# An hp-DISCONTINUOUS GALERKIN METHOD <br> FOR THE TIME-DEPENDENT MAXWELL'S EQUATIONS 

A. Zaghdani, Christian Daveau

## - To cite this version:

A. Zaghdani, Christian Daveau. An hp-DISCONTINUOUS GALERKIN METHOD FOR THE TIME-DEPENDENT MAXWELL'S EQUATIONS. 2007. hal-00150119

## HAL Id: hal-00150119

## https://hal.science/hal-00150119

Preprint submitted on 30 May 2007

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# An $h p$-DISCONTINUOUS GALERKIN METHOD FOR THE TIME-DEPENDENT MAXWELL'S EQUATIONS 

ZAGHDANI A $\dagger$ and DAVEAU C $\dagger$<br>$\dagger$ Université de Cergy-Pontoise, Laboratoire de Mathematiques.<br>95302, Cergy-Pontoise, France

May 21, 2007


#### Abstract

One discontinuous Galerkin method for the numerical approximation for the time-dependant Maxwell's equations in "stable medium" with supraconductive boundary, is introduced and analysed. its $h p$-analysis is carried out and error estimates that are optimal in the meshsize $h$ and slightly suboptimal in the approximation degree $p$ are obtained.


## 1 Introduction

The problem considered for the most of this Note is the initial-boundary value problem derived from Maxwell's equations in "stable medium" with supraconductive boundary

$$
\begin{align*}
& \frac{\partial^{2} u}{\partial t^{2}}+c^{2} \nabla \times(\nabla \times u)=f, \quad \nabla \cdot u=0 \quad \text { in } \Omega \times I  \tag{1}\\
& n \times u(x, t)=0 \quad \text { on } \quad \partial \Omega \times I, \quad u(x, 0)=u_{0}(x), \quad \frac{\partial u}{\partial t}(x, 0)=u_{1}(x) \quad \text { on } \Omega \tag{2}
\end{align*}
$$

Here $\Omega$ is a convex polyhedron included in $\mathrm{IR}^{3}, I=\left[0, t^{*}\right] \subset \mathrm{IR}, u_{0}$ and $u_{1}$ are in $H_{0}(\nabla \times, \Omega) \cap$ $H(\nabla \cdot 0, \Omega)$ and $f$ is defined on $\Omega \times I$. Physically, $u$ is the electric field, $f$ is related to a current density. $\mu_{0} \varepsilon_{0} c^{2}=1$ where $\mu_{0} \approx 4 \pi 10^{-7} \mathrm{H} . \mathrm{m}^{-1}$ and $\varepsilon_{0} \approx\left(36 \pi 10^{9}\right)^{-1} \mathrm{~F} . \mathrm{m}^{-1}$ are the magnetic permeability and the electric permittivity in vacuum, respectively. If we assume that the domain $\Omega$ is "stable medium" with supraconductive boundary and if $u$ is the exact solution of Maxwell problem then $u$ and $\nabla \times u$ belong to $H^{1}(\Omega)^{3}$. For the notations, if $I$ is an interval, $X$ is one function space and $\phi$ is a function on $\Omega \times I$ then $\|\phi\|_{L^{p}(I, X)}$ denote the norm in $L^{p}(I)$ of the function $t \rightarrow\|\phi(\cdot, t)\|_{X} . L^{p}(X)$ is short for $L^{p}(I, X)$. Let $\Pi_{h}$ be a partition into tetrahedra for $\Omega$ and consider the same spaces and notations in [2].
Finite element spaces: Let $p=\left(p_{K}\right)_{K \in \Pi_{h}}$ be a degree vector that assigns to each element $K \in \Pi_{h}$ a polynomial approximation order $p_{K} \geq 1$. The generic $h p$-finite element space of piecewise polynomials is given by $S^{p}\left(\Pi_{h}\right):=\left\{u \in L^{2}(\Omega): u_{\mid K} \in S^{p_{K}}(K) \quad \forall K \in \Pi_{h}\right\}$ where $S^{p_{K}}(K)$ is the space of real polynomials of degree at most $p_{K}$ in $K$. We also set $\Sigma_{h}:=S^{p}\left(\Pi_{h}\right)^{3}$.
Now, fix a face $e \subset F_{h}$ and define the local parameters $h, p$ by $h:=\min \left(h_{K}, h_{K^{\prime}}\right), p=$ $\max \left(p_{K}, p_{K^{\prime}}\right)$ in the case of interior faces and $h:=h_{K}, p=p_{K}$ in the case of boundary faces [1].

## 2 Formulation for the Maxwell problem

In order to derive a weak formulation of (1)-(2), we note that formulas (1)-(2) in [2] implies for any $u$ with $\nabla \times u \in H(\nabla \times, \Omega)$

$$
c^{2}(\nabla \times(\nabla \times u), v)=c^{2}(\nabla \times u, \nabla \times v)+a(u, v)
$$

where we have denoted by

$$
a(u, v)=c^{2}<n \times(\nabla \times u), v>-c^{2} \sum_{e \subset F_{h}^{I}}<[v]_{T},\{\nabla \times u\}>_{e} .
$$

Now, we introduce the penalty term via the form

$$
J_{0}(u, v)=J(u, v)+J^{\sigma}(u, v)-a(v, u) \text { with } J(u, v)=(\nabla \cdot u, \nabla \cdot v)
$$

and

$$
J^{\sigma}(u, v)=\sum_{e \subset F_{h}^{I}}<\sigma[u]_{N},[v]_{N}>_{e}+\sum_{e \subset F_{h}}<\sigma[u]_{T},[v]_{T}>_{e} \quad u, v \in H^{1}\left(\nabla \times, \Pi_{h}\right)^{3}
$$

where $\sigma:=\kappa p^{2} / h$ is a stabilization parameter and $\kappa$ is a constant supposed $\geq 1$. We also define
$A(u, v)=c^{2}(\nabla \times u, \nabla \times v)+a(v, u)-a(u, v)+J(u, v)$ and $B(u, v)=A(u, v)+J^{\sigma}(u, v)$.

### 2.1 Properties of the bilinear form

### 2.1.1 Mesh-dependant norm

We now, introduce norm associated with the bilinear form $B$ and set for $u \in H^{1}\left(\nabla \times, \Pi_{h}\right)$

$$
\begin{aligned}
\|u\|_{h}^{2}=\|u\|^{2}+\|\nabla \times u\|^{2} & +\|\nabla \cdot u\|^{2}+\left\|\frac{1}{\sqrt{\sigma}}<\nabla \times u>\right\|_{0, F_{h}}^{2} \\
+ & \left\|\sqrt{\sigma}[u]_{N}\right\|_{0, F_{h}^{I}}^{2}+\left\|\sqrt{\sigma}[u]_{T}\right\|_{0, F_{h}}^{2} .
\end{aligned}
$$

We start by studying the continuity of the bilinear forms introduced above. We have :
Proposition $2.1 \forall v, u \in H^{1}\left(\nabla \times, \Pi_{h}\right)$ there exists a constant $C$ independent of $h$ and $p$ such that

$$
|A(u, v)| \leq C\|u\|_{h}\|v\|_{h} \text { and }\left|J^{\sigma}(u, v)\right| \leq C\|u\|_{h}\|v\|_{h} .
$$

Proof: The proof is easly deduced from the definition of $A, J^{\sigma},\|\cdot\|_{h}$ and the CauchySchwarz inequality.
In order to study the coercivity of the bilinear form $B$, we start by introducing the following inequality of Poincarré-Friedrichs type valid for $u \in H^{1}\left(\Pi_{h}\right)^{3}$.

Lemma 2.1 Let $u \in H^{1}\left(\Pi_{h}\right)^{3}$. Then there exists $C$ independent of $h$ and $p$ such that

$$
\|u\|^{2} \leq C\left(\|\nabla \times u\|^{2}+\|\nabla \cdot u\|^{2}+\sum_{e \subset F_{h}}\left\|\sqrt{\sigma}[u]_{T}\right\|_{0, e}^{2}+\sum_{e \subset F_{h}^{I}}\left\|\sqrt{\sigma}[u]_{N}\right\|_{0, e}^{2}\right)
$$

Proof: The proof follows immediately from Lemma 3.1 in [2] since $\kappa p^{2} \geq 1$.
Now, the following coercivity result holds.

Proposition 2.2 There exists two constants $\alpha>0$ and $\tilde{C}>0$ independent of $h$ and $p$ such that

$$
B(v, v) \geq \alpha\|v\|_{h}^{2}+\tilde{C} J^{\sigma}(v, v) \quad \forall v \in \Sigma_{h}
$$

Proof: The proof is easly deduced from the Poincaré inequality given in the previous Lemma. Indeed, $\forall v \in \Sigma_{h}$ and $\forall \alpha>0$ we have

$$
\begin{aligned}
B(v, v)-\alpha\|v\|_{h}^{2} & =(1-\alpha) A(v, v)+(1-\alpha) J^{\sigma}(v, v)-\alpha\|v\|^{2} \\
& \geq(1-\alpha) A(v, v)+(1-\alpha) J^{\sigma}(v, v)-\alpha C A(v, v)-\alpha C J^{\sigma}(v, v) \\
& \geq(1-\alpha-\alpha C)\left(A(v, v)+J^{\sigma}(v, v)\right)
\end{aligned}
$$

Now, the following $h p$-approximation result to interpolate scalar function holds (see [1]).
Proposition 2.3 Let $K \in \Pi_{h}$ and suppose that $u \in H^{t_{K}}(K), t_{K} \geq 1$. Then there exists $a$ sequence of polynomials $\pi_{p_{K}}^{h_{K}}(u) \in S^{p_{K}}(K), p_{K}=1,2 \ldots$ satisfying, $\forall 0 \leq q \leq t_{K}$
$\left\|u-\pi_{p_{K}}^{h_{K}}(u)\right\|_{q, K} \leq C \frac{h_{K}^{\min \left(p_{K}+1, t_{K}\right)-q}}{p_{K}^{t_{K}-q}}\|u\|_{t_{K}, K} \quad$ and $\left\|u-\pi_{p_{K}}^{h_{K}}(u)\right\|_{0, \partial K} \leq C \frac{h_{K}^{\min \left(p_{K}+1, t_{K}\right)-\frac{1}{2}}}{p_{K}^{t_{K}-\frac{1}{2}}}\|u\|_{t_{K}, K}$.
The constant $C$ is independent of $u, h_{K}$ and $p_{K}$, but depends on the shape regularity of the mesh.

In order to interpolate vector function, we define
Definition 2.1 For $u=\left(u_{1}, u_{2}, u_{3}\right)$ we define
$\Pi_{p}^{h}: H^{t}\left(\nabla \times, \Pi_{h}\right) \longrightarrow \Sigma_{h} \quad$ by $\quad \Pi_{p}^{h}(u)=\left(\pi_{p}^{h}\left(u_{1}\right), \pi_{p}^{h}\left(u_{2}\right), \pi_{p}^{h}\left(u_{3}\right)\right.$ with $\pi_{p}^{h}$ is defined by $\pi_{p}^{h}(u)_{\mid K}=\pi_{p_{K}}^{h_{K}}\left(u_{\mid K}\right)$ where $\pi_{p_{K}}^{h_{K}}$ is given by the previous Proposition.

### 2.2 Model problem

The interior penalty finite element approximation to $u$ is to find $U: I \longrightarrow \Sigma_{h}$ such that

$$
\begin{equation*}
\left(U_{t t}, v\right)+B(U, v)=(f, v) \quad \forall v \in \Sigma_{h}, \quad U(0)=\boldsymbol{\Pi}_{p}^{h}\left(u_{0}\right), \quad U_{t}(0)=\boldsymbol{\Pi}_{p}^{h}\left(u_{1}\right) \tag{3}
\end{equation*}
$$

Upon choice of a basis for $\Sigma_{h}$ and the data $f$, (3) determines $U$ as the only solution to an initial value problem for a linear system of ordinary differential equations. Note that, if $u$ is the exact solution of (1)-(2), then $u$ satisfies the first equation in (3) and thus the problem is consistent.
We now analyse the proposed procedure by the method of energy estimates.

### 2.2.1 A priori error estimate

In this Section, $u$ denotes the exact solution of (1)-(2) and $U$ the discrete solution of (3). $C$ is generic constant independent of $h$ and $p$ which takes different values at the different places and depends on $\alpha, \tilde{C}$ the coercivity constants of the form $B, t^{*}$ and $\Omega$.
Let $\zeta=U-u$, then $\zeta$ satisfies

$$
\left(\zeta_{t t}, v\right)+B(\zeta, v)=0 \quad \forall v \in \Sigma_{h}
$$

Decompose $\zeta$ as $\mu-\nu$ where $\mu=\boldsymbol{\Pi}_{p}^{h}(u)-u$ and $\nu=\boldsymbol{\Pi}_{p}^{h}(u)-U$. Thus

$$
\left(\nu_{t t}, v\right)+B(\nu, v)=\left(\mu_{t t}, v\right)+B(\mu, v) \quad \forall v \in \Sigma_{h}
$$

Since $\nu_{t}(t) \in \Sigma_{h}$, we can set $v=\nu_{t}(t)$, obtaining

$$
\begin{aligned}
\frac{1}{2} \frac{d}{d t}\left\|\nu_{t}(t)\right\|^{2}+\frac{1}{2} \frac{d}{d t} B(\nu(t), \nu(t)) & =\left(\mu_{t t}(t), \nu_{t}(t)\right)+B\left(\mu(t), \nu_{t}(t)\right) \\
& \leq \frac{1}{2}\left\|\mu_{t t}(t)\right\|^{2}+\frac{1}{2}\left\|\nu_{t}(t)\right\|^{2}+B\left(\mu(t), \nu_{t}(t)\right) .
\end{aligned}
$$

So

$$
\frac{d}{d t}\left\|\nu_{t}(t)\right\|^{2}+\frac{d}{d t} B(\nu(t), \nu(t)) \leq\left\|\mu_{t t}(t)\right\|^{2}+\left\|\nu_{t}(t)\right\|^{2}+2 B\left(\mu(t), \nu_{t}(t)\right)
$$

Since $\nu_{t}(0)=\nu(0)=0$, integration over $[0, t] \subset I$, yields

$$
\left\|\nu_{t}(t)\right\|^{2}+B(\nu(t), \nu(t)) \leq\left\|\mu_{t t}\right\|_{L^{2}\left(L^{2}\right)}^{2}+\int_{0}^{t}\left\|\nu_{t}(t)\right\|^{2} d t+2 \int_{0}^{t} B\left(\mu(t), \nu_{t}(t)\right) d t
$$

The final term may be integrated by parts in time. Hence,

$$
2 \int_{0}^{t} B\left(\mu(t), \nu_{t}(t)\right) d t \leq 2|B(\mu(t), \nu(t))|+2 \int_{0}^{t}\left|B\left(\mu_{t}(t), \nu(t)\right)\right| d t
$$

Therefore, we can apply the coercivity and continuity of $B$ to get

$$
\begin{aligned}
& \left\|\nu_{t}(t)\right\|^{2}+\alpha\|\nu(t)\|_{h}^{2}+\tilde{C} J^{\sigma}(\nu(t), \nu(t)) \\
& \leq\left\|\mu_{t t}\right\|_{L^{2}\left(L^{2}\right)}^{2}+\int_{0}^{t}\left\|\nu_{t}(t)\right\|^{2} d t+C\|\nu(t)\|_{h}\|\mu(t)\|_{h}+2 \int_{0}^{t}\left|B\left(\mu_{t}(t), \nu(t)\right)\right| d t \\
& \leq\left\|\mu_{t t}\right\|_{L^{2}\left(L^{2}\right)}^{2}+\int_{0}^{t}\left\|\nu_{t}(t)\right\|^{2} d t+C\|\mu(t)\|_{h}^{2}+\frac{\alpha}{2}\|\nu(t)\|_{h}^{2}+C \int_{0}^{t}\left(\left\|\mu_{t}(t)\right\|_{h}^{2}+\|\nu(t)\|_{h}^{2}\right) d t \\
& \leq C\left(\left\|\mu_{t t}\right\|_{L^{2}\left(L^{2}\right)}^{2}+\sup _{t \in I}\|\mu(t)\|_{h}^{2}+\int_{0}^{t^{*}}\left\|\mu_{t}(t)\right\|_{h}^{2} d t\right)+\frac{\alpha}{2}\|\nu(t)\|_{h}^{2}+C \int_{0}^{t}\left(\left\|\nu_{t}(t)\right\|^{2}+\|\nu(t)\|_{h}^{2}\right) d t .
\end{aligned}
$$

In particular,

$$
\begin{aligned}
\left\|\nu_{t}(t)\right\|^{2} & +\|\nu(t)\|_{h}^{2} \\
& \left.\leq C\left(\left\|\mu_{t t}\right\|_{L^{2}\left(L^{2}\right)}^{2}+\sup _{t \in I}\|\mu(t)\|_{h}^{2}+\int_{0}^{t^{*}}\left\|\mu_{t}(t)\right\|_{h}^{2} d t\right)+C \int_{0}^{t}\left(\left\|\nu_{t}(t)\right\|^{2}+\|\nu(t)\|_{h}^{2}\right) d t\right) .
\end{aligned}
$$

As this holds for all $t \in I$, Gronwall's Lemma implies that

$$
\left\|\nu_{t}(t)\right\|^{2}+\|\nu(t)\|_{h}^{2} \leq C\left(\left\|\mu_{t t}\right\|_{L^{2}\left(L^{2}\right)}^{2}+\sup _{t \in I}\|\mu(t)\|_{h}^{2}+\int_{0}^{t^{*}}\left\|\mu_{t}(t)\right\|_{h}^{2} d t\right)
$$

Since $\zeta=\mu-\nu$,
$\left\|\zeta_{t}(t)\right\|^{2}+\|\zeta(t)\|_{h}^{2} \leq C\left(\left\|\mu_{t t}\right\|_{L^{2}\left(L^{2}\right)}^{2}+\sup _{t \in I}\|\mu(t)\|_{h}^{2}+\int_{0}^{t^{*}}\left\|\mu_{t}(t)\right\|_{h}^{2} d t+\left\|\mu_{t}\right\|_{L^{\infty}\left(L^{2}\right)}^{2}\right)$.
Thus, error bounds for the finite element approximation to the true solution reduce to the error bounds for the piecewise polynomial interpolant. Thus, we start by estimating $\left\|u-\boldsymbol{\Pi}_{p}^{h}(u)\right\|_{h}$ where $\boldsymbol{\Pi}_{p}^{h}$ is defined after Proposition 2.4. By using Proposition 2.4 and the definition of $\|\cdot\|_{h}$, we obtain the following estimates
$\left\|u-\Pi_{p}^{h}(u)\right\|_{h}^{2} \leq C \sum_{K \in \Pi_{h}} \frac{h_{K}^{2 \mu_{K}-2}}{p_{K}^{2 t_{K}-3}}\|u\|_{t_{K}, K}^{2} \quad$ and $\left\|u-\pi_{p_{K}}^{h_{K}}(u)\right\|_{q, K} \leq C \frac{h_{K}^{\mu_{K}-q}}{p_{K}^{t_{K}-q}}\|u\|_{t_{K}, K} \quad \forall 0 \leq q \leq t_{K}$.
By using the previous estimates, we can get the following result

Proposition 2.4 Let $\mu_{K}=\min \left(p_{K}+1, t_{K}\right)$ and $u$ be the exact solution of (1)-(2). Suppose that $u_{\mid K} \in C^{2}\left(I, H^{t_{K}}(K)^{3}\right), \forall K \in \Pi_{h}$ with $t_{K} \geq 2$. Let $U$ the discrete solution of (3). Then, the error $\zeta=U-u$ satisfies

$$
\begin{aligned}
& \left\|\zeta_{t}(t)\right\|^{2}+\|\zeta(t)\|_{h}^{2} \\
& \leq C \sum_{K \in \Pi_{h}} \frac{h_{K}^{2 K_{K}-2}}{p_{K}^{2 t}-3}\left(\left\|u_{t t}\right\|_{L^{2}\left(H^{t} K(K)^{3}\right)}^{2}+\|u\|_{L^{\infty}\left(H^{\left.t_{K}(K)^{3}\right)}\right.}^{2}+\left\|u_{t}\right\|_{L^{2}\left(H^{\left.t_{K}(K)^{3}\right)}\right.}^{2}+\left\|u_{t}\right\|_{L^{\infty}\left(H^{\left.t_{K}(K)^{3}\right)}\right.}^{2}\right) .
\end{aligned}
$$

## References

[1] I. Perugia and D. Schötzau, "The $h p$-Local Discontinuous Galerkin method for the Low-Frequency Time-Harmonic Maxwell's Equations." Math. Comp,2003, no. 243 pp 1179-1214.
[2] A. Zaghdani and C. Daveau "Two new discrete inequalities of Poincaré-Friedrichs on discontinuous spaces for Maxwell's equations." C. R. Acad. Sci. Paris. Ser.I. 342, 2006, pp 29-32.

