidth around 1.53 μm . Work is now in progress for the fabrication of channel waveguides and their characterization.

A second class is constituted by soda-lime-aluminosilicate glasses, that combine the very good chemical durability of silicates with the better spectroscopic properties induced by the presence of a relatively high percentage of aluminium. An increase up to 60% of the effective fluorescence bandwidth was observed varying alumina content from 1% to 20% molar, and an high value of quantum efficiency was calculated for these glasses. $\text{Ag}^+ \Leftrightarrow \text{Na}^+$ ion-exchanged waveguides with good characteristics were also successfully produced. Both spectroscopic and ion-exchange results have shown that there is a significant change in the behaviour of the glass when the aluminum atomic concentration is of the same order as that of sodium: these aspects are now being investigated in greater detail. Moreover, integrated optical amplifiers in these glasses are being produced and tested.

Microspherical WGM resonators were produced in other Er^{3+} -doped oxide glasses (a modified silicate and a phosphate), with the aim of producing low-threshold

microlasers. Preliminary studies have shown that the transformation of a glass into a microsphere could actually change the environment of erbium ions, and therefore its spectroscopic properties. This issue is being further investigated. Lasing action was obtained in both the investigated glasses.

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