



HAL
open science

Determination of the body force of a two-dimensional isotropic elastic body

Dang Duc Trong, Alain Pham Ngoc Dinh, Phan Thanh Nam, Truong Trung Tuyen

► **To cite this version:**

Dang Duc Trong, Alain Pham Ngoc Dinh, Phan Thanh Nam, Truong Trung Tuyen. Determination of the body force of a two-dimensional isotropic elastic body. 2007. hal-00146732v1

HAL Id: hal-00146732

<https://hal.science/hal-00146732v1>

Preprint submitted on 15 May 2007 (v1), last revised 15 Jul 2008 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Determination of the body force of a two-dimensional isotropic elastic body

DANG DUC TRONG^a PHAM NGOC DINH ALAIN^b,
PHAN THANH NAM^a TRUONG TRUNG TUYEN^c AND

^a*Mathematics Department, HoChiMinh City National University, Viet Nam*

^b*Mathematics Department, Mapmo UMR 6628, BP 67-59, 45067 Orleans cedex, France*

^c*Department of mathematics, Indiana University, Rawles Hall , Bloomington, IN 47405*

Abstract

Let Ω represent a two-dimensional isotropic elastic body. We consider the problem of determining the body force F whose form $\varphi(t)(f_1(x), f_2(x))$ with φ be given inexactly. The problem is nonlinear and ill-posed. Using the Fourier transform, the methods of Tikhonov's regularization and truncated integration, we construct a regularized solution from the data given inexactly and derive the explicitly error estimate.

MSC 2000: 35L20, 35R30, 42B10, 70F07, 74B05.

Key words: body force, elastic body, Fourier transform, ill-posed problem, Tikhonov's regularization, truncated integration.

1. Introduction

Let $\Omega = (0, 1) \times (0, 1)$ represent a two-dimensional isotropic elastic body. For each $x := (x_1, x_2) \in \Omega$, we denote by $u = (u_1(x, t), u_2(x, t))$ the displacement, where u_j is the displacement in the x_j - direction, for all $j \in \{1, 2\}$. As known, u satisfies the Lamé system (see, e.g., [1, 2])

$$\frac{\partial^2 u}{\partial t^2} = \mu \Delta u + (\lambda + \mu) \nabla (\operatorname{div}(u)) + F$$

where $F := (F_1, F_2)$ is the body force, $\operatorname{div}(u) = \nabla \cdot u = \partial u_1 / \partial x_1 + \partial u_2 / \partial x_2$, and λ, μ are Lamé constants. We shall assume that the boundary of the elastic body is clamped and the initial conditions are given.

In this paper, we shall consider the problem of determining the body force F . The problem is a kind of inverse source problems. The inverse source problems are investigated in many aspects such as the uniqueness, the stability and the regularization. There are many papers devoted to the uniqueness and the stability problem. In [7], Isakov discussed the problem of finding a pair of functions (u, f) satisfying

$$cu_{tt} - \Delta u = f$$

where f is independent of t . He proved that using some preassumptions on f , from the *final overdetermination*

$$u(x, T) = h(x)$$

, we get the uniqueness of (u, f) .

As shown in [9], the body force (in the form $\phi(t)f(x)$) will be defined uniquely from an observation of surface stress (the *lateral overdetermination*) given on a suitable boundary of $\Omega \times (0, T)$. In the paper, the authors also gave an abstract formula of reconstruction.

Another inverse source problem is one of finding the heat source $F(x, t, u)$ satisfying

$$u_t - \Delta u = F.$$

The problem was considered intensively in the last century. The problem with the final overdetermination was studied by Tikhonov in 1935 (see [8]). He proved the uniqueness of problem with prescribed lateral and final data. In the last three decades, the problem is considered by many authors (see [3, 4, 11, 12, 13, 14]). Although we have many works on the uniqueness and the stability of inverse source problems, the literature on the regularization problem is quite scarce. Very recently, in [3, 4], the authors considered the regularization problem under both the lateral and the final overdetermination. The ideas of using the Fourier transform and truncated integration in the two papers are used in the present paper. We also consider the regularization problem under the final data and prescribed surface stress.

To get the lateral overdetermination, some mechanical arguments are in order. Let σ_1, σ_2, τ be the stresses (see [1, 2]) defined by

$$\begin{aligned} \tau &= \mu \left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) \\ \sigma_j &= \lambda \operatorname{div}(u) + 2\mu \frac{\partial u_j}{\partial x_j}, \quad j \in \{1, 2\} \end{aligned}$$

We shall assume that the surface stress is given on the boundary of the body, i.e.,

$$\begin{pmatrix} \sigma_1 & \tau \\ \tau & \sigma_2 \end{pmatrix} \begin{pmatrix} n_1 \\ n_2 \end{pmatrix} = \begin{pmatrix} X_1 \\ X_2 \end{pmatrix}$$

where $X = (X_1, X_2)$ is given on $\partial\Omega$, and $n = (n_1, n_2)$ is the outward unit normal vector of $\partial\Omega$.

As discussed, our problem is severely ill-posed. Hence, to simplify the problem, a preassumption on the form of the body force is needed. We shall use the separable form force as in [9]

$$(F_1(x, t), F_2(x, t)) = \varphi(t)(f_1(x), f_2(x))$$

where φ is given inexactly. The form is issued from an approximated model for elastic wave generated from a point dislocation source (see, e.g., [9, 10]). But, since φ is inexact, our problem is nonlinear. Moreover, the problem is still ill-posed because the measured data is

not only inexact but also non-smooth.

Precisely, we consider the problem of identifying a pair of functions (u, f) satisfying the system:

$$\frac{\partial^2 u_j}{\partial t^2} = \mu \Delta u_j + (\lambda + \mu) \frac{\partial}{\partial x_j} \operatorname{div}(u) + \varphi(t) f_j(x), \forall j \in \{1, 2\} \quad (1)$$

for $(x, t) \in \Omega \times (0, T)$, where μ, λ are real constants satisfying $\mu > 0$ and $\lambda + 2\mu > 0$. Since the boundary of the elastic body is clamped, the displacement $u = (u_1, u_2)$ satisfies the boundary condition

$$(u_1(x, t), u_2(x, t)) = (0, 0), \quad x \in \partial\Omega \quad (2)$$

In addition, the initial and final displacement are given in Ω

$$\begin{cases} (u_1(x, 0), u_2(x, 0)) = (u_{01}(x), u_{02}(x)) \\ \left(\frac{\partial u_1}{\partial t}(x, 0), \frac{\partial u_2}{\partial t}(x, 0) \right) = (u_{01}^*(x), u_{02}^*(x)) \\ (u_1(x, T), u_2(x, T)) = (u_{T1}(x), u_{T2}(x)) \end{cases} \quad (3)$$

Finally, the surface stress is given on $\partial\Omega$

$$\begin{cases} n_1 \sigma_1 + n_2 \tau = X_1 \\ n_2 \sigma_2 + n_1 \tau = X_2 \end{cases} \quad (4)$$

We shall assume that the data of the system (1) – (4)

$$I = (\varphi, X, u_0, u_0^*, u_T) \in (L^1(0, T), (L^1(0, T, L^1(\partial\Omega)))^2, (L^1(\Omega))^2, (L^1(\Omega))^2, (L^1(\Omega))^2)$$

are given inexactly since they are results of experimental measurements. The system (1)–(4) usually has no solution; moreover, even if the solution exists, it does not depend continuously on the given data. Hence, a regularization is in order. Denoting by I_{ex} the exact data, which are probably unknown, corresponding to an exact solution (u_{ex}, f_{ex}) of the system (1) – (4), from the inexact data I_ε approximating I_{ex} , we shall construct a regularized solution f_ε approximating f_{ex} .

In fact, using the Fourier transform, we shall reduce our problem to finding the solutions of the binomial equations whose binomial term is an entire function (see Lemma 1). In this case, the problem is unstable in the neighborhood of zeros of the entire function. The zeroes can be seen as singular values. Using the method of Tikhonov's regularization and truncated integration, we shall eliminate the singular values to regularize our problem. Error estimates are given.

The remainder of the paper is divided into two sections. In Section 2, we shall set some notations and state our main results. In Section 3, we give the proofs of the results.

2. Notations and main results

We recall that $\Omega = (0, 1) \times (0, 1)$. We always assume that the data $I = (\varphi, X, u_0, u_T, u_T^*)$ belong to

$$(L^1(0, T), (L^1(0, T, L^1(\partial\Omega)))^2, (L^1(\Omega))^2, (L^1(\Omega))^2, (L^1(\Omega))^2)$$

For all $\xi = (\xi_1, \xi_2), \zeta = (\zeta_1, \zeta_2) \in \mathbb{R}^2$, we set $\xi \cdot \zeta = \xi_1\zeta_1 + \xi_2\zeta_2$ and $|\xi| = \sqrt{\xi \cdot \xi}$.

We first have the following lemma.

Lemma 1. *If $u \in (C^2([0, T]; L^2(\Omega)) \cap L^2(0, T; H^2(\Omega)))^2$, $f \in (L^2(\Omega))^2$ satisfy (1) – (4) corresponding the data I , then for all $\alpha = (\alpha_1, \alpha_2) \in \mathbb{R}^2 \setminus \{0\}$, we have*

$$2D(I) \cdot \int_{\Omega} f_j(x) \cdot \cos(\alpha \cdot x) dx = g_j(I), \quad \forall j \in \{1, 2\}$$

where

$$D(I) = D_1(I) \cdot D_2(I), g_j(I) = \frac{2}{|\alpha|^2} (\alpha_j D_2(I) h_0 + D_1(I) h_j)$$

with

$$D_1(I) = \int_0^T \varphi(T-t) \sin(\sqrt{\lambda + 2\mu} |\alpha| t) dt, D_2(I) = \int_0^T \varphi(T-t) \sin(\sqrt{\mu} |\alpha| t) dt$$

$$\begin{aligned} h_0(I) &= -\sin(\sqrt{\lambda + 2\mu} |\alpha| T) \cdot \int_{\Omega} (\alpha \cdot u_0^*) \cdot \cos(\alpha \cdot x) dx \\ &\quad + \sqrt{\lambda + 2\mu} \cdot |\alpha| \cdot \int_{\Omega} (\alpha \cdot u_T) \cdot \cos(\alpha \cdot x) dx \\ &\quad - \sqrt{\lambda + 2\mu} \cdot |\alpha| \cdot \cos(\sqrt{\lambda + 2\mu} |\alpha| T) \cdot \int_{\Omega} (\alpha \cdot u_0) \cdot \cos(\alpha \cdot x) dx \\ &\quad - \int_0^T \int_{\partial\Omega} \sin(\sqrt{\lambda + 2\mu} |\alpha| (T-t)) (\alpha \cdot X) \cdot \cos(\alpha \cdot x) d\omega dt \end{aligned}$$

$$\begin{aligned} h_j(I) &= -\sin(\sqrt{\mu} |\alpha| T) \cdot \int_{\Omega} (|\alpha|^2 u_{0j}^* - \alpha_j (\alpha \cdot u_0^*)) \cdot \cos(\alpha \cdot x) dx \\ &\quad + \sqrt{\mu} \cdot |\alpha| \cdot \int_{\Omega} (|\alpha|^2 u_{Tj} - \alpha_j (\alpha \cdot u_T)) \cdot \cos(\alpha \cdot x) dx \\ &\quad - \sqrt{\mu} \cdot |\alpha| \cdot \cos(\sqrt{\mu} |\alpha| T) \cdot \int_{\Omega} (|\alpha|^2 u_{0j} - \alpha_j (\alpha \cdot u_0)) \cdot \cos(\alpha \cdot x) dx \\ &\quad - \int_0^T \int_{\partial\Omega} \sin(\sqrt{\mu} |\alpha| (T-t)) (|\alpha|^2 X_j - \alpha_j (\alpha \cdot X)) \cdot \cos(\alpha \cdot x) d\omega dt, \forall j \in \{1, 2\}. \end{aligned}$$

From Lemma 1, we consider the function

$$D(I) = \int_0^T \varphi(T-t) \sin(\sqrt{\lambda + 2\mu} |\alpha| t) dt. \int_0^T \varphi(T-t) \sin(\sqrt{\mu} |\alpha| t) dt$$

The problem is unstable in the neighborhood of zeros of this function. However, from the properties of analytic function, we can show that if $\varphi \not\equiv 0$ then this function differ from 0 for almost every where in R^3 . Furthermore, using the idea of Theorem 4 in [5], we get the following lemma.

Lemma 2. *Let τ, q be positive constants, $\varphi_0 \in L^1(0, T) \setminus \{0\}$ and $D(\varphi_0, \tau) : R^2 \rightarrow R$*

$$D(\varphi_0, \tau)(\alpha) = \int_0^T \varphi_0(t) \sin(\sqrt{\tau} |\alpha| t) dt$$

Then $D(\varphi_0, \tau) \neq 0$ for a.e $\alpha \in R^2$. Moreover, if we put

$$R_\varepsilon = \frac{q}{9eT} \cdot \frac{\ln(\varepsilon^{-1})}{\ln(\ln(\varepsilon^{-1}))}, \quad \forall \varepsilon > 0$$

then the Lebesgue measure of the set

$$B(\varphi_0, \tau, \varepsilon) = \{\alpha \in B(0, R_\varepsilon), |D(\varphi_0, \tau)(\alpha)| \leq \varepsilon^q\}$$

is less than R_ε^{-1} for $\varepsilon > 0$ small enough, where $B(0, R_\varepsilon)$ is the open ball in R^2 .

Lemma 1 and Lemma 2 imply immediately the uniqueness result.

Theorem 1. *Let $u, u^* \in (C^2([0, T]; L^2(\Omega)) \cap L^2(0, T; H^2(\Omega)))^2$, $f, f^* \in (L^2(\Omega))^2$. If (u, f) , (u^*, f^*) satisfy (1) – (4) corresponding the same data I , and $\varphi \not\equiv 0$, then*

$$(u, f) = (u^*, f^*)$$

Let (u_{ex}, f_{ex}) be the exact solution of the system (1) – (4) corresponding the exact data $I_{ex} = (\varphi_{ex}, X_{ex}, u_0^{ex}, u_0^{*ex}, u_T^{ex})$. Notice that, if we assume

$$u_{ex} \in (C^2([0, T]; L^2(\Omega)) \cap L^2([0, T]; H^2(\Omega)))^2, f_{ex} \in (L^2(\Omega))^2, \varphi_{ex} \in L^1(0, T) \setminus \{0\} \quad (5)$$

then for all $j \in \{1, 2\}$,

$$F(\tilde{f}_{jex})(\alpha) = 2 \int_{\Omega} f_{jex}(x) \cos(\alpha \cdot x) dx = \frac{g_j(I_{ex})}{D(I_{ex})}$$

for a.e $\alpha \in R^2$, where g_j, D are defined by Lemma 1, $\tilde{f}_{jex} : R^2 \rightarrow R$ is defined by $\tilde{f}_{jex}(x) = \chi(\Omega)f_{jex}(x) + \chi(-\Omega)f_{jex}(-x)$, and F is the Fourier transform in R^2 .

From approximate data $I_\varepsilon = (\varphi, X, u_0, u_0^*, u_T)$ satisfying

$$\begin{aligned} \|\varphi - \varphi_{ex}\|_{L^1(0,T)} &\leq \varepsilon, \|X_j - X_j^{ex}\|_{L^1(0,T,L^1(\partial\Omega))} \leq \varepsilon, \|u_{0j} - u_{0j}^{ex}\|_{L^1(\Omega)} \leq \varepsilon \\ \|u_{0j}^* - u_{0j}^{*ex}\|_{L^1(\Omega)} &\leq \varepsilon, \|u_{Tj} - u_{Tj}^{ex}\|_{L^1(\Omega)} \leq \varepsilon, \quad \forall j \in \{1, 2\} \end{aligned} \quad (6)$$

, we construct a regularized solution $f_\varepsilon = (f_{1\varepsilon}, f_{2\varepsilon})$ whose Fourier transform is

$$F(f_{j\varepsilon})(\alpha) = \chi(B(0, R_\varepsilon)) \cdot \frac{g_j(I_\varepsilon) \cdot D(I_\varepsilon)}{\delta_\varepsilon + (D(I_\varepsilon))^2}, \forall \alpha \in R^2 \setminus \{0\}$$

where

$$q = \frac{1}{7}, \delta_\varepsilon = \varepsilon^{\frac{1+6q}{2}}, R_\varepsilon = \frac{q}{9eT} \cdot \frac{\ln(\varepsilon^{-1})}{\ln(\ln(\varepsilon^{-1}))} \quad (7)$$

We have two regularization results.

Theorem 2. *Let (u_{ex}, f_{ex}) be the exact solution of the system (1) – (4) corresponding the exact data I_{ex} , and (5) hold. Then from the given data I_ε satisfying (6), we can construct a regularized solution $f_\varepsilon \in (C(\overline{\Omega}))^2$ such that*

$$\lim_{\varepsilon \rightarrow 0} \|f_{j\varepsilon} - f_{jex}\|_{L^2(\Omega)} = 0, \quad \forall j \in \{1, 2\}$$

If we assume, in addition, that $f_{ex} \in (H^1(\Omega))^2$, then

$$\|f_{j\varepsilon} - f_{jex}\|_{L^2(\Omega)}^2 \leq 63eT \left(66 \|f_{jex}\|_{H^1(\Omega)}^2 + (2\pi)^{-2} \right) \cdot \frac{\ln(\ln(\varepsilon^{-1}))}{\ln(\varepsilon^{-1})}, \quad \forall j \in \{1, 2\}$$

for $\varepsilon > 0$ small enough.

Theorem 3. *Let (u_{ex}, f_{ex}) be the exact solution of the system (1) – (4) corresponding the exact data I_{ex} , and (5) hold. We assume, in addition, that*

$$\int_{R^2} \left| \int_{\Omega} f_{jex}(x) \cdot \cos(\alpha \cdot x) dx \right| d\alpha < \infty, \quad \forall j \in \{1, 2\}$$

Then from the given data I_ε satisfying (6), we can construct a regularized solution $f_\varepsilon \in (C(\overline{\Omega}))^2$, which coincides the one in Theorem 2, such that

$$\lim_{\varepsilon \rightarrow 0} \|f_{j\varepsilon} - f_{jex}\|_{L^\infty(\Omega)} = 0, \quad \forall j \in \{1, 2\}$$

3. Proofs of the results

Proof of Lemma 1

Proof. Let $\alpha = (\alpha_1, \alpha_2) \in R^2$ and $G = \cos(\alpha \cdot x)$. Notice that the j -th equation of the system (1) can rewrite

$$\frac{\partial^2 u_j}{\partial t^2} = \frac{\partial \sigma_j}{\partial x_j} + \frac{\partial \tau}{\partial x_k} + \varphi(t) f_j(x), \quad \{j, k\} = \{1, 2\}$$

Getting the inner product (in $L^2(\Omega)$) of the equation and G and using the condition (2), for $\{j, k\} = \{1, 2\}$, we get

$$\begin{aligned} \frac{d}{dt^2} \int_{\Omega} u_j G &= \int_{\partial\Omega} (n_j \sigma_j + n_k \tau) G d\omega - \int_{\Omega} \sigma_j \frac{\partial G}{\partial x_j} dx - \int_{\Omega} \tau \frac{\partial G}{\partial x_k} dx + \varphi(t) \int_{\Omega} f_j G dx \\ &= \int_{\partial\Omega} X_j G d\omega - \mu |\alpha|^2 \int_{\Omega} u_j G dx - (\lambda + \mu) \alpha_j \int_{\Omega} (\alpha \cdot u) G dx + \varphi(t) \int_{\Omega} f_j G dx \end{aligned} \quad (8)$$

Multiplying (8) by α_j , then getting the sum for $j = 1, 2$, we obtain

$$\frac{d}{dt^2} \int_{\Omega} (\alpha \cdot u) G dx = \int_{\partial\Omega} (\alpha \cdot X) G d\omega - (\lambda + 2\mu) |\alpha|^2 \int_{\Omega} (\alpha \cdot u) G dx + \varphi(t) \int_{\Omega} (\alpha \cdot f) G dx \quad (9)$$

Multiplying (8) by $|\alpha|^2$ and multiplying (9) by $-\alpha_j$, then getting the sum of them, we have

$$\begin{aligned} \frac{d}{dt^2} \int_{\Omega} \left(|\alpha|^2 u_j - \alpha_j \cdot (\alpha \cdot u) \right) G dx &= \int_{\partial\Omega} \left(|\alpha|^2 X_j - \alpha_j \cdot (\alpha \cdot X) \right) G dx \\ -\mu |\alpha|^2 \int_{\Omega} \left(|\alpha|^2 u_j - \alpha_j \cdot (\alpha \cdot u) \right) G dx &+ \varphi(t) \int_{\Omega} \left(|\alpha|^2 f_j - \alpha_j \cdot (\alpha \cdot f) \right) G dx \end{aligned} \quad (10)$$

We consider (9) and (10) as the differential equations whose form

$$y'' + \eta^2 y = h(t) \quad (11)$$

where η is a real constant and $y(0)$, $y'(0)$, $y(T)$ are given. Getting the inner product (in $L^2(0, T)$) of (11) and $\sin(\eta(T-t))$, we have

$$-y'(0) \sin(\eta T) + \eta y(T) - \eta y(0) \cos(\eta T) = \int_0^T h(T-t) \sin(\eta t) dt \quad (12)$$

Applying (12) to (9) with $\eta = \sqrt{(\lambda + 2\mu)} |\alpha|$ and $y = \int_{\Omega} (\alpha \cdot u) \cdot G dx$, we get

$$D_1(I) \cdot \int_{\Omega} (\alpha \cdot f) \cdot G dx = h_0(I) \quad (13)$$

where $D_1(I), h_0(I)$ are defined by Lemma 1.

Similarly, applying (12) to (10) with $\eta = \sqrt{\mu}|\alpha|$ and $y = \int_{\Omega} (|\alpha|^2 u_j - \alpha_j \cdot (\alpha \cdot u)) \cdot G dx$, we get

$$D_2(I) \cdot \int_{\Omega} (|\alpha|^2 f_j - \alpha_j (\alpha \cdot f)) \cdot G dx = h_j(I), \quad \forall j \in \{1, 2\} \quad (14)$$

where $D_2(I), h_j(I)$ are defined by Lemma 1.

Multiplying (13) by $\alpha_j D_2(I)$ and multiplying (14) by $D_1(I)$, then getting the sum of them, we obtain the result of Lemma 1. \square

Proof of Lemma 2

Proof. Put $\tilde{\varphi}_0 : R \rightarrow R$

$$\tilde{\varphi}_0(t) = \frac{1}{2} \begin{cases} \varphi_0(t) & t \in (0, T) \\ -\varphi_0(-t) & t \in (-T, 0) \\ 0 & t \notin (-T, T) \end{cases}$$

and $\phi : C \rightarrow C$

$$\phi(z) = \int_{-\infty}^{\infty} e^{-itz} \tilde{\varphi}_0(t) dt = \int_{-T}^T e^{-itz} \tilde{\varphi}_0(t) dt$$

Then ϕ is an entire function and $D(\varphi_0, \tau)(\alpha) = i\phi(\sqrt{\tau}|\alpha|)$. Because $\tilde{\varphi}_0 \not\equiv 0$, its Fourier transform (in R) does not coincide 0. Therefore, there exists $z_0 \in R$ such that $|\phi(z_0)| = C_1 > 0$. Thus $\phi \not\equiv 0$. Since ϕ is an entire function, its zeros set is either finite or countable. Consequently, $D(\varphi_0, \tau)(\alpha) \neq 0$ for a.e $\alpha \in R^2$.

To estimate the measure of $B(\varphi_0, \tau, \varepsilon)$, we shall use the following result (see Theorem 4 of §11.3 in [6]).

Lemma 3. *Let $f(z)$ be a function analytic in the disk $\{z : |z| \leq 2eR\}$, $|f(0)| = 1$, and let η be an arbitrary small positive number. Then the estimate*

$$\ln |f(z)| > -\ln\left(\frac{15e^3}{\eta}\right) \cdot \ln(M_f(2eR))$$

is valid everywhere in the disk $\{z : |z| \leq R\}$ except a set of disks (C_j) with sum of radii $\sum r_j \leq \eta R$. Where $M_f(r) = \max_{|z|=r} |f(z)|$.

Returning Lemma 2, we put $\phi_1 : C \rightarrow C$

$$\phi_1(z) = \frac{\phi(z + z_0)}{C_1}$$

Then ϕ_1 is an entire function, $\phi_1(0) = 1$, and for all $z \in C, |z| \leq 2eR$,

$$C_1 |\phi_1(z)| = \left| \int_{-T}^T e^{-it(z+z_0)} \tilde{\varphi}_0(t) dt \right| \leq e^{2eRT} \cdot \int_{-T}^T |\tilde{\varphi}_0(t)| dt = e^{2eRT} \|\varphi_0\|_{L^1(0,T)}$$

For $\varepsilon > 0$ small enough, applying Lemma 3 with $R = \frac{4}{3}R_\varepsilon$ and $\eta = \frac{\sqrt{\tau}}{8\pi R_\varepsilon^3}$, we get

$$\begin{aligned} \ln |\phi_1(z)| &> - \left[3 \ln R_\varepsilon + \ln\left(\frac{8\pi}{\sqrt{\tau}}\right) + \ln(15e^3) \right] \cdot \left[\frac{8}{3} \cdot eTR_\varepsilon + \ln\left(\frac{\|\varphi_0\|_{L^1(0,T)}}{C_1}\right) \right] \\ &> -\frac{17}{2}T \cdot R_\varepsilon \ln R_\varepsilon > -q \ln(\varepsilon^{-1}) - \ln(C_1) = \ln\left(\frac{\varepsilon^q}{C_1}\right) \end{aligned}$$

for all $|z| \leq \frac{4}{3}R_\varepsilon$ except a set of disks $\{B(z_j, r_j)\}_{j \in J}$ with sum of radii $\sum r_i \leq \eta R = \frac{\sqrt{\tau}}{6\pi R_\varepsilon^2}$.

Consequently, for $\varepsilon > 0$ small enough, we have $|z_0| < \frac{1}{3}R_\varepsilon$ and $|\phi(z)| = C_1 \cdot |\phi_1(z - z_0)| \geq \varepsilon^q$ for all $|z| \leq R_\varepsilon$ except the set $\cup_{j \in J} B(z_j + z_0, r_j)$. Hence, $B(\varphi_0, \tau, \varepsilon)$ is contained in the set $\cup_{j \in J} B_j$, where

$$B_j = \{\alpha \in B(0, R_\varepsilon), |\sqrt{\tau}|\alpha - y_j| \leq r_j\}$$

with $y_j = Re(z_j + z_0)$.

If $y_j > \sqrt{\tau}R_\varepsilon + r_j$ then $B_j = \emptyset$. If $y_j \leq r_j$ then $B_j \subset B(0, \frac{2r_j}{\sqrt{\tau}})$, so $m(B_j) \leq \frac{4\pi r_j^2}{\tau}$. If $r_j < y_j \leq \sqrt{\tau}R_\varepsilon + r_j$ then

$$B_j \subset B(0, \frac{y_j + r_j}{\sqrt{\tau}}) \setminus B(0, \frac{y_j - r_j}{\sqrt{\tau}})$$

hence

$$m(B_j) \leq \frac{\pi(y_j + r_j)^2}{\tau} - \frac{\pi(y_j - r_j)^2}{\tau} = \frac{4\pi y_j r_j}{\tau} \leq \frac{4\pi(\sqrt{\tau}R_\varepsilon + r_j)r_j}{\tau}$$

Thus we get

$$\begin{aligned} m(B(\varphi, \tau, \varepsilon)) &\leq \sum \frac{4\pi(\sqrt{\tau}R_\varepsilon + r_j)r_j}{\tau} + \sum \frac{4\pi r_j^2}{\tau} \\ &\leq \frac{4\pi R_\varepsilon}{\sqrt{\tau}} \sum r_j + \frac{8\pi}{\tau} (\sum r_j)^2 \leq \frac{4\pi R_\varepsilon}{\sqrt{\tau}} \cdot \frac{\sqrt{\tau}}{6\pi R_\varepsilon^2} + \frac{8\pi}{\tau} \cdot \left(\frac{\sqrt{\tau}}{6\pi R_\varepsilon^2}\right)^2 < \frac{1}{R_\varepsilon} \end{aligned}$$

for $\varepsilon > 0$ small enough. The proof of Lemma 2 is completed. \square

Proof of theorem 1

Proof. Put $w = u - u^*$ and $v = f - f^*$ then (w, v) satisfies (1) – (4) corresponding the data

$$I = (\varphi, (0, 0), (0, 0), (0, 0), (0, 0))$$

Let $\tilde{v}_j : R^2 \rightarrow R$ be defined by $\tilde{v}_j(x) = \chi(\Omega)v_j(x) + \chi(-\Omega)v_j(-x)$. Lemma 1 implies that, for all $j \in \{1, 2\}$, for all $\alpha \in R^2 \setminus \{0\}$, we get

$$D(I).F(\tilde{v}_j)(\alpha) = 2D(I). \int_{\Omega} v_j(x) \cos(\alpha \cdot x) dx = g_j(I) = 0$$

Applying Lemma 2 with $\varphi_0(t) = \varphi(T - t)$, we get $D(I) \neq 0$ for a.e $\alpha \in R^2$. Therefore, $F(\tilde{v}_j) \equiv 0$, and it implies that $\tilde{v}_j \equiv 0$. Thus $v \equiv (0, 0)$. Hence, w satisfies that

$$\frac{\partial^2 w}{\partial t^2} = \mu \Delta w + (\lambda + \mu) \nabla (\operatorname{div}(w)) \quad (15)$$

Getting the inner product (in $(L^2(\Omega))^2$) of (15) and $\partial w / \partial t$, we have

$$\frac{1}{2} \cdot \frac{d}{dt} \sum_{j=1}^2 \left\| \frac{\partial w_j}{\partial t} \right\|_{L^2(\Omega)}^2 = -\frac{\mu}{2} \cdot \frac{d}{dt} \sum_{j=1}^2 \|\nabla w_j\|_{L^2(\Omega)}^2 - \frac{\lambda + \mu}{2} \cdot \frac{d}{dt} \|\operatorname{div}(w)\|_{L^2(\Omega)}^2$$

Integrating this equality in $(0, t)$, we get

$$\sum_{j=1}^2 \left\| \frac{\partial w_j}{\partial t} \right\|_{L^2(\Omega)}^2 + \mu \sum_{j=1}^2 \|\nabla w_j\|_{L^2(\Omega)}^2 + (\lambda + \mu) \|\operatorname{div}(w)\|_{L^2(\Omega)}^2 = 0 \quad (16)$$

for all $t \in (0, T)$. Using the condition (2), we have

$$\begin{aligned} \|\operatorname{div}(w)\|_{L^2(\Omega)}^2 &= \sum_{j=1}^2 \left\| \frac{\partial w_j}{\partial x_j} \right\|_{L^2(\Omega)}^2 + 2 \int_{\Omega} \frac{\partial w_1}{\partial x_1} \cdot \frac{\partial w_2}{\partial x_2} = \sum_{j=1}^2 \left\| \frac{\partial w_j}{\partial x_j} \right\|_{L^2(\Omega)}^2 + 2 \int_{\Omega} \frac{\partial w_1}{\partial x_2} \cdot \frac{\partial w_2}{\partial x_1} \\ &\leq \sum_{j=1}^2 \left\| \frac{\partial w_j}{\partial x_j} \right\|_{L^2(\Omega)}^2 + \left(\left\| \frac{\partial w_1}{\partial x_2} \right\|_{L^2(\Omega)}^2 + \left\| \frac{\partial w_2}{\partial x_1} \right\|_{L^2(\Omega)}^2 \right) = \sum_{j=1}^2 \|\nabla w_j\|_{L^2(\Omega)}^2 \end{aligned}$$

Since $\mu > 0$ and $\lambda + 2\mu > 0$, the above inequality implies that

$$\mu \sum_{j=1}^2 \|\nabla w_j\|_{L^2(\Omega)}^2 + (\lambda + \mu) \|\operatorname{div}(w)\|_{L^2(\Omega)}^2 \geq 0$$

From (16), we obtain $\partial w / \partial t = (0, 0)$. Since $w(x, 0) = (0, 0)$, the proof is completed. \square

To prove two main regularization results, we state and prove some preliminary lemmas.

Lemma 4. Let (u_{ex}, f_{ex}) be the exact solution of (1) – (4) corresponding the exact data I_{ex} satisfying (5), and the given data I_ε satisfying (6). Using notations of (7), we put

$$G_j(I_\varepsilon) = \chi(B(0, R_\varepsilon)) \cdot \frac{g_j(I_\varepsilon)D(I_\varepsilon)}{\delta_\varepsilon + (D(I_\varepsilon))^2}$$

Then for all $j \in \{1, 2\}$, we have $G_j(I_\varepsilon) \in L^1(R^2) \cap L^2(R^2)$; moreover, there exists a constant C_0 depend only on I_{ex} such that for all $\varepsilon \in (0, e^{-e})$,

$$\begin{aligned} \left| G_j(I_\varepsilon) - F(\tilde{f}_{jex}) \right| &\leq \chi(B(0, R_\varepsilon)) C_0 R_\varepsilon \varepsilon^{\frac{1-6q}{2}} \\ &\quad + 2\chi(B_\varepsilon) \|f_{jex}\|_{L^2(\Omega)} + \chi(R^2 \setminus B(0, R_\varepsilon)) \left| F(\tilde{f}_{jex}) \right| \end{aligned}$$

where $B_\varepsilon = \{\alpha \in B(0, R_\varepsilon), |D(I_{ex})(\alpha)| \leq \varepsilon^{2q}\}$.

Proof. First, we show that there exists a constant $C_2 > 0$ depend only on I_{ex} such that for all $\varepsilon \in (0, e^{-e})$, $r > r_0 = q/(9T)$, $j \in \{1, 2\}$,

$$\begin{aligned} \|D(I_\varepsilon)\|_{L^\infty(R^2)} &\leq C_2, \|D(I_\varepsilon) - D(I_{ex})\|_{L^\infty(R^2)} \leq C_2\varepsilon \\ \|g_j(I_{ex})\|_{L^\infty(B(0,r))} &\leq C_2r, \|g_j(I_\varepsilon) - g_j(I_{ex})\|_{L^\infty(B(0,r))} \leq C_2r\varepsilon \end{aligned}$$

Recall that $D_1(I), D_2(I), h_0(I), h_j(I)$ are defined by Lemma 1. For all $\alpha \in R^3$ we have

$$|D_k(I_{ex})| \leq \|\varphi_{ex}\|_{L^1(0,T)}, |D_k(I_\varepsilon) - D_k(I_{ex})| \leq \|\varphi_\varepsilon - \varphi_{ex}\|_{L^1(0,T)} \leq \varepsilon$$

for all $k \in \{1, 2\}$. Hence, $|D(I_{ex})| \leq \|\varphi_{ex}\|_{L^1(0,T)}^2$ and

$$\begin{aligned} |D(I_\varepsilon) - D(I_{ex})| &= |D_1(I_\varepsilon) - D_1(I_{ex})| \cdot |D_2(I_\varepsilon)| + |D_1(I_{ex})| \cdot |D_2(I_\varepsilon) - D_2(I_{ex})| \\ &\leq \varepsilon \cdot (\|\varphi_{ex}\|_{L^1(0,T)} + \varepsilon) + \|\varphi_{ex}\|_{L^1(0,T)} \cdot \varepsilon \leq (2\|\varphi_{ex}\|_{L^1(0,T)} + e^{-e}) \cdot \varepsilon \end{aligned}$$

A straightforward calculation show that, for all $\alpha \in B(0, r) \setminus \{0\}$, we have

$$\begin{aligned} |\alpha_j h_0(I_{ex})| &\leq C_3r |\alpha|^2, |\alpha_j (h_0(I_\varepsilon) - h_0(I_{ex}))| \leq C_3r |\alpha|^2 \varepsilon, \\ |h_j(I_{ex})| &\leq C_3r |\alpha|^2, |h_j(I_\varepsilon) - h_j(I_{ex})| \leq C_3r |\alpha|^2 \varepsilon \end{aligned}$$

for all $j \in \{1, 2\}$, where C_3 is a positive constant depending only on I_{ex} . Therefore,

$$|g_j(I_{ex})| \leq \frac{|\alpha_j h_0(I_{ex})|}{|\alpha|^2} \cdot |D_2(I_{ex})| + \frac{|h_j(I_{ex})|}{|\alpha|^2} \cdot |D_1(I_{ex})| \leq 2C_3 \|\varphi_{ex}\|_{L^1(0,T)} r$$

and

$$\begin{aligned} |g_j(I_\varepsilon) - g_j(I_{ex})| &\leq \frac{|\alpha_j (h_0(I_\varepsilon) - h_0(I_{ex}))|}{|\alpha|^2} \cdot |D_2(I_\varepsilon)| + \frac{|\alpha_j h_0(I_{ex})|}{|\alpha|^2} \cdot |D_2(I_\varepsilon) - D_2(I_{ex})| \\ &\quad + \frac{|h_j(I_\varepsilon) - h_j(I_{ex})|}{|\alpha|^2} \cdot |D_1(I_\varepsilon)| + \frac{|h_j(I_{ex})|}{|\alpha|^2} \cdot |D_1(I_\varepsilon) - D_1(I_{ex})| \\ &\leq C_3r\varepsilon \cdot (\|\varphi_{ex}\|_{L^1(0,T)}^2 + \varepsilon) + C_3r \cdot \varepsilon + C_3r\varepsilon \cdot (\|\varphi_{ex}\|_{L^1(0,T)}^2 + \varepsilon) + C_3r \cdot \varepsilon \\ &\leq 2C_3 \left(\|\varphi_{ex}\|_{L^1(0,T)}^2 + e^{-e} + 1 \right) r\varepsilon \end{aligned}$$

Returning Lemma 4, for all $j \in \{1, 2\}$, we get $G_j(I_\varepsilon) \in L^1(R^2) \cap L^2(R^2)$ because the support of $G_j(I_\varepsilon)$ is contained in $\overline{B(0, R_\varepsilon)}$ and $G_j(I_\varepsilon) \in L^\infty(R^2)$. Moreover,

$$\begin{aligned} \left| G_j(I_\varepsilon) - F(\tilde{f}_{jex}) \right| &\leq \chi(B(0, R_\varepsilon)) \left| \frac{g_j(I_\varepsilon) D(I_\varepsilon)}{\delta_\varepsilon + (D(I_\varepsilon))^2} - \frac{g_j(I_{ex}) D(I_{ex})}{\delta_\varepsilon + (D(I_{ex}))^2} \right| \\ &\quad + \chi(B(0, R_\varepsilon)) \left| \frac{g_j(I_{ex}) D(I_{ex})}{\delta_\varepsilon + (D(I_{ex}))^2} - \frac{g_j(I_{ex})}{D(I_{ex})} \right| + \chi(R^2 \setminus B(0, R_\varepsilon)) \cdot |F(\tilde{f}_{jex})| \end{aligned}$$

We shall estimate each of the terms of the right-hand side. We have

$$\begin{aligned} \left| \frac{g_j(I_\varepsilon) D(I_\varepsilon)}{\delta_\varepsilon + (D(I_\varepsilon))^2} - \frac{g_j(I_{ex}) D(I_{ex})}{\delta_\varepsilon + (D(I_{ex}))^2} \right| &\leq \frac{\delta_\varepsilon |g_j(I_\varepsilon) D(I_\varepsilon) - g_j(I_{ex}) D(I_{ex})|}{\left(\delta_\varepsilon + (D(I_\varepsilon))^2 \right) \left(\delta_\varepsilon + (D(I_{ex}))^2 \right)} \\ &\quad + \frac{|D(I_\varepsilon)| \cdot |D(I_{ex})| \cdot |g_j(I_\varepsilon) D(I_{ex}) - g_j(I_{ex}) D(I_\varepsilon)|}{\left(\delta_\varepsilon + (D(I_\varepsilon))^2 \right) \left(\delta_\varepsilon + (D(I_{ex}))^2 \right)} \\ &\leq \frac{|g_j(I_\varepsilon) D(I_\varepsilon) - g_j(I_{ex}) D(I_{ex})|}{\delta_\varepsilon} + \frac{|g_j(I_\varepsilon) D(I_{ex}) - g_j(I_{ex}) D(I_\varepsilon)|}{\delta_\varepsilon} \end{aligned}$$

If $\varepsilon \in (0, e^{-e})$ then $R_\varepsilon > r_0$, so for all $\alpha \in B(0, R_\varepsilon)$ we get

$$\begin{aligned} &|g_j(I_\varepsilon) D(I_\varepsilon) - g_j(I_{ex}) D(I_{ex})| \\ &\leq |g_j(I_\varepsilon) - g_j(I_{ex})| \cdot |D(I_\varepsilon)| + |g_j(I_{ex})| \cdot |D(I_\varepsilon) - D(I_{ex})| \\ &\leq C_2 R_\varepsilon \varepsilon \cdot (C_2 + \varepsilon) + C_2 R_\varepsilon \varepsilon \leq (C_2 + 1)^2 R_\varepsilon \varepsilon \end{aligned}$$

and similarly,

$$|g_j(I_\varepsilon) D(I_{ex}) - g_j(I_{ex}) D(I_\varepsilon)| \leq (C_2 + 1)^2 R_\varepsilon \varepsilon$$

Consequently, for all $\varepsilon \in (0, e^{-e})$, we can estimate the first term

$$\chi(B(0, R_\varepsilon)) \left| \frac{g_j(I_\varepsilon) D(I_\varepsilon)}{\delta_\varepsilon + (D(I_\varepsilon))^2} - \frac{g_j(I_{ex}) D(I_{ex})}{\delta_\varepsilon + (D(I_{ex}))^2} \right| \leq \chi(B(0, R_\varepsilon)) \cdot \frac{2(C_2 + 1)^2 R_\varepsilon \varepsilon}{\delta_\varepsilon}$$

Considering the second term, we have

$$\left| \frac{g_j(I_{ex}) D(I_{ex})}{\delta_\varepsilon + (D(I_{ex}))^2} - \frac{g_j(I_{ex})}{D(I_{ex})} \right| = \frac{\delta_\varepsilon |g_j(I_{ex})|}{\left(\delta_\varepsilon + (D(I_{ex}))^2 \right) \cdot |D(I_{ex})|}$$

We always have

$$\frac{\delta_\varepsilon |g_j(I_{ex})|}{\left(\delta_\varepsilon + (D(I_{ex}))^2 \right) \cdot |D(I_{ex})|} \leq \left| \frac{g_j(I_{ex})}{D(I_{ex})} \right| = 2 \left| \int_{\Omega} f_{jex}(x) \cos(\alpha \cdot x) dx \right| \leq 2 \|f_{jex}\|_{L^2(\Omega)}$$

Furthermore, if $\alpha \in B(0, R_\varepsilon) \setminus B_\varepsilon$ then

$$\frac{\delta_\varepsilon |g_j(I_{ex})|}{\left(\delta_\varepsilon + (D(I_{ex}))^2 \right) \cdot |D(I_{ex})|} \leq \frac{\delta_\varepsilon |g_j(I_{ex})|}{|D(I_{ex})|^3} \leq \frac{\delta_\varepsilon C_2 R_\varepsilon}{\varepsilon^{6q}}$$

Therefore, for all $\varepsilon \in (0, e^{-e})$, we can estimate the second term

$$\chi(B(0, R_\varepsilon)) \left| \frac{g_j(I_{\varepsilon x}) D(I_{\varepsilon x})}{\delta_\varepsilon + (D(I_{\varepsilon x}))^2} - \frac{g_j(I_{\varepsilon x})}{D(I_{\varepsilon x})} \right| \leq 2\chi(B_\varepsilon) \|f_{j\varepsilon x}\|_{L^2(\Omega)} + \chi(B(0, R_\varepsilon)) \frac{\delta_\varepsilon C_2 R_\varepsilon}{\varepsilon^{6q}}$$

Thus, for all $\varepsilon \in (0, e^{-e})$, we have

$$\begin{aligned} \left| G_j(I_\varepsilon) - F(\tilde{f}_{j\varepsilon x}) \right| &\leq \chi(B(0, R_\varepsilon)) \left(\frac{2(C_2 + 1)^2 R_\varepsilon \varepsilon}{\delta_\varepsilon} + \frac{\delta_\varepsilon C_2 R_\varepsilon}{\varepsilon^{6q}} \right) \\ &\quad + 2\chi(B_\varepsilon) \|f_{j\varepsilon x}\|_{L^2(\Omega)} + \chi(R^2 \setminus B(0, R_\varepsilon)) \left| F(\tilde{f}_{j\varepsilon x}) \right| \end{aligned}$$

Choosing $\delta_\varepsilon = \varepsilon^{\frac{6q+1}{2}}$ and $C_0 = 2(C_2 + 1)^2 + C_2$, we complete the proof. \square

It is obvious that, for all $j \in \{1, 2\}$, by Lebesgue's dominated convergence theorem, $\chi(R^2 \setminus B(0, R_\varepsilon)) \left| F(\tilde{f}_{j\varepsilon x}) \right|$ converges to 0 in $L^2(R^2)$ when $\varepsilon \rightarrow 0$. However, to get an explicitly estimate for it, some a-priori information about $f_{\varepsilon x}$ must be assume.

Lemma 5. *Let $a \in R$, Q be an measurable subset of R^n ($n \geq 1$), and $w \in L^1(Q) \cap L^2(Q)$. Then*

$$\int_{R^n} \left| \int_Q w(x) \sin\left(a + \sum_{k=1}^n \alpha_k x_k\right) dx \right|^2 d\alpha = 2^{n-1} \pi^n \|w\|_{L^2(Q)}^2$$

Proof. We first prove in the case $a = 0$. Put $\tilde{w} : R^n \rightarrow R$

$$\tilde{w}(x) = \chi(Q)w(x) - \chi(-Q)w(-x)$$

Then $\tilde{w} \in L^1(R^n) \cap L^2(R^n)$ and

$$F_n(\tilde{w})(\alpha) = 2i \int_Q w(x) \sin\left(\sum_{k=1}^n \alpha_k x_k\right) dx$$

where F_n is the Fourier transform in R^n . Using Paserval equality, we get

$$\int_{R^n} \left| \int_Q w(x) \sin\left(\sum_{k=1}^n \alpha_k x_k\right) dx \right|^2 d\alpha = \frac{1}{4} \|F_n(\tilde{w})\|_{L^2(R^n)}^2 = \frac{(2\pi)^n}{4} \|\tilde{w}\|_{L^2(R^n)}^2 = 2^{n-1} \pi^n \|w\|_{L^2(Q)}^2$$

Similarly, we also have

$$\int_{R^n} \left| \int_Q w(x) \cos\left(\sum_{k=1}^n \alpha_k x_k\right) dx \right|^2 d\alpha = 2^{n-1} \pi^n \|w\|_{L^2(Q)}^2$$

Now, we notice that

$$\begin{aligned} \left| \int_Q w(x) \sin\left(a + \sum_{k=1}^n \alpha_k x_k\right) dx \right|^2 &= (\cos(a))^2 \left| \int_Q w(x) \sin\left(\sum_{k=1}^n \alpha_k x_k\right) dx \right|^2 \\ &+ (\sin(a))^2 \left| \int_Q w(x) \cos\left(\sum_{k=1}^n \alpha_k x_k\right) dx \right|^2 + v(\alpha) \end{aligned}$$

where

$$v(\alpha) = \sin(2a) \cdot \int_{\Omega} w(x) \sin\left(\sum_{k=1}^n \alpha_k x_k\right) dx \cdot \int_{\Omega} w(x) \cos\left(\sum_{k=1}^n \alpha_k x_k\right) dx$$

Since $v(-\alpha) = -v(\alpha)$ for all $\alpha \in R^n$, we get $\int_{R^n} v(\alpha) d\alpha = 0$. Thus

$$\begin{aligned} \int_{R^n} \left| \int_Q w(x) \sin\left(a + \sum_{k=1}^n \alpha_k x_k\right) dx \right|^2 d\alpha \\ = (\cos(a))^2 \cdot 2^{n-1} \pi^n \|w\|_{L^2(Q)}^2 + (\sin(a))^2 \cdot 2^{n-1} \pi^n \|w\|_{L^2(Q)}^2 = 2^{n-1} \pi^n \|w\|_{L^2(Q)}^2 \end{aligned}$$

The proof is completed. □

Using Lemma 5, we have the following result.

Lemma 6. *Let $w \in H^1(\Omega)$ and $r > \pi/(2\sqrt{2})$. Then*

$$\int_{R^2 \setminus B(0,r)} \left| \int_{\Omega} w(x) \cos(\alpha \cdot x) dx \right|^2 d\alpha \leq \frac{72\sqrt{2}\pi}{r} \|w\|_{H^1(\Omega)}^2$$

Proof. Since

$$\int_{R^2 \setminus B(0,r)} \left| \int_{\Omega} w(x) \cos(\alpha \cdot x) dx \right|^2 d\alpha \leq \sum_{j=1}^2 \int_{|\alpha_j| \geq r/\sqrt{2}} \left| \int_{\Omega} w(x) \cos(\alpha \cdot x) dx \right|^2 d\alpha$$

, the proof will be completed if we show that, for all $j \in \{1, 2\}$,

$$\int_{|\alpha_j| \geq r/\sqrt{2}} \left| \int_{\Omega} w(x) \cos(\alpha \cdot x) dx \right|^2 d\alpha \leq \frac{24\sqrt{2}\pi}{r} \left(\|w\|_{L^2(\Omega)}^2 + 2 \left\| \frac{\partial w}{\partial x_j} \right\|_{L^2(\Omega)}^2 \right)$$

We will prove for the case $j = 1$, and the other cases are similar. We have

$$\int_{\Omega} w(x) \cos(\alpha \cdot x) dx = \int_0^1 \left[w(x) \frac{\sin(\alpha \cdot x)}{\alpha_1} \right]_{x_1=0}^{x_1=1} dx_2 - \int_{\Omega} \frac{\partial w}{\partial x_1} \cdot \frac{\sin(\alpha \cdot x)}{\alpha_1} dx$$

so

$$\begin{aligned} \left| \int_{\Omega} w(x) \cos(\alpha \cdot x) dx \right|^2 &\leq \frac{3}{\alpha_1^2} \left| \int_0^1 w(1, x_2) \sin(\alpha_1 + \alpha_2 x_2) dx_2 \right|^2 \\ &+ \frac{3}{\alpha_1^2} \left| \int_0^1 w(0, x_2) \sin(\alpha_2 x_2) dx_2 \right|^2 + \frac{3}{\alpha_1^2} \left| \int_{\Omega} \frac{\partial w}{\partial x_1} \cdot \sin(\alpha \cdot x) dx \right|^2 \end{aligned}$$

Therefore,

$$\begin{aligned} \int_{|\alpha_1| \geq r/\sqrt{2}} \left| \int_{\Omega} w(x) \cos(\alpha \cdot x) dx \right|^2 d\alpha &\leq \frac{6}{r^2} \int_{R^2} \left| \int_{\Omega} \frac{\partial w}{\partial x_1}(x) \cdot \sin(\alpha \cdot x) dx \right|^2 d\alpha \\ &+ \int_{|\alpha_1| \geq r/\sqrt{2}} \frac{3}{\alpha_1^2} d\alpha_1 \cdot \int_{-\infty}^{\infty} \left| \int_0^1 w(1, x_2) \sin(\alpha_1 + \alpha_2 x_2) dx_2 \right|^2 d\alpha_2 \\ &+ \int_{|\alpha_1| \geq r/\sqrt{2}} \frac{3}{\alpha_1^2} d\alpha_1 \cdot \int_{-\infty}^{\infty} \left| \int_0^1 w(0, x_2) \sin(\alpha_2 x_2) dx_2 \right|^2 d\alpha_2 \\ &= \frac{12\pi^2}{r^2} \left\| \frac{\partial w}{\partial x_1} \right\|_{L^2(\Omega)}^2 + \frac{6\sqrt{2}\pi}{r} \|w(1, \cdot)\|_{L^2(0,1)}^2 + \frac{6\sqrt{2}\pi}{r} \|w(0, \cdot)\|_{L^2(0,1)}^2 \end{aligned}$$

Noting that

$$w(1, x_2) = \int_0^1 \frac{\partial}{\partial x_1} (x_1 w(x)) dx_1 = \int_0^1 \left(w(x) + x_1 \frac{\partial w}{\partial x_1}(x) \right) dx_1$$

, we get

$$|w(1, x_2)|^2 \leq \int_0^1 \left(2|w(x)|^2 + 2 \left| \frac{\partial w}{\partial x_1}(x) \right|^2 \right) dx_1$$

Hence,

$$\int_0^1 |w(1, x_2)|^2 dx_2 \leq 2 \|w\|_{L^2(\Omega)}^2 + 2 \left\| \frac{\partial w}{\partial x_1} \right\|_{L^2(\Omega)}^2$$

Similarly,

$$\begin{aligned} \int_0^1 |w(0, x_2)|^2 dx_2 &= \int_0^1 \left| \int_0^1 \frac{\partial}{\partial x_1} ((1-x_1)w(x)) dx_1 \right|^2 dx_2 \\ &\leq \int_0^1 \int_0^1 \left(2|w(x)|^2 + 2 \left| \frac{\partial w}{\partial x_1}(x) \right|^2 \right) dx_1 dx_2 = 2 \|w\|_{L^2(\Omega)}^2 + 2 \left\| \frac{\partial w}{\partial x_1} \right\|_{L^2(\Omega)}^2 \end{aligned}$$

Thus, we have

$$\begin{aligned} \int_{|\alpha_1| \geq r/\sqrt{2}} \left| \int_{\Omega} w(x) \cos(\alpha \cdot x) dx \right|^2 d\alpha &\leq \frac{12\pi^2}{r^2} \left\| \frac{\partial w}{\partial x_1} \right\|_{L^2(\Omega)}^2 + \\ &+ \frac{24\sqrt{2}\pi}{r} \left(\|w(1, \cdot)\|_{L^2(\Omega)}^2 + \left\| \frac{\partial w}{\partial x_1} \right\|_{L^2(\Omega)}^2 \right) \leq \frac{24\sqrt{2}\pi}{r} \left(\|w(1, \cdot)\|_{L^2(\Omega)}^2 + 2 \left\| \frac{\partial w}{\partial x_1} \right\|_{L^2(\Omega)}^2 \right) \end{aligned}$$

The proof is completed. \square

Remark 1. By the same way, we can show that, if $w \in H^1(\Omega)$ and $r > \pi/(2\sqrt{2})$ then

$$\int_{R^2 \setminus B(0, r)} \left| \int_Q w(x_1, x_2) \cos(\alpha_1 x_1) \cos(\alpha_2 x_2) dx \right|^2 d\alpha \leq \frac{16\sqrt{2}\pi}{r} \|w\|_{H^1(Q)}^2$$

This result improves immediately the results of [4].

Proof of theorem 2

Proof. Recall that $q, \delta_\varepsilon, R_\varepsilon$ are defined by (7), and $G_j(I_\varepsilon), B_\varepsilon$ are defined by Lemma 4. For all $j \in \{1, 2\}$, we define $f_{j\varepsilon} : R^2 \rightarrow R$

$$f_{j\varepsilon}(\xi) = \frac{1}{4\pi^2} \int_{R^2} G_j(I_\varepsilon)(\alpha) e^{i(\xi \cdot \alpha)} d\alpha$$

Appying Lemma 4, we have $G_j(I_\varepsilon) \in L^1(R^2) \cap L^2(R^2)$, so $f_{j\varepsilon} \in C(R^2) \cap L^2(R^2)$ and $F(f_{j\varepsilon}) = G_j(I_\varepsilon)$. Applying Lemma 4 again, for all $\varepsilon \in (0, e^{-e})$, we get

$$\begin{aligned} \left| F(f_{j\varepsilon}) - F(\tilde{f}_{j\varepsilon}) \right| &\leq \chi(B(0, R_\varepsilon)) C_0 R_\varepsilon \varepsilon^{\frac{1-6q}{2}} \\ &+ 2\chi(B_\varepsilon) \|f_{j\varepsilon}\|_{L^2(\Omega)} + \chi(R^2 \setminus B(0, R_\varepsilon)) \left| F(\tilde{f}_{j\varepsilon}) \right| \end{aligned} \quad (17)$$

where C_0 is a positive constant depending only I_{ex} . It implies that

$$\begin{aligned} \left| F(f_{j\varepsilon}) - F(\tilde{f}_{jex}) \right|^2 &\leq 2\chi(B(0, R_\varepsilon)) C_0^2 R_\varepsilon^2 \varepsilon^{1-6q} \\ &+ 4\chi(B_\varepsilon) \|f_{jex}\|_{L^2(\Omega)}^2 + 2\chi(R^2 \setminus B(0, R_\varepsilon)) \left| F(\tilde{f}_{jex}) \right|^2 \end{aligned}$$

Hence,

$$\left\| F(f_{j\varepsilon}) - F(\tilde{f}_{jex}) \right\|_{L^2(R^2)}^2 \leq 2C_0^2 \pi R_\varepsilon^4 \varepsilon^{1-6q} + 4m(B_\varepsilon) \|f_{jex}\|_{L^2(\Omega)}^2 + 2 \int_{R^2 \setminus B(0, R_\varepsilon)} \left| F(\tilde{f}_{jex}) \right|^2 d\alpha$$

It is obvious that $2C_0^2 \pi R_\varepsilon^4 \varepsilon^{1-6q} \leq R_\varepsilon^{-1}$ for $\varepsilon > 0$ small enough. Moreover, since

$$B_\varepsilon \subset (\{\alpha \in B(0, R_\varepsilon), |D_1(I_{ex})(\alpha)| \leq \varepsilon^q\} \cup \{\alpha \in B(0, R_\varepsilon), |D_2(I_{ex})(\alpha)| \leq \varepsilon^q\})$$

, we apply Lemma 2 (with $\varphi_0(t) = \varphi_{ex}(T - t)$) to get that $m(B_\varepsilon) \leq 2R_\varepsilon^{-1}$ for $\varepsilon > 0$ small enough. Thus, for $\varepsilon > 0$ small enough, we get

$$\left\| F(f_{j\varepsilon}) - F(\tilde{f}_{jex}) \right\|_{L^2(R^2)}^2 \leq \frac{1}{R_\varepsilon} + \frac{8}{R_\varepsilon} \|f_{jex}\|_{L^2(\Omega)}^2 + 2 \int_{R^2 \setminus B(0, R_\varepsilon)} \left| F(\tilde{f}_{jex}) \right|^2 d\alpha d\beta$$

By Paserval equality, we have

$$\begin{aligned} \|f_{j\varepsilon} - f_{jex}\|_{L^2(\Omega)}^2 &\leq \|f_{j\varepsilon} - \tilde{f}_{jex}\|_{L^2(R^2)}^2 = \frac{1}{4\pi^2} \left\| F(f_{j\varepsilon}) - F(\tilde{f}_{jex}) \right\|_{L^2(R^2)}^2 \\ &\leq \frac{1}{4\pi^2} \left(\frac{1}{R_\varepsilon} + \frac{8}{R_\varepsilon} \|f_{jex}\|_{L^2(\Omega)}^2 + 2 \int_{R^2 \setminus B(0, R_\varepsilon)} \left| F(\tilde{f}_{jex}) \right|^2 d\alpha \right) \end{aligned} \quad (18)$$

for $\varepsilon > 0$ small enough. Since $F(\tilde{f}_{jex}) \in L^2(R^2)$, we obtain that

$$\lim_{\varepsilon \rightarrow 0} \|f_{j\varepsilon} - f_{jex}\|_{L^2(\Omega)} = 0$$

If $f_{jex} \in H^1(\Omega)$ then using (18) and Lemma 6, we get

$$\begin{aligned} \|f_{j\varepsilon} - f_{jex}\|_{L^2(\Omega)}^2 &\leq \frac{1}{4\pi^2} \left(\frac{1}{R_\varepsilon} + \frac{8}{R_\varepsilon} \|f_{jex}\|_{L^2(\Omega)}^2 + 2.4 \cdot \frac{72\sqrt{2}\pi}{R_\varepsilon} \|f_{jex}\|_{H^1(\Omega)}^2 \right) \\ &\leq \left(66 \|f_{jex}\|_{H^1(\Omega)}^2 + \frac{1}{4\pi^2} \right) \cdot \frac{1}{R_\varepsilon} = 63eT \left(66 \|f_{jex}\|_{H^1(\Omega)}^2 + \frac{1}{4\pi^2} \right) \cdot \frac{\ln(\ln(\varepsilon^{-1}))}{\ln(\varepsilon^{-1})} \end{aligned}$$

for $\varepsilon > 0$ small enough. This complete the proof. \square

Proof of Theorem 3

Proof. We shall use the notations of the proof of Theorem 2. Notice that the assumption

$$\int_{R^2} \left| \int_{\Omega} f_{jex}(x) \cdot \cos(\alpha \cdot x) dx \right| d\alpha < \infty,$$

is equivalent to $F(\tilde{f}_{jex}) \in L^1(R^2)$. Since $\tilde{f}_{jex}, F(\tilde{f}_{jex}) \in L^1(R^2) \cap L^2(R^2)$, we get

$$\tilde{f}_{jex}(\xi) = \frac{1}{4\pi^2} \int_{R^2} F(\tilde{f}_{jex})(\alpha) e^{i(\alpha \cdot \xi)} d\alpha$$

Therefore,

$$4\pi^2 \|f_{j\varepsilon} - f_{jex}\|_{L^\infty(\Omega)} \leq 4\pi^2 \|f_{j\varepsilon} - \tilde{f}_{jex}\|_{L^\infty(R^2)} \leq \|F(f_{j\varepsilon}) - F(\tilde{f}_{jex})\|_{L^1(R^2)} \quad (19)$$

From (17), we have

$$\|F(f_{j\varepsilon}) - F(\tilde{f}_{jex})\|_{L^1(R^2)} \leq C_0 \pi R_\varepsilon^3 \varepsilon^{\frac{1-3q}{2}} + 2m(B_\varepsilon) \|f_{jex}\|_{L^2(\Omega)} + \int_{R^2 \setminus B(0, R_\varepsilon)} |F(\tilde{f}_{jex})| d\alpha$$

For $\varepsilon > 0$ small enough, we have $C_0 \pi R_\varepsilon^3 \varepsilon^{\frac{1-3q}{2}} \leq R_\varepsilon^{-1}$ and $m(B_\varepsilon) \leq 2R_\varepsilon^{-1}$. Thus, from (19), for $\varepsilon > 0$ small enough, for all $j \in \{1, 2\}$, we get

$$4\pi^2 \|f_{j\varepsilon} - f_{jex}\|_{L^\infty(\Omega)} \leq \frac{1}{R_\varepsilon} + \frac{4}{R_\varepsilon} \|f_{jex}\|_{L^2(\Omega)} + \int_{R^2 \setminus B(0, R_\varepsilon)} |F(\tilde{f}_{jex})| d\alpha$$

Since $F(\tilde{f}_{jex}) \in L^1(R^2)$, we obtain that $\lim_{\varepsilon \rightarrow 0} \|f_{j\varepsilon} - f_{jex}\|_{L^\infty(\Omega)} = 0$ for all $j \in \{1, 2\}$. \square

References

- [1] Timoshenko, S. and J. N. Goodier, *Theory of Elasticity*, New York, Mc Graw-Hill 1970.
- [2] Martin H. Sadd, *Elasticity Theory, Applications, and Numerics*, Elsevier 2005.
- [3] Dang Duc Trong, Nguyen Thanh Long, Pham Ngoc Dinh Alain, *Nonhomogeneous heat equation: Identification and regularization for the inhomogeneous term*, J. Math. Anal. Appl. 312 (2005), 93-104.
- [4] D.D.Trong, P.H.Quan, P.N.Dinh Alain, *Determination of a two-dimensional heat source: Uniqueness, regularization and error estimate*, J. Comp. Appl. Math, 191(2006), 50-67.
- [5] Dang Duc Trong and Truong Trung Tuyen, *Error of Tikhonov's regularization for intergral convolution equation*, arXiv:Math.NA/0610046 v1 1 Oct 2006.

- [6] B.Ya.Levin, *Lectures on Entire Functions*, Trans Math Monographs, Vol.150, AMS, Providence, Rhode Island, 1996.
- [7] Isakov, V., *Inverse source problems*, Math. surveys and monographs series, Vol.34, AMS, Providence, Rhode Island, 1990, chap.7, page. 166
- [8] Tikhonov A. N., *Théorèmes d'unicité pour l'équation de la chaleur*, Math. Sborn.42 (1935), 199-216
- [9] M. Grasselli, M. Ikehata, M. Yamamoto, *An inverse source problem for the Lamé system with variable coefficients*, Applicable Analysis. 84(4)(2005), 357-375.
- [10] Aki, K. and Richards, P.G., 1980, *Quantitative Seismology Theory and Methods*, Vol. I, New York, Freeman.
- [11] M.I. Ivanchov, *The inverse problem of determining the heat source power for a parabolic equation under arbitrary boundary conditions*, J. Math. Sci (New York) 88(3)(1998), 432-436.
- [12] M.I. Ivanchov, *Inverse problem for a multidimensional heat equation with an unknown source function*, Mat. Stud. 16(1)(2001), 93-98.
- [13] D.U. Kim, *Construction of the solution of a certain system of heat equations with heat sources that depend on the temperature*, Izv. Akad. Nauk. Kazak. SSR Ser. Fiz-Mat. (1)(1971), 49-53.
- [14] G.S. Li, L.Z. Zhang, *Existence of a nonlinear heat source in inverse heat conduction problems*, Hunan Ann. Math. 17(2)(1997), 19-24