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Formulary of geodesics of the projected averaged Kepler Hamiltonian*

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Abstract

This short note gives the quadratures of the geodesics of the averaged Hamiltonian of the controlled Kepler equation (energy criterion) projected on S^2 . The endpoints of the corresponding cut locus are also deduced, as well as the injectivity radius of the associated Riemannian metric on the 2-sphere.

Keywords. Averaged Kepler equation, geodesics on S^2 , cut locus.

Classification AMS. 49K15, 70Q05

The Hamiltonian is

$$H = \frac{1}{2} \left(\frac{p_{\theta}^2}{G(\varphi)} + p_{\varphi}^2 \right)$$

with

$$G(\varphi) = \frac{\sin^2 \varphi}{1 - (1 - \mu^2)\sin^2 \varphi}.$$

Herebefore, μ is a parameter which is equal to $1/\sqrt{5}$ in the case of the Kepler equation [1, 2, 3, 4]. The system is clearly integrable, and the geodesics on the level set $\{H=1/2\}$ of the corresponding Riemannian metric are as parameterized by p_{θ} (note that θ is cyclic) as follows:

$$\varphi = a\sin\sqrt{\frac{(1+b) - (1-b)\cos(a(t-t_1))}{2}},$$

and

$$\theta = \operatorname{sign}(p_{\theta}) \left[\operatorname{atan} \left(\frac{\tan \left(a(t - t_1)/2 \right)}{\sqrt{b}} \right) \right]_0^t - (1 - \mu^2) p_{\theta} t.$$

These two quadratures are valid for t in $[t_1, t_1 + T/4]$ with nonegative p_{φ_0} , and extended on $[t_1, t_1 + T]$ by using the obvious symmetries of the Hamiltonian (a $-2t_1$ time translation gives the case when p_{φ_0} is negative). The coordinate

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 φ and the time derivative $\dot{\theta}$ are *T*-periodic and hence clearly extended to the whole real line. The period *T*, the time t_1 , and the parameters *a* and *b* have the following expression (θ_0 is assumed to be zero by symmetry):

$$\begin{array}{rcl} a & = & 2\sqrt{1+(1-\mu^2)p_{\theta}^2}, \\ b & = & p_{\theta}^2/(1+(1-\mu^2)p_{\theta}^2), \\ t_1 & = & (1/a)\left[-\pi/2-\mathrm{asin}\left(\frac{2\sin^2\varphi_0-(1+b)}{1-b}\right)\right], \\ T & = & 4\pi/a. \end{array}$$

The cut locus [2] of the Riemannian metric is then easily deduced. For symmetry reasons, it is included in the antipodal parallel $\{\pi - \varphi_0\}$, and the coordinates of its left endpoint are

$$\theta_l = \pi(1 - (1 - \mu^2)\sin\varphi_0),$$

$$\varphi_l = \pi - \varphi_0.$$

The time, that is the distance to the cut locus is

$$t_l = \pi \sqrt{1 - (1 - \mu^2) \sin^2 \varphi_0},$$

which gives the injectivity radius of the Riemannian metric $G(\varphi)d\theta^2 + d\varphi^2$ on \mathbf{S}^2 , that is the infimum of distances from a point to its cut locus (clearly reached on the equator, for $\varphi_0 = \pi/2$),

$$i(\mathbf{S}^2) = \mu \pi.$$

Another consequence of the computations is the existence of a simple closed geodesic which is non-trivial, that is different from the equator $\{\varphi=\pi/2\}$ or a meridian $\{\theta=\mathrm{cst}\}$, for $p_\theta=5/4$.

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