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# Crystal isomorphisms for irreducible highest weight $\mathcal{U}_{v}\left(\widehat{\mathfrak{s}}_{e}\right)$-modules of higher level 

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#### Abstract

We study the crystal graphs of irreducible $\mathcal{U}_{v}\left(\widehat{\mathfrak{s}}_{e}\right)$-modules of higher level $\ell$. Generalizing works of the first author, we obtain a simple description of the bijections between the classes of multipartitions which naturally label these graphs: the Uglov multipartitions. This is achieved by expliciting an embedding of the $\mathcal{U}_{v}\left(\widehat{\mathfrak{s}}_{e}\right)$-crystals of level $\ell$ into $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-crystals associated to highest weight modules.


## 1 Introduction

Let $\mathcal{U}_{v}\left(\widehat{\mathfrak{s}}_{e}\right)$ be the affine quantum group of type $A_{e-1}^{(1)}$ and $\left\{\Lambda_{i} \mid i \in \mathbb{Z} / e \mathbb{Z}\right\}$ its set of fundamental weights. Consider $\ell \in \mathbb{N}$ and $\underline{\mathbf{s}}=\left(\mathbf{s}_{0}, \ldots ., \mathbf{s}_{\ell-1}\right) \in(\mathbb{Z} / e \mathbb{Z})^{\ell}$. Denote by $V_{e}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ the irreducible $\mathcal{U}_{v}\left(\widehat{\mathfrak{s}}_{e}\right)$-module of highest weight $\Lambda_{\underline{\mathbf{s}}}=\sum_{i=0}^{\ell-1} \Lambda_{\mathbf{s}_{i}}$. The general theory of Kashiwara provides a crystal basis and a global basis for $V_{e}\left(\Lambda_{\underline{\mathbf{s}}}\right)$. The crystal basis of $V_{e}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ comes equipped with the crystal graph $B_{e}\left(\Lambda_{\underline{s}}\right)$ which encodes many informations on the module structure.
There are different possible realizations of $V_{e}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ depending on the choice of a representant $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right) \in \mathbb{Z}^{\ell}$ in the class $\underline{\mathbf{s}} \in(\mathbb{Z} / e \mathbb{Z})^{\ell}$. Indeed, to $\underline{s}$ is associated a Fock space $\mathfrak{F}_{e}^{s}$ which provides an explicit construction of $V_{e}\left(\Lambda_{\underline{\mathbf{s}}}\right)$. We will denote it $V_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ and write $B_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ for the corresponding crystal graph. The crystals $B_{\bar{e}}^{s}\left(\Lambda_{\underline{s}}\right)$ with $\underline{s} \in \underline{\mathbf{s}}$ are thus all isomorphic to the abstract crystal $B_{e}\left(\Lambda_{\underline{\mathbf{s}}}\right)$.

The purpose of the paper is to explicit the isomorphisms between the crystals $B_{e}^{s}\left(\Lambda_{\mathbf{s}}\right)$. This study is also motivated by the works of Lascoux-Leclerc-Thibon [16] and Ariki [2] on the modular representation theory of Ariki-Koike algebras.
To be more precise, let $\eta$ be a primitive $e^{\text {th }}$-root of unity. Write $\mathcal{H}=\mathcal{H}\left(\eta ; \eta^{\mathbf{s}_{0}}, \ldots, \eta^{\mathbf{s}_{\ell-1}}\right)$ for the Ariki-Koike algebra defined over analgebraically closed field $F$ of characteristic 0 . This algebra is generated by $T_{0}, \cdots, T_{n-1}$ subject to the relations $\left(T_{0}-\eta^{\mathbf{s o g}_{0}}\right) \ldots\left(T_{0}-\eta^{\mathrm{s}_{\ell-1}}\right)=0$, $\left(T_{i}-\eta\right)\left(T_{i}+1\right)=0$, for $1 \leq i \leq n$ and the type $B$ braid relations. The algebra $\mathcal{H}$ is not semisimple in general, and, by a deep Theorem of Ariki, its representation theory is intimately connected

[^0]to the global bases of the irreducible $\mathcal{U}_{v}\left(\widehat{\mathfrak{s l}}_{e}\right)$-modules. In particular, the simple modules of $\mathcal{H}$ are labelled by the vertices of any crystal $B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ such that $\underline{s} \in \underline{\mathbf{s}}$. In fact the Fock space $\mathfrak{F}_{e}^{\underline{s}}$ admits a crystal basis indexed by multipartitions of length $\ell$. This implies that the vertices of the crystal $B_{e}^{\frac{s}{e}}\left(\Lambda_{\mathbf{s}}\right)$ can be identified with certain multipartitions of length $\ell$ which are called the Uglov multipartitions. Note that when $\ell=2$, remarkable results of Geck [8] show that these multipartitions also naturally appear in the context of Kazhdan-Lusztig theory and cellular structure of Hecke algebras. Other connections between Uglov multipartitions and modular representation theory of Ariki-Koike algebras are also known to hold for $\ell>2$ (11]).
Although nice properties have been given in particular cases (see [3] and [4]), the combinatorics of the Uglov multipartitions and their associated crystal graphs are not really well understood. For example, we do not even have a non recursive characterization of the Uglov multipartitions for any $\underline{s} \in \underline{\mathbf{s}}$.

This paper extends the results obtained in 10 for $\ell=2$ by giving a combinatorial description of the isomorphisms between the crystals $B_{e}^{\frac{s}{( }}\left(\Lambda_{\underline{\mathbf{s}}}\right)$. Nevertheless, the ideas we use are quite different. Indeed, we show that the crystals $B_{\bar{e}}^{s}\left(\Lambda_{\underline{s}}\right)$ can be embedded in crystals corresponding to irreducible highest weight $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-modules. This result allows us to prove that most of the crystal isomorphisms in type $A_{e-1}^{(1)}$ can be derived from crystal isomorphisms in type $A_{\infty}$. Now the combinatorial description of the isomorphisms of $A_{\infty}$-crystals is very closed to that of the isomorphisms of $A_{r}$-crystal in finite rank $r$. This permits us to use some elegant results of Nakayashiki and Yamada [20] on combinatorial $R$-matrices in type $A_{r}$. One of the advantages of this new method is to avoid cumbersome case by case verifications unavoidable in [10].
We would like to mention also that there is another way to realize the abstract crystals $B_{e}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ by using Fock spaces of level $\ell$ which are tensor products of Fock spaces of level 1. The crystal $B_{e}^{a}\left(\Lambda_{\underline{s}}\right)$ so obtained notably appears in the works by Ariki (see [2]). The vertices of $B_{e}^{a}\left(\Lambda_{\underline{s}}\right)$ are parametrized by multipartitions called the "Kleshchev multipartitions". The crystals $B_{e}^{s}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ and $B_{e}^{a}\left(\Lambda_{\underline{s}}\right)$ does not coincide in general. In particular, our method does not permit to embed $B_{e}^{a}\left(\Lambda_{\underline{\mathbf{s}}}\right)$ in a crystal of type $A_{\infty}$ (but see the remark after Theorem 4.2.2).

The present paper is organized as follows. The second section is devoted to the combinatorial description of certain isomorphisms of $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-crystals. In section 3, we recall basic results on $\mathcal{U}_{v}\left(\widehat{\mathfrak{s}}_{e}\right)$-crystals. By using two natural parametrizations of the Dynkin diagram in type $A_{e-1}^{(1)}$, we link in particular the two usual presentations of the crystals $B \frac{s}{e}\left(\Lambda_{\underline{s}}\right)$ which appear in the literature. We then show in Section 4 that the crystals $B_{e}^{s}\left(\Lambda_{\underline{s}}\right)$ can be embedded in crystals corresponding to irreducible highest weight $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-modules. This embedding allows us to give in Section 5 a description of the isomorphisms between the crystals $B_{e}^{s}\left(\Lambda_{\underline{s}}\right)$ and to obtain another characterization for the sets of Uglov multipartitions. This characterization does not necessitate an induction on the sum of the parts of the Uglov multipartitions contrary to the original one.

## 2 Crystals in type $A_{\infty}$

In this section, we study crystal isomorphisms in type $A_{\infty}$.

### 2.1 Background on $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$

Let $\mathfrak{s l}_{\infty}$ be the Lie algebra associated to the doubly infinite Dynkin diagram in type $A_{\infty}$ (see (13] and (14).

$$
\begin{equation*}
A_{\infty}: \cdots-{ }_{\circ}^{-m}-{ }_{\circ}^{1-m}-\cdots-\stackrel{-1}{\circ}_{\circ}-0_{\circ}^{0}-\stackrel{1}{\circ}-\cdots-{ }_{\circ}^{m-1}-{ }_{\circ}^{m}-\cdots . \tag{1}
\end{equation*}
$$

We denote by $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$ the corresponding quantum group and we write $\omega_{i}, i \in \mathbb{Z}$ for the fundamental weights of the corresponding root system. We associate to the sequence $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right) \in \mathbb{Z}^{\ell}$ the dominant weight $\omega_{\underline{s}}=\sum_{i=0}^{\ell-1} s_{i} \omega_{i}$. Then the irreducible highest weight $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-modules are parametrized by the sequences $\underline{s}$ of arbitrary length $\ell$. We denote by $V_{\infty}\left(\omega_{\underline{s}}\right)$ the irreducible $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-module of highest weight $\omega_{\underline{s}}$. The module $V_{\infty}\left(\omega_{\underline{s}}\right)$ admits a crystal basis. We refer the reader to [9] for a complete review on crystal bases. We write $B\left(\omega_{\underline{s}}\right)$ for the crystal graph corresponding to $V_{\infty}\left(\omega_{\underline{s}}\right)$. When $\underline{s}=\left(s_{0}\right)$ we write for short $V_{\infty}\left(\omega_{s_{0}}\right)$ and $B_{\infty}\left(\omega_{s_{0}}\right)$ instead of $V\left(\omega_{\underline{s}}\right)$ and $B\left(\omega_{\underline{s}}\right)$.

In addition to the irreducible highest weight modules $V_{\infty}\left(\omega_{\underline{s}}\right)$, it will be convenient to consider also irreducible modules $V_{\infty}(k)$ indexed by nonnegative integers which are not of highest weight. By a column of height $k$, we mean a column shaped Young diagram

$$
C=\begin{array}{|c|}
\hline x_{1}  \tag{2}\\
\hline \cdot \\
\hline \cdot \\
\hline x_{k} \\
\hline
\end{array}
$$

of height $k$ filled by integers $x_{i} \in \mathbb{Z}$ such that $x_{1}>\cdots>x_{k}$. When $a \in C$ and $b \notin C$, we write $C-\{a\}+\{b\}$ for the column obtained by replacing in $C$ the letter $a$ by the letter $b$. The module $V_{\infty}(k)$ is defined as the vector space with basis $\mathcal{B}_{k}=\left\{v_{C} \mid C\right.$ is a column of height $\left.k\right\}$. The actions of the Chevalley generators $e_{i}, f_{i}, k_{i}, i \in \mathbb{Z}$ of $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$ are given by

$$
\begin{aligned}
& f_{i}\left(v_{C}\right)=\left\{\begin{array}{l}
0 \text { if } \delta_{i}=0 \text { or } \delta_{i+1}=1 \\
v_{C^{\prime}} \text { with } C^{\prime}=C-\{i\}+\{i+1\} \text { otherwise }
\end{array},\right. \\
& e_{i}\left(v_{C}\right)=\left\{\begin{array}{l}
0 \text { if } \delta_{i}=1 \text { or } \delta_{i+1}=0 \\
v_{C^{\prime}} \text { with } C^{\prime}=C-\{i+1\}+\{i\} \text { otherwise }
\end{array},\right. \\
& k_{i}\left(v_{C}\right)=v^{\delta_{i}(C)-\delta_{i+1}(C) v_{C} .}
\end{aligned}
$$

where for any $i \in \mathbb{Z}, \delta_{i}(C)=1$ if $i \in C$ and $\delta_{i}(C)=0$ otherwise.
Remark: Consider $a, b$ two integers such that $a<b$. Denote by $[a, b[$ the set of integers $i$ such that $a \leq i \leq b-1$. Write $\mathcal{U}_{v}\left(\mathfrak{s l}_{a, b}\right)$ for the subalgebra of $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$ generated by the Chevalley generators $e_{i}, f_{i}, k_{i}$ with $i \in\left[a, b\left[\right.\right.$. Then $\mathcal{U}_{v}\left(\mathfrak{s l}_{a, b}\right)$ can be identified with the quantum group associated to the Dynkin diagram obtained by deleting the nodes $i \notin[a, b[$ in (11). In particular $\mathcal{U}_{v}\left(\mathfrak{s l}_{a, b}\right)$ is isomorphic to the quantum group $\mathcal{U}_{v}\left(\mathfrak{s l}_{d}\right)$ with $d=b-a+1$. The subspace $V_{a, b}(k)$ of $V_{\infty}(k)$ generated by the basis vectors $v_{C}$ such that $C$ contains only letters $x_{i}$ with $a \leq x_{i} \leq b$ has the structure of a $\mathcal{U}_{v}\left(\mathfrak{s l}_{a, b}\right)$-module. Moreover $V_{a, b}(k)$ is isomorphic to the $k^{\text {th }}$ fundamental module of $\mathcal{U}_{v}\left(\mathfrak{s l}_{a, b}\right)$.

We follow the convention of 15 and consider $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$ as a Hopf algebra with coproduct given by

$$
\begin{gathered}
\Delta\left(k_{i}\right)=k_{i} \otimes k_{i} \\
\Delta\left(e_{i}\right)=e_{i} \otimes k_{i}+1 \otimes e_{i} \text { and } \Delta\left(f_{i}\right)=f_{i} \otimes 1+k_{i}^{-1} \otimes f_{i}
\end{gathered}
$$

Given $M_{1}$ and $M_{2}$ two $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$ with crystal graphs $B_{1}$ and $B_{2}$, the crystal graph structure on $B_{1} \otimes B_{2}$ is then given by

$$
\begin{align*}
& \tilde{f}_{i}(u \otimes v)=\left\{\begin{array}{l}
\widetilde{f}_{i}(u) \otimes v \text { if } \varphi_{i}(v) \leq \varepsilon_{i}(u) \\
u \otimes \widetilde{f}_{i}(v) \text { if } \varphi_{i}(v)>\varepsilon_{i}(u)
\end{array}\right.  \tag{3}\\
& \widetilde{e}_{i}(u \otimes v)=\left\{\begin{array}{l}
u \otimes \widetilde{e}_{i}(v) \text { if } \varphi_{i}(v) \geq \varepsilon_{i}(u) \\
\widetilde{e}_{i}(u) \otimes v \text { if } \varphi_{i}(v)<\varepsilon_{i}(u)
\end{array}\right. \tag{4}
\end{align*}
$$

### 2.2 The crystals graphs of irreducible $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-modules

For any $s \in \mathbb{Z}$, the $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-module $V_{\infty}\left(\omega_{s}\right)$ can be obtained as an irreducible component of the Fock space $F_{\infty}\left(\omega_{s}\right)$ defined by considering semi-infinite wedge products of $V_{\infty}(1)$. The crystal $B_{\infty}\left(\omega_{s}\right)$ is identified with the graph whose vertices are the infinite columns

$$
\mathcal{C}=\begin{array}{|c|}
\hline x_{1} \\
\hline \cdot \\
\hline \cdot \\
\hline x_{k} \\
\hline s-k+1 \\
\hline s-k+2 \\
\hline \cdots \\
\hline
\end{array}
$$

that is, the infinite columns shaped Young diagrams filled by decreasing integers $x_{i}$ from top to bottom and such that $x_{k}=s-k+1$ for $k$ sufficiently large. Given $\mathcal{C}_{1}$ and $\mathcal{C}_{2}$ in $B_{\infty}\left(\omega_{s}\right)$, we have an arrow $\mathcal{C}_{1} \xrightarrow{i} \mathcal{C}_{2}$ if and only if $i \in \mathcal{C}_{1}, i+1 \notin \mathcal{C}_{2}$ and $\mathcal{C}_{2}=\mathcal{C}_{1}-\{i\}+\{i+1\}$. The highest weight vertex of $B_{\infty}\left(\omega_{s}\right)$ is the column $\mathcal{C}^{(s)}$ such that $x_{k}=s-k+1$ for any $k \geq 1$.
We associate to the $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-module $V_{\infty}(k)$ a graph $B_{\infty}(k)$. Its vertices are the columns of height $k$ (see (2)) and we draw an oriented arrow $C_{1} \xrightarrow{i} C_{2}$ if and only if $i \in C_{1}, i+1 \notin C_{2}$ and $C_{2}=C_{1}-\{i\}+\{i+1\}$. The graph $B_{\infty}(k)$ can be regarded as the crystal graph of $V_{\infty}(k)$. In fact, one can define a notion of crystal basis for the module $V_{\infty}(k)$ despite it is not of highest weight. This can be done essentially by considering the direct limit of the direct system formed by the crystal bases of the $\mathcal{U}_{v}\left(\mathfrak{s l}_{a, b}\right)$-modules $V_{a, b}(k)$ 18]. The corresponding crystal graph then coincide with $B_{\infty}(k)$. In the sequel we will only need the crystal $B_{\infty}(k)$ and not the whole crystal basis of $B_{\infty}(k)$.
Consider an infinite column $\mathcal{C} \in B_{\infty}\left(\omega_{s}\right)$ with letters $x_{i}, i \geq 1$. We write $\pi_{a}(\mathcal{C})$ for the finite column obtained by deleting the infinite sequence of letters $x_{i}$ such that $x_{i}<a$ in $\mathcal{C}$.
Let $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right) \in \mathbb{Z}^{\ell}$. The crystal $B_{\infty}\left(\omega_{\underline{s}}\right)$ is identified with the connected component of $B_{\infty}(\underline{s})=B_{\infty}\left(\omega_{0}\right) \otimes \cdots \otimes B_{\infty}\left(\omega_{s_{\ell-1}}\right)$ with highest weight $b^{(\underline{s})}=\mathcal{C}^{\left(s_{0}\right)} \otimes \cdots \otimes \mathcal{C}^{\left(s_{\ell-1}\right)}$. Consider
$b=\mathcal{C}_{0} \otimes \cdots \otimes \mathcal{C}_{\ell-1} \in B_{\infty}\left(\omega_{\underline{s}}\right)$ with $\mathcal{C}_{k} \in B_{\infty}\left(\omega_{s_{k}}\right), k=0, \ldots, \ell-1$. For any integer $a$ and any $k=0, \ldots, \ell-1$, let $h_{k}$ be the height of the finite column $\pi_{a}\left(\mathcal{C}_{k}\right)$. We set

$$
\pi_{a}(b)=\pi_{a}\left(\mathcal{C}_{0}\right) \otimes \cdots \otimes \pi_{a}\left(\mathcal{C}_{\ell-1}\right) \in B_{\infty}\left(h_{0}\right) \otimes \cdots \otimes B_{\infty}\left(h_{\ell-1}\right)
$$

Consider $b_{1}$ and $b_{2}$ two vertices of $B_{\infty}\left(\omega_{\underline{s}}\right)$. Let $\widetilde{K}$ be any path between $b_{1}$ and $b_{2}$ in $B_{\infty}\left(\omega_{\underline{s}}\right)$, that is any sequence of crystal operators such that $b_{1}=\widetilde{K}\left(b_{2}\right)$.

Lemma 2.2.1 With the previous notation, for any integer $a \geq 1$ sufficiently large, we have $\pi_{a}\left(b_{1}\right) \in \underset{\sim}{B}\left(h_{0}\right) \otimes \cdots \otimes B_{\infty}\left(h_{\ell-1}\right)$ and $\pi_{a}\left(b_{2}\right) \in B_{\infty}\left(h_{0}\right) \otimes \cdots \otimes B_{\infty}\left(h_{\ell-1}\right)$. In this case $\pi_{a}\left(b_{1}\right)=\widetilde{K}\left(\pi_{a}\left(b_{2}\right)\right)$ in $B_{\infty}\left(h_{0}\right) \otimes \cdots \otimes B_{\infty}\left(h_{\ell-1}\right)$.

Proof. We can choose $a$ sufficiently large so that the infinite columns appearing in any vertex $b$ of the path joining $b_{1}$ to $b_{2}$ contain all the letters $x<a$. Then, for each $k=0, \ldots, \ell-1$, the height of the column obtained by deleting the letters $x<a$ in the $k$-th column of $b$ does not depend on $b$. Denote it $h_{k}$. The deleted letters $x<a$ do not interfere during the computation of the crystal operators defining the path between $b_{1}$ and $b_{2}$. This implies that the path joining $\pi_{a}\left(b_{1}\right)$ to $\pi_{a}\left(b_{2}\right)$ obtained by applying the same sequence $\widetilde{K}$ of crystal operators also exists in $B_{\infty}\left(h_{0}\right) \otimes \cdots \otimes B_{\infty}\left(h_{\ell-1}\right)$. Thus $\pi_{a}\left(b_{1}\right)=\widetilde{K}\left(\pi_{a}\left(b_{2}\right)\right)$.

The vertices of $B_{\infty}\left(\omega_{s}\right)$ are also naturally labelled by partitions. Recall that a partition $\lambda=\left(\lambda_{1}, \ldots, \lambda_{p}\right)$ of length $p$ is a decreasing sequence $\lambda_{1} \geq \cdots \geq \lambda_{p}$ of nonnegative integers called the parts of $\lambda$. In the sequel we identify the partitions having the same nonzero parts and denote by $\mathcal{P}$ the set of all partitions. It is convenient to represent $\lambda$ by its Young diagram. In the sequel we use the french convention for the Young diagram (see example below). To any pair ( $a, b$ ) of integers, we associate the node $\gamma=(a, b)$ which is the box obtained after $a$ vertical moves and $b$ horizontal moves starting from the south-west box of $\lambda$. Observe that the node $\gamma$ does not necessarily belong to $\lambda$. We say that $\gamma$ is removable when $\gamma=(a, b) \in \lambda$ and $\lambda \backslash\{\gamma\}$ is a partition. Similarly $\gamma$ is said addable when $\gamma=(a, b) \notin \lambda$ and $\lambda \cup\{\gamma\}$ is a partition.
We associate to the infinite column $\mathcal{C} \in B_{\infty}\left(\omega_{s}\right)$ with letters $x_{k}$ the partition $\lambda$ such that $\lambda_{k}=$ $x_{k}-s+k-1$. Since $\lambda_{k}=0$ for $k$ sufficiently large, this permits to label $B_{\infty}\left(\omega_{s}\right)$ by $\mathcal{P}$. Let $\gamma=(a, b)$ be a node for $\lambda \in B_{\infty}\left(\omega_{s}\right)$. The contain of $\gamma$ is

$$
\begin{equation*}
c(\gamma)=b-a+s \tag{5}
\end{equation*}
$$

In $B_{\infty}\left(\omega_{s}\right)$, we have an arrow $\lambda \xrightarrow{i} \mu$ if and only if $\mu \backslash \lambda=\gamma$ where $\gamma$ is addable in $\lambda$ with $c(\gamma)=i$. For any partitions $\lambda=\left(\lambda_{1}, \ldots, \lambda_{p}\right)$ we set $|\lambda|=\lambda_{1}+\cdots+\lambda_{p}$. Then $|\lambda|$ is equal to the length of the path joining $\emptyset$ to $\lambda$ in the crystal $B_{\infty}\left(\omega_{s}\right)$.

## Example 2.2.2


for the node $\gamma=(3,4)$ is addable in $\lambda$ with contain $c(\gamma)=4-3+3=4$.
Since $B_{\infty}\left(\omega_{\underline{s}}\right)$ is labelled by the partitions, the crystal $B_{\infty}(\underline{s})$ is labelled by multipartitions $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$. More precisely, to each vertex $b=\mathcal{C}_{0} \otimes \cdots \otimes \mathcal{C}_{\ell-1} \in B_{\infty}(\underline{s})$ such that $\mathcal{C}_{k} \in B\left(\omega_{s_{k}}\right)$ for any $k=0, \ldots, \ell-1$, corresponds the multipartition $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$ where $\lambda^{(k)}$ is obtained from $\mathcal{C}_{k}$ as described above. The action of the operators $\widetilde{e}_{i}$ and $\widetilde{f}_{i}$ on $b$ are deduced from (3) and (4). One verifies easily that they can be also obtained by applying the following algorithm. Let $w_{i}$ be the word obtained by reading the letters $i$ or $i+1$ successively in the columns $\mathcal{C}_{0}, \mathcal{C}_{1}, \ldots, \mathcal{C}_{\ell-1}$ of $b$. Encode in $w_{i}$ each integer $i$ by a symbol + and each integer $i+1$ by a symbol - . Choose any factor -+ and delete it. Repeat this procedure until no factor -+ can be deleted. The final sequence $\widetilde{w}_{i}$ is uniquely determined and has the form $\widetilde{w}_{i}=+^{p}{ }_{-}^{q}$. Then $\widetilde{f}_{i}(b)$ is obtained from $b$ by replacing the integer $i$ corresponding to the rightmost symbol + in $\widetilde{w}_{i}$ by $i+1$. Similarly, $\widetilde{e}_{i}(b)$ is obtained from $\lambda$ by replacing the integer $i+1$ corresponding to the leftmost symbol + in $\widetilde{w}_{i}$ by $i$.
One can describe similarly the action of $\widetilde{e}_{i}$ and $\widetilde{f}_{i}$ on $b$ considered as the multipartition $\underline{\lambda}=$ $\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$. This time, one considers the word $W_{i}$ obtained by reading the addable and removable nodes of $\underline{\lambda}$ with contain $i$ successively in the partitions $\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}$. This reading is well defined since a partition cannot both contain addable and removable nodes with contain $i$. Encode in $W_{i}$ each removable node by $R$ and each addable node by $A$. Let $\widetilde{W}_{i}$ be the word obtained by deleting the factors $R A$ successively in the encoding. One can write

$$
\begin{equation*}
\widetilde{W}_{i}=A^{p} R^{q} \tag{6}
\end{equation*}
$$

Then $\widetilde{f}_{i}(b)$ is obtained from $\underline{\lambda}$ by adding the node $\gamma$ corresponding to the rightmost symbol $A$ in $\widetilde{W}_{i}$.

### 2.3 The crystal isomorphism between $B_{\infty}\left(\omega_{k}\right) \otimes B_{\infty}\left(\omega_{l}\right)$ and $B_{\infty}\left(\omega_{l}\right) \otimes B_{\infty}\left(\omega_{k}\right)$

Consider $k$ and $l$ two integers. We denote by $\psi_{k, l}$ the crystal graph isomorphism

$$
\begin{equation*}
\psi_{k, l}: B_{\infty}\left(\omega_{k}\right) \otimes B_{\infty}\left(\omega_{l}\right) \stackrel{\cong}{\rightrightarrows} B_{\infty}\left(\omega_{l}\right) \otimes B_{\infty}\left(\omega_{k}\right) . \tag{7}
\end{equation*}
$$

To give the explicit combinatorial description of $\psi_{k, l}$, we start by considering the crystal graph isomorphism $\theta_{k, l}$ between $B_{\infty}(k) \otimes B_{\infty}(l)$ and $B_{\infty}(l) \otimes B_{\infty}(k)$. Consider $C_{1} \otimes C_{2}$ in $B_{\infty}(k) \otimes B_{\infty}(l)$. We are going to associate to $C_{1} \otimes C_{2}$ a vertex $C_{2}^{\prime} \otimes C_{1}^{\prime} \in B_{\infty}(l) \otimes B_{\infty}(k)$.
Suppose first $k \geq l$. Consider $x_{1}=\min \left\{t \in C_{2}\right\}$. We associate to $x_{1}$ the integer $y_{1} \in C_{1}$ such that

$$
y_{1}=\left\{\begin{array}{l}
\max \left\{z \in C_{1} \mid z \leq x_{1}\right\} \text { if } \min \left\{z \in C_{1}\right\} \leq x  \tag{8}\\
\max \left\{z \in C_{1}\right\} \text { otherwise }
\end{array} .\right.
$$

We repeat the same procedure to the columns $C_{1}-\left\{y_{1}\right\}$ and $C_{2}-\left\{x_{1}\right\}$. By induction this yields a sequence $\left\{y_{1}, \ldots, y_{l}\right\} \subset C_{1}$. Then we define $C_{2}^{\prime}$ as the column obtained by reordering the integers of $\left\{y_{1}, \ldots, y_{l}\right\}$ and $C_{1}^{\prime}$ as the column obtained by reordering the integers of $C_{1}-\left\{y_{1}, \ldots, y_{l}\right\}+C_{2}$. Now, suppose $k<l$. Consider $x_{1}=\min \left\{t \in C_{2}\right\}$. We associate to $x_{1}$ the integer $y_{1} \in C_{1}$ such that

$$
y_{1}=\left\{\begin{array}{l}
\min \left\{z \in C_{1} \mid x_{1} \leq z\right\} \text { if } \max \left\{z \in C_{1}\right\} \geq x  \tag{9}\\
\min \left\{z \in C_{1}\right\} \text { otherwise }
\end{array} .\right.
$$

We repeat the same procedure to the columns $C_{1}-\left\{y_{1}\right\}$ and $C_{2}-\left\{x_{1}\right\}$ and obtain a sequence $\left\{y_{1}, \ldots, y_{l}\right\} \subset C_{1}$. Then we define $C_{2}^{\prime}$ and $C_{1}^{\prime}$ as in the previous case.
We denote by $\theta_{k, l}$ the map defined from $B_{\infty}(k) \otimes B_{\infty}(l)$ to $B_{\infty}(l) \otimes B_{\infty}(k)$ by setting $\theta_{k, l}\left(C_{1} \otimes\right.$ $\left.C_{2}\right)=C_{2}^{\prime} \otimes C_{1}^{\prime}$ 。

Example 2.3.1 Consider $C_{1}=$\begin{tabular}{|c|}
\hline 9 <br>
\hline 8 <br>
\hline 7 <br>
\hline 7 <br>
\hline 5 <br>
\hline 4 <br>
\hline 2 <br>
\hline

 and $C_{2}=$

\hline 7 <br>
\hline 6 <br>
\hline 6 <br>
\hline 5 <br>
\hline 3 <br>
\hline 1. <br>
\hline 1. <br>
\hline

. We obtain $\left\{y_{1}, y_{2}, y_{3}, y_{4}, y_{5}\right\}=\{9,2,4,5,7\}$. This gives $C_{2}^{\prime}=$

\hline 9 <br>
\hline 7 <br>
\hline 5 <br>
\hline 4 <br>
\hline 2 <br>
\hline 2 <br>
\hline

 and $C_{1}^{\prime}=$

\hline 8 <br>
\hline 7 <br>
\hline 6 <br>
\hline 5 <br>
\hline \hline 3 <br>
\hline 1 <br>
\hline
\end{tabular}. . One can then verify that $\theta_{5,6}\left(C_{2}^{\prime} \otimes C_{1}^{\prime}\right)=C_{1} \otimes C_{2}$.

Proposition 2.3.2 The map $\theta_{k, l}$ is an isomorphism of $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-crystals that is, for any integer $i$ and any vertex $C_{1} \otimes C_{2} \in B_{\infty}(k) \otimes B_{\infty}(l)$, we have

$$
\begin{equation*}
\theta_{k, l} \circ \widetilde{e}_{i}\left(C_{1} \otimes C_{2}\right)=\widetilde{e}_{i}\left(C_{2}^{\prime} \otimes C_{1}^{\prime}\right) \text { and } \theta_{k, l} \circ \widetilde{f}_{i}\left(C_{1} \otimes C_{2}\right)=\widetilde{f}_{i}\left(C_{2}^{\prime} \otimes C_{1}^{\prime}\right) . \tag{10}
\end{equation*}
$$

Proof. Choose $a<b$ two integers such that the letters of $C_{1}$ and $C_{2}$ belong to $[a, b]$. It follows from Proposition 3.21 of [20], that the isomorphism between the finite $\mathcal{U}_{v}\left(\mathfrak{s l}_{a-1, b+2}\right)$ crystals $B_{a-1, b+2}(k) \otimes B_{a-1, b+2}(l)$ and $B_{a-1, b+2}(l) \otimes B_{a-1, b+2}(k)$ is given by the map $\theta_{k, l}$. Since the actions of the crystal operators $\widetilde{e}_{i}$ and $\widetilde{f}_{i}$ with $i \in\left[a-1, b+2\left[\right.\right.$ on the crystals $B_{a-1, b+2}(l)$ and $\left.B_{a-1, b+2}(k)\right)$ can be obtained from the crystal structures of $B_{\infty}(k)$ and $B_{\infty}(k)$, this implies the commutation relations (10) for any $i \in[a-1, b+2[$. Now if $i \notin[a-1, b+2[$ we have $\widetilde{e}_{i}\left(C_{1} \otimes C_{2}\right)=\widetilde{f}_{i}\left(C_{1} \otimes C_{2}\right)=0$ because the letters of $C_{1}$ and $C_{2}$ belong to [a,b]. Similarly $\tilde{e}_{i}\left(C_{2}^{\prime} \otimes C_{1}^{\prime}\right)=C_{2}^{\prime} \otimes C_{1}^{\prime}=0$ because the letters of $C_{1} \cup C_{2}$ are the same as those of $C_{1}^{\prime} \cup C_{2}^{\prime}$. Hence (10) holds for any integer $i$.

Now consider $\mathcal{C}_{1} \otimes \mathcal{C}_{2} \in B_{\infty}\left(\omega_{k}\right) \otimes B_{\infty}\left(\omega_{l}\right)$. Let $a$ be any integer such that $\mathcal{C}_{1}$ and $\mathcal{C}_{2}$ both contain all the integers $x<a$. Set $C_{1}=\pi_{a}\left(\mathcal{C}_{1}\right), C_{2}=\pi_{a}\left(\mathcal{C}_{2}\right)$ and $\theta_{k, l}\left(C_{1} \otimes C_{2}\right)=C_{2}^{\prime} \otimes C_{1}^{\prime}$. Since $C_{1}$ and $C_{2}$ both contain only letters $x \geq a$, it is also the case for the columns $C_{1}^{\prime}$ and $C_{2}^{\prime}$ by the previous combinatorial description of the isomorphism $\theta_{k, l}$. Write $\mathcal{C}_{1}^{\prime}$ and $\mathcal{C}_{2}^{\prime}$ for the infinite columns obtained respectively from $C_{1}^{\prime}$ and $C_{2}^{\prime}$ by adding boxes containing all the letters $x<a$. Then $\mathcal{C}_{1}^{\prime} \in B_{\infty}\left(\omega_{k}\right)$ and $\mathcal{C}_{2}^{\prime} \in B_{\infty}\left(\omega_{l}\right)$. Indeed $\mathcal{C}_{1} \in B_{\infty}\left(\omega_{k}\right), \mathcal{C}_{2} \in B_{\infty}\left(\omega_{l}\right)$ and $C_{1}, C_{1}^{\prime}$ (resp. $\left.C_{2}, C_{2}^{\prime}\right)$ have the same height.

Corollary 2.3.3 (of Proposition 2.3.2) For any $\mathcal{C}_{1} \otimes \mathcal{C}_{2} \in B_{\infty}\left(\omega_{k}\right) \otimes B_{\infty}\left(\omega_{l}\right)$ we have with the above notation $\psi_{k, l}\left(\mathcal{C}_{1} \otimes \mathcal{C}_{2}\right)=\mathcal{C}_{2}^{\prime} \otimes \mathcal{C}_{1}^{\prime}$.

Proof. Recall that $\mathcal{C}^{(k)} \otimes \mathcal{C}^{(l)}$ is the highest weight vertex of $B_{\infty}\left(\omega_{k}\right) \otimes B_{\infty}\left(\omega_{l}\right)$. Write $\mathcal{C}_{1} \otimes \mathcal{C}_{2}=\widetilde{F}\left(\mathcal{C}^{(k)} \otimes \mathcal{C}^{(l)}\right)$ where $\widetilde{F}$ is a sequence of crystal operators $\widetilde{f}_{i}, i \in \mathbb{Z}$. The crystal graph isomorphism $\psi_{k, l}$ must send the highest vertex of $B_{\infty}\left(\omega_{k}\right) \otimes B_{\infty}\left(\omega_{l}\right)$ on the highest weight vertex of $B_{\infty}\left(\omega_{l}\right) \otimes B_{\infty}\left(\omega_{k}\right)$. Thus, we have $\psi_{k, l}\left(\mathcal{C}^{(k)} \otimes \mathcal{C}^{(l)}\right)=\mathcal{C}^{(l)} \otimes \mathcal{C}^{(k)}$. In particular, the Corollary is true for $\mathcal{C}_{1} \otimes \mathcal{C}_{2}=\mathcal{C}^{(k)} \otimes \mathcal{C}^{(l)}$.
Now consider $\mathcal{C}_{1} \otimes \mathcal{C}_{2} \in B_{\infty}\left(\omega_{k}\right) \otimes B_{\infty}\left(\omega_{l}\right)$ and choose $a$ an integer such that $\mathcal{C}_{1}$ and $\mathcal{C}_{2}$ both contain all the integers $x<a$. Set $C^{(k)}=\pi_{a}\left(\mathcal{C}^{(k)}\right)$ and $C^{(l)}=\pi_{a}\left(\mathcal{C}^{(l)}\right)$. Similarly write $C_{1}=$ $\pi_{a}\left(\mathcal{C}_{1}\right), C_{2}=\pi_{a}\left(\mathcal{C}_{2}\right)$. Since $\mathcal{C}_{1} \otimes \mathcal{C}_{2}=\widetilde{F}\left(\mathcal{C}^{(k)} \otimes \mathcal{C}^{(l)}\right)$ we obtain from Lemma 2.2.1 the equality $C_{1} \otimes C_{2}=\widetilde{F}\left(C^{(k)} \otimes C^{(l)}\right)$. By definition of the crystal isomorphism $\theta_{k, l}$, we thus have $C_{2}^{\prime} \otimes C_{1}^{\prime}=$ $\widetilde{F}\left(C^{(l)} \otimes C^{(k)}\right)$. The columns $C^{(k)}, C^{(l)}, C_{1}^{\prime}$ and $C_{2}^{\prime}$ contains only letters $x \geq a$. Hence we can write $\widetilde{F}=\widetilde{f_{i_{1}}} \cdots \widetilde{f}_{i_{r}}$ with $i_{m} \geq a$ for any $m \in\{1, \ldots, r\}$. Now the infinite columns $\mathcal{C}^{(l)}, \mathcal{C}^{(k)}, \mathcal{C}_{1}^{\prime}$ and $\mathcal{C}_{2}^{\prime}$ are respectively obtained from $C^{(k)}, C^{(l)}, C_{1}^{\prime}$ and $C_{2}^{\prime}$ by adding all the letters $x<a$. We can thus deduce from the equality $C_{2}^{\prime} \otimes C_{1}^{\prime}=\widetilde{F}\left(C^{(l)} \otimes C^{(k)}\right)$, the equality $\mathcal{C}_{2}^{\prime} \otimes \mathcal{C}_{1}^{\prime}=\widetilde{F}\left(\mathcal{C}^{(l)} \otimes \mathcal{C}^{(k)}\right)$. This implies that $\psi_{k, l}\left(\mathcal{C}_{1} \otimes \mathcal{C}_{2}\right)=\mathcal{C}_{2}^{\prime} \otimes \mathcal{C}_{1}^{\prime}$.

Example 2.3.4 Consider $\lambda^{(1)}=(5,5,5,4,4,3) \in B_{\infty}\left(\omega_{4}\right)$ and $\lambda^{(2)}=(4,4,4,3,2) \in B_{\infty}\left(\omega_{3}\right)$. The
columns $\mathcal{C}_{1}, \mathcal{C}_{2}$ such that $\mathcal{C}_{1} \otimes \mathcal{C}_{2}=\left(\lambda^{(1)}, \lambda^{(2)}\right) \in B_{\infty}\left(\omega_{4}\right) \otimes B_{\infty}\left(\omega_{3}\right)$ are

We deduce from Example 2.3 .1 that

$$
\mathcal{C}_{1}^{\prime}=\begin{array}{|c|}
\hline \begin{array}{|c|}
\hline 8 \\
\hline 7 \\
\hline 6 \\
\hline 5 \\
\hline 3 \\
\hline 1 \\
\hline \overline{2} \\
\hline \cdots \\
\hline
\end{array} \text { and } \mathcal{C}_{2}^{\prime}=\begin{array}{|c|}
\hline 9 \\
\hline 7 \\
\hline 5 \\
\hline \hline 4 \\
\hline 2 \\
\hline \overline{2} \\
\hline \cdots \\
\hline
\end{array} . \\
\hline
\end{array}
$$

Thus we can write $\psi_{4,3}\left(\lambda^{(1)}, \lambda^{(2)}\right)=\left(\mu^{(2)}, \mu^{(1)}\right)$ with $\mu^{(1)}=(4,4,4,4,3,2)$ and $\mu^{(2)}=(6,5,4,4,3)$.
Remark: The columns $C_{1}$ and $C_{2}$ obtained in the previous algorithm depend on the integer $a$ considered. Nevertheless, this is not the case of the resulting infinite columns $\mathcal{C}_{1}^{\prime}$ and $\mathcal{C}_{2}^{\prime}$.

## 3 Crystals in type $A_{e-1}^{(1)}$

We now turn to the problem of studying the crystals in type $A_{e-1}^{(1)}$. In this section, we recall and show basic facts on their combinatorial descriptions.

### 3.1 Background on $\mathcal{U}_{v}\left(\widehat{\mathfrak{s l}}_{e}\right)$

In order to link the different labellings of the crystal graphs in type $A_{e-1}^{(1)}$ appearing in the literature, we shall need the two following Dynkin diagrams

Let $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s l}}_{e}\right)$ and $\mathcal{U}_{v}^{-}\left(\widehat{\mathfrak{s l}}_{e}\right)$ be the affine quantum groups defined respectively from the root systems in type $A_{e-1}^{(1),+}$ and $A_{e-1}^{(1),-}$ (see for example [19], Chapter 6). Write $\left\{\Lambda_{0}^{+}, \Lambda_{1}^{+}, \ldots, \Lambda_{e-1}^{+}\right\}$and $\left\{\Lambda_{0}^{-}, \Lambda_{1}^{-}, \ldots, \Lambda_{e-1}^{-}\right\}$for the dominant weights of the root systems $A_{e-1}^{(1),+}$ and $A_{e-1}^{(1),-}$. We have
then $\Lambda_{k}^{+}=\Lambda_{e-k}^{-}$. Write $\left\{E_{i}^{+}, F_{i}^{+}, K_{i}^{+} \mid i=0, \ldots, e-1\right\}$ and $\left\{E_{i}^{-}, F_{i}^{-}, K_{i}^{-} \mid i=0, \ldots, e-1\right\}$ for the sets of Chevalley generators respectively in $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s l}}_{e}\right)$ and $\mathcal{U}_{v}^{-}\left(\widehat{\mathfrak{s l}}_{e}\right)$. Let $v^{d^{+}} \in \mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s l}}_{e}\right)$ and $v^{d^{-}} \in \mathcal{U}_{v}^{-}\left(\widehat{\mathfrak{s l}}_{e}\right)$ be the quantum derivation operators. Then the map

$$
\iota:\left\{\begin{array}{l}
\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s l}}_{e}\right) \rightarrow \mathcal{U}_{v}^{-}\left(\widehat{\mathfrak{s}}_{e}\right)  \tag{12}\\
v^{d^{+}} \longmapsto v^{d-} \\
E_{i}^{+} \longmapsto E_{e-i}^{-}, F_{i}^{+} \longmapsto F_{e-i}^{-} \text {and } K_{i}^{+} \longmapsto K_{e-i}^{-}
\end{array}\right.
$$

is an isomorphism of algebras.
We associate to $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right) \in \mathbb{Z}^{\ell}$ the dominants weights $\Lambda_{\underline{\mathbf{s}}}^{+}=\sum_{k=0}^{\ell-1} \Lambda_{s_{k}(\bmod e)}^{+}$and $\Lambda_{\underline{\mathbf{s}}}^{-}=$ $\sum_{k=0}^{\ell-1} \Lambda_{s_{k}(\bmod e)}^{-}$. Let $V_{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$(resp. $V_{e}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$) be the irreducible $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s l}}_{e}\right)$-module (resp. $\mathcal{U}_{v}^{-}\left(\widehat{\mathfrak{s}}_{e}\right)-$ module) of highest weight $\Lambda_{\underline{\mathbf{s}}}^{+}$(resp. $\Lambda_{\underline{\mathbf{s}}}^{-}$). These modules can be constructed by using the Fock space representation $\mathfrak{F}^{\underline{s}}$ of level $\ell$. Let $\Pi_{\ell, n}$ be the set of multipartitions $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$ of length $\ell$ with rank $n$, i.e. such that $\left|\lambda^{(0)}\right|+\cdots\left|\lambda^{(\ell-1)}\right|=n$. The Fock space $\mathfrak{F}^{\frac{s}{e}}$ is defined as the $\mathbb{C}(v)$-vector space generated by the symbols $|\underline{\lambda}, s\rangle$ with $\underline{\lambda} \in \Pi_{\ell, n}$

$$
\mathfrak{F} \frac{s}{e}=\bigoplus_{n \geq 0} \bigoplus_{\underline{\lambda} \in \Pi_{\ell, n}} \mathbb{C}(v)|\underline{\lambda}, s\rangle
$$

The Fock space $\mathfrak{F}^{\frac{s}{e}}$ can be endowed with the structure of a $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s l}}_{e}\right)$-module or equivalently with the structure of a $\mathcal{U}_{v}^{-}\left(\widehat{\mathfrak{s l}}_{e}\right)$-module. These actions are defined by introducing total orders $\prec_{\underline{s}}^{+}$and $\prec_{\underline{s}}^{-}$on the $i$-nodes of the multipartitions. Consider $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$ a multipartition and suppose $\underline{s}$ (called the multicharge) fixed. Then the nodes of $\underline{\lambda}$ can be identified with the triplet $\gamma=(a, b, c)$ where $c \in\{0, \ldots, \ell-1\}$ and $a, b$ are respectively the row and column indices of the node $\gamma$ in $\lambda^{(c)}$. The contain of $\gamma$ is the integer $c(\gamma)=b-a+c$ and the residue res $(\gamma)$ of $\gamma$ is the element of $\mathbb{Z} / e \mathbb{Z}$ such that

$$
\begin{equation*}
\operatorname{res}(\gamma) \equiv c(\gamma)(\bmod e) \tag{13}
\end{equation*}
$$

We will say that $\gamma$ is an $i$-node of $\underline{\lambda}$ when $\operatorname{res}(\gamma) \equiv i(\bmod e)$. Let $\gamma_{1}=\left(a_{1}, b_{1}, c_{1}\right)$ and $\gamma_{2}=$ $\left(a_{2}, b_{2}, c_{2}\right)$ be two $i$-nodes in $\underline{\lambda}$. We define the total order $\prec_{\underline{s}}^{+}$and $\prec_{\underline{s}}^{-}$on the $i$-nodes of $\underline{\lambda}$ by setting

$$
\begin{align*}
& \gamma_{1} \prec_{\underline{s}}^{+} \gamma_{2} \Longleftrightarrow\left\{\begin{array}{l}
c\left(\gamma_{1}\right)<c\left(\gamma_{2}\right) \text { or } \\
c\left(\gamma_{1}\right)=c\left(\gamma_{2}\right) \text { and } c_{1}<c_{2}
\end{array}\right.  \tag{14}\\
& \gamma_{1} \prec_{\underline{s}}^{-} \gamma_{2} \Longleftrightarrow\left\{\begin{array}{l}
c\left(\gamma_{1}\right)<c\left(\gamma_{2}\right) \text { or } \\
c\left(\gamma_{1}\right)=c\left(\gamma_{2}\right) \text { and } c_{1}>c_{2}
\end{array}\right. \tag{15}
\end{align*}
$$

Using these orders, it is possible to define an action of $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s l}}_{e}\right)$-module and an action of $\mathcal{U}_{v}^{-}\left(\widehat{\mathfrak{s l}}_{e}\right)-$ module on $\mathfrak{F}_{e}^{\underline{s}}$. These modules will be denoted by $\mathfrak{F}_{e}^{\underline{s},+}$ and $\mathfrak{F}_{e}^{\underline{s},-}$. For these actions the empty multipartition $\underline{\emptyset}=(\emptyset, \ldots, \emptyset)$ is a highest weight vector respectively of highest weight $\Lambda_{\underline{\mathbf{s}}}^{+}$and $\Lambda_{\underline{\mathbf{s}}}^{-}$. We denote by $V_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$and $V_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$the irreducible component of highest weight $\underline{\emptyset}$ in $\mathfrak{F} e$ and $\mathfrak{F}_{e}^{\underline{s},-}$. We refer the reader to [12] for a detailed exposition. Note that the modules $V_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$ (resp. $\left.V_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)\right)$such that $\underline{s} \in \underline{\mathbf{s}}$ are all isomorphic to the abstract module $V_{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)\left(\operatorname{resp} . V_{e}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)\right.$). However, the corresponding actions of the Chevalley operators do not coincide in general. This will yield different parametrizations of the associated crystals.

### 3.2 Crystal graph of the Fock space $\mathfrak{F}_{e}^{s,+}$ and $\mathfrak{F}_{e}^{s,-}$

The modules $\mathfrak{F}_{e}^{s,+}$ and $\mathfrak{F}_{e}^{s,-}$ are integrable modules, thus by the general theory of Kashiwara, they admit crystal bases. Write $B_{e}^{s,+}$ and $B_{e}^{s,-}$ for the crystal graphs corresponding to the action of $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s}}_{e}\right)$ and $\mathcal{U}_{v}^{-}\left(\widehat{\mathfrak{s}}_{e}\right)$ on $\mathfrak{F}_{e}^{s}$. These crystals are labelled by multipartitions and their crystal graph structure can be explicitly described by using the total orders $\prec_{\underline{s}}^{+}$and $\prec_{\underline{s}}^{-}$. Consider two multipartitions $\underline{\lambda}, \underline{\mu}$ and an integer $i \in\{0, \ldots, e-1\}$. The crystal graph $\bar{B} \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)\left(\right.$resp. $\left.B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)\right)$of $V_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$(resp. $\left.V_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)\right)$is the connected component of $B_{e}^{s,+}$ (resp. $B_{e}^{\underline{s},-}$ ) whose highest weight vertex is the empty multipartition.

### 3.2.1 Crystal structure on $B_{e}^{s,+}$

Consider the set of addable and removable $i$-nodes of $\underline{\lambda}$ (see Section 2.2 for the definition of removable and addable nodes). Let $w_{i}$ be the word obtained by writing the $i$-nodes of $\underline{\lambda}$ in decreasing order with respect to $\prec_{\underline{s}}^{+}$, next by encoding each addable $i$-node by the letter $A$ and each removable $i$-node by the letter $R$. Write $\widetilde{w}_{i}=A^{p} R^{q}$ for the word obtained from $w_{i}$ by deleting the factors $R A$ as many as possible. If $p>0$, let $\gamma$ be the rightmost addable $i$-node in $\widetilde{w}_{i}$. The node $\gamma$ is called the good $i$-node. Then we have an arrow $\underline{\lambda} \xrightarrow{i} \underline{\mu}$ in $B_{e}^{s,+}$ if and only if $\underline{\mu}$ is obtained from $\underline{\lambda}$ by adding the good $i$-node $\gamma$.

### 3.2.2 Crystal structure on $B_{e}^{s,-}$

Consider the set of addable and removable $i$-nodes of $\underline{\lambda}$. Let $w_{i}$ be the word obtained by writing the $i$-nodes of $\underline{\lambda}$ in increasing order with respect to $\prec_{\underline{s}}^{-}$, next by encoding each addable $i$-node by the letter $A$ and each removable $i$-node by the letter $R$. Write $\widetilde{w}_{i}=A^{p} R^{q}$ for the word obtained from $w_{i}$ by deleting the factors $R A$ as many as possible. If $p>0$, let $\gamma$ be the rightmost addable $i$-node in $\widetilde{w}_{i}$. Then we have an arrow $\underline{\lambda} \xrightarrow{i} \underline{\mu}$ in $B_{e}^{s,-}$ if and only if $\underline{\mu}$ is obtained from $\underline{\lambda}$ by adding the good $i$-node $\gamma$.

### 3.2.3 Link between the crystals $B_{e}^{s,+}$ and $B_{e}^{s,-}$

For any multicharge $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right)$, we denote by $\underline{s}^{*}$ the multicharge $\underline{s}^{*}=\left(e-s_{0}, \ldots, e-\right.$ $s_{\ell-1}$ ). According to the isomorphism (12), there should exist a bijection $\xi: B_{e}^{s,+} \rightarrow B_{e}^{s^{*},-}$ such that, given any multipartitions $\underline{\lambda}, \underline{\mu}$,

$$
\xrightarrow{\boldsymbol{\lambda}} \xrightarrow{i} \underline{\mu} \text { in } B B_{e}^{s,+} \Longleftrightarrow \xi(\underline{\lambda}) \stackrel{e-i}{\rightarrow} \xi(\underline{\mu}) \text { in } B \stackrel{s}{e}_{s^{*},-} .
$$

Proposition 3.2.1 Consider a multipartition $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$ in $B_{e}^{s,+}$. Then we have $\xi(\underline{\lambda})=\left(\lambda^{\prime(0)}, \ldots, \lambda^{\prime(\ell-1)}\right)$ where for any $k=0, \ldots, \ell-1, \lambda^{\prime(k)}$ is the conjugate partition of $\lambda^{(k)}$.

Proof. Set $\underline{\lambda}^{\prime}=\left(\lambda^{\prime(0)}, \ldots, \lambda^{(\ell-1)}\right)$. Consider $\gamma$ a node appearing at the right end of the $a$-th row of $\lambda^{(c)}$ in $\underline{\lambda}$. One associates to $\gamma$ the node $\gamma^{\prime}$ appearing on the top of the $b$-th column of $\lambda^{\prime(k)}$ in $\underline{\lambda}^{\prime}$. Now, consider $\underline{\lambda}$ and $\underline{\lambda}^{\prime}$ as vertices respectively of the crystals $\mathfrak{F}_{e}^{s,+}$ and $\mathfrak{F} e_{e}^{s^{*},-}$. Let us
write $\gamma=(a, b, c)$. Then by definition of the node $\gamma^{\prime}$, we have $\gamma^{\prime}=(b, a, e-c)$. This gives for the contain of the nodes $\gamma$ and $\gamma^{\prime}$

$$
c(\gamma)=b-a+c \text { and } c\left(\gamma^{\prime}\right)=a-b+e-c .
$$

Now observe that $\gamma$ is a removable (resp. addable) node if and only if $\gamma^{\prime}$ is a removable (resp. addable) node. We have thus

$$
\begin{equation*}
\gamma \text { is an } A(\text { resp. } R) i \text {-node } \Longleftrightarrow \gamma^{\prime} \text { is an } A(\text { resp. } R)(e-i) \text {-node. } \tag{16}
\end{equation*}
$$

To prove our proposition, we proceed by induction on $|\underline{\lambda}|$. When $|\underline{\lambda}|=0$, we have $\underline{\lambda}=\underline{\emptyset}=\underline{\lambda}^{\prime}$, hence our statement holds. Suppose the statement true for all the multipartitions $\underline{\lambda}$ with $|\underline{\lambda}| \leq n$ where $n$ is a positive integer. Consider $\underline{\mu} \in B_{e}^{s,+}$ with $|\underline{\mu}|=n+1$. One can find $\underline{\lambda} \in B_{e}^{s,+}$ and $i \in\{0, \ldots, e-1\}$ such that $\widetilde{f}_{i}(\underline{\lambda})=\underline{\mu}$. Let $w_{i}$ be the word obtained by writing the addable or removable $i$-nodes of $\underline{\lambda}$ in decreasing ${ }^{-}$order with respect to $\prec_{\underline{s}}^{+}$as in Section 3.2.1. Similarly, let $w_{e-i}^{\prime}$ be the word obtained by writing the addable or removable ( $e-i$ )-nodes of $\underline{\lambda}^{\prime}$ in increasing order with respect to $\prec_{\underline{s}}^{+}$as in Section 3.2.2. Write $w_{i}=\gamma_{1} \cdots \gamma_{r}$ where for any $m=1, \ldots, r, \gamma_{m}$ is an $i$-node addable or removable. Then by definition of the order $\prec_{\underline{s}}^{+}$and $\prec_{\underline{s}}^{-}$(see (14) and (15)), we have

$$
w_{e-i}^{\prime}=\gamma_{1}^{\prime} \cdots \gamma_{r}^{\prime}
$$

Write $\widetilde{w}_{i}=A^{p} R^{q}$ for the word obtained from $w_{i}$ by the cancellation process of the pairs $R A$ . We deduce from (16) that the word $\widetilde{w}_{e-i}^{\prime}=\left(A^{\prime}\right)^{p}\left(R^{\prime}\right)^{q}$ coincide with the word obtained from $w_{e-i}^{\prime}$ by the cancellation process of the letters $R A$. In particular, $\gamma$ is the good $i$-node for $\underline{\lambda}$ if and only if $\gamma^{\prime}$ is the good $(e-i)$-node for $\underline{\lambda}^{\prime}$. This implies that $\tilde{f}_{e-i}\left(\underline{\lambda}^{\prime}\right)=\underline{\mu}^{\prime}$ and suffices to show our proposition by induction.

Since $V_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$and $V_{e}^{\underline{s}^{*}}\left(\Lambda_{\underline{\mathbf{s}}^{*}}^{-}\right)$are the irreducible components of highest weight vector $\underline{\emptyset}$ in $\mathfrak{F}_{e}^{s,+}$ and $\mathfrak{F}_{e}^{\underline{s},-}$, their crystals graphs $B_{e}^{s}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$and $B_{e}^{s^{*}}\left(\Lambda_{\underline{\mathbf{s}}^{*}}^{-}\right)$can be realized as the connected components of highest weight vertex $\underline{\emptyset}$ in $B_{e}^{s,+}$ and $B_{e}^{s_{e}^{*},-}$. Since $\xi(\underline{\emptyset})=\underline{\emptyset}$, we derive from the previous proposition the equivalence

$$
\begin{equation*}
\underline{\lambda} \xrightarrow{i} \underline{\mu} \text { in } B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right) \Longleftrightarrow \xi(\underline{\lambda}) \xrightarrow{e-i} \xi(\underline{\mu}) \text { in } B \frac{s^{s^{*}}}{}\left(\Lambda_{\underline{s}^{*}}^{-}\right) . \tag{17}
\end{equation*}
$$

Definition 3.2.2 The vertices of $B_{\bar{e}}^{\frac{s}{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$and $B_{\bar{e}}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$are called the Uglov multipartitions.
Example 3.2.3 Suppose $\ell=1$ and $\underline{s}=(s)$.

- The vertices of $B_{e}^{\frac{s}{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$are the e-restricted partitions, that is the partitions $\lambda=\left(\lambda_{1}, \ldots, \lambda_{p}\right)$ such that $\lambda_{i}-\lambda_{i+1} \leq e-1$ for any $i=1, \ldots, p-1$.
- The vertices of $B_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$are the e-regular partitions, that is the partitions $\lambda$ with at most $e-1$ parts equal.


## Remarks:

(i) : The crystals $B_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$are essentially those used in [2] whereas the crystals $B{ }_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$appear in (7] and 11].
(ii) : Consider $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right)$ and $\underline{s}^{\prime}=\left(s_{0}^{\prime}, \ldots, s_{\ell-1}^{\prime}\right)$ two multicharges such that $s_{k} \equiv s_{k}^{\prime}$ for any $k=0, \ldots, \ell-1$. Then $\Lambda_{\underline{\mathbf{s}}}^{+}=\Lambda_{\underline{\mathbf{s}^{\prime}}}^{+}$, thus the crystals $B B_{\underline{s}}^{\frac{s}{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$and $B^{\frac{s^{\prime}}{}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$are isomorphic. The combinatorial description of this isomorphism is complicated in general (but see Section 5). Nevertheless, in the case when there exists an integer $d$ such that $s_{k}=s_{k}^{\prime}+d e$ for any $k=0, \ldots, \ell-1$, one derives easily from the description of the crystal structure on $B_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$that the relevant isomorphism coincide with the identity map. The situation is the same for the crystals $B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$and $B_{\bar{e}}^{s^{\prime}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$.

When the multicharge $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right)$ verifies $0 \leq s_{0} \leq \cdots \leq s_{\ell-1} \leq e-1$, there exists a combinatorial description of the Uglov multipartitions labelling $B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$due to Foda, Leclerc, Okado, Thibon and Welsh.

Proposition 3.2.4 [6]When $0 \leq s_{0} \leq \cdots \leq s_{\ell-1} \leq e-1$, the multipartition $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$ belongs to $B_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$if and only if:

1. $\underline{\lambda}$ is cylindric, that is, for every $k=0, \ldots, \ell-2$ we have $\lambda_{i}^{(k)} \geq \lambda_{i+s_{k+1}-s_{k}}^{(k+1)}$ for all $i>1$ (the partitions are taken with an infinite numbers of empty parts) and $\lambda_{i}^{(\ell-1)} \geq \lambda_{i+e+s_{0}-s_{\ell-1}}^{(0)}$ for all $i>1$
2. for all $r>0$, among the residues appearing at the right ends of the length $r$ rows of $\underline{\lambda}$, at least one element of $\{0,1, \ldots, e-1\}$ does not appear.

## Remarks:

(i) : By using (17), it is easy to deduce a combinatorial characterization of the Uglov multipartitions appearing in $B^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$when $0 \leq s_{\ell-1} \leq \cdots \leq s_{0} \leq e-1$. These multipartitions are called the FLOTW multipartitions.
(ii) : In the sequel we will essentially consider the crystals $B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$because they can be embedded in a natural way in $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-crystals. Thanks to (17), it will be easy to translate our statements to make them compatible with the crystals $B_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{-}\right)$.
(iii) : There exists also another realization $B_{e}^{a}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$of the abstract crystal $B_{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$which is in particular used by Ariki [2]. This realization is obtained by defining the Fock space of level $\ell$ as a tensor product of Fock spaces of level 1. It is suited to the Specht modules theory for Ariki-Koike algebras introduced by Dipper, James and Mathas [5]. Note that in the level 2 case, Geck [8] has generalized this theory by defining Specht modules adapted to the definition of the Fock spaces used in this paper. It is expected that similar results hold in higher level.
(iv) : The multipartitions which label the vertices of the crystals $B_{e}^{a}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$are called the Kleshchev multipartitions. For any nonnegative integer $n$, the subgraph of the crystal $B_{e}^{a}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$containing all the Kleshchev multipartitions of rank $m \leq n$ coincide with the corresponding subgraph of $B_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$when the multicharge $\underline{s}$ verifies $s_{i}-s_{i+1}>n-1$ for any $i=0, \ldots, \ell-2$. As a consequence Kleshchev multipartitions are particular cases of Uglov multipartitions.

## 4 Embedding of $B \frac{s}{e}\left(\Lambda_{\underline{s}}^{+}\right)$in $B_{\infty}\left(\omega_{\underline{s}^{\diamond}}\right)$

The aim of this section is to show that we have an embedding of crystals from $B_{e}^{s}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$to $B_{\infty}\left(\omega_{\underline{s}^{\diamond}}\right)$ with $\underline{s}^{\diamond}=\left(s_{\ell-1}, \ldots, s_{0}\right)$.

### 4.1 Embedding of $B \frac{s}{e}\left(\Lambda_{\underline{s}}^{+}\right)_{\leq n}$ in $B_{f}^{\frac{s}{s}}\left(\Lambda_{\underline{s}}^{+}\right)_{\leq n}$

In the sequel we denote by $\widetilde{E}_{i}$ and $\widetilde{F}_{i}, i=0, \ldots, e-1$ the crystal operators corresponding to the $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s}}_{e}\right)$-crystals. Consider $\underline{\lambda}$ and $\underline{\mu}$ two multipartitions in $\mathfrak{F}_{e}^{s,+}$ and $i \in\{0, \ldots, e-1\}$ such that $\widetilde{F}_{i}(\underline{\lambda})=\underline{\mu}$. Then $\underline{\mu}$ is obtained by adding an $i$-node $\gamma$ in $\underline{\lambda}$. Set $j=c(\gamma)$. In the sequel we slightly abuse the notation and write $\widetilde{F}_{j}(\underline{\lambda})=\underline{\mu}$ although $j$ does not belong to $\{0, \ldots, e-1\}$ in general. Note that $j$ is uniquely determined from the equality $\widetilde{F}_{i}(\underline{\lambda})=\underline{\mu}$. Observe that with this convention, an arrow in $\mathfrak{F}_{e}^{s,+}$ can be labelled by any integer. To recover the original labelling of a $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s l}}_{e}\right)$-crystal, it suffices to read these labels modulo $e$.
The following proposition is a generalization in level $\ell$ of Proposition 4.1 in [11].
Proposition 4.1.1 Consider $\underline{\lambda}$ a multipartition in $B_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$such that $\underline{\lambda}=\widetilde{F}_{j_{1}} \cdots \widetilde{F}_{j_{r}}(\underline{\emptyset})$. Set $|\underline{\lambda}|=n$. Then for any integer $f>2 n+\max (\underline{s})$, we have $\underline{\lambda}=\widetilde{F}_{j_{1}} \cdots \widetilde{F}_{j_{r}}(\underline{\emptyset})$ in $B \frac{s}{f}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$.

Proof. We proceed by induction on $m \leq n$. When $m=0$, the proposition is immediate. Suppose our statement true for any multipartition $\underline{\lambda}$ with $|\underline{\lambda}|=m$. Consider $\underline{\mu} \in B_{e}^{s}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$with $|\underline{\mu}|=m+1 \leq n$. With the above convention, there exists an integer $j$ and a multipartition $\underline{\lambda} \in B_{\bar{e}}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$such that $\widetilde{F}_{j}(\underline{\lambda})=\underline{\mu}$. Let $i \in\{0, \ldots, e-1\}$ be such that $i \equiv j(\bmod e)$. Let $w_{i}$ be the word obtained by writing the addable or removable $i$-nodes of $\underline{\lambda}$ in decreasing order with respect to $\prec_{\underline{s}}^{+}$as in Section 3.2.1. Similarly let $w_{j}$ be the word obtained by writing the addable or removable $j$-nodes of $\underline{\lambda}$ in decreasing order with respect to $\prec_{\underline{s}}^{+}$. Observe that each $j$-node $\gamma$ verifies $c(\gamma)=\operatorname{res}(\gamma)$ when its residue is computed modulo $f$ for we have supposed $f>2 n+\max (\underline{s})$. Hence the nodes in $w_{j}$ are the addable and removable nodes of $\underline{\lambda}$ with contain $j$. This implies that $w_{j}$ is a subword of $w_{i}$, that is each node appearing in $w_{j}$ also appears in $w_{i}$. If $\gamma$ is a node in $w_{i}$ which does not belong to $w_{j}$, we must thus have $c(\gamma)>j$ or $c(\gamma)<j$. Since the nodes of $w_{i}$ are ordered according to $\preceq_{e}^{+}$, we deduce that $w_{j}$ is a factor of $w_{i}$ (i.e., there is no node $\gamma$ such that $c(\gamma) \neq j$ between two nodes of $\left.w_{j}\right)$.
Write $\widehat{w}_{i}$ and $\widehat{w}_{j}$ respectively for the words obtained from $w_{i}$ and $w_{j}$ by the cancellation process of the factors $R A$. Since these words does not depend on the choice of the factors $R A$ which are deleted, $\widehat{w}_{j}$ is a factor of $\widehat{w}_{i}$. Let $\gamma$ be the good $i$-node in $\widehat{w}_{i}$. Suppose that $\gamma$ is not a node of $\widehat{w}_{j}$. Then, there exists a node $\delta$ in $w_{j}$ which can be paired with $\gamma$ to make a factor $R A$. In this case, $\gamma$ and $\delta$ are nodes of $w_{i}$ because $w_{j}$ is a factor of $w_{i}$. This gives a contradiction. Indeed, the word $w_{i}$ does not depend on the cancellation process for the factors $R A$. So, if we can pair $\gamma$ and $\delta$ together to obtain a factor $R A$ in $w_{j}$, we can do the same in $w_{i}$ (for $w_{j}$ is a factor of $w_{i}$ ) and $\gamma$ cannot be the good $i$-node in $\widehat{w}_{i}$. Hence we have shown that $\gamma$ is a $j$-node of $\widehat{w}_{j}$. Since $\widehat{w}_{j}$ is a factor of $\widehat{w}_{i}, \gamma$ is the rightmost addable $j$-node in $\widehat{w}_{j}$, thus is the good $j$-node for $w_{j}$. By the induction hypothesis, we know that $\underline{\lambda}$ belongs to $B_{f}^{\underline{s}}\left(\Lambda_{\underline{s}}^{+}\right)$. Moreover, we have $\widetilde{F}_{j}(\underline{\lambda})=\underline{\mu}$ in $B_{f}^{\frac{s}{f}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$by the previous arguments. This permits to derive our proposition by induction.

For any fixed nonnegative integer $n$, write $B_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)_{\leq n}=\left\{\underline{\lambda} \in B_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)| | \underline{\mid} \mid \leq n\right\}$. We deduce from the above proposition, that the identity map $\iota_{e, f}: \underline{\lambda} \longmapsto \underline{\lambda}$ from $B_{e}^{\frac{s}{e}}\left(\Lambda_{\underline{s}}^{+}\right)_{\leq n}$ to $B_{f}^{\frac{s}{f}}\left(\Lambda_{\underline{s}}^{+}\right)_{\leq n}$ is an embedding of crystals.

### 4.2 Embedding of $B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$in $B_{\infty}\left(\omega_{\underline{s}}\right)$

For any multicharge $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right)$ we write $\underline{s}^{\diamond}=\left(s_{\ell-1}, \ldots, s_{0}\right)$.
Proposition 4.2.1 Consider a multicharge $\underline{s}$ and $f, n$ two nonnegative integers such that $f>$ $2 n+\max (\underline{s})$. Let $\underline{\lambda}$ be a multipartition in $B_{f}^{\underline{s}}\left(\Lambda_{\underline{s}}^{+}\right)$with $|\underline{\lambda}|=n$. Suppose $\underline{\lambda}=\widetilde{F}_{j_{1}} \cdots \widetilde{F}_{j_{r}}(\underline{\emptyset})$ in $B_{f}^{\frac{s}{f}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$. Then we have $\underline{\lambda}=\widetilde{f}_{j_{1}} \cdots \widetilde{f}_{j_{r}}(\underline{\emptyset})$ in the $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-crystal $B_{\infty}\left(\omega_{\underline{s} \diamond}\right)$.

Proof. We proceed by induction on $m \leq n$. When $m=0$, the proposition is immediate. Suppose our statement true for any multipartition $\underline{\lambda}$ with $|\underline{\lambda}|=m$. Consider $\underline{\mu} \in B_{f}^{s}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$with $|\underline{\mu}|=m+1 \leq n$. There exists an integer $j$ and a multipartition $\underline{\lambda} \in B_{\underline{e}}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$such that $\widetilde{F}_{j}(\underline{\lambda})=\underline{\mu}$. Since $f>2 n$, we have res $(\gamma)=c(\gamma)$ for each node $\gamma$ in $\underline{\lambda}$. Hence the word $w_{j}$ obtained by writing the addable or removable $j$-nodes of $\underline{\lambda}$ in decreasing order with respect to $\prec_{\underline{s}}^{+}$contains exactly the addable and removable nodes of $\underline{\lambda}$ with contain equal to $j$. Set $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$. By definition of the order $\prec_{\underline{s}}^{+}$(see 14), this means that $w_{j}$ is the word obtained by reading the addable and removable nodes with contain equal to $j$ successively in the parts $\lambda^{(\ell-1)}, \lambda^{(\ell-2)}, \ldots, \lambda^{(0)}$ of $\underline{\lambda}$. Now observe that each part $\lambda^{(k)}, k=0, \ldots, \ell-1$ contains at most one addable or removable node with contain $j$ because we must have $\left|\lambda^{(k)}\right| \leq|\underline{\lambda}| \leq n<f / 2$. By the induction hypothesis, we know that $\underline{\lambda}$ belongs to $B_{\infty}\left(\omega_{\underline{s}}{ }^{\triangleright}\right)$. The previous arguments show that the action of $\widetilde{F}_{j}$ on $\underline{\lambda}$ considered as a vertex of the $\mathcal{U}_{v}\left(\widehat{\mathfrak{s}}_{f}\right)$-crystal $B \frac{s}{f}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$is the same as the action of $\widetilde{f}_{j}$ on $\underline{\lambda}$ considered as a vertex of the $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-crystal $B_{\infty}\left(\omega_{\underline{s} \diamond}\right)$ (see (6). The proposition then follows by induction.
Set $B_{\infty}\left(\omega_{\underline{s}}\right)_{\leq n}=\left\{\underline{\lambda} \in B_{\infty}\left(\omega_{\underline{s}}\right)| | \underline{\lambda} \mid \leq n\right\}$. We deduce from the previous proposition that the map

$$
\varphi_{f, \infty}^{(n)}:\left\{\begin{array}{l}
B_{\bar{e}}^{s}\left(\Lambda_{\mathbf{s}}^{+}\right)_{\leq n} \rightarrow B_{\infty}\left(\omega_{s^{\diamond}}\right) \leq n \\
\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right) \stackrel{( }{\longmapsto} \lambda^{(\ell-1)} \otimes \cdots \otimes \lambda^{(0)}
\end{array}\right.
$$

is an embedding of crystals for any $f>2 n+\max (\underline{s})$.
Theorem 4.2.2 Given any positive integer $e$ and any multicharge $\underline{s}$ the map

$$
\varphi_{e, \infty}:\left\{\begin{array}{l}
B_{\bar{e}}^{s}\left(\Lambda_{\mathbf{s}}^{+}\right) \rightarrow B_{\infty}\left(\omega_{s^{\diamond}}\right) \\
\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right) \longmapsto \lambda^{(\ell-1)} \otimes \cdots \otimes \lambda^{(0)}
\end{array}\right.
$$

is an embedding of crystals : for any $\underline{\lambda} \in B_{e}^{\frac{s}{e}}\left(\Lambda_{\mathbf{s}}^{+}\right)$we have

$$
\underline{\lambda}=\widetilde{F}_{j_{1}} \cdots \widetilde{F}_{j_{r}}(\underline{\emptyset}) \text { in } B \frac{s}{e}\left(\Lambda_{\underline{s}}^{+}\right) \Longrightarrow \lambda^{(\ell-1)} \otimes \cdots \otimes \lambda^{(0)}=\widetilde{f}_{j_{1}} \cdots \widetilde{f}_{j_{r}}(\emptyset \otimes \cdots \otimes \emptyset) \text { in } B_{\infty}\left(\omega_{\underline{s}}\right) .
$$

Proof. Consider $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right) \in B_{\bar{e}}^{s}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$such that $\underline{\lambda}=\widetilde{F}_{j_{1}} \cdots \widetilde{F}_{j_{r}}(\underline{\emptyset})$ and set $|\underline{\lambda}|=$ $n$. By Proposition 4.1.1 we have $\underline{\lambda}=\widetilde{F}_{j_{1}} \cdots \widetilde{F}_{j_{r}}(\underline{\emptyset})$ in any crystal $B_{f}^{\underline{s}}\left(\Lambda_{\underline{s}}^{+}\right)$with $f>2 n+\max (\underline{s})$. Fix such an integer $f$. We derive from Proposition 4.2.1 that

$$
\lambda^{(\ell-1)} \otimes \cdots \otimes \lambda^{(0)}=\widetilde{f}_{j_{1}} \cdots \widetilde{f}_{j_{r}}(\emptyset \otimes \cdots \otimes \emptyset) \text { in } B_{\infty}\left(\omega_{\underline{s}}\right) .
$$

Clearly the map $\varphi_{e, \infty}$ is injective. Thus it is an embedding from the $\mathcal{U}_{v}\left(\widehat{\mathfrak{s}}_{e}\right)$-crystal $B_{e}^{\frac{s}{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$to the $\mathcal{U}_{v}\left(\mathfrak{s l}_{\infty}\right)$-crystal $B_{\infty}\left(\omega_{s^{\wedge}}\right)$.

Remark: According to the previous theorem, the crystal $B_{e}^{\frac{s}{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$can be embedded in $B_{\infty}\left(\omega_{\underline{s}^{\diamond}}\right)$. The situation is more complicated for the crystal $B_{e}^{a}\left(\Lambda_{\underline{s}}^{+}\right)$labelled by Kleshchev multipartitions (see Remark (iv) after Proposition 3.2.4). Indeed, the sububrystal $B_{e}^{a}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)_{\leq n}$ can be embedded in a crystal $B_{\infty}\left(\omega_{\underline{s}(n)}\right)$ where the multicharge $\underline{s}(n)=\left(s_{0}(n), \ldots, s_{1}(n)\right)$ verifies $s_{k}(n)-s_{k+1}(n)>n-1$ for any $k=0, \ldots, e-2$. Since $s(n)$ depends on $n$, this procedure cannot provide an embedding of the whole crystal $B_{e}^{a}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$in a crystal $B_{\infty}\left(\omega_{\underline{t}}\right)$ where $\underline{t}$ is a fixed multicharge.

## 5 Isomorphism class of a multipartition

We can now use the above embedding to obtain a simple characterization of the set of Uglov multipartitions.

### 5.1 The extended affine symmetric group $\widehat{S}_{\ell}$

We write $\widehat{S}_{\ell}$ for the extended affine symmetric group in type $A_{\ell-1}$. The group $\widehat{S}_{\ell}$ can be regarded as the group generated by the elements $\sigma_{1}, \ldots, \sigma_{\ell-1}$ and $y_{0}, \ldots ., y_{\ell-1}$ together with the relations

$$
\sigma_{i} \sigma_{i+1} \sigma_{i}=\sigma_{i+1} \sigma_{i} \sigma_{i+1}, \quad \sigma_{i} \sigma_{j}=\sigma_{j} \sigma_{i} \text { for }|i-j|>1, \quad \sigma_{i}^{2}=1
$$

with all indices in $\{1, \ldots, \ell-1\}$ and

$$
y_{i} y_{j}=y_{j} y_{i}, \quad \sigma_{i} y_{j}=y_{j} \sigma_{i} \text { for } j \neq i, i+1, \quad \sigma_{i} y_{i} \sigma_{i}=y_{i+1} .
$$

We identify the subgroup of $\widehat{S}_{\ell}$ generated by the transpositions $\sigma_{i}, i=1, \ldots, \ell-1$ with the symmetric group $S_{\ell}$ of rank $\ell$. For any $i \in\{0, \ldots, \ell-1\}$, we set $z_{i}=y_{1} \cdots y_{i}$. Write also $\xi=\sigma_{\ell-1} \cdots \sigma_{1}$ and $\tau=y_{\ell} \xi$. Since $y_{i}=z_{i-1}^{-1} z_{i}, \widehat{S}_{\ell}$ is generated by the transpositions $\sigma_{i}$ with $i \in\{1, \ldots, \ell-1\}$ and the elements $z_{i}$ with $i \in\{1, \ldots, \ell\}$. Observe that for any $i \in\{1, \ldots, \ell-1\}$, we have

$$
\begin{equation*}
z_{i}=\xi^{(\ell-i)} \tau^{(i)} . \tag{18}
\end{equation*}
$$

This implies that $\widehat{S}_{\ell}$ is generated by the transpositions $\sigma_{i}$ with $i \in\{1, \ldots, \ell-1\}$ and $\tau$.
Consider $e$ a fixed positive integer. We obtain a faithful action of $\widehat{S}_{\ell}$ on $\mathbb{Z}^{\ell}$ by setting for any $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right) \in \mathbb{Z}^{\ell}$

$$
\sigma_{i}(\underline{s})=\left(s_{0}, \ldots, s_{i}, s_{i-1}, \ldots, s_{\ell-1}\right) \text { and } y_{i}(\underline{s})=\left(s_{0}, \ldots, s_{i-1}, s_{i}+e, \ldots, s_{\ell-1}\right)
$$

Then $\tau(\underline{s})=\left(s_{1}, s_{2}, \ldots, s_{\ell-1}, s_{0}+e\right)$. We denote by $\mathcal{C}(\underline{s})$ the class of the multicharge $\underline{s}$ under the action of $\widehat{S}_{\ell}$ on $\mathbb{Z}^{\ell}$. Clearly each class $\mathcal{C}(\underline{s})$ contains a unique multicharge $\underline{\widetilde{s}}=\left(\widetilde{s}_{0}, \ldots, \widetilde{s}_{\ell-1}\right)$ such that

$$
\begin{equation*}
0 \leq \widetilde{s}_{\ell-1} \leq \cdots \leq \widetilde{s}_{0} \leq e-1 . \tag{19}
\end{equation*}
$$

Hence the classes $\mathcal{C}(\underline{s})$ are parametrized by the multicharge verifying (19). Given any multicharge $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right) \in \mathbb{Z}^{\ell}$, it is easy to determinate $w \in \widehat{S}_{\ell}$ such that $\underline{\widetilde{s}}=w(\underline{s})$. Choose $k \in \mathbb{N}$
minimal to have $s_{i}+k e \geq 0$ for any $i=0, \ldots, \ell-1$. Then $z_{\ell-1}^{k}(\underline{s}) \in \mathbb{Z}^{\ell}$. Consider $\sigma \in S_{\ell}$ such that the coordinates of $\underline{s}^{(\ell-1)}=\sigma z_{\ell-1}^{k}(\underline{s})$ weakly decrease. Write $r_{\ell-2}$ for the quotient of the division of $\underline{s}_{\ell-2}^{(\ell-2)}$ by $e$ and set $\underline{s}^{(\ell-3)}=z_{\ell-2}^{-r_{\ell-2}}\left(\underline{g}^{(\ell-2)}\right)$. By induction one can compute a sequence $\underline{s}^{(\ell-2)}, \ldots, \underline{s}^{(0)}$ such that, for any $i=1, \ldots, \ell-2, \underline{s}^{(i-1)}=z_{i}^{-r_{i}}\left(\underline{s}^{(i)}\right)$ where $r_{i}$ is the quotient of the division of $\underline{s}_{i}^{(i)}-\underline{s}_{i+1}^{(i)}$ by $e$. We have then $\underline{\widetilde{s}}=\underline{s}^{(0)}$ and

$$
\begin{equation*}
\underline{\tilde{s}}=w(\underline{s})=z_{0}^{-r_{0}} \cdots z_{\ell-2}^{-r_{\ell-2}} \sigma z_{\ell-1}^{k}(\underline{s}) . \tag{20}
\end{equation*}
$$

### 5.2 Action of the transformations $s_{i}$ and $\tau$ on a multipartition

Consider a multicharge $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right)$ and $w$ an element of the extended affine symmetric group. Set $\underline{s}^{\prime}=w(\underline{s})$. Since the indices of the fundamental weights of $\mathcal{U}_{v}^{+}\left(\widehat{\mathfrak{s}}_{e}\right)$ belong to $\mathbb{Z} / e \mathbb{Z}$, we have $\Lambda_{\underline{\mathbf{s}}^{\prime}}^{+}=\Lambda_{\underline{\mathbf{s}}}^{+}$. This implies that the crystals graphs $B_{e}^{s}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$and $B_{e}^{s_{e}^{\prime}}\left(\Lambda_{\underline{\underline{s}}^{\prime}}^{+}\right)$are isomorphic. Write $\Gamma_{\underline{s}, \underline{s}^{\prime}}$ for the isomorphism between $B \frac{\underline{e}}{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$and $B_{e}^{s^{\prime}}\left(\Lambda_{\underline{\underline{s}}^{\prime}}^{+}\right)$. Given a multipartition $\underline{\lambda}$ in $B \frac{s}{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$, we are going to see how it is possible to determinate $\mu=\Gamma_{\underline{s}, \mathbf{s}^{\prime}}(\underline{\lambda})$ in a non-inductive way, that is, without computing a path joining $\underline{\lambda}$ to $\underline{\emptyset}$ in $B_{e}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$. According to Section 5.1, $w$ decomposes as a product of the elements $\sigma_{i} i=1, \ldots, \ell-1$ and $\tau$. We write for short $\Xi_{\underline{s}}$ and $\Sigma_{\underline{s}, i}$ respectively for the crystal graph isomorphisms $\Gamma_{\underline{s}, \tau(\underline{s})}$ and $\Gamma_{\underline{s}, \sigma_{i}(\underline{s})}$. The following proposition is a generalization of [11, Prop. 3.1].

Proposition 5.2.1 Consider $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$ a multipartition and $\underline{s}$ a multicharge. Then

$$
\Xi_{\underline{\underline{s}}}(\underline{\lambda})=\left(\lambda^{(1)}, \ldots, \lambda^{(\ell-1)}, \lambda^{(0)}\right) .
$$

Proof. Set $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right)$. Then $\tau(\underline{s})=\left(s_{1}, \ldots, s_{\ell-1}, s_{0}+e\right)$. Let $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$ a multipartition and put $\underline{\lambda}^{\prime}=\left(\lambda^{(1)}, \ldots, \lambda^{(\bar{\ell}-1)}, \lambda^{(0)}\right)$. Consider $i \in\{0,1, \ldots, e-1\}$ and $\gamma_{1}=$ $\left(a_{1}, b_{1}, c_{1}\right), \gamma_{2}=\left(a_{2}, b_{2}, c_{2}\right)$ two $i$-nodes of $\underline{\lambda}$. Then $\gamma_{1}^{\prime}=\left(a_{1}, b_{1}, c_{1}-1(\bmod e)\right)$ and $\gamma_{2}^{\prime}=$ $\left(a_{2}, b_{2}, c_{2}-1(\bmod e)\right)$ are two $i$-nodes of $\underline{\lambda}^{\prime}$. We then easily check that $\gamma_{2} \prec_{\underline{s}}^{+} \gamma_{1}$ if and only if $\gamma_{2}^{\prime} \prec_{\tau(\underline{s})}^{+} \gamma_{1}^{\prime}$. This implies that $\Xi_{\underline{s}}(\underline{\lambda})=\underline{\lambda}^{\prime}$.

Consider $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$ a multipartition and $\underline{s}$ a multicharge. Let $i$ be any integer of $\{1, \ldots, n\}$. Set

$$
\begin{equation*}
\psi_{s_{i}, s_{i-1}}\left(\lambda^{(i)}, \lambda^{(i-1)}\right)=\left(\widetilde{\lambda}^{(i-1)}, \widetilde{\lambda}^{(i)}\right) \tag{21}
\end{equation*}
$$

where $\psi_{s_{i}, s_{i-1}}$ is the crystal graph isomorphism defined in (7).
Proposition 5.2.2 With the above notation, we have

$$
\Sigma_{\underline{s}, i}(\underline{\lambda})=\left(\lambda^{(0)}, \ldots, \tilde{\lambda}^{(i)}, \tilde{\lambda}^{(i-1)}, \ldots \lambda^{(\ell-1)}\right)
$$

that is $\Sigma_{\underline{s}, i}(\underline{\lambda})$ is obtained by replacing in $\underline{\lambda}, \lambda^{(i-1)}$ by $\widetilde{\lambda}^{(i)}$ and $\lambda^{(i)}$ by $\widetilde{\lambda}^{(i-1)}$.
Proof. We have to prove that the diagram

$$
\begin{array}{lll}
B_{\bar{e}}^{\underline{s}}\left(\Lambda_{\mathbf{s}}^{+}\right) & \xrightarrow{\varphi_{e, \infty}} & B_{\infty}\left(\omega_{s_{\ell-1}}\right) \otimes \cdots \otimes B_{\infty}\left(\omega_{s_{i}}\right) \otimes B_{\infty}\left(\omega_{s_{i-1}}\right) \otimes \cdots \otimes B_{\infty}\left(\omega_{s_{0}}\right) \\
\Sigma_{\underline{s}, i} \downarrow  \tag{22}\\
B_{e}^{\sigma_{i}(\underline{s})}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right) & \stackrel{\downarrow \psi_{s_{i}, s_{i-1}}}{ } & \\
& B_{\infty}\left(\omega_{s_{\ell-1}}\right) \otimes \cdots \otimes B_{\infty}\left(\omega_{s_{i-1}}\right) \otimes B_{\infty}\left(\omega_{s_{i}}\right) \otimes \cdots \otimes B_{\infty}\left(\omega_{s_{0}}\right)
\end{array}
$$

commutes. Consider a multipartition $\underline{\lambda} \in B_{\underline{e}}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$. Set $\underline{\lambda}=\left(\lambda^{(0)}, \ldots, \lambda^{(\ell-1)}\right)$. Let $\widetilde{F}_{i_{1}}, \ldots, \widetilde{F}_{i_{r}}$ be a sequence of crystal operators such that $\underline{\lambda}=\widetilde{F}_{i_{1}} \cdots \widetilde{F}_{i_{r}}(\underline{\emptyset})$. Then we can consider the multipartition $\underline{\mu}=\widetilde{F}_{i_{1}}, \ldots, \widetilde{F}_{i_{r}}(\underline{\emptyset})$ in the crystal $B_{e}^{\sigma_{i}(\underline{s})}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$. Observe first that

$$
\begin{equation*}
\psi_{s_{i}, s_{i-1}} \circ \varphi_{e, \infty}(\underline{\emptyset})=\varphi_{e, \infty} \circ \Sigma_{\underline{s}, i}(\underline{\emptyset}) . \tag{23}
\end{equation*}
$$

Moreover, the maps $\varphi_{e, \infty}, \Sigma_{\underline{s}, i}$ and $\psi_{s_{i}, s_{i-1}}$ commutes with the crystal operators. This permits to write

$$
\psi_{s_{i}, s_{i-1}} \circ \varphi_{e, \infty}(\underline{\lambda})=\psi_{s_{i}, s_{i-1}} \circ \varphi_{e, \infty}\left(\widetilde{F}_{i_{1}} \cdots \widetilde{F}_{i_{r}}(\underline{\emptyset})\right)=\widetilde{f}_{i_{1}} \cdots \widetilde{f}_{i_{r}}\left(\psi_{s_{i}, s_{i-1}} \circ \varphi_{e, \infty}(\underline{\emptyset})\right) .
$$

One the other hand we have

$$
\varphi_{e, \infty} \circ \Sigma_{\underline{s}, i}(\underline{\lambda})=\varphi_{e, \infty} \circ \Sigma_{\underline{s}, i}\left(\widetilde{F}_{i_{1}} \cdots \widetilde{F}_{i_{r}}(\underline{\emptyset})\right)=\widetilde{f}_{i_{1}} \cdots \widetilde{f}_{i_{r}}\left(\varphi_{e, \infty} \circ \Sigma_{\underline{s}, i}(\underline{\emptyset})\right) .
$$

Hence we derive the equality $\psi_{s_{i}, s_{i-1}} \circ \varphi_{e, \infty}(\underline{\lambda})=\varphi_{e, \infty} \circ \Sigma_{\underline{s}, i}(\underline{\lambda})$ from (23). This shows that the diagram (22) commutes and establish our proposition.

Example 5.2.3 Suppose $\underline{s}=(4,0,1)$ and $\underline{\lambda}=\left(\lambda^{(0)}, \lambda^{(1)}, \lambda^{(2)}\right)$ with $\lambda^{(0)}=(4,3,3,2), \lambda^{(1)}=$ $(3,3,1)$ and $\lambda^{(2)}=(5,3,2)$. Let us compute $\Sigma_{\underline{s}, 2}(\underline{\lambda})$. The infinite columns associated to $\lambda^{(1)}$ and $\lambda^{(2)}$ are respectively

$$
\mathcal{C}_{1}=\begin{array}{|c|}
\hline 3 \\
\hline 2 \\
\hline \overline{1} \overline{3} \\
\hline \overline{3} \\
\hline \overline{4} \\
\hline \cdots \\
\hline
\end{array} \text { and } \mathcal{C}_{2}=\begin{array}{|c|}
\hline 6 \\
\hline 3 \\
\hline 1 \\
\hline \overline{2} \\
\hline \overline{3} \\
\hline \overline{4} \\
\hline \cdots \\
\hline
\end{array} .
$$

For $a=\overline{3}$ the corresponding finite columns are

$$
C_{1}=\begin{array}{|l|}
\hline 3 \\
\hline 2 \\
\hline \overline{1} \\
\hline \overline{3} \\
\hline
\end{array} \text { and } C_{2}=\begin{array}{|l|}
\hline 6 \\
\hline 3 \\
\hline 1 \\
\hline \overline{2} \\
\hline \overline{3} \\
\hline
\end{array} .
$$

We have to determinate the image of $C_{2} \otimes C_{1}$ under the isomorphism $\theta_{5,4}$ of Proposition 2.3.5.

We obtain $\left\{y_{1}, y_{2}, y_{3}, y_{4}\right\}=\{3,1, \overline{2}, \overline{3}\}$. This gives $\theta_{5,4}\left(C_{2} \otimes C_{1}\right)=C_{1}^{\prime} \otimes C_{2}^{\prime}$ with $C_{1}^{\prime}=$| 3 |
| :--- |
| 1 |
| $\overline{2}$ |
| $\overline{3}$ |

$C_{2}^{\prime}=$| 6 |
| :---: |
| 3 |
| 2 |
| $\overline{1}$ |
| $\overline{3}$ |. Hence $\psi_{1,0}\left(\mathcal{C}_{2} \otimes \mathcal{C}_{1}\right)=\mathcal{C}_{1}^{\prime} \otimes \mathcal{C}_{2}^{\prime}$ where

$$
\mathcal{C}_{1}^{\prime}=\begin{array}{|c|}
\hline 3 \\
\hline 1 \\
\hline \overline{2} \overline{3} \\
\hline \overline{4} \\
\hline \cdots \\
\hline \cdots \\
\hline
\end{array} \text { and } \mathcal{C}_{2}^{\prime}=\begin{array}{|c|}
\hline 6 \\
\hline 3 \\
\hline 2 \\
\hline \overline{1} \\
\hline \overline{3} \\
\hline \overline{4} \\
\hline \cdots \\
\hline
\end{array} .
$$

Finally we derive $\Sigma_{\underline{s}, 2}(\underline{\lambda})=\left(\lambda^{(0)}, \widetilde{\lambda}^{(2)}, \widetilde{\lambda}^{(1)}\right)$ with $\widetilde{\lambda}^{(1)}=(3,2)$ and $\widetilde{\lambda}^{(2)}=(5,3,3,1)$.

Remark: Assume that $\underline{\lambda}=\left(\lambda^{(0)}, \lambda^{(1)}\right)$ is a bipartition such that $\underline{\lambda}$ belongs to $B_{\bar{e}}^{s}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$where the multicharge $\underline{s}=\left(s_{0}, s_{1}\right)$ verifies $s_{0}<s_{1}$. Then the combinatorial procedure illustrated by the previous example which permits to compute the crystal isomorphisms $\Sigma_{\underline{s}, i}$, essentially reduces, up to a renormalization due to the change of labelling of the Dynkin diagram in type $A_{e-1}^{(1)}$ (see (11)), to the algorithm depicted in Theorem 4.6 of [11].

### 5.3 A non recursive characterization of the Uglov multipartitions

Consider a multicharge $\underline{s}$ and define the multicharge $\underline{\widetilde{s}}$ as in (19). Then the crystals $B_{\bar{e}}^{s}\left(\Lambda_{\underline{s}}^{+}\right)$and $B_{\bar{e}}^{\underline{s}}\left(\Lambda_{\underline{s}}^{+}\right)$are isomorphic. For any multipartition $\underline{\lambda} \in B_{\bar{e}}^{\underline{s}}\left(\Lambda_{\underline{s}}^{+}\right)$, write $F(\underline{\lambda}) \in B_{\underline{\tilde{s}}}^{\underline{\tilde{s}}}\left(\Lambda_{\underline{s}}^{+}\right)$for its image under this crystal isomorphism. It is possible to obtain $F(\underline{\lambda})$ from $\underline{\lambda}$ by using results of $\S \overline{5} .2$. We keep the notations of $\oint 5.1$. For any $i=0, \ldots, \ell-1$, we have $z_{i}=\xi^{(\ell-i)} \tau^{(i)}$. This permits to compute $\Gamma_{\underline{s}, z_{i}(\underline{s})}(\underline{\lambda})$ by using Propositions 5.2.1 and 5.2.2. We have then

$$
F(\underline{\lambda})=\Gamma_{\underline{s}, w(\underline{s})}(\underline{\lambda})
$$

with the notation (2G).
Conversely, given any FLOTW multipartition $\underline{\mu}$ and any multicharge $\underline{s}$, one can compute the multipartition $\underline{\lambda} \in B_{\bar{e}}^{\underline{s}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$such that $F(\underline{\lambda})=\underline{\mu}$. Indeed, we have then $\underline{\lambda}=\Gamma_{\widetilde{\underline{s}}, w^{-1}(\underline{\mathfrak{s}})}(\underline{\mu})$. By remark (i) following Proposition 3.2.4, we thus derive a non recursive combinatorial description of the Uglov multipartitions labelling $B_{e}^{+,(n)}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$.

Proposition 5.3.1 For any multicharge s

$$
B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)_{\leq n}=\left\{\Gamma_{\tilde{\underline{\underline{s}}}, w^{-1}(\underline{\underline{( }})}(\underline{\mu}) \mid \underline{\mu} \in B_{\bar{e}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)_{\leq n}\right\}
$$

where $w$ is obtained from $\underline{s}$ as in (2d).

### 5.4 Isomorphism class of a multipartition

Suppose that $e$ is a fixed positive integer and $\underline{s}$ a multicharge of level $\ell$. Consider $\underline{\lambda}$ a multipartition in $B \frac{s}{e}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$. The isomorphism class of $\underline{\lambda}$ is the set

$$
\mathfrak{C}(\underline{\lambda})=\left\{\Gamma_{\underline{s}, \underline{s}^{\prime}}(\underline{\lambda}) \mid s^{\prime} \in \mathcal{C}(\underline{s})\right\} .
$$

Thus $\mathfrak{C}(\underline{\lambda})$ is the set of all multipartitions $\underline{\mu}$ which appear at the same place as $\underline{\lambda}$ in a crystal $B_{\bar{e}}^{s^{\prime}}\left(\Lambda_{\underline{\mathbf{s}}}^{+}\right)$where $\underline{s}^{\prime}$ is a multicharge of the orbit of $\underline{s}$ under the action of $\widehat{S}_{\ell}$. Then $\mathfrak{C}(\underline{\lambda})$ can be determined from $\underline{\lambda}$ by applying successive elementary transformations using Propositions 5.2.1 and 5.2.2. Observe that for any $\underline{\mu} \in \mathfrak{C}(\underline{\lambda})$ we must have $|\underline{\lambda}|=|\underline{\mu}|$. This implies in particular that $\mathfrak{C}(\underline{\lambda})$ is finite. The cardinality of $\mathfrak{C}(\underline{\lambda})$ is in general rather complicated to evaluate without computing the whole class $\mathfrak{C}(\underline{\lambda})$. Nevertheless, we are going to see in Theorem 5.4.2, that it is possible to obtain an upper bound for $\operatorname{card}(\mathfrak{C}(\underline{\lambda}))$ and to determinate a finite subset $S_{\underline{\boldsymbol{\lambda}}}$ of $\widehat{S}_{\ell}$ such that

$$
\mathfrak{C}(\underline{\lambda})=\left\{\Gamma_{\underline{s}, \underline{s}^{\prime}}(\underline{\lambda}) \mid s^{\prime} \in S_{\underline{\lambda}} \cdot \underline{s}\right\} .
$$

Lemma 5.4.1 Let $\underline{\lambda}$ be a multipartition of rank $n$. Assume that $\underline{s}$ is a multicharge of level $\ell$ verifying :

$$
\begin{equation*}
s_{j}-s_{j+1}>n-1 \tag{24}
\end{equation*}
$$

for $j=0, \ldots, \ell-2$. Then for any $k \in\{0, \ldots, \ell-1\}$ we have $\Gamma_{\underline{s}, z_{k}(\underline{s}}(\underline{\lambda})=\underline{\lambda}$.
Proof. Let $\underline{s}=\left(s_{0}, \ldots, s_{\ell-1}\right)$ be such that $s_{j}-s_{j+1}>n-1$ for $j=0, \ldots, \ell-2$ and let $\underline{\mu}$ be a multipartition in $B_{e}^{s,+}$ such that $|\mu| \leq n$. Let $i \in\{0,1, \ldots, e-1\}$ and consider $\gamma_{1}=\left(a_{1}, b_{1}, c_{1}\right)$ and $\gamma_{2}=\left(a_{2}, b_{2}, c_{2}\right)$ two $i$-nodes in $\mu$ such that $c_{1}<c_{2}$. We have $b_{1}-a_{1}+s_{c_{1}}-\left(b_{2}-a_{2}+s_{c_{2}}\right)>$ $b_{1}-a_{1}-\left(b_{2}-a_{2}\right)+n-1 \geq 0$. Hence, the contains of $\gamma_{1}$ and $\gamma_{2}$ considered as nodes of $\underline{\mu} \in B_{e}^{s},+$ are such that $c\left(\gamma_{1}\right)>c\left(\gamma_{2}\right)$. Hence we have $\gamma_{2} \prec_{\underline{s}}^{+} \gamma_{1}$.

Now, put $\left(s_{0}^{\prime}, \ldots, s_{l-1}^{\prime}\right):=z_{k}(\underline{s})=\left(s_{0}+e, \ldots, s_{k}+e, s_{k+1}, \ldots, s_{\ell-1}\right)$. As we have $s_{j}^{\prime}-s_{j+1}^{\prime}>n-1$ for $j=0, \ldots, \ell-2$, the above discussion shows that the order $\prec_{\underline{+}}^{+}$and $\prec_{z_{k}(\underline{s})}^{+}$on the $i$-nodes of $\underline{\mu}$ coincide.

Theorem 5.4.2 Suppose that $e$ is a fixed positive integer and $\underline{s}$ a multicharge of level $\ell$ such that

$$
0 \leq s_{\ell-1} \leq \cdots \leq s_{0}<e
$$

Consider $\underline{\lambda}$ a FLOTW multipartition in $B_{\bar{e}}^{\underline{s}}\left(\Lambda_{\underline{s}}^{+}\right)$of order $n$. For any $j \in\{0, \ldots, \ell-2\}$, let $p_{j}$ be the minimal nonnegative integer such that

$$
s_{j}+p_{j} e-s_{j+1}>n-1 .
$$

Then we have :

$$
\mathfrak{C}(\underline{\lambda})=\left\{\Gamma_{\underline{s}, z_{0}^{a_{0}} \ldots z_{\ell-2}^{\alpha_{\ell-2}} \sigma(\underline{s})}(\underline{\lambda}) \mid \sigma \in S_{\ell} \text { and } 0 \leq a_{j} \leq p_{j} \text { for any } j \in\{0, \ldots, \ell-2\}\right\} .
$$

In particular, $\mathfrak{C}(\underline{\lambda})$ is finite and

$$
\operatorname{card}(\mathfrak{C}(\underline{\lambda})) \leq \ell!\prod_{j=0}^{\ell-2}\left(p_{j}+1\right)
$$

Proof. Consider $\underline{\mu} \in \mathfrak{C}(\underline{\lambda})$. According to (20), one can write

$$
\underline{\mu}=\Gamma_{\underline{s}, w^{-1}(\underline{s})}(\underline{\lambda}) \text { with } w^{-1}=z_{\ell-1}^{-k} \sigma z_{\ell-2}^{r_{\ell-2}} \cdots z_{0}^{r_{0}} .
$$

By Remark (ii) following Example 3.2.3, we have $\underline{\mu}=\Gamma_{\underline{s}, w^{-1}(\underline{s})}(\underline{\lambda})=\Gamma_{\underline{s}, u(\underline{s})}(\underline{\lambda})$ where $u=$ $\sigma z_{\ell-2}^{r_{\ell-2}} \cdots z_{0}^{r_{0}}$. By Lemma 5.4.1, we can thus derive

$$
\mathfrak{C}(\underline{\lambda})=\left\{\Gamma_{\underline{s}, z_{0}^{a_{0}} \ldots z_{\ell-2}^{a_{\ell-2} \sigma(\underline{s})}}(\underline{\lambda}) \mid \sigma \in S_{\ell} \text { and } 0 \leq a_{j} \leq p_{j} \text { for any } j \in\{0, \ldots, \ell-2\}\right\}
$$

and our theorem follows.
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